# The Pristine survey - VI. The first three years of medium-resolution follow-up spectroscopy of Pristine EMP star candidates 

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#### Abstract

We present the results of a 3 -yr long, medium-resolution spectroscopic campaign aimed at identifying very metal-poor stars from candidates selected with the CaHK , metallicitysensitive Pristine survey. The catalogue consists of a total of 1007 stars, and includes 146 rediscoveries of metal-poor stars already presented in previous surveys, 707 new very metalpoor stars with $[\mathrm{Fe} / \mathrm{H}]<-2.0$, and 95 new extremely metal-poor stars with $[\mathrm{Fe} / \mathrm{H}]<-3.0$. We provide a spectroscopic $[\mathrm{Fe} / \mathrm{H}]$ for every star in the catalogue, and $[\mathrm{C} / \mathrm{Fe}]$ measurements for a subset of the stars ( 10 per cent with $[\mathrm{Fe} / \mathrm{H}]<-3$ and 24 per cent with $-3<[\mathrm{Fe} / \mathrm{H}]<-2$ ) for which a carbon determination is possible, contingent mainly on the carbon abundance, effective temperature and signal-to-noise ratio of the stellar spectra. We find an average carbon enhancement fraction ( $[\mathrm{C} / \mathrm{Fe}] \geq+0.7$ ) of $41 \pm 4$ per cent for stars with $-3<[\mathrm{Fe} / \mathrm{H}]<-2$ and $58 \pm 14$ per cent for stars with $[\mathrm{Fe} / \mathrm{H}]<-3$, and report updated success rates for the Pristine survey of 56 per cent and 23 per cent to recover stars with $[\mathrm{Fe} / \mathrm{H}]<-2.5$ and $<-3$, respectively. Finally, we discuss the current status of the survey and its preparation for providing targets to upcoming multi-object spectroscopic surveys such as William Herschel Telescope Enhanced Area Velocity Explorer.


Key words: stars: abundances - Galaxy: evolution - Galaxy: formation - Local Group - dark ages, reionization, first stars - early Universe.

## 1 INTRODUCTION

The current picture of Galactic chemical enrichment is based on the production of elements heavier than He in the interiors of stars, their subsequent release into the interstellar medium (ISM) through supernova explosions, and their eventual reintegration into ensuing stellar generations. Apart from a few exceptions, such as mass transfer binaries, the current elemental compositions of stars are expected to maintain the chemical imprint of their birth environments, which in turn reflect this enrichment process. Based

[^0]on this principle, it is possible to use stars with primitive elemental abundance patterns, also known as very metal-poor (VMP: [Fe/H] $<-2$ ), to study the early Universe.

One issue that hampers our ability to study the detailed abundance trends of metal-poor stars, is their scarcity in our local environment with respect to the younger, more metal-rich populations. However, metal-poor stars are more abundant in certain Galactic environments, making them promising searching grounds. Cosmological simulations demonstrate that the outer regions of the Galaxy are the most dominated by old and/or metal-poor stars (see for recent studies using hydrodynamical simulations Starkenburg et al. 2017a and El-Badry et al. 2018). If one has a good method to efficiently distinguish metal-poor from more metal-rich populations and is
interested in the oldest stars among the most metal-poor, then the Galaxy's inner regions and some of its satellites are also promising hunting grounds (e.g. White \& Springel 2000; Tumlinson 2010; Starkenburg et al. 2017a).
Naturally, a substantial amount of effort has gone into finding and studying these rare stars, and they remain a strong focus of current and future surveys dedicated to Galactic Archaeology. Given that they are so rare among the far more numerous foreground populations, there are two options when searching for metal-poor stars: (i) observing a large sample of stars from general science purpose surveys to find the few metal-poor stars among them, or (ii) targeted searches which aim for these stars specifically. The former approach has been quite successful and has contributed significantly to the current sample of the most metal-poor stars (e.g. Caffau et al. 2013; Aoki et al. 2013; Allende Prieto et al. 2015; Aguado et al. 2016; Li et al. 2015; Aguado et al. 2017a,b, 2018a,b), mostly with the help of large spectroscopic surveys such as the Sloan Digital Sky Survey (SDSS, York et al. 2000), the Sloan Extension for Galactic Understanding and Exploration (SEGUE, Yanny et al. 2009), the Baryonic Oscillations Spectroscopic Survey (BOSS, Eisenstein et al. 2011; Dawson et al. 2013), and more recently the large Sky Area Multi-Object Fiber Spectroscopic Telescope (LAMOST, Deng et al. 2012). More targeted searches have also been in use for many years, from early efforts using a Ca H \& K objective-prism technique, such as the HK survey (Beers, Preston \& Shectman 1985, 1992) and the Hamburg ESO survey (Christlieb, Wisotzki \& Graßhoff 2002), to more recent efforts using targeted narrow-/medium-band photometry at blue wavelengths, like the SkyMapper survey (Keller et al. 2007; Wolf et al. 2018; Casagrande et al. 2019; Huang et al. 2019), and the Pristine survey (Starkenburg et al. 2017b).

Future metal-poor star searches will be even more effective by combining both of these strategies. The upcoming generation of multi-object spectroscopic (MOS) surveys such as the William Herschel Telescope Enhanced Area Velocity Explorer (WEAVE, Dalton et al. 2018), the Dark Energy Spectroscopic Instrument (DESI, Levi et al. 2013), the 4-metre Multi-Object Spectroscopic Telescope (4MOST, de Jong et al. 2019), the Galactic Archaeology with HERMES (Buder et al. 2018), the SDSS-V (Kollmeier et al. 2017), and the Maunakea Spectroscopic Explorer (McConnachie et al. 2016; McConnachie 2019) will provide, together with Gaia (Gaia Collaboration et al. 2018), an unprecedented number of spectra over the whole sky. Although these surveys will have the capability to observe tens of millions of stars, it will still be necessary to target metal-poor stars specifically in order to maximize the output for Galactic Archaeology studies. When used in tandem with pre-selection surveys such as SkyMapper, and Pristine, it will be possible to obtain high-quality observations of metal-poor stars across an unprecedented range of magnitudes, wavelengths, and Galactic environments. In order for the target pre-selection from such surveys to be maximally effective, they must be validated beforehand by dedicated spectroscopic follow-up programs.
In this paper, we present the results of the first three years of spectroscopic follow-up for the Pristine survey, using low- and medium-resolution spectroscopic facilities. This not only provides a detailed understanding of the selection of candidates to target with future MOS surveys, but also has the added value of providing the Galactic Archaeology community with a sizeable catalogue of new, metal-poor stars, a subset of which also have measurements of carbon abundances.

