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PRIMARY SPACES, MACKEY'S OBSTRUCTION, AND THE GENERALIZED BARYCENTRIC DECOMPOSITION

PATRICK IGLESIAS-ZEMMOUR AND FRANÇOIS ZIEGLER

We call a hamiltonian N-space primary if its moment map is onto a single coadjoint orbit. The question has long been open whether such spaces always split as (homogeneous) \times (trivial), as an analogy with representation theory might suggest. For instance, Souriau's $barycentric\ decomposition\ theorem$ asserts just this when N is a Heisenberg group. For general N, we give explicit examples which do not split, and show instead that primary spaces are always flat bundles over the coadjoint orbit. This provides the missing piece for a full "Mackey theory" of hamiltonian G-spaces, where G is an overgroup in which N is normal.

Il est bien entendu que chaque moment, confondu qu'il est en tous les autres, demeure pourtant en lui-même différencié.

—André Breton, Les vases communicants

Introduction

Let G be a finite group and let N be a normal subgroup, so that we have an exact sequence

$$1 \longrightarrow N \longrightarrow G \longrightarrow G/N \longrightarrow 1.$$

Questions about G then often reduce to similar ones about N and G/N. A classic example is the calculation of the dual $Irr(G) = \{irreducible unitary G-modules\}/isomorphism.$ This, as Clifford [C37] showed, boils down in 3 steps to finding Irr(N) and the projective duals of subgroups of G/N:

- (1) When restricted to N, every $X \in Irr(G)$ decomposes into the irreducibles belonging to a single G-orbit $G(U) \subset Irr(N)$. Here the action of G on Irr(N) is contragredient to its conjugation action on N.
- (2) Write $G(U) = G/G^1$. Then we have $X = \operatorname{Ind}_{G^1}^G X^1$ for some $X^1 \in \operatorname{Irr}(G^1)$ whose restriction to N is a multiple of U. Leaving out the superscripts,

- we are reduced to the *primary case* where U is G-invariant and X itself restricts to a multiple of U.
- (3) In that case we have a factorization $X = U \otimes T$, where the action of G on each factor is only projective. On U, it uniquely extends the linear action of N. On T, it comes from a projective action of G/N. Both give rise to central extensions of G/N by the circle, and these add up to zero in $H^2(G/N, S^1)$.

In [M58] Mackey famously extended this analysis to locally compact groups, and the central extension attached to $U \in Irr(N)^G$ has since been known as the *Mackey obstruction of* U. (For an exposition see [F88, pp. 1263, 1285].)

Now assume that G, N, etc. are Lie groups with Lie algebras \mathfrak{g} , \mathfrak{n} , etc. In an agenda pioneered by Kirillov, Kostant, and Souriau, one expects propositions about (not only irreducible) unitary G-modules to be reflected by propositions about hamiltonian G-spaces (i.e., symplectic manifolds on which our groups act with equivariant moment maps). Thus for instance, one expects that

- the role of decomposition into irreducibles is played by the decomposition of the image of the moment map into coadjoint orbits [K62];
- the role of multiplicities (such as T above) is played by Marsden-Weinstein reduced spaces [M74; G82];
- the role of tensor product is played by the cartesian product of hamiltonian G-spaces [W77, p. 32];
- the role of unitary induction is played by the symplectic induction construction of Kazhdan-Kostant-Sternberg [K78, p. 498];
- the role of Mackey's imprimitivity theorem is played by the symplectic imprimitivity theorem of [**Z96**, Thm 2.9; **L06**, Thm 4]; etc.

Using these ingredients, a complete geometric parallel of steps (1) and (2) was established in [Z96, Prop 3.8]. This reduces the study of a hamiltonian G-space X, when G has a closed normal subgroup N, to the *primary case* where the moment map for N is onto a single coadjoint orbit U. (We can think of this case as lying at the opposite end of the spectrum from the *multiplicity-free case* of [G84a], where all reduced spaces are points [G84b, 45.8].)

