

Missouri University of Science and Technology [Scholars' Mine](https://scholarsmine.mst.edu/) 

[Chemistry Faculty Research & Creative Works](https://scholarsmine.mst.edu/chem_facwork) [Chemistry](https://scholarsmine.mst.edu/chem) Chemistry

01 Nov 2000

# Cerium LIII-Edge X-ray Absorption Study of the Ce<sub>xFe4-y</sub>Co<sub>y</sub>Sb<sub>12</sub> **Skutterudites**

Fernande Grandjean Missouri University of Science and Technology, grandjeanf@mst.edu

Gary J. Long Missouri University of Science and Technology, glong@mst.edu

Robert Cortès

Donald T. Morelli

et. al. For a complete list of authors, see [https://scholarsmine.mst.edu/chem\\_facwork/874](https://scholarsmine.mst.edu/chem_facwork/874) 

Follow this and additional works at: [https://scholarsmine.mst.edu/chem\\_facwork](https://scholarsmine.mst.edu/chem_facwork?utm_source=scholarsmine.mst.edu%2Fchem_facwork%2F874&utm_medium=PDF&utm_campaign=PDFCoverPages)

**C** Part of the Chemistry Commons

# Recommended Citation

F. Grandjean et al., "Cerium LIII-Edge X-ray Absorption Study of the Ce<sub>xFe4-y</sub>Co<sub>y</sub>Sb<sub>12</sub> Skutterudites," Physical Review B (Condensed Matter), vol. 62, no. 19, pp. 12569-12572, American Physical Society (APS), Nov 2000. The definitive version is available at <https://doi.org/10.1103/PhysRevB.62.12569>

This Article - Journal is brought to you for free and open access by Scholars' Mine. It has been accepted for inclusion in Chemistry Faculty Research & Creative Works by an authorized administrator of Scholars' Mine. This work is protected by U. S. Copyright Law. Unauthorized use including reproduction for redistribution requires the permission of the copyright holder. For more information, please contact [scholarsmine@mst.edu](mailto:scholarsmine@mst.edu).

## **Cerium**  $L_{\text{III}}$ -edge x-ray absorption study of the  $Ce_xFe_{4-y}Co_ySb_{12}$  skutterudites

Fernande Grandjean

*Institute of Physics, B5, University of Lie`ge, B-4000 Sart-Tilman, Belgium*

Gary J. Long

*Department of Chemistry, University of Missouri-Rolla, Rolla, Missouri 65409-0010*

Robert Cortes

*Laboratoire pour l'Utilisation du Rayonnement Electromagne´tique, Universite´ Paris-Sud, F-91405 Orsay, France*

Donald T. Morelli

*Materials Group, Delphi Automotive Systems Central R&D, Warren, Michigan 48090*

Gregory P. Meisner

*Materials and Processes Laboratory, Mail Code 480-106-224, General Motors Global Research and Development Operations, Warren, Michigan 48090*

(Received 7 July 2000)

The cerium *L*<sub>III</sub>-edge x-ray absorption near-edge spectra of the  $Ce_xFe_{4-y}Co_ySb_{12}$  compounds have been obtained at 295 K and unambiguously indicate that cerium is in the  $4f<sup>1</sup>$  electronic ground state for all values of  $0.22 \le x \le 0.98$  and  $0.0 \le y \le 3.5$ . This stable trivalent state of cerium is in agreement with the proposed  $(CeFe_4Sb_{12})_{1-\alpha}$  ( $\Box Co_4Sb_{12}$ )<sub>a</sub>, solid solution structure, in which the cerium atoms are always surrounded by twelve antimony first neighbors and six iron second neighbors, the observed magnetic properties of  $Cefe_4Sb_{12}$ and  $Ce_{0.9}Fe_3CoSb_{12}$ , and the electronic structure of  $CeFe_4Sb_{12}$  obtained from band-structure calculations.