Carbon abundance is a well-studied quantity in metal-poor stars, and has important implications for understanding the earliest stellar generations. First, the carbon abundance of a star influences the cooling channels and may allow for low-mass star formation (Bromm \& Loeb 2003a). Secondly, as discussed at length in Beers \& Christlieb (2005), Yong et al. (2013), Bonifacio et al. (2015), Yoon et al. (2016), and Chiaki \& Wise (2019), the increase in carbon enhancement with decreasing metallicity in Extremely Metal-Poor (EMP) stars allow us to make a phenomenological taxonomy of ancient stars.
There are two definitions for carbon-enhanced metal-poor (CEMP) stars currently presented in the literature. Beers \& Christlieb (2005) propose a definition of CEMP stars as stars with $[\mathrm{C} / \mathrm{Fe}]>+1.0,{ }^{1}$ while Aoki et al. (2007) use $[\mathrm{C} / \mathrm{Fe}]>+0.7$ with an additional correction depending on the luminosity. These different values do not reflect theoretical studies but still provide a useful quantitative classification. On the other hand, the original critical carbon abundance from Bromm \& Loeb (2003b) $\left([\mathrm{C} / \mathrm{H}]_{\text {crit }} \simeq-3.5 \pm 0.1\right)$ has recently been improved to include the effect of the silicate grains in cooling processes allowing for fragmentation of the proto-stellar clouds (Chiaki, Tominaga \& Nozawa 2017). These studies propose three regions in the $\mathrm{A}(\mathrm{C})-[\mathrm{Fe} / \mathrm{H}]$ plane: the carbon-dominated, the silicate-dominated area, and the forbidden area due the insufficient dust cooling. So far only one star, J1029 +1729 , belonging to the $[\mathrm{Fe} / \mathrm{H}]<-4.5$ regime is clearly carbon normal (Caffau et al. 2011) with $[\mathrm{C} / \mathrm{Fe}]<+0.7$. $\mathrm{J} 1029+1729$ is still the most metal-poor star known but remains in the silicate-dominated region well below $[\mathrm{C} / \mathrm{Fe}]=+2.3$ line. Discovered by Starkenburg et al. (2018) and included in this work, Pristine $221.8781+9.7844$ is the second most metal-poor star also in the silicate dominated region with $[\mathrm{C} / \mathrm{Fe}]<+1.76$ and could also potentially be a carbon-normal ultra-metal-poor (UMP) star. All 11 other stars from the literature with $[\mathrm{Fe} / \mathrm{H}]<-4.5$ show a clear enhancement in carbon (see e.g. Bonifacio et al. 2018b; Yoon et al. 2019, and references therein). Larger samples of extremely metal-poor stars, especially those with robust carbon measurements, are important in order to better understand these trends.

The paper is organized as follows. In Section 2, we summarize the data set, observations, and reduction methods. In Section 3, the analysis of the data using the FERRE code is explained. In Section 4, we present the spectroscopic follow-up catalogue, including a discussion of the updated success rates for finding EMP and VMP stars of the Pristine survey. In Section 5, we look at the future of Pristine and its synergies with other upcoming surveys, and we conclude the paper in Section 6.

## 2 DATA AND OBSERVATIONS

As discussed in detail in Starkenburg et al. (2017b), one of the main aims of the Pristine project is to enlarge the number of metalpoor stars currently known in our Galaxy and characterize them to better understand the Galactic halo. Fig. 1 shows the current Pristine footprint which covers a total of $\sim 5000 \mathrm{deg}^{2}$ in the Northern Galactic halo. The targets selected for follow-up spectroscopy are shown in cyan, and were selected form a $\sim 2500 \mathrm{deg}^{2}$ region of the total footprint.

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Figure 1. The current footprint of the Pristine survey (black) covering $\sim 5000 \mathrm{deg}^{2}$. The stars making up the spectroscopic follow-up sample are plotted as cyan points and were selected from $\sim 2500 \mathrm{deg}^{2}$ of the total region. The Galactic plane is shown as the black line.


Figure 2. Distribution of $V$ magnitudes for the full follow-up spectroscopic sample of 1007 Pristine stars.

### 2.1 Observations

The spectroscopic data presented here were collected over a period of six semesters, from 2016 March to 2019 February. Fig. 2 shows the $V$-band magnitude distribution of the spectroscopic followup sample, totalling 1008 stars. Due to the wide range in target brightness, three different facilities were used to conduct follow-up observations of EMP candidates selected from the Pristine survey: the Intermediate Dispersion Spectrograph (IDS) on the $2.5-\mathrm{m}$ Isaac Newton Telescope (INT), the Intermediate-dispersion Spectrograph and Imaging System (ISIS, Jorden 1990) on the $4.2-\mathrm{m}$ William Herschel Telescope (WHT), and the ESO Faint Object Spectrograph and Camera (EFOSC2, Buzzoni et al. 1984) on the 3.6-m New Technology Telescope (NTT). The selected mode in all cases was long slit providing low- and medium-resolution spectroscopy (see Table 1 for further technical details).

Fainter targets ( $g>16.2$ ) were observed with the larger aperture WHT and NTT telescopes, while brighter targets ( $g<16.2$ ) were observed with the INT. The total number of observing nights were 182 ( 145 with IDS, 25 with ISIS, and 12 with EFOSC). Although the ISIS observations were shared with another program so that the resulting equivalent observing nights came out to $\sim 10$.

### 2.2 Observational strategy

The minimum desired signal-to-noise ( $\mathrm{S} / \mathrm{N}$ ) ratio per pixel for the observations was $\sim 15-25$ in the calcium $\mathrm{H} \& \mathrm{~K}$ spectral region ( $\sim 3950 \AA$ ), depending on the effective temperature of the specific star. Therefore, the average exposure time for a single integration was 1500,900 , and 1500 s , for the INT, ISIS, and EFOSC observations, respectively. Naturally, exposure times varied slightly for each individual object depending on the target brightness and the visibility conditions. The observational strategy was designed to maximize the ratio between the number of observed candidates and the reliability of the derived parameters. However, stars that were identified as UMP candidates during an observing run were subsequently followed-up with more exposures to achieve a higher S/N. Stars that still seemed highly interesting at this stage were then followed up with larger telescopes at higher resolution. This observing strategy was designed to maximize the detection of very low-metallicity stars, and has yielded the discovery of Pristine $221.8781+9.7844$, an UMP sub-giant star with $[\mathrm{Fe} / \mathrm{H}]=$ $-4.66 \pm 0.13$ and $[\mathrm{C} / \mathrm{Fe}]<1.76$. The detailed analysis of this star with high-resolution follow-up with the Ultraviolet and Visual Echelle Spectrograph (UVES) at the Very Large Telescope (VLT) is described in Starkenburg et al. (2018).

### 2.3 Data reduction

The spectral data reduction included bias substraction, flat-fielding, and wavelength calibration - using $\mathrm{CuNe}+\mathrm{CuAr}$ lamps for IDS and ISIS, and $\mathrm{He}+\mathrm{Ar}$ for EFOSC--, and was performed using the ONESPEC package in IRAF (Tody 1993). At the moderate S/N levels

Table 1. Technical information for facilities used in this analysis.

| Instrument | Telescope | Detector | Grating | Dispersion <br> $\left(\AA\right.$ pixel $\left.^{-1}\right)$ | Range $(\AA)$ | Resolution at <br> $\sim 4500 \AA$ | Slit (arcsec) |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| IDS | $2.5-\mathrm{m}$ INT | EEV10 | R900V | 0.69 | $3600-5200$ | 3300 | 1.0 |
| ISIS | $4.2-\mathrm{m}$ WHT | EEV12 | R600B | 0.45 | $3600-5100$ | 2400 | 1.0 |
| EFOSC2 | $3.6-\mathrm{m}$ NTT | CCD40 | 600 | 0.95 | $3600-5200$ | 930 | $0.7-1.0$ |

required for this program and at medium resolution, the contribution of the ISM in the $\mathrm{CaH} \& \mathrm{~K}$ area is, in general, not resolved (see e.g. Aguado et al. 2016, 2017a). In order to reduce the uncertainties from the spectral analysis, we remove the bluest part of the spectrum most affected by noise, considering only the region redder than 3700 Å.

