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CORRELATION OF DRAIN BREAKDOWN WITH EXCESS NOISE AND OTHER SURFACE-RELATED PHENOMENA

IN ENHANCEMENT MOSFETS

ΒY

JERRY J. RIJ

A THESIS

PRESENTED IN PARTIAL FULFILLMENT OF

THE REQUIREMENTS FOR THE DEGREE

0F

MASTER OF SCIENCE IN ELECTRICAL ENGINEERING

ΑT

NEWARK COLLEGE OF ENGINEERING

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> Newark, New Jersey 1973

APPROVAL OF THESIS

CORRELATION OF DRAIN BREAKDOWN WITH EXCESS NOISE AND OTHER SURFACE-RELATED PHENOMENA

IN ENHANCEMENT MOSFETS

ΒY

JERRY J. RIJ

FOR

DEPARTMENT OF ELECTRICAL ENGINEERING NEWARK COLLEGE OF ENGINEERING

ΒY

FACULTY COMMITTEE

APPROVED:

NEWARK, NEW JERSEY SEPTEMBER, 1973

ABSTRACT

AN INVESTIGATION OF SEMICONDUCTOR SURFACE - RELATED PHENOMENA HAS BEEN UNDERTAKEN TO CORRELATE 1/F NOISE WITH DRAIN BREAKDOWN IN P-CHANNEL ENHANCEMENT MOSFET'S. | N -CREASES IN THE INTENSITY OF DRAIN CURRENT FLUCTUATIONS AT 10 Hz AND AT 1 KHz, PARTICULARLY AT THE HIGHER FREQUENCY, WAS OBSERVED FOLLOWING ACCELERATED LIFE-TESTING FOR THRESH-OLD VOLTAGE DRIFT. IT WAS CONCLUDED THAT MOBILE IONS NEAR THE OXIDE-SEMICONDUCTOR INTERFACE PRODUCE A SHARP INCREASE IN FAST SURFACE STATES AND THAT THESE STATES MAY BE REGARDED AS FAST TRAPPING CENTERS. IT WAS ALSO FOUND THAT THE 1/F NOISE INTENSITY CONSISTENTLY PEAKED AT THE THRESHOLD OF DRAIN BREAKDOWN AND THAT IT STEADILY DECREASED WITH FUR-THER INCREASES IN DRAIN CURRENT. OF ALL THE TRANSISTORS TESTED, THOSE WITH RELATIVELY LOW NOISE INTENSITY WERE FOUND TO EXHIBIT SHARPER BREAKDOWN CHARACTERISTICS AND HIGHER BREAKDOWN VOLTAGES WHILE THOSE TRANSISTORS WITH HIGH NOISE SHOWED SOFT BREAKDOWN CHARACTERISTICS. IT WAS THEREFORE CONCLUDED THAT LOW NOISE MOS TRANSISTORS ARE SUPERIOR TO THOSE WITH RELATIVELY HIGH NOISE AND THAT THE 1/F NOISE - DRAIN CURRENT CHARACTERISTICS MAY BE USED IN NONDESTRUCTIVE TESTING TO DETERMINE THE APPROXIMATE DRAIN BREAKDOWN VOLTAGE.

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Ρ	A	G	Е
-	-	-	-

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LIST OF SYMBOLS

Α	AREA OF GATE METALIZATION IN CM ²
Å	Angstrom unit, 10 ⁻¹⁰ meters, 10 ⁻⁸ cm
В	BANDWIDTH IN CYCLES PER SECOND, HERTZ, HZ
c	CAPTURE RATE PROBABILITY
¢ _s	CAPTURE RATE PROBABILITY AT THE SI-SIO ₂ INTERFACE
С _в	Impurity concentration in the substrate per cm 3
Co	Capacitance of the oxide layer in Farads/cm ²
Cs	Capacitance of the semiconductor space-charge region in Farads/cm ²
E	Surface-state energy in electron volts, eV
Е _в	BARRIER POTENTIAL TO BE EXCEEDED IN THE TUNNELING MECHANISM OF HOLES
Ec	CONDUCTION-BAND ENERGY LEVEL
EF	FERMI ENERGY LEVEL
Ε _ν	VALENCE-BAND ENERGY LEVEL
E,	TRANSVERSE ELECTRIC FIELD IN THE CHANNEL DEPLETION REGION BETWEEN THE EXPOP (EXTRAPOLATED PINCH-OFF POINT) AND THE DRAIN-TO-SUBSTRATE JUNCTION
F, f	FREQUENCY IN CYCLES PER SECOND
f _T	QUASI-FERMI FUNCTION OF SURFACE STATES
f_{TP}	$1 - t_r$
G	The rate that surface states of energy E located at a distance x into the oxide layer gain holes
Gм	Transconductance in ohms ⁻¹ , mhos
9dsat	Small-signal saturation drain conductance in ohms ⁻¹
ħ	Planck's constant, 4.136×10 ⁻¹⁵ eV·second

- *i*² *i*² MEAN SQUARED NOISE CURRENT WITH DRAIN SHORTED IN MA²/cycle
- DRAIN CURRENT
- DSAT SATURATION DRAIN CURRENT
- I_{DSAT} * SATURATION DRAIN CURRENT WHEN $V_D = V_{DSAT}$
- DRAIN LEAKAGE CURRENT WHEN THE GATE IS RETURNED TO THE SOURCE, THE SOURCE AND SUBSTRATE ARE COM-MON, AND THE DRAIN-TO-SOURCE VOLTAGE IS 10 VOLTS
- J CURRENT DENSITY
- k BOLTZMANN'S CONSTANT, 8.617×10⁻⁵eV·°K⁻¹
- K MATERIAL CONSTANT
- Ko Relative permittivity of SiO2, 3.54×10⁻¹³Farad/cm
- K_S Relative permittivity of silicon
- kT ENERGY LEVEL, 0.0258EV at $T = 300^{\circ}\text{K}$
- Length of Diffusion of surface states from the SI-SIO2 interface
- L EFFECTIVE CHANNEL LENGTH (IN THE DIRECTION OF CURRENT FLOW)
- ΔL . Length by which the conducting channel is shortened when V_{DS} exceeds $V_{DS\,\Delta\,T}$
- M^{*} EFFECTIVE MASS OF A HOLE, 9.108×10⁻³¹Kg
- NA ACCEPTOR DOPING DENSITY PER CM³
- N_D DONOR DOPING DENSITY PER CM³
- $N_{T}(E)$ Surface-state density of free holes and electrons in cm⁻³·eV⁻¹
- h_{si} Free electron concentration at the surface when the Fermi level is at the surface-state level
- hso STEADY-STATE VALUE OF hsi
- $h_{\tau}(\epsilon)$ Empty surface states at energy E located at a distance x from the SI-SIO₂ interface
- P⁺ HEAVY P-TYPE DIFFUSION (10¹⁸ to 10²⁰ cm⁻³)

- P_s Hole density in the inverted channel in cm⁻³.eV⁻¹
- P_{s_i} Free hole concentration at the surface when the Fermi level is at the surface-state level
- PSO STEADY-STATE VALUE OF PSI

Q_{SS} Surface-state charge in coulombs

9 MAGNITUDE OF ELECTRONIC CHARGE, 1.602×10⁻¹⁹coul

R RESISTANCE IN OHMS

- R RATE THAT SURFACE STATES OF ENERGY E LOCATED AT A DISTANCE X INTO THE OXIDE LAYER LOSE HOLES
- Si(w) SHORT-CIRCUIT DRAIN CURRENT NOISE SPECTRAL INTEN-SITY (DRAIN CURRENT FLUCTUATIONS IN MOSFET'S)
- $S_{p\tau(\omega)}$ Spectral intensity of the trapped hole density fluctuation
- T TEMPERATURE IN DEGREES KELVIN, °K
- V VOLUME
- VD DRAIN VOLTAGE
- VDS DRAIN-TO-SOURCE VOLTAGE
- V_{DSAT} Drain voltage which causes channel pinch-off
- V_{FB} FLAT-BAND VOLTAGE NECESSARY TO COUNTER BALANCE THE WORK FUNCTION DIFFERENCE BETWEEN THE METAL-OXIDE BARRIER ENERGY AND THE SILICON-OXIDE BAR-RIER IN AN MOS TRANSISTOR
- VG GATE VOLTAGE
- VG' EFFECTIVE GATE VOLTAGE
- V_{RS} GATE-TO-SOURCE VOLTAGE
- VGSUB GATE-TO-SUBSTRATE VOLTAGE
- VT THRESHOLD VOLTAGE
- x DISTANCE FROM THE SI-SIO₂ INTERFACE
- X_O THICKNESS OF THE OXIDE LAYER
- Z PARABOLIC RATE CONSTANT FOR GROWTH OF SIO2

- α CAPTURE RATE TIME CONSTANT
- β GAIN TERM
- $\mathfrak{F}, \mathfrak{I}, \mathfrak{O}, \mathfrak{A}$ Exponents of $\mathbb{R}, \mathfrak{F}, \mathfrak{I}, \mathfrak{v}$ in Bell's general equation for frequency-dependent noise
- δ INFINITESIMALLY SMALL INCREMENT
- €. PERMITTIVITY OF FREE SPACE, 8.86×10⁻¹⁴Farad/cm
- 5,5 FIELD FRINGING FACTORS EXPERIMENTALLY FOUND TO BE 0.2 AND 0.6, RESPECTIVELY
- MOBILITY OF ELECTRONS IN CM²/VOLT.SECOND
- MOBILITY OF HOLES IN CM²/VOLT.SECOND
- *π* P1, 3.1415926
- P RESISTIVITY
- S SHAPE FACTOR, EG. CROSS-SECTION/UNIT LENGTH
- TIME CONSTANT FOR SURFACE-STATE CHARGE DENSITY FLUCTUATION
- $r_{\rm s}$ Time constant for surface states located right at the SI-SIO₂ interface
- ELECTROSTATIC POTENTIAL WITH REFERENCE TO THE BULK SEMICONDUCTOR MATERIAL
- ϕ_F ELECTROSTATIC POTENTIAL AT THE FERMI LEVEL; THE AMOUNT BY WHICH THE FERMI LEVEL IS DISPLACED FROM THE INTRINSIC LEVEL OR THE CENTER OF THE GAP (AS MEASURED IN THE BULK)
- \$\$ ELECTROSTATIC POTENTIAL AT THE SURFACE; THE AMOUNT BY WHICH THE INTRINSIC FERMI LEVEL, AT THE SURFACE, HAS BEEN CHANGED WITH RESPECT TO THE FERMI LEVEL
- ω ANGULAR FREQUENCY IN RADIANS PER SECOND

CHAPTER 1

INTRODUCTION AND THEORY OF MOSFET OPERATION

1.1 INTRODUCTION TO BASIC MOS DEVICES

THE BASIC METAL-OXIDE-SEMICONDUCTOR FIELD-EFFECT TRANSISTOR (MOSFET) IS COMPOSED OF TWO HEAVILY-DOPED REGIONS, THE <u>SOURCE</u> AND <u>DRAIN</u>, DIFFUSED INTO A SEMICON-DUCTOR <u>SUBSTRATE</u> AND SEPARATED BY A <u>CHANNEL</u>. AN <u>INSULA-</u> <u>TED METAL GATE</u> ABOVE THE CHANNEL MODULATES <u>MAJORITY CAR-</u> <u>RIER</u> CONDUCTION FROM THE SOURCE TO THE DRAIN BY FIELD EFFECTS.

Two types of devices are commercially available: the p-channel MOSFET (in which holes are the majority carriers) and the n-channel MOSFET (in which the majority carriers are electrons). In addition, there are two common modes of operation: the <u>enhangement-mode</u> (in which the device is normally off and channel conduction is <u>enhanced</u> by the appropriate gate voltage) and the <u>depletion-mode</u> (in which the channel normally conducts and a gate voltage <u>depletes</u> it of majority carriers). Until recently, it was custom ary to produce primarily either p-channel enhancement- or n-channel depletion-mode MOS transistors. However, complimentary MOS integrated circuits consisting of pairs of both p- and n-channel transistors on a single chip are replacing a large portion of the bipolar devices in use today.

1.2 THEORY OF MOSFET OPERATION

MODULATION OF CHANNEL CONDUCTION. FIGURE 1-1 SHOWS THE GEOMETRIC CONSTRUCTION OF A P-CHANNEL ENHANCEMENT -MODE MOS TRANSISTOR PROPERTLY BIASED FOR CHANNEL CONDUC-TION. SUCH DEVICES REQUIRE A SUFFICIENTLY NEGATIVE GATE VOLTAGE BEFORE THE CHANNEL BECOMES <u>INVERTED</u>. THIS THEN FORMS A P TYPE CONDUCTING PATH BETWEEN THE SOURCE AND THE DRAIN, THE CONDUCTIVITY OF WHICH IS <u>ENHANCED</u> WITH INCREAS-INGLY NEGATIVE GATE BIAS. N-CHANNEL ENHANCEMENT MOSFET'S HAVE N⁺ SOURCE AND DRAIN REGIONS AND REQUIRE A DRAIN AND GATE BIAS OF OPPOSITE POLARITY TO THOSE OF P-CHANNEL MOS TRANSISTORS IN ORDER TO OPERATE. DEPLETION-MODE DEVICES HAVE A BUILT-IN CONDUCTING CHANNEL BENEATH THE OXIDE OF THE SAME TYPE OF MATERIAL USED FOR THEIR RESPECTIVE DRAIN AND SOURCE REGIONS.

FIGURE 1-2A IS A LATERAL VIEW TAKEN IN THE CHANNEL BETWEEN THE SOURCE AND DRAIN OF A P-CHANNEL ENHANCEMENT MOSFET. V_{GS} is the gate bias applied between the gate and source terminals. With the source and substrate at an electrical ground, there are three gate bias conditions of interest.

FIGURE 1-2B ILLUSTRATES THE SEMICONDUCTOR ENERGY BANDS and characteristic "band bending" when the applied V_{GS} is greater than zero. An abrupt difference is assumed between the oxide (SiO₂) and the semiconductor substrate. The energy

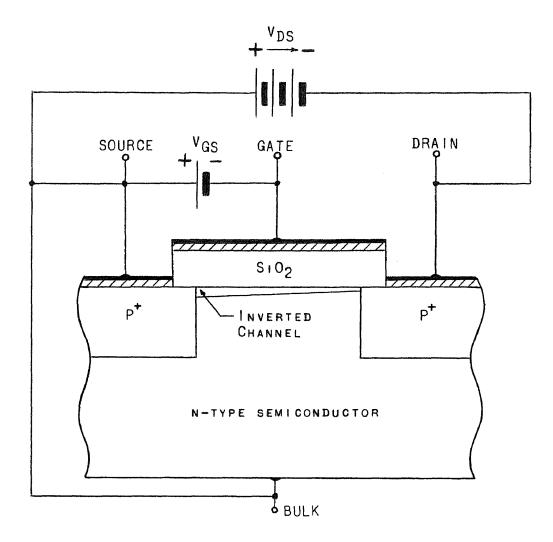


FIGURE 1-1. GEOMETRICAL CONSTRUCTION OF A P-CHANNEL EN-HANCEMENT MOSFET BLASED FOR CHANNEL CONDUCTION. SUCH DEVICES HAVE ISOLATED P⁺ SOURCE AND DRAIN DIFFUSIONS WITHIN AN N-TYPE SEMICONDUCTOR SUBSTRATE. AN OXIDE LAYER, USUALLY SIO₂, IN-SULATES THE GATE FROM THE SUBSTRATE AND PASSIVATES THE SUR-FACE OF THE DRAIN AND SOURCE P-N JUNCTIONS. TERMINALS ARE PROVIDED FOR THE DRAIN, GATE, AND SOURCE AND, IN SOME CASES, A SEPARATE TERMINAL FROM THE SUBSTRATE BULK MAY ALSO BE EX-TERNALLY AVAILABLE. OTHERWISE, THE SUBSTRATE IS CONNECTED INTERNALLY TO THE SOURCE AND THE DEVICE IS THEN USUALLY IN-TENDED TO BE OPERATED WITH ONLY ONE OF THE P⁺ REGIONS AS THE DRAIN.

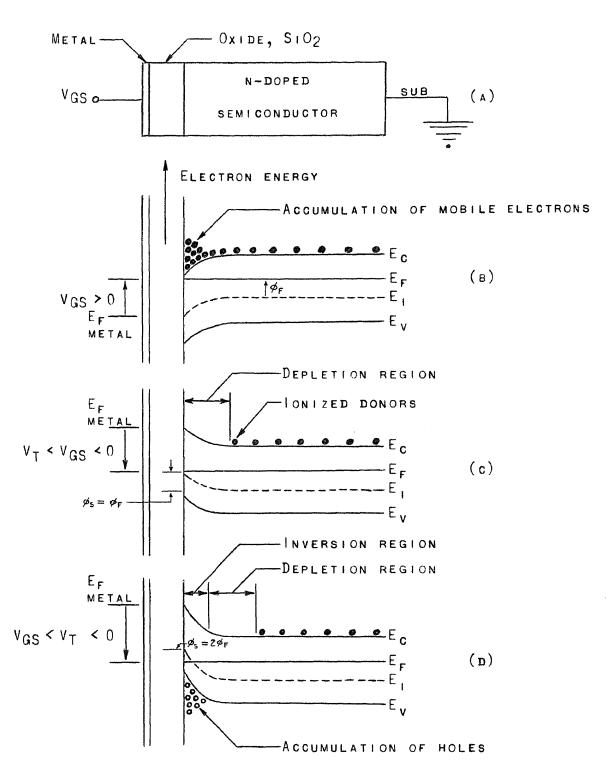


FIGURE 1-2. ENERGY BAND-BENDING NEAR THE SEMICONDUCTOR SURFACE UNDER THREE CONDITIONS OF GATE BIAS FOR A P-CHANNEL ENHANCEMENT MOSFET! (B) POSITIVE V_{GS} RESULTS IN AN ELECTRON ACCUMULATION BENEATH THE OXIDE; (C) SMALL $|-V_{GS}|$ RESULTS IN DEPLETION REGION; AND (D) LARGER $|-V_{GS}|$ CAUSES SURFACE INVERSION AND AN ACCUMULATION OF HOLES.

BANDS ARE IDENTIFIED AS FOLLOWS:

- E ELECTRON ENERGY AT THE EDGE OF THE CONDUCTION BAND
- E ELECTRON ENERGY AT THE FERMI LEVEL, DEFINED AS THE ENERGY LEVEL AT WHICH "THE PROBABILITY OF OCCUPATION OF AN ENERGY STATE BY AN ELECTRON IS EXACTLY ONE-HALF."²
- E, ELECTRON ENERGY AT THE INTRINSIC FERMI LEVEL, THE FERMI LEVEL IN AN INTRINSIC SEMICONDUCTOR
- Ev Electron energy at the edge of the valence band.

A POSITIVE GATE BIAS LOWERS THE ENERGY LEVEL ASSOCI-ATED WITH THE SURFACE STATES DOWN NEAR THE FERMI LEVEL IN THE SEMICONDUCTOR MATERIAL. THE PROBABILITY OF THE STATES BEING OCCUPIED BY AN ELECTRON IS GREATLY INCREASED AND AN ELECTRON ACCUMULATION RESULTS BENEATH THE OXIDE.

FIGURE 1-2c shows the effect of applying a small negative gate bias. As the states are elevated, their probability of occupation by electrons is decreased. The channel near the oxide becomes depleted of electrons. This results in a <u>depletion region</u> which consists of mostly bound ions of impurity atoms.

WHEN THE GATE BIAS IS SUFFICIENTLY NEGATIVE, THE PRO-BABILITY OF OCCUPATION APPROCHES ZERO. ENOUGH ELECTRONS ARE THEN DISPLACED FROM THE SEMICONDUCTOR SURFACE (HEREAFTER REFERRED TO SIMPLY AS THE <u>SURFACE</u>) THAT THEY BECOME THE MIN-ORITY CARRIERS, WHILE HOLES BECOME DOMINANT. FIGURE 1-2D SHOWS A NET ACCUMULATION OF HOLES IN THE SUBSTRATE DIRECTLY BENEATH THE OXIDE. WHEREAS THE CHANNEL WAS HITHERTO COM-POSED OF N-DOPED SILICON, THE SURFACE IS NOW WITHIN AN <u>INVERSION REGION</u>, THEREBY ALLOWING A HOLE CURRENT TO FLOW FROM THE SOURCE TO THE DRAIN WHEN THE DRAIN IS BIASED AT A POTENTIAL NEGATIVE WITH RESPECT TO THE SOURCE.

The negative gate voltage which permits a current of ten microamperes to flow with an applied $-V_{DS}$ equal to $-V_{GS}$ is defined to be the <u>threshold voltage</u>, V_T . It is so-named because it is at this potential that the surface is at the threshold of inversion.

It is evident that the depth, and so the conductance, of the conducting channel may be easily modulated by variations in the applied gate voltage. The MOS transistor acquires its property of amplification from the fact that a relatively small variation in $V_{\mbox{GS}}$ induces a large variation in drain current.

<u>Operation beyond channel pinch-off</u>. Figure 1-3a shows a cross-section of a p-channel enhancement-mode transistor with gate voltage large enough to establish a conducting channel between the source and drain. In the common-source configuration it is customary to ground the source and substrate bulk and apply a negative voltage to the drain. A small $|-V_{DS}|$ results in the bias condition formerly illustrated in figures 1-1 and 1-2D above.

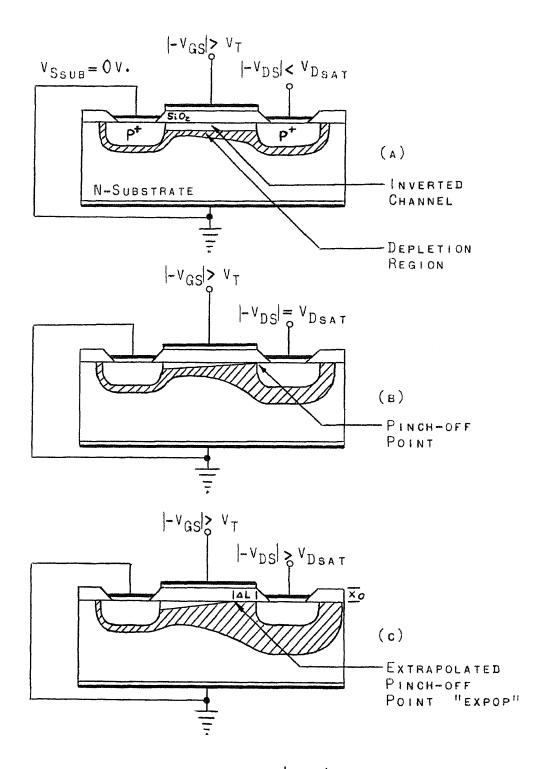


Figure 1-3. Effects of increasing $|-V_{\rm DS}|$ on the inversion region of a p-channel enhancement MOSFET. (a) Condition with small $|-V_{\rm DS}|$; (b) Pinch-off condition; and (c) Condition beyond channel pinch-off.

