## Observation of the narrow state $X(3872)->\mathrm{J} / \mathrm{psi} \mathrm{pi}(+) \mathrm{pi}(-)$ in (p)over-barp collisions at root $\mathrm{s}=1.96$ TeV

D. Acosta, T. Affolder, M. H. Ahn, T. Akimoto, M. G. Albrow, D. Ambrose, S. Amerio, D. Amidei, A. Anastassov, K. Anikeev, A. Annovi, J. Antos, M. Aoki, G. Apollinari, J. F. Arguin, T. Arisawa, A. Artikov, T. Asakawa, W. Ashmanskas, A. Attal, F. Azfar, P. Azzi-Bacchetta, N. Bacchetta, H. Bachacou, W. Badgett, S. Bailey, A. Barbaro-Galtieri, G. Barker, V. E. Barnes, B. A. Barnett, S. Baroiant, M. Barone, G. Bauer, F. Bedeschi, S. Behari, S. Belforte, W. H. Bell, G. Bellettini, J. Bellinger, D. Benjamin, A. Beretvas, A. Bhatti, M. Binkley, D. Bisello, M. Bishai, R. E. Blair, C. Blocker, K. Bloom, B. Blumenfeld, A. Bocci, A. Bodek, G. Bolla, A. Bolshov, P. S. L. Booth, D. Bortoletto, J. Boudreau, S. Bourov, C. Bromberg, E. Brubaker, J. Budagov, H. S. Budd, K. Burkett, G. Busetto, P. Bussey, K. L. Byrum, S. Cabrera, P. Calafiura, M. Campanelli, M. Campbell, A. Canepa, D. Carlsmith, S. Carron, R. Carosi, M. Casarsa, A. Castro, P. Catastini, D. Cauz, A. Cerri, C. Cerri, L. Cerrito, J. Chapman, C. Chen, Y. C. Chen, M. Chertok, G. Chiarelli, G. Chlachidze, F. Chlebana, K. Cho, D. Chokheli, M. L. Chu, J. Y. Chung, W. H. Chung, Y. S. Chung, C. I. Ciobanu, M. A. Ciocci, A. G. Clark, D. Clark, M. N. Coca, A. Connolly, M. E. Convery, J. Conway, M. Cordelli, G. Cortiana, J. Cranshaw, R. Culbertson, C. Currat, D. Cyr, D. Dagenhart, S. DaRonco, S. D'Auria, P. de Barbaro, S. De Cecco, G. De Lentdecker, S. Dell'Agnello, M. Dell'Orso, S. Demers, L. Demortier, M. Deninno, D. De Pedis, P. F. Derwent, C. Dionisi, J. R. Dittmann, P. Doksus, A. Dominguez, S. Donati, M. D'Onofrio, T. Dorigo, V. Drollinger, K. Ebina, N. Eddy, R. Ely, R. Erbacher, M. Erdmann, D. Errede, S. Errede, R. Eusebi, H. C. Fang, S. Farrington, I. Fedorko, R. G. Feild, M. Feindt, J. P. Fernandez, C. Ferretti, R. D. Field, I. Fiori, G. Flanagan, B. Flaugher, L. R. Flores-Castillo, A. Foland, S. Forrester, G. W. Foster, M. Franklin, H. Frisch, Y. Fujii, I. Furic, A. Gaijar, A. Gallas, M. Gallinaro, J. Galyardt, M. Garcia-Sciveres, A. F. Garfinkel, C. Gay, H. Gerberich, E. Gerchtein, D. W.

Gerdes, S. Giagu, P. Giannetti, A. Gibson, K. Gibson, C. Ginsburg, K. Giolo, M. Giordani, G. Giurgiu, V. Glagolev, D. Glenzinski, M. Gold, N. Goldschmidt, D. Goldstein, J. Goldstein, G. Gomez, G. Gomez-Ceballos, M. Goncharov, I. Gorelov, A. T. Goshaw, Y. Gotra, K. Goulianos, A. Gresele, C. Grosso-Pilcher, M. Guenther, J. Guimaraes da Costa, C. Haber, K. Hahn, S. R. Hahn, E. Halkiadakis, C. Hall, R. Handler, F. Happacher, K. Hara, M. Hare, R. F. Harr, R. M. Harris, F. Hartmann, K. Hatakeyama, J. Hauser, C. Hays, H. Hayward, E. Heider, B. Heinemann, J. Heinrich, M. Hennecke, M. Herndon, C. Hill, D. Hirschbuehl, A. Hocker, K. D. Hoffman, A. Holloway, S. Hou, M. A. Houlden, Y. Huang, B. T. Huffman, R. E. Hughes, J. Huston, K. Ikado, J. Incandela, G. Introzzi, M. Iori, Y. Ishizawa, C. Issever, A. Ivanov, Y. Iwata, B. Iyutin, E. James, D. Jang, J. Jarrell, D. Jeans, H. Jensen, M. Jones, S. Y. Jun, T. Junk, T. Kamon, J. Kang, M. K. Unel, P. E. Karchin, S. Kartal, Y. Kato, Y. Kemp, R. Kephart, U. Kerzel, V. Khotilovich, B. Kilminster, B. J. Kim, D. H. Kim, H. S. Kim, J. E. Kim, M. J. Kim, M. S. Kim, S. B. Kim, S. H. Kim, T. H. Kim, Y. K. Kim, B. T. King, M. Kirby, L. Kirsch, S. Klimenko, B. Knuteson, B. R. Ko, H. Kobayashi, P. Koehn, K. Kondo, J. Konigsberg, K. Kordas, A. Korn, A. Korytov, K. Kotelnikov, A. V. Kotwal, A. Kovalev, J. Kraus, I. Kravchenko, A. Kreymer, J. Kroll, M. Kruse, V. Krutelyov, S. E. Kuhlmann, N. Kuznetsova, A. T. Laasanen, S. Lai, S. Lami, S. Lammel, J. Lancaster, M. Lancaster, R. Lander, K. Lannon, A. Lath, G. Latino, R. Lauhakangas, I. Lazzizzera, Y. Le, C. Lecci, T. LeCompte, J. Lee, J. Lee, S. W. Lee, N. Leonardo, S. Leone, J. D. Lewis, K. Li, C. S. Lin, M. Lindgren, T. M. Liss, D. O. Litvintsev, T. Liu, Y. Liu, N. S. Lockyer, A. Loginov, J. Loken, M. Loreti, P. Loverre, D. Lucchesi, P. Lukens, L. Lyons, J. Lys, D. MacQueen, R. Madrak, K. Maeshima, P. Maksimovic, L. Malferrari, G. Manca, R. Marginean, A. Martin, M. Martin, V. Martin, M. Martinez, T. Maruyama, H. Matsunaga, M. Mattson, P. Mazzanti, K. S. McFarland, D. McGivern, P. M. McIntyre, P. McNamara, R. McNulty, S. Menzemer, A. Menzione, P. Merkel, C. Mesropian, A. Messina, A. Meyer, T. Miao, N. Miladinovic, L. Miller, R. Miller, J. S. Miller, R. Miquel, S. Miscetti, M. Mishina, G. Mitselmakher, A. Miyamoto, Y. Miyazaki, N. Moggi, R. Moore, M. Morello, T. Moulik, A. Mukherjee, M. Mulhearn, T. Muller, R. Mumford, A. Munar, P. Murat, J. Nachtman, S. Nahn, I. Nakamura, I. Nakano, A. Napier, R. Napora, V. Necula, F. Niell, J. Nielsen, C. Nelson, T. Nelson, C. Neu, M. S. Neubauer, C. Newman-Holmes, A. S. Nicollerat, T. Nigmanov, L. Nodulman, K. Oesterberg, T. Ogawa, S. Oh, Y. D. Oh, T. Ohsugi, R. Oishi, T. Okusawa, R. Oldeman, R. Orava, W. Orejudos, C. Pagliarone, F. Palmonari, R. Paoletti, V. Papadimitriou, S. Pashapour, J. Patrick, G. Pauletta, M. Paulini, T. Pauly, C. Paus, D. Pellett, A. Penzo, T. J. Phillips, G. Piacentino, J. Piedra, K. T. Pitts, A. Pompos, L. Pondrom, G. Pope, O. Poukhov, F. Prakoshyn, T. Pratt, A. Pronko, J. Proudfoot, F. Ptohos, G. Punzi, J. Rademacker, A. Rakitine, S. Rappoccio, F. Ratnikov, H. Ray, A. Reichold, V. Rekovic, P. Renton, M. Rescigno, F. Rimondi, K. Rinnert, L. Ristori, W. J. Robertson, A. Robson, T. Rodrigo, S. Rolli, L. Rosenson, R. Roser, R. Rossin, C. Rott, J. Russ, A. Ruiz, D. Ryan, H. Saarikko, A. Safonov, R. St Denis, W. K. Sakumoto, D. Saltzberg, C. Sanchez, A. Sansoni, L. Santi, S. Sarkar, K. Sato, P. Savard, A.