## 3 ANALYSIS WITH ferre

The entire sample of the spectroscopic data has been analysed using the grid of synthetic stellar spectra computed with the ASSET code (Koesterke, Allende Prieto \& Lambert 2008) and published in Aguado et al. (2017b, hereafter DA17). The model atmospheres were computed with the KURUCZ codes, and are described in Mészáros et al. (2012). We use the FERRE ${ }^{2}$ code (Allende Prieto et al. 2006) to search for the best fit to the observed spectrum by simultaneously deriving the main three stellar atmospheric parameters (effective temperature $T_{\text {eff }}$, surface gravity $\log g$, metallicity $[\mathrm{Fe} / \mathrm{H}]$ ), and carbon abundance $[\mathrm{C} / \mathrm{Fe}]$. FERRE is able to interpolate between the nodes of the grid and provide a synthetic spectrum for each set of derived parameters. A similar analysis in implemented in Youakim et al. (2017, hereafter KY17) used the Powel's truncated Newton algorithm to find the best-fitting solution. However, for this work we also use a Markov Chain Monte Carlo (MCMC) algorithm based on self-adaptive randomized subspace sampling (Vrug et al. 2009), which provides the added advantage of deriving uncertainties by sampling the probability distribution function. The grid of synthetic spectra spans the space $-6 \leq[\mathrm{Fe} / \mathrm{H}] \leq-2$, $-1 \leq[\mathrm{C} / \mathrm{Fe}] \leq 5,4750 \mathrm{~K} \leq \mathrm{T}_{\text {eff }} \leq 7000 \mathrm{~K}$, and $1.0 \leq \log g \leq 5.0$. Although we targeted objects in the $[\mathrm{Fe} / \mathrm{H}]_{\text {Pristine }}<-2$ regime, there were some stars that were observed with higher metallicities. Those targets were re-analysed with a more generic grid, suitable for higher metallicities, and described in Allende Prieto et al. (2018).
In order to cross-validate our analysis methods, we observed a number of well-known EMP stars from the literature that have robust stellar parameter determinations from high-resolution analyses. Comparing those stellar parameters with the ones measured in this work, we find a median deviation of $177 \mathrm{~K}, 0.86$, and 0.27 dex for $T_{\text {eff }}, \log g$, and $[\mathrm{Fe} / \mathrm{H}]$, respectively. Table 2 summarizes the FERRE analysis performed on this sample and demonstrates that our derived metallicities are in very good agreement with those from the literature, thus demonstrating the ability of our method to derive precise metallicities using medium-resolution spectra. More comparisons of stellar parameter determination with the FERRE code and standard stars in the literature can be found in DA17. Fig. 3 shows a subsample of the observed spectra together with the bestfitting synthetic spectrum as determined by FERRE for each of the three different instruments as well as three more well-known metalpoor stars.

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### 3.1 Stellar parameters

To simultaneously derive $T_{\text {eff, }} \log g,[\mathrm{Fe} / \mathrm{H}]$, and $[\mathrm{C} / \mathrm{Fe}]$, we smooth the grid of models and resample them to the appropriate resolving power corresponding to each instrument (see Table 1). We then normalize both the synthetic models and the observed spectra using a running-mean filter with a 30 -pixel window (see DA17 for further details). Finally, FERRE derives the set of parameters assuming $[\alpha / \mathrm{Fe}]=+0.4$ and a fixed value of the microturbulence of $2.0 \mathrm{~km} \mathrm{~s}^{-1}$.
$T_{\text {eff }}$ is obtained by fitting the entire spectrum, although the derived $T_{\text {eff }}$ is largely influenced by the Balmer lines present in the spectral range $\left(\mathrm{H}_{\beta}-4861 \AA, \mathrm{H}_{\gamma}-4340 \AA, \mathrm{H}_{\delta}-4101 \AA, \mathrm{H}_{\epsilon}-3970 \AA, \mathrm{H}_{\zeta}-\right.$ $3889 \AA, \mathrm{H}_{\eta}-3835 \AA, \mathrm{H}_{\theta}-3797 \AA, \mathrm{H}_{\imath}-3770 \AA, \mathrm{H}_{\kappa}-3750 \AA, \mathrm{H}_{\lambda}-3734 \AA$, $\left.\mathrm{H}_{\mu}-3721 \AA\right)$. The temperature determination method relies on the broadening theory of the Balmer lines which is described in Barklem, Piskunov \& O'Mara (2000). The running mean normalization reduces the dependence on the specific determination of the continuum, allowing improved temperature determinations based on the shape of each H line, even with a moderate $\mathrm{S} / \mathrm{N}$ ( $\sim 15-20$ ). DA17 consider a systematic uncertainty for deriving temperatures of $\delta T_{\text {eff }}=100 \mathrm{~K}$, which is then combined quadratically with the statistical error from the MCMC method. Referring back to Table 2, the derived effective temperatures are fully compatible with those from previous works. In Fig. 4, we show the relation between the photometric temperatures derived using the SDSS ( $g$ $-i)$-temperature relation ${ }^{3}$ and the temperatures derived from FERRE using the spectroscopic data from IDS, ISIS, and EFOSC.
Measuring $\log g$ values from medium-resolution spectra when no $\mathrm{Fe}_{\text {II }}$ lines are available is a challenge. Particularly at moderate $\mathrm{S} / \mathrm{N}(\sim 15-25)$, the shape of the Balmer lines alone do not allow for it to be derived precisely. However, a coarse classification between the dwarf/giant regimes is possible with FERRE. Robust $\log g$ determination in metal-poor stars using Gaia data is possible, but good quality parallax measurements are required. Since this is not available for most of our sample, particularly the fainter objects, we use the spectroscopic values from FERRE and assume the same systematic error as were used in DA17 of $\delta \log g=0.5$.
The deepest metallic absorption in the optical range is caused by the calcium $\mathrm{H} \& \mathrm{~K}$ resonant lines at 3933 and $3968 \AA$, respectively. Using these features as robust indicators to infer metallicities in EMP stars with low-/med-resolution spectroscopy is a longstanding method (see e.g. Beers et al. 1985, 1992; Ryan \& Norris 1991; Carney et al. 1996), and is still used today in large spectroscopic surveys such as SEGUE, BOSS, and LAMOST (see e.g. Caffau et al. 2013; Aguado et al. 2016, 2017a; Li et al. 2015, 2018; François et al. 2018, and references therein). However, there is additional information present in this spectral range, such as the

[^3]Table 2. FERRE analysis for a sample of well-known EMP stars. Uncertainties include both systematic and statistical errors. Values from the literature derived from high-resolution analyses are also shown.