The purpose of this paper is to unravel the structure of primary spaces and thereby to complete step (3), which has proved more elusive. Here the representation theoretic analogy might lead one to expect a splitting $X = \bar{U} \times T$ where \bar{U} is homogeneous (a covering of U) and T a trivial N-space. The earliest results of this kind are those of Souriau [S70], which can be viewed as asserting just such a splitting when N is a Heisenberg group; see §4.1 below. The case N = SO(3) was dealt with in [I84; I91, Thm 4.3], and further generalizations occur in [L92; L94; Z96, §4], but always under hypotheses that guarantee a splitting of the above form. To our knowledge,

the literature contains neither a theory nor in fact a single example of the non-split situation, which we will see does actually occur.

The paper is divided into four sections. In section 1 we disregard the role of any ambient group G (or, what amounts to the same, assume that G = N) and show that while perhaps not split, primary N-spaces over a given coadjoint orbit U are always flat bundles $\tilde{U} \times_{\Gamma} V$ associated to a certain covering of U. (As Alan Weinstein kindly points out, this is a general property of Poisson maps to symplectic manifolds: [C99, §7.6].) Section 2 shows how the extra structure arising from the hamiltonian action of a group G normalizing N is encoded in the fiber V. Section 3 then spells out three corollaries which generalize Souriau's theorems to an arbitrary Lie group N, and comprise our symplectic "translation" of the Mackey obstruction step (3). Finally Section 4 shows that our theory is not empty by exhibiting examples which indeed do not split.

We conclude this introduction with two remarks of principle. First, while our notion of primary space arises naturally from the geometrical problem we set out to solve, one should not overestimate its parallelism with primary representations. Indeed several representations (parametrized by the characters of Γ) may correspond to an orbit U, and our non-split examples seem to correspond to the situation where one representation of G decomposes into many of them. Secondly, we have scrupulously avoided any connectedness assumptions, as these would spoil the recursive applicability of our symplectic "Mackey machine"; this explains why all coverings are constructed by hand, since one knows that the universal covering of a disconnected group or homogeneous space is not in general a group or homogeneous space [T54].

1. Decomposition of a primary N-space

Let N be a Lie group and let (X, ω, Π) be a hamiltonian N-space, that is, a symplectic manifold X with an action of N which preserves its 2-form ω , and an equivariant moment map $\Pi: X \to \mathfrak{n}^*$. If ω or Π are understood, we may drop them from the notation or use subscripts such as ω_X , Π_X , etc.

(1.1) **Definition.** We say that X is *primary* if the image of the moment map is a single coadjoint orbit U of N. When N (or U) needs emphasis we say N-*primary* (*over* U).

Our purpose is to describe and classify primary N-spaces over a given orbit U. We do not distinguish between two spaces X_1 , X_2 which are *isomorphic*, i.e., related by an N-equivariant diffeomorphism which transforms ω_1 into ω_2 and Π_1 into Π_2 .

(1.2) Example (Homogeneous spaces; split spaces). Any homogeneous hamiltonian N-space is primary: indeed one knows that its moment map is a covering map, $\bar{U} \to U$, onto a coadjoint orbit [K70; S70]. More generally

it is clear that the product

$$(1.3)$$
 $\bar{U} \times T$

of such a covering by any trivial hamiltonian N-space, is primary over U. Here by a *trivial* hamiltonian N-space T we mean one with trivial N-action and zero moment map, $T \to \{0\}$; and we recall that the *product* $X \times Y$ of two hamiltonian N-spaces is the product manifold with 2-form $\omega_X + \omega_Y$, action n(x,y) = (n(x), n(y)), and moment map $\Pi_X + \Pi_Y$.

- (1.4) **Definition.** We say that a primary space over U *splits* if it has the form (1.3) ($\bar{\mathbf{U}}$ homogeneous, T trivial). We further characterize the splitting as *trivial* or not according as the covering $\bar{\mathbf{U}} \to \mathbf{U}$ is trivial (bijective) or not.
- (1.5) Example (Flat bundles). Split spaces, it will turn out, do not exhaust primary spaces. Instead we have to resort to the next simplest construction, which we now present. First observe that the class of homogeneous N-spaces over U contains a maximal object, $\tilde{\mathbf{U}}$, which we may describe as follows. Write K for the stabilizer \mathbf{N}_c of a point $c \in \mathbf{U}$ which we fix once and for all, \mathbf{K}^o for its identity component, and

$$\tilde{\mathbf{U}} \xrightarrow{\rho} \mathbf{U}$$

for the covering $N/K^o \longrightarrow N/K$ with principal group

(1.7)
$$\Gamma = K/K^{o}$$

(a.k.a. the "N-equivariant fundamental group" of U [C93]). Endowing \tilde{U} with the pull-back $\rho^*\omega_U$ of the Kirillov-Kostant-Souriau 2-form of U, we obtain a primary N-space with moment map ρ .