Savoy-Navarro, P. Schemitz, P. Schlabach, E. E. Schmidt, M. P. Schmidt, M. Schmitt, L. Scodellaro, A. Scribano, F. Scuri, A. Sedov, S. Seidel, Y. Seiya, F. Semeria, L. Sexton-Kennedy, I. Sfiligoi, M. D. Shapiro, T. Shears, P. F. Shepard, M. Shimojima, M. Shochet, Y. Shon, A. Sidoti, M. Siket, A. Sill, P. Sinervo, A. Sisakyan, A. Skiba, A. J. Slaughter, K. Sliwa, J. R. Smith, F. D. Snider, R. Snihur, S. V. Somalwar, J. Spalding, M. Spezziga, L. Spiegel, F. Spinella, M. Spiropulu, P. Squillacioti, H. Stadie, B. Stelzer, O. Stelzer-Chilton, J. Strologas, D. Stuart, A. Sukhanov, K. Sumorok, H. Sun, T. Suzuki, A. Taffard, R. Tafirout, S. F. Takach, H. Takano, R. Takashima, Y. Takeuchi, K. Takikawa, P. Tamburello, M. Tanaka, R. Tanaka, N. Tanimoto, S. Tapprogge, M. Tecchio, P. K. Teng, K. Terashi, R. J. Tesarek, S. Tether, J. Thom, A. S. Thompson, E. Thomson, R. Thurman-Keup, P. Tipton, V. Tiwari, S. Tkaczyk, D. Toback, K. Tollefson, D. Tonelli, M. Tonnesmann, S. Torre, D. Torretta, W. Trischuk, J. Tseng, R. Tsuchiya, S. Tsuno, D. Tsybychev, N. Turini, M. Turner, F. Ukegawa, T. Unverhau, S. Uozumi, D. Usynin, L. Vacavant, T. Vaiciulis, A. Varganov, E. Vataga, S. Vejcik, G. Velev, G. Veramendi, T. Vickey, R. Vidal, I. Vila, R. Vilar, I. Volobouev, M. von der Mey, R. G. Wagner, R. L. Wagner, W. Wagner, N. Wallace, T. Walter, Z. Wan, M. J. Wang, S. M. Wang, A. Warburton, B. Ward, S. Waschke, D. Waters, T. Watts, M. Weber, W. Wester, B. Whitehouse, A. B. Wicklund, E. Wicklund, H. H. Williams, P. Wilson, B. L. Winer, P. Wittich, S. Wolbers, M. Wolter, M. Worcester, S. Worm, T. Wright, X. Wu, F. Wurthwein, A. Wyatt, A. Yagil, T. Yamashita, K. Yamamoto, U. K. Yang, W. Yao, G. P. Yeh, K. Yi, J. Yoh, P. Yoon, K. Yorita, T. Yoshida, I. Yu, S. Yu, Z. Yu, J. C. Yun, L. Zanello, A. Zanetti, I. Zaw, F. Zetti, J. Zhou, A. Zsenei, and S. Zucchelli

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## Observation of the Narrow State $X(3872) \rightarrow J / \psi \pi^{+} \pi^{-}$in $\bar{p} p$ Collisions at $\sqrt{s}=1.96 \mathrm{TeV}$

