

# FORESAIL-1 cubesat mission to measure radiation belt losses and demonstrate de-orbiting

M. Palmroth<sup>1,2</sup>, J. Praks<sup>3</sup>, R. Vainio<sup>4</sup>, P. Janhunen<sup>2</sup>, E. K. J. Kilpua<sup>1</sup>, N. Yu. Ganushkina<sup>2,5</sup>, A. Afanasiev<sup>4</sup>, M. Ala-Lahti<sup>1</sup>, A. Alho<sup>3</sup>, T. Asikainen<sup>6</sup>, E. Asvestari<sup>1</sup>, M. Battarbee<sup>1</sup>, A. Binios<sup>3</sup>, A. Bosser<sup>3</sup>, T. Brito<sup>1</sup>, J. Envall<sup>2</sup>, U. Ganse<sup>1</sup>, H. George<sup>1</sup>, J. Gieseler<sup>4</sup>, S. Good<sup>1</sup>, M. Grandin<sup>1</sup>, S. Haslam<sup>2</sup>, H.-P. Hedman<sup>7</sup>, H. Hietala<sup>4</sup>, N. Jovanovic<sup>3</sup>, S. Kakakhel<sup>7</sup>, M. Kalliokoski<sup>1</sup>, V. V. Kettunen<sup>3</sup>, T. Koskela<sup>1,4</sup>, E. Lumme<sup>1</sup>, M. Meskanen<sup>2</sup>, D. Morosan<sup>1</sup>, M. Rizwan Mughal<sup>3</sup>, P. Niemelä<sup>3</sup>, S. Nyman<sup>3</sup>, P. Oleynik<sup>4</sup>, A. Osmane<sup>1</sup>, E. Palmerio<sup>1</sup>, Y. Pfau-Kempf<sup>1</sup>, J. Peltonen<sup>4</sup>, J. Plosila<sup>7</sup>, J. Polkko<sup>2</sup>, S. Poluianov<sup>6,8</sup>, J. Pomoell<sup>1</sup>, D. Price<sup>1</sup>, A. Punkkinen<sup>4</sup>, R. Punkkinen<sup>7</sup>, B. Riwanto<sup>3</sup>, L. Salomaa<sup>7</sup>, A. Slavinskis<sup>3,9</sup>, T. Sääntti<sup>7</sup>, J. Tammi<sup>7</sup>, H. Tenhunen<sup>7</sup>, P. Toivanen<sup>2</sup>, J. Tuominen<sup>7</sup>, L. Turc<sup>1</sup>, E. Valtonen<sup>4</sup>, P. Virtanen<sup>4</sup>, T. Westerlund<sup>7</sup>

<sup>1</sup>University of Helsinki, Department of Physics, Helsinki, Finland

<sup>2</sup>Finnish Meteorological Institute, Space and Earth Observation Centre, Helsinki, Finland

<sup>3</sup>Aalto University, School of Electrical Engineering, Espoo, Finland

<sup>4</sup>University of Turku, Department of Physics and Astronomy, Turku, Finland

<sup>5</sup>University of Michigan, Department of Climate and Space Sciences and Engineering, Ann Arbor, USA

<sup>6</sup>University of Oulu, Space Climate Research Unit, Oulu, Finland

<sup>7</sup>University of Turku, Department of Future Technologies, Turku, Finland

<sup>8</sup>University of Oulu, Sodankylä Geophysical Observatory (Oulu Unit), Oulu, Finland

<sup>9</sup>University of Tartu, Tartu Observatory, Tõravere, Estonia

## Key Points:

- FORESAIL-1 mission measures energetic electron precipitation and solar energetic neutral atom flux
- We will demonstrate a cost-efficient de-orbiting and orbit manoeuvring technology without propellants
- The goal of the mission is to contribute significantly to the sustainable utilisation of space

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Corresponding author: Minna Palmroth, [minna.palmroth@helsinki.fi](mailto:minna.palmroth@helsinki.fi)

## Abstract

Today, the near-Earth space is facing a paradigm change as the number of new spacecraft is literally sky-rocketing. Increasing numbers of small satellites threaten the sustainable use of space, as without removal, space debris will eventually make certain critical orbits unusable. A central factor affecting small spacecraft health and leading to debris is the radiation environment, which is unpredictable due to an incomplete understanding of the near-Earth radiation environment itself and its variability driven by the solar wind and outer magnetosphere. This paper presents the FORESAIL-1 nanosatellite mission, having two scientific and one technological objectives. The first scientific objective is to measure the energy and flux of energetic particle loss to the atmosphere with a representative energy and pitch angle resolution over a wide range of magnetic local times. To pave the way to novel model - *in situ* data comparisons, we also show preliminary results on precipitating electron fluxes obtained with the new global hybrid-Vlasov simulation Vlasiator. The second scientific objective of the FORESAIL-1 mission is to measure energetic neutral atoms (ENAs) of solar origin. The solar ENA flux has the potential to contribute importantly to the knowledge of solar eruption energy budget estimations. The technological objective is to demonstrate a satellite de-orbiting technology, and for the first time, make an orbit manoeuvre with a propellantless nanosatellite. FORESAIL-1 will demonstrate the potential for nanosatellites to make important scientific contributions as well as promote the sustainable utilisation of space by using a cost-efficient de-orbiting technology.

## 1 Introduction

Unprecedented numbers of new spacecraft are now being launched into Earth orbit to satisfy the growing demand from the scientific, commercial, and military sectors. Most of these new spacecraft need to survive in the radiation belts (RBs; van Allen & Frank, 1959), which are regions of trapped energetic charged particles. The RBs are critical in terms of space weather, as the radiation ages the spacecraft and deteriorates hardware. All new satellites contribute to the already existing large number of orbital debris, if they are not actively removed at the end of the mission. This section introduces the state of the art in the three scientific and technological objectives of the FORESAIL-1 mission: measurements of energetic particle precipitation and solar energetic neutral atoms (ENAs), and de-orbiting technologies.

### 1.1 State of the Art: Electron precipitation observations

The RBs are produced by a complex balance of particle source and loss processes that vary both temporally and spatially (e.g., Tverskoy, 1969; Walt, 1996; Chen et al., 2007; Shprits et al., 2008; Thorne, 2010). Significant variations in electron fluxes occur over various time scales as a function of both energy and distance, driven by solar-magnetospheric interactions and internal magnetospheric processes (e.g., Li et al., 1997; Elkington et al., 2003; Reeves et al., 2003; Shprits et al., 2006; Baker & Kanekal, 2008). Effective losses from the outer radiation belts consist of 1) loss through the outer edge of the magnetosphere (magnetopause shadowing (e.g., West et al., 1972; Ukhorskiy et al., 2006; Saito et al., 2010; Matsumura et al., 2011; Turner et al., 2014)), 2) radial outward displacement of the electrons due to waves (Mann et al., 2016), and weakening of the Earth's magnetic field (the  $D_{st}$  effect (McIlwain (1966); Kim & Chan (1997); Millan & Thorne (2007))), and 3) wave-particle interactions resulting in scattering of particles into the loss cone (Kennel & Petschek, 1966; Thorne & Kennel, 1971; Thorne, 1974). There are no comprehensive estimates about which of these processes is most important during different conditions, while it is clear that particle losses play a central part in regulating the RBs.

80 The wave-particle interactions leading to losses from the RBs vary on time scales  
81 ranging from 100 milliseconds to several minutes (Millan & Thorne, 2007). Balloon ex-  
82 periments have historically been the earliest method to determine this loss category by  
83 measuring the X-rays from bremsstrahlung radiation induced by the interaction of pre-  
84 cipitating particles with the neutrals in the upper atmosphere (Barcus et al., 1973; Pytte  
85 et al., 1976). Latest such observations are provided by the BARREL mission (Millan &  
86 the BARREL Team, 2011; Woodger et al., 2015). All balloon missions are constrained  
87 to balloon-reachable altitudes and thus only allow indirect observation of the precipitat-  
88 ing particles.

89 Energetic particle precipitation has also been observed from the ground, as precip-  
90 itating electrons with energies over several tens of keV cause enhanced ionisation in the  
91 ionospheric D-region at an altitude of about 90 km. Relative ionospheric opacity meters  
92 (riometers) (Hargreaves, 1969) are ground-based passive radars measuring the so-called  
93 cosmic noise absorption (Shain, 1951), which corresponds to absorbed radio wave power  
94 in the ionosphere resulting from enhanced D-region electron density. Recently, interfer-  
95 ometric riometry has been developed to produce all-sky maps (e.g. McKay-Bukowski et  
96 al., 2015). Ground-based observations of energetic electron precipitation can also be achieved  
97 using incoherent scatter radars, which can accurately measure D-region electron density  
98 down to about 70 km altitude (e.g. Miyoshi et al., 2015). However, the intrinsically in-  
99 direct ground-based observations do not allow inferring precipitating particle fluxes and  
100 energies, even using newly developed approaches such as spectral riometry (Kero et al.,  
101 2014). Hence having direct measurements of precipitating fluxes from space-borne in-  
102 struments is critical for radiation belt loss studies.

103 One of the first satellite missions to study energetic electron precipitation was the  
104 Solar, Anomalous, and Magnetospheric Particle Explorer (SAMPEX, 1992 - 2012) used  
105 in a number of studies (Li et al., 2001; Tu et al., 2010; Nakamura et al., 2000). DEME-  
106 TER microsatellite observed electron fluxes at energies between 70 keV and 2.5 MeV with  
107 high energy resolution (256 channels) on a 700 km orbit (Sauvaud et al., 2006). These  
108 observations have been used to infer energetic electron precipitation (Graf et al., 2009),  
109 although DEMETER viewed primarily trapped particles. The main data set of direct  
110 measurements of precipitating energetic particles comes from the Medium Energy Pro-  
111 ton and Electron Detector (MEPED) onboard NOAA/POES satellites (D. S. Evans &  
112 Greer, 2000). MEPED consists of two telescopes, the  $0^\circ$  telescope designed to measure  
113 precipitating particle fluxes and the  $90^\circ$  telescope for trapped particle fluxes, measur-  
114 ing electrons in three energy channels ( $>30$  keV,  $>100$  keV, and  $>300$  keV) and protons  
115 in five energy channels ( $>30$  keV,  $>80$  keV,  $>250$  keV,  $>800$  keV, and  $>2.5$  MeV). How-  
116 ever, the NOAA/POES particle data suffer from two issues. First, the  $0^\circ$  telescope only  
117 partially views the bounce loss cone and does not offer any angular resolution inside its  
118 viewcone leading to poor pitch angle resolution. This leads to an underestimation of the  
119 precipitating fluxes (Rodger et al., 2013). Second, the electron channels are affected by  
120 proton contamination; partly corrected by a new data set by Asikainen & Mursula (2013).

121 Particle precipitation is a key element in magnetosphere–ionosphere–thermosphere  
122 coupling, and therefore a crucial objective for research in numerical models, especially  
123 as there is an increasing demand for space weather forecasting capabilities. The first at-  
124 tempts to model precipitating particle fluxes relied on statistical patterns inferred from  
125 satellite observations. McDiarmid et al. (1975) produced a model for precipitating elec-  
126 tron flux within 0.15–200 keV as a function of magnetic local time (MLT) and invari-  
127 ant latitude based on about 1100 passes of the ISIS 2 spacecraft. Using data measured  
128 by the Low Energy Electron experiment onboard the Atmosphere Explorer C and D satel-  
129 lites, Spiro et al. (1982) parametrised precipitating electron energy flux and average en-  
130 ergy as a function of MLT, geomagnetic latitude, and geomagnetic activity measured by  
131 the  $Kp$  and  $AE$  indices. One of the reference models for auroral-energy electron precip-  
132 itation is the Hardy et al. (1985) model, empirically derived by compiling several years

133 of observations from the Defense Meteorological Satellite Program and Satellite (DMSP)  
134 Test Program spacecraft. The Hardy model is parametrised by the  $Kp$  index, and it is  
135 still used to provide precipitation input in the 50 eV–20 keV energy range to state-of-  
136 the-art ionospheric models (e.g., Marchaudon & Blelly, 2015). At higher energies (30 keV–  
137 1 MeV), the recent van de Kamp et al. (2016) model provides energy-flux spectra of pre-  
138 cipitating electrons as a function of  $L$  parameter and geomagnetic activity rendered with  
139 the  $Ap$  index. This empirical model was obtained from a statistical analysis of 11 years  
140 of NOAA/POES energetic electron precipitation observations and is averaged across all  
141 MLTs in its present version.