Now suppose that V is any primary Γ -space, or in other words, any symplectic manifold with an ω_V -preserving action of Γ . (This is hamiltonian with moment map $V \to \{0\}$.) We may then form the associated "flat bundle"

$$(1.8) \tilde{\mathbf{U}} \times_{\Gamma} \mathbf{V} \xrightarrow{\Pi} \mathbf{U}$$

and observe that its total space (the set of orbits $[\tilde{u}, v]$ of the product action of Γ on $\tilde{U} \times V$) is naturally a primary N-space over U: the 2-form is the one obtained from $\rho^*\omega_U + \omega_V$ by passage to the quotient, N acts by $n([\tilde{u}, v]) = [n(\tilde{u}), v]$, and the moment map is $\Pi([\tilde{u}, v]) = \rho(\tilde{u})$.

Conversely:

(1.9) Theorem (Generalized Barycentric Decomposition). (i) Every primary N-space (X, ω, Π) over U is obtained in this way. More precisely one always has

(1.10)
$$X = \tilde{U} \times_{\Gamma} V$$
 where $V = \Pi^{-1}(c)$.

(ii) Two primary N-spaces $(X_i, \omega_i, \Pi_i)_{i=1,2}$ over U are isomorphic iff the fibers $V_i = \Pi_i^{-1}(c)$ are isomorphic as primary Γ -spaces.

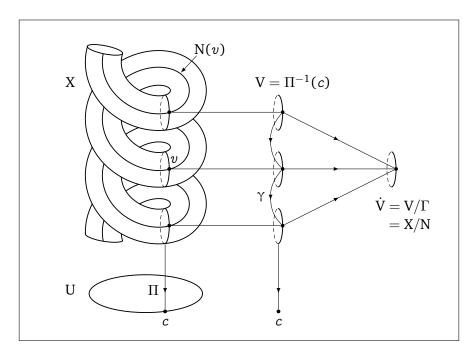


Figure 1. Barycentric decomposition of the primary space X.

Implied here is the assertion that $\Pi^{-1}(c)$ is always naturally a primary Γ -space. This is ensured by the following lemma and corollary, which give the decomposition at the infinitesimal level.

(1.11) Lemma (Fig. 1). Let (X, ω, Π) be a primary N-space over U = N(c). Then the fiber $V = \Pi^{-1}(c)$ is a symplectic submanifold of X and we have, for each $v \in V$, the direct sum decomposition into symplectically orthogonal subspaces

(1.12)
$$T_vX = T_vN(v) \oplus T_vV.$$

Proof. We know that $\Pi: X \to \mathfrak{n}^*$ remains smooth when regarded as a map $X \to U$ where U = N/K has the quotient manifold structure [O04, 2.3.12(i)]. In particular it is (by equivariance) a submersion: hence, its fiber V is a submanifold. Moreover, by a well-known property of moment maps [S70, 12.152; M74, Lemma (ii)], the tangent space to V at v is

(1.13)
$$T_v V = \ker(D\Pi(v)) = \mathfrak{n}(v)^{\omega},$$

the orthogonal of the tangent space $T_v N(v) = \mathfrak{n}(v)$ relative to the 2-form ω . But $\mathfrak{n}(v)$ and $\mathfrak{n}(v)^{\omega}$ span $T_v X$: indeed, given $\delta v \in T_v X$ we can find a $\nu \in \mathfrak{n}$ such that $D\Pi(v)(\delta v) = \nu(c)$, and then δv is the sum of

(1.14)
$$\nu(v) \in \mathfrak{n}(v)$$
 and $\delta v - \nu(v) \in \ker(\mathrm{D}\Pi(v)),$

where the last inclusion is by equivariance of Π . Now, taking orthogonals on both sides of the relation $T_vX = \mathfrak{n}(v) + \mathfrak{n}(v)^{\omega}$, we get $\{0\} = \mathfrak{n}(v) \cap \mathfrak{n}(v)^{\omega}$. So the sum is direct, and each summand is a symplectic subspace of T_vX . \square

(1.15) Corollary. In the setting of (1.11), the action of the stabilizer $K = N_c$ on V induces the trivial action of K^o , so that V carries a natural action of $\Gamma = K/K^o$.