D. Acosta, ${ }^{14}$ T. Affolder, ${ }^{7}$ M. H. Ahn, ${ }^{25}$ T. Akimoto, ${ }^{52}$ M. G. Albrow, ${ }^{13}$ D. Ambrose, ${ }^{41}$ S. Amerio, ${ }^{40}$ D. Amidei, ${ }^{31}$ A. Anastassov, ${ }^{48}$ K. Anikeev, ${ }^{29}$ A. Annovi, ${ }^{42}$ J. Antos, ${ }^{1}$ M. Aoki, ${ }^{52}$ G. Apollinari, ${ }^{13}$ J-F. Arguin, ${ }^{30}$ T. Arisawa, ${ }^{54}$ A. Artikov, ${ }^{11}$ T. Asakawa, ${ }^{52}$ W. Ashmanskas, ${ }^{2}$ A. Attal, ${ }^{6}$ F. Azfar, ${ }^{39}$ P. Azzi-Bacchetta, ${ }^{40}$ N. Bacchetta, ${ }^{40}$ H. Bachacou, ${ }^{26}$ W. Badgett, ${ }^{13}$ S. Bailey, ${ }^{18}$ A. Barbaro-Galtieri, ${ }^{26}$ G. Barker, ${ }^{23}$ V. E. Barnes, ${ }^{44}$ B. A. Barnett, ${ }^{22}$ S. Baroiant, ${ }^{5}$ M. Barone, ${ }^{15}$ G. Bauer, ${ }^{29}$ F. Bedeschi, ${ }^{42}$ S. Behari, ${ }^{22}$ S. Belforte, ${ }^{51}$ W. H. Bell, ${ }^{17}$ G. Bellettini, ${ }^{42}$ J. Bellinger, ${ }^{56}$ D. Benjamin, ${ }^{12}$ A. Beretvas, ${ }^{13}$ A. Bhatti, ${ }^{46}$ M. Binkley, ${ }^{13}$ D. Bisello, ${ }^{40}$ M. Bishai, ${ }^{13}$ R. E. Blair, ${ }^{2}$ C. Blocker, ${ }^{4}$ K. Bloom, ${ }^{31}$ B. Blumenfeld, ${ }^{22}$ A. Bocci, ${ }^{46}$ A. Bodek, ${ }^{45}$ G. Bolla, ${ }^{44}$ A. Bolshov, ${ }^{29}$ P. S. L. Booth, ${ }^{27}$ D. Bortoletto, ${ }^{44}$ J. Boudreau, ${ }^{43}$ S. Bourov, ${ }^{13}$ C. Bromberg, ${ }^{32}$ E. Brubaker, ${ }^{26}$ J. Budagov, ${ }^{11}$ H. S. Budd, ${ }^{45}$ K. Burkett, ${ }^{13}$ G. Busetto, ${ }^{40}$ P. Bussey, ${ }^{17}$ K. L. Byrum, ${ }^{2}$ S. Cabrera, ${ }^{12}$ P. Calafiura, ${ }^{26}$ M. Campanelli, ${ }^{16}$ M. Campbell, ${ }^{31}$ A. Canepa, ${ }^{44}$ D. Carlsmith, ${ }^{56}$ S. Carron, ${ }^{12}$ R. Carosi, ${ }^{42}$ M. Casarsa, ${ }^{51}$ A. Castro, ${ }^{3}$ P. Catastini, ${ }^{42}$ D. Cauz, ${ }^{51}$ A. Cerri, ${ }^{26}$ C. Cerri, ${ }^{42}$ L. Cerrito, ${ }^{21}$ J. Chapman, ${ }^{31}$ C. Chen, ${ }^{41}$ Y. C. Chen, ${ }^{1}$ M. Chertok, ${ }^{5}$ G. Chiarelli, ${ }^{42}$ G. Chlachidze, ${ }^{11}$ F. Chlebana, ${ }^{13}$ K. Cho, ${ }^{25}$ D. Chokheli, ${ }^{11}$ M. L. Chu, ${ }^{1}$ J. Y. Chung, ${ }^{36}$ W-H. Chung, ${ }^{56}$ Y. S. Chung, ${ }^{45}$ C. I. Ciobanu, ${ }^{21}$ M. A. Ciocci, ${ }^{42}$ A. G. Clark, ${ }^{16}$ D. Clark, ${ }^{4}$ M. N. Coca, ${ }^{45}$ A. Connolly, ${ }^{26}$ M. E. Convery, ${ }^{46}$ J. Conway, ${ }^{48}$ M. Cordelli, ${ }^{15}$ G. Cortiana, ${ }^{40}$ J. Cranshaw, ${ }^{50}$ R. Culbertson, ${ }^{13}$ C. Currat, ${ }^{26}$ D. Cyr, ${ }^{56}$ D. Dagenhart, ${ }^{4}$ S. DaRonco, ${ }^{40}$ S. D'Auria, ${ }^{17}$ P. de Barbaro, ${ }^{45}$ S. De Cecco, ${ }^{47}$ G. De Lentdecker, ${ }^{45}$ S. Dell'Agnello, ${ }^{15}$ M. Dell'Orso, ${ }^{42}$ S. Demers, ${ }^{45}$ L. Demortier, ${ }^{46}$ M. Deninno, ${ }^{3}$ D. De Pedis, ${ }^{47}$ P. F. Derwent, ${ }^{13}$ C. Dionisi, ${ }^{47}$ J. R. Dittmann, ${ }^{13}$ P. Doksus, ${ }^{21}$ A. Dominguez, ${ }^{26}$ S. Donati, ${ }^{42}$ M. D'Onofrio, ${ }^{16}$ T. Dorigo, ${ }^{40}$ V. Drollinger, ${ }^{34}$ K. Ebina, ${ }^{54}$ N. Eddy, ${ }^{21}$ R. Ely, ${ }^{26}$ R. Erbacher, ${ }^{13}$ M. Erdmann, ${ }^{23}$ D. Errede, ${ }^{21}$ S. Errede, ${ }^{21}$ R. Eusebi, ${ }^{45}$ H-C. Fang, ${ }^{26}$ S. Farrington, ${ }^{27}$ I. Fedorko, ${ }^{42}$ R. G. Feild, ${ }^{57}$ M. Feindt, ${ }^{23}$ J. P. Fernandez, ${ }^{44}$ C. Ferretti, ${ }^{31}$ R. D. Field, ${ }^{14}$ I. Fiori, ${ }^{42}$ G. Flanagan, ${ }^{32}$ B. Flaugher, ${ }^{13}$ L. R. Flores-Castillo, ${ }^{43}$ A. Foland, ${ }^{18}$ S. Forrester, ${ }^{5}$ G.W. Foster, ${ }^{13}$ M. Franklin, ${ }^{18}$ H. Frisch, ${ }^{10}$ Y. Fujii, ${ }^{24}$ I. Furic, ${ }^{29}$ A. Gaijar, ${ }^{27}$ A. Gallas, ${ }^{35}$ M. Gallinaro, ${ }^{46}$ J. Galyardt, ${ }^{9}$ M. Garcia-Sciveres, ${ }^{26}$ A. F. Garfinkel, ${ }^{44}$ C. Gay, ${ }^{57}$ H. Gerberich, ${ }^{12}$ E. Gerchtein, ${ }^{9}$ D.W. Gerdes, ${ }^{31}$ S. Giagu, ${ }^{47}$ P. Giannetti, ${ }^{42}$ A. Gibson, ${ }^{26}$ K. Gibson, ${ }^{9}$ C. Ginsburg, ${ }^{56}$ K. Giolo, ${ }^{44}$ M. Giordani, ${ }^{51}$ G. Giurgiu, ${ }^{9}$ V. Glagolev, ${ }^{11}$ D. Glenzinski, ${ }^{13} \mathrm{M}$. Gold, ${ }^{34} \mathrm{~N}$. Goldschmidt, ${ }^{31}$ D. Goldstein, ${ }^{6}$ J. Goldstein, ${ }^{39}$ G. Gomez, ${ }^{8}$ G. Gomez-Ceballos, ${ }^{29}$ M. Goncharov, ${ }^{49}$ I. Gorelov, ${ }^{34}$ A. T. Goshaw, ${ }^{12}$ Y. Gotra, ${ }^{43}$ K. Goulianos, ${ }^{46}$ A. Gresele, ${ }^{3}$ C. Grosso-Pilcher, ${ }^{10}$ M. Guenther, ${ }^{44}$ J. Guimaraes da Costa, ${ }^{18}$ C. Haber, ${ }^{26}$ K. Hahn, ${ }^{41}$ S. R. Hahn, ${ }^{13}$ E. Halkiadakis, ${ }^{45}$ C. Hall, ${ }^{18}$ R. Handler, ${ }^{56}$ F. Happacher, ${ }^{15}$ K. Hara, ${ }^{52}$ M. Hare, ${ }^{53}$ R. F. Harr, ${ }^{55}$ R. M. Harris, ${ }^{13}$ F. Hartmann, ${ }^{23}$ K. Hatakeyama, ${ }^{46}$ J. Hauser, ${ }^{6}$ C. Hays, ${ }^{12}$ H. Hayward, ${ }^{27}$ E. Heider, ${ }^{53}$ B. Heinemann, ${ }^{27}$ J. Heinrich, ${ }^{41}$ M. Hennecke, ${ }^{23}$ M. Herndon, ${ }^{22}$ C. Hill, ${ }^{7}$ D. Hirschbuehl, ${ }^{23}$ A. Hocker, ${ }^{45}$ K. D. Hoffman, ${ }^{10}$ A. Holloway, ${ }^{18}$ S. Hou, ${ }^{1}$ M. A. Houlden, ${ }^{27}$ Y. Huang, ${ }^{12}$ B. T. Huffman, ${ }^{39}$ R. E. Hughes, ${ }^{36}$ J. Huston, ${ }^{32}$ K. Ikado, ${ }^{54}$ J. Incandela, ${ }^{7}$ G. Introzzi, ${ }^{42}$ M. Iori, ${ }^{47}$ Y. Ishizawa, ${ }^{52}$ C. Issever, ${ }^{7}$ A. Ivanov, ${ }^{45}$ Y. Iwata, ${ }^{20}$ B. Iyutin, ${ }^{29}$ E. James, ${ }^{13}$ D. Jang, ${ }^{48}$ J. Jarrell, ${ }^{34}$ D. Jeans, ${ }^{47}$ H. Jensen, ${ }^{13}$ M. Jones, ${ }^{44}$ S. Y. Jun, ${ }^{9}$ T. Junk, ${ }^{21}$ T. Kamon, ${ }^{49}$ J. Kang, ${ }^{31}$ M. Karagoz Unel, ${ }^{35}$ P. E. Karchin, ${ }^{55}$ S. Kartal, ${ }^{13}$ Y. Kato, ${ }^{38}$ Y. Kemp, ${ }^{23}$ R. Kephart, ${ }^{13}$ U. Kerzel,,$^{23}$ V. Khotilovich, ${ }^{49}$ B. Kilminster, ${ }^{36}$ B. J. Kim, ${ }^{25}$ D. H. Kim, ${ }^{25}$ H. S. Kim, ${ }^{21}$ J. E. Kim, ${ }^{25}$ M. J. Kim, ${ }^{9}$ M. S. Kim, ${ }^{25}$ S. B. Kim, ${ }^{25}$ S. H. Kim, ${ }^{52}$ T. H. Kim, ${ }^{29}$ Y. K. Kim, ${ }^{10}$ B. T. King, ${ }^{27}$ M. Kirby, ${ }^{12}$ L. Kirsch, ${ }^{4}$ S. Klimenko, ${ }^{14}$ B. Knuteson, ${ }^{29}$ B. R. Ko, ${ }^{12}$ H. Kobayashi, ${ }^{52}$ P. Koehn, ${ }^{36}$ K. Kondo, ${ }^{54}$ J. Konigsberg, ${ }^{14}$ K. Kordas, ${ }^{30}$ A. Korn, ${ }^{29}$ A. Korytov, ${ }^{14}$ K. Kotelnikov, ${ }^{33}$ A. V. Kotwal, ${ }^{12}$ A. Kovalev, ${ }^{41}$ J. Kraus, ${ }^{21}$ I. Kravchenko, ${ }^{29}$ A. Kreymer, ${ }^{13}$ J. Kroll, ${ }^{41}$ M. Kruse, ${ }^{12}$ V. Krutelyov, ${ }^{49}$ S. E. Kuhlmann, ${ }^{2}$ N. Kuznetsova, ${ }^{13}$ A. T. Laasanen, ${ }^{44}$ S. Lai, ${ }^{30}$