Since both the drain and source regions are of P^+ type silicon (N_A between 10^{18} and 10^{20} cm⁻³) in an n-doped semiconductor substrate (N_D $\simeq 10^{15}$ cm⁻³), there is a natural <u>Recombination</u> process as electrons fill holes near the p-n junctions. With the drain and source heavily doped, recombination extends much further into the substrate than into either the source or the drain. This results in a <u>space-charge depletion region</u> around the source and drain junctions. Furthermore, since the substrate is electrically connected to the source, any applied $-V_{DS}$ reverse biases the drain-to-substrate p-n junction, thereby further increasing the width of that particular depletion region.³

IF $|-V_{DS}|$ is small compared to $|-V_{GS}|$, there is a linear drop in potential along the channel as hole current proceeds from the source to the drain. In this condition, the drain current $-I_D$ (negative to indicate that the majority carriers are holes) is linearly proportional to $|-V_{DS}|$ and the transistor is said to be operating in the <u>linear portion</u> of its I-V (current-voltage) characteristics.

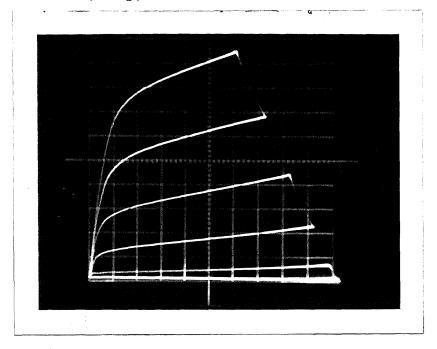
It is evident that as $|-V_{DS}|$ is increased, the difference in potential between the gate and drain terminals is lessened. Since the inversion region is dependent upon the electric field across the oxide, which is a function of the potential difference across each differential element of the oxide along the channel, the width of the inversion region is

DIRECTLY PROPORTIONAL TO THE ELECTRIC FIELD INTENSITY IN THAT PORTION OF THE OXIDE DIRECTLY ABOVE IT. CONSEQUENTLY, AS $|-V_{DS}|$ increases, the field in the oxide at the drain-end of the channel decreases. Eventually, if $|-V_{DS}|$ is made high enough, the field intensity is unable to support channel inversion near the drain-to-substrate junction.⁴ Figure 1-3B illustrates this effect at the onset of channel <u>pinch-off</u>, a condition marked by essentially zero channel depth at the drain-end.^{*}

FIGURE 1-3c shows the effect of further increasing $|-V_{DS}|$. The length of the inverted channel decreases by an amount ΔL and the <u>expop</u> (extrapolated pinch-off point at the tip of the inversion region) recedes toward the source. If the resistivity of the inverted region is averaged over its entire length, any shortening yields a decrease in channel resistance and a corresponding increase in drain current.

IF IT IS ASSUMED THAT A DIFFERENTIAL INCREASE IN DRAIN VOLTAGE BEYOND THAT REQUIRED FOR SATURATION IS ACCOMPANIED BY A DIFFERENTIAL INCREASE IN DRAIN CURRENT DUE ENTIRELY TO A DIFFERENTIAL REDUCTION IN CHANNEL LENGTH, THEN THE INSTAN-

^{*}Actually, the channel depth must be some finite dimension in order for drain current to flow. "Zero channel depth" is an approximation used for the convience of defining the channel pinch-off condition. Refer to sources 1969-2 and 1965-6 in the BIBLIOGRAPHY for a more detailed explanation of the saturation phenomena.



DRAIN CURRENT, -ID, 2MA/DIVISION

DRAIN VOLTAGE, -VDS, 5 VOLTS/DIVISION

Gate Voltage in one volt increments starting from $V_{\mbox{GS}}=-6$ volts

FIGURE 1-4. PHOTOGRAPH SHOWING THE CURRENT-VOLTAGE CHARACTERISTICS OF A TYPICAL P - CHANNEL ENHANCEMENT MOSFET. The curves were taken from a Tetronix Type 527 Transistor Curve Tracer with a slight modification which allowed for Variations in applied gate voltage. TANEOUS SLOPE OF THE I-V CHARACTERISTICS IN THE SATURATION REGION YIELDS THE SMALL-SIGNAL SATURATION DRAIN CONDUCTANCE. This quantity may be expressed by the following:⁵

$$9_{dsat} = \frac{d | Dsat}{dV_D} = \frac{d | Dsat}{dL} \cdot \frac{dL}{dV_D}$$
 1-1

FIGURE 1-4 SHOWS THE 1-V CHARACTERISTICS OF A TYPICAL P-CHANNEL ENHANCEMENT MOSFET. THE LINEAR AND SATURATION REGIONS ARE LABELED. IT IS EVIDENT THAT CHANNEL SHORTEN-ING YIELDS A FINITE OUTPUT RESISTANCE FOR MOS TRANSISTORS.

1.3 PREVIOUS INVESTIGATIONS OF THE SATURATION PHENOMENON

FROHMAN-BENTCHKOWSKY AND GROVE⁶ AND OTHERS^{7,8} Have studied MOSFET saturation conductance in terms of device parameters. Appendix A reviews the analytical analysis of the effects of oxide thickness, electric field intensity and associated fringing effects, and the semiconductor impurity concentration on drain conductance. In general, the saturation drain conductance can be expected to increase with both <u>increased oxide thickness</u> (see equation A-12) and <u>decreased impurity concentration</u> (see equation A-9) in the substrate.⁹

Devices which exhibit abnormally high drain conductance in the saturation region can therefore also be expected to have a thicker oxide layer, a lower substrate impurity concentration, or both. This investigation suggests that the saturation drain conductance is also a function of the surface-

STATE DENSITY AT THE OXIDE-SEMICONDUCTOR INTERFACE.

Reddi and Sah¹⁰ have shown that the drain-to-source voltage required for pinch-off decreases with both increasing oxide thickness and higher bulk doping. Since small increases in the number of impurities used to dope the substrate yield great increases in the number of free carriers according to the familiar equation, ¹¹

$$h p = h_i^2 = K^2 T^3 E x P (-E_c/kT)$$
 1-2

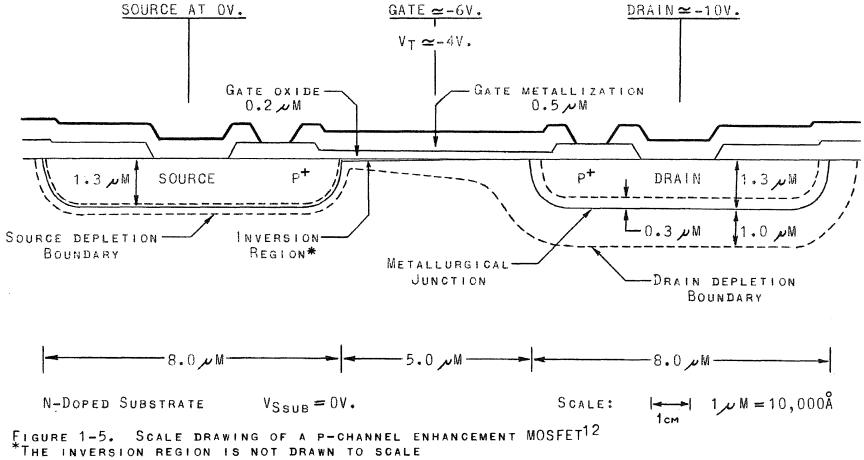
where $E_{g} = E_{c} - E_{v} = Forbidden gap in eV$

- k = BOLTZMANN'S CONSTANT
- $K = MATERIAL CONSTANT, 3.87 \times 10^{22} \text{ for silicon}$ and 1.76×10²² for germanium at 25°C
- $h = \text{Concentration of electrons per cm}^3$
- a = Concentration of holes per cm³
- $T = TEMPERATURE IN \circ K$,

IT IS REASONABLE TO EXPECT THAT A LARGER GATE BIAS IS NEC-ESSARY TO DEPLETE, AND FINALLY INVERT THE CHANNEL WHEN THE SUBSTRATE IMPURITY CONCENTRATION IS HIGH.

FIGURE 1-5 ILLUSTRATES A P-CHANNEL ENHANCEMENT MOSFET DRAWN TO SCALE. THE CONTOURS OF THE DEPLETION REGIONS AND THE VARIOUS DIMENSIONS APPROXIMATE THEORETICAL BOUNDARIES BASED ON CUSTOMARY IMPURITY DOPING CONCENTRATIONS AND APPLIED TERMINAL VOLTAGES. SINCE THE WIDTH OF THE INVERSION REGION IS TYPICALLY LESS THAN 100Å, CHANNEL INVERSION AND DRAIN CONDUCTANCE ARE BOTH ESSENTIALLY SURFACE EFFECTS. THEREFORE, THE CHARACTERISTICS OF THE OXIDE-SEMICONDUCTOR INTERFACE LARGELY DETERMINE THE OBSERVED PROPERTIES OF MOS TRANSIS-TORS.

SUCCEEDING CHAPTERS DISCUSS THE SI-SIO₂ INTERFACE IN MORE DETAIL. CORRELATIONS ARE MADE BETWEEN SURFACE-STATE DENSITY, SEMICONDUCTOR IMPURITY CONCENTRATION, AND BOTH THE EXCESS NOISE AND SATURATION DRAIN CONDUCTANCE PHENOM-ENA AND HOW THEY RELATE TO THE TYPE AND EXTENT OF DRAIN BREAKDOWN IN ENHANCEMENT-MODE MOS TRANSISTORS.



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³D. Frohman-Bentchkowsky and A. S. Grove, "Conductance of MOS Transistors in Saturation," <u>IEEE Transactions on</u> <u>Electron Devices</u> (1969), p. 110. Ref. 1969-2.

⁴Op. Cit., Gosling et al., p. 39.

⁵<u>Op</u>. <u>Cit</u>., Bentchkowsky and Grove, p. 110.

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⁸J. A. Guerst, "Theory of Insulated-Gate Field-Effect Transistors Near and Beyond Pinch-Off," <u>Solid-State Elec-</u> <u>tronics</u>(GB) (1966), pp. 129-42. Ref. 1966-5.

⁹<u>Op</u>. <u>Cit</u>., Bentchkowsky and Grove, p. 110.

¹⁰V. G. K. REDDI AND C. T. SAH, "Source to Drain Resistance Beyond Pinch-Off in Metal-Oxide-Semiconductor Transistors (MOST)," <u>IEEE Transactions on Electron Devices</u>(1965), P. 140. Ref. 1965-6.

¹¹HARRY E. STEWART, <u>ENGINEERING</u> <u>ELECTONICS</u> (BOSTON, 1969), PP. 28-29. Ref. 1969-8.

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CHAPTER 2

FABRICATION AND PROPERTIES OF THE SI-SIO2 INTERFACE

2.1 INTRODUCTION

MANY OPERATING CHARACTERISTICS OF MOS FIELD-EFFECT TRANSISTORS, PARTICULARLY THOSE WHICH FUNCTION IN THE EN-HANCEMENT MODE, MAY BE TRACED DIRECTLY TO THE PROPERTIES OF THE SILICON-SILICON DIOXIDE INTERFACE. THE FOLLOWING DISCUSSIONS QUALITATIVELY INVESTIGATE THESE PROPERTIES AND REVIEW THE MECHANISMS BY WHICH THE INSULATING LAYER IS GROWN ON SILICON. EXPERIMENTAL EVIDENCE IS OFFERED TO EXPLAIN THRESHOLD VOLTAGE DRIFT IN TERMS OF A REDIS-TRIBUTION OF MOBILE IMPURITY IONS WITHIN THE OXIDE. THE MAGNITUDE OF SUCH DRIFT IS CORRELATED WITH THE INTENSITY OF DRAIN CURRENT FLUCTUATIONS AND IT IS SHOWN THAT TRAN-SISTORS WHICH EXHIBIT HIGH NOISE ALSO TEND TO SHOW HIGH THRESHOLD VOLTAGE DRIFT. MEASUREMENTS OF ¹DSS CONFIRMED THAT THERE IS NO MEANINGFUL CORRELATION BETWEEN LEAKAGE AT THE DRAIN P-N JUNCTION AND THRESHOLD VOLTAGE DRIFT.

2.2 BACKGROUND INFORMATION

<u>Growth and oxidation of silicon</u>. Crystal defects may be introduced during the growth of n- or p-doped monocrystaline silicon. Crystal defects may be thermally stress-induced near the liquid-solid interface. During epitaxial growth, vapor-phase reduction of silicon-tetrachloride introduces defects as the temperature of growth is REDUCED.

THE OVERALL REVERSIBLE CHEMICAL REACTION¹

$$SICL_4 + 2H_2 \implies SI(SOLID) + 4HCL 2-1$$

SHOWS THAT HYDROCHLORIC ACID CONTAINED IN THE CARRIER GAS ENTERING THE REACTOR MAY RESULT IN REMOVAL RATHER THAN GROWTH OF SILICON.

THE COMPETING REACTION²

$$S_1CL_4 + S_1(SOLID) = 2S_1CL_2$$
 2-2

SUGGESTS THAT "IF THE [SILANE] CONCENTRATION IS TOO HIGH, ETCHING OF THE SILICON WILL TAKE PLACE EVEN IN THE AB-SENCE OF A SIGNIFICANT CONCENTRATION OF HCL."³ SUCH ETCHING DESTROYS MONO-CRYSTALLINE SURFACE PERIODICITY DUE TO LOSS OF SURFACE ATOMS. HIGH TEMPERATURES DURING THE GROWTH PROCESS ENABLES NEIGHBORING ATOMS TO EASILY MIGRATE TO REDUCE THE ENERGY LEVEL AT STRESS POINTS CAUSED BY VACANCIES. THIS INDUCES EVEN MORE LATTICE DE-FORMATION WHICH MAY LEAVE THE SUBSTRATE WITH AN EXCESS OF HIGH ENERGY STATES AS A RESULT OF WEAK OR MISSING CO-ORDINATIVE BONDS NEAR THE SEMICONDUCTOR SURFACE. IN CONTRAST WITH RELATIVELY RIGID TETRAHEDRONS OF NEAR PER-FECT CRYSTALLINE SILICON, THE HIGH ENERGY STATES MAY FACILITATE HIGHER DIFFUSIVITY OF IMPURITIES INTO THE SUBSTRATE. THERMAL OXIDATION OF SILICON MAY PROCEED BY EITHER OF TWO REACTIONS DEPENDING ON THE NATURE OF THE OXIDATION PROCESS USED:⁴

DRY
$$S_{1(SOLID)} + O_2 \xrightarrow{S_1O_2(SOLID)} 2-3$$

WET
$$S_{1(SOLID)} + H_{20} = S_{10} (SOLID)^{+2H_{2}} - 4$$

IT HAS BEEN SHOWN THAT SOME OF THE SURFACE EPITAXIAL SILICON IS CONSUMED DURING OXIDATION:

"From the densities and molecular weights of silicon dioxide, a layer of silicon .45 x_0 thick is consumed [where x_0 is the net oxide thickness]. It has been demonstrated by the use of radioactive tracers that oxidation proceeds by an inward motion of the oxidizing species through the oxide layer rather than by the opposite process of the outward motion of silicon to the outer surface of the oxide."⁵

Accordingly, the oxidizing species must undergo three consecutive steps during oxidation: 6

1. THEY MUST BE TRANSPORTED FROM THE BULK OF THE GAS TO THE GAS-OXIDE INTERFACE;

2. THEY MUST DIFFUSE ACROSS THE OXIDE LAYER ALREADY PRESENT: AND

3. THEY MUST REACT AT THE SILICON SURFACE.

IT IS NOW CLEAR THAT OXIDE-MASKING STEPS DURING PRO-DUCTION OF MOS DEVICES INDUCE CRYSTAL STRESS WHEN THE OXIDIZING SPECIES MAKE THEIR WAY THROUGH EXISTING LAYERS OF OXIDE AND THE SEMICONDUCTOR BENEATH THEM. OTHER IONIC IMPURITIES SUCH AS SODIUM AND HYDROGEN MAY REACT WITH THE SILICON AND SILICON DIOXIDE AT THE SI-SIO₂ INTERFACE AND IN-TRODUCE A LARGE NUMBER OF STATES WITH ENERGY LEVELS WITHIN THE FORBIDDEN GAP. SUCH STATES ARE CALLED <u>SURFACE STATES</u> AND WILL BE SHOWN TO PLAY A VITAL PART IN NOISE PHENOMENA ASSOCIATED WITH MOS TRANSISTORS.

IT IS A WELL KNOWN FACT THAT "IN EQUILIBRIUM, THE CON-CENTRATION OF A SPECIES WITHIN A SOLID IS PROPORTIONAL TO THE PARTIAL PRESSURE OF THE SPECIES IN THE SURROUNDING GAS."⁷ EXPERIMENTAL DATA UNDER WIDELY RANGING CONDITIONS OF TEM-PERATURE, PARTIAL PRESSURE OF OXIDENTS, AND USING EITHER OXYGEN OR WATER VAPOR AS THE OXIDIZING SPECIES CONFIRM THAT FOR LARGE TIMES T, THE THICKNESS OF THE OXIDE FILM MAY BE EXPRESSED AS⁸

$$x_0^2 \simeq Z \cdot T$$
 2-5

HERE Z IS A PARABOLIC RATE CONSTANT FOR THE PROCESS AND IS PROPORTIONAL TO THE PARTIAL PRESSURE OF THE OXIDENT IN THE GAS AND RELATED TO BOTH THE DIFFUSIVITY AND CONCENTRATION OF THE OXIDIZING SPECIES IN THE OXIDE LAYER.

EXPERIMENTAL INVESTIGATIONS HAVE INDICATED THAT "FOR OXIDATION BOTH WITH OXYGEN AND WITH WATER VAPOR, THE OXI-DIZING SPECIES MOVING THROUGH THE OXIDE LAYER ARE APPARENTLY <u>MOLECULAR</u>."⁹ AT LEAST ONE THEORY OF 1/F NOISE¹⁰ MAKES USE OF A "SLOW RELAXATION" IN THE TRANSISTION FROM MOLECULAR

CHEMISORBED WATER TO LARGER CHEMISORBED COMPLEXES.*

<u>Properties of the SI-SIO2 interface</u>. The siliconsilicon dioxide interface in MOS transistors may be characterized by <u>fixed surface-state charge</u> in the oxide, <u>space charges</u> within the insulating layer due to <u>mobile</u> <u>impurity ions and traps</u> ionized thermally or by radiation, and <u>surface states</u>. Each of these phenomena are treated individually in the following discussions.

1. Fixed surface-state charge. It is theorized that the unexpected existence of a fixed surface-state charge, Q_{SS} , located at or near the SI-SIO₂ interface is attributed to excess unreacted ionic silicon present in the oxide during the oxidation process. The characteristics of this charge have been extensively investigated by Deal et al¹¹ and have been summarized by Grove:¹²

IT IS FIXED AND CAN BE NEITHER CHARGED NOR DIS-CHARGED OVER A WIDE VARIATION OF BENDING OF THE SILICON ENERGY BANDS.

IT IS NOT AFFECTED UNDER CONDITIONS LEADING TO SODIUM MIGRATION IN THE OXIDE.

IT IS LOCATED WITHIN APPROXIMATELY 200Å OF THE SI-SIO2 INTERFACE.

THE CHARGE DENSITY IS NOT AFFECTED BY EITHER THE OXIDE THICKNESS OR THE TYPE AND CONCENTRATION OF IMPURITY PRESENT.

*O. JANTSCH¹³ AND REVESZ¹⁴ HAVE SHOWN THAT THE "SLOW RELAXATION" RESULTS IN A "MODULATION OF SURFACE RECOMBINA-TION" WHICH PRODUCES DRAIN CURRENT FLUCTUATIONS IN MOSFETS. THE CHARGE IS A FUNCTION OF THE OXIDATION AND ANNEALING CONDITIONS AS WELL AS THE ORIENTA-TION OF THE SILICON CRYSTAL.

THE FIXED SURFACE-STATE CHARGE RESULTS IN A "PARA-LLEL TRANSLATION OF THE CAPACITANCE-VOLTAGE CHARACTER-ISTICS ALONG THE VOLTAGE AXIS.¹¹¹⁵ IT ALSO INCREASES THE EFFECTIVE GATE VOLTAGE BY¹⁶

$$V_{G}' = V_{G} + Q_{SS}/C_{O} \qquad 2-6$$

AND DECREASES THE FLAT-BAND VOLTAGE ACCORDING TO¹⁷

$$V_{FB} = \Phi_{MS} - Q_{SS} \left(\frac{x_0}{K_0 \epsilon_0} \right)$$
 2-7

where V_{FB} is due to the difference in work functions between the metal and the semiconductor material and Φ_{MS} is the metal-semiconductor work function difference with-out the influence of Q_{SS} .

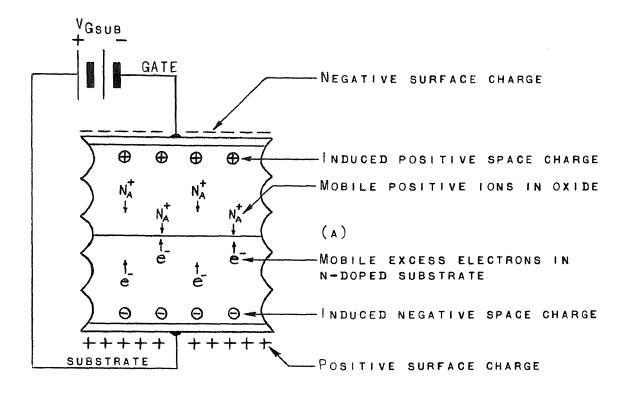
2. Space charges due to ionic contamination. A major problem with early MOS devices has been traced to sodium contamination and ionic drift under the influence of an electric field. This causes a "rearrangement of ionic space-charge distribution in the oxide."¹⁸

FIGURE 2-1 PRESENTS A MODEL OF THE SILICON-SILICON DIOXIDE INTERFACE WITH BIAS APPLIED AT ELEVATED TEMPERA-TURES. FIGURE 2-1A SHOWS THE APPLICATION OF A -V_{GSUB} AND THE INDUCED CHARGES AT THE METAL-GATE AND SUBSTRATE-METAL INTERFACES. THE DIAGRAM SHOWS POSITIVE SODIUM ATOMS, REPELLED BY THE INDUCED OXIDE CHARGE AND HAVING SUFFICIENT THERMAL ENERGY TO MIGRATE TOWARD THE SUBSTRATE, DRIFTING TOWARD THE SI-SIO₂ INTERFACE. THIS CAUSES MOBILE ELECTRONS IN THE N-DOPED SUBSTRATE TO DRIFT TOWARD THE OXIDE AND CHANNEL REGION IN P-CHANNEL ENHANCEMENT MOSFET'S.

FIGURE 2-1B ILLUSTRATES THE FOUR RESIDUAL SPACE CHARGES WHICH REMAIN IF THE DRIFT PROCESS IS ALLOWED TO CONTINUE OVER A LONG PERIOD OF TIME. THE RESULT IS THE SAME AS IF A POSITIVE VOLTAGE IS APPLIED TO THE GATE AS WAS SHOWN IN FIGURE 1-2B. THE CHANNEL APPEARS TO BE MORE HEAVILY N-DOPED THAN BEFORE IONIC DRIFT. THE NEGATIVE GATE VOLTAGE REQUIRED FOR SURFACE INVERSION, V_T , MUST NOW BE INCREASED, THEREBY ACCOUNTING FOR APPARENT NEGATIVE THRESH-OLD DRIFT.