#### **I. INTRODUCTION**

The valence of cerium in the filled skutterudite  $CeFe<sub>4</sub>X<sub>12</sub>$ compounds, where *X* is P, As, or Sb, has been the subject of many discussions and much recent speculation. A cerium valence of four in  $CeFe_4P_{12}$  and of three in  $CeFe_4Sb_{12}$  has been postulated<sup>1</sup> from their unit-cell volumes. A trivalent state for cerium in CeFe<sub>4</sub>Sb<sub>12</sub> was also predicted<sup>2</sup> from bandstructure calculations. In contrast, a cerium valence of 3.74 at 4.2 K was deduced<sup>3</sup> from electrical conductivity measurements. The concept of the cerium valence in these compounds requires some explanation.<sup>2</sup> In no case is the cerium atom oxidized to the cerium $(III)$  or cerium $(IV)$  oxidation state in these intermetallic compounds. Rather, the concept of the cerium valence refers to the electronic ground state of the cerium atom; the  $4f^0$  and  $4f^1$  ground states are usually referred to as tetravalent and trivalent cerium, respectively.

In order to better understand and predict the technologically important thermoelectric properties<sup>4</sup> of the filled skutterudite compounds, it is necessary to obtain a good description of their electronic structure and, particularly, of the electronic configuration of cerium in these materials. X-ray absorption spectroscopy at the cerium *L*III edge is a unique technique<sup>5</sup> for obtaining this information, because it is element and electronic transition specific. Recently, the cerium *L*III-edge x-ray absorption near-edge structure spectrum (XANES spectrum) of CeRu<sub>4</sub>P<sub>12</sub> at room temperature has been reported,<sup>6</sup> and the results indicate that cerium is predominately trivalent, i.e., cerium is in the  $4f<sup>1</sup>$  electronic ground state in this filled skutterudite. We present herein a cerium  $L_{III}$ -edge XANES study of the  $Ce_{x}Fe_{4-y}Co_{y}Sb_{12}$  series of compounds and we discuss the results in the framework of their electronic structure.

## **II. EXPERIMENT**

The samples studied herein are identical to those prepared<sup>7</sup> and studied<sup>8,9</sup> previously. The powder absorbers had thicknesses of  $\sim$ 15 mg/cm<sup>2</sup>. The cerium  $L_{\text{III}}$ -edge XANES spectra were obtained at 295 K on the D21 station at LURE. The ring was operating at 1.85 GeV and a stored current of 200 mA. Detailed preliminary measurements indicated that better XANES spectra were obtained in transmission mode than in fluorescence detection mode. Thus the measurements were carried out in transmission mode with a  $Si(311)$  double-crystal monochromator using ion chambers filled with a mixture of nitrogen and helium, both in front and behind the sample. Harmonic contamination was rejected by using a total reflection mirror with a cut off at 9 keV. The resolution calculated from the width of the two slits located in front of the monochromator and in front of the first ion chamber is estimated to be 1.4 eV. The spectra were recorded with a step of 2 eV below 5710 eV and of 0.2 eV between 5710 and 5760 eV. Because of the extensive scattering produced by the presence of large amounts of antimony in the absorbers, the jump in absorption at the cerium  $L_{\text{III}}$  edge is smaller than desirable and the quality of the XANES spectra is somewhat lower than usually spectra is somewhat lower than usually obtained $6,10,11$  with other cerium-containing intermetallic compounds. The *K*-edge absorption spectrum of a thin metallic foil of chromium was used for energy calibration.



FIG. 1. The 295 K cerium *I*<sub>III</sub>-edge XANES spectra of  $Ce<sub>x</sub>Fe<sub>4-y</sub>Co<sub>y</sub>Sb<sub>12</sub>$ , with  $x=0.93$ , 0.71, 0.47, 0.22 and their fits as described in the text.