S. Lami, ${ }^{46}$ S. Lammel, ${ }^{13}$ J. Lancaster, ${ }^{12}$ M. Lancaster, ${ }^{28}$ R. Lander, ${ }^{5}$ K. Lannon, ${ }^{36}$ A. Lath, ${ }^{48}$ G. Latino, ${ }^{34}$ R. Lauhakangas, ${ }^{19}$ I. Lazzizzera, ${ }^{40}$ Y. Le, ${ }^{22}$ C. Lecci, ${ }^{23}$ T. LeCompte, ${ }^{2}$ J. Lee, ${ }^{25}$ J. Lee, ${ }^{45}$ S. W. Lee, ${ }^{49}$ N. Leonardo, ${ }^{29}$ S. Leone, ${ }^{42}$ J. D. Lewis, ${ }^{13}$ K. Li, ${ }^{57}$ C. S. Lin, ${ }^{13}$ M. Lindgren, ${ }^{6}$ T. M. Liss, ${ }^{21}$ D. O. Litvintsev, ${ }^{13}$ T. Liu, ${ }^{13}$ Y. Liu, ${ }^{16}$ N. S. Lockyer, ${ }^{41}$ A. Loginov, ${ }^{33}$ J. Loken, ${ }^{39}$ M. Loreti, ${ }^{40}$ P. Loverre, ${ }^{47}$ D. Lucchesi, ${ }^{40}$ P. Lukens, ${ }^{13}$ L. Lyons, ${ }^{39}$ J. Lys, ${ }^{26}$ D. MacQueen, ${ }^{30}$ R. Madrak, ${ }^{18}$ K. Maeshima, ${ }^{13}$ P. Maksimovic, ${ }^{22}$ L. Malferrari, ${ }^{3}$ G. Manca, ${ }^{39}$ R. Marginean, ${ }^{36}$ A. Martin, ${ }^{57}$ M. Martin, ${ }^{22}$ V. Martin, ${ }^{35}$ M. Martinez, ${ }^{13}$ T. Maruyama, ${ }^{10}$ H. Matsunaga, ${ }^{52}$ M. Mattson, ${ }^{55}$ P. Mazzanti, ${ }^{3}$ K. S. McFarland, ${ }^{45}$ D. McGivern, ${ }^{28}$ P. M. McIntyre, ${ }^{49}$ P. McNamara, ${ }^{48}$ R. McNulty, ${ }^{27}$ S. Menzemer, ${ }^{29}$ A. Menzione, ${ }^{42}$ P. Merkel, ${ }^{13}$ C. Mesropian, ${ }^{46}$ A. Messina, ${ }^{47}$ A. Meyer, ${ }^{13}$ T. Miao, ${ }^{13}$ N. Miladinovic, ${ }^{4}$ L. Miller, ${ }^{18}$ R. Miller, ${ }^{32}$ J. S. Miller, ${ }^{31}$ R. Miquel, ${ }^{26}$ S. Miscetti, ${ }^{15}$ M. Mishina, ${ }^{13}$ G. Mitselmakher, ${ }^{14}$ A. Miyamoto, ${ }^{24}$ Y. Miyazaki, ${ }^{38}$ N. Moggi, ${ }^{3}$ R. Moore, ${ }^{13}$ M. Morello, ${ }^{42}$ T. Moulik, ${ }^{44}$ A. Mukherjee, ${ }^{13}$ M. Mulhearn, ${ }^{29}$ T. Muller, ${ }^{23}$ R. Mumford, ${ }^{22}$ A. Munar, ${ }^{41}$ P. Murat, ${ }^{13}$ J. Nachtman, ${ }^{13}$ S. Nahn, ${ }^{57}$ I. Nakamura, ${ }^{41}$ I. Nakano, ${ }^{37}$ A. Napier, ${ }^{53}$ R. Napora, ${ }^{22}$ V. Necula, ${ }^{14}$ F. Niell, ${ }^{31}$ J. Nielsen, ${ }^{26}$ C. Nelson, ${ }^{13}$ T. Nelson, ${ }^{13}$ C. Neu, ${ }^{36}$ M. S. Neubauer, ${ }^{29}$ C. Newman-Holmes, ${ }^{13}$ A-S. Nicollerat, ${ }^{16}$
T. Nigmanov,,${ }^{43}$ L. Nodulman, ${ }^{2}$ K. Oesterberg, ${ }^{19}$ T. Ogawa, ${ }^{54}$ S. Oh, ${ }^{12}$ Y. D. Oh,,$^{25}$ T. Ohsugi, ${ }^{20}$ R. Oishi, ${ }^{52}$ T. Okusawa, ${ }^{38}$ R. Oldeman, ${ }^{41}$ R. Orava, ${ }^{19}$ W. Orejudos, ${ }^{26}$ C. Pagliarone, ${ }^{42}$ F. Palmonari, ${ }^{42}$ R. Paoletti, ${ }^{42}$ V. Papadimitriou, ${ }^{50}$ S. Pashapour, ${ }^{30}$ J. Patrick, ${ }^{13}$ G. Pauletta,,${ }^{51}$ M. Paulini, ${ }^{9}$ T. Pauly, ${ }^{39}$ C. Paus, ${ }^{29}$ D. Pellett, ${ }^{5}$ A. Penzo, ${ }^{51}$ T. J. Phillips,,${ }^{12}$ G. Piacentino, ${ }^{42}$ J. Piedra, ${ }^{8}$ K. T. Pitts, ${ }^{21}$ A. Pompoš, ${ }^{44}$ L. Pondrom, ${ }^{56}$ G. Pope,,${ }^{43}$ O. Poukhov, ${ }^{11}$ F. Prakoshyn, ${ }^{11}$ T. Pratt, ${ }^{27}$ A. Pronko, ${ }^{14}$ J. Proudfoot, ${ }^{2}$ F. Ptohos, ${ }^{15}$ G. Punzi, ${ }^{42}$ J. Rademacker, ${ }^{39}$ A. Rakitine,,${ }^{29}$ S. Rappoccio, ${ }^{18}$ F. Ratnikov, ${ }^{48}$ H. Ray, ${ }^{31}$ A. Reichold, ${ }^{39}$ V. Rekovic, ${ }^{34}$ P. Renton, ${ }^{39}$ M. Rescigno, ${ }^{47}$ F. Rimondi, ${ }^{3}$ K. Rinnert, ${ }^{23}$ L. Ristori, ${ }^{42}$ W. J. Robertson, ${ }^{12}$ A. Robson, ${ }^{39}$ T. Rodrigo, ${ }^{8}$ S. Rolli, ${ }^{53}$ L. Rosenson, ${ }^{29}$ R. Roser, ${ }^{13}$ R. Rossin,,$^{40}$ C. Rott, ${ }^{44}$ J. Russ, ${ }^{9}$ A. Ruiz, ${ }^{8}$ D. Ryan, ${ }^{53}$ H. Saarikko, ${ }^{19}$ A. Safonov, ${ }^{5}$ R. St. Denis, ${ }^{17}$ W. K. Sakumoto, ${ }^{45}$ D. Saltzberg, ${ }^{6}$ C. Sanchez, ${ }^{36}$ A. Sansoni, ${ }^{15}$ L. Santi, ${ }^{51}$ S. Sarkar, ${ }^{47}$ K. Sato, ${ }^{52}$ P. Savard, ${ }^{30}$ A. Savoy-Navarro, ${ }^{13}$ P. Schemitz, ${ }^{23}$ P. Schlabach, ${ }^{13}$ E. E. Schmidt, ${ }^{13}$ M. P. Schmidt, ${ }^{57}$ M. Schmitt, ${ }^{35}$ L. Scodellaro, ${ }^{40}$ A. Scribano, ${ }^{42}$ F. Scuri, ${ }^{42}$ A. Sedov, ${ }^{44}$ S. Seidel, ${ }^{34}$ Y. Seiya, ${ }^{52}$ F. Semeria, ${ }^{3}$ L. Sexton-Kennedy, ${ }^{13}$ I. Sfiligoi, ${ }^{15}$ M. D. Shapiro, ${ }^{26}$ T. Shears, ${ }^{27}$ P. F. Shepard, ${ }^{43}$ M. Shimojima, ${ }^{52}$ M. Shochet, ${ }^{10}$ Y. Shon, ${ }^{56}$ A. Sidoti, ${ }^{42}$ M. Siket, ${ }^{1}$ A. Sill, ${ }^{50}$ P. Sinervo, ${ }^{30}$ A. Sisakyan, ${ }^{11}$ A. Skiba, ${ }^{23}$ A. J. Slaughter, ${ }^{13}$ K. Sliwa, ${ }^{53}$ J. R. Smith, ${ }^{5}$ F. D. Snider, ${ }^{13}$ R. Snihur, ${ }^{30}$ S.V. Somalwar, ${ }^{48}$ J. Spalding, ${ }^{13}$ M. Spezziga, ${ }^{50}$ L. Spiegel, ${ }^{13}$ F. Spinella, ${ }^{42}$ M. Spiropulu, ${ }^{7}$ P. Squillacioti, ${ }^{42}$ H. Stadie, ${ }^{23}$ B. Stelzer, ${ }^{30}$ O. Stelzer-Chilton, ${ }^{30}$ J. Strologas, ${ }^{34}$ D. Stuart, ${ }^{7}$ A. Sukhanov, ${ }^{14}$ K. Sumorok, ${ }^{29}$ H. Sun, ${ }^{53}$ T. Suzuki, ${ }^{52}$ A. Taffard, ${ }^{21}$ R. Tafirout, ${ }^{30}$ S. F. Takach,,${ }^{55}$ H. Takano, ${ }^{52}$ R. Takashima, ${ }^{20}$ Y. Takeuchi, ${ }^{52}$ K. Takikawa,,${ }^{52}$ P. Tamburello, ${ }^{12}$ M. Tanaka, ${ }^{2}$ R. Tanaka, ${ }^{37}$ N. Tanimoto, ${ }^{37}$ S. Tapprogge, ${ }^{19}$ M. Tecchio, ${ }^{31}$ P. K. Teng, ${ }^{1}$ K. Terashi, ${ }^{46}$ R. J. Tesarek, ${ }^{13}$ S. Tether, ${ }^{29}$ J. Thom, ${ }^{13}$ A. S. Thompson, ${ }^{17}$ E. Thomson, ${ }^{36}$ R. Thurman-Keup, ${ }^{2}$ P. Tipton, ${ }^{45}$ V. Tiwari, ${ }^{9}$ S. Tkaczyk, ${ }^{13}$ D. Toback, ${ }^{49} \mathrm{~K}$. Tollefson, ${ }^{32}$ D. Tonelli, ${ }^{42} \mathrm{M}$. Tönnesmann, ${ }^{32}$ S. Torre, ${ }^{42}$ D. Torretta, ${ }^{13} \mathrm{~W}$. Trischuk, ${ }^{30}$ J. Tseng,,${ }^{29}$ R. Tsuchiya, ${ }^{54}$ S. Tsuno, ${ }^{52}$ D. Tsybychev, ${ }^{14} \mathrm{~N}$. Turini, ${ }^{42}$ M. Turner, ${ }^{27}$ F. Ukegawa, ${ }^{52}$ T. Unverhau, ${ }^{17}$ S. Uozumi, ${ }^{52}$ D. Usynin, ${ }^{41}$ L. Vacavant, ${ }^{26}$ T. Vaiciulis, ${ }^{45}$ A. Varganov, ${ }^{31}$ E. Vataga, ${ }^{42}$ S. Vejcik III, ${ }^{13}$ G. Velev, ${ }^{13}$ G. Veramendi, ${ }^{26}$ T. Vickey, ${ }^{21}$ R. Vidal, ${ }^{13}$ I. Vila, ${ }^{8}$ R. Vilar, ${ }^{8}$ I. Volobouev, ${ }^{26}$ M. von der Mey, ${ }^{6}$ R. G. Wagner, ${ }^{2}$ R. L. Wagner, ${ }^{13}$ W. Wagner, ${ }^{23}$ N. Wallace, ${ }^{48}$ T. Walter, ${ }^{23}$ Z. Wan, ${ }^{48}$ M. J. Wang, ${ }^{1}$ S. M. Wang, ${ }^{14}$ A. Warburton, ${ }^{30}$ B. Ward, ${ }^{17}$ S. Waschke, ${ }^{17}$ D. Waters, ${ }^{28}$ T. Watts, ${ }^{48}$ M. Weber, ${ }^{26}$ W. Wester, ${ }^{13}$ B. Whitehouse, ${ }^{53}$ A. B. Wicklund, ${ }^{2}$ E. Wicklund, ${ }^{13}$ H. H. Williams, ${ }^{41}$ P. Wilson, ${ }^{13}$ B. L. Winer, ${ }^{36}$ P. Wittich, ${ }^{41}$ S. Wolbers, ${ }^{13}$ M. Wolter, ${ }^{53}$ M. Worcester, ${ }^{6}$ S. Worm, ${ }^{48}$ T. Wright, ${ }^{31}$ X. Wu, ${ }^{16}$ F. Würthwein, ${ }^{29}$ A. Wyatt, ${ }^{28}$ A. Yagil, ${ }^{13}$ T. Yamashita,,${ }^{37}$ K. Yamamoto, ${ }^{38}$ U. K. Yang, ${ }^{10}$ W. Yao, ${ }^{26}$ G. P. Yeh, ${ }^{13}$ K. Yi,,${ }^{22}$ J. Yoh, ${ }^{13}$ P. Yoon, ${ }^{45}$ K. Yorita, ${ }^{54}$ T. Yoshida, ${ }^{38}$ I. Yu, ${ }^{25}$ S. Yu, ${ }^{41}$ Z. Yu, ${ }^{57}$ J. C. Yun, ${ }^{13}$ L. Zanello, ${ }^{47}$ A. Zanetti, ${ }^{51}$ I. Zaw, ${ }^{18}$ F. Zetti, ${ }^{42}$ J. Zhou, ${ }^{48}$ A. Zsenei, ${ }^{16}$ and S. Zucchelli ${ }^{3}$
(CDF II Collaboration)