3. Space charges due to ionized traps. Exposure of silicon dioxide films to x-ray, gamma ray, and low-and highenergy irradiation results in the formation of electron-hole pairs. When in the presence of an electric field, the electrons and holes separate.¹⁹

Consider the model in figure 2-1 again. If a $-V_{GS}$ is applied and, instead of heat the device is brought into contact with x-radiation, optically excited electrons drift toward the oxide and holes drift toward the semiconductor. Since the mobility of electrons in silicon dioxide is very



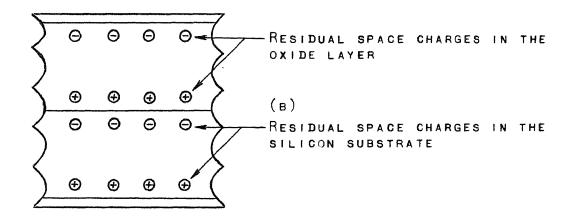


FIGURE 2-1. MODEL OF THE SI-SIO₂ INTERFACE IN A P-CHANNEL ENHANCEMENT MOSFET SHOWING THE RESULTS OF IONIC DRIFT IN THE OXIDE UNDER THE INFLUENCE OF AN ELECTRIC FIELD. (A) SODIUM IONS ARE ATTRACTED TOWARD THE SEMICONDUCTOR WHILE MOBILE ELECTRONS IN THE N-DOPED SUBSTRATE DRIFT TO-WARD THE OXIDE. (B) RESIDUAL SPACE-CHARGES IN THE OXIDE AND IN THE SEMICONDUCTOR MATERIAL. THE APPARANT HIGHER DOPING IN THE CHANNEL REGION CAUSES A NEGATIVE SHIFT IN THE THRESHOLD VOLTAGE, THE GATE VOLTAGE REQUIRED TO IN-VERT THE SEMICONDUCTOR SURFACE. LOW, MAJORITY CARRIERS IN THE N-DOPED SUBSTRATE CANNOT RECOMBINE IN THE BULK INSULATOR AT THE SI-SIO₂ INTERFACE. THEREFORE, THERE IS A BUILD-UP OF SPACE CHARGE IN THE OXIDE AS HOLES ARE TRAPPED NEXT TO THE SUBSTRATE.

Continued irradiation in the presence of an electric field causes a larger fraction of the total applied gate voltage to be dropped across the space charge region, thereby further increasing the electric field across the space-charge region due to the accumulation of holes. If the process is allowed to continue, the field in the rest of the oxide gradually approaches zero and a steady-state highly localized electric field condition is brought about in the oxide at the silicon-silicon dioxide interface.²⁰

4. Fast surface states. Figure 1-2 illustrated the energy band bending which takes place in the semiconductor near: the SI-SIO₂ interface when the gate is held at some positive or negative potential with respect to the substrate. Due to a disruption of the periodicity of the semiconductor surface during oxidation, there are many surface energy states within the forbidden gap. Consequently, variations in the energy bands relative to the Fermi level affect the "probability of occupation of those states."²¹

By DEFINITION, "THOSE STATES WHOSE CHARGE CAN BE READILY EXCHANGED WITH THE SEMICONDUCTOR ARE CALLED SURFACE

STATES.¹¹²² THE PRESENCE OR ABSENCE OF ELECTRONS IN SUCH STATES IS A FUNCTION OF THE SURFACE POTENTIAL WHICH IS ITSELF A FUNCTION OF THE APPLIED GATE VOLTAGE.^{*}

2.3 PREVIOUS EXPERIMENTAL INVESTIGATIONS

Revesz²³ has investigated the SI-SIO₂ interface in enhancement MOSFET's and has shown that the silicon surface state is a <u>donor</u>, that is, positively ionized silicon provides an excess of donor electrons. During oxidation in oxygen at relatively high temperatures it is theorized that O⁻ and O⁻⁻ consume electrons as they diffuse into the silicon substrate. Revesz has concluded that oxidation in water vapor involves significantly more surface states than does oxidation in oxygen alone and that the oxidation rate is increased while the density of free states is reduced.²⁴

IT IS WELL KNOWN THAT STATES HAVE A MUCH GREATER PROB-ABILITY OF EXISTING AT DISLOCATIONS AND BOUNDARY IMPERFEC-TIONS THAN ANYWHERE ELSE AND THAT THEY GIVE RISE TO THE FIXED SURFACE-STATE CHARGE PHENOMENON IN THE OXIDE. IT HAS BEEN THEORIZED THAT DONOR ELECTRONS PRESENT IN THIS MANNER ACT AS <u>RECOMBINATION CENTERS</u>, ENABLE THE DIFFUSION OF SUCH

^{*}SURFACE-STATE FLUCTUATIONS RESULTING IN VARIATIONS IN SURFACE POTENTIAL MODULATE CHANNEL CONDUCTANCE WHICH IN TURN CAUSES VARIATIONS IN DRAIN CURRENT. THESE DRAIN CUR-RENT FLUCTUATIONS MAY BE OBSERVED AS NOISE.

IMPURITIES AS HIGHLY MOBILE SODIUM ION, Na⁺, Hydrogen ION, H^+ , and other <u>acceptors</u> which may be present.²⁵

WHEN THE SEMICONDUCTOR SURFACE OF A P-CHANNEL ENHANCE-MENT MOSFET IS INVERTED, THE FREE HOLE CONCENTRATION IS VERY LARGE WHILE THE FREE ELECTRON DENSITY IS NEGLIGIBLE. SINCE IONIZED SILICON, SI⁺³ and SI⁺⁴, IS ALSO AN ACCEPTOR, IT HAS BEEN THEORIZED THAT THE SURFACE STATES IN P-CHANNEL DEVICES ARE ALMOST EXCLUSIVELY <u>FAST TRAPPING CENTERS</u> IN THE CONDUC-TION CHANNEL.^{*}

2.4 INVESTIGATION OF THRESHOLD VOLTAGE DRIFT

EXPERIMENTAL METHOD. IN ORDER TO QUANTITATIVELY IN-VESTIGATE THE MAGNITUDE OF THRESHOLD VOLTAGE DRIFT DUE TO IONIC CONTAMINATION IN THE INSULATING LAYER AND IN ORDER TO CORRELATE SUCH DRIFT WITH EXCESS NOISE AND OTHER SURFACE-RELATED PHENOMENA, LIFE-TESTING OF SEVERAL ENHANCEMENT MOSFET'S WAS UNDERTAKEN.

A LARGE NUMBER OF MEDIUM CONDUCTANCE P-CHANNEL FIELD-EFFECT TRANSISTORS WERE OBTAINED. ** THESE DEVICES WERE

*Such other chemisorbed species as H₂O⁺ may also be considered fast trapping centers. O. Jantsch claims they figure prominently in the production of excess noise in MOS transistors. See reference 1967-8 for more details.

** THE DEVICES WERE MANUFACTURED BY AMERICAN MICRO-SYS-TEMS, INC. OF SANTA CLARA, CALIFORNIA AND WERE PART OF JOINT RELIABILITY STUDIES UNDERTAKEN BY NEWARK COLLEGE OF ENGIN-EERING AND PICATINNY ARSENAL DURING THE YEARS 1968-72. PURPORTED TO BE DESIGNED PRIMARILY FOR "LINEAR WIDEBAND AMPLIFIERS AND HIGH SPEED SWITCHING AND COMMUTATING APPLI-CATIONS." THEY FEATURED A HIGH GAIN-BANDWIDTH PRODUCT, TYPICALLY 195MHZ, AND A THRESHOLD VOLTAGE RANGE OF FROM -3.5 TO -5.5V. Two CONFIGURATIONS AND THREE TYPES WERE AVAILABLE: TYPE DDO?P UNITS HAD A BUILT-IN ZENER DIODE VOLTAGE LIMITER BETWEEN THE GATE AND THE SUBSTRATE WHILE TYPES DDO8P AND DD08K HAD NO SUCH GATE PROTECTION. ALL WERE FOUR LEAD TRANSISTORS (ONE LEAD FOR THE DRAIN, GATE, SOURCE AND SUBSTRATE BULK) AND HERMETICALLY SEALED IN OR-DINARY TO-72 METAL CANS. THERE WERE NO SIGNIFICANT DIF-FERENCES IN THE CONSTRUCTION OF THE THREE TYPES OF MOSFET'S.

TWENTY DEVICES, TEN EACH OF TYPES DD07P AND DD08K, WERE SELECTED FOR LIFE-TESTING FOR THRESHOLD VOLTAGE DRIFT ON THE BASIS OF THEIR LOW FREQUENCY NOISE INTENSITY.* ACCORDINGLY, FIVE HIGH- AND FIVE LOW-NOISE UNITS WERE SELECTED FROM EACH OF THE TWO TYPES.

Teflon sockets were used to connect all transistors with common drain, source, and substrate. All type DD07P devices were biased with $V_{\rm GS}=-25V$. and all type DD08K units had -50V. on the gate. The complete assembly was wired inside an oven latter maintained at 130±5°C for 500 hours.

^{*}Noise measurements were made at 10Hz. and at 1KHz., both with a bandwidth of 1 cycle. The noise measurements are detailed in section 3.5.

EXPERIMENTAL RESULTS AND DISCUSSION. APPENDIX C LISTS THE PERCENTAGE NEGATIVE THRESHOLD VOLTAGE DRIFT, $\% \Delta V_T$, FOR ALL TWENTY UNITS TESTED. TYPE DD07P TRANSISTORS EXHIBITED RANGES OF DRIFT OF FROM 1.75 TO 2.66 PERCENT AND FROM 3.16 TO 7.15 PERCENT FOR LOW- AND HIGH-NOISE UNITS, RESPECTIVELY. RANGES OF 6.88 TO 9.11 AND 11.5 TO 23.6 PERCENT DRIFT WERE OBSERVED FOR THE LOW- AND HIGH+ NOISE TYPE DD08K UNITS.

It is apparant that threshold voltage drift is directly proportional to the electric field across the oxide, as evidenced by the higher percent drift for those units which were blased at the higher gate voltage. Statistical analysis confirmed that the type DD08K units exhibited substantial ΔV_T while drift in the DD07P transistors was very small.

OF PARTICULAR IMPORTANCE WAS THE FACT THAT HIGHER THRESHOLD VOLTAGE DRIFT WAS OBSERVED IN THE HIGH NOISE UNITS THAN IN THOSE WITH LOW NOISE INTENSITY. IT WAS THEREFORE CONCLUDED THAT THERE IS A DEFINITE CORRELATION BETWEEN 1/F NOISE AND THRESHOLD VOLTAGE DRIFT AND THAT, FOR P-CHANNEL ENHANCEMENT MOSFET'S, IONIC DRIFT IN THE OXIDE DEFINITELY AFFECTS THE SEMICONDUCTOR SURFACE.

^{*}Previous investigations showed that the type DD08K devices exhibited unusually high ΔV_T . It was therefore of particular importance that such high drift units also be included in this investigation.

^{**}Chapter 3 shows that 1/F noise is a surface-related phenomena and that ionic Drift increases the surface-state density.

APPENDIX E DESCRIBES THE PROCEDURES USED TO MEASURE THE SHORT-CIRCUIT DRAIN LEAKAGE CURRENT (WITH ZERO GATE BIAS) AND LISTS THE RESULTS OBTAINED BOTH BEFORE AND AFTER LIFE-TESTING FOR THRESHOLD VOLTAGE DRIFT. IT WAS FOUND THAT THE LEAKAGE CURRENT INCREASED BY A FACTOR OF FROM THREE TO FOUR BUT THAT THIS WAS INSUFFICIENT TO CONCLUDE THAT A SIGNIFICANT AMOUNT OF DRAIN JUNCTION DEGRADATION HAD TAKEN PLACE AFTER 500 HOURS OF BIAS AT 130±5°C.

Furthermore, since there was no significant difference between the amount by which I_{DSS} increased for units having either relatively high or relatively low noise, it was concluded that drain leakage was not a factor which results in the production of drain current fluctuations. Also, since the increase in leakage for transistors which showed appreciable threshold voltage drift was similar to that for units which exhibited only a small amount of drift, it was concluded that there is no correlation between I_{DSS} and threshold voltage drift and the mechanism of ionic drift at the Si-SiO₂ interface which produces the drift.

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20 IBID. PP. 341-2.

²¹<u>IBID</u>., P. 282. ²²<u>IBID</u>., P. 283. ²³<u>OP</u>. <u>Cit</u>., Revesz, p. 98. ²⁴<u>IBID</u>., P. 98. ²⁵<u>OP</u>. <u>Cit</u>., O. Jantsch, p. 268.

CHAPTER 3

THE SURFACE-STATE RELATED NOISE PHENOMENA

IN MOS TRANSISTORS

3.1 INTRODUCTION

IN THE ABSENCE OF EXTERNALLY APPLIED POWER, VIRTUALLY ALL RESISTIVE ELEMENTS MAY BE THOUGHT OF AS NOISE GENERA-TORS. HOWEVER, MOS TRANSISTORS ARE UNIQUE IN THAT THEY PRODUCE AN EXCEPTIONALLY HIGH INTENSITY OF <u>EXCESS NOISE</u> WHEN POWER IS APPLIED. THIS NOISE HAS BEEN SHOWN TO BE DIRECTLY PROPORTIONAL TO CURRENT DENSITY AND INVERSELY PROPORTIONAL TO THE FREQUENCY AT WHICH IT IS DETECTED.¹ FURTHERMORE, NUMEROUS INVESTIGATIONS HAVE SHOWN THAT THIS "1/F" NOISE IS STRONGLY DEPENDENT UPON THE GENERAL SURFACE CONDITIONS AT THE SILICON-SILICON DIOXIDE INTERFACE IN MOS TRANSISTORS AND ON THE SURFACE STATE DENSITY IN PARTICULAR.

THE FOLLOWING DISCUSSIONS PRESENT A BRIEF BACKGROUND ON THE MAJOR TYPES OF ELECTRICAL NOISE WITH PARTICULAR EM-PHASIS ON LOW FREQUENCY EXCESS NOISE INTENSITY AND ITS USE IN AN ANALYSIS OF COMPONENT RELIABILITY. EXPERIMENTAL RE-SULTS ARE THEN PRESENTED WHICH QUALITATIVELY CORRELATE SUR-FACE-STATE RELATED 1/F NOISE INTENSITY WITH THE SURFACE-STATE DENSITY, SEMICONDUCTOR IMPURITY CONCENTRATION, AND OBSERVED DIFFERENCES IN SATURATION DRAIN CONDUCTANCE BE-TWEEN HIGH- AND LOW-NOISE ENHANCEMENT MOSFET'S.

3.2 BASIC NOISE THEORY

ELECTRICAL NOISE MAY BE DEFINED AS SPURIOUS, UNWANTED SIGNALS WHICH ARE RANDOM IN NATURE AND WHICH TEND TO BE APERIODIC. NOISE MAY BE NAMED FOR ITS APPARANT CAUSE AND EFFECT (IE. THERMAL EXCITATION, RANDOM RECOMBINATION, OR RANDOM COLLISIONS) OR BY ITS FREQUENCY DOMAIN (EG. LOW-FREQUENCY NOISE). ACCORDINGLY, THERE ARE THREE MAJOR TYPES OF NOISE COMMONLY FOUND IN ELECTRICAL COMPONENTS, EACH OF WHICH IS DISCUSSED INDIVIDUALLY BELOW.

THERMAL (JOHNSON) NOISE. THE RANDOM VELOCITIES OF CARRIERS BETWEEN COLLISIONS CONSTITUTES CURRENTS WHICH CAUSE A FLUCTUATING VOLTAGE TO APPEAR ACROSS THE TERMINALS OF ANY RESISTANCE R. FOR MOS TRANSISTORS, THE NYQUIST RELATION²

$$\overline{V}^2 = 4kTR\Delta f \qquad 3-1$$

EXPRESSES THE MEAN-SQUARED NOISE VOLTAGE IN TERMS OF THE TOTAL DRAIN-TO-SOURCE RESISTANCE, THE ABSOLUTE TEMPERATURE, AND THE BANDWIDTH IN CYCLES PER SECOND.

THE SPECTRAL INTENSITY OF THERMAL NOISE IS "WHITE," THAT IS, ITS MAGNITUDE IS PRACTICALLY UNIFORM WITH RESPECT TO FREQUENCY. CONSEQUENTLY, THE LARGER THE INCLUDED BAND-WIDTH, THE GREATER WILL BE THE INTENSITY OF NOISE MEASURED. ALL PRACTICAL ELECTRONIC COMPONENTS POSSESS A FINITE RESIS-TANCE WHICH CONTRIBUTES THERMAL NOISE. SHOT NOISE. DIRECT CURRENTS IN SEMICONDUCTOR MATER-IALS ARE BY WAY OF A CONTINUOUS PROCESS OF DISSOCIATION AND RECOMBINATION IN A SPASMODIC "JERKY" FASHION. THIS RESULTS IN INSTANTANEOUS CURRENT FLUCTUATIONS, THE MEAN-SQUARED VALUE OF WHICH IS GIVEN BY³

$$\overline{i}^2 = 29 \mathrm{I}\Delta f \qquad 3-2$$

where I is the direct current measured, \mathscr{G} is the carrier charge, and Δf is the included bandwidth.

Previous investigations^{4,5,6} have shown that shot noise is essentially white for low frequencies. It may be considered independent of frequency over the range of interest for MOSFET operation.

Excess (1/F. MODULATION, FLICKER) NOISE. THE ORIGIN OF EXCESS NOISE IS "THOUGHT TO LIE IN SEMICONDUCTOR CRYSTAL IMPERFECTIONS AND SURFACE EFFECTS."⁷ By passivating silicon Devices with silicon dioxide to reduce the role of surface states, defects are necessarily introduced at the SI-SIO₂ interface. The presence of fast surface states at the semiconductor surface leads to a fluctuation of surface charge in MOS transistors and may produce surface-state related noise which may be observed in Either one or both the gate and the drain circuit in MOSFET's. Bell⁸ has proposed two mechanisms which result in observed excess noise in MOS devices: 1. <u>Gate noise</u>^{*} due to fluctuations in total surface charge as carriers move to and from the oxide, and

2. <u>Drain current fluctuations</u>** resulting from variations in channel conductance as a result of variations in V_{GS} due to the charge fluctuations which produce gate noise.

Bell offers the following generalized expression for frequency-dependent noise:⁹

$$\frac{V^2B}{B} = K(J^{\theta}R^{\eta}v^{\lambda}s/f^{\eta})B \qquad 3-3$$

where B = BANDWIDTH in cycles per second

J = CURRENT DENSITY K = MATERIAL CONSTANT R = RESISTANCE IN OHMS S = SHAPE FACTOR V = VOLUME 7 + 8 = 2 to balance the equation in time B = 2 to balance the equation in charge

*For any ΔQ into the oxide from the semiconductor, there must be a ΔQ of opposite polarity onto the gate and vice versa. This leads to gate charge fluctuation which is observed as a noise voltage in the gate circuit.

^{**}Variations of charge between the Depletion Region, the inversion region, and the surface states alters the instantaneous value of $V_{\rm GS}$. Recall from section 1.2 that the channel conductance, and therefore the Drain current, is easily modulated by variations in gate voltage. Hence the observed drain current fluctuations in MOS transistors. THE GENERAL EXPRESSION MAY BE SIMPLIFIED TO¹⁰

$$\frac{V^{2}B}{R} = K(J^{2}R^{2-\eta}v^{1-\varkappa}s/f)B \qquad 3-4$$

where $x = (\gamma - 1)/3$. For an exact 1/F spectrum, let $\gamma = 1$, 1-x = 1, 2- $\eta = 1$, and let s be an area/unit length. Then The product R.s is in units of resistivity, ρ .

The noise power density for a 1/F spectral intensity (in which B is one cycle at an angular frequency of ω) may be stated as¹¹

$$P_{N(\omega)} = K(J^2 q V/f) d\omega \qquad 3-5$$

WHICH IS DIRECTLY PROPORTIONAL TO THE DC POWER APPLIED (AND THE CURRENT DENSITY SQUARED) AND INVERSELY PROPOR-TIONAL TO FREQUENCY.*

3.3 PREVIOUS STUDIES OF THE EXCESS NOISE SPECTRAL INTENSITY

SURFACE-STATE RELATED 1/F NOISE HAS BEEN ATTRIBUTED TO A MODULATION OF SURFACE RECOMBINATION. THIS REASONING IS BASED ON McWhorter's¹² theory of trapping centers at the SILICON-SILICON DIOXIDE INTERFACE:

1. THE INTRODUCTION OF TRAPPING CENTERS INTO THE SEMICONDUCTOR MATERIAL RESULTS IN A "DIRECT CHANGE IN CURRENT DUE TO A CHANGE IN THE NUMBER OF FREE CARRIERS AND AN INDIRECT CHANGE IN CURRENT DUE TO A CHANGE IN THE NUMBER OF TRAPPED CARRIERS."13

*1/F NOISE IS USUALLY MASKED BY THERMAL NOISE AT ALL FREQUENCIES ABOVE 20KHZ AND THEREFORE MUST BE OBSERVED AT RELATIVELY LOW FREQUENCIES FOR TRUE 1/F DEPENDENCE. 2. THE RANDOM OCCUPANCY OF SURFACE STATES BY BROWNIAN MOTION "MODULATES THE SURFACE POTEN-TIAL, WHICH CAUSES A <u>FLUCTUATION IN SURFACE</u> <u>RECOMBINATION VELOCITY</u> AND A <u>MODULATION OF THE</u> <u>FREE CARRIER DENSITY</u> IN THE SPACE-CHARGE REGION NEAR THE SILICON SURFACE."¹⁴

3. SINCE THE DRIFT OF CHARGED CARRIERS CONSTI-TUTES A DIFFUSION CURRENT, ANY FLUCTUATION IN THE FREE CARRIER DENSITY NECESSARILY INDICATES A "FLUCTUATION IN DIFFUSION CURRENT."¹⁵

THE 1/F NOISE CAN THEN BE EXPLAINED BY "CONDUCTION CHANNEL CHARGE DENSITY FLUCTUATION CAUSED BY THE MODULA-TION OF SURFACE POTENTIAL DUE TO THE RANDOM OCCUPANCY OF SURFACE STATES"¹⁶ which supports the theory of Drain CURRENT FLUCTUATIONS PUT FORTH BY BELL.