| Object | Values from FERRE |  |  |  |  | Values from the literature |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $T_{\text {eff }}(\mathrm{K})$ | $\log g$ | [ $\mathrm{Fe} / \mathrm{H}$ ] | S/N | Inst. | $T_{\text {eff }}(\mathrm{K})$ | $\log g$ | [Fe/H] | Ref. |
| HE 0057-5959 | $5333 \pm 118$ | $1.72 \pm 0.72$ | $-3.69 \pm 0.27$ | 39 | EFOSC | $5257 \pm 100$ | $1.72 \pm 0.30$ | $-4.08 \pm 0.30$ | 1 |
| SDSS J0723 + 3637 | $5258 \pm 212$ | $2.70 \pm 1.33$ | $-3.41 \pm 0.21$ | 28 | IDS | $5150 \pm 150$ | $2.20 \pm 0.50$ | $-3.32 \pm 0.20$ | 2 |
| HD 84937 | $6379 \pm 109$ | $4.75 \pm 0.50$ | $-2.19 \pm 0.21$ | 181 | IDS | $6431 \pm 100$ | $4.08 \pm 0.30$ | $-2.14 \pm 0.20$ | 3 |
| SDSS J1004 + 3442 | $6002 \pm 140$ | $2.84 \pm 0.95$ | $-2.83 \pm 0.25$ | 13 | IDS | $6100 \pm 150$ | $4.00 \pm 0.50$ | $-3.09 \pm 0.20$ | 2 |
| SDSS J1036 + 1212 | $6052 \pm 102$ | $1.26 \pm 0.50$ | $-3.24 \pm 0.21$ | 34 | IDS | $5850 \pm 150$ | $4.00 \pm 0.50$ | $-3.47 \pm 0.20$ | 2 |
| SDSS J1108 + 1747 | $5930 \pm 104$ | $4.89 \pm 0.50$ | $-3.07 \pm 0.21$ | 35 | IDS | $6050 \pm 150$ | $4.00 \pm 0.50$ | $-3.17 \pm 0.20$ | 2 |
| SDSS J1128 + 3841 | $6416 \pm 126$ | $4.61 \pm 0.61$ | $-3.28 \pm 0.22$ | 39 | IDS | $6550 \pm 150$ | $4.00 \pm 0.50$ | $-2.82 \pm 0.20$ | 2 |
| HE 1207-3108 | $5545 \pm 156$ | $3.11 \pm 0.87$ | $-3.01 \pm 0.22$ | 93 | EFOSC | $5294 \pm 100$ | $2.85 \pm 0.30$ | $-2.70 \pm 0.30$ | 1 |
| HE 1320-2952 | $5658 \pm 123$ | $4.09 \pm 0.59$ | $-3.13 \pm 0.22$ | 50 | EFOSC | $5106 \pm 100$ | $2.26 \pm 0.30$ | $-3.69 \pm 0.30$ | 1 |
| HE 1327-2326 | $6400 \pm 109$ | $4.82 \pm 0.50$ | $-5.40 \pm 0.43$ | 30 | IDS | $6180 \pm 80$ | $4.50 \pm 0.50$ | $-5.70 \pm 0.20$ | 4 |
| G64-12 | $6435 \pm 105$ | $4.97 \pm 0.50$ | $-3.24 \pm 0.22$ | 80 | IDS | $6550 \pm 100$ | $4.68 \pm 0.30$ | $-3.21 \pm 0.20$ | 3 |
| CS 30336-0049 | $5194 \pm 161$ | $2.60 \pm 1.14$ | $-3.97 \pm 0.22$ | 51 | EFOSC | $4725 \pm 100$ | $1.19 \pm 0.30$ | $-4.10 \pm 0.30$ | 1 |
| HE 2047-5612 | $6281 \pm 122$ | $4.64 \pm 0.55$ | $-2.94 \pm 0.22$ | 41 | EFOSC | $6128 \pm 100$ | $3.68 \pm 0.30$ | $-3.14 \pm 0.30$ | 1 |
| SDSS J2206-0925 | $5210 \pm 100$ | $1.01 \pm 0.50$ | $-2.66 \pm 0.20$ | 29 | IDS | $5100 \pm 150$ | $2.10 \pm 0.50$ | $-3.17 \pm 0.20$ | 2 |
| BD+174708 | $6100 \pm 106$ | $3.90 \pm 0.50$ | $-1.80 \pm 0.21$ | 120 | IDS | $6085 \pm 50$ | $4.10 \pm 0.10$ | $-1.60 \pm 0.10$ | 5 |
| SDSS J2338-0902 | $5052 \pm 101$ | $1.03 \pm 0.50$ | $-2.62 \pm 0.20$ | 32 | IDS | $4900 \pm 150$ | $1.90 \pm 0.50$ | $-3.12 \pm 0.20$ | 2 |

Notes: References: $1=$ Yong et al. (2013); $2=$ Aoki et al. (2013); $3=$ Ishigaki, Chiba \& Aoki (2012); $4=$ Frebel et al. (2005); and 5=Gratton et al. (2003)

Mg Ib triplet and some weak Fe I and Sr II lines, and these features can also contribute to the derivation of metallicities, provided that the $\mathrm{S} / \mathrm{N}$ is high enough to resolve them. Reassuringly, we find good agreement between our $[\mathrm{Fe} / \mathrm{H}]$ values and those from highresolution analyses even with the relatively low-resolution EFOSC2 instrument ( $R \sim 1000$, see Table 2 and section 4.1 in DA17). DA17 assumed a systematic uncertainty in metallicity of 0.1 dex. However, due to the significantly lower $\mathrm{S} / \mathrm{N}$ of the current sample, the ISM contribution to the $\mathrm{Ca} \mathrm{H} \& \mathrm{~K}$ absorption lines is largely unresolved for most of the spectra. Therefore, we assume a more conservative value of $\delta[\mathrm{Fe} / \mathrm{H}]=0.2$ dex and add it to the derived statistical uncertainty.

### 3.2 Carbon abundance

Due to the lack of spectral features in EMP stars, particularly at higher $T_{\text {eff }}$, it is not always possible to derive a reliable carbon abundance, particularly with medium-resolution spectroscopy (Bonifacio et al. 2015). DA17 provide some reference levels regarding our ability to measure carbon, but as previously discussed, the average $\mathrm{S} / \mathrm{N}$ in the current sample is significantly lower. With the aim to constrain the confidence levels with which it is possible to derive carbon abundances without important systematic effects, we performed the following theoretical exercise.

A set of synthetic spectra with the same coverage as our IDS/ISIS/EFOSC data ( $3600-5200 \AA$ ) were computed with ASSET $R=3000$. A total number of 5670 spectra covering different ranges of $T_{\text {eff }}, \log g$, and absolute carbon abundance, $\mathrm{A}(\mathrm{C})$, were analysed using 10 Markov Chains of 1000 experiments each for different values of S/N ranging from 8 to 200. In total, 30854 spectra were analysed with FERRE. We then compared the synthetic absolute carbon abundance $\mathrm{A}(\mathrm{C})$ and the corresponding $[\mathrm{C} / \mathrm{Fe}]$ derived value. We marked a given trial as successful if it was able to recover the theoretical value provided by the synthetic grid, where $\mid A(C)_{\text {in }}$ $-A(C)_{\text {out }} \mid<\epsilon_{A(C)}$, with $\epsilon_{A(C)}$ the assumed systematic uncertainty of 0.2 dex as estimated in DA17. Fig. 5 shows all the ratios versus $\mathrm{S} / \mathrm{N}$ for different effective temperatures and carbon abundances. For this work, we consider the reliable areas of the plot to be those where the correct value is recovered with a frequency that is higher
than 68 per cent. For example, at solar $(A(C)=8.39$, Asplund, Grevesse \& Sauval 2005) or higher carbon abundance (red line), we are able to measure $[\mathrm{C} / \mathrm{Fe}]$ at any temperature with $\mathrm{S} / \mathrm{N}>$ 10 , while for values below $A(C)=4.4$, it is not likely to be able to detect carbon at this resolution, regardless of the $T_{\text {eff }}$. Table 3 summarizes the approximate $\mathrm{S} / \mathrm{N}$ required to detect the $G$ band in Pristine spectroscopic data. As expected, lower temperatures allow for a better carbon determination due the larger absorption of the $G$ band. We apply these cuts to the sample and only provide carbon abundance values for the 169 (i.e. 18 per cent) stars that satisfy these criteria. We note that we are able to measure carbon in 10 per cent of the stars with $[\mathrm{Fe} / \mathrm{H}]<-3$, and 24 per cent of the stars with -3 $<[\mathrm{Fe} / \mathrm{H}]<-2$.

In order to better understand the systematics involved in the determination of carbon, we assess its correlation with the determination of $\log g$. In Fig. 6, we compare the derived $[\mathrm{C} / \mathrm{Fe}]$ with those values we find if we fix $\log g=2.5$ as a function of $\log g$. The points are also coloured according to the $T_{\text {eff }}$. The most relevant part of the plot is the giant regime since this is where the majority of the sample with good $[\mathrm{C} / \mathrm{Fe}]$ determinations are located, due not only to the lower temperatures (see Table 3), but also because at these $\log g$ values carbon is more likely to be overestimated and therefore considered to be reliably determined based on the criteria in Table 3. As a result, for the stars for which we derive carbon abundances, we have systematic uncertainties which are large but well delimited, especially at $\mathrm{S} / \mathrm{N}<25$. Therefore, we assume systematic uncertainties from 0.2 up to 0.6 dex, depending on the $\mathrm{S} / \mathrm{N}$ of the spectrum and subsequent reliability of the $\log g$.

## 4 RESULTS

### 4.1 Comparison of the photometric and spectroscopic metallicities

Photometric metallicities were derived using the Pristine narrowband photometry and the SDSS broad-band photometry. The detailed methods of this procedure are described in Starkenburg et al. (2017b).


Figure 3. A subsample of the spectra of the Pristine observed targets with IDS/INT (red), ISIS/WHT (blue), and EFOSC/3.6 NNT (green) together with the best fit derived with FERRE. Three well-known metal-poor stars are shown for comparison, SDSS J0723+3637, G64-12, and CS 30336-0049. The main stellar parameters are also displayed.