[^0]${ }^{26}$ Ernest Orlando Lawrence Berkeley National Laboratory, Berkeley, California 94720, USA<br>${ }^{27}$ University of Liverpool, Liverpool L69 7ZE, United Kingdom<br>${ }^{28}$ University College London, London WC1E 6BT, United Kingdom<br>${ }^{29}$ Massachusetts Institute of Technology, Cambridge, Massachusetts 02139, USA<br>${ }^{30}$ Institute of Particle Physics, McGill University, Montréal, Canada H3A 2T8,<br>and University of Toronto, Toronto, Canada M5S 1A7<br>${ }^{31}$ University of Michigan, Ann Arbor, Michigan 48109, USA<br>${ }^{32}$ Michigan State University, East Lansing, Michigan 48824, USA<br>${ }^{33}$ Institution for Theoretical and Experimental Physics, ITEP, Moscow 117259, Russia<br>${ }^{34}$ University of New Mexico, Albuquerque, New Mexico 87131, USA<br>${ }^{35}$ Northwestern University, Evanston, Illinois 60208, USA<br>${ }^{36}$ The Ohio State University, Columbus, Ohio 43210, USA<br>${ }^{37}$ Okayama University, Okayama 700-8530, Japan<br>${ }^{38}$ Osaka City University, Osaka 588, Japan<br>${ }^{39}$ University of Oxford, Oxford OX1 3RH, United Kingdom<br>${ }^{40}$ Universitá di Padova, Istituto Nazionale di Fisica Nucleare, Sezione di Padova-Trento, I-35131 Padova, Italy<br>${ }^{41}$ University of Pennsylvania, Philadelphia, Pennsylvania 19104, USA<br>${ }^{42}$ Istituto Nazionale di Fisica Nucleare, University and Scuola Normale Superiore of Pisa, I-56100 Pisa, Italy<br>${ }^{43}$ University of Pittsburgh, Pittsburgh, Pennsylvania 15260, USA<br>${ }^{44}$ Purdue University, West Lafayette, Indiana 47907, USA<br>${ }^{45}$ University of Rochester, Rochester, New York 14627, USA<br>${ }^{46}$ The Rockefeller University, New York, New York 10021, USA<br>${ }^{47}$ Instituto Nazionale de Fisica Nucleare, Sezione di Roma, University di Roma I, "La Sapienza," I-00185 Roma, Italy<br>${ }^{48}$ Rutgers University, Piscataway, New Jersey 08855, USA<br>${ }^{49}$ Texas A\&M University, College Station, Texas 77843, USA<br>${ }^{50}$ Texas Tech University, Lubbock, Texas 79409, USA<br>${ }^{51}$ Istituto Nazionale di Fisica Nucleare, Universities of Trieste and Udine, Italy<br>${ }^{52}$ University of Tsukuba, Tsukuba, Ibaraki 305, Japan<br>${ }^{53}$ Tufts University, Medford, Massachusetts 02155, USA<br>${ }^{54}$ Waseda University, Tokyo 169, Japan<br>${ }^{55}$ Wayne State University, Detroit, Michigan 48201, USA<br>${ }^{56}$ University of Wisconsin, Madison, Wisconsin 53706, USA<br>${ }^{57}$ Yale University, New Haven, Connecticut 06520, USA<br>(Received 4 December 2003; published 9 August 2004)