Since the Life-time of the fast surface states is typically about 10⁻⁹ second, the recombination noise produced is of very high frequency. The observed low-frequency noise must therefore be a <u>modulation of surface</u> Recombination.¹⁷

Heiman and Warfield¹⁸ have estimated that surface states are located within about 20Å from the the $SI-SIO_2$ interface. Since the oxide thickness in MOS transistors is typically 600-2000Å (and since the oxide thickness must be orders of magnitude greater than the inversion layer), the resident distance of surface states is negligibly small and the <u>1/F</u> noise may be considered a surface phenomenon. Hsu¹⁹ has analytically predicted the noise spectral intensity of low-frequency excess noise in p-channel enhancement MOSFET's and a review of his derivation is included in Appendix B. The <u>spectral intensity of the short</u>-<u>circuit drain current fluctuations</u> is given by²⁰

$$S_{i(\omega)} = \left[\frac{g G_{M}}{A C_{0}}\right]^{2} \frac{4 k T A N_{T}(E) \gamma_{s}}{\alpha} (e^{\alpha \ell} - 1) \qquad 3-6 A$$
FOR $\omega \gamma_{s} e^{\alpha \ell} \ll 1$

$$S_{i}(\omega) = \left[\frac{g^{G}G}{AC_{0}}\right]^{2} \frac{kTAN_{T}(E)}{\alpha f} \qquad 3-6B$$
FOR $\omega \gamma_{s} \ll 1 \ll \omega \gamma_{s} e^{\alpha \ell}$

$$S_{i(\omega)} = \left[\frac{\varphi G_{M}}{AC_{0}}\right]^{2} \frac{4kTAN_{T}(E)}{\alpha \omega^{2} \gamma_{s}} (1 - e^{-\alpha \ell}) \qquad 3-6c$$
FOR $1 \ll \omega \gamma_{s}$

where A = Effective area of the gate in cm² $C_{O} = \frac{K_{O} \epsilon_{O}}{x_{O}} \text{ is the oxide capacitance}$ $G_{M} = \frac{\partial I_{D}}{\partial V_{G}} | \text{ is the MOST transconductance}$ $K = B_{O} \text{ transformed} \text{ for } V_{D} = \text{constant}$ $k = B_{O} \text{ transformed} \text{ to be 20Å}$ $N_{T}(E) = D_{E} \text{ nsity of trapped surface states per unit}$ $K = A_{B} \text{ solute temperature in } K$ $\frac{2}{s} = \text{Time constant of surface states located at}$ $\omega = A_{N} \text{ sullar frequency in Rad/sec.}$

AND WHERE α is experimentally determined to be 2×10⁸ cm⁻¹ and exp($\alpha \ell$) = 10¹⁷.²¹

From the above expressions it is apparant that very Low frequency noise is independent of frequency, a true 1/F spectral intensity is found only within the range of $(2\pi\gamma_{\rm s}e^{\alpha\ell})^{-1}$ to $(2\pi\gamma_{\rm s})^{-1}$ (a range theoretically predicted to be about 17 orders of magnitude), and that noise at higher frequencies has a $1/F^2$ dependence.

OF PARTICULAR INTEREST IS THE DEPENDENCE OF NOISE INTENSITY ON THE SURFACE-STATE DENSITY. SINCE ONLY THOSE STATES "WITHIN $\pm 2kT$ from the surface-state Fermi level contribute noise to the device, if the surface states are arbitrarily distributed in energy, the value $N_T(E)$ represents the surface-state density within $\pm 2kT$ of the Fermi level."²² Processes which introduce states into this range of energy therefore result in an increase in the number of surface states and may be observed by an increase in noise intensity.

3.4 ELECTRICAL NOISE AND COMPONENT RELIABILITY

AN EXCEPTIONALLY HIGH NOISE INTENSITY MAY BE REGARDED AS AN INDICATION OF SOME ABNORMALITY IN AN ELECTRICAL DEVICE. THOSE DEVICES EXHIBITING HIGH NOISE CHARACTERISTICS MAY THEREFORE BE CLASSIFIED AS POTENTIAL <u>RELIABILITY RISKS</u>. NOISE MEASUREMENTS ARE THEREFORE A <u>RELIABILITY TOOL</u> USED TO COMPARE DEVICES WITH AN ACCEPTED STANDARD.

ANY MECHANISM WHICH "FOCUSES" OR "FUNNELS" A SIGNIFI-CANT AMOUNT OF CURRENT THROUGH A SMALL CROSS-SECTION RESULTS IN A HIGH LOCALIZED CURRENT DENSITY AND ACCOMPANYING LOCAL-IZED HEATING IN A DEVICE. DEFECTS IN THE CRYSTAL STRUCTURE AT P-N JUNCTIONS, CRACKS, FLAWS, OR INCLUSIONS IN THE SEMI-CONDUCTOR MATERIAL MAY LEAD TO "CURRENT FUNNELING" AND HIGH 1/F NOISE INTENSITY.²³

SINCE ENHANCEMENT MOSFET'S CONDUCT THROUGH A RELATIVELY NARROW INVERSION REGION AT THE SEMICONDUCTOR SURFACE, ALL DRAIN CURRENT IS CONCENTRATED WITHIN AN EXTREMELY LIMITED PORTION OF THE CHANNEL. AT EVEN MODERATE CURRENT LEVELS, SUCH DEVICES POSSESS AN INHERENTLY HIGH CURRENT DENSITY NEAR THE SI-SIO₂ INTERFACE (WHERE INCLUSIONS SUCH AS SIO₂ MOLECULES ARE PRESENT) AND CONSEQUENTLY EXHIBIT MUCH HIGHER 1/F NOISE INTENSITY THAN BIPOLAR TRANSISTORS.

THEREFORE, UNUSUALLY HIGH NOISE UNITS MAY BE ELIMINATED FROM PRODUCTION BATCHES IN ANTICIPATION OF A CORRESPONDINGLY HIGH FAILURE RATE PROBABILITY DUE TO HIGH CURRENT DENSITY. FURTHERMORE, THE RELATIVE 1/F NOISE INTENSITY SHOULD BE AN EXCELLENT INDICATOR OF SURFACE CONDITIONS PRESENT IN EACH DEVICE AND THE SURFACE-STATE DENSITY IN PARTICULAR. THIS REASONING WAS APPLIED TO ALL ANALYISES OF NOISE PHENEMENA ENCOUNTERED IN THIS RESEARCH AND WAS ONE OF THE REASONS IT WAS THOUGHT THAT NOISE MEASUREMENTS WOULD BE BENEFICIAL IN A STUDY OF BREAKDOWN IN ENHANCEMNENT MOSFET'S.

3.5 NOISE MEASUREMENTS UNDERTAKEN

DESCRIPTION OF EQUIPMENT. NOISE MEASUREMENTS WERE PERFORMED USING THE QUAN-TECH MODEL 327 DIODE NOISE ANALYZER. FIGURE 3-1 SHOWS THE INTERNAL CONNECTION BLOCK DIAGRAM AND EXTERNAL GATE SUPPLY OF SUCH AN INSTRUMENT MODIFIED FOR USE WITH P-CHANNEL ENHANCEMENT MOSFET'S.

THE DEVICE CONSISTED OF A BUILT-IN VARIABLE DC SUPPLY, A SHIELDED TEST JIG AND SOCKET, CURRENT-LIMITING RESISTORS, A BROADBAND AMPLIFIER FOR AMPLIFYING THE NOISE VOLTAGE ACROSS A SINGLE, SELECTED, FIXED RESISTOR, TWO ACTIVE FIL-TERS RESONANT AT TEN HERTZ AND ONE KILOHERTZ WITH A BAND-WIDTH OF ONE CYCLE EACH, AND ANALOG OUTPUTS WHICH READ AT THE APPROPRIATE FREQUENCY IN RMS VOLTS.

NICKEL-CADMIUM BATTERIES WERE USED TO BIAS THE GATE OF EACH MOSFET UNDER TEST IN ORDER TO MINIMIZE NOISE INTRO-DUCED INTO THE GATE CIRCUIT WHICH MIGHT IN TURN CAUSE UN-DESIREABLE DRAIN CURRENT FLUCTUATIONS.

ALL NOISE MEASUREMENTS WERE CONDUCTED WITHIN A SHIELDED ROOM AT A TEMPERATURE OF ABOUT 25°C AND NOT MORE THAN 30% HUMIDITY.

EXPERIMENTAL METHOD. THE P-CHANNEL MOSFET'S UNDER TEST WERE INDIVIDUALLY CONNECTED WITH SOURCE AND SUBSTRATE TO THE MAIN POWER SUPPLY POSITIVE TERMINAL, DRAIN TO THE NEGATIVE TERMINAL, AND GATE TO THE EXTERNAL VARIABLE SUPPLY.

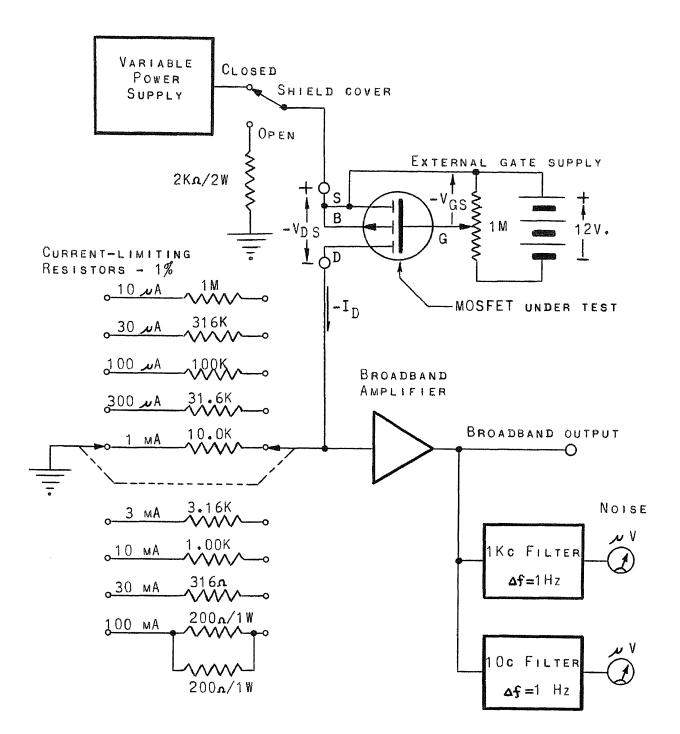


FIGURE 3-1. CONNECTION BLOCK DIAGRAM OF THE QUAN-TECH Model 327 Diode Noise Analyzer modified for measuring Noise voltages of p-channel enhancement MOS transistors.²⁴

THE MAIN DC SUPPLY AND THE CURRENT-LIMITING RESISTORS PROVIDED CONTINUOUS VARIATION OF EITHER $-V_{DS}$ or $-I_D$. With use of the external gate supply, noise voltages^{*} could be obtained anywhere on the MOSFET I-V characteristics.

APPROXIMATELY 80 TRANSISTORS WERE SELECTED FROM A LARGE BATCH OF P-CHANNEL MOSFET'S.^{**} All devices were FIRST SCREENED FOR PROPER OPERATION UP TO AND SLIGHTLY BE-YOND THE THRESHOLD OF DRAIN BREAKDOWN. THEN, 25 TRANSIS-TORS OF EACH TYPE WERE SELECTED FOR ACTUAL NOISE MEASURE-MENTS.

Noise voltages were obtained at 10Hz and at 1KHz with all units biased at a constant $V_{\rm GS}$ of -5.5 volts and four different values of drain voltage: -3V., -10V., -20V., and at -25V. This permitted noise voltages to be made in both the linear and saturation regions of operation for each unit. The data provided a profile of the mean noise voltage and standard deviation as a function of drain voltage for each type of transistor.

**Section 2.4 described the types DD07P, DD08P, and DD08K P-CHANNEL ENHANCEMENT MOSFET'S USED IN THIS RESEARCH. THOSE DEVICES ON WHICH NOISE MEASUREMENTS WERE MADE WERE SELECTED AT RANDOM FROM A BATCH ESTIMATED TO BE OVER 300 UNITS.

^{*}Actual Drain current fluctuations were computed from the measured noise voltages by dividing the observed quantity by the known, fixed resistance across which it appeared. The noise intensity (of drain current fluctuations) was then the square of the calculated noise current at a specified operating point.

Using the built-in voltmeter and ammeter on the Quan-Tech noise measuring instrument, saturation drain conductance measurements were made at an operating point of -5.5V. on the gate and $V_{\rm DS} = -20V$. The relative noise voltage at 10Hz was correlated with the drain conductance for each individual transistor.

On the basis of the relative noise intensity at ten cycles, the five highest and five lowest noise transistors of types DD07P and DD08K were selected for life-testing for threshold voltage drift. Noise voltage versus drain current measurements at constant gate voltage were made before and after 500 hours of bias^{*} at $130 \pm 5^{\circ}$ C.

Noise voltage versus drain current characteristics at several gate voltages were plotted for three high- and three low-noise transistors of type DD07P. A curve tracer was used to visually verify the plotted I-V characteristics.

All noise, current, and voltage measurements were checked and rechecked to ensure accurate readings. A <u>Fluke Model</u> <u>8000A Digital Voltmeter</u> was used to determine the precise threshold voltage for each transistor as well as for purposes of double checking the voltage readings on the Quan-Tech instrument.

^{*}Type DD07P units were biased at $V_{\rm GS}=-25V.$ while type DD08K units were biased at $V_{\rm GS}=-50V.$ The drain, source, and substrate bulk were common for both types of devices.

3.6 EXPERIMENTAL RESULTS AND DISCUSSION

<u>COMPARISON OF NOISE INTENSITIES</u>. FIGURE 3-2 REPRE-SENTS THE NOISE VOLTAGE PROFILE OF THREE SETS OF DEVICES AS A FUNCTION OF DRAIN VOLTAGE BEFORE LIFE-TESTING FOR THRESHOLD VOLTAGE DRIFT. FIGURE 3-2A SHOWS THE MEAN AND STANDARD DEVIATION OF NOISE VOLTAGES AT TEN CYCLES WHILE FIGURE 3-2B GIVES THE CORRESPONDING INFORMATION AT ONE-THOUSAND CYCLES.^{*}

IT WAS NOT THE INTENTION OF THIS RESEARCH TO ESTAB-LISH THE EXACT FREQUENCY DEPENDENCE OF EXCESS NOISE IN MOS TRANSISTORS. THIS HAS BEEN DONE BEFORE AND IS READILY AVAILABLE IN THE LITERATURE.^{25,26,27} Rather, it may be concluded that the low frequency excess noise intensity was found to be approximately inversely proportional to the frequency at which it was measured <u>before</u> threshold voltage drift. The 1/F noise theory developed earlier therefore applies to all subsequent discussions regarding the noise intensity of the devices tested.

^{*}The figures represent RMS noise voltage in microvolts at two frequencies. Comparisons of mean noise intensities between any two types of devices or between two frequencies for the same device must reflect the ratio of the squares of the noise voltage readings used. Example: Compare the mean noise intensity of type DD08P at 10Hz and VDS =-3V. and at the same drain voltage and 1KHz. In the first case the mean noise voltage was $10 \times 10^{-6}V$. And in the second it was approximately $0.82 \times 10^{-6}V$. The ratio of the square of the first reading to that of the second shows that the noise intensity at ten cycles was in fact more that 100 times that at one kilohertz.

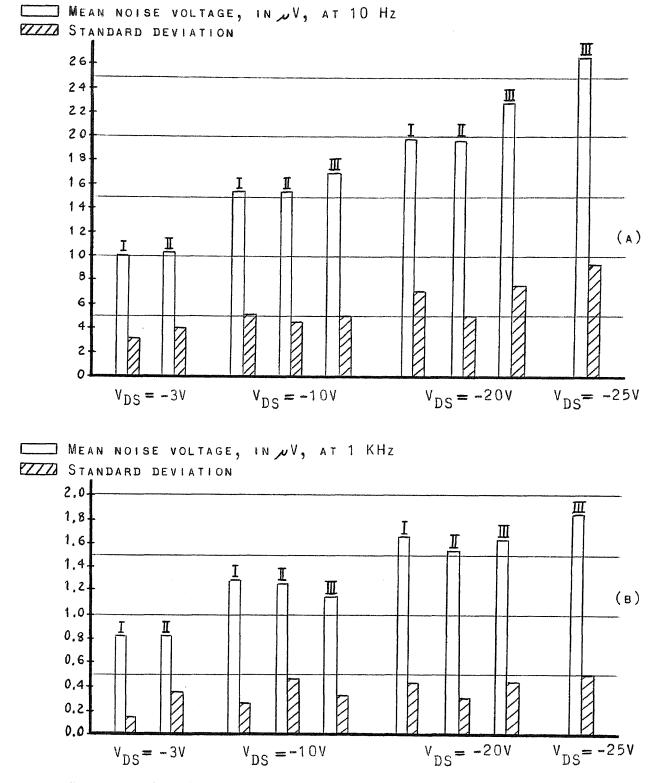


FIGURE 3-2. COMPARISON OF MEAN EXCESS NOISE VOLTAGE FOR THREE TYPES OF P-CHANNEL ENHANCEMENT MOSFET'S. MEASUREMENTS WERE MADE AT FOUR DRAIN VOLTAGES WITH CONSTANT VGS OF -5.5V. TYPEI: DD08P TYPEII: DD08K TYPE III: DD07P

As expected, the noise voltage was found to increase with increasing drain voltage when the gate was held at a constant potential. Increasing V_{DS} in this manner also increases the drain current which flows in a channel of fairly constant dimensions. The approximate 1/F frequency dependence was found to hold for both the linear and saturation regions of operation.

FIGURES 3-3 THROUGH 3-5 COMPARE MEAN NOISE VOLTAGE at 10Hz with saturation drain conductance at $V_{DS} = -10V$. and $V_{GS} = -5.5V$. For the three types of units tested. Two observations are immediately apparant:

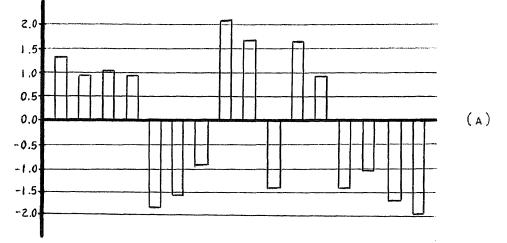
1. THERE WAS A WIDE RANGE OF RELATIVE NOISE INTENSITY WITHIN EACH GROUP OF MOSFET'S, AND

2. IN EVERY CASE, THOSE UNITS HAVING A HIGH NOISE IN-TENSITY SHOWED A CORRESPONDINGLY HIGHER SATURATION DRAIN CONDUCTANCE THAN DID THE LOW-NOISE UNITS.

Two THEORIES ARE OFFERED TO EXPLAIN THE CORRELATION BETWEEN 1/F NOISE INTENSITY AND SATURATION DRAIN CONDUC-TANCE.

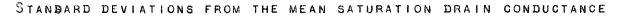
1. VARIATIONS OF SEMICONDUCTOR IMPURITY CONCENTRATION.

Recall from section 1.2 that high saturation drain conductance is synonymous with relatively thick oxide layers and/or low substrate impurity concentrations. Under the condition of <u>uniform oxidation</u> with <u>low surface-state den-</u>



STANDARD DEVIATIONS FROM THE MEAN NOISE VOLTAGE AT 10 Hz

UNIT NO. 'S: 45 47 48 54 56 57 59 60 61 70 72 75 77 78 79 80



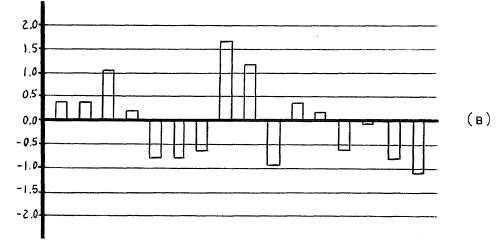
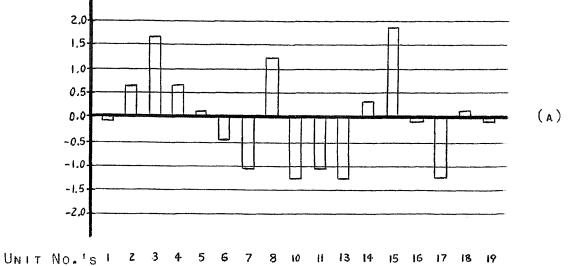


Figure 3-3. Comparison of noise voltage at 10Hz with saturation drain conductance for selected units of type DD07P MOSFET's. VDS = -20volts and VGS = -5.5volts.



STANDARD DEVIATIONS FROM THE MEAN NOISE VOLTAGE AT 10 Hz

STANDARD DEVIATIONS FROM THE MEAN SATURATION DRAIN CONDUCTANCE

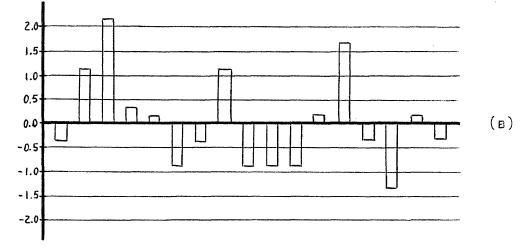
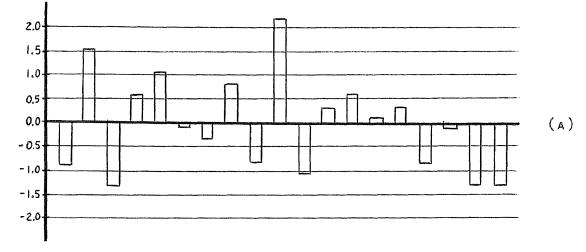
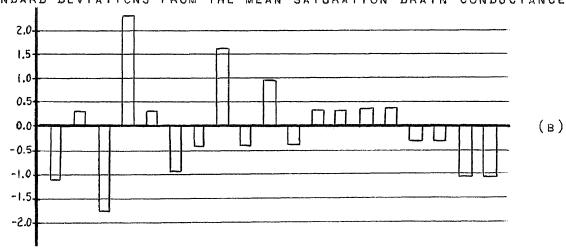


Figure 3-4. Comparison of noise voltage at 10Hz with saturation drain conductance for selected units of type DD08P MOSFET's. $V_{DS} = -20$ volts and $V_{GS} = -5.5$ volts.



STANDARD DEVIATIONS FROM THE MEAN NOISE VOLTAGE AT 10 Hz

UNIT NO. 'S 101 102 103 104 105 106 107 109 110 111 112 113 114 115 116 118 119 120 121



STANDARD DEVIATIONS FROM THE MEAN SATURATION DRAIN CONDUCTANCE

Figure 3-5. Comparison of noise voltage at 10Hz with saturation drain conductance for selected units of type DD08K MOSFET's. $V_{DS} = -20$ volts and $V_{GS} = -5.5$ volts.

SITY, THE DRAIN CONDUCTANCE IS SIMPLY A FUNCTION OF THE RELATIVE BULK DOPING CONCENTRATION. IT THEREFORE FOLLOWS THAT, WITH CONSTANT GATE VOLTAGE AND OPERATION IN THE SATURATION REGION AT CONSTANT DRAIN VOLTAGE, A RELATIVELY HIGH DRAIN CONDUCTANCE IS INDICATIVE OF A LOW SEMICONDUC-TOR IMPURITY CONCENTRATION.

SINCE THE DRAIN CURRENT IN THE SATURATION REGION IS DEPENDENT UPON THE AMOUNT BY WHICH THE EXPOP RECEDES TO-WARD THE SOURCE (WHICH IS ITSELF A FUNCTION OF TH BULK DOPING CONCENTRATION), THE RELATIVE CURRENT DENSITY IN THE TRANSISTOR IS DIRECTLY RELATED TO BOTH THE SEMICON-DUCTOR IMPURITY CONCENTRATION AND THE SATURATION DRAIN CONDUCTANCE.