Figure 4. Comparison of the photometric temperatures to those derived spectroscopically with FERRE. Stars that fall outside of the plotted region are plotted on the edge, and marked with an arrow. Histograms of both distributions are shown to indicate the density of the points.


Figure 5. The successful ratio measuring carbon abundance in a set of 30854 synthetic spectra covering a wide range of effective temperatures, absolute carbon abundance, and different levels of $\mathrm{S} / \mathrm{N}$ ratio. Dashed line is 68 per cent of successful ratio.

Fig. 7 shows the relation between the photometric and spectroscopic metallicities. In the left-hand panel, we show the total parameter space occupied by the data, and in the right-hand panel is a zoomed-in view to better show the details of the plot. We only plot the 863 stars for which there are reliable FERRE and Pristine metallicity determinations. For the former, these are stars flagged with ' $X$ ' in Table 4, described in Section 4.4. For the later,


Figure 6. An analysis over the Pristine subsample observed with IDS on the INT. The vertical axis show the difference between the original [C/Fe] determination and the one we derive assuming a fixed $\log g=2.5$ versus the derived $\log g$ colour coded by $T_{\text {eff }}$.

Table 3. The minimum $\mathrm{S} / \mathrm{N}$ values needed to detect the carbon $G$ band with $R \sim 3000$ as a function of $T_{\text {eff }}$ and $\mathrm{A}(\mathrm{C})$.

| $\mathrm{A}(\mathrm{C})$ | $<5300 \mathrm{~K}$ | $5300-5900 \mathrm{~K}$ | $5900-6450 \mathrm{~K}$ | $>6450 \mathrm{~K}$ |
| :--- | :---: | :---: | :---: | :---: |
| $\geq \odot$ | 8 | 8 | 8 | 12 |
| $8.4-7.4$ | 8 | 10 | 25 | 60 |
| $7.4-6.4$ | 10 | 25 | 100 | 800 |
| $6.4-5.4$ | 15 | 100 | - | - |
| $5.4-4.4$ | 45 | - | - | - |
| $4.4-3.4$ | 900 | - | - | - |
| $>3.4$ | - | - | - |  |

we have removed stars that exhibit variability, that may be white dwarfs, that are identified as non-point sources in their point spread functions, and that are flagged as being problematic in their SDSS $g$ or $i$ broad-band magnitudes (mainly bright sources that show some saturation). These criteria are described in greater detail in the list below fig. 3 in KY17. Here, we have omitted all criteria based on metallicity but keep all criteria pertaining to photometric quality. Many of the removed stars were observed early on in the follow-up campaign, as we were improving our selection of targets. They are, however, still included in the full table for completeness since the derived spectroscopic metallicities are not affected by the problematic photometry.

In Fig. 7, most of the stars are clustered at $-3.5<[\mathrm{Fe} / \mathrm{H}]<$ -2.0 due to our follow-up strategy of the best metal-poor candidates first. Since there are more metal-rich stars than metal-poor stars, the metal-rich stars will scatter into the metal-poor regime with a higher frequency than the other way around, and the relative contamination will be higher at the metal-poor end. The combination of this effect and the photometric selection function from the follow-up strategy produces the offset from the 1 to 1 relation (black dotted line). The right-hand panel also shows a fairly significant dispersion, but given that the uncertainties are on the order of $\sim 0.2$ dex for both the vertical and horizontal axes, it is not surprising to see a dispersion of $\sim 0.5$ dex, although the scatter seems more severe due to the small range in metallicities covered by the data, and the outliers


Figure 7. Photometric metallicities derived with Pristine ( $[\mathrm{Fe} / \mathrm{H}]_{\text {Pristine }}$ ) versus spectroscopic metallicities derived with FERRE ( $[\mathrm{Fe} / \mathrm{H}]$ FERRE) for the total sample. The red star represents Pristine $221.8781+9.7844$ from Starkenburg et al. (2018). The left-hand panel shows the full metallicity space covered by the follow-up sample, and the right-hand panel shows a zoom-in of the highest density region around $-4<[\mathrm{Fe} / \mathrm{H}]<-2$.

Table 4. Metallicities, temperatures, and carbon abundances of the Pristine spectroscopic sample. Uncertainties include systematic and statistical errors. We only include a small sample of 9 of the 1007 stars observed, to illustrate the structure of the table. The columns are described in more detail in the text. The full table is available online. The adopted solar abundances are those from Asplund et al. (2005).

| Name <br> Origin <br> Units | $V^{\dagger}$ SDSS mag | CaHK <br> Pristine mag | $[\mathrm{Fe} / \mathrm{H}]$ <br> Pristine | [ $\mathrm{Fe} / \mathrm{H}$ ] <br> FERRE | $T_{\text {eff }}$ <br> FERRE <br> K | $\log g$ <br> FERRE | [C/Fe] <br> FERRE | S/N <br> FERRE pixel $^{-1}$ | $\begin{aligned} & \text { Flag } \\ & \text { Q,C } \end{aligned}$ | Previously Observed |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| P138.xxxx+16.xxxx* | 16.06 | 16.90 | -2.84 | $-3.4 \pm 0.2$ | $5266 \pm 122$ | $4.8 \pm 0.5$ | $0.3 \pm 0.4$ | 34 | X, -1 | - |
| P149.1350+15.0447 | 16.06 | 16.38 | -2.72 | $-2.5 \pm 0.2$ | $6047 \pm 104$ | $1.1 \pm 0.5$ | $0.5 \pm 1.0$ | 49 | X, -1 | - |
| P151.4987+13.9300 | 16.67 | 16.90 | -2.84 | $-2.9 \pm 0.6$ | $5380 \pm 176$ | $1.0 \pm 0.5$ | $1.0 \pm 0.5$ | 26 | X, -1 | LAMOST,SEGUE |
| P184.xxxx+43.xxxx* | 15.92 | 16.71 | -2.94 | $-3.7 \pm 0.2$ | $5509 \pm 103$ | $4.9 \pm 0.5$ | $0.4 \pm 0.5$ | 24 | X, -1 | - |
| P185.8616+41.3093 | 15.83 | 16.71 | -2.87 | $-2.8 \pm 0.2$ | $5264 \pm 104$ | $1.4 \pm 0.6$ | $0.8 \pm 0.4$ | 25 | X,1 | SEGUE |
| P218.6977+15.5932 | 15.57 | 15.89 | -2.93 | $-3.1 \pm 0.2$ | $6305 \pm 111$ | $4.4 \pm 0.5$ | $0.7 \pm 0.8$ | 28 | X, -1 | - |
| P220.7009+13.1405 | 16.88 | 17.49 | -3.40 | $-3.3 \pm 0.2$ | $5464 \pm 111$ | $3.5 \pm 0.5$ | $1.1 \pm 0.4$ | 26 | X, -1 | - |
| P235.0067+07.1438 | 16.88 | 17.12 | -3.10 | $-3.0 \pm 0.2$ | $6216 \pm 108$ | $4.9 \pm 0.5$ | $0.0 \pm 0.7$ | 27 | X, -1 | SDSS |
| P256.0374+17.0031 | 16.51 | 17.12 | -2.97 | $-2.6 \pm 0.2$ | $5552 \pm 117$ | $3.6 \pm 0.5$ | $0.3 \pm 0.2$ | 45 | $\mathrm{X},-1$ | - |

Notes: * Coordinates of select stars have been removed as they are the subject of an ongoing high-resolution follow-up study (Kietly et al., in preparation).
${ }^{\dagger}$ Derived using SDSS $g$ and $r$ according to https://www.sdss3.org/dr8/algorithms/sdssUBVRITransform.php.
at $[\mathrm{Fe} / \mathrm{H}]_{\text {FERRE }}<-2$. We also note that it is not crucial to have a tight relation in this space, because a coarse differentiation of stars as EMP or VMP is enough to identify promising candidates for follow-up, as well as for much of the interesting ancillary science cases.