#### **III. RESULTS AND DISCUSSION**

The cerium *L*<sub>III</sub>-edge XANES spectra of  $Ce_xFe_{4-y}Co_ySb_{12}$  are shown in Fig. 1 for  $x=0.93, 0.71$ , 0.47, and 0.22 and the *y* values given in Table I. The spectra obtained for  $x=0.98, 0.82, 0.60,$  and 0.35, are not shown herein but are very similar. All the spectra have been analyzed in the fashion discussed below. It is obvious that all the spectra are characterized by a main absorption peak at  $\sim$ 5725 eV and a minor broader absorption peak at  $\sim$ 5735 eV. Further, these spectra are very similar to that observed<sup>o</sup> for  $CeRu_4P_{12}$ . Hence, a simple visual inspection indicates a

TABLE I. Cerium *L*<sub>III</sub>-edge XANES energies, and core-hole lifetimes.  $\Gamma$  in eV, and the cerium valence for  $Ce_{x}Fe_{4-y}Co_{y}Sb_{12}$ .

| $\boldsymbol{x}$ | y   | $E_0$  | $E_{1}$ | Δ    | Г    | Ce valence |
|------------------|-----|--------|---------|------|------|------------|
| 0.22             | 3.5 | 5721.7 | 5725.5  | 2.58 | 3.40 | 3.00       |
| 0.35             | 3.0 | 5721.4 | 5725.3  | 2.76 | 3.62 | 3.02       |
| 0.47             | 2.5 | 5721.7 | 5725.3  | 2.70 | 3.74 | 3.02       |
| 0.60             | 2.0 | 5721.8 | 5725.2  | 2.80 | 3.84 | 3.04       |
| 0.71             | 1.5 | 5722.2 | 5725.0  | 2.71 | 4.06 | 3.04       |
| 0.82             | 1.0 | 5722.3 | 5725.0  | 2.79 | 4.12 | 3.06       |
| 0.93             | 0.5 | 5722.1 | 5725.1  | 2.86 | 4.00 | 3.07       |
| 0.98             | 0.0 | 5722.1 | 5725.4  | 2.62 | 3.94 | 3.03       |

predominately trivalent state, i.e., a  $4f<sup>1</sup>$  electronic ground state for cerium in  $Ce<sub>v</sub>Fe<sub>4-v</sub>Co<sub>v</sub>Sb<sub>12</sub>$ .

A more detailed analysis of the spectra shown in Fig. 1 and a search for any trend with composition requires a more sophisticated fit of the absorption profile and several different models for fitting the cerium *L*<sub>III</sub>-edge XANES spectra have been reported.<sup>6,10–12</sup> The common points between these models are the following. The continuum absorption is represented by an arctangent function superimposed on a linear background and the absorption lines corresponding to the  $2p_{3/2}$   $\rightarrow$  5*d* transitions are represented by Lorentzian lines. Hence, the cerium  $L_{\text{III}}$ -edge XANES spectrum of CeAl<sub>2</sub>, a typical trivalent cerium compound, is well fitted<sup>11</sup> by equation.

$$
\mu(E) = B_0 + B_1 E + \left\{ 0.5 + \frac{1}{\pi} \arctan\left(\frac{E - (E_1 + \delta)}{\Gamma/2}\right) \right\} + \frac{A(\Gamma/2)^2}{(E - E_1)^2 + (\Gamma/2)^2},
$$
\n(1)

where  $B_0$  and  $B_1$  are the linear background parameters,  $E_1$  is the energy at which the  $2p_{3/2}$   $\rightarrow$  5*d* transition occurs.  $\Gamma$  is the core-hole lifetime for the transition. *A* is the area of the Lorentzian peak, and  $\delta$  is the shift between the onset of the continuum and bound-state transitions. For the spectrum of CeAl<sub>2</sub>,  $\delta$  is positive.