Proof. The relation $D\Pi(v)(\mathfrak{k}(v)) = \mathfrak{k}(c) = 0$ shows that the tangent space $T_vK(v) = \mathfrak{k}(v)$ lies in the intersection $\mathfrak{n}(v) \cap \ker(D\Pi(v))$, which we have just seen is zero. Thus \mathfrak{k} acts trivially on V, and hence so does the identity component K^o .

Proof of the theorem. (i) We claim that the flat bundle (1.8) built from $V = \Pi^{-1}(c)$ is isomorphic to (X, ω, Π) . To see this, let us denote it by (Y, ω_Y, Π_Y) . We have N-equivariant submersions α, β :

(1.16)
$$X \stackrel{\alpha}{\longleftarrow} N \times V \stackrel{\beta}{\longrightarrow} Y$$

$$n(v) =: x \longleftarrow (n, v) \longmapsto y := n([eK^{o}, v])$$

where N acts on $N \times V$ via its left action on N. Now since

$$\beta(n', v') = \beta(n, v) \Leftrightarrow (n'K^{o}, v') = (nK^{o}\gamma^{-1}, \gamma(v)) \qquad \text{for some } \gamma \in \Gamma$$

$$\Leftrightarrow (n', v') = (nk^{-1}, k(v)) \qquad \text{for some } k \in K$$

$$\Leftrightarrow n'(v') = n(v) \text{ and } n^{-1}n' \in K$$

$$\Leftrightarrow \alpha(n', v') = \alpha(n, v),$$

we see that the fibers of α and β define the same partition of N × V. So (1.16) induces an N-equivariant bijection X \rightarrow Y, which is a diffeomorphism by [B67, 5.9.6], and which pulls $\Pi_{\rm Y}$ back to Π since $\Pi(\alpha(n,v)) = n(c) = \Pi_{\rm Y}(\beta(n,v))$. We must still show that this diffeomorphism is symplectic. To this end, write $\delta x \longleftrightarrow (\delta n, \delta v) \longleftrightarrow \delta y$ for the diagram tangent to (1.16), i. e.

(1.18)
$$\delta x = n_*(\nu(v) + \delta v) \quad \text{and} \quad \delta y = n_*([\nu(e\mathbf{K}^{\mathrm{o}}), \delta v])$$

where $\nu = n^{-1}\delta n$ and the bracket denotes the derived projection $T\tilde{U}\times TV\to TY$. Then we have

$$\omega(\delta x, \delta' x) = \omega(\nu(v) + \delta v, \nu'(v) + \delta' v) \qquad \text{since } n^* \omega = \omega$$

$$= \omega(\nu(v), \nu'(v)) + \omega(\delta v, \delta' v) \qquad \text{by (1.12)}$$

$$= \langle c, [\nu', \nu] \rangle + \omega_{V}(\delta v, \delta' v) \qquad \text{by [S70, 11.17\sharp]}$$

$$= \omega_{Y}(\delta y, \delta' y) \qquad \text{by definition,}$$

as was to be shown. This completes the proof of (i).

(ii) Isomorphic X_i give rise to isomorphic V_i . Indeed, let an isomorphism $F: X_1 \to X_2$ be given. Then the relation $\Pi_1 = \Pi_2 \circ F$ implies that F maps V_1

onto V_2 and induces a diffeomorphism f between them. Since the V_i have by construction the 2-form and action of $\gamma \in \Gamma$ induced by the ω_i and by the action of some $k \in \gamma$ on the X_i , it is clear that f is another isomorphism.