There is a distinct population of stars for which the photometric metallicities from Pristine are highly discrepant from the spectroscopic metallicities. These are seen in the left-hand panel of Fig. 7, as the tail of stars extending to $[\mathrm{Fe} / \mathrm{H}]_{\text {FERRE }}>-2$. The criteria for selecting stars for spectroscopic follow-up was investigated and summarized in detail in KY17. Despite ensuring good quality photometry, cleaning white dwarfs (cutting all stars with $\left(u_{0}-g_{0}\right)<$ $0.6)$ and variable stars, there are still 12 per cent of stars predicted to have $[\mathrm{Fe} / \mathrm{H}]_{\text {Pristine }}<-2.5$ that have $[\mathrm{Fe} / \mathrm{H}]_{\text {FERRE }}>-2$. This number rises to 18 per cent for $[\mathrm{Fe} / \mathrm{H}]_{\text {Pristine }}<-3$ (see Table 5). Many of these stars have a large temperature discrepancy between spectroscopy and photometry $\left(\left|\Delta T_{\text {eff }}\right|>500 \mathrm{~K}\right.$ for $\sim 40$ per cent
of these stars), which probably indicates problems with the SDSS broad-band photometry for these stars. This would, in turn, affect the colour, and thus the measured photometric metallicity. In addition, some of this contamination may be attributable to long-period variable stars that were not detected in the Panoramic Survey Telescope and Rapid Response System (Pan-STARRS1) variability catalogue, non-stellar objects, or chromospherically active stars with Ca H\&K in emission (although we note that only nine such objects with peculiar spectra were identified in the followup spectroscopy).
At the lowest metallicities of $[\mathrm{Fe} / \mathrm{H}]_{\text {Pristine }}<-3.5$, the percentage of stars with spectroscopic $[\mathrm{Fe} / \mathrm{H}]>-2$ rises to 57 per cent. This clearly indicates an increasing contamination fraction with decreasing metallicity. Although the slope in this region of the Milky Way metallicity distribution function is not well constrained, it is known to be quite steep, such that stars at these metallicities are incredibly rare with respect to stars of higher

Table 5. Numbers of stars with photometric predictions $[\mathrm{Fe} / \mathrm{H}]_{\text {Pristine }}$ below -2.5 and -3.0 , the numbers of stars that are spectroscopically confirmed below those metallicities, and the success rates, given for all stars with $\mathrm{S} / \mathrm{N}>25$, with the selection criteria applied (described partially in Section 4.1, and in detail in KY17), and the sample of stars with $[\mathrm{Fe} / \mathrm{H}]_{\text {Pristine }} \leq-3.0$.

|  | All stars | Selection criteria | $[\mathrm{Fe} / \mathrm{H}]_{\text {Pristine }} \leq-3.0$ |
| :--- | :---: | :---: | :---: |
| $\mathrm{~S} / \mathrm{N}>25$ | $\mathrm{~S} / \mathrm{N}>25$ | $\mathrm{~S} / \mathrm{N}>25$ |  |
| Total number | 344 | 331 | 129 |
| $[\mathrm{Fe} / \mathrm{H}]_{\text {Pristine }} \leq-2.5$ | $325 / 344(94$ per cent $)$ | $315 / 331(95$ per cent $)$ | $129 / 129(100$ per cent $)$ |
| $[\mathrm{Fe} / \mathrm{H}]_{\text {Pristine }} \leq-3.0$ | $132 / 344(38$ per cent $)$ | $129 / 331(39$ per cent $)$ | $129 / 129(100$ per cent $)$ |
| $[\mathrm{Fe} / \mathrm{H}]_{\text {FERRE }} \leq-2.5$ | $184 / 344(53$ per cent $)$ | $180 / 331(54$ per cent $)$ | $76 / 129(59$ per cent $)$ |
| $[\mathrm{Fe} / \mathrm{H}]_{\text {FERRE }} \leq-3.0$ | $48 / 344(14$ per cent $)$ | $47 / 331(14$ per cent $)$ | $30 / 129(23$ per cent $)$ |
| $[\mathrm{Fe} / \mathrm{H}]_{\text {FERRE }} \geq-2.0$ | $45 / 344(13$ per cent $)$ | $39 / 331(12$ per cent $)$ | $23 / 129(18$ per cent $)$ |
| Success $[\mathrm{Fe} / \mathrm{H}] \leq-2.5$ | $178 / 325(55$ per cent $)$ | $175 / 315(56$ per cent $)$ | - |
| Success $[\mathrm{Fe} / \mathrm{H}] \leq-3.0$ | $30 / 132(23$ per cent $)$ | $30 / 129(23$ per cent $)$ | $30 / 129(23$ per cent $)$ |

Table 6. Number of candidate stars in different magnitude bins and metallicity ranges. The first number in each cell is the number of stars followed up with spectroscopy from the sample in this paper, and the second is the total number of candidates as of the time of publication over the $\sim 2500 \mathrm{deg}^{2}$ of the Pristine footprint used to select candidates that are the focus of this paper. $[\mathrm{Fe} / \mathrm{H}]$ values shown are photometric Pristine metallicities.

|  | Number of <br> candidates | $[\mathrm{Fe} / \mathrm{H}] \leq-2.5$ | $[\mathrm{Fe} / \mathrm{H}] \leq-3.0$ |
| :--- | :---: | :---: | :---: |
| $V<15$ | $169 / 509$ | $139 / 293$ | $66 / 92$ |
| $15<V<16$ | $536 / 1809$ | $475 / 989$ | $160 / 206$ |
| $16<V<17$ | $246 / 5423$ | $238 / 2785$ | $148 / 540$ |
| $17<V<18$ | $57 / 14682$ | $56 / 7321$ | $43 / 1393$ |
| $18<V<19$ | $0 / 35036$ | $0 / 16887$ | $0 / 3977$ |
| Total | $1008 / 57459$ | $908 / 28275$ | $417 / 6208$ |

metallicity. As a result, even a small number of interloping higher metallicity stars can dominate the candidate sample at these low metallicities.

### 4.2 Updated purity and success rates of the Pristine survey

The success rates of the Pristine survey were reported after the first year of spectroscopic follow-up using a sample of 205 stars observed at medium resolution at the WHT and INT (KY17). Due to the small size of that sample, the success rates for finding metalpoor stars computed from them were preliminary estimates. Now that we have a larger follow-up sample of nearly five times as many stars, we can update these numbers with better statistics. In order to remain consistent and to allow for an easy comparison, we will use the same metrics to quantify the purity and success rates as were used in KY17, namely:

$$
\text { success rate per cent }=\frac{[\mathrm{Fe} / \mathrm{H}]_{\mathrm{FERRE}}<\mathrm{X}}{[\mathrm{Fe} / \mathrm{H}]_{\text {Pristine }}<\mathrm{X}} \times 100
$$

where FERRE refers to the spectroscopically derived $[\mathrm{Fe} / \mathrm{H}]$, Pristine to the photometric prediction by Pristine, and X the metallicity limit of interest.