THEREFORE, A HIGH CURRENT DENSITY, GIVING A HIGH 1/F NOISE INTENSITY, IS A DIRECT CONSEQUENCE OF RELATIVELY LOW BULK DOPING IN THE SUBSTRATE. NOISE MEASUREMENTS ARE THEREFORE USEFUL IN AN ANALYSIS OF DOPING PROFILES FOR MOS TRANSISTORS. HOWEVER, THE VARIATION IN NOISE INTEN-SITY MAY BE EXPECTED TO BE RELATIVELY SMALL DUE TO THE SMALL VARIATION SEMICONDUCTOR DOPING LEVELS BETWEEN UNITS OF THE SAME BATCH.

2. HIGH SURFACE-STATE DENSITY. RECALL FROM SECTION 3.3 THAT 1/F NOISE INTENSITY IS DIRECTLY PROPORTIONAL TO THE SURFACE-STATE DENSITY WITHIN AN MOS DEVICE. IT THERE-FORE FOLLOWS THAT, WITH <u>CONSTANT</u> <u>SEMICONDUCTOR</u> <u>IMPURITY</u>

CONCENTRATION AND OPERATING POINT, THOSE UNITS HAVING A RELATIVELY HIGH 1/F NOISE INTENSITY MAY ALSO BE EXPECTED TO POSSESS A HIGH DENSITY OF SURFACE STATES.

MEASUREMENTS OF IDSS CONFIRMED THAT HIGH DRAIN-DIODE LEAKAGE WAS NOT A FACTOR IN THOSE UNITS HAVING AN ABNOR-MALLY HIGH DRAIN CONDUCTANCE AND THEREFORE SUCH CONDUC-TANCE MUST NECESSARILY BE A SEMICONDUCTOR SURFACE EFFECT. IT IS THEN REASONABLE TO SUSPECT THAT SURFACE STATES CON-TRIBUTE TO SURFACE CONDUCTION AND THAT SURFACE-STATE RE-LATED 1/F NOISE INTENSITY IS THEREFORE HIGHER IN THOSE TRANSISTORS HAVING HIGH SATURATION DRAIN CONDUCTANCE.

SECTION 2.4 PROPOSED THAT THE MAJORITY OF FASE SUR-FACE STATES IN P-CHANNEL ENHANCEMENT MOSFET'S ARE TRAPPING CENTERS, THAT IS HOLES NEAR THE SI-SIO₂ INTERFACE. SINCE THE OXIDIZING SPECIES CONSUME ELECTRONS DURING DIFFUSION INTO THE SILICON, IT IS PLAUSIBLE TO EXPECT THAT A LOW SEMICONDUCTOR IMPURITY CONCENTRATION RESULTS IN A GREAT MANY MORE TRAPPING CENTERS THAN WOULD BE THE CASE WITH HIGH BULK DOPING. THE DENSITY OF FAST TRAPPING CENTERS IN P-CHANNEL DEVICES WOULD THEN BE AN INVERSE FUNCTION OF THE IMPURITY CONCENTRATION IN THE SUBSTRATE. THEREFORE, BOTH THEORIES 1 AND 2 ARE RELATED AND THE RELATIVE 1/F NOISE INTENSITY MAY BE AN INDICATION OF BOTH THE DOPING LEVEL AND THE SURFACE-STATE DENSITY. FURTHERMORE, THE SATURATION DRAIN CONDUCTANCE MIGHT BE PROPORTIONAL TO

THE SURFACE-STATE DENSITY AS WELL.

NOISE VOLTAGE AND THRESHOLD VOLTAGE DRIFT. THE NOISE VOLTAGES VERSUS DRAIN CURRENT AT CONSTANT GATE VOLTAGE FOR THE TWENTY UNITS LIFE-TESTED FOR THRESHOLD VOLTAGE DRIFT ARE LISTED IN APPENDIX D. ALL UNITS INCREASED IN NOISE INTENSITY FOLLOWING 500 HOURS OF BIAS AT AN ELEVATED TEM-PERATURE. OF PARTICULAR IMPORTANCE IS THE APPARANT LOSS OF 1/F FREQUENCY DEPENDENCE DUE TO UNUSUALLY HIGH INCREASES IN NOISE VOLTAGES AT 1KHZ. IT WAS CONCLUDED THAT IONIC DRIFT AT THE SILICON-SILICON DIOXIDE INTERFACE PROBABLY RE-SULTED IN A HIGH DENSITY OF <u>SHALLOW TRAPS</u> AT THE SEMICON-DUCTOR SURFACE. IT IS BELIEVED THAT SUCH TRAPS HAVE TIME CONSTANTS WHICH ARE MUCH SHORTER THAN THE TIME CONSTANTS OF OTHER TRAPPING CENTERS AND THAT MODULATION OF SURFACE RECOMBINATION UTILIZING SHALLOW TRAPS ACCOUNTS FOR THE GREATLY INCREASED NOISE INTENSITY AT THE HIGHER FREQUENCY.

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¹³S. T. HSU, "SURFACE-STATE RELATED 1/F NOISE IN MOS TRANSISTORS," <u>Solid-State Electronics</u>(GB) (1970), p. 1452. Ref. 1970-4.

¹⁴<u>Івір</u>., р. 1452. ¹⁵<u>Івір</u>., р. 1452. ¹⁶<u>Івір</u>., р. 1452.

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CHAPTER 4

CORRELATION OF 1/F NOISE WITH DRAIN BREAKDOWN

4.1 INTRODUCTION

DRAIN BREAKDOWN IN MOS TRANSISTORS IS SHOWN TO BE A COMPOSITE OF IMPACT IONIZATION IN THE CHANNEL AND GATE-CONTROLLED DRAIN-DIODE BREAKDOWN AT THE DRAIN-TO-SUBSTRATE P-N JUNCTION. THE SEMICONDUCTOR IMPURITY CONCENTRATION, THE SURFACE-STATE DENSITY, THE SATURATION DRAIN CONDUCTANCE, AND THE APPLIED GATE VOLTAGE ARE CRITICAL FACTORS INFLUENC-ING WHEN AND TO WHAT EXTENT DRAIN BREAKDOWN OCCURS AT HIGH DRAIN VOLTAGES. LOW FREQUENCY EXCESS NOISE MEASUREMENTS WERE CONDUCTED ON SEVERAL HIGH- AND LOW-NOISE TRANSISTORS AT VARIOUS STAGES OF BREAKDOWN AND THE EXPERIMENTAL RESULTS WERE CORRELATED WITH SEMICONDUCTOR SURFACE EFFECTS.

4.2 THEORY OF ELECTRICAL BREAKDOWN AT P-N JUNCTIONS

THE CUMULATIVE DRAIN BREAKDOWN PHENOMENA IN ENHANCEMENT MOSFET'S IS ESSENTIALLY AN EXTENSION OF BREAKDOWN THEORY AT ANY REVERSE-BLASED P-N JUNCTION. TWO TYPES OF BREAKDOWN IN SEMICONDUCTOR DEVICES ARE DISCUSSED BELOW.

<u>Avalanche breakdown</u>. Under the influence of a high electric field, ionization of semiconductor materials results in the creation of electron-hole pairs. When the electrons gain sufficient kinetic energy while being accelerated in the electric field, their collisions with the CRYSTAL LATTICE FRACTURE SI-SI BONDS. EACH SUCH COLLISION PRODUCES SEVERAL ADDITIONAL ELECTRON-HOLE PAIRS WHICH IN TURN MAY RESULT IN MORE COLLISIONS AND EVEN FURTHER IONIZA-TION. SUCH IMPACT IONIZATION RESULTING IN RAPID MULTIPLI-CATION OF AVAILABLE CARRIERS HAS BEEN TERMED AVALANCHE BREAKDOWN.¹

ZENER BREAKDOWN. A VERY HIGH ELECTRIC FIELD ACROSS THE SPACE-CHARGE REGION OF A REVERSE-BLASED P-N JUNCTION DIS-TORTS AND "TEARS APART" COVALENT SI-SI BONDS. ELECTRON-HOLE PAIRS ARE THEN IMMEDIATELY ELEVATED TO THE CONDUCTION BAND. SUCH PENETRATION OF THE FORBIDDEN GAP IS REFERRED TO AS <u>TUNNELING</u> AND LEADS TO THE SO-CALLED ZENER BREAKDOWN EFFECT.²

The electric field required for tunneling in silicon is on the order of 10⁶volts/cm. Zener breakdown is therefore usually preceeded by avalanche breakdown which occurs at a lower electric field potential. The two types of breakdown are not totally independent and one may be responsible for the onset of the other.³

4.3 PREVIOUS INVESTIGATION OF DRAIN BREAKDOWN

<u>CHANNEL IMPACT IONIZATION</u>. COBBOLD⁴ HAS SHOWN THAT CHANNEL BREAKDOWN IN MOS TRANSISTORS IS CHARACTERIZED BY HIGH SATURATION DRAIN CONDUCTANCE AND A GRADUAL INCREASE IN DRAIN CURRENT AT HIGH DRAIN VOLTAGES WHEN THE CHANNEL IS STRONGLY INVERTED.

Recall that as the expop recedes toward the source when the drain voltage is increased following the onset of channel pinch-off, that protion of the channel between the expop and the drain junction becomes a space-charge depletion region. Impact ionization occurs within this depletion region when the transverse electric field^{*} reaches a critical value. Cobbold has found that <u>soft</u> <u>Drain breakdown</u> in MOS transistors is accompanyied by avalanche multiplication at the drain-end of the channel where the electric field is the strongest.⁵

KRESSEL⁶ HAS INVESTIGATED THE EFFECT OF INCLUSIONS IN SEMICONDUCTORS HAVING DIELECTRIC CONSTANTS DIFFERENT FROM THE MEDIUM IN WHICH THEY ARE IMMERSED. IN PARTICULAR, HE HAS FOUND THAT

"IN THE CASE OF A SPHERICAL PRECIPITATE SUCH AS SIO₂ in Si, it may readily be shown that the maximum electric field at the surface of the sphere is approximately 1.5 times larger than the average field in the surrounding medium."⁷

IT HAS LONG BEEN RECOGNIZED THAT INCLUSIONS IN SEMICON-DUCTOR MATERIALS CONTRIBUTES TO ELECTRICAL BREAKDOWN AT P-N JUNCTIONS BY IMPOSING UNUSUALLY HIGH LOCALIZED ELECTRIC FIELDS WHICH ARE ALSO REGIONS OF HIGH CURRENT DENSITY.

^{*}A COMPLETE EXPRESSION FOR THE TRANSVERSE ELECTRIC FIELD MAY BE FOUND IN APPENDIX A, EQUATION A-5. FURTHER INFORMA-TION AND A DERIVATION OF THIS EXPRESSION MAY BE FOUND IN REF-ERENCE 1969-2.

DRAIN-DIODE BREAKDOWN. COBBOLD⁸ HAS SHOWN THAT DRAIN-DIODE BREAKDOWN IN MOS TRANSISTORS IS CHARACTERIZED BY REL-ATIVELY SHARP INCREASES IN DRAIN CURRENT WITH INCREASING VDS. THIS <u>HARD DRAIN BREAKDOWN</u> PHENOMENON HAS BEEN FOUND TO BE A FUNCTION OF THE POTENTIAL DIFFERENCE BETWEEN THE DRAIN AND GATE TERMINALS AND IS THEREFORE <u>GATE-CONTROLLED</u>. THIS FORM OF DRAIN BREAKDOWN HAS BEEN OBSERVED TO BE PAR-TICULARLY PREVALENT IN ENHANCEMENT-MODE DEVICES WHICH ARE BIASED AT OR BELOW CUTOFF AND THAT IT BECOMES LESS APPARANT AS THE CHANNEL IS INVERTED.⁹

FIGURE 4-1 ILLUSTRATES THE MECHANISM BY WHICH GATE-CONTROLLED DRAIN-DIODE BREAKDOWN TAKES PLACE IN AN N-CHANNEL MOSFET. WHEN THE GATE VOLTAGE IS INSUFFICIENT TO ESTABLISH A CONDUCTING CHANNEL, THE TRANSISTOR IS SAID TO BE IN <u>CUTOFF</u>. WITH THE SOURCE AND SUBSTRATE GROUNDED AND THE DRAIN AT A POSITIVE POTENTIAL, A SPACE-CHARGE DEPLETION REGION EXISTS AROUND THE DRAIN JUNCTION. THE POTENTIAL DIFFERENCE BETWEEN THE DRAIN AND GATE TERMINALS AND BE-TWEEN THE DRAIN AND SUBSTRATE CAUSES AN ELECTRIC FIELD TO EXIST IN BOTH THE OXIDE AND IN THE DRAIN DEPLETION REGION.

Since the depletion region is narrowest near the semiconductor surface, the electric field in this <u>corner region</u> is the strongest, as indicated by the density of the electric field lines. Furthermore, since the depletion region extends along and directly beneath the oxide, the depletion

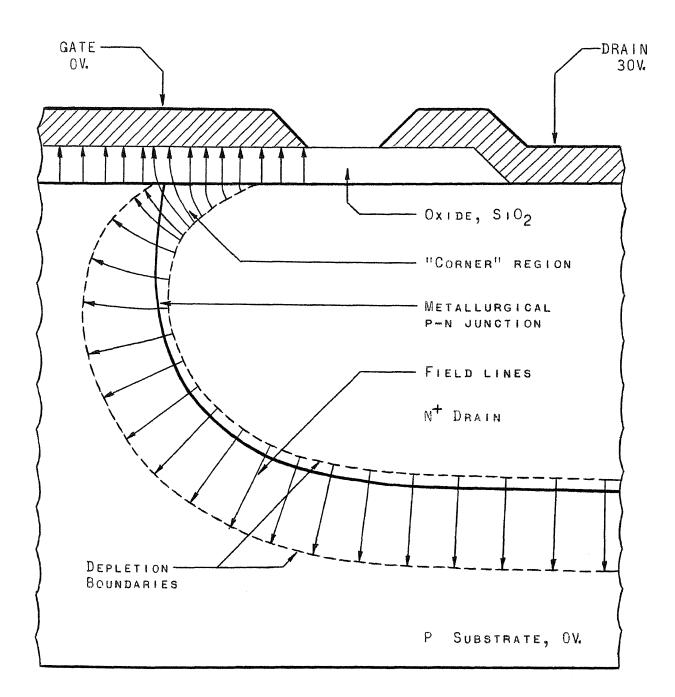


FIGURE 4-1. DRAIN-PORTION OF AN N-CHANNEL ENHANCEMENT MOSFET showing the mechanism of drain-diode breakdown in the "corner" region.

REGION FIELD LINES JOIN WITH AND REINFORCE THE FIELD LINES IN THE INSULATOR. THIS RESULTS IN A VERY HIGH FIELD INTEN-SITY IN THE CORNER REGION AND BRINGS ABOUT IMPACT IONIZA-TION AND AVALANCHE DRAIN-DIODE BREAKDOWN.¹⁰

4.4 INVESTIGATION OF DRAIN BREAKDOWN UNDER CONDITIONS OF STRONG CHANNEL INVERSION IN ENHANCEMENT MOS TRANSISTORS

EXPERIMENTAL METHOD. IN ORDER TO EVALUATE THE EXTENT TO WHICH SOFT DRAIN BREAKDOWN OCCURS AND IN ORDER TO COR-RELATE IT WITH OTHER SURFACE-RELATED PHENOMENA IN P-CHANNEL ENHANCEMENT MOSFET'S, DRAIN BREAKDOWN CHARACTERISTICS OF BOTH RELATIVELY HIGH- AND LOW-NOISE TRANSISTORS WERE OBTAINED.

Using the Quan-Tech Model 327 Diode Noise Analyzer, drain breakdown and noise voltage versus drain current characteristics were plotted for each unit life-tested for threshold voltage drift. I-V characteristics were plotted up to and slightly beyond the threshold of drain breakdown.*

Noise voltage versus drain current at several gate voltages was plotted in order to determine the effects of

^{*}DRAIN BREAKDOWN WAS ARBITRARILY DEFINED AS OCCURRING WHEN THE INSTANTANEOUS SATURATION DRAIN CONDUCTANCE REACHED APPROXIMATELY 100 MICROMHOS. THIS LIMIT WAS IMPOSED IN ORDER TO ENSURE THAT NO UNIT EXCEEDED ITS RATED POWER DISSIPATION LEVEL DURING THE ACCOMPANYING NOISE MEASUREMENTS. ALL NOISE MEASUREMENTS MADE AT HIGH DRAIN VOLTAGES WERE ACCOMPLISHED IN THE SHORTEST TIME INTERVAL POSSIBLE WITHOUT SACRIFICING ACCURACY. DUE TO THE CONFIGURATION OF THE NOISE MEASURING INSTRUMENT, IT WAS IMPOSSIBLE TO USE ANY KIND OF HEAT SINK.

VARIATIONS IN CHANNEL DEPTH AND ELECTRIC FIELD INTENSITY ON BOTH THE NOISE INTENSITY AND THE AMOUNT OF BREAKDOWN WHICH TOOK PLACE.

A CURVE TRACER WAS USED TO VISUALLY INVESTIGATE THE EFFECTS OF VARYING V_{GS} when the devices were operated at a drain voltage producing drain breakdown.

<u>Analysis of Breakdown characteristics</u>. Differences between the drain breakdown characteristics of high- and low-noise transistors is most easily seen by comparing the high voltage I-V characteristics of several devices biased at the same gate voltage.

FIGURE 4-2 ILLUSTRATES THE CHARACTERISTICS OF FIVE HIGH- AND FIVE LOW-NOISE TYPE DD07P MOSFET'S. ALTHOUGH ALL OF THE TRANSISTORS EXHIBITED VARIOUS DEGREES OF SOFT DRAIN BREAKDOWN AT $V_{GS} = -5.5$ volts, three observations are IMMEDIATELY APPARANT:

1. ALL HIGH NOISE UNITS SHOWED CONSIDERABLY HIGHER SATURATION DRAIN CONDUCTANCE THAN DID THE UNITS HAVING A LOW NOISE INTENSITY;

2. THE BREAKDOWN CHARACTERISTICS OF THE HIGH NOISE TRANSISTORS ARE VERY SOFT COMPARED TO THOSE WITH LOW NOISE; AND

3. Some high noise units appear to reach the threshold of drain breakdown at a drain voltage only about half

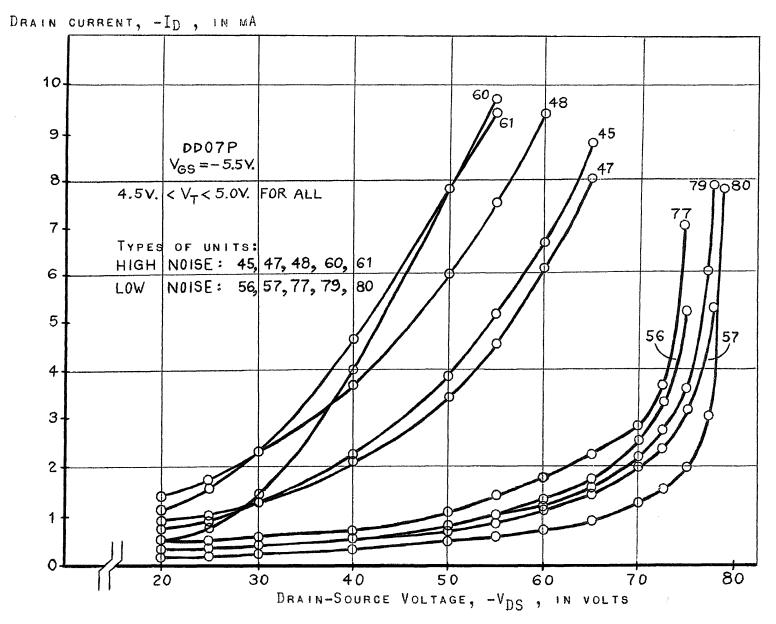


FIGURE 4-2. DRAIN BREAKDOWN CHARACTERISTICS OF FIVE HIGH- AND FIVE LOW-NOISE TYPE DD07P MOSFET'S.

OF THAT REQUIRED BY THEIR LOW NOISE COUNTERPARTS.

Recall that drain leakage measurements conducted with zero gate bias confirmed that saturation drain conductance was a strictly surface effect. Since the magnitude of ¹DSS was approximately the same for both high- and low-noise units, it may be concluded that drain-diode leakage was not a factor influencing the extent of drain breakdown observed. Rather, it may be stated that avalanche multiplication in the channel--primarily in the depletion region between the expop and the drain junction--was the dominant breakdown mechanism when the transistors were biased with strongly inverted channels.

SINCE THE RELATIVE INTENSITY OF 1/F NOISE IS RELATED TO THE SURFACE-STATE DENSITY, NOISE MEASUREMENTS AT CONSTANT DRAIN CURRENT AND GATE VOLTAGE APPEAR TO BE EXCELLENT INDI-CATORS OF THE EXPECTED 'SOFTNESS'' OF DRAIN BREAKDOWN. UNIT-TO UNIT VARIATIONS IN THE CHANNEL SEMICONDUCTOR IMPURITY CONCENTRATION AND/OR SURFACE-STATE DENSITY (BOTH OF WHICH HAVE BEEN SHOWN TO BE RELATED TO THE SATURATION DRAIN CON-DUCTANCE) MAY BE RESPONSIBLE FOR THE GREAT DIFFERENCES IN DRAIN CHARACTERISTICS BETWEEN TRANSISTORS WITH RELATIVELY HIGH NOISE AND THOSE WITH RELATIVELY LOW NOISE INTENSITY.

APPENDIX D LISTS NOISE VOLTAGES VERSUS DRAIN CURRENT FOR ALL TRANSISTORS LIFE-TESTED FOR THRESHOLD VOLTAGE DRIFT AND GIVES CORRESPONDING DRAIN VOLTAGES FOR 1-V CHARACTERISTICS.

CORRELATION OF 1/F NOISE WITH DRAIN BREAKDOWN. FIG-URES 4-3 AND 4-4 SHOW THE NOISE VOLTAGE-DRAIN CURRENT CHARACTERISTICS AT CONSTANT GATE VOLTAGE FOR THREE LOW-AND THREE HIGH-NOISE DD07P MOSFET'S, RESPECTIVELY. By COMPARING THESE CURVES WITH THE DRAIN CHARACTERISTICS SHOWN IN FIGURE 4-2, IT IS POSSIBLE TO MAKE FOUR OBSER-VATIONS:

1. THERE IS AN APPROXIMATE 1/F FREQUENCY DEPENDENCE FOR THE NOISE VOLTAGE OVER THE FULL RANGE OF EACH I-V CHARACTERISTIC;

2. THE NOISE CHARACTERISTICS ARE DEFINITELY PEAKED, INCREASING WITH THE ONSET OF DRAIN BREAKDOWN AND DECREASING WITH FURTHER INCREASES IN DRAIN CURRENT;

3. THE PEAK NOISE VOLTAGE IS LOWEST IN THE LOW NOISE TRANSISTORS AND OCCURS AT A LOWER DRAIN CURRENT BUT CORRES-PONDINGLY HIGHER DRAIN VOLTAGE THAN FOR THOSE UNITS HAVING A HIGH NOISE INTENSITY; AND

4. THE SATURATION DRAIN CONDUCTANCE AT THE THRESHOLD OF DRAIN BREAKDOWN IS APPROXIMATELY IDENTICAL FOR EACH UNIT AT THE DRAIN CURRENT FOR WHICH ITS RESPECTIVE NOISE VOLTAGE IS A MAXIMUM.