142 Modelling particle precipitation using first-principle models is not easy, given that  
143 many processes operating at spatial and temporal scales spanning many orders of mag-  
144 nitude are at play in the inner magnetosphere (energisation, and loss-cone scattering pro-  
145 cesses, among others). The emergence of global kinetic magnetospheric codes may en-  
146 able addressing this issue in the near future. Recently, a preliminary run was performed  
147 using the Vlasiator code (von Althaus et al., 2014; Palmroth et al., 2018), in which elec-  
148 trons were added as a modelled species during a substorm-time, polar-plane global mag-  
149 netospheric run throughout the magnetospheric simulation box. Figure 1 shows an ex-  
150 ample of these preliminary results of 0.1 – 60 keV electron precipitation estimation ob-  
151 tained from the analysis of the nightside velocity distribution functions of electrons at  
152 a single time step of this simulation. The top panel shows the differential number flux  
153 of precipitating electrons as a function of  $L$  shell (blue shading), as well as the mean pre-  
154 cipitating energy (black line) in the same units as typical spacecraft data. The bottom  
155 panel shows the integral energy flux as a function of  $L$ . The integral energy flux was ob-  
156 tained by multiplying the differential number flux by the corresponding energies, and in-  
157 tegrating across energies. The mean precipitating energy was calculated by dividing the  
158 integral energy flux by the total number flux (i.e., the differential number flux integrated  
159 across energies). The Hardy et al. (1985) model predicts a maximum integral energy flux  
160 of the order of  $108 - 109 \text{ keV cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$  in the midnight sector, reached at geomag-  
161 netic latitudes comprised between  $62^\circ$  and  $69^\circ$  (translating into  $L$  values between 4.5 and  
162 7.8), depending on geomagnetic activity given by the  $Kp$  index. The preliminary results  
163 from Vlasiator in Fig. 1 are therefore in reasonable agreement with those values.

164 With the expanding human activity in space, it is increasingly important to mea-  
165 sure precipitating particle fluxes *in situ* and predict the precipitation by modelling, in  
166 order to understand the Earth’s radiation environment. While previous satellite missions  
167 have provided a plethora of observations of the physical processes within the precipita-  
168 tion environment, none of the missions were designed specifically with a clear focus on  
169 precipitation. A number of new cubesat missions recently launched or being built focus  
170 on precipitation. These include, for example, the CSSWE mission (Kohnert et al., 2011),  
171 the ELFIN mission (Shprits et al., 2018), the Firebird mission Crew et al. (2016), and  
172 the AMICal Sat mission (Barthelemy et al., 2018). FORESAIL-1 will be a complemen-  
173 tary mission, improving the spatial and temporal resolution of precipitation that may  
174 be offered by these missions together.

## 175 1.2 State of the Art: Solar energetic neutral atom observations

176 The energy budgets of the solar corona and solar eruptions are major unsolved ques-  
177 tions in solar physics. Coronal heating leads to an abundance of suprathermal particles  
178 in the corona. Suprathermal ions are important for estimating the energy budget of an  
179 eruption (e.g., Emslie et al., 2004), however, they do not produce measurable amounts  
180 of electromagnetic radiation, and thus their abundance is difficult to measure. Direct in-  
181 situ measurements of suprathermal ions will be provided by the recently launched NASA  
182 Parker Solar Probe mission, but only within the outermost reaches of the solar corona.  
183 During solar eruptions these suprathermal ions can be driven to participate in charge  
184 exchange processes with neutral atoms, resulting in the formation of solar energetic neu-



**Figure 1.** Preliminary results of Vlasiator modelling of electron precipitation. (top) Differential flux of precipitating electrons as a function of  $L$  shell, in the same units as usually measured by telescopes onboard spacecraft. The black line indicates the mean precipitating energy. (bottom) Total precipitating energy flux as a function of  $L$  shell.

185 tral atoms (ENAs). So far, ENAs have been measured only during one single event (Mewaldt  
186 et al., 2009) with the IMPACT/LET instrument onboard the twin the Solar TERrestrial  
187 RELations Observatory (STEREO) spacecraft close to the beginning of the mission. Even  
188 these results may be questionable, as they have been disputed by Simnett (2011), who  
189 suggested the ENA observations could be explained by an earlier precursor event, de-  
190 tected as an electron burst.

### 191 **1.3 State of the Art: Space debris removal technologies**

192 The number of spacecraft in low-Earth orbit (LEO) is rapidly increasing, as tech-  
193 nological and regulatory changes of launchers have allowed smaller satellites. These so-  
194 called nanosatellites typically do not have propulsion systems requiring bulky or volatile  
195 propellants for orbit control or de-orbiting, making them a significant source of future  
196 orbital debris. Furthermore, if they cannot handle high-radiation environments, these  
197 nanosatellites will fail early, thus on one hand contributing to existing debris and on the  
198 other hand defeating potential plans for active deorbiting at the end of mission. The sus-  
199 tainable use of the near-Earth space has become of great interest (e.g., Bastida Virgili  
200 et al., 2016). To avoid low-Earth orbits becoming unusable in the future, also nanosatel-  
201 lites should be removed at end-of-life, otherwise the amount of space debris will increase  
202 exponentially due to collisions with bigger objects (Kessler & Cour-Palais, n.d.; Klinkrad,  
203 1993; Bradley & Wein, 2009; Bonnal et al., 2013), creating a potential danger to all later  
204 space missions. International standards have been developed to impose requirements on  
205 space missions to mitigate space debris production (e.g., European Space Agency, 2014;  
206 for Standardization, 2011). Thus, it is crucial to develop robust instruments for both con-  
207 trolling the small satellite orbits as well as for removing them from orbit after the end  
208 of the mission.

209 Apart from technological solutions for reduction of space debris that are inherent  
210 to the satellite design, efforts for third-party orbit removal techniques are ongoing. Ap-  
211 proaches include spacecraft that perform automatic rendez-vous, attachment and joint  
212 de-orbiting of larger space debris items. Using high-powered lasers (either ground- or satellite-  
213 based) to exert radiation pressure, or directly ablate the surface material of space de-  
214 bris (a so-called "Laser Broom") has been a research project in both civilian (Bekey, 1997;  
215 Phipps et al., 2012) and military (Campbell, 2000) projects. Compact propulsion meth-  
216 ods possibly allowing de-orbiting of nanosatellites include miniaturised pulsed plasma  
217 and Hall-effect thrusters, which have reached commercial technological readiness, but  
218 still require propellants and a sizeable energy budget (Tummala & Dutta, 2017). Pho-  
219 tonic solar sails have been investigated for propellantless propulsion and used success-  
220 fully for interplanetary missions (Tsuda et al., 2013), as well as multiple nanosatellite  
221 missions (Lappas et al., 2011) with mixed success. Meanwhile, electric sails, in which elec-  
222 trically charged structures interact with the ion environment, have been proposed (Jan-  
223 hunen, 2004), and suitable packages have been implemented for nanosatellites, but suc-  
224 cessful experimental verification is still outstanding.

## 25 **2 Science goals**

### 226 **2.1 Mission statement**

227 FORESAIL-1 is the first nanosatellite mission designed to measure the energy-dependent  
228 pitch angle spectra of the precipitating radiation belt particles, and solar ENA flux. Fur-  
229 ther, FORESAIL-1 will demonstrate the effectiveness of the plasma brake as a means  
230 of manipulating the spacecraft orbit in operation and lowering the spacecraft altitude  
231 to speed up de-orbiting at the end of the mission, thus addressing the sustainability of  
232 LEO space operations.



## 2.2 Science objectives

The FORESAIL-1 mission answers the following science questions: What are the pitch-angle and energy signatures of precipitation events as a function of MLT? How is the precipitation pitch-angle and energy distribution dependent on geomagnetic activity and driving from the solar wind? Thus, the FORESAIL-1 mission aims to perform precise directional measurements of electron and proton precipitation, as well as the energy spectrum and particle fluxes over a wide energy range (tens of keV to approximately one MeV). The spacecraft orbit shall drift in magnetic local time to allow determining precipitation as a function of MLT. The time resolution needs to be good enough to describe the lower bound on precipitation budget due to wave-particle interactions between chorus waves and electrons. Combining several measurements of at least three energy channels into a full pitch-angle resolution throughout the LEO region, with a time resolution of at least 15 s, will enable research of loss processes from the RBs.

To understand the energy budget of solar eruptions, the second science goal of this mission is to measure solar ENAs. This requires novel observations in an energy range well exceeding the magnetospheric ENA range. For this purpose, we use the geomagnetic field as a filter of solar particles and measure the ENA flux, thus inferring the flux that originates from the solar direction. Solar ENAs can only be measured reliably at energies exceeding the ring-current ion energies.

## 2.3 Technological objectives

In addition to the science objectives outlined above, FORESAIL-1 has a technological objective to ensure the sustainable use of space and to set a precedent for maintaining clean orbits. The objective is to test the plasma brake technology and achieve at least a 100 km lowering of the spacecraft altitude at mid-to-high altitude LEO. The consequences of this lowering of the orbit are 1) the orbital drift of the mission allowing monitoring the science objectives as a function of MLT, and 2) demonstrating that the technology can be used to de-orbit spacecraft. We will observe efficiency and performance of the plasma brake during the experiment to determine general information about the orbit lowering process. The success of the plasma brake experiment (and thus the completion of the mission's sustainability goals) is dependent on the reliable operation of the avionics, making reliability a primary design driver for FORESAIL-1. Radiation effects are identified as a major potential source of failures, hence radiation hardening techniques are used in the design.

## 3 Requirements

The study of precipitating electrons from the RBs is intrinsically coupled to the characteristic energy ranges of the electron seed populations there. Scientific and operational requirements are as follows:

1. NOAA/POES, for which the energy resolution is  $(E_2 - E_1)/E_1 = 3$  (where  $E_1$  and  $E_2$  are the upper and lower limits for consecutive integral channels), provides the lower reference bar in terms of energy resolution (Evans & Greer, 2006). The nominal energy resolution for FORESAIL-1 is 0.4 between the upper and lower limits of a channel.
2. The particle detector must have a sufficient discriminating ability between electrons and protons, such that the electron channel does not suffer from proton contamination. For lower energy channels there is no contamination, while in the higher energy channels we allow a small background, however, this should be so small that the electrons are discernible.
3. The orbit must drift to cover several MLTs.

- 281 4. For electron precipitation and solar ENA measurements, the orbital altitude should  
282 lie between 400 km and 800 km.
- 283 5. The electron pitch angle resolution should include at least three bins measured  
284 every 15 s.
- 285 6. The mission profile must allow for the use of the plasma brake for orbital and al-  
286 titude control.
- 287 7. The mission profile must provide the ability for daily updates of measurements  
288 of orbital parameters to assess the effect of the plasma brake, once activated.

289 In order to achieve the above scientific requirements, the spacecraft spin axis must  
290 be oriented with  $3^\circ$  accuracy, with spin of  $4.00 \pm 0.04$  rpm. Attitude determination sys-  
291 tem must supply the magnetic field vector with  $1^\circ$  accuracy and the satellite position  
292 must be known with 5-km accuracy. There must be a  $\sim 1$  kbs $^{-1}$  data downlink.

## 293 4 Description of the Mission

### 294 4.1 General Concept

295 FORESAIL-1 is a nanosatellite mission in LEO designed to answer the science ob-  
296 jectives outlined in Section 2. The payload consists of a PArticle Telescope (PATE) and  
297 a Plasma Brake (PB). PATE will measure energetic electrons in the energy range 80–  
298 800 keV as well as  $H^+$  ions (protons) and neutral atoms in the energy range 0.3–10 MeV.  
299 PB consists of a tether that will be used to lower the spacecraft altitude. The spacecraft  
300 is constrained under the CubeSat 3U standard to fit the two payloads.

### 301 4.2 Mission Timeline and Orbit

302 The manufacturing of the FORESAIL-1 payload and spacecraft started in 2018 and  
303 the manufacturing and integration will continue throughout 2019 until launch. The space-  
304 craft is scheduled to be launched in 2020 into a Sun-synchronous orbit at or lower than  
305 700 km altitude. Available launch opportunities are sought in 2019. Following the suc-  
306 cessful launch of the mission, the 1-month commissioning phase is scheduled to start im-  
307 mediately. After the commissioning phase, the mission's primary science phase is sched-  
308 uled for 4 months at the initial Sun-synchronous orbit. After the primary phase, the plasma  
309 brake will be applied to lower the spacecraft by more than 100 km such that 1) the plasma  
310 braking force is fully demonstrated and 2) from the lowered orbit the spacecraft will nat-  
311 urally de-orbit after 25 years. The lowering of the orbit will inject the spacecraft into  
312 a drifting polar orbit, allowing precipitation measurements in different MLTs. Follow-  
313 ing the successful orbital manoeuvring of the spacecraft, the science phase continues with  
314 detecting particle precipitation in the drifting orbit for at least 1–2 years. After this will  
315 be the ENA measurement phase. There is a possibility of an extended science phase that  
316 will be scheduled depending on the health of the spacecraft.