We report the observation of a narrow state decaying into $J / \psi \pi^{+} \pi^{-}$and produced in $220 \mathrm{pb}^{-1}$ of $\bar{p} p$ collisions at $\sqrt{s}=1.96 \mathrm{TeV}$ in the CDF II experiment. We observe $730 \pm 90$ decays. The mass is measured to be $3871.3 \pm 0.7$ (stat) $\pm 0.4$ (syst) $\mathrm{MeV} / c^{2}$, with an observed width consistent with the detector resolution. This is in agreement with the recent observation by the Belle Collaboration of the $X(3872)$ meson.

DOI: 10.1103/PhysRevLett. 93.072001
PACS numbers: $14.40 . \mathrm{Gx}, 12.39 . \mathrm{Mk}$

The study of bound states of charm-anticharm quarks revolutionized our understanding of hadrons beginning with the discovery of the $J / \psi$ meson in 1974 [1]. Although numerous charmonium ( $c \bar{c}$ ) states are now known, others should be observable. Recently, the Belle Collaboration reported a new particle $X(3872)$ observed in exclusive decays of $B$ mesons produced in $e^{+} e^{-}$collisions [2]. This particle has a mass of $3872 \mathrm{MeV} / c^{2}$ and decays into $J / \psi \pi^{+} \pi^{-}$. A natural interpretation of this particle would be a previously unobserved charmonium state, but there are no such states predicted to lie at or near the observed mass with the right quantum numbers to decay into $J / \psi \pi^{+} \pi^{-}[3,4]$. Within the framework of QCD, mesons may also arise from more complex systems than the conventional quark-antiquark bound state [5]. The proximity of the $X(3872)$ mass to the sum of the $D^{0}$ and $D^{* 0}$ masses suggests that $X(3872)$ may be a weakly bound deuteronlike "molecule" composed of a
$D$ and $\bar{D}^{*}$. Another possibility is that $X(3872)$ is a $c \bar{c} g$ hybrid meson-a $c \bar{c}$ system possessing a valence gluon. These novel interpretations have excited great interest in $X(3872)$ [6]. Whether it is a new form of hadronic matter or a conventional $c \bar{c}$ state in conflict with theoretical models, $X(3872)$ is an important object of study. Here we report the observation of a $J / \psi \pi^{+} \pi^{-}$resonance produced inclusively in $\bar{p} p$ collisions and which is consistent with $X(3872)$.

The analysis uses a data sample of $\bar{p} p$ collisions at $\sqrt{s}=1.96 \mathrm{TeV}$ with an integrated luminosity of $220 \mathrm{pb}^{-1}$ collected with the upgraded collider detector (CDF II) at the Fermilab Tevatron between February 2002 and August 2003. The important components of the CDF II detector for this analysis include a tracking system composed of a silicon-strip vertex detector (SVX II) [7] surrounded by an open-cell drift chamber system called the central outer tracker (COT) [8]. The

SVX II detector comprises five concentric layers of double-sided sensors located at radii between 2.5 and 10.6 cm . On one side of the sensors, axial strips measure positions in the plane transverse to the beam line. Strips on the other side are used for stereo measurements. The latter strips are tilted with respect to the axial strips: on one layer by $+1.2^{\circ}$, another by $-1.2^{\circ}$, and three by $90^{\circ}$. The active volume of the COT is a 3.1 m long cylinder covering radii from 43 to 132 cm with 8 superlayers of 12 wires each. In order to provide three-dimensional tracking, superlayers of axial wires alternate with superlayers of $+2^{\circ}$ stereo angle wires and superlayers of $-2^{\circ}$ stereo angle wires. The central tracking system is immersed in a 1.4 T solenoidal magnetic field for the measurement of charged particle momenta transverse to the beam line, $p_{T}$. The outermost detection system consists of planes of multilayer drift chambers for detecting muons [9]. The central muon system (CMU) covers $|\eta| \leq$ 0.6 , where $\eta \equiv-\ln [\tan (\theta / 2)]$ and $\theta$ is the angle of the particle with respect to the direction of the proton beam. Additional muon chambers (CMX) extend the rapidity coverage to $|\eta|=1.0$.