THE FIRST OBSERVATION INDICATES THAT 1/F NOISE THEORY IS APPROPRIATE, EVEN AFTER DRAIN BREAKDOWN HAS RESULTED IN SUBSTANTIAL DRAIN CURRENT. HOWEVER, THE PEAKING AND STEADY DECREASE IN NOISE VOLTAGE FOLLOWING DRAIN BREAKDOWN WAS NOT

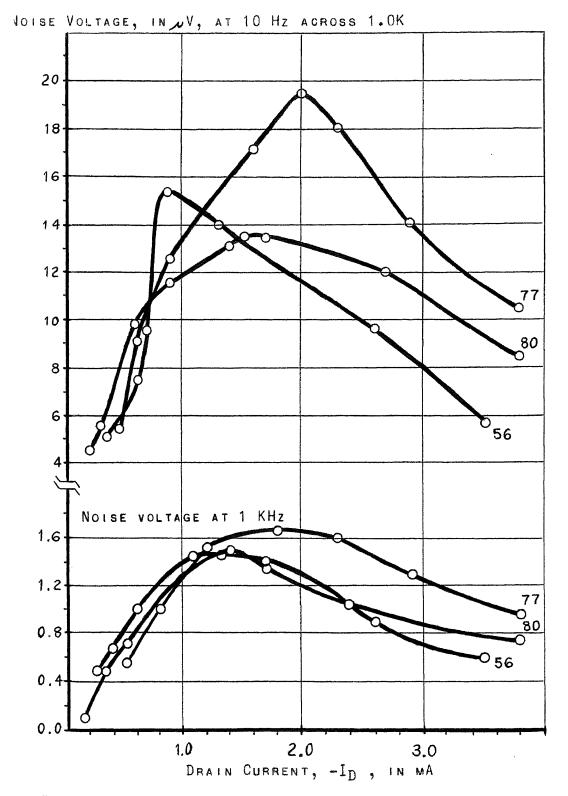
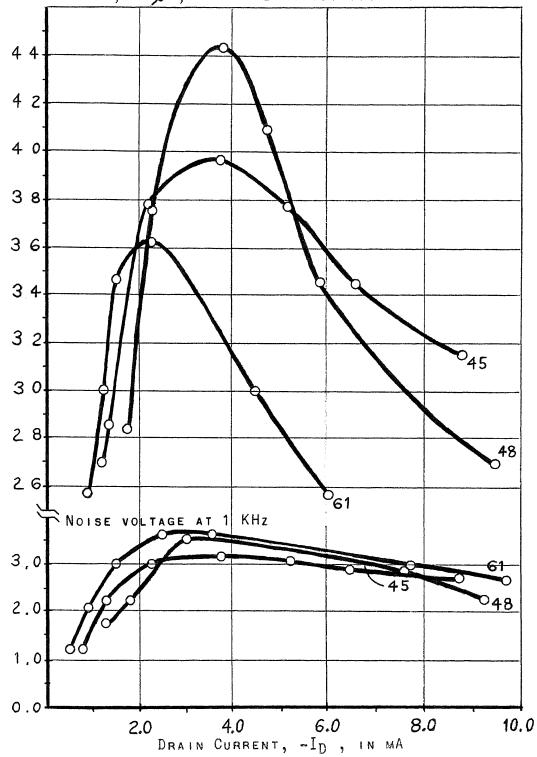


Figure 4-3. Noise voltage-drain current characteristics for three low noise type DD07P MOSFET's. V_{GS} =-5.5V. and V_{DS} ranged from -20 to -78V.



Noise Voltage, IN V, AT 10 Hz ACROSS 1.0 K

Figure 4-4. Noise voltage-drain current characteristics for three high noise type DD07P MOSFET's. $V_{\rm GS}=-5.5$ V. and $V_{\rm DS}$ ranged from -20 to -65V.

EXPECTED, PARTICULARLY IN VIEW OF THE FACT THAT THE NOISE INTENSITY DECREASES WHILE THE DRAIN CURRENT RISES ALMOST EXPONENTIALLY WITH FURTHER INCREASES IN DRAIN VOLTAGE.

The rapid increase in noise voltage at the onset of drain breakdown may be explained on the basis of the relationship between 1/F noise and current density. By increasing $|-V_{DS}|$ with constant gate voltage, the expop recedes toward the source, thereby permitting increased drain current. When the drain voltage is sufficiently large, the transverse electric field in the channel depletion region causes impact ionization resulting in avalanche breakdown. Increased drain current within a relatively constant crosssection of the channel (due to the constant gate voltage) results in increased current density and correspondingly increased 1/F noise intensity.

As the drain voltage is further increased, ionization takes place further along the drain p-n junction. As breakdown progresses, the effective cross-section through which carriers are injected into the drain increases. This lowers the current density and the noise intensity decreases proportionately.

FIGURE 4-5 SHOWS THE NOISE VOLTAGE-DRAIN CURRENT CHAR-ACTERISTICS AT SEVERAL GATE VOLTAGES FOR A SINGLE HIGH NOISE TYPE DD07P MOSFET. IT IS EVIDENT THAT THE PEAK NOISE INTEN-SITY IS PROPORTIONAL TO THE APPLIED GATE VOLTAGE AND THAT

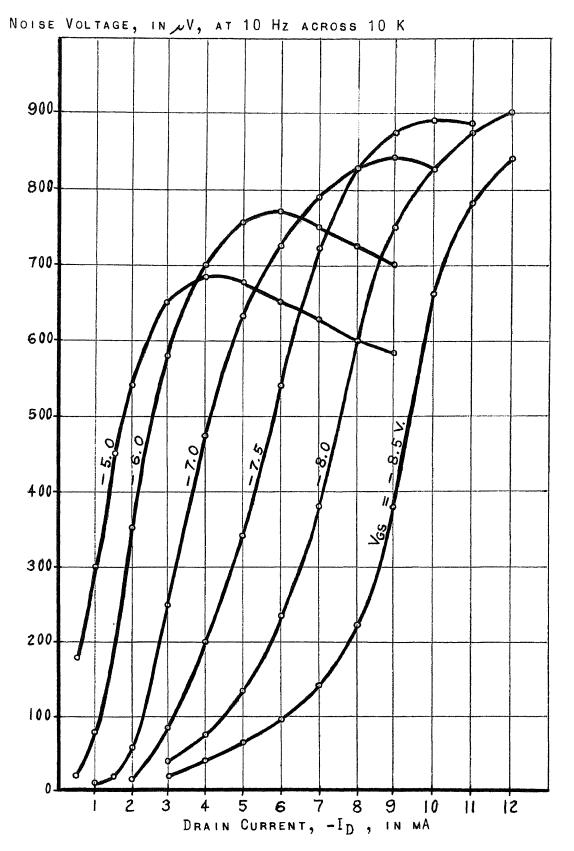


FIGURE 4-5. NOISE VOLTAGE-DRAIN CURRENT CHARACTERISTICS AT SEVERAL GATE VOLTAGES FOR A SINGLE HIGH NOISE TYPE DD07P MOSFET. UNIT NO. 61.

THE PEAKS OCCUR AT SUCCESSIVELY HIGHER VALUES OF DRAIN CURRENT WITH INCREASING $|-V_{\rm GS}|$. This is in excellent agreement with the findings of ${\rm Hsu}^{11}$ who states that "low frequency noise increases monotonically with negative gate bias on p-channel enhancement MOSFET's"¹² at the same drain voltage, indicating that "the surface-state density increases toward the edge of the valence band."¹³ It was also found that the surface-state density increased with increased surface potential. Therefore, 1/F noise measurements at constant drain current should indicate the relative surface-state density as a function of gate bias.

THE THIRD OBSERVATION--DIFFERENCES IN THE MAGNITUDE OF THE PEAK NOISE VOLTAGE BETWEEN THE LOW NOISE TRANSISTORS AND THOSE HAVING HIGH NOISE INTENSITY--MAY BE TRACED TO DIFFERENCES IN SURFACE-STATE DENSITY BETWEEN THE VARIOUS DEVICES. SINCE 1/F NOISE IS DIRECTLY PROPORTIONAL TO SUR-FACE-STATE DENSITY, IT FOLLOWS THAT THOSE UNITS HAVING A RELATIVELY LOW NOISE INTENSITY ALSO POSSESS FEWER SURFACE STATES THAN HIGH NOISE TRANSISTORS.

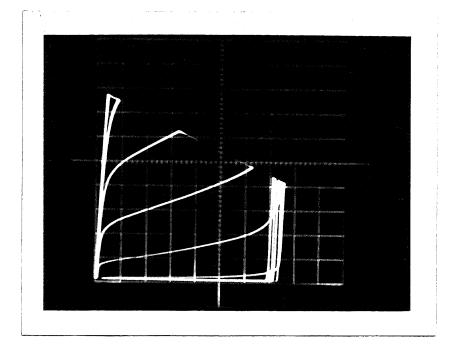
THE FOURTH OBSERVATION INDICATES THAT EXCESS NOISE MEASUREMENTS MAY BE USED WITH SOME DEGREE OF ACCURACY TO DETERMINE THE APPROXIMATE BREAKDOWN VOLTAGE OF AN MOS TRAN-SISTOR WITHOUT ACTUALLY BREAKING DOWN THE DEVICE. SUCH NONDESTRUCTIVE TESTING COULD PROVE INVALUABLE WHEN IT IS DESIRED TO KNOW THE BREAKDOWN VOLTAGE OF A DEVICE OR THAT

AN ENTIRE MOS INTEGRATED CIRCUIT. THE RELATIVE PEAK VOL-TAGE THAT CAN BE SAFELY HANDLED BY A DEVICE COULD BE DE-TERMINED WITHOUT EXHAUSTIVE MEASUREMENTS WHICH, IF NOT DONE PROPERLY, COULD SUBJECT THE CHIP TO UNUSUALLY HIGH LEVELS OF POWER DISSIPATION.

4.5 INVESTIGATION OF GATE-CONTROLLED DRAIN-DIODE BREAKDOWN UNDER CUTOFF CONDITIONS

EXPERIMENTAL METHOD. DRAIN BREAKDOWN CHARACTERISTICS AT SEVERAL GATE VOLTAGES EXTENDING DOWN BELOW CUTOFF WERE PLOTTED FOR TRANSISTORS HAVING RELATIVELY HIGH- AND RELATIVELY LOW-NOISE INTENSITY. THE QUAN-TECH MODEL 327 DIODE NOISE ANALYZER WAS USED TO OBTAIN 1-V CURVES FOR THE TYPE DD07P TRANSISTORS WHICH HAD UNDERGONE 500 HOURS OF BIAS AT $130 \pm 5^{\circ}$ C. Gate voltages ranged from $V_{\rm GS} = -8.5$ to +1 volt. In addition, a curve tracer was used to visually investigate drain-diode breakdown and to obtain pictures showing extremes in this phenomena between high and low Noise TRANSISTORS.

EXPERIMENTAL RESULTS AND DISCUSSION. FIGURES 4-6 AND 4-7 ILLUSTRATE THE DRAMATIC DIFFERENCE IN DRAIN BREAKDOWN CHARACTERISTICS BETWEEN P-CHANNEL TRANSISTORS HAVING RELA-TIVELY HIGH- AND LOW-NOISE INTENSITY. FIGURE 4-6 SHOWS THE TYPICAL MODERATE SATURATION DRAIN CONDUCTANCE AND SOFT CHAN-NEL BREAKDOWN AT INTERMEDIATE AND HIGHLY NEGATIVE GATE VOL-TAGES AND THE INCREASINGLY HARDER DRAIN-DIODE BREAKDOWN AT

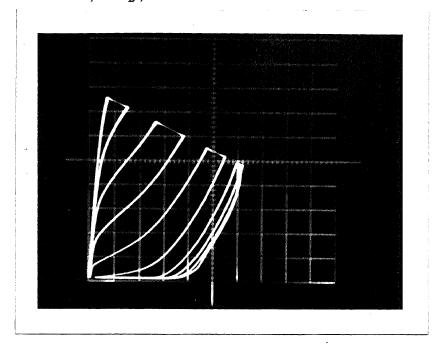


DRAIN CURRENT, $-I_D$, 2MA/DIVISION

DRAIN VOLTAGE, -VDS, 10 VOLTS/DIVISION

Gate Voltage in one volt increments starting from $-\mathrm{V}_{GS}=0$ volts

Figure 4-6. Photograph showing the current-voltage drain breakdown characteristics of a relatively low noise p-channel enhancement MOSFET. Hard drain - diode breakdown is evident at low $|-V_{GS}|$ and $-V_{DS} = 70$ volts.



DRAIN CURRENT, -ID, 2MA/DIVISION

DRAIN VOLTAGE, -VDS, 10 VOLTS/DIVISION

Gate Voltage in one volt increments starting from -V_{GS} = 0 volts

FIGURE 4-7. PHOTOGRAPH SHOWING THE CURRENT-VOLTAGE DRAIN BREAKDOWN CHARACTERISTICS OF A RELATIVELY HIGH NOISE P-CHAN-NEL ENHANCEMENT MOSFET. UNUSUALLY HIGH SATURATION DRAIN CONDUCTANCE AND VERY SOFT BREAKDOWN INDICATES THAT CHANNEL PUNCH-THROUGH HAS OCCURED. LOW GATE VOLTAGES FOR A VERY LOW NOISE DEVICE. FIGURE 4-7 shows the characteristics of a technically poor, high noise MOS transistor.

IT IS EVIDENT THAT INCREASING THE MAGNITUDE OF THE GATE VOLTAGE LESSENS THE POTENTIAL DIFFERENCE BETWEEN THE DRAIN AND GATE TERMINALS, THEREBY REDUCING THE ELECTRIC FIELD INTENSITY IN THE CORNER REGION AND SOFTENING THE BREAKDOWN CHARACTERISTICS. ON THE OTHER HAND, AS THE GATE VOLTAGE IS BROUGHT TO ZERO, OR EVEN MADE POSITIVE,^{*} THE ELECTRIC FIELD AT THE SEMICONDUCTOR SURFACE NEAR THE DRAIN JUNCTION IS GREATLY INCREASED AND THE BREAKDOWN CHARACTER-ISTICS BECOME VERY SHARP.

The very high saturation drain conductance (which gives the curves in figure 4-7 the appearance of triode characteristics) found in the relatively high noise transistor is indicative of <u>drain punch-through</u>, whereby the channel depletion region extends from the drain to the source and the drain current is space-charge limited. Frohman-Bentchkowsky and Grove¹⁴ have investigated the mechanisms by which MOS transistors breakdown by the punch-through phenomenon. They have concluded that it is characteristic of devices having short channel lengths and/or unusually low bulk doping.

^{*}THE DD07P MOSFET'S POSSESSED ZENER DIODE VOLTAGE LIM-ITING BETWEEN THE GATE AND SOURCE. APPLICATION OF A POSI-TIVE GATE VOLTAGE GREATER THAN 1V. WOULD HAVE DRAWN TOO MUCH CURRENT THROUGH THIS DIODE, CAUSING EXCESSIVE HEATING.

However, under the assumption that both the semiconductor impurity concentration and channel length were nearly identical for both the high and low noise transistors, and on the basis of the earlier correlation between 1/F noise and surface-state density, there is strong evidence to suggest that conditions at the silicon-silicon dioxide interface are important determining factors in the type of drain breakdown which occurs under a given bias condition. Since saturation drain conductance and 1/F noise are both surface-related, it was concluded that the relative noise intensity was indicative of either soft or hard drain breakdown.

IT IS CUSTOMARY TO DESIGN HIGH FREQUENCY TRANSISTORS (INCLUDING THE TYPES DD07P, DD08P, AND DD08K MOSFET'S USED IN THIS RESEARCH) WITH RELATIVELY SHORT CHANNELS. THERE-FORE, 1/F NOISE MEASUREMENTS MAY BE USEFUL IN "WEEDING OUT" DEVICES HAVING BREAKDOWN CHARACTERISTICS SIMILAR TO THOSE OF FIGURE 4-7.

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CHAPTER 5

CONCLUSIONS AND DISCUSSION OF RESULTS

This investigation has found a definite correlation Between excess noise intensity and drain breakdown in p-channel enhancement-mode MOS field-effect transistors. In particular, it was found that transistors having relatively low noise intensity exhibited sharper breakdown characteristics at higher drain voltages than devices having high noise. It is believed that the surfacestate density is related to the semiconductor impurity concentration and that transistors having low bulk doping and/or high surface-state density may be expected to exhibit high saturation drain conductance and soft breakdown characteristics. Finally, it was concluded that 1/F noise measurements may be used in nondestructive testing to determine the approximate breakdown voltage of individual devices or entire MOS integrated circuits.

Previous investigations have traced the origin of 1/F noise to modulation of surface recombination and have shown that the noise intensity is directly proportional to both the surface-state density and the current density. Measurements of drain current fluctuations at 10Hz and at 1KHz were found to be consistent with 1/F noise theory. The frequency dependence was found to exist over a wide range of operating conditions, including at the onset and follow-

ING DRAIN BREAKDOWN. THEREFORE, 1/F NOISE THEORY WAS APPLIED IN THIS INVESTIGATION.

Results from life-testing for threshold voltage drift were consistent with accepted theory of ionic drift in the oxide. Excess noise intensity, particularly at 1KHz, showed a substantial increase following negative gate bias at $130 \pm 5^{\circ}$ C for 500 hours. This was indicative of increased surface-state density and confirms that cations at the SI-SiO₂ interface contribute to an increase in the number of fast trapping states. It was concluded that ionic drift near the semiconductor surface resulted in a large number of shallow traps, the fast time constants of which accounted for an unusually high increase in noise intensity at higher frequencies.

IT IS WELL KNOWN THAT SATURATION DRAIN CONDUCTANCE IS DIRECTLY PROPORTIONAL TO THE OXIDE THICKNESS AND INDIRECTLY PROPORTIONAL TO THE SEMICONDUCTOR IMPURITY CONCENTRATION. CURRENT-VOLTAGE CHARACTERISTICS OF SEVERAL TRANSISTORS SHOWED THAT HIGH DRAIN CONDUCTANCE CORRESPONDED TO HIGH 1/F NOISE INTENSITY. THIS ESTABLISHED A CORRELATION BETWEEN DRAIN CONDUCTANCE, SURFACE-STATE DENSITY, AND THE LEVEL OF BULK DOPING.

SINCE THERE WAS NO CORRELATION BETWEEN IDS AND OBSERVED DIFFERENCES BETWEEN TRANSISTORS HAVING RELATIVELY HIGH AND LOW NOISE, IT WAS CONCLUDED THAT HIGH SATURATION DRAIN CON-

DUCTANCE IS STRICTLY A SURFACE EFFECT AND THAT THE RELA-TIVE 1/F NOISE INTENSITY MAY BE USED TO CORRELATE DRAIN CONDUCTANCE WITH THE SURFACE-STATE DENSITY.

IT WAS FOUND THAT TRANSISTORS HAVING RELATIVELY HIGH EXCESS NOISE CONSISTENTLY EXHIBITED SOFT DRAIN BREAKDOWN AND HIGH DRAIN CONDUCTANCE WHILE BREAKDOWN IN RELATIVELY LOW NOISE DEVICES UNDER SIMILAR BIAS CONDITIONS WAS ALWAYS SIGNIFICANTLY MORE ABRUPT AND WAS FOUND TO OCCUR AT UP TO TWICE THE DRAIN VOLTAGE.

IN A FEW EXTREME CASES, VERY HIGH NOISE TRANSISTORS BLASED IN CUTOFF SHOWED BREAKDOWN CHARACTERISTICS WHICH WERE INDICATIVE OF PUNCH-THROUGH. IT WAS CONCLUDED THAT DRAIN BREAKDOWN IS RELATED TO THE SEMICONDUCTOR IMPURITY CONCENTRATION AND THE SURFACE-STATE DENSITY AND THAT LOW BULK DOPING AND/OR AN UNUSUALLY HIGH DENSITY OF SURFACE STATES FACILITATES EXPOP RECESSION TOWARD THE SOURCE WHEN THE TRANSISTOR IS OPERATED IN SATURATION.

OF PARTICULAR SIGNIFICANCE WAS THE FACT THAT THE NOISE INTENSITY CONSISTENTLY PEAKED AT THE THRESHOLD OF DRAIN BREAKDOWN AND STEADILY DECREASED WITH INCREASING DRAIN CURRENT. SINCE THE NOISE INTENSITY IS DIRECTLY PROPOR-TIONAL TO THE CURRENT DENSITY WITHIN THE DEVICE, IT WAS REASONED THAT THE EFFECTIVE CHANNEL DEPTH AT THE DRAIN JUNCTION INCREASES AS IONIZATION PROGRESSES FOLLOWING THE ONSET OF DRAIN BREAKDOWN. THIS MAY BE EXPLOITED IN NON-

DESTRUCTIVE TESTING OF MOS DEVICES, INCLUDING ENTIRE MOS INTEGRATED CIRCUITS, BY MEASURING THE RELATIVE 1/F NOISE INTENSITY TO DETERMINE THE APPROXIMATE DRAIN BREAKDOWN VOLTAGE WITHOUT ACTUALLY BREAKING DOWN THE DEVICE.

IN CONCLUSION, AN EXCEPTIONALLY HIGH 1/F NOISE INTEN-SITY IS INDICATIVE OF A HIGH SURFACE-STATE DENSITY AND/OR LOW SEMICONDUCTOR IMPURITY CONCENTRATION. MOS TRANSISTORS HAVING RELATIVELY LOW NOISE INTENSITY MAY BE EXPECTED TO EXHIBIT LOW SATURATION DRAIN CONDUCTANCE AND SHARP DRAIN BREAKDOWN CHARACTERISTICS AT HIGHER DRAIN VOLTAGES THAN HIGH NOISE DEVICES.

APPENDIX A

THE INFLUENCE OF DEVICE PARAMETERS

FROHMAN-BENTCHKOWSKY AND GROVE¹ AND OTHERS²,³ have investigated the drain conductance of MOS transistors in the saturation region and have found it to be "a sensitive function of the oxide thickness as well as the substrate impurity concentration."⁴ The following is a model presented which evaluates the saturation drain conductance of an n-channel MOS transistor on the basis of device parameters.