Conversely, isomorphic V_i give rise to isomorphic X_i . Indeed, let an isomorphism $f: V_1 \to V_2$ be given. Then in the diagram

$$(1.20) \qquad \begin{array}{c} \tilde{\mathrm{U}} \times \mathrm{V}_1 & \xrightarrow{\mathrm{id} \times f} & \tilde{\mathrm{U}} \times \mathrm{V}_2 \\ & & \downarrow & & \downarrow \\ & \text{mod} \, \Gamma \\ & \tilde{\mathrm{U}} \times_{\Gamma} \mathrm{V}_1 & \xrightarrow{\mathrm{F}} & \tilde{\mathrm{U}} \times_{\Gamma} \mathrm{V}_2 \end{array}$$

one verifies without trouble that the top arrow is an isomorphism of primary $N \times \Gamma$ -spaces, relative to the 2-forms $\rho^* \omega_U + \omega_{V_i}$ and actions $(n, \gamma)(\tilde{u}, v_i) = (n\tilde{u}\gamma^{-1}, \gamma(v_i))$. So it descends to an isomorphism F between the quotients, of which we already know by (1.10) that they are isomorphic to X_1 and X_2 . \square

(1.21) Corollary 1. If the pair (N, U) is in any of the following cases, then every primary N-space over U splits trivially as $X = U \times V$:

- (i) points of U have connected stabilizers in N;
- (ii) N is connected and U is simply connected;
- (iii) N is connected and compact or abelian or nilpotent or exponential;
- (iv) N is connected, complex semisimple and U is a semisimple orbit.

Proof. (i) If K is connected, then $\Gamma = \{1\}$ and (1.10) reduces to $X = U \times V$; (ii) implies (i) by the homotopy exact sequence $\pi_1(U) \to \pi_0(K) \to \pi_0(N)$; (iii, iv) each imply (i) or (ii): see [B72, p. 4; B87, 8.20; C93, 2.3.4].

(1.22) Remark. The component groups Γ are also well known when N is connected solvable or simple: in the former case they are lattices \mathbf{Z}^d [A71]; in the latter, finite products of groups \mathbf{Z}_d , S₃, S₄, S₅ [K92; C93]. Lest the reader conclude that they are always dull, we note here that every finite group occurs for some coadjoint orbit of some connected Lie group [P78, Prop. 1(iii)].

As soon as $\Gamma \neq \{1\}$ it becomes an issue to detect which primary spaces do or don't split. For this we have the following simple criteria.

(1.23) Corollary 2. Let X be primary with fiber V. (i) If X splits trivially, then Γ acts trivially on V. (ii) If X splits nontrivially, then V is not connected.

Proof. (i) The hypothesis means that $(X, \omega, \Pi) \cong (U \times T, \omega_U + \omega_T, \operatorname{pr}_1)$ for some T on which N acts trivially. By (1.9ii) it follows that $V \cong \{c\} \times T$, on which Γ acts trivially. (ii) Likewise if $(X, \omega, \Pi) \cong (\overline{U} \times T, \pi^* \omega_U + \omega_T, \pi \circ \operatorname{pr}_1)$ for some nontrivial covering π then $V \cong \pi^{-1}(c) \times T$, which is not connected.

- (1.24) Remarks. (i) We emphasize that thanks essentially to the rigidity of the moment map these criteria are very much simpler than the topological arguments one would need to decide whether $U \times_{\Gamma} V$ is trivial as a bundle, or whether it is homeomorphic to a product $\bar{U} \times T$.
- (ii) It may also be worth noting that Theorem (1.9) places any primary space half-way between two products:

$$(1.25) \qquad \overset{\tilde{U}\times V}{\longrightarrow} X \overset{\text{mod }\Gamma}{\longrightarrow} U\times \dot{V},$$

where \dot{V} is the quotient $X/N = V/\Gamma$ (Fig. 1). However, the upper product here does not naturally carry an action of N; and the lower one is only a (trivially split) primary space insofar as \dot{V} is a manifold, which fails in general: see e.g. Example (4.21) below, where \dot{V} is not even an orbifold.