In this analysis, $J / \psi \rightarrow \mu^{+} \mu^{-}$decays are recorded using a dimuon trigger. The CDF II detector has a three-level trigger system. The level-1 trigger uses tracks in the muon chambers with a clear separation in azimuth from neighboring tracks. The extremely fast tracker (XFT) [10] uses information from the COT to select tracks based on $p_{T}$. XFT tracks with $p_{T} \geq 1.5 \mathrm{GeV} / c$ ( $p_{T} \geq 2.0 \mathrm{GeV} / c$ ) are extrapolated into the CMU (CMX) muon chambers and compared with the positions of muon tracks. If there are two or more XFT tracks with matches to muon tracks, the event passes the level-1 trigger. Dimuon triggers have no requirements at level 2. At level 3, the full tracking information from the COT is used to reconstruct a pair of opposite-sign muon candidates in the mass range from 2.7 to $4.0 \mathrm{GeV} / c^{2}$. Events passing the level-3 trigger are recorded for further analysis.

The offline analysis makes use of the best available calibrations of the tracking system for reconstructing events. Well-reconstructed tracks are selected by accepting only those with $\geq 3$ axial SVX II hits and $>20$ axial and $>16$ stereo COT hits. Tracks are refit to take into account the ionization energy loss appropriate for the particle hypotheses under consideration [11]. Dimuon candidates are selected in the mass range from 2.8 to $3.2 \mathrm{GeV} / c^{2}$ after being constrained to originate from a common point in a three-dimensional vertex fit. The resulting signal-to-background ratio for $J / \psi$ candidates is about 5 to 1 [12]. Pairs of charged tracks, both having $p_{T} \geq 0.35 \mathrm{GeV} / c$ and assumed to be pions, are then fit with the dimuon candidates to a common vertex. In this three-dimensional vertex fit, the dimuon mass is constrained to be the world average $J / \psi$ mass [13]. We require that the $\chi^{2}$ for the $J / \psi \pi \pi$ vertex fit must be less than 40 for 6 degrees of freedom.

The number of $J / \psi \pi \pi$ candidates per event passing the above preselection requirements can be quite large for events with a high multiplicity of charged tracks. These events contribute a large amount of combinatorial background relative to a small potential signal. We reject events that have more than 12 preselection candidates with masses below $4.5 \mathrm{GeV} / c^{2}$. A large number of candidates are accepted at this stage. However, after the final selection the average number of $J / \psi \pi^{+} \pi^{-}$candidates within the mass window of $3.65-4.0 \mathrm{GeV} / c^{2}$ is less than 1.2 per event for events with at least one accepted candidate. The specific number of preselection candidates allowed per event is determined by the optimization procedure described below.

In order to suppress $J / \psi \pi^{+} \pi^{-}$backgrounds, we tighten the selection criteria to $\chi^{2}<15$ for the 1 degree of freedom dimuon vertex fit, dimuon invariant mass within $60 \mathrm{MeV} / c^{2}(\sim 4$ standard deviations) of the world average $J / \psi$ mass, $p_{T}(J / \psi) \geq 4 \mathrm{GeV} / c, \chi^{2}<25$ for the $J / \psi \pi \pi$ vertex fit, $p_{T}(\pi) \geq 0.4 \mathrm{GeV} / c$, and $\Delta R \leq 0.7$ for both pions. Here $\Delta R$ is defined as $\sqrt{(\Delta \phi)^{2}+(\Delta \eta)^{2}}$, where $\Delta \phi$ is the difference in azimuthal angle between the pion and the $J / \psi \pi \pi$ candidate and $\Delta \eta$ is the difference in pseudorapidity.

The values used in the above selection criteria are determined by an iterative optimization procedure in which the significance $S / \sqrt{S+B}$ is maximized. The quantities $S$ and $B$ respectively represent the numbers of signal and background candidates obtained as a function of the values of the selection parameters. $B$ is available from background fits of the data in a window around $3872 \mathrm{MeV} / c^{2}$. We use $\psi(2 S) \rightarrow J / \psi \pi^{+} \pi^{-}$decays to model the $X(3872)$ yield $S$ as the selection is varied. The $\psi(2 S)$ signal is much larger than that of the $X(3872)$ and must therefore be scaled down for the significance calculation. The scale factor is determined such that $S$ matches the observed $X$ yield from a reference selection. Because the $X(3872)$ signal is considerably smaller than the background, the denominator of the significance ratio is dominated by $B$, and the optimization is not sensitive to the precise value of the scaling.

The $J / \psi \pi^{+} \pi^{-}$mass distribution of the selected candidates is displayed in Fig. 1. Besides the large peak showing the $\psi(2 S)$, a small peak is observed at a $J / \psi \pi^{+} \pi^{-}$mass around $3872 \mathrm{MeV} / c^{2}$. To fit the mass distribution, we model each peak by a single Gaussian and use a quadratic polynomial to describe the background. A binned maximum likelihood fit of the mass spectrum between 3.65 and $4.0 \mathrm{GeV} / c^{2}$ is also shown in Fig. 1. The fit yields signals of $5790 \pm 140 \psi(2 S)$ candidates and $580 \pm 100 X(3872)$ candidates.

The "wrong-sign" $J / \psi \pi^{ \pm} \pi^{ \pm}$mass distribution is also shown in Fig. 1, and no significant structures are apparent. We examine the hypothesis that the 3872 peak may originate from another state by incorrect assignment of the pion mass. The masses of $J / \psi \pi^{+} \pi^{-}$candidates in a window around the 3872 peak are recomputed for the


FIG. 1 (color online). The mass distributions of $J / \psi \pi^{+} \pi^{-}$ and $J / \psi \pi^{ \pm} \pi^{ \pm}$candidates passing the selection described in the text. A large peak for the $\psi(2 S)$ is seen in the $J / \psi \pi^{+} \pi^{-}$ distribution as well as a small signal near a mass of $3872 \mathrm{MeV} / c^{2}$. The curve is a fit using two Gaussians and a quadratic background to describe the data. The inset shows an enlargement of the $J / \psi \pi^{+} \pi^{-}$data and fit around $3872 \mathrm{MeV} / c^{2}$.
alternate hypotheses $J / \psi h_{1}^{+} h_{2}^{-}$, where $h_{1}^{+} h_{2}^{-}$are $\pi^{+} K^{-}$, $K^{+} K^{-}, p \pi^{-}, p K^{-}$, and $p \bar{p}$ (and charge conjugates). This results in broad mass distributions with no peaklike structures. Thus, the 3872 peak is not an artifact of some other state, known or unknown, decaying into a $J / \psi$ and a pair of hadrons in which one or both hadrons are misassigned as pions.

The $X(3872)$ signal reported by the Belle Collaboration favors large $\pi^{+} \pi^{-}$masses. Our data support this conclusion as well. We divide the data into two subsamples: candidates with dipion masses greater, or less, than $0.5 \mathrm{GeV} / c^{2}$. From the Belle results, this is a large enough value to probe the high-mass behavior of the $X(3872)$ candidates and yet not eliminate all the $\psi(2 S)$ reference signal from the high-mass subsample. Figure 2 shows the resulting $J / \psi \pi^{+} \pi^{-}$mass distributions. The prominence of the $X(3872)$ peak is enhanced over the background in the high-mass sample, and no peak is apparent for low masses. Fitting the high-mass spectrum between 3.65 and $4.0 \mathrm{GeV} / c^{2}$ gives $3530 \pm 100 \psi(2 S)$ candidates and $730 \pm 90 X(3872)$ candidates. The fitted mass and width of the $\psi(2 S)$ are $3685.65 \pm 0.09$ (stat) $\mathrm{MeV} / c^{2}$ and $3.44 \pm$ 0.09 (stat) $\mathrm{MeV} / c^{2}$, respectively. For $X(3872)$ we obtain a mass of $3871.3 \pm 0.7$ (stat) $\mathrm{MeV} / c^{2}$ and a width of $4.9 \pm$ $0.7 \mathrm{MeV} / c^{2}$. The latter value is consistent with detector resolution. Our mass is in good agreement with the Belle result of $3872.0 \pm 0.6$ (stat) $\pm 0.5$ (syst) $\mathrm{MeV} / c^{2}$ [2].