THE DRAIN VOLTAGE AT THE ONSET OF SATURATION MAY BE EX-

$$V_{DSAT} = V_{G} - V_{FB} - 2\phi_{F} + \frac{K_{s}\epsilon_{0}gC_{B}}{C_{0}^{2}}$$

$$\cdot \left[1 - \sqrt{\frac{1 + 2C_{0}^{2}(V_{G} - V_{FB})}{K_{s}\epsilon_{0}gC_{B}}}\right]$$
A-1

where $C_B = Impurity$ concentration of the substrate

 $C_0 = K_0 \varepsilon_0 / x_0 = CAPACITANCE OF THE OXIDE PER UNIT AREA$ $<math>K_0, K_s = Dielectric constants of the oxide and semicon$ ductor, respectively $<math>V_G = MOS$ transistor gate voltage $V_{FB} = Flat-band$ voltage g = Magnitude of electronic charge

 $x_o = Thickness of the oxide$

 $\epsilon_o = Permittivity$ of free space

 $\phi_{\rm F}$ = Fermi potential in the substrate

The corresponding drain current when the drain voltage is greater than $V_{D,\text{sat}}$ is given by 6

$$\begin{vmatrix} D_{SAT} \\ V_{D} = V_{DSAT} \end{vmatrix} = \frac{|D_{SAT}^{*}|}{1 - \Delta L/L} A^{-2}$$

where I_{DSAT}^* is the drain current when $V_D^-V_{DSAT}$, L is the effective length of the inverted channel before saturation, ΔL is the length by which the channel is shortened when the expop has moved toward the source, and where it is assumed that an increase in drain current after the onset of saturation is attributed entirely to a reduction of channel length. If the expop recedes only slightly and $\Delta L \ll L$,

$$I_{\text{DSAT}} \simeq I_{\text{DSAT}}^*$$
 A-3

OVER A WIDE RANGE OF DRAIN VOLTAGES IN THE SATURATION RE-GION OF THE CURRENT-VOLTAGE CHARACTERISTICS.

THE LENGTH OF THE CHANNEL DEPLETION REGION BETWEEN THE EXPOP AND THE DRAIN JUNCTION IS GIVEN BY⁷

$$\Delta L = \frac{V_{\rm D} - V_{\rm D_{SAT}}}{\varepsilon_{\rm T}} \qquad A-4$$

where $V_{\rm D}$ is the potential at the drain, $V_{\rm DSAT}$ is the poten-tial at the expop, and where $\Xi_{\rm T}$, the average transverse electric field in the channel depletion region, is composed of components associated with the following:

1. THE ELECTRIC FIELD DUE TO THE FIXED CHARGE IN THE REVERSE BLASED DRAIN-TO-SUBSTRATE P-N JUNCTION;

2. The fringing electric field due to the drain-togate potential drop, $V_D - V_G'$; and

3. The fringing electric field due to the potential difference $V_{G}' - V_{DSAT}$ between the gate and the expop and may be expressed as⁸

$$\mathcal{E}_{T} = \frac{1}{K} (V_{D} - V_{DSAT})^{\frac{1}{2}} + \frac{5}{K_{0}} \frac{(V_{D} - V_{G'})}{K_{s} - x_{0}} + \frac{\frac{5}{K_{0}} \frac{(V_{G'} - V_{DSAT})}{K_{s} - x_{0}}$$

where $V_G' = V_G + Q_{SS}/C_0$ = the effective gate voltage $Q_{SS} = F_{IXED}$ surface-state charge density per unit area $5, \xi = F_{IELD}-F_{RINGING}$ factors experimentally found to BE 0.2 AND 0.6, RESPECTIVELY⁹

AND WHERE¹⁰

$$\mathbf{K} = \left(\frac{2\kappa_{\rm s}\epsilon_{\rm o}}{\mathscr{P}^{\rm C_{\rm B}}}\right)^{\frac{1}{2}} \qquad A-6$$

IS A CONSTANT INVERSELY PROPORTION TO THE SUBSTRATE IMPUR-ITY CONCENTRATION.

After substituting for \mathcal{E}_{T} in equation A-4, the length of the depletion region is given by 11

$$\Delta L = K(V_{D} - V_{DSAT})^{-\frac{1}{2}} + \frac{K_{S} x_{O}}{K_{O}} \frac{(V_{D} - V_{DSAT})}{\alpha(V_{D} - V_{G}') + \beta(V_{G}' - V_{DSAT})}$$
A-7

which expresses ΔL by the sum of two semi-autonomous terms. The first is an inverse function of the semiconductor impurity concentration (based on the expression for K) and the second is directly proportional to the oxide thickness. Accordingly, there are two limiting cases which emphasize the dependencies of the saturation drain conductance on the device parameters.

Case 1. $x_0 \gg K(V_D - V_{DSAT})^{\frac{1}{2}}$. With the assumption of a high substrate impurity concentration, thick oxide layer, or both, the reduction in channel length becomes¹²

$$\Delta L \simeq K (v_D - v_{D_{SAT}})^{\frac{1}{2}} \qquad A-8$$

AND THE SATURATION DRAIN CONDUCTANCE DUE TO A DIFFERENTIAL REDUCTION IN CHANNEL LENGTH FROM EQUATION 1-1 BECOMES¹³

$$g_{d_{SAT}} \simeq \frac{K |D_{SAT}^*}{2L(1 - \Delta L/L)^2 \cdot (V_D - V_{DSAT})^{\frac{1}{2}}}$$
 A-9

which is dependent only on the level of bulk doping, ${\rm C}^{}_{\rm B}{\scriptstyle \bullet}$

IT IS EVIDENT THAT THE SMALL-SIGNAL SATURATION DRAIN CONDUCTANCE INCREASES AS THE SQUARE-ROOT OF THE IMPURITY CONCENTRATION DECREASES. THIS AGREES WELL WITH RESULTS FOUND BY CRAWFORD¹⁴ who has expressed the length of the CHANNEL DEPLETION REGION AS A FUNCTION OF V_T, the threshold VOLTAGE, AND THE DONOR IMPURITY CONCENTRATION¹⁵

$$\Delta L = \sqrt{\frac{2\epsilon_{s} \left[V_{D} - (V_{G} - V_{T})\right]}{g^{N_{D}}}} \qquad A-10$$

CASE 2. $x_0 \ll K(V_D - V_{DSAT})^{\frac{1}{2}}$. With low bulk doping, thin oxide, or both, ΔL is given by ¹⁶

$$\Delta L \simeq \frac{K_{s} x_{0}}{K_{0}} \frac{(V_{D} - V_{DsAT})}{\alpha (V_{D} - V_{G}') + \beta (V_{G}' - V_{DsAT})}$$
 A-11

AND THE CORRESPONDING SMALL-SIGNAL SATURATION DRAIN CON-DUCTANCE IS EXPRESSED AS¹⁷

$$g_{A_{SAT}} = \frac{K_{S} x_{O}}{K_{O} L} \frac{(\beta - \alpha) (V_{G}' - V_{DSAT})}{[\alpha (V_{D} - V_{G}') + \beta (V_{G}' - V_{DSAT})]^{2}} A-12$$

$$\cdot \frac{I_{DSAT}^{*}}{(1 - \Delta L/L)^{2}}$$

which is a function of the oxide thickness only. It is evident that variations in the oxide thickness produce direct variations in drain conductance. The term x_0 is therefore slightly more dominant an influencing factor than C_B .

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APPENDIX B

A DERIVATION OF THE EQUATION OF THE SPECTRAL INTENSITY

OF THE SHORT-CIRCUIT DRAIN CURRENT FLUCTUATIONS

IN METAL-OXIDE-SEMICONDUCTOR TRANSISTORS

THE ORIGIN OF EXCESS NOISE IN MOS TRANSISTORS CAN BE TRACED TO "CONDUCTION CHANNEL CHARGE DENSITY FLUCTUATIONS CAUSED BY THE MODULATION OF SURFACE POTENTIAL DUE TO RAN-DOM OCCUPANCY OF SURFACE STATES."¹ Hsu² has undertaken a DERIVATION OF THE 1/F NOISE SPECTRAL INTENSITY OF P-CHANNEL ENHANCEMENT MOSFET'S BASED ON MODULATION OF SURFACE-STATE DENSITY WITH THE FOLLOWING FIVE ASSUMPTIONS:³

1. ONLY A ONE-DIMENTIONAL PROBLEM IS CONSIDERED FOR CARRIER DISTRIBUTION AT THE FIELD-INDUCED DE-PLETION REGION.

2. THE SURFACE STATES ARE UNIFORMLY DISTRIBUTED IN SPACE AND IN ENERGY AT THE CENTRAL PORTION OF THE SILICON FORBIDDEN GAP.

3. SINCE THE OXIDE THICKNESS IS ON THE ORDER OF 600 to 2000Å, and since it has been shown that the great majority of surface states reside within no more than 20Å from the SI-SIO2 interface, ⁴ all surface states are located at the surface or interface of the oxide and silicon substrate.

4. THE RELAXATION TIME OF FREE HOLES AND ELECTRONS INSIDE THE SURFACE DEPLETION REGION IS NEGLIGIBLY SMALL COMPARED TO THAT OF THE SURFACE STATES.

5. For a p-channel MOSFET, when the surface is inverted, the free hole concentration is very large while the free electron density is negligibly small in the conduction channel. Therefore, the surface states act as trapping centers instead of recombination centers. Using the gain-loss process of Burgess^{5,6} and of van Vliet and Fassett,⁷ the rate that surface states of energy E, located at a distance x into the silicon dioxide layer, gain holes can be expressed as⁸

$$G = c \left[p_{S} \eta_{T}(E) + \eta_{SI} \eta_{T}(E) \right] A \Delta E \Delta x \qquad B-1$$

WHERE A = EFFECTIVE AREA OF THE GATE

- $\eta_{\tau}(E) = Number of empty surface states at energy E$ and located at a distance x from the SI-SIO₂interface
 - $\gamma_{\rm SI}={\sf FREE}$ electron concentration at the surface when the Fermi level is at the surface-state level
 - P_s = Hole density in the inverted channel of a p-channel enhancement MOSFET

AND WHERE C IS THE HOLE CAPTURE PROBABILITY OF SURFACE STATES, A FUNCTION OF ENERGY AND DISTANCE INTO THE OXIDE AND STRONGLY DEPENDENT ON THE LOCATION OF THE TRAPPING CENTER, WHICH MAY BE EXPRESSED AS⁹

$$c = c_s e^{-\alpha x}$$
 B-2

Here c_s is the capture probability at a distance x = 0 from the SI-SIO₂ interface and the factor α is given by¹⁰

$$\alpha = \frac{4\pi}{\hbar} \sqrt{2M^* E_B} \qquad B-3$$

where $E_B = E_NERGY$ barrier to be tunneled by holes

h = Planck's constant $m^* = Effective mass of a hole$ IF THE CAPTURE PROBABILITY OF ELECTRONS IS ASSUMED TO BE THE SAME AS THAT FOR HOLES, THEN THE RATE THAT THE SURFACE STATES OF ENERGY E LOSE HOLES AT A DISTANCE X INTO THE OXIDE IS GIVEN BY ¹¹

$$\mathcal{R} = c \left[\eta_{s} \rho_{T}(E) + \rho_{s} \rho_{T}(E) \right] A \Delta E \Delta x \qquad B-4$$

- where $\eta_{\rm S}={\rm Electron\ density\ in\ the\ inverted\ channel\ of\ a}$ p-channel enhancement MOSFET
 - $p_{s_1} = F_{REE}$ hole concentration at the surface when the Fermi level is at the surface-state level
 - $P_{\rm T}({\rm E}) = {\rm Density}$ of trapped holes at energy E and located at a distance x from the SI-SIO₂ interface

The time constant for the surface-state occupancy fluctuation, obtained by differentiating equations B-1 and B-4 with respect to the number of holes in the surface states ($p_T A \Delta E \Delta x$), is given by¹²

$$\gamma = \frac{1}{c \left[p_{so} + \eta_{s1} + p_{s1} + \eta_{so} \right]}$$
 B-5

where p_{so} and γ_{so} are the steady-state values of p_s and γ_s , respectively.

When the surface is inverted, assumption (5) reduces the expression for γ to¹³

$$\gamma = \frac{1}{c(\varphi_{so} + \gamma_{s_i})}$$
 B-6

"The variance of the fluctuation of the total number of occupied surface states in an energy ΔE and at a distance Δx is given by $G_0 \gamma$, "¹⁴ where G_0 is the steady-state rate of formation of trapping centers (the rate of gaining holes) which is assumed to be equal to \mathcal{R}_0 , the steadystate rate of losing holes. In the inversion region of a p-channel enhancement MOSFET virtually all surface states can be thought of as trapping centers and the variance of the trapped hole density fluctuation is¹⁵

$$V_{AR} p_{\tau} = \frac{N_{\tau} p_{sc} \eta_{si}}{(p_{sc} + \eta_{si})^2} \qquad B-7$$

where N_{τ} is the density of surface states per cm^3 per eV.

The spectral intensity of the trapped hole density fluctuation must be evaluated over a length ℓ from the SI-SIO₂ interface and over all possible energy levels and is given by ¹⁶

$$S_{\mathcal{P}_{\tau}}(\omega) = \int_{-\infty}^{+\infty} \int_{0}^{\infty} \frac{4 V_{AR} \mathcal{P}_{\tau} \cdot A_{\tau}}{1 + (\omega \tau)^{2}} dx dE \qquad B-8$$

AFTER SUBSTITUTING FOR VAR 10, AND 7, AND AFTER EVAL-UATING THE DOUBLE INTEGRAL, THE SPECTRAL INTENSITY OF THE TRAPPED HOLE DENSITY FLUCTUATION MAY BE EXPRESSED IN THE FORM¹⁷

$$S_{\rho_{\tau}}(\omega) = \frac{4kTAN_{\tau}(E)}{\alpha\omega} \left[TAN^{-1}\omega_{\tau_{s}} e^{\alpha \ell} - TAN^{-1}\omega_{\tau_{s}} \right] \qquad B-9$$

where $N_T(E)$ is the combined density of free holes and electrons in cm⁻³/eV and where γ_s is the time constant of the surface states located right at the oxide-semiconductor interface.

The flow of mobile charge in the channel is dependent on the transconductance, Gm, the oxide capacitance per unit area, C_0 , and on the effective area of the gate, A, and is given by 18

$$\begin{bmatrix} \frac{\varphi G M}{A C_0} \end{bmatrix} = \begin{bmatrix} \frac{(\text{coul.})(\text{coul.}^2/\text{joule} \cdot \text{sec.})}{(\text{cm}^2)(\text{coul.}^2/\text{joule} \cdot \text{cm}^2)} \end{bmatrix}$$
$$= \begin{bmatrix} \frac{(\text{coul.}^3/\text{joule} \cdot \text{sec.})}{(\text{coul.}^2/\text{joule})} \end{bmatrix} = \frac{\text{coul.}}{\text{sec.}} \qquad B-10$$

WHICH IS IN UNITS OF CURRENT.

CHARGE DENSITY FLUCTUATION CAUSING A VARIATION IN THE SURFACE POTENTIAL RESULTS IN DRAIN CURRENT FLUCTUATIONS IN THE DRAIN CIRCUIT. THESE MAY BE OBSERVED AS NOISE. THERE-FORE, THE SPECTRAL INTENSITY OF THE SHORT-CIRCUIT DRAIN CUR-RENT FLUCTUATIONS MAY BE EXPRESSED AS¹⁹

$$S_{i(\omega)} = \left[\frac{q_{f}G_{M}}{AC_{o}}\right]^{2} \cdot S_{p_{r}(\omega)} = \overline{i}_{df}^{2} \qquad B-11$$

WHERE $\overline{\dot{l}_{df}}^2$ is the mean-squared current fluctuation.

THE NATURE OF THE FREQUENCY DEPENDENCE OF THE NOISE SPECTRAL INTENSITY IS DETERMINED BY EVALUATING THE EXPAN-SION OF THE ARCTANGENT TERMS IN EQUATION B-9. <u>Case 1.</u> Very low frequency: $\omega T_5 \stackrel{\alpha \ell}{\in} < 1$. After substituting for \mathcal{A} , M^* , $E_B = 4eV$, $\mathcal{L} \simeq 20 \hat{A}$, and solving for in equation B-3, $\stackrel{\alpha \ell}{\cong} \simeq 10^{17}$. Then, if $\omega T_5 \stackrel{\alpha \ell}{\cong}$ is assumed to be less than unity, $(\omega T_5 \stackrel{\alpha \ell}{\cong})^2 \ll 1$. The appropriate expansion for the arctangent terms in equation B-9 when $x^2 < 1$ (where $x = \omega T_5 \stackrel{\alpha \ell}{\equiv}$), is given By^{20}

$$TAN^{-1}(x) = x - \frac{x^3}{3} + \frac{x^5}{5} - \frac{x^7}{7} + \cdots - \cdots B^{-12}$$

SUBSTITUTING FOR X IN EQUATION B-12 GIVES THE FOLLOW-ING PARTIAL SERIES EXPANSION FOR THE BRACKETED TERM IN EQUATION B-9:

$$\left[\omega\gamma_{s}\overset{al}{\in} - \frac{(\omega\gamma_{s}\overset{al}{\in})^{3}}{3} + \cdots\right] - \left[\omega\gamma_{s} - \frac{(\omega\gamma_{s})^{3}}{3} + \cdots\right] B - 13$$

AFTER NEGLECTING THE SMALL TERMS, THE SPECTRAL INTEN-SITY OF THE SHORT-CIRCUIT DRAIN CURRENT FLUCTUATIONS AT LOW FREQUENCIES IS GIVEN BY²¹

$$S_{i(\omega)} \doteq \left[\frac{2^{GM}}{AC_{0}}\right]^{2} \cdot \frac{4kTAN_{T}(E)\gamma_{s}}{\alpha} (e^{\alpha \ell} - 1)$$
 B-14

IN WHICH IT IS EVIDENT THAT VERY LOW FREQUENCY NOISE IS WITHOUT FREQUENCY DEPENDENCE AND IS PROBABLY WHITE. THIS IS IN EXCELLENT AGREEMENT WITH OTHER NOISE THEORIES.^{22,23} FURTHERMORE, SINCE $\gamma_s \simeq 10^{-9}$ second, it is apparent that the FLAT NOISE SPECTRUM BEGINS AT $f < 10^{-9}$ Hz. Since this is far LOWER THAN IT IS POSSIBLE TO MEASURE, THE LOWER LIMIT OF FREQUENCY-DEPENDENT NOISE WILL NO DOUBT REMAIN A PURELY THEO-RETICAL LIMIT.

CASE 2. INTERMEDIATE FREQUENCY:
$$\omega \gamma_5 < 1 < \omega \gamma_5 e^{\alpha \ell}$$
.
The expansion for tan⁻¹(x) when x >1 (where $x = \omega \gamma_5 e^{\alpha \ell}$) is given by²⁴

$$TAN^{-1}(x) = \frac{n}{2} - \frac{1}{x} + \frac{1}{3x^3} - \frac{1}{5x^5} + \cdots - \cdots B - 15$$

WITH $e^{\alpha \ell} = 10^{17}$, ω_{7} remains much less than unity. Therefore, substituting $\omega_{7} e^{\alpha \ell}$ for x in equation B-15 and ω_{7} for x in equation B-12 gives the following partial series expansion for the arctangent terms in equation B-9:

$$\left[\frac{\overline{n}}{2} - \frac{1}{\omega \gamma_{s}} e^{\alpha t}\right] - \left[\omega \gamma_{\overline{s}} - \frac{(\omega \gamma_{\overline{s}})^{3}}{3} + \cdots\right] \qquad B-16$$

AFTER NEGLECTING THE SMALL TERMS, THE SPECTRAL INTEN-SITY OF THE SHORT-CIRCUIT DRAIN CURRENT FLUCTUATIONS AT INTERMEDIATE FREQUENCIES IS GIVEN BY²⁵

$$S_{i(\omega)} \doteq \left[\frac{g_{GM}}{AC_{0}}\right]^{2} \cdot \frac{kTAN_{T}(E)}{\alpha f}$$
 B-17

which is inversely proportional to frequency. With $\gamma_{\rm S} \simeq 10^{-9}$ second, the full range of "1/F" noise is within $10^{-9} < f < 10^8$. Measurements made at ten cycles and at one kilocycles were therefore well within the accepted range of frequencies. As for equation B-14, the noise intensity is directly proportional to the surface-state density. This fact has been used to correlate saturation drain conductance with 1/F noise intensity in this research. <u>Case 3.</u> High frequency: $\omega \gamma_5 > 1$. Using the expression of equation B-15 for both $\omega \gamma_5 \in e^{\alpha \ell}$ and $\omega \gamma_5$ gives the following partial series expansion for the arctangent terms in equation B-9:

$$\left[\frac{n}{2} - \frac{1}{\omega \tau_{s}} e^{\alpha \ell} + \cdots - \cdots\right] - \left[\frac{n}{2} - \frac{1}{\omega \tau_{s}} + \cdots - \cdots\right] \quad B-18$$

FROM WHICH THE SPECTRAL INTENSITY OF THE SHORT-CIRCUIT DRAIN CURRENT FLUCTUATIONS AT HIGH FREQUENCIES MAY BE GIVEN BY²⁶

$$S_{i(\omega)} \doteq \left[\frac{g G_M}{A C_0}\right]^2 \cdot \frac{4kTAN_T(E)}{\alpha \omega^2 \gamma_s} (1 - e^{-\alpha \ell})$$
 B-19

HIGH FREQUENCY NOISE IS THEREFORE EXPECTED TO BE INVERSELY PROPORTIONAL TO THE SQUARE OF THE FREQUENCY AT WHICH IT IS MEASURED AND, AS WITH THE NOISE INTENSITY AT INTERMEDIATE AND LOW FREQUENCIES, IS DIRECTLY PROPORTIONAL TO THE SUR-FACE-STATE DENSITY.

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²⁴<u>Op</u>. <u>Cit</u>., <u>Tables</u>, p. 433. ²⁵<u>Op</u>. <u>Cit</u>., Hsu, "MOS Transistors," p. 1453. ²⁶<u>IBID</u>., p. 1453.

APPENDIX C

THRESHOLD VOLTAGE DRIFT

IN P-CHANNEL ENHANCEMENT MOSFETS

SPECIFICATIONS

	DD07P	DD08K
Low noise units	56 5 7 77 79 80	103 109 110 112 121
High noise units	45 47 48 60 61	102 104 105 114 120
Rated V _{GS} Gate-to-source	-30V. Zener diode protected	±70V. No gate protection
Applied V _{GS} Gate-to-source Drain, source, and Substrate shorted	-25V.	-50V.
Oven Temperature	130±5°C	130±5°C
TIME DURATION	500 н <mark>ours</mark>	500 Hours

CONNECTION DIAGRAM

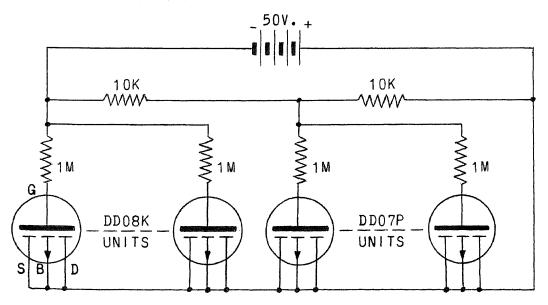


FIGURE C-1. CONNECTION DIAGRAM FOR BLASING P-CHANNEL MOS TRANSISTORS FOR NEGATIVE THRESHOLD VOLTAGE DRIFT.