The cerium *L*<sub>III</sub>-edge XANES spectra of intermediatevalent cerium compounds show $10,12-15$  two Lorentzian absorption peaks separated by the replicate splitting of 9 to 10 eV, and they have been analyzed<sup>12,14</sup> with two arctangent functions and two Lorentzian peaks representing the continuum absorption and the  $2p_{3/2}$   $\rightarrow$  5*d* transitions for the trivalent and tetravalent cerium states, respectively. In the case of the CeFe<sub>11</sub>TiH<sub>y</sub> compounds,<sup>12</sup> it should be noted that the fit shown is poor and the parameter  $\delta$  is negative for both valence states, in disagreement with the analysis of Croft *et al.*<sup>11</sup> for CeAl<sub>2</sub>. In this model the cerium XANES spectra have been fitted with the equation

$$
\mu(E) = B_0 + B_1 E + \frac{A_1}{A_1 + A_2} \left\{ 0.5 + \frac{1}{\pi} \arctan\left(\frac{E - (E_1 + \delta)}{\Gamma/2}\right) \right\} \n+ \frac{A_2}{A_1 + A_2} \left\{ 0.5 + \frac{1}{\pi} \arctan\left(\frac{E - (E_2 + \delta)}{\Gamma/2}\right) \right\} \n+ \frac{(\Gamma/2)^2 A_1}{(E - E_1)^2 + (\Gamma/2)^2} + \frac{(\Gamma/2)^2 A_2}{(E - E_2)^2 + (\Gamma/2)^2},
$$
\n(2)

where  $E_1$  and  $E_2$  are the energies at which the  $2p_{3/2} \rightarrow 5d$ transitions occur for the trivalent and the tetravalent cerium, respectively.  $\Gamma$  is the core hole lifetime for the transition.  $A_1$ and  $A_2$  are the areas of the Lorentzian peaks, and  $\delta$  is the energy shift between the onset of the continuum and the bound state. All our attempts to fit the spectra in Fig. 1 with Eq.  $(2)$  lead to a poor fit of the main peak and we conclude, in agreement with Lee *et al.*,<sup>6</sup> that the peak at  $\sim$ 5725 eV is not a simple Lorentzian peak.

Hence, after subtraction of a linear background, the cerium  $L_{III}$ -edge XANES spectra of  $Ce_xFe_{4-y}Co_ySb_{12}$  have been fitted with a model that uses two arctangent functions to fit the continuum and three Lorentzian peaks, two for the main peak at  $\sim$ 5725 eV and one for the minor peak at  $\sim$  5735 eV, i.e., with the equation

$$
\mu(E) = (A_1 + A_2)/(A_1 + A_2 + A_3) \left\{ 0.5 + \frac{1}{\pi} \arctan\left(\frac{E - E_0}{\Gamma/2}\right) \right\}
$$

$$
+ \frac{A_1(\Gamma/2)^2}{(E - E_1)^2 + (\Gamma/2)^2} + \frac{A_2(\Gamma/2)^2}{(E - E_1 - \Delta)^2 + (\Gamma/2)^2}
$$

$$
+ A_3/(A_1 + A_2 + A_3) \left\{ 0.5 + \frac{1}{\pi} \arctan\left(\frac{E - E_0 - 10}{\Gamma/2}\right) \right\}
$$

$$
+ \frac{A_3(\Gamma/2)^2}{(E - E_1 - 10)^2 + (\Gamma/2)^2}, \tag{3}
$$

where  $E_0$  is the energy of the onset of the continuum.  $E_1$  and  $E_1 + \Delta$  are the energies at which the  $2p_{3/2} \rightarrow 5d$  transitions occur for the trivalent cerium.  $\Gamma$  is the core-hole lifetime for the transition, and  $A_1$ ,  $A_2$ , and  $A_3$  are the areas of the Lorentzian peaks. Because the peak at  $\sim$ 5235 eV is very weak, a replicate splitting of 10 eV has been assumed in Eq.  $(3).$ 