### 317 4.3 Spacecraft Conjunctions

318 FORESAIL-1 can be used in conjunction with various other spacecraft in order to  
319 determine the origin of precipitating particles observed. Spacecraft that can provide con-  
320 text to the FORESAIL-1 observations include the Solar and Heliospheric Observatory  
321 (SOHO), STEREO, the Advanced Composition Explorer (ACE), Wind, DSCOVR, the  
322 Geostationary Operational Environmental Satellites (GOES), and the Parker Solar Probe.  
323 These missions directly monitor solar wind conditions, coronal mass ejections and so-  
324 lar energetic particles.

325 The conditions encountered by FORESAIL-1 will depend strongly on the state of  
326 the other regions in the magnetosphere. Simultaneous observations from satellites such  
327 as Cluster, the Magnetospheric Multi-Scale mission, the Time History of Events and Macroscale



Interactions during Substorms spacecraft (THEMIS), and the Geomagnetic Tail Lab (GEO-TAIL) will be invaluable for understanding the global context in which the FORESAIL-1 measurements are made. In particular, when located in the relevant region, these spacecraft can monitor substorms occurring in the magnetotail and the associated fast earthward plasma flows which are also the sources of the energetic particle precipitations to be measured by FORESAIL-1, in addition to the RB source. They can also provide information about the wave activity in the magnetosphere, which will be key to interpreting the FORESAIL-1 observations. In the regions closer to Earth, data from the recently-launched Arase mission in the radiation belts will be of particular importance. Direct complementary observations to FORESAIL-1 will be provided by the POES satellites, which will provide precipitating particle data at similar energy ranges, however these data are often problematic and require corrections.

#### 4.4 Coordinated Ground-Based Observations

Coordinated observation will also be possible with ground-based instrumentation. Whenever suitable conjunctions with riometer chains such as the Finnish chain operated by Sodankylä Geophysical Observatory or the Canadian chain (NORSTAR) take place, it will be possible to compare energetic particle precipitation patterns to those inferred from cosmic noise absorption measurements. The special case of the Kilpisjärvi Atmospheric Imaging Receiver Array (KAIRA) is of particular interest, as this instrument, which may be used as a multibeam, multifrequency riometer could allow to finely study energetic precipitation along the FORESAIL-1 overpass. Indeed, KAIRA can provide 1 s time resolution observations of cosmic noise absorption with beams of  $10^{\circ}30^{\circ}$  width, depending on the considered frequency, which is accurate enough to study, e.g., individual patches of pulsating aurora with KAIRA (Grandin et al., 2017), suggesting that it may be possible to study mesoscale ( $<100$  km) structures in energetic precipitation using FORESAIL-1 and KAIRA conjunctions.

During conjunctions, it can also prove valuable to combine FORESAIL-1 data with observations of other space weather manifestations. For instance, auroral precipitation can be inferred from observations by all-sky imagers such as the MIRACLE network in Fennoscandia, and the geomagnetic context of FORESAIL-1 measurements can be given more accurately during conjunctions with ground-based magnetometer networks as well as by making use of polar cap convection maps derived from Super Dual Auroral Radar Network observations (SUPERDARN). Finally, FORESAIL-1 precipitation data could prove useful in studies focusing on electron density enhancements in the ionosphere using, e.g., ionosondes or incoherent scatter radars such as the European Incoherent Scatter radars (EISCAT). The latter ones also enable the study of effects such as ionospheric Joule heating or ion outflow, as they measure ion and electron temperatures and ion line-of-sight velocity, and when in a specific configuration they also allow estimating electric fields (Nygrén et al., 2011).

## 5 Payloads

The FORESAIL-1 mission will carry two science payloads, the Particle Telescope (PATE) and the Plasma Brake (PB).

### 5.1 Particle Telescope (PATE)

PATE measures energetic electrons,  $H^+$  ions (protons) and neutral H atoms. The targeted energy range of hydrogen measurement is 0.3–10 MeV, which covers the energies above the typical ring-current proton energies. This is to avoid the neutral hydrogen background originating from the interaction of the ring current with the geocorona. The primary energy range for electrons is 80–800 keV. In addition, the instrument is



**Figure 2.** a) Mechanical design concept of PATE. b) Schematic of the anti-coincidence (AC) and main detector (D) stack of each telescope.

376 sensitive to electrons at higher and lower energies in channels, where reliable particle iden-  
 377 tification cannot be performed, but especially the high-energy integral channel is still valu-  
 378 able since the contamination from heavier species comes only from relativistic protons,  
 379 which have low fluxes compared to relativistic electrons in typical conditions.

380 PATE consists of two particle telescopes with identical stacks of detectors (see Fig.  
 381 2). The longer Telescope 1 is directed along the long axis of the spacecraft, that is, per-  
 382 pendicular to the rotation axis, and thus, it scans the directions in a plane perpendic-  
 383 ular to the rotation axis of the spacecraft. The shorter Telescope 2 is directed along the  
 384 rotation axis, so it can maintain a stable orientation. When the rotation axis is pointed  
 385 towards the Sun, the telescope is able to measure the neutral hydrogen flux from the Sun.  
 386 Note that the instrument itself does not determine the hydrogen charge state but relies  
 387 on the geomagnetic field as a rigidity filter and on the measurement of angular distri-  
 388 bution to disentangle the neutral flux from the solar direction. The motivation for the  
 389 use of longer collimator in Telescope 1 is to improve the pitch angle angular resolution  
 390 to better than 10 degrees for the scanning telescope.

391 Both telescopes have a mechanical collimator defining the aperture, consisting of  
 392 an aluminum tube housing 18 (12) ( $500\ \mu\text{m}$  Al +  $500\ \mu\text{m}$  Ta) collimator rings in Tele-  
 393 scope 1 (Telescope 2), followed by a stack of silicon detectors, D1, D2, and D3, measur-  
 394 ing the energy losses of the particles in adjacent layers. The measured signals allow the  
 395 determination of particle energy and the identification of particle species (electron / H).  
 396 The thicknesses of the D detectors are  $20\ \mu\text{m}$ ,  $350\ \mu\text{m}$ , and  $350\ \mu\text{m}$ , respectively. The D1

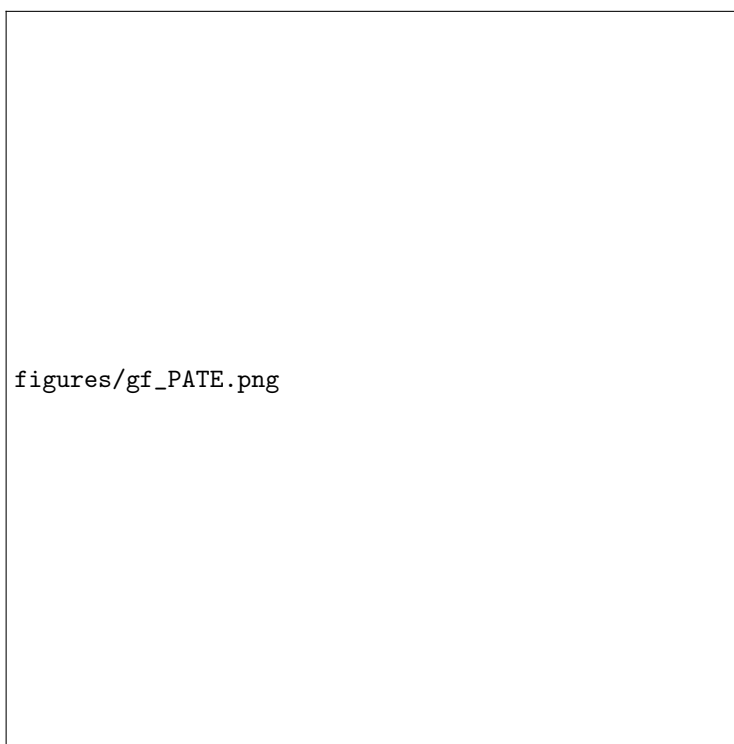
397 detectors have a bias voltage of 5 V, while the other ones are biased at 70 V. D1 and D2  
398 detectors are segmented so that the central elements have diameters of 5.2 mm while the  
399 total active-area diameters of both are 16.4 mm. D3 has a single active area of 16.4 mm  
400 diameter. Both detector stacks are covered at the bottom of the collimator by two thin  
401 (nominally 0.5  $\mu\text{m}$  each, 1 mm apart) Ni foils preventing low-energy ( $< 250$  keV) ions  
402 and soft ( $< 500$  eV) photons from entering the detector stack. Each aperture is limited  
403 from above by an annular anti-coincidence detector AC1 (300  $\mu\text{m}$  thickness) with a cir-  
404 cular hole of 14.0 mm diameter in the middle and an outer active-area diameter of  
405 33.8 mm. Another single-element circular anti-coincidence detector AC2 (350  $\mu\text{m}$  thick-  
406 ness, 33.8 mm active-area diameter) at the bottom of the stack flags particles penetrat-  
407 ing the whole D detector stack. Note that while the AC2 is operated in anti-coincidence  
408 with the D detectors for the nominal energy range of the instrument (particles stopping  
409 in the D stack), PATE still analyses the D detector pulse heights for particles trigger-  
410 ing AC2 but not AC1 to provide integral flux channels above the nominal energy ranges.  
411 The distances from upper surface to upper surface in the detector stack AC1–D1–D2–  
412 D3–AC2 are 2.5 mm, 2.0 mm, 2.5 mm, and 2.5 mm. The lower Ni foil is 3.2 mm above  
413 the upper surface of AC1.

14 Signal processing is based on continuous sampling and digitization of the analog  
15 signals and on digital filtering and pulse height analysis. The signal processing board con-  
416 tains 16 Analog to Digital Converters (ADCs) and a *Microsemi ProAsic3L* Field Pro-  
417 grammable Gate Array (FPGA) handling the signal processing for both telescopes. The  
18 signal sampling rate in each ADC channel is 10 MHz, and the digital streams are pro-  
419 cessed by the FPGA, which is running at 40 MHz. The logic analyzes the incoming dig-  
420 ital data streams, detects pulses, and identifies the particle for each valid coincidence event,  
21 counting and rejecting from further analysis any events not matching validity criteria.  
422 Valid particle events are then counted in separate counters based on their detection time,  
23 species and measured energy, forming the bulk of the science data of the instrument. The  
424 electron (hydrogen) spectrum consists of seven (ten) energy channels, log-spaced in mea-  
425 sured energy.

### 426 5.1.1 Instrument Performance

27 The performance of the PATE electron measurement has been simulated using GEANT4  
428 (Agostinelli et al., 2003). The simulation is performed for an isotropic electron distri-  
429 bution launched from a (15-cm radius) spherical surface surrounding PATE. The simu-  
430 lated pulse heights of all D-detector signals are analysed to separate electrons and hy-  
431 drogen (ions/ENAs) and to measure particle energies, as in the FPGA-based on-board  
432 analysis. Particles producing a hit (i.e., an energy deposit  $> 50$  keV) only in D2 are iden-  
433 tified as electrons and particles producing a hit only in D1 are identified as hydrogen.  
34 Hit levels can be set freely in the logic, but values lower than 50 keV in D1a, D1c, AC1  
435 and AC2 (see Fig. 2) should not be used as the simulated RMS noise levels in those pads  
436 are around 9–14 keV. While electrons are able to produce hits in D1 as well, there is only  
37 a minor level of electron contamination in the nominal hydrogen energy channels, which  
438 require the energy deposit in D1 to exceed 110 keV. If more than two adjacent D lay-  
439 ers detect a pulse, the Delta E – E method is used for clean species separation. The ge-  
440 ometric factor of the seven electron channels as a function of electron energy for (the shorter)  
441 Telescope 2 is shown in Fig. 3 (Oleynik & Vainio, 2019). The high-energy tails of the  
442 response functions are due to the inevitable scattering of electrons off the detectors and  
443 other structures inside the telescope, which prevents the full energy of the electron to  
444 be absorbed in active detector elements. The internal energy resolution of the instru-  
445 ment is much better for hydrogen than for electrons and the response functions are close  
446 to boxcar type within the nominal energy range.

447 The geometric factors of both telescopes are mostly determined by the uppermost  
448 collimator ring (with an inner diameter of 21.5 mm) and the hole of the AC1 detector



**Figure 3.** The geometric factor of PATE as a function of energy for the seven electron energy channels within the nominal energy range (80–800 keV), simulated using GEANT4 (Agostinelli et al., 2003). The range above 800 keV is additionally covered by a penetrating particle flux channel.