For all of the stars included in Table 4, we did not make a cut in $\mathrm{S} / \mathrm{N}$, but rather checked by eye the goodness of the fit for the synthetic spectrum by FERRE. The reason for this is because stars that are cooler and more metal-rich have larger absorption lines, and are therefore easier to identify at lower $\mathrm{S} / \mathrm{N}$ than stars that are hotter and more metal-poor. As a result, we successfully determine the metallicities for more stars at higher metallicities $([\mathrm{Fe} / \mathrm{H}]>$
-2 ) with low $\mathrm{S} / \mathrm{N}$ values ( $\mathrm{S} / \mathrm{N}<15$ ), rather than to cut these stars out with an $\mathrm{S} / \mathrm{N}$ cut. However, for the calculation of the success rates, this would bias our sample with more metal-rich stars and fewer metal-poor stars. Therefore, we compute the success rates using only stars with $\mathrm{S} / \mathrm{N}>25$, the regime in which we can reliably measure metallicities, even at $[\mathrm{Fe} / \mathrm{H}]<-3$. Taking this sample, we find a success rate of 23 per cent for finding stars with $[\mathrm{Fe} / \mathrm{H}]$ $<-3.0$, and 56 per cent for finding stars with $[\mathrm{Fe} / \mathrm{H}]<-2.5$. In KY17, we reported a success rate of 22 per cent for $[\mathrm{Fe} / \mathrm{H}]<$ -3.0 , and 70 per cent for $[\mathrm{Fe} / \mathrm{H}]<-2.5$. This discrepancy can be attributed to the cut at $\mathrm{S} / \mathrm{N}>25$. If we make the same cut in the KY17 sample, this decreases the success rates to 20 per cent and 58 per cent for $[\mathrm{Fe} / \mathrm{H}]<-3$ and $<-2.5$, respectively, meaning that these values are fully compatible with what we find in this work. In KY17, we did not originally make a cut at $\mathrm{S} / \mathrm{N}>25$ when computing the success rates as this would have reduced the sample from 205 down to 62 stars, leading to uncertainties of low number statistics. In the current work, making this cut still leaves 331 stars, and still allows for a robust determination of the success rates.

We therefore update the success rates of the Pristine survey to 23 per cent for $[\mathrm{Fe} / \mathrm{H}]<-3.0$, and 56 per cent for $[\mathrm{Fe} / \mathrm{H}]$ $<-2.5$. These values, along with other diagnostics, such as the contamination rate (fraction of stars with $[\mathrm{Fe} / \mathrm{H}]>-2$ ) are summarized in Tables 5 and 6.

### 4.3 The carbon-enhancement present in the sample

Fig. 8 shows the distribution of absolute carbon ( $\mathrm{A}(\mathrm{C})$; bottom panel) and $[\mathrm{C} / \mathrm{Fe}]$ abundances as a function of metallicity for the 169 stars for which we are able to make a reliable carbon determination (Section 3.2). Both CEMP reference lines at $[\mathrm{C} / \mathrm{Fe}]=1.0$ (Beers \& Christlieb 2005) and $[\mathrm{C} / \mathrm{Fe}]=+0.7$ (Aoki et al. 2007) are plotted as solid and dashed lines, respectively. The high-resolution carbon abundance value for Pristine $221.8781+9.7844$ is also included as the red star.

To compute the CEMP fractions, we first draw a new sample of values for the $[\mathrm{C} / \mathrm{Fe}]$, and $[\mathrm{Fe} / \mathrm{H}]$ measurements, taking into account both the statistical and systematic uncertainties of each. We then compute the fraction of stars with $[\mathrm{C} / \mathrm{Fe}]$ above the two limits of $[\mathrm{C} / \mathrm{Fe}]=+1.0$ and +0.7 , and repeat this exercise $10^{6}$ times in a Monte Carlo fashion. The resulting distributions are approximately Gaussian in shape, and are therefore reasonably well described by a mean and standard deviation. For the $[\mathrm{Fe} / \mathrm{H}]<-3$ sample, we compute CEMP fractions of $58 \pm 14$ per cent and $43 \pm 13$ per cent for the $[\mathrm{C} / \mathrm{Fe}]>+0.7$ and +1.0 , respectively. For the $-3<[\mathrm{Fe} / \mathrm{H}]$


Figure 8. Carbon versus iron for $[\mathrm{C} / \mathrm{Fe}]$ (top) and absolute carbon (bottom). The red star represents Pristine $221.8781+9.7844$ with parameters derived from the analysis of a high-resolution UVES spectrum (Starkenburg et al. 2018). The dashed-dotted line at $[\mathrm{C} / \mathrm{Fe}]=0$ shows the solar carbon abundance, the dashed and solid lines at $[\mathrm{C} / \mathrm{Fe}]=+0.7$ and +1.0 show the thresholds for carbon enhancement defined in Aoki et al. (2007) and Beers \& Christlieb (2005), respectively, and the dotted line at $[\mathrm{C} / \mathrm{Fe}]=2.3$ shows the boundary of the silicate-dominated region, as described in Chiaki et al. (2017). The error bars shown in the top left of each panel represent the median errors of the sample.
$<-2$ sample, we compute CEMP fractions of $41 \pm 4$ per cent and $23 \pm 3$ per cent for $[\mathrm{C} / \mathrm{Fe}]>+0.7$ and +1.0 , respectively.

Placco et al. (2014) find 43 per cent of stars to have $[\mathrm{C} / \mathrm{Fe}] \geq$ +0.7 and $-4<[\mathrm{Fe} / \mathrm{H}]<-3$, a value which differs at the $1 \sigma$ level compared to the $58 \pm 14$ per cent derived in this work, and is therefore not statistically significant. Furthermore, Norris \& Yong (2019) perform a rigorous analysis of the 3D and NLTE (non-local thermodynamic equilibrium) corrections relevant for the carbon abundance determinations, and demonstrate a significant decrease in the carbon content for a number of CEMP stars from the literature when full 3D-NLTE corrections are taken into account. The CEMPno group are stars that do not show significant enrichment in neutron-capture elements (s- and r-processes), and are the most numerous subgroup among CEMP stars. As a result of those 3D and NLTE corrections in Norris \& Yong (2019), a significant number of CEMP-no stars become carbon-normal. However, we do not know the fraction CEMP-no stars in our current sample, but if we consider that a similar fraction of them likely are, as is the case
in the literature, it is likely that the computed CEMP fractions would decrease considerably. It is therefore difficult to draw firm conclusions from this current sample of CEMP stars, but further, more detailed follow-up - particularly targeting carbon and the neutron capture elements in the EMP stars - could potentially be a very nice sample with which to investigate this further.

### 4.4 The full sample

In this paper we present a full catalogue from three years of followup spectroscopy of Pristine candidates. The full table, consisting of 1007 stars is available online. An abbreviated version of the full table showing the provided columns as well as a sample of nine rows is shown in Table 4. The column CaHK is the magnitude obtained from the Pristine narrow-band filter, the column $[\mathrm{Fe} / \mathrm{H}]_{\text {Pristine }}$ is the photometric metallicity determined using the $(g-i)_{0}$ SDSS colours and Pristine photometry (described in Starkenburg et al. 2017 b, section 3.2). The next two columns are the spectroscopic metallicities, effective temperatures, and surface gravities derived from FERRE and their associated uncertainties. Column $\mathrm{S} / \mathrm{N}$ is the signal-to-noise ratio of the analysed spectrum. We also provide a Q-flag, representing the reliability of the spectroscopic metallicity determination. An entry of ' X ' indicates that the synthetic spectral fit was reliable and that the given $[\mathrm{Fe} / \mathrm{H}]$ FERRE value can be trusted to within the provided uncertainties ( 93 per cent of the sample have this flag). In order to provide as much information as possible, we also provide tentative metallicity values for stars for which the $\mathrm{S} / \mathrm{N}$ is too low for a robust determination of stellar parameters, but that still have some information in the observed spectrum. These stars are given a flag of $T$ ( 6 per cent of the sample), and are good candidates to be re-observed with higher $\mathrm{S} / \mathrm{N}$ and at higher resolution facilities. The C-flag shows if the carbon determination is reliable (value 1) or not (value -1 ), and was derived based on $\mathrm{S} / \mathrm{N}$ and temperature criteria described in Section 3.2. The last column indicates whether the object was already spectroscopically observed by other surveys. Finally, the object coordinates are contained in the name, but we provide these explicitly as RA and Dec. in degrees on the online version of the table.