(1.26) Example (Reduced spaces of Kazhdan *et al.*). An important source of primary spaces is the construction of Kazhdan-Kostant-Sternberg [K78, p. 482]: if (Y, σ, Ψ) is an arbitrary hamiltonian N-space, then under appropriate transversality hypotheses one can form the *reduced space at* U,

(1.27)
$$X = \Psi^{-1}(U) / \ker(\sigma)$$

(the quotient of $\Psi^{-1}(U)$, which is coisotropic, by its characteristic foliation). This is naturally a primary N-space over U, to which the above applies. In this setting we note that (1.11), (1.21i) and a version of (1.9i) are found respectively in [K78, Prop. 1.1; G84b, Thm 26.6; O84, Thm 1]. (Our treatment avoids the assumption, made in these references, that $U \hookrightarrow \mathfrak{n}^*$ be an embedding.) Also, the spaces \dot{V} and $U \times \dot{V}$ of (1.25) are none other than the quotients

(1.28)
$$\Psi^{-1}(c)/N_c \qquad \text{and} \qquad \Psi^{-1}(U)/R$$

of Marsden-Weinstein [M74] and Marle [M76] (to which we refer for the definition of the equivalence relation R). In particular Marle's quotient always splits trivially, beyond the cases elucidated in [M76, Cor. p. 257].

2. Extensions of a primary N-space

From here on we assume again that N is a closed normal subgroup of a Lie group G. Our aim is to describe those N-primary spaces (X, ω, Π) over U that arise by *restriction* of an action of the larger group — i.e. X admits a hamiltonian G-action, with equivariant moment map Φ , such that the

following diagrams commute:

(2.1)
$$\bigcap_{N \text{ N-action}} G \text{ G-action}$$

$$X \bigcap_{\Pi \to \mathfrak{n}^*} \mathfrak{g}^*$$

Suppose for a moment that such an extended action exists. Because it normalizes N, G has coadjoint actions on both \mathfrak{g}^* and \mathfrak{n}^* , relative to which the second diagram (2.1) is equivariant. In particular G preserves the image of Π , so U must belong to the set $(\mathfrak{n}^*/N)^G$ of G-invariant coadjoint orbits of N. Under this assumption we have the following result, which in effect reduces the extension problem to the case where U is a point.

(2.2) Theorem. Let N be normal in G, pick an orbit $U = N(c) \in (\mathfrak{n}^*/N)^G$, and write K and L for the stabilizers N_c and G_c . Then the correspondence $X \rightleftharpoons V$ of Theorem (1.9) naturally induces another between

N-primary hamiltonian G-spaces
$$(X, \omega, \Phi)$$
 over U \longleftrightarrow K-primary hamiltonian L-spaces (V, ω_V, Ψ) over the point $\{c_{|\mathfrak{k}}\}$.

The proof will make repeated use of the following elementary observation, which essentially describes the trivial case where $U = \{0\}$.

(2.3) Lemma. If N is normal in G, the annihilator $\operatorname{ann}_{\mathfrak{g}^*}(\mathfrak{n}) = \ker(\mathfrak{g}^* \to \mathfrak{n}^*)$ is

- (i) invariant under the coadjoint action of G on g*,
- (ii) pointwise fixed by the action of N,
- (iii) as a G/N-module, canonically isomorphic to $(\mathfrak{g}/\mathfrak{n})^*$.

Proof. Write $\underline{l}(g) = lgl^{-1}$ for the conjugation action of $l \in G$ on elements (or subsets) of G. Since N is normal, \underline{l} maps each coset gN to another such coset,

(2.4)
$$\underline{l}(gN) = lgNl^{-1} = \underline{l}(g)N.$$

So we have an induced conjugation action on G/N which is just the lift of its conjugation action on itself, for (2.4) can also be written $(lN)gN(lN)^{-1}$. In other words, the exact sequence

$$(2.5) 1 \longrightarrow \underset{G}{\overset{N}{\longrightarrow}} \underset{G}{\overset{G}{\longrightarrow}} \underset{G}{\overset{G/N}{\longrightarrow}} 1$$

is G-equivariant, and the third G-action is in fact an action of G/N (so N acts trivially). Deriving at e and dualizing, we get exactness and equivariance of

$$(2.6) \qquad 0 \longrightarrow (\mathfrak{g/n})^* \longrightarrow \mathfrak{g}^* \longrightarrow \mathfrak{n}^* \longrightarrow 0$$

relative to the 3 coadjoint actions, of which the first is in fact lifted from G/N (so N acts trivially). This is precisely what the lemma claims.