FIG. 2 (color online). The mass distributions of $J / \psi \pi^{+} \pi^{-}$ candidates with $m\left(\pi^{+} \pi^{-}\right)>0.5 \mathrm{GeV} / c^{2} \quad$ (points) and $m\left(\pi^{+} \pi^{-}\right)<0.5 \mathrm{GeV} / c^{2}$ (open circles). The curve is a fit with two Gaussians and a quadratic background. The inset shows an enlargement of the high dipion-mass data and fit.

Requiring $M\left(\pi^{+} \pi^{-}\right)>0.5 \mathrm{MeV} / c^{2}$ reduces the background by almost a factor of 2 and apparently increases the amount of fitted $X(3872)$ signal. A significant part of the additional signal is attributable to an increase in the fitted width. The original fit over all dipion masses returns a smaller but consistent width of $4.2 \pm 0.8 \mathrm{MeV} / c^{2}$. We conclude that the $X(3872)$ signal yield after the dipion requirement is unchanged within statistics, and thus there is little signal with dipion masses below $0.5 \mathrm{GeV} / c^{2}$. The same conclusion is reached by direct examination of the low dipion-mass distribution shown in Fig. 2. We use the high-mass sample for measuring the $X(3872)$ mass as the improved signal-to-noise ratio reduces the statistical uncertainty.

The fit displayed in Fig. 2 has a $\chi^{2}$ of 74.9 for 61 degrees of freedom, which corresponds to a probability of $10.9 \%$. To estimate the significance of the signal, we first count the number of candidates in the three bins centered on the peak, i.e., 3893. The three-bin background is estimated from the fit to be 3234 candidates, leaving a signal of 659 candidates. In a Gaussian approach, this corresponds to a significance of $659 / \sqrt{3234}=11.6$ standard deviations. The Poisson probability for 3234 to fluctuate up to or above 3893 is in good agreement with the Gaussian estimate, considering the approximations of each method.

The systematic uncertainty on the mass scale is related to the momentum scale calibration, the various tracking systematics, and the vertex fitting. These effects
were studied in detail for our measurement of the mass difference $m\left(D_{s}^{+}\right)-m\left(D^{+}\right)$[11], where the systematic uncertainty was $\pm 0.21 \mathrm{MeV} / c^{2}$. A larger systematic uncertainty arises for our $X(3872)$ mass determination because it is an absolute measurement. We use the $\psi(2 S)$ mass to gauge our systematic uncertainty. With the dipion mass requirement, the $\psi(2 S)$ mass is measured to be $0.3 \mathrm{MeV} / c^{2}$ below the world average mass of $3685.96 \pm$ 0.09 [13], a difference substantially larger than the statistical uncertainty of $0.1 \mathrm{MeV} / c^{2}$. However, studies of the stability of the $\psi(2 S)$ mass for different selection requirements indicate that an uncertainty of $0.4 \mathrm{MeV} / c^{2}$ should be assigned. Variations of the fit model and fit range have negligible effect on the mass.

In summary, we report the observation of a state consistent with $X(3872)$ decaying into $J / \psi \pi^{+} \pi^{-}$. From a sample of $730 \pm 90$ candidates we measure the $X(3872)$ mass to be $3871.3 \pm 0.7$ (stat) $\pm 0.4$ (syst) $\mathrm{MeV} / \mathrm{c}^{2}$ and find that the observed width is consistent with the detector resolution. This is in agreement with the measurement by the Belle Collaboration using $B^{ \pm}$decays [2]. The average mass from the two experiments, assuming uncorrelated systematic uncertainties, is $3871.7 \pm$ $0.6 \mathrm{MeV} / c^{2}$. Our large sample of this new particle opens up avenues for future investigations, such as production mechanisms, the dipion mass distribution, and spinparity analysis.

We thank the Fermilab staff and the technical staffs of the participating institutions for their vital contributions. This work was supported by the U.S. Department of Energy and National Science Foundation; the Italian Istituto Nazionale di Fisica Nucleare; the Ministry of Education, Culture, Sports, Science and Technology of Japan; the Natural Sciences and Engineering Research Council of Canada; the National Science Council of the Republic of China; the Swiss National Science Foundation; the A.P. Sloan Foundation; the Bundesministerium für Bildung und Forschung, Germany; the Korean Science and Engineering Foundation and the Korean Research Foundation; the Particle Physics and Astronomy Research Council and
the Royal Society, United Kingdom; the Russian Foundation for Basic Research; the Comisión Interministerial de Ciencia y Tecnología, Spain; the European Community's Human Potential Programme under Contract No. HPRN-CT-20002, Probe for New Physics; and the Research Fund of Istanbul University Project No. 1755/ 21122001.
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[^0]:    ${ }^{1}$ Institute of Physics, Academia Sinica, Taipei, Taiwan 11529, Republic of China
    ${ }^{2}$ Argonne National Laboratory, Argonne, Illinois 60439, USA
    ${ }^{3}$ Istituto Nazionale di Fisica Nucleare, University of Bologna, I-40127 Bologna, Italy
    ${ }^{4}$ Brandeis University, Waltham, Massachusetts 02254, USA
    ${ }^{5}$ University of California at Davis, Davis, California 95616, USA
    ${ }^{6}$ University of California at Los Angeles, Los Angeles, California 90024, USA
    ${ }^{7}$ University of California at Santa Barbara, Santa Barbara, California 93106, USA
    ${ }^{8}$ Instituto de Fisica de Cantabria, CSIC-University of Cantabria, 39005 Santander, Spain
    ${ }^{9}$ Carnegie Mellon University, Pittsburgh, Pennsylvania 15213, USA
    ${ }^{10}$ Enrico Fermi Institute, University of Chicago, Chicago, Illinois 60637, USA
    ${ }^{11}$ Joint Institute for Nuclear Research, RU-141980 Dubna, Russia
    ${ }^{12}$ Duke University, Durham, North Carolina 27708, USA
    ${ }^{13}$ Fermi National Accelerator Laboratory, Batavia, Illinois 60510, USA
    ${ }^{14}$ University of Florida, Gainesville, Florida 32611, USA
    ${ }^{15}$ Laboratori Nazionali di Frascati, Istituto Nazionale di Fisica Nucleare, I-00044 Frascati, Italy
    ${ }^{16}$ University of Geneva, CH-1211 Geneva 4, Switzerland
    ${ }^{17}$ Glasgow University, Glasgow G12 8QQ, United Kingdom
    ${ }^{18}$ Harvard University, Cambridge, Massachusetts 02138, USA
    ${ }^{19}$ The Helsinki Group, Helsinki Institute of Physics, and Division of High Energy Physics, Department of Physical Sciences, University of Helsinki, FIN-00014 Helsinki, Finland
    ${ }^{20}$ Hiroshima University, Higashi-Hiroshima 724, Japan
    ${ }^{21}$ University of Illinois, Urbana, Illinois 61801, USA
    ${ }^{22}$ The Johns Hopkins University, Baltimore, Maryland 21218, USA
    ${ }^{23}$ Institut für Experimentelle Kernphysik, Universitüt Karlsruhe, 76128 Karlsruhe, Germany
    ${ }^{24}$ High Energy Accelerator Research Organization (KEK), Tsukuba, Ibaraki 305, Japan
    ${ }^{25}$ Center for High Energy Physics, Kyungpook National University, Taegu 702-701, Seoul National University, Seoul 151-742, and SungKyunKwan University, Suwon 440-746, Korea