The fits with Eq.  $(3)$  are good and the resulting parameters are given in Table I. The predominant peak at  $\sim$ 5725 eV is well fitted with two Lorentzian peaks at  $\sim$  5725 and 5728 eV. This double Lorentzian absorption is assigned to the  $2p_{3/2}4f^{1}\rightarrow 4f^{1}5d$  transitions in which the 5*d* orbitals are split into  $e_g$  and  $t_{2g}$  orbitals by the crystal field of the twelve antimony atoms surrounding the cerium atom. A splitting,  $\Delta$ , of  $\sim$ 2.6 eV corresponds to a 10*Dq* of 21 000 cm<sup>-1</sup>, a very reasonable value for this type of electronic environment. The Lorentzian peak at  $\sim$  5735 eV is assigned to the 2*p*<sub>3/2</sub>4 $f$ <sup>0</sup>  $\rightarrow$  4 $f^0$ 5*d* transition. As indicated in Table I, the linewidth  $\Gamma$ or the core-hole lifetime is essentially independent of composition and is similar to that measured by Lee *et al.*<sup>6</sup> in  $CeRu<sub>4</sub>P<sub>12</sub>$ .

The cerium valence obtained from the areas  $A_1$ ,  $A_2$ , and *A*<sup>3</sup> is also given in Table I and is shown in Fig. 2 as a function of the cerium content. The samples for  $x=0.93$  and 0.98 have been measured at two different times on the same absorber and the two values are both plotted in Fig. 2. These duplicate measurements give an estimate of  $\pm 0.03$  for the error on the cerium valence. Hence, we conclude that the cerium valence is  $3.03 \pm 0.03$  and cerium is in the  $4 f<sup>1</sup>$  ground state in all the  $Ce_xFe_{4-y}Co_ySb_{12}$  compounds. The slight increase in cerium valence with cerium content is may be due to a trivial amount of oxidation of surface cerium atoms in the grains.

## **IV. CONCLUSIONS**

Our experimental measurements indicate that cerium is trivalent in all the  $Ce_xFe_{4-y}Co_ySb_{12}$  compounds and is independent of both *x* and *y*, an independence which is in agreement with theoretical predictions from band-structure calculations of Nordström and Singh<sup>2</sup> on CeFe<sub>4</sub>Sb<sub>12</sub>. Unfortunately, the band-structure calculations extend at most





FIG. 2. The cerium valence in  $Ce_xFe_{4-y}Co_ySb_{12}$  as a function of cerium content. It should be noted that the  $\pm 0.03$  error bars are almost the size of the vertical axis of the plot and that no true correlation between the cerium valence and the cerium content exists.

to only 4 eV above the Fermi level, and hence the energies of the unoccupied 5*d* levels probed by XANES cannot be compared with the band-structure calculations. The independence of the cerium valence on composition agrees with the proposal<sup>8</sup> that the  $Ce_xFe_{4-y}Co_ySb_{12}$  compounds consist of solid solutions  $(CeFe_4Sb_{12})_{1-\alpha}(\square Co_4Sb_{12})_{\alpha}$ , in which the cerium atoms are surrounded by twelve antimony first neighbors and six iron second neighbors. The effective magnetic moments<sup>16</sup> of  $3.8\mu_B$  and  $2.6\mu_B$  in CeFe<sub>4</sub>Sb<sub>12</sub> and  $Ce<sub>0.9</sub>Fe<sub>3</sub>CoSb<sub>12</sub>$  are also in agreement with the trivalent state of cerium whose free ion magnetic moment is  $2.54\mu_B$ .