Data source	Data rate [bit/s]	Data amount per day [kiB]
Rotating telescope	993	10500
Solar pointing telescope	60	633
Housekeeping	13	135
Total	1066	11268

**Table 1.** Summary of the PATE telemetry budget during science operations.

(diameter of 14 mm), which are at a vertical separation of 12.0 cm (7.0 cm) in Telescope 1 (Telescope 2). The nominal value of the geometric factor is 0.037 cm<sup>2</sup> sr (0.11 cm<sup>2</sup> sr) for Telescope 1 (Telescope 2), but especially electrons have somewhat lower values (see Fig. 3) due to scattering off the Ni foils and the D1 detector, causing trajectories to miss the D2 detector. The nominal angular widths of the acceptance cones (half width at half the on-axis value) for the two telescopes are 4.6° and 7.9°, respectively. The instrument can also be operated in a mode where only the central segments of the D1 and D2 are included in the coincidence logic while the rim areas are logically included in the anti-coincidence. This allows to decrease the geometric factors of the telescopes by a factor of about 7. Any detector element can also be switched off from the logic entirely.

### 5.1.2 Mass, power and telemetry budgets

The mass of PATE is 1.2 kg, consisting of the instrument box and mechanical support structures (435 g), detector and pre-amplified board housings (290 g), the collimators (180 g), cables (115 g), and the rest (180 g), including the back-end electronics stack. The power budget for PATE is 2.5 W, half of which is consumed by the FPGA, with an additional margin of 20%, mainly required to account for the final FPGA power consumption.

The telemetry budget of PATE is summarized in Table 1. Spectral resolution of the data products for both electrons and hydrogen is on average  $\Delta E/E \approx 40\%$  within the nominal energy ranges, which means that the spectral counter data consists of eight (ten) differential and one (two) integral channels for electrons (hydrogen). The rotation period, nominally 15 s, of the satellite equals the time resolution of the instrument. This main time frame is further broken in 36 angular sectors for the rotating telescope to provide the angular distribution measurement. Both telescopes deliver, in addition to the spectral counter data, also pulse height data samples, which allow an accurate in-flight calibration and health monitoring of the detection system.

## 5.2 Plasma Brake (PB)

### 5.2.1 Operating Principle

The Plasma Brake instrument is designed to measure the Coulomb drag, i.e., the braking force caused by the ionospheric plasma ram flow to an electrically charged tether (Figure 4). When interacting with the surrounding plasma, the negative tether gathers positive ions, which tends to neutralize the tether. To maintain the charge, the tether is connected by a high-voltage power system to a conducting surface (deployable booms for FORESAIL-1) that acts as an electron sink to close the current system through the plasma (Janhunen, 2011). The braking force can be measured in two modes. One is to monitor the system spin rate while charging the tether synchronously with the tether rotation (PB Measurement). The other is to maintain a constant charging and monitor the spacecraft velocity and orbital elements (PB Mode).

figures/PB\_operating.png

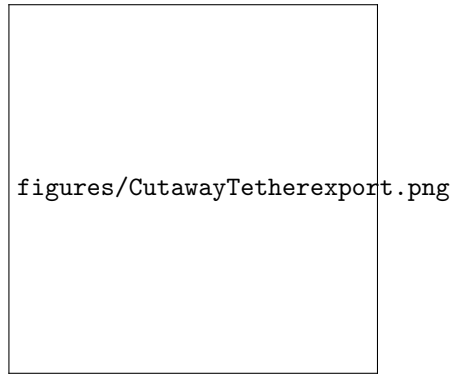
**Figure 4.** Operating principle of the PB: an electric current system and a net thrust exerted on the negative tether by the plasma ram flow.

We employ  $-1$  kV voltage, which is the maximum without risking electron field emission from micrometeoroid-struck parts of the tether wires. At this voltage, the expected nominal plasma brake thrust per tether length is  $58$  nN/m when the tether is perpendicular to the ram flow. This value is obtained by using Equation 1 of Janhunen (2014) and assuming plasma density of  $3 \cdot 10^{10} \text{ m}^{-3}$ , mean ion mass of 10 proton masses, ram flow speed of  $7.8$  km/s and tether's effective electric radius of  $1$  mm. The tether's collected ion current is small, nominally  $30 \mu\text{A}$  for a  $300$  m long tether.

### 5.2.2 Design

The tether is deployed from a chamber (blue) by the centrifugal force affecting the tether tip mass (gray button inside the red collar) (Figure 5). The mechanical interface through the satellite side panel is provided by an anti-static collar (red) to avoid triple junctions of plasma, insulator, and high voltage tether. The tip mass is locked during the launch by two launch locks located on opposite sides of the tether chamber opening. The tether reel (dark gray) and the stepper motor that turns it (orange) are nested





**Figure 5.** PB payload structure.

inside the tether chamber. The tether high-voltage contact is through the conducting reel and the slider (copper brown) attached straight to the high-voltage converter (orange) board behind the tether chamber.

The PB tether is made of thin metallic wires that are periodically bonded to each to withstand the micrometeoroid and space debris collisions (e.g., Seppänen et al., 2011). The single wire thickness is a few tens of micrometers to minimise the ion current to the tether and thus the mass of the power system and size of the electron gathering surface.

The nominal FORESAIL-1 tether length is 300 meters requiring  $\sim 20$  Nms of the total angular momentum to deploy it. The initial angular speed of  $180^\circ/\text{s}$  is sufficient to provide a centrifugal force of 0.4 cN which would safely pull out a tip mass weighing 2.5 g without breaking the tether. The angular momentum is provided by several consecutive satellite spin-up and reel-out maneuvers to compensate for the decreasing spin rate associated with the increasing moment of inertia. Since magnetorquers are used for attitude control, the angular acceleration is low to avoid the tether winding around the satellite. It also indicates that a considerable amount of time is needed to provide the angular momentum. However, after measuring the Coulomb drag force with a few tens of meters of the tether, the force can be used to spin up the satellite by modulating the tether voltage in synchronization with the rotation (charging downstream to increase the spin rate). After deploying the tether and taking PB measurements, the PB mode will start by continuously charging the tether which in turn will lower the orbit and start the satellite drift in MLT. When reaching  $\sim 600$ -km altitude and gaining a sufficient drift in MLT, the satellite will be prepared for PATE observations – the tether will be discarded to allow the spin axis being pointed towards the Sun. It can be done by attempting to reel in the tether which might break because it would be partially broken by micrometeoroids. A broken tether would deorbit in a few months thanks to its large area/mass ratio. If the tether does not break, the attitude control system and/or PB itself will have to provide an angular momentum to compensate for an increasing spin rate.

### 5.2.3 Mass, Power and Telemetry Budgets

The mass of the PB is 0.6 kg including a margin of 25%. The structure (frame, tether chamber, and motor mounting shaft) contributes 67% to the mass budget. The size of the system is  $67 \times 84 \times 96$  mm. For the two measurement modes of the PB, the power budget for PB is 600 mW. For the tether reeling, the motor and the controller consumes 7 W continuously. In case the reel-out power cannot be provided continuously by the spacecraft, the operation can be done in stages. The telemetry budget of the PB is summarised in Table 2. To reduce the overall telemetry budget, the long-duration routine PB mode

Mode	Data rate [bit/s]	Data amount per day [kiB]
Reel-Out/In	128	1350
PB Mode	19	200
PB Measurement	256	2700
Standby	64	675

**Table 2.** Summary of the PB telemetry budget for the operation modes.

Subsystem	Planned mass (g)	Mass with contingency (g)	Fraction (%)
OBC	80	96	2
EPS	926	1093	23
Magnetorquers	240	288	6
UHF Transceiver	80	96	2
Antennas	50	60	1
Structure	1033	1240	26
PATE	1000	1200	26
PB	600	660	14
Total	4009	4733	100

**Table 3.** FORESAIL-1 mass budget.

uses a lower telemetry rate. Frequent housekeeping data are not required because in the PB mode, altitude change over weeks to months will show the success of experiment.

## 6 Spacecraft platform

The platform has been designed to accommodate the payloads and to provide data, power and mechanical interfaces. The overall mission tree for the FORESAIL-1 is presented in Figure 6. Since one of the key technological drivers for the mission is reliability, the avionics stack has several designs to this end. The avionics stack is enclosed in a vault providing substantially better shielding than what is typically seen on CubeSats (around 4 mm equivalent aluminium, instead of the more typical 2 mm), thus enhancing system tolerance to total dose. Single-event effects will be mitigated using dual cold redundancy, hardware overcurrent protection, minimization of the software footprint, and systematic data integrity checks. Finally, the FORESAIL-1 components will be submitted to radiation test campaigns; the radiation response data will be made available publicly in order to benefit from the broader field of (small) satellite technology and help other designers addressing this issue.

The avionics stack consists of the Electrical Power System (EPS) for power collection, storage and distribution, Communication System for telemetry, On-Board Computer (OBC) for telemetry handling and mission and payload management and data storage, and Attitude Determination and Control System (ADCS) for maintaining the attitude modes during different operation phases. The mechanical structure satisfies dimensional limitations of the CubeSat standard and ensures modular configuration of the spacecraft's subsystems. The configuration of the platform is shown in Figure 7. The total mass budget is shown in Table 3.



**Figure 6.** FORESAIL-1 mission product tree.



**Figure 7.** Spacecraft structure including the subsystems without shielding.

Parameter	X,Y axes	Z axis
Number of wire turns	184	952
Nominal current, mA	26.2	19.3
Dipole moment, Am <sup>2</sup>	0.2	0.2
Power, mW	51.7	56.5

**Table 4.** Quantitative parameters for the air core magnetorquers.

## 6.1 Platform subsystems

Each payload has different requirements for attitude. PATE needs to be oriented towards the Sun, while the detector with longer collimator scans the environment in the directions perpendicular to the Earth–Sun vector. The PB needs the spinning control for deploying and maintaining the tension of the tether. While the PB does not need specific pointing direction, its tether deployment introduces a mass distribution change that will require a proper control of the spacecraft angular momentum.

The attitude determination and control system (ADCS) is divided into Attitude Determination System (ADS) and Attitude Control System (ACS). The ADS is equipped with gyroscope (L3GD20H), magnetometer (LIS3MDL), and in-house built sun sensors (Noman et al., 2017). The angular velocity of the spacecraft is a crucial parameter for the payload attitude modes, thus gyroscope is necessary. Sun sensors are required for a precise sun pointing control, and magnetometers are needed for properly using with attitude control. The outputs will provide a full attitude information through using an unscented Kalman filter algorithm on all sensors.

The ACS changes the orientation of the spacecraft by using magnetorquers. There is a closed loop feedback, which ensures the maintenance of desired attitude by repeating the torque command until the desired orientation is achieved. The mission requirements for PATE and PB impose specific constraints on the attitude control. The spacecraft uses the following attitude control modes:

- Detumble mode: To stabilize the spacecraft after deployment
- Spin control mode: In order to deploy the PB tether, this mode spins up the spacecraft preferably with the spin axis being aligned with the earth pole. After deorbiting, spin down might be required during tether reel in if the tether does not break.
- Sun pointing mode: This mode continuously points Telescope 2 (at  $-Y$  axis) towards the sun while spinning in order for Telescope 1 (at  $Z$ -axis) to scan the sky.

Magnetorquers are sufficient to provide all necessary control modes for the manipulation of the attitude and angular velocity. They are designed in-house, in form of air coils, so that they can be integrated to the solar panels (for the  $X$  and  $Y$  axes), and on a small factor magnetorquer driver board for the  $Z$  axis. All axes have two magnetorquers connected in parallel. The designed magnetorquers are driven through a custom-built coil driver to optimize either consumed power or the generated magnetic moment. The quantitative parameters of air core magnetorquers are given in Table 4.

The EPS of the satellite consist of solar panels, power conditioning and power distribution (Ali et al., 2014), (Mughal et al., 2014). The solar panels are mounted on every long side of the satellite. The solar panels at other sides than where PATE is located consist of 7 solar cells connected in series, whereas the panel at the PATE side consists

596 of 6 cells in order to provide space for the telescope. The power conditioning consists of  
597 switching buck converters to convert the incoming solar power to the battery voltage level.  
598 A perturb and observe based algorithm extracts the maximum power from the solar cells.  
599 Each subsystem houses a linear regulator to effectively convert the voltage down to a sta-  
600 ble 3.3 V with a low ripple factor. The theoretical efficiency of the EPS is above 85 %.  
601 All losses have been accounted for in the calculation of power budget. The maximum  
602 input power is 7 W, whereas the average power consumption in the nominal mode 3.7  
603 W.