There are a small number of stars for which the Pristine metallicity classification fails, meaning that from photometry the object was expected to be a metal-poor star, but from spectroscopy it was determined to be some other type of object. These could be stars with CaHK in emission, non-stellar objects, or various other objects with unusual behaviour in the CaHK region. However, this only occurs for nine of the observed objects ( $<1$ per cent of the sample), indicating that the sample is well cleaned. We remove these nine objects from the catalogue since both their photometric and spectroscopic metallicities are unreliable, but consider them in the sample when computing the success rates since they do contribute to the contamination.

## 5 FUTURE OF THE SURVEY

In addition to hunting for the most metal-poor stars in the Galaxy, the photometric metallicities that are produced by the narrow-band photometry of the Pristine survey can be used for several other interesting science cases. For instance, Longeard et al. (2018) conducted an in-depth study of the metallicity distribution and velocity dispersion of the faint Milky Way satellite Draco II using CaHK photometry, and work is ongoing on a similar analysis to characterize the properties of many other nearby satellites (Longeard et al. 2019). Another study by Starkenburg et al. (in
preparation) demonstrated the powerful capabilities of the Pristine narrow-band filter to identify blue horizontal branch stars and disentangle them from the contaminating blue straggler population, providing a uniquely clean sample of distance indicators with which to study the outer reaches of the Galactic halo. Finally, Arentsen et al. (in preparation) are studying the metal-poor component of the Galactic bulge with the Pristine Inner Galaxy Survey.

### 5.1 Pristine and Gaia

The highly anticipated Gaia data have initiated a revolution in the study of galactic archaeology and it is changing our understanding of the Galaxy (Gaia Collaboration et al. 2018). The latest data release provided for high-precision astrometry measurements and threefilter photometry for over 1.3 billion. The range of possibilities for using Gaia photometry together with more than 5 million Pristine metallicity determinations are broad, and open the door to an unprecedented mapping of the Galaxy using the full 6D phasespace plus metallicity information. For example, work is ongoing using Gaia and Pristine to study the substructures present in and around our Galaxy and their dependence on metallicity, as well as an analysis of the metallicity distribution function of the halo at the lowest metallicities (Youakim et al. in preparation). On the other hand, exquisite Gaia parallaxes, proper motions, and photometry allow us to derive surface gravities, effective temperatures, and orbits for EMP stars (see e.g. Bonifacio et al. 2018a; Frebel et al. 2019; Sestito et al. 2019). The dynamics of the most ancient stars of the Milky Way could be a crucial piece of information for understanding the formation and evolution of the Galactic halo. For example, recent work by Sestito et al. (2019), demonstrated that an important fraction of the known UMP stars seem to have orbits that are confined to Galactic plane, suggesting interesting new scenarios for their origins. In addition, a complete kinematical analysis of the sample presented in this paper will be presented in Sestito et al. (2019). Finally, Bonifacio et al. (2019) combined Gaia parallaxes and Pristine photometry to derive photometric metallicities, effective temperatures, and surface gravities. These authors also studied the chemical composition and ages of 40 metalpoor stars with the SOPHIE high-resolution spectrograph.

### 5.2 Pristine and WEAVE

The impending arrival of the new large spectroscopic surveys will nicely complement the still ongoing Gaia project. A new, deeper view - not only kinematically but also chemically - of the Milky Way halo, will shed light on the formation and evolution of the Galaxy. This unprecedented amount of high-quality data will greatly expand the capabilities of the Galactic archaeology community thanks to surveys like 4MOST, DESI, or WEAVE. The success rates presented in Table 5 demonstrate that the Pristine filter is one of the best ways to pre-select EMP candidates to observe in those surveys. In particular, the WEAVE project will devote up to 20 fibers per WEAVE $3.14 \mathrm{deg}^{2}$ field of view to Pristine-selected EMP candidates in the magnitude range $15<G<19$, in the lowresolution Galactic archaeology survey of high Galactic latitudes (Jin et al., in preparation). Over the planned $\sim 8500 \mathrm{deg}^{2}$ of the survey, of which we anticipate $\geq 5000 \mathrm{deg}^{2}$ will be in common with the Pristine footprint at the time they are observed in WEAVE, this adds up to up to $\sim 30000$ candidate EMP stars, of which according to Table 5, $\sim 5000-7000$ would turn out to be $[\mathrm{Fe} / \mathrm{H}]<$ -3 . This would increase the number of spectroscopically confirmed EMP and UMP stars with known chemical signatures by one order
of magnitude. After five years of observing we expect to have measured the chemical abundances such as $\mathrm{C}, \mathrm{Na}, \mathrm{Mg}, \mathrm{Al}, \mathrm{Si}, \mathrm{Ca}$, Ti, and Fe , for about $\sim 3000$ stars with $[\mathrm{Fe} / \mathrm{H}]<-3.0$ and $\sim 150-$ 200 stars with $[\mathrm{Fe} / \mathrm{H}]<-4.0$, including $\sim 5-10$ hyper metal-poor stars $([\mathrm{Fe} / \mathrm{H}]<-5.0)$, doubling the samples currently available from several decades of efforts. Additionally, WEAVE Galactic archaeology high-resolution (HR) survey will be able to measure the full suite of chemical signatures for the brightest part of the Pristine sample ( $g \leq 15.5$ ) where it overlaps with the WEAVE HR survey, although the density of such bright targets will be much lower.

## 6 CONCLUSIONS

Expanding upon the previous work conducted in Starkenburg et al. (2017b) and KY17, we have presented a sample consisting of 1008 stars, representing three years of follow-up of medium- and low-resolution spectroscopy of EMP candidates from the Pristine survey. The number of stars followed-up spectroscopically has increased by a factor of 5, allowing for the success rate of stars with $[\mathrm{Fe} / \mathrm{H}]<-2.5$ and $<-3.0$ to be updated to 56 per cent and 23 per cent, respectively. This is a relevant milestone in the field of Galactic archaeology, demonstrating the utility of the Pristine filter to select EMP candidates for the next generation of spectroscopic surveys such as WEAVE. The recent discovery of Pristine $221.8781+9.7844$ (Starkenburg et al. 2018), the second most metal-poor star yet discovered, shows that Pristine photometry is also effective in finding UMP stars in the most interesting and poorly populated regime of $[\mathrm{Fe} / \mathrm{H}]<-4.5$. In addition, we demonstrated that the FERRE code is capable of deriving stellar parameters even at relatively low-resolution, namely with the stars observed with EFOSC2. Furthermore, we show for the first time in the Pristine project that we are able to provide individual carbon abundances from measurements of the $G$ band with moderate $\mathrm{S} / \mathrm{N}$ in medium-resolution spectra for 169 stars, or $\sim 20$ per cent of the total sample, although lower average $\mathrm{S} / \mathrm{N}$ as compared to DA17 results in higher overall uncertainties in the carbon measurements than previously achieved. With this medium-resolution follow-up spectroscopy sample (along with the previous analysis of KY17), we have been able to thoroughly characterize the photometric selection and success rates of the Pristine survey in this magnitude range, and future follow-up is planned to mostly be done with MOS facilities such as WEAVE. More observations with low- and mediumresolution spectroscopic facilities of metal-poor candidates selected from Pristine are highly desirable with the aim of increasing the number of UMP/hyper metal-poor stars, but also to provide a larger sample of CEMP and carbon-normal EMP stars.

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## SUPPORTING INFORMATION

Supplementary data are available at MNRAS online.
Table 4. Efective temperatures, surface gravities, metallicities and carbon abundances of the Pristine spectectroscopic sample.

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[^1]:    ${ }^{1}$ More recently Bonifacio et al. (2018b) proposed to establish a fixed A(C) > 5.5 reference value for stars with $[\mathrm{Fe} / \mathrm{H}]<-4.0$ to be CEMP stars.

[^2]:    ${ }^{2}$ FERRE is available from http://github.com/callendeprieto/ferre

[^3]:    ${ }^{3}$ For the equation used to compute the photometric temperatures, see the InfraRed Flux Method, https://www.sdss.org/dr12/spectro/sspp_irfm/, $[\mathrm{Fe} / \mathrm{H}]=-2.5$ was assumed.