The stable trivalent state of cerium in  $Ce<sub>x</sub>Fe<sub>4-y</sub>Co<sub>y</sub>Sb<sub>12</sub>$ contrasts with the intermediate-valent state<sup>17</sup> of ytterbium in  $YbFe<sub>4</sub>Sb<sub>12</sub>$ . Steric arguments may be used to understand this difference. The unit cell volume of  $YbFe<sub>4</sub>Sb<sub>12</sub>$  is larger than that of  $CeFe<sub>4</sub>Sb<sub>12</sub>$  and hence the void in which the ytterbium atom resides is larger than that available for the cerium. In addition, because of the lanthanide contraction ytterbium, with a 12-coordinate metallic radius of 1.75 Å, is smaller than cerium, with a 12-coordinate metallic radius of 1.81 Å. Hence, in order to better fill in the available void in the structure, ytterbium tends to be divalent.

## **ACKNOWLEDGMENTS**

This research was supported in part by the U.S. National Science Foundation through Grants No. DMR95-21739 and No. INT-9815138 and by the European Commission for Human Capital and Mobility, Contract No. ERBHXCT930360.

- <sup>1</sup>W. Jeitschko and D. Braun, Acta Crystallogr., Sect. B: Struct. Crystallogr. Cryst. Chem. 33, 3401 (1977).
- $2$ L. Nordström and D. J. Singh, Phys. Rev. B **53**, 1103 (1996).
- $3$  D. T. Morelli and G. P. Meisner, J. Appl. Phys. **77**, 3777 (1995).
- 4B. C. Sales, D. Mandrus, and R. K. Williams, Science **272**, 1325  $(1996).$
- 5A. Bianconi, in *X-ray Absorption: Principles, Applications, Techniques of EXAFS, SEXAFS, and XANES*, edited by D. C. Koningsberger and R. Prins (Wiley, New York, 1988), Chap. 11.
- 6C. H. Lee, H. Oyanagi, C. Sekine, I. Shirotani, and M. Ishii, Phys. Rev. B 60, 13 253 (1999).
- $<sup>7</sup>D$ . T. Morelli, G. P. Meisner, B. Chen, S. Hu, and C. Uher, Phys.</sup> Rev. B 56, 7376 (1997).
- 8G. P. Meisner, D. T. Morelli, S. Hu, J. Yang, and C. Uher, Phys. Rev. Lett. **80**, 3551 (1998).
- <sup>9</sup>G. J. Long, D. Hautot, F. Grandjean, D. T. Morelli, and G. P. Meisner, Phys. Rev. B 60, 7410 (1999).
- 10D. Vandormael, F. Grandjean, V. Briois, D. P. Middleton, K. H.
- J. Buschow, and G. J. Long, Phys. Rev. B 56, 6100 (1997).
- 11M. Croft, R. Neifeld, C. U. Segre, S. Raaen, and R. D. Parks, Phys. Rev. B 30, 4164 (1984).
- <sup>12</sup> J. Chaboy, A. Marcelli, L. Bozukov, F. Baudelet, E. Dartyge, A. Fontaine, and S. Pizzini, Phys. Rev. B 51, 9005 (1995).
- 13T. W. Capehart, R. K. Mishra, G. P. Meisner, C. D. Fuerst, and J. F. Herbst, Appl. Phys. Lett. **63**, 3642 (1993).
- <sup>14</sup> J. Chaboy, J. Garcia, and A. Marcelli, Appl. Phys. Lett. **65**, 3149  $(1994).$
- 15T. W. Capehart, R. K. Mishra, G. P. Meisner, C. D. Fuerst, and J. F. Herbst, Appl. Phys. Lett. **65**, 3151 (1994).
- 16D. A. Gajewski, N. R. Dilley, E. D. Bauer, E. J. Freeman, R. Chau, M. B. Maple, D. Mandrus, B. C. Sales, and A. H. Lacerda, J. Phys: Condens. Matter 10, 6973-6985 (1998).
- 17A. Leithe-Jasper, D. Kaczorowski, P. Rogl, J. Bogner, M. Reissner, W. Steiner, G. Wiesinger, and C. Godart, Solid State Commun. 109, 395 (1999).