604 The on-board computer (OBC) is the satellite's main computer responsible for compu-  
605 tational and data storage needs, running the ADCS algorithms, and storing the teleme-  
606 try and housekeeping data logs. A safety-critical ARM Cortex-R4 based processor is se-  
607 lected as the central OBC of the spacecraft. The processor features 256 k data rapid-  
608 access memory and 3 MB of program flash. There are external flash memory devices also  
609 interfaced with the processor. To facilitate possible faults due to radiation, the OBC houses  
610 two cold redundant symmetric processors. Only one of the processors is active and pow-  
611 ered. The arbiter switches the control to the redundant processor in case of failure. The  
612 OBC is responsible for operational work during the PB and PATE operations and col-  
613 lects all relevant telemetry data for downlink.

614 The UHF transceiver onboard the FORESAIL-1 consist of CC1125 transceiver with  
615 maximum output power of 15 dBm (30 mW); an external power amplifier RF5110G (gain:  
616 31.5 dB, maximum output power 34.5 dBm) to amplify the power to desired 1.5-watt  
617 power in the transmit chain. In the receive chain, it consist of a Low Noise Amplifier and  
618 a bandpass filter.

## 619 7 Ground Segment and Operations

620 The primary ground station used for satellite operation is located in Aalto Univer-  
621 sity Campus in Espoo, Finland. The ground station operates mainly as a radio amateur  
622 satellite station and has capabilities for operation on radio amateur UHF, VHF and S-  
623 bands. Due to its northern location the ground station has an average link time to po-  
624 lar orbit up to 90 minutes per day for all passes above the horizon. Ground station ra-  
625 dio systems are built based on the Software Defined Radio (SDR) architecture which fa-  
626 cilitates easy satellite-specific customization. The ground station infrastructure and equip-  
627 ment are designed and operated by students and serve also for educational purposes for  
628 the Aalto University. The Aalto University ground station also operates as the mission  
control centre.

630 Typically, each satellite pass provides 10-20 min of link time between satellite and  
631 ground station. The data rate requirements for FORESAIL-1 in science mode require  
632 the downlink data rate to be approximately 8 kbps. Since the data rate requirements  
633 are not stringent in order to accomplish both the missions, the ground station at Aalto  
634 University easily handles the data rate requirements.

## 635 8 Data Products

636 FORESAIL-1 data products are outlined in Table 5. The Aalto University ground  
637 station is responsible for downlinking the L0/L1 data containing general spacecraft house-  
638 keeping data, PATE raw data, PB housekeeping and mission log information. These low  
639 level data products are shared forward using file-based web interface. For PATE data  
640 processing, Level 1 raw data files combined to ADCS metadata and orbital information  
641 are used to produce calibrated and geolocated measurements of the particle fluxes.

642 These data products will be used to demonstrate that the PB de-orbits the satel-  
643 lite at a measurable amount, as a function of time. The PATE data will be used to in-



Provider	Data product	Details
Mission	List of data availability	
Spacecraft	Orbital data, Attitude	Includes position and altitude required to estimate the PB performance
PB	Electric current in the tether L0 telemetry stream L1 raw data L2 calibrated data L3 derived products	Plasma density L2 and plasma density including position
PATE	Flux L0 telemetry stream L1 raw data L2 calibrated data (fluxes) L3 derived products	L2 including location Precipitation maps (varying time resolutions) Angular distribution data (spin resolution) Event catalogues (mission duration)

**Table 5.** FORESAIL-1 data products. PATE fluxes are given as function of energy, time and pointing azimuth.

for the precipitation energy spectrum in time and place, addressing the science objectives. The science data are open for everyone, with the nominal rules-of-the-road typical in the field of space physics.

## 9 Summary and Discussion

The increasing number of small satellites launched into Earth's orbit raises timely concerns about the sustainable use of space. Small, rapidly built and launched satellites have a large future potential for scientific and commercial use. However, the satellites will become debris sooner rather than later, if they have poor radiation tolerance and if they are not de-orbited at the end of mission. The Finnish centre Of excellence in Research of Sustainable space (FORESAIL) funded by the Academy of Finland tackles this issue by focusing both on science of the near-Earth radiation environment and on novel technological solutions related to building more resilient instruments and debris removal.

The first nanosatellite designed and built by the Centre of Excellence, FORESAIL-1, is a 3U cubesat operating at polar LEO orbit at and below 700 km which will produce energy-dependent pitch angle spectra of electrons and protons that precipitate from the RBs into the ionosphere. Further, it will measure energetic neutral atoms (ENAs) from the Sun and test the PB technology to lower the spacecraft altitude and manage its orbit in space.

Today, nanosatellites can address significant scientific questions. This requires focused measurements and innovative technological approaches, as well as coordination with the other spacecraft and facilities operating simultaneously. The FORESAIL-1 PATE instrument will make unprecedented and precise measurements of precipitating electrons and protons with high temporal resolution. It will be able to discriminate between electrons and protons, over a wide energy range (80 – 800 keV for electrons and 0.3 – 10 MeV for protons and neutral atoms). The large coverage of the polar cap at different orbital planes is achieved by operating the PB at the beginning of the mission, which sets the

671 spacecraft orbital plane to drift in magnetic local time. This is the first time such a ma-  
 672 noeuvre is attempted with a nanosatellite, and if successful, it will open new avenues for  
 673 controlling the orbits of propellantless spacecraft, expanding their operational and ob-  
 674 servational ranges.

675 FORESAIL-1 is targeted to make important advances in radiation belt physics. With  
 676 simultaneous observations in the solar wind, magnetosphere and from the ground, FORESAIL-  
 677 1 will allow quantifying the role of solar wind and outer magnetospheric driving, and the  
 678 role of different plasma waves in the inner magnetosphere in the precipitation process.  
 679 The primary science phase allows connecting precipitation signatures and mechanisms  
 680 to geomagnetic activity levels and solar wind conditions. Together with the novel Vlasi-  
 681 ator model, a global hybrid-Vlasov simulation (Palmroth et al., 2018), it will be possi-  
 682 ble to tie precipitation measurements to global processes in the outer magnetosphere for  
 683 the first time, as Vlasiator begins its full six-dimensional operation before the launch of  
 684 FORESAIL-1. The six dimensions refer to three dimensions in the ordinary space and  
 685 three in the velocity space to describe the full particle distribution function, which is used  
 686 to infer the energy spectrum and pitch angle in time. At the time of writing, Vlasiator  
 687 allows already 2D electron precipitation calculations (see Fig. 1), which are in reason-  
 688 able agreement with previous estimates (Hardy et al., 1985). Once the model is fully 6D,  
 689 *in situ* observations of electron precipitation can be directly compared to kinetic pro-  
 690 cesses anywhere in the magnetosphere, without limitations as to whether a spacecraft  
 691 traverses particular regions. This unprecedented plan will likely open new avenues in space  
 692 physics.

693 Successful observations of ENAs from the Sun will allow estimating for the first time  
 694 comprehensively estimating the suprathermal coronal ion population indirectly. This is  
 695 the key knowledge for improved understanding of the charged particle acceleration pro-  
 696 cesses at the CME-driven shock waves close to the Sun and of the CME energy budget.  
 697 The solar ENA flux has been measured only once in a very fortuitous event (Mewaldt  
 698 et al., 2009). Observing the ENA flux from the Sun on a regular basis as a function of  
 699 time and solar activity is unprecedented.

700 The demonstration of the altitude manoeuvre of FORESAIL-1 will bring poten-  
 701 tial for future applications for the PB as a standard tool for removing satellites from their  
 702 orbits. This is in particular a compelling solution considering the possibly upcoming reg-  
 703 ulations for including debris mitigation techniques in newly launched spacecraft.

704 FORESAIL-1 is at the forefront of scientific nanosatellites. The advances we have  
 705 made will be particularly important in demonstrating the usefulness of nanosatellites in  
 706 making relevant physics and discovery measurements (ENA), whose spatio-temporal res-  
 707 olution could be brought to a new level using fleets of nanosatellites. Technological so-  
 708 lutions of FORESAIL-1 have particularly far reaching impact for future debris removal  
 709 solutions and orbit control. All these aspects are expected to pave the way for the sus-  
 710 tainable use of space.

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(Palmroth, 2019)) is distributed under the GPL-2 open source license at <https://github.com/fmihpc/vlasiator/> (Palmroth & the Vlasiator team, 2019). Vlasiator uses a data structure developed in-house (Sandroos, 2019), which is compatible with the VisIt visualization software (Childs et al., 2012) using a plugin available at the VLSV repository. The Analysator software (<https://github.com/fmihpc/analysator/>, (Hannuksela & the Vlasiator team, 2019) was used to produce the presented figures. The run described here takes several terabytes of disk space and is kept in storage maintained within the CSC IT Center for Science. Data presented in this paper can be accessed by following the data policy on the Vlasiator web site.

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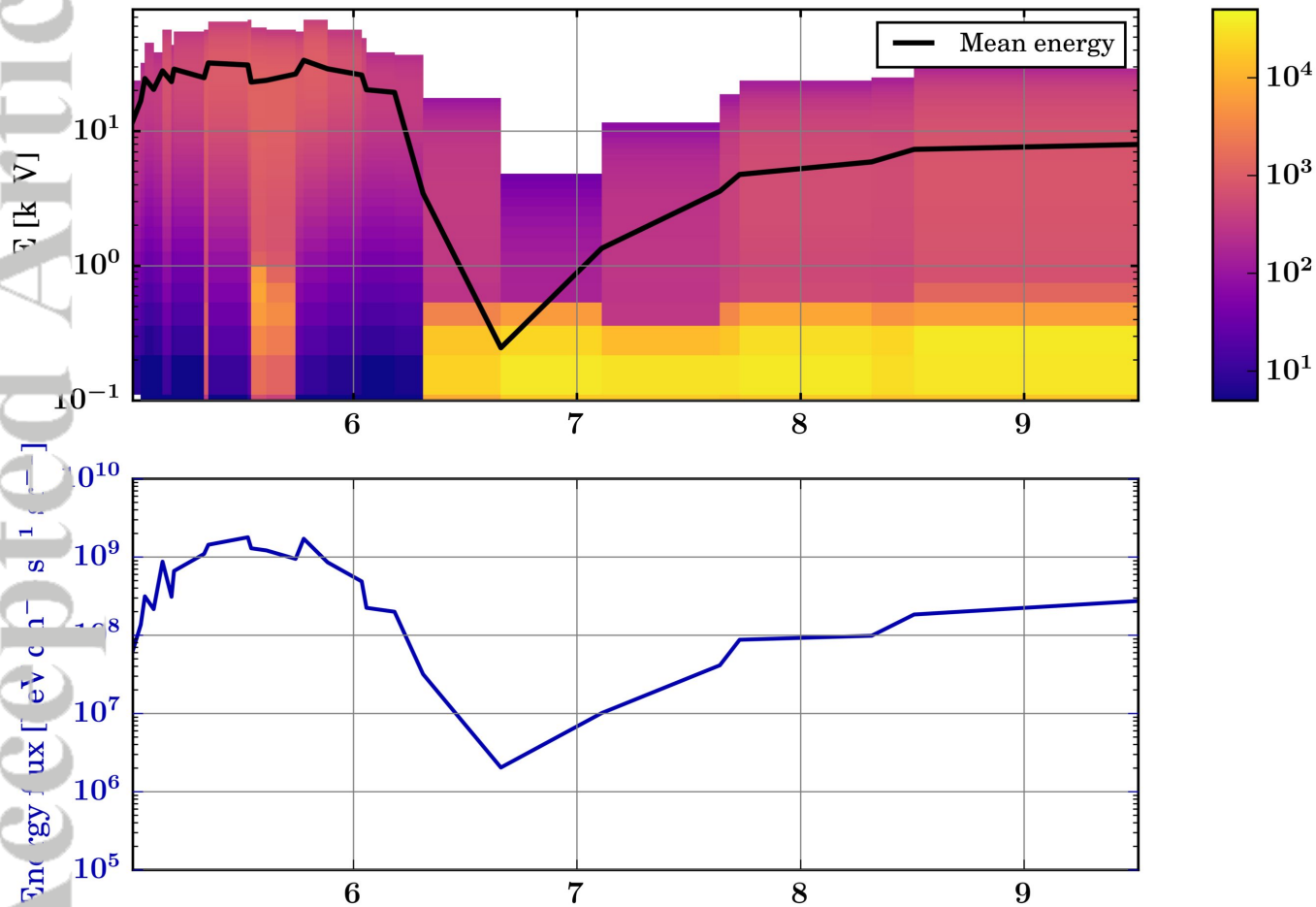
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Figure 1.