LIFE-TEST DATA FOR DD07P

-V _{T1}	THRESHOLD	VOLTAGE	AS	OF	6/25/73
-V _{T2}	THRESHOLD	VOLTAGE	AS	0 F	7/17/73

UNIT NO.	-V _{T1}	-V _{T2}	% ∆V _T
*****	VOLTS	VOLTS	an a
56	4.60	4.72	2.61
57	4.51	4.63	2.66
77	4.49	4.59	2.23
79	4.56	4.64	1.75
80	4.70	4.81	2.34
45	4.52	4.76	5.32
47	4.58	*	*
48	4.44	4.58	3.16
60	4.87	5.20	6.78
61	4.62	4.95	7.15

* UNIT FAILED DURING BIAS AT ELEVATED TEMPERATURE

INTERPRETATION OF DATA

MEAN $-V_{T1} = 4.59V$. Standard deviation =	= 0.037V.
MEAN $-V_{T2} = 4.63V$. Standard deviation =	= 0.196V.
ABSOLUTE DIFFERENCE BETWEEN MEANS:	0.04V.
STANDARD ERROR OF THE DIFFERENCE BETWEEN MEANS:	0.08 87V .

Since the absolute difference between means is less than the standard error of the difference between means, it can be concluded that there was no significant change in threshold voltage for type DD07P MOSFET's after bias at $130\pm5^{\circ}$ C for 500 hours.

LIFE-TEST DATA FOR DD08K

-V _{T1}	THRESHOLD	VOLTAGE	AS	OF	6/25/73
-V _{T2}	THRESHOLD	VOLTAGE	AS	OF	7/17/73

UNIT NO.	-V _{T1} volts	-VT2 VOLTS	%∆V _T
103	4.27	4.59	7.50
109	4.17	4.55	9.11
110	4.19	*	*
112	4.21	4.50	6.88
121	4.31	4.61	6.95
102	4.11	*	*
104	3.82	4.26	11.5
105	4.16	4.66	12.0
114	4.15	5.13	23.6
120	4.17	4.76	14.1

* UNIT FAILED DURING BIAS AT ELEVATED TEMPERATURE

INTERPRETATION OF DATA

MEAN $-V_{T1} = 4.16V$. Standard deviation =	= 0.044V.
MEAN $-V_{T2} = 4.62V$. Standard deviation =	= 0.073V.
ABSOLUTE DIFFERENCE BETWEEN MEANS:	0.46V.
Standard error of the difference between means:	0.0945V.

Since three times the standard error of the difference between means is less than the absolute difference between means, the population after threshold voltage drift has less than 0.27% chance of being part of the original population. Therefore, it can be concluded that the type DD08K MOSFET's showed substantial negative threshold voltage drift after bias at 130±5°C for 500 hours.

APPENDIX D

EXCESS NOISE-DRAIN CURRENT MEASUREMENTS

AT CONSTANT GATE VOLTAGE

SPECIFICATIONS

	<u>DD07P</u>	DD08K
Low noise units	56 57 77 79 80	103 109 110 112 121
HIGH NOISE UNITS	45 47 48 60 61	102 104 105 114 120
Rated VGS Gate-to-source	-30V. Zener diode protected	±70V. No gate protection
APPLIED VGS Gate-to-source	-5.5V.	-5.5V.
RATED IDS Drain-to-source VGS = VDS = -15V.	-44.0mA	-44.0mA
Max. Applied IDS V _{GS} = -5.5V. VDS variable VSsub = 0V.	-8.OmA	-8.ОмА
RATED DISSIPATION $T_A = 25^{\circ} C$	225м₩	225mW
AMBIENT TEMPERATURE	25±5°C	25±5°C
INSTANTANEOUS <i>ga</i> sat at drain breakdown	100 мноѕ	100 мноз

CONNECTION DIAGRAM

REFER TO FIGURE 3-1 ON PAGE 42 OF TEXT.

		DUVIN	i conne		IA FU	טעע חי	(E MUS	FEIS	(6/25/
Noise	Volt	AGE I An	VOL D Spec	TS AT	I n d i Freq	CATED UENCIE	Drain Is	Curr	ENTS
		0.5м	A		1.0мА	4		1.5м/	Ą
UNIT NO.	10Hz	1KHz	-V _{DS}	10Hz	1KHz	-V _{DS}	10Hz	1 KHz	-V _{DS}
45 47 48 60 61	44 42 35 72 47	5.7 5.4 5.5 7.2 6.0	15.5 13.0 6.5 19.5 13.5	69 63 66 110 76	8.2	24.0 23.5 14.0 26.0 20.5	100 79 76 120 92	10 11 9.8 12 9.2	
56 57 77 79 80	28 31 31 28 31	6.6 6.3 6.0 6.3 6.6	39.0 36.5 22.5 34.0 44.0	51 44	9.8 8.7	58.0 60.0 44.5 57.0 64.5	60 66 77 52 50	11 12 12 12 11	69.5
		2.0м/	Ą		З.ОмА	۱.		4.0n	١A
45 47 48 60 61	130 110 110 130 120	14 13 13 14 13	36.0 26.5	140 130 130 140 125		41.5 43.0 33.0 36.0 33.0	110 80 110 130 120	15 17	49.0 56.0 40.0 42.5 43.0
56 57 77 79 80	70 60 62 73 60	11 12 11 13 12	70.0 71.0 65.0 76.5 77.0	60 62 55 60 54	10 11 10 10 10	75.5 78.0 74.5 80.0 81.0	50 45 43 47 39	9 10 9 8 8	80.0 83.5 78.0 83.0 82.0
		6.0мА	N		8.0мА				
45 47 48 60 61	100 70 100 100 95	15 13 16 16 15	53.0 58.0 46.0 46.5 52.0	88 50 80 79 81	14 12 14 15 14	60.5 63.0 54.5 51.5 57.0			
56 57 77 79 80	17 13 13 12 11		83.5 87.0 80.5 84.0 84.0						

EXCESS NOISE-DRAIN CURRENT DATA FOR DD07P MOSFETS (6/25/73)

EXCESS NOISE-DRAIN CURRENT DATA FOR DD07P MOSFETS (7/17/73)

Noise Voltage in wolts at Indicated Drain Currents and Specified Frequencies

		0.5мА	١		1.0мА	١		1.5мА	
UNIT NO.	10Hz	1KHz	-V _{DS}	10Hz	1KHz	-V _{DS}	10Hz	1KHz	-V _{DS}
45 47 48 60 61	90 *Unit 60 100 60	45 FAIL 50 80 50	22.5 ED AFT 7.0 23.0 16.0	145 ER BI 140 170 100	90 ASAT 110 130 80	32.0 130±5 18.5 29.0 24.5	180 •C Fo 180 200 120	110 R 500 140 140 90	36.0 HOURS 26.0 31.0 30.0
56 57 77 79 80	80 80 60 80	60 60 45 60	39.5 38.0 23.5 36.0 47.0	1 30 1 20 1 40 1 40 1 40	100 100 110 100 110	59.0 68.0 45.0 61.5 68.0	180 160 180 180 190	140 130 140 130 160	69.0 73.5 57.0 71.0 75.0

	2.OmA			3.0mA			4 • 0 m A		
45	210	120	39.0	220	130	44.5	225	145	49.5
47	*	*	*	*	*	*	*	*	*
48	210	160	30.0	250	180	37.0	270	200	40.0
60	210	160	34.0	225	170	38.0	230	185	42.0
61	130	100	33.0	140	110	38.0	130	100	44.0
56	210	160	73.0	230	190	78.5	140	110	81.5
57	200	160	76.5	210	170	79.0	90	80	83.0
77	210	150	65.0	240	180	74.0	200	100	77.0
79	200	160	76.5	230	180	80.0	130	90	82.0
80	210	150	78.0	200	150	80.0	110	100	81.5

		6.0м	Ą	8.0mA		
45	205	100	54.0	190	90	63.0
47	*	*	*	*	*	*
48	240	170	46.5	200	140	55.0
60	240	1 90	46.0	200	110	50.5
61	120	90	51.0	90	70	59.0
56 57 77 79 80	32 24 26 22 32	16 15 19 18 17	82.0 84.0 79.0 83.0 83.5			

EXCESS NO	01 SE-	DRAIN	CURRE	NT DA	TA FC	R DDOE	K MOS	FETS	(6/25/7
Noise	Volt	AGE I AN	N سر D Spec	TS AT	INDI Freq	CATED	DRAIN s	Curr	ENTS
		0.5м/	4		1.0м	A		1.5м	A
UNIT NO.	10Hz	1KHz	-V _{DS}	10Hz	1 KHz	-VDS	10Hz	1 KHz	-VDS
102 104 105 114 120	55 30 50 36 55	6.4 5.2 4.8 4.4 4.7	0.60 1.00 5.00 7.50 18.0	75 56 71 60 80	14 6.8 9.5 9.3 9.3	23.0 14.5 28.0 64.0 62.0	110 71 82 64 110	17 11 13 12 12	43.0 27.0 46.5 69.0 73.0
103 109 110 112 121			25.0 9.50 20.0 8.00 16.5	72 80 76 68 67	11 11 10 10 10	57.0 21.0 38.5 24.0 48.5	83 98 92 90 87	12 12 12 13 11	72.5 30.0 55.5 38.0 64.5
		2.0м/	A		3.Ом/	Ą		4.Ом/	Ą
102 104 105 114 120	130 110 120 100 125	20 12 17 18 16	56.0 30.5 58.0 74.5 75.0	150 130 140 145 135	25 15 19 22 17	65.0 46.5 70.0 76.0 77.5	155 145 150 160 130	24 16 18 21 16	69.0 58.0 73.0 77.0 79.0
103 109 110 112 121	76 84 110 100 95	12 13 10 12 13	74.5 38.0 69.0 47.0 72.0	49 34 78 60 50	10 11 7.9 7.1 12	77.0 50.0 76.0 61.0 77.0	36 22 60 30 33	8.2 9.3 6.2 4.2 10	77.5 57.0 77.0 70.5 79.0
		6.ОмА	l.		8.0мА	Ą			
102 104 105 114 120	145 140 140 155 120	20 15 16 18 13	75.0 70.0 75.0 78.5 79.5	110 105 100 110 90	17 14 13 16 11	77.5 73.0 76.0 79.0 80.0			
103 109 110 112 121	32 15 40 14 12	4.7 5.1 4.3 3.9 8.1	78.0 65.0 78.5 76.0 80.0	43 32 50 27 29	6.8 7.2 5.8 5.0 8.5	78.5 69.0 79.5 76.5 80.0			

EXCESS	NOI	SE-DRAIN	CURRENT	DATA	FOR	DD08K	MOSFETS	(6/25/73)
					the second states and second			التجاجب وعدران بجرائل سوالتهمية الجنوان والمتحدة

EXCESS NOISE-DRAIN	CURRENT	DATA	FOR	DD08K	MOSFETS ((7)	/17	/73)
--------------------	---------	------	-----	-------	-----------	-----	-----	-----	---

Noise	VOLT		VOL مر D Spec					Curr	ENTS
		0.5мА	١		1.ОмА	١		1.5мА	N .
UNIT NO.	10Hz	1KHz	-V _{DS}	10Hz	1KHz	-V _{DS}	10Hz	1KHz	-V _{DS}
102 104 105 114 120	*Unit 100 110 130 120	FAIL 80 100 110 90	ED AFT 5.50 13.0 10.5 34.0	ER BI 160 200 220 180	AS AT 130 160 170 110	130±5 18.5 38.0 66.0 71.0	о ^о С го 210 240 250 220	R 500 170 200 190 140	HOURS 29.5 51.0 71.5 76.0
103 109 110 112 121	100 100 100 90 100	90 90 80 55 60	27.5 11.0 23.0 9.00 19.0	180 160 180 170 160	130 110 160 100 110	60.0 22.5 35.0 26.0 50.0	200 210 220 190 180	150 150 180 180 140	72.0 30.0 49.0 38.5 66.0

		2.0мА	١		3.Ом	A		4.0м	٩
102 104 105 114 120	* 250 270 300 270	* 200 210 220 200	* 32.0 60.5 78.0 78.0	* 300 310 320 240	* 230 240 230 160	* 46.0 70.5 79.5 78.5	* 290 270 280 170	* 200 230 200 100	* 57.0 74.0 78.0 79.0
103 109 110 ** 112 121	180 200 Unit 210 120	120 130 DEFE 160 90	75.0 39.0 CTIVE 46.5 73.0	110 110 AT HI 170 100	80 75 gher 100 70	77.0 49.0 Drain 61.5 78.0	78 41 curre 90 76	55 24 NTS 60 45	78.0 57.5 70.0 79.5

		6.0м/	4	8 .0mA			
102 104 105 114 120	* 250 200 210 140	* 190 180 110 90	* 76.0 79.0 79.5	* 170 120 90 89	* 140 80 50 65	* 75.0 77.0 79.5 80.5	
103 109 110 112 121	48 30 ** 21 32	18 20 ** 19 22	80.0 69.0 ** 75.0 80.0	34 17 ** 18 22	12 16 ** 12 12	80.5 72.0 ** 76.0 80.5	

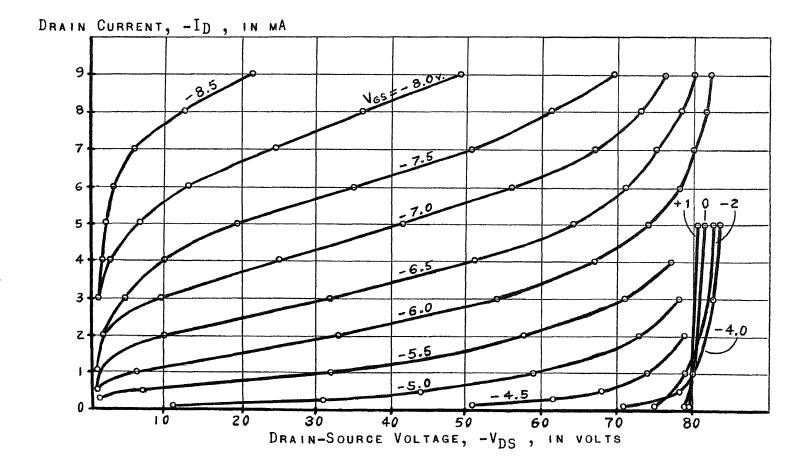


FIGURE D-1. CURRENT-VOLTAGE CHARACTERISTICS FOR A RELATIVELY LOW NOISE DD07P MOSFET DRAIN-DIODE BREAKDOWN IS EVIDENT WHEN THE DEVICE IS BIASED IN CUTOFF. UNIT NO. 44.

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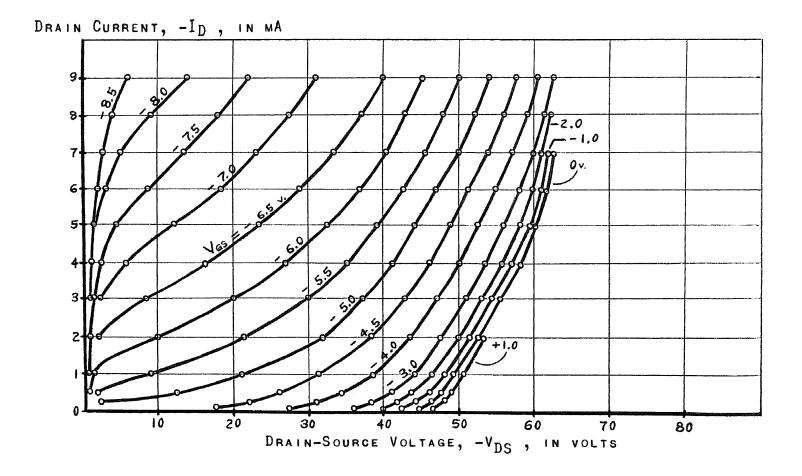


FIGURE D-2. CURRENT-VOLTAGE CHARACTERISTICS FOR A RELATIVELY HIGH NOISE DD07P MOSFET HIGH SATURATION DRAIN CONDUCTANCE AND EVIDENCE OF DRAIN PUNCH-THROUGH WERE TYPICAL OF such high noise transistors. Unit no. 69.

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APPENDIX E

DRAIN LEAKAGE CURRENT MEASUREMENTS

DEFINITION OF IDSS

The short-circuit drain leakage current, I_{DSS} , of pchannel enhancement MOSFET's is defined at $V_{DS} = -10$ volts and $V_{GS} = V_{SSUB} = 0$ volts. It is a measure of the leakage around the entire drain-to-substrate p-n junction and is usually no more than a few nanoamperes $(10^{-9} \text{ amperes})$ in most enhancement-mode devices.

EXPERIMENTAL METHOD

IN ORDER TO ASSESS THE POSSIBLE EFFECTS OF DRAIN LEAK-AGE ON THE SATURATION I-V CHARACTERISTICS OF P-CHANNEL EN-HANCEMENT MOSFET'S AND IN ORDER TO CORRELATE THIS LEAKAGE WITH CHANGES IN THE SURFACE-STATE DENSITY AND OTHER CHARAC-TERISTICS OF THE SI-SIO₂ INTERFACE, MEASUREMENTS OF I_{DSS} WERE CONDUCTED ON ALL UNITS LIFE-TESTED FOR THRESHOLD VOL-TAGE DRIFT, BOTH BEFORE AND AFTER BIAS AT 130±5°C FOR 500 HOURS.

Figure E-1 shows a simplified schematic of the <u>General</u> <u>Radio Type 1230-A DC Amplifier and Electrometer</u> which was used to measure I_{DSS} for the units tested. The instrument consisted of variable resistors R_A (adjustable from 10^4 to 10^{11} ohms in eight orders of magnitude) and R_B (preset for voltage measurements) and electrometer tube V_1 connected as

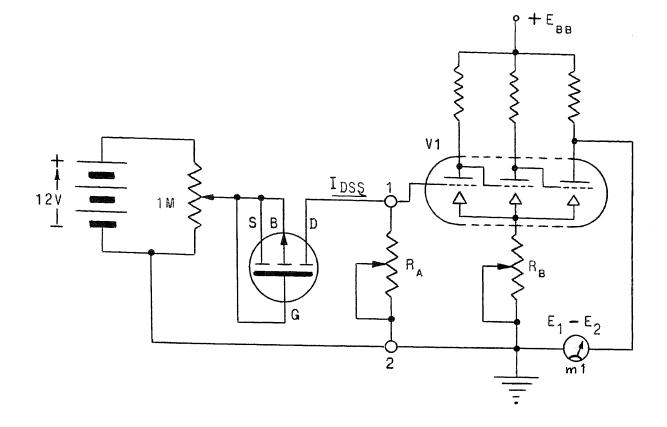


FIGURE E-1. SIMPLIFIED SCHEMATIC OF A GENERAL RADIO TYPE 1230-A DC AMPLIFIER AND ELECTROMETER AND EXTERNAL DC SUP-PLY USED TO MEASURE IDSS OF P-CHANNEL ENHANCEMENT MOSFET'S. TUBE V1 WAS CONNECTED AS A THREE-STAGE, DIRECT-COUPLED, VOLTAGE AMPLIFIER WHICH, DUE TO THE COMMON-CATHODE CON-FIGURATION, PROVIDED AN INPUT RESISTANCE GREATER THAN 10^{14} OHMS AND WHICH DROVE METER M1. DUE TO THE EXTREMELY HIGH TRANSCONDUCTANCE OF THE ELECTROMETER AMPLIFIER, VOLTAGE VARIATIONS ACROSS R_A (ADJUSTABLE BETWEEN 10^4 AND 1011_Ω). PRODUCED ESSENTIALLY IDENTICAL VARIATIONS ACROSS R_B . THE LEAKAGE CURRENT PRODUCED A POTENTIAL DROP ACROSS TERMINALS 1 AND 2 AND WAS FOUND TO BE $(E_1-E_2)/R_A$. A THREE-STAGE, DIRECT-COUPLED, CATHODE-FOLLOWER WHICH DROVE A 1MÅ METER AND WHICH PROVIDED AN INPUT RESISTANCE GREATER THAN 10^{14} ohms (open-grid). Tube V₁ was treated with General Electric Dri-Film SC-87 and all insulation was teflon for low leakage. A special shielded enclosure was used to cover the test jig and all measurements were conducted within a shielded room at a temperature of $25\pm5^{\circ}$ C with not more than 30% humidity.

Two SIX-VOLT NICKEL-CADMIUM BATTERIES AND A POTENTIO-METER WERE USED TO PROVIDE THE PROPER DRAIN BIAS. THE MOSFET UNDER TEST WAS CONNECTED WITH GATE, SOURCE AND SUB-STRATE TO THE WIPER-ARM OF THE POTENTIOMETER AND WITH THE DRAIN TO THE TERMINAL 1 VOLTAGE INPUT OF THE METER.

When measuring $I_{\rm DSS}$, the leakage current which flows developes a voltage across input resistance $R_{\rm A}$. Since the overall transconductance of the electrometer amplifier was so great, any voltage change across cathode resistance $R_{\rm B}$ was essentially the same as that across $R_{\rm A}$ and was read directly on the meter. The unknown leakage current was then found from $(E_1 - E_2)/R_{\rm A}$, where $R_{\rm A}$ was adjusted to provide a large deflection on the 100mV scale of meter M1.

DRAIN LEAKAGE CURRENT (I_{DSS}) MEASUREMENTS FOR TEN TYPE DD07P and ten type DD08K MOSFET'S is listed on the following pages.

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DSS1	Drain leakage Drain leakage		
UNIT NO.	IDSS1 PA	DSS2 PA	
56	1000	4000	
57	1000	4200	
77	1100	4300	RELATIVELY LOW NOISE
79	1000	4300	TRANSISTORS
80	900	4100	
45	1100	3900	
47	1600	4300	
48	1000	4000	RELATIVELY HIGH NOISE
60	1200	4200	TRANSISTORS
61	11 0 0	4100	

DRAIN LEAKAGE DATA FOR DD07P MOSFET'S

I DSS1 I DSS2	Drain leakage Drain leakage		
UNIT NO.	IDSS1 PA	IDSS2 PA	
103	1200	4900	
109	1100	4800	
110	900	60 A*	
112	800	4600	TRANSISTORS
121	900	4700	
102	1100	.6 A*	
104	1600	4600	
105	1100	4500	RELATIVELY HIGH NOISE
114	1600	4400	TRANSISOTRS
120	1 400	4900	

DRAIN LEAKAGE DATA FOR DDO8K MOSFET'S

*Units 102 and 110 failed during bias at 130±5°C for 500 hours.

APPENDIX F

EQUIPMENT USED IN THIS RESEARCH



FIGURE F-1. PHOTOGRAPH SHOWING THE PRINCIPAL INSTRUMENTS USED IN THIS RESEARCH. AT CENTER ON TOP: DIGITAL MULTI-METER (FLUKE, MODEL 8000-A) AND ELECTRONIC VOLTMETER (GENERAL RADIO, TYPE 1808-A). LEFT TO RIGHT: TRANSISTOR CURVE TRACER (TETRONIX, TYPE 527 WITH ADAPTOR FOR P-CHAN-NEL MOS TRANSISTORS AND SHOWING TYPICAL CURRENT-VOLTAGE CHARACTERISTICS ON SCREEN); DC AMPLIFIER AND ELECTROMETER (GENERAL RADIO, TYPE 1230-A WITH SHIELDED ENCLOSURE AND TEST JIG NOT VISIBLE); AND DIODE NOISE ANALYZER (QUAN-TECH LABORATORIES, MODEL 327). FRONT, AT CENTER: VARIABLE BIAS POTENTIOMETER FOR GATE CIRCUIT.

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