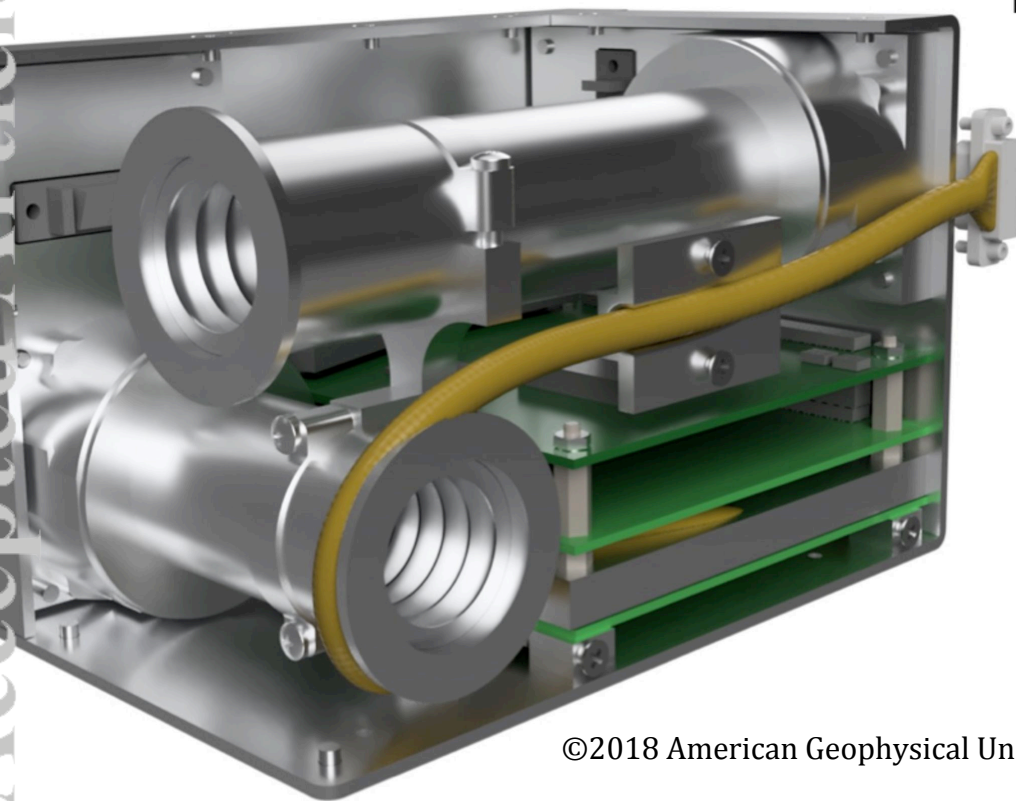
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# Electron precipitation – 00 MLT

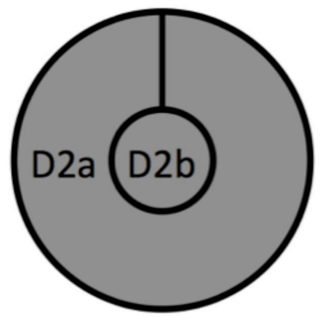
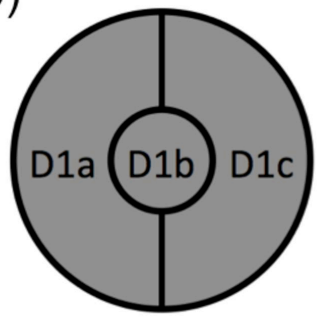
Differential  
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Accepted Article



b)



AC1

AC1



D1a D1b D1c

D2a D2b D2a

D3



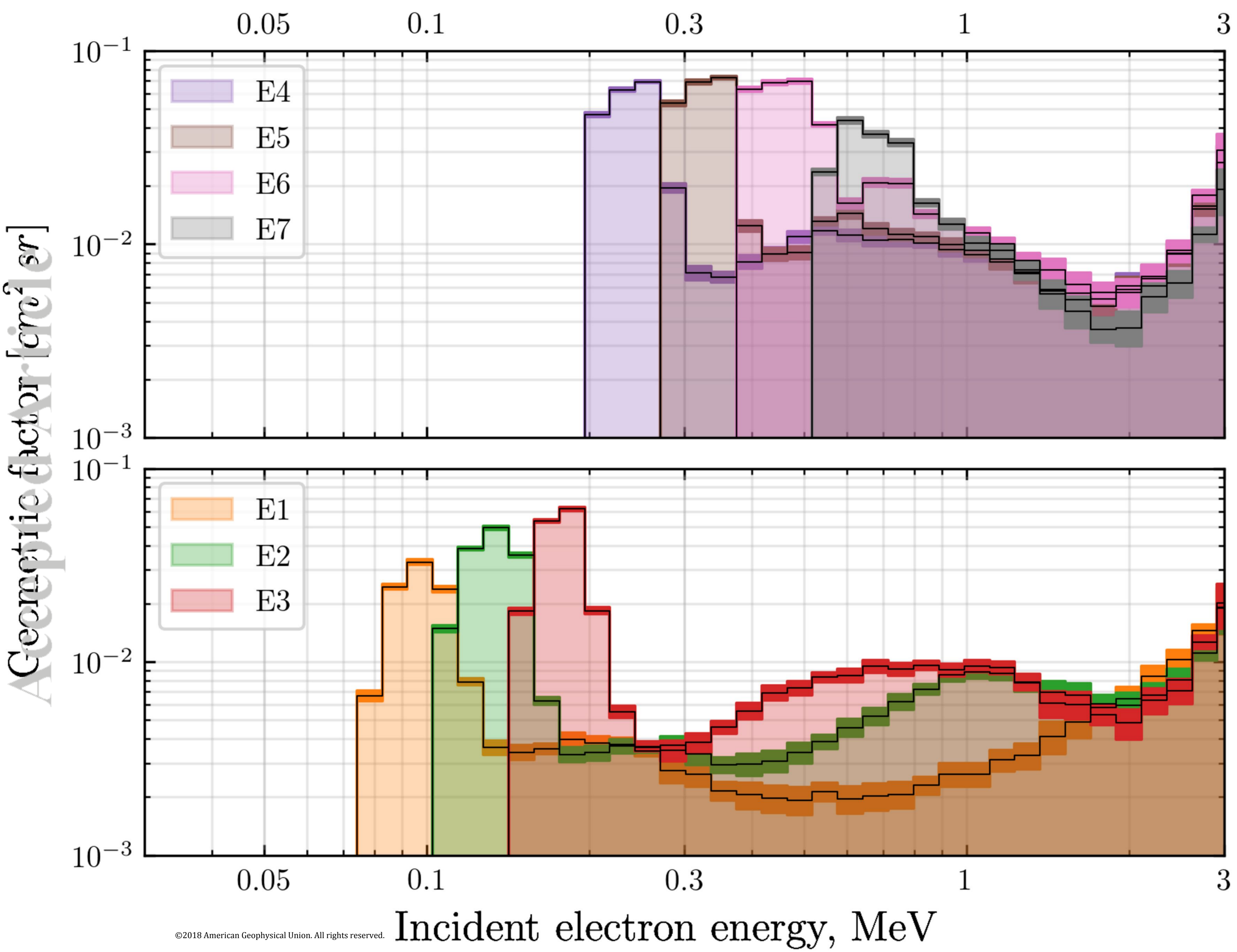
AC2



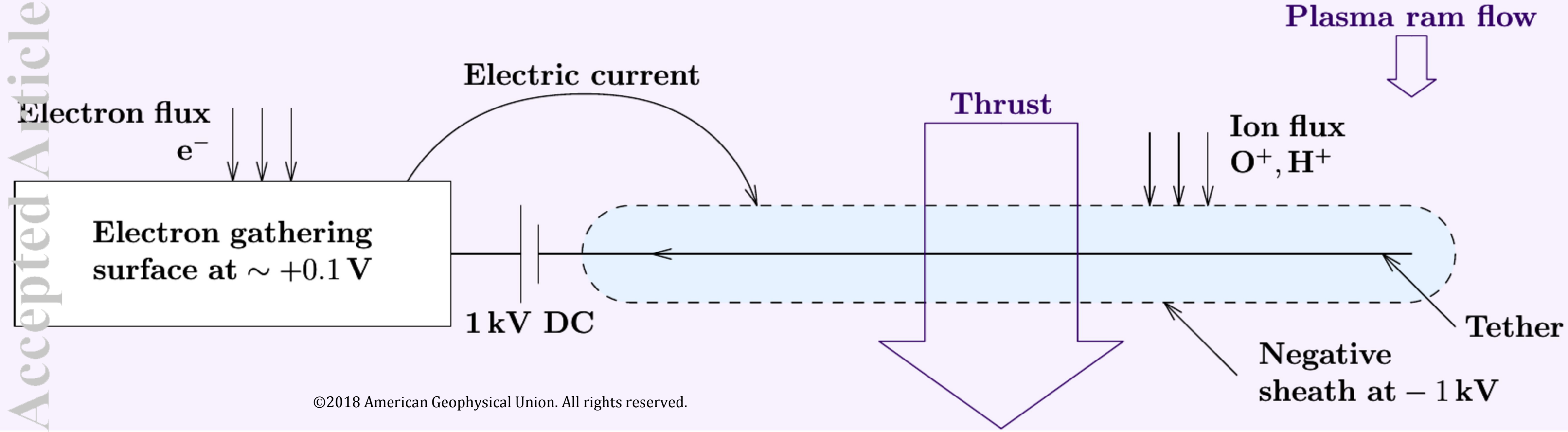
Figure 3.

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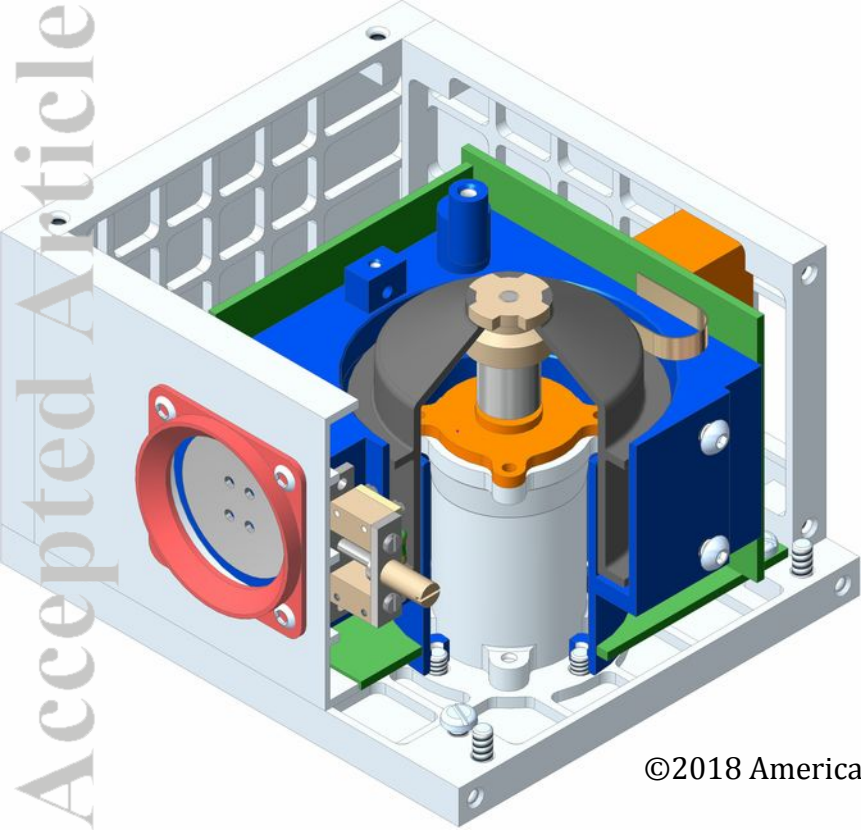




Accepted Article

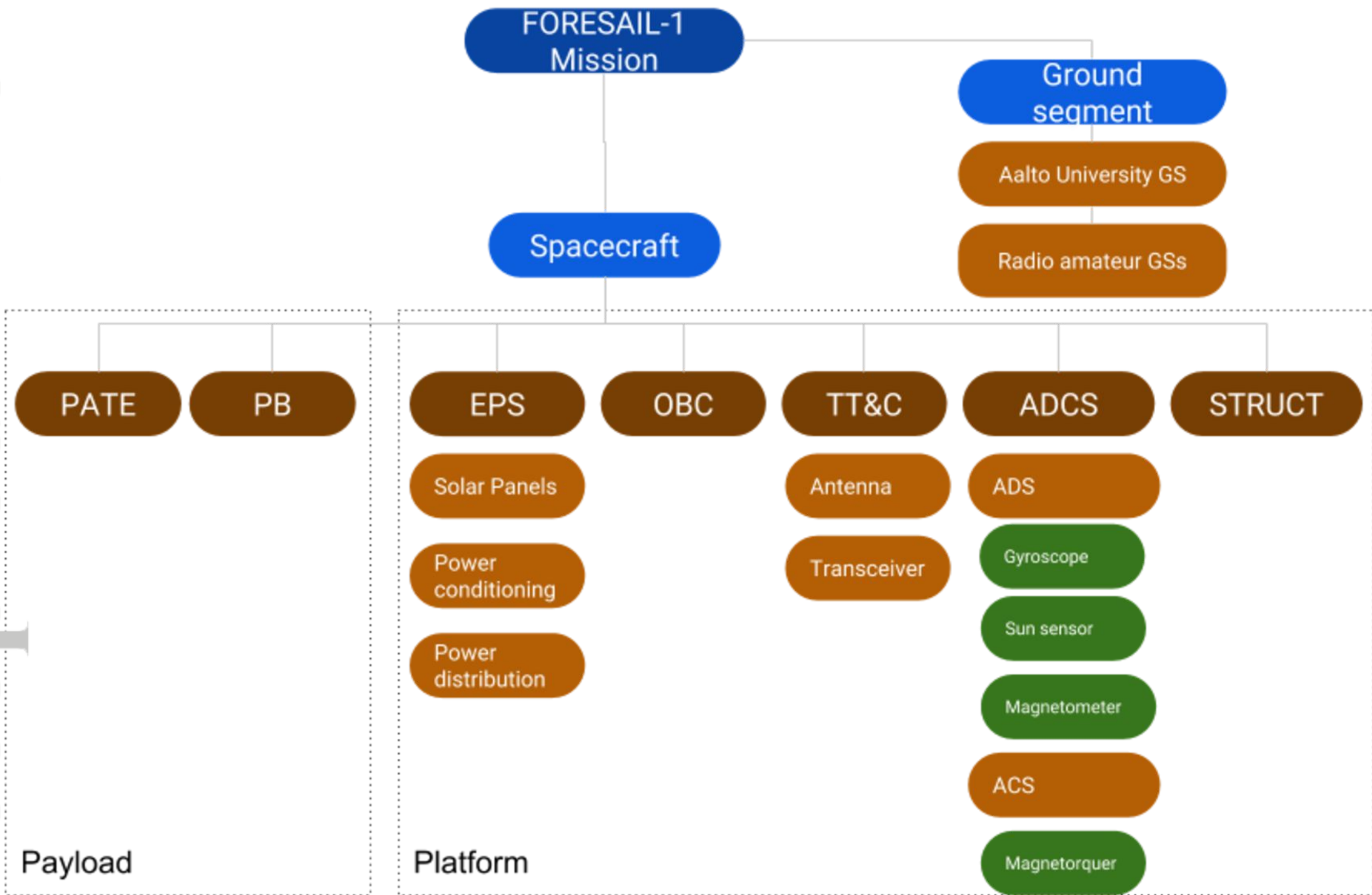


Accepted Article

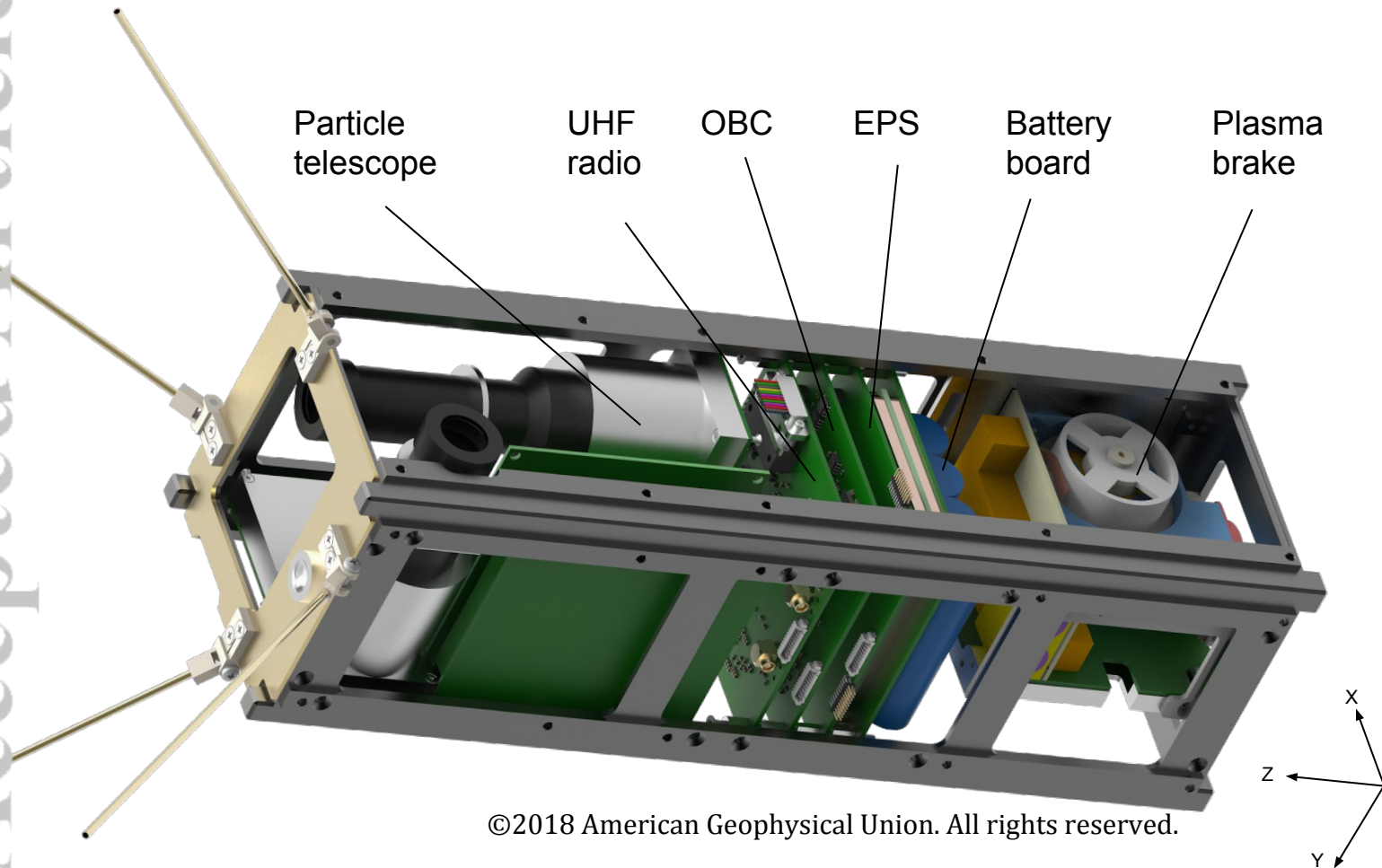


Accepted Article



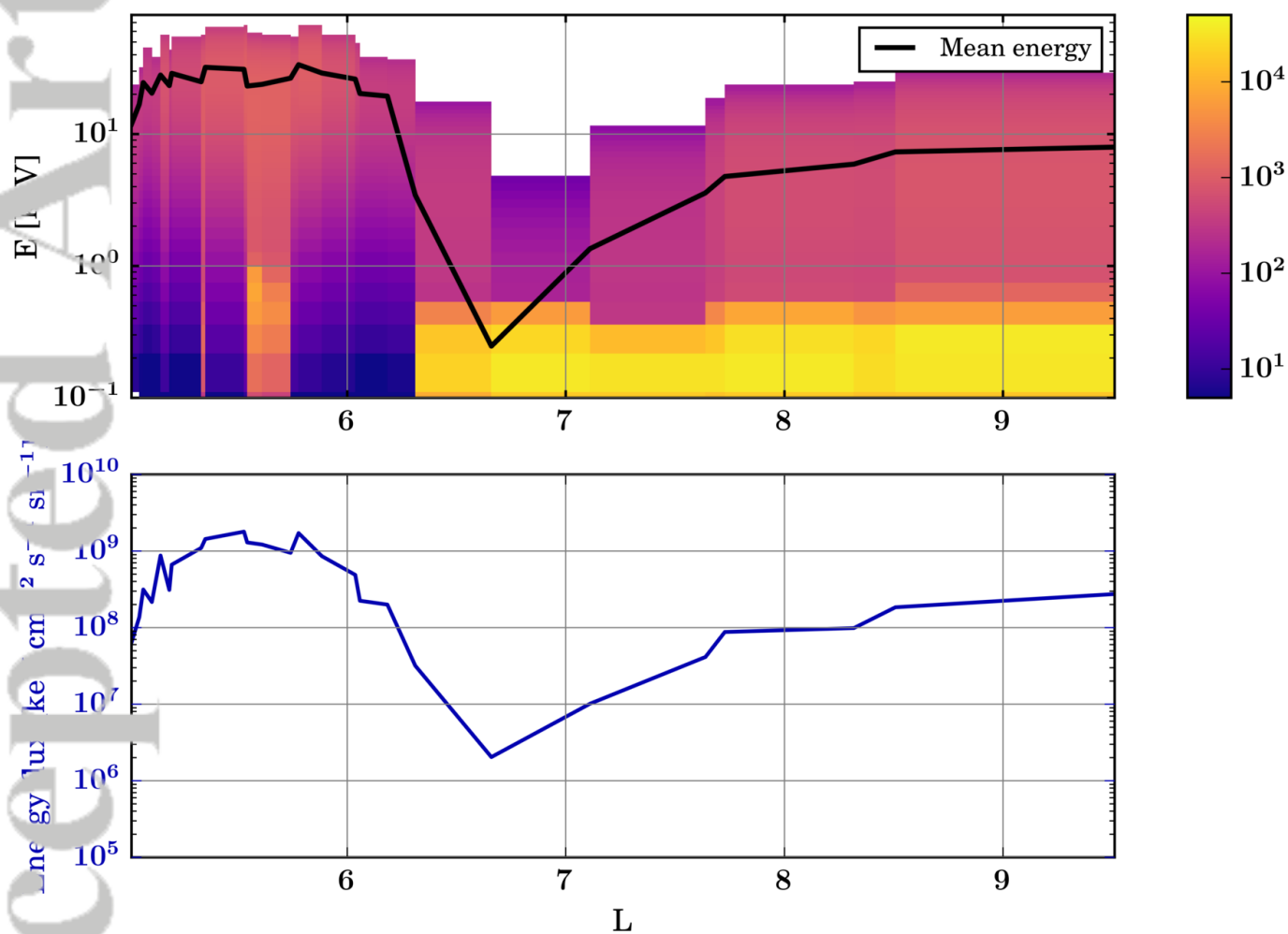


Accepted Article



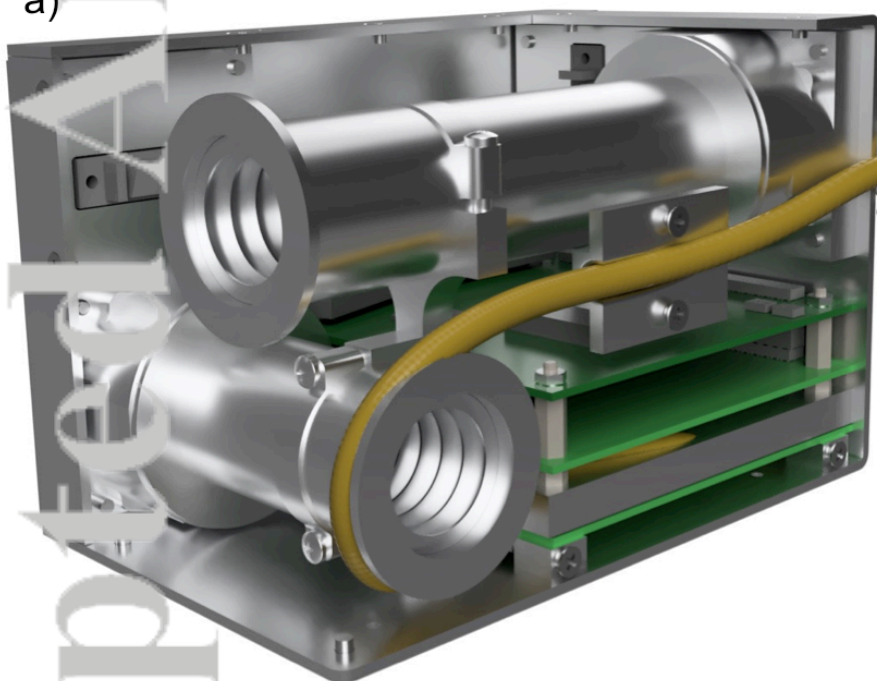
### Electron precipitation – 00 MLT

Differential number flux  
[ $\text{el cm}^{-2} \text{s}^{-1} \text{sr}^{-1} \text{eV}^{-1}$ ]

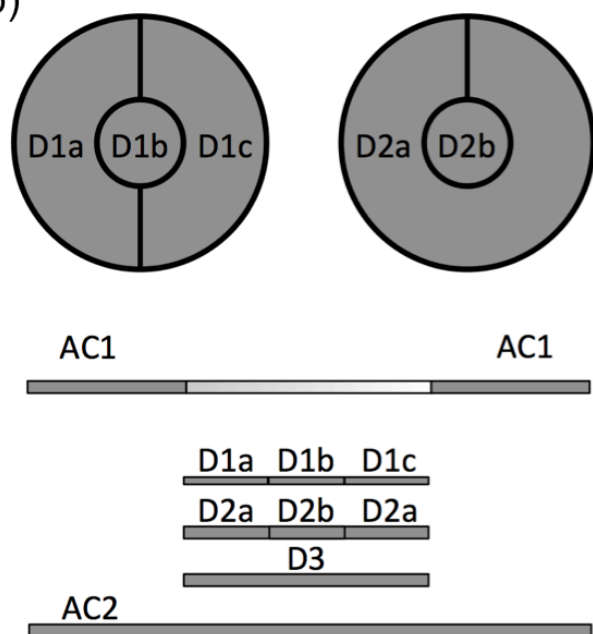


2018JA026354-f01-z-.png

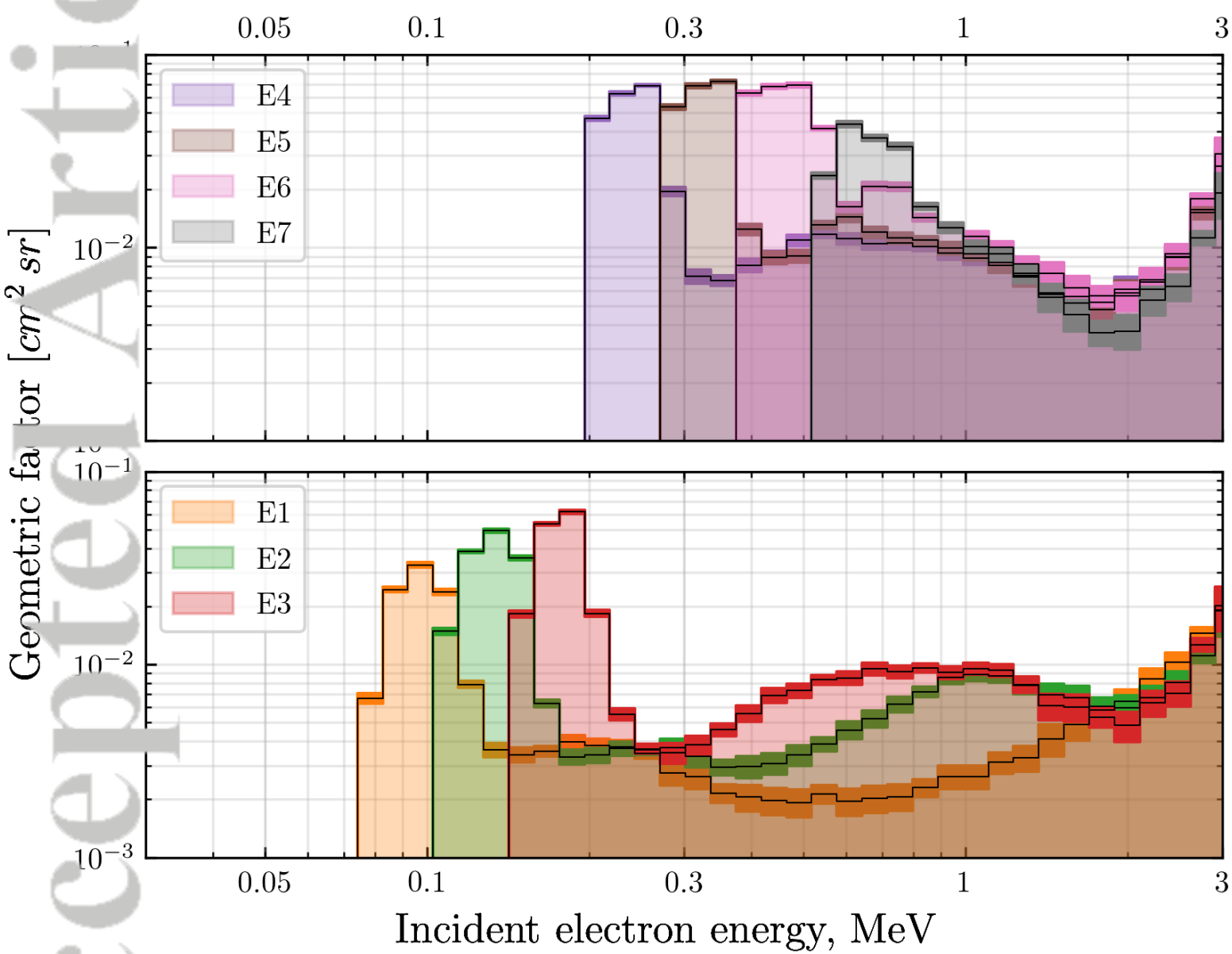
a)



b)

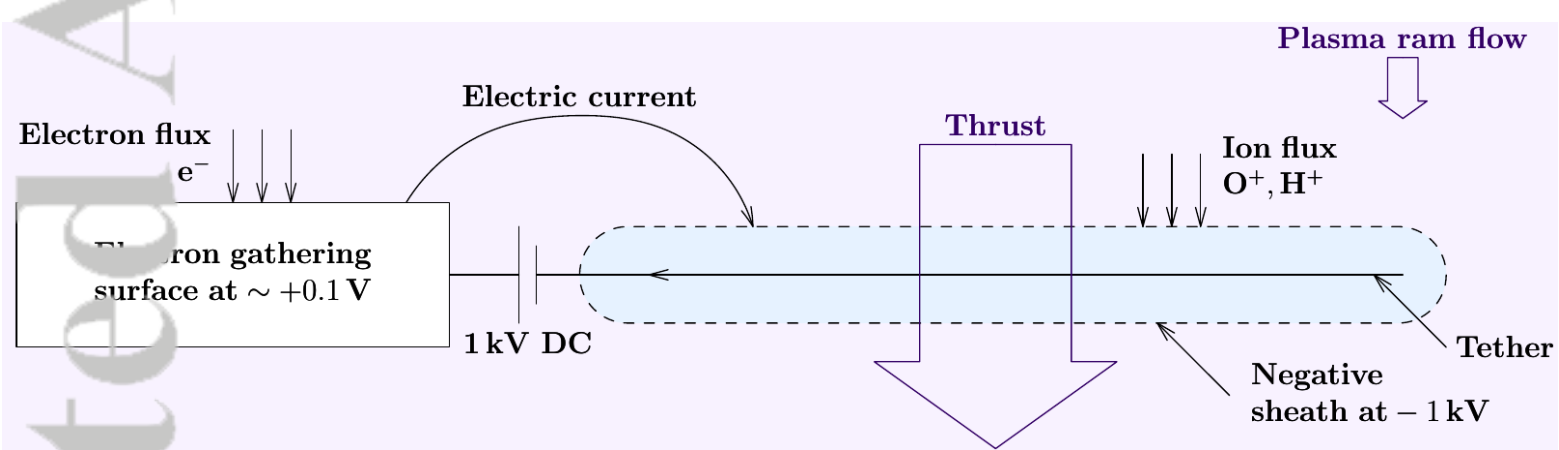


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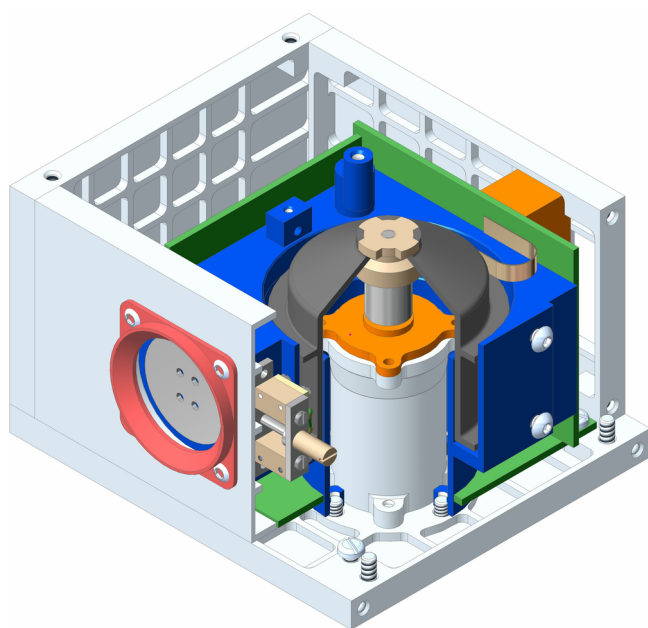


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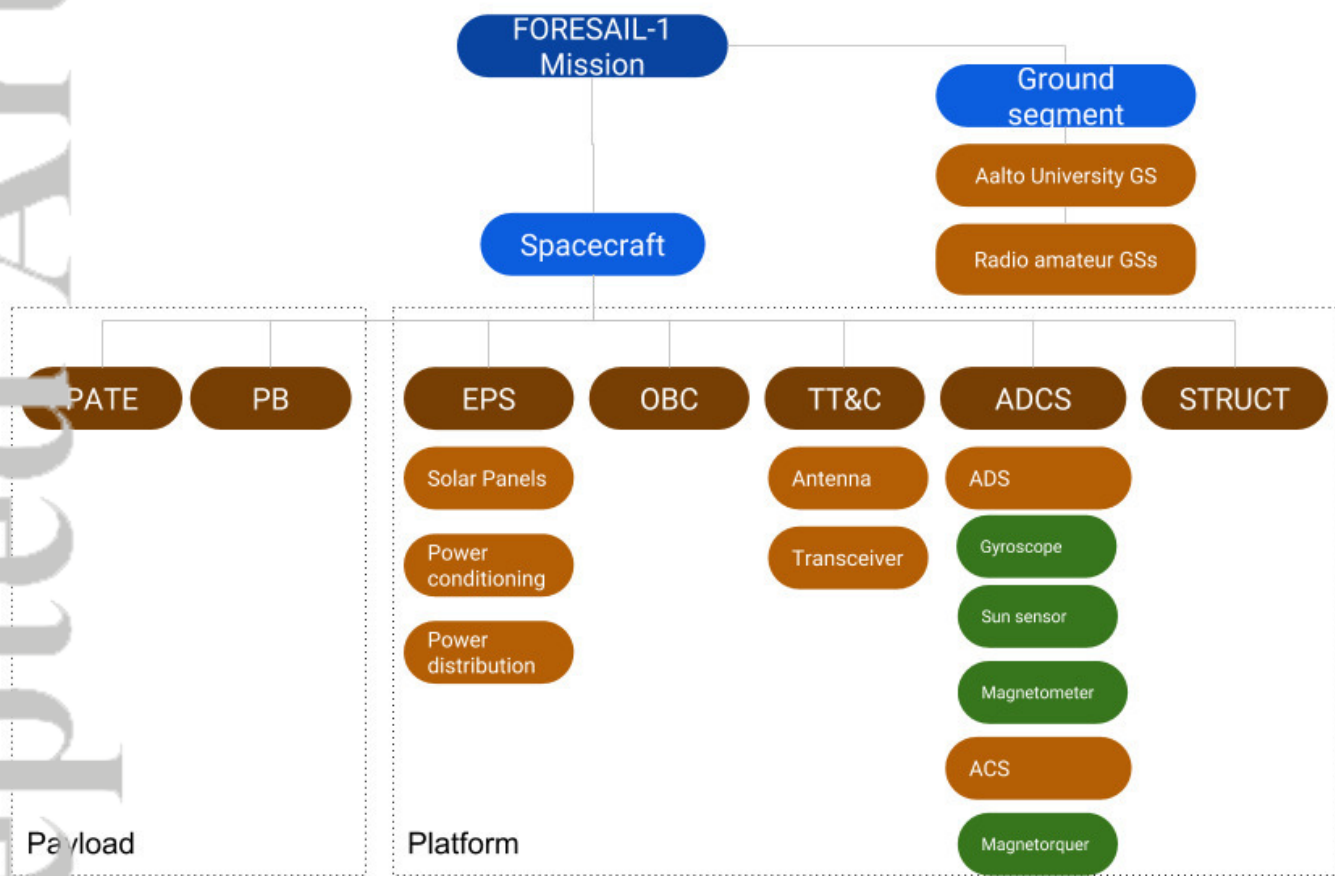




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2018JA026354-f06-z-.jpg