Proof of the theorem. Suppose that (X, ω, Φ) is given. The G-equivariance of Π ensures that the stabilizer L of c preserves $V = \Pi^{-1}(c)$. So V becomes a hamiltonian L-space, whose moment map Ψ we can define by the commutativity of

(2.7)
$$X \xrightarrow{\Phi} \mathfrak{g}^* \longrightarrow \mathfrak{n}^* \ni c$$

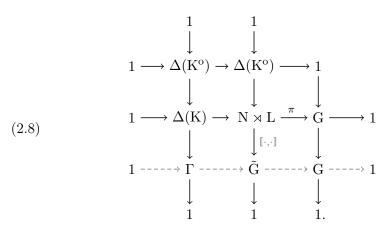
$$\downarrow \qquad \qquad \downarrow$$

$$V \xrightarrow{\Psi} \mathfrak{f}^* \longrightarrow \mathfrak{k}^* \ni c_{|\mathfrak{k}}.$$

Because V sits above $c \in \mathfrak{n}^*$ by construction, it follows that it sits above $c_{|\mathfrak{k}}$ under the composition $V \to \mathfrak{l}^* \to \mathfrak{k}^*$, which is the moment map for the induced K-action; in other words V is K-primary over $\{c_{|\mathfrak{k}}\}$, as claimed.

Conversely suppose that (V, ω_V, Ψ) is given. We must construct on $X := \tilde{U} \times_{\Gamma} V$ a G-action and moment map Φ satisfying (2.1). The rough (but wrong) idea is that G acts separately on U and V, with moment maps ϕ and ψ that add up to an equivariant Φ . To make this correct we must in fact take a detour via \tilde{U} , where ϕ will be defined, and an attendant \tilde{G} which actually acts on \tilde{U} and V. We do this in the following steps, where $N \rtimes L$ denotes the semidirect product with law $(n,l)(n',l') = (n\underline{l}(n'),ll')$ (notation (2.4)).

1. We can define a covering extension G of G by Γ by the following commutative diagram of exact sequences, where $\Delta(k) = (k^{-1}, k)$ and $\pi(n, l) = nl$:



Indeed, π clearly has kernel $\Delta(K)$ and is onto: given $g \in G$, we can find $n \in N$ such that g(c) = n(c), whence $g = \pi(n, n^{-1}g)$. So the top and middle row are exact. In particular $N \rtimes L$ normalizes $\Delta(K)$, hence also its identity component

 $\Delta(K^{\circ})$ [H63, 7.16]. Writing \tilde{G} for the resulting quotient and

(2.9)
$$(n,l) \mapsto [n,l]$$
 and $\tilde{g} \mapsto g$

for the projections $N \rtimes L \to \tilde{G}$ and $\tilde{G} \to G$, we have exact columns too, so the nine lemma (or the second isomorphism theorem [H63, 5.35]) completes the diagram. Before continuing we observe that $[\![\cdot,e]\!]$ and $[\![e,\cdot]\!]$ map N and L injectively onto isomorphic copies which we shall denote

$$\tilde{\mathbf{N}} = [\![\mathbf{N}, e]\!], \qquad \qquad \tilde{\mathbf{L}} = [\![e, \mathbf{L}]\!].$$

(The gain inside \tilde{G} is that unlike $N \cap L$ (= K), the intersection $\tilde{N} \cap \tilde{L}$ (= $[\![K^o,e]\!]$) is always connected.)

2. We claim that the coadjoint action of G on U lifts to an action of \tilde{G} on \tilde{U} . To see this, we note that (2.8) contains the relations G = NL and $\tilde{G} = \tilde{N}\tilde{L}$. Therefore the first isomorphism theorem [H63, 5.33] implies that all horizontal maps in the following commutative diagrams are bijections:

Now the first diagram identifies our N-equivariant covering (1.6), on the left, with the \tilde{G} -equivariant covering on the right. This equips \tilde{U} with an action of \tilde{G} which fits the bill and works out, since $[n, l][n', L] = [n\underline{l}(n'), L]$, as

(2.12)
$$[n, l](\tilde{u}) = n\underline{l}(\tilde{u}), \qquad \tilde{u} = n'K^{o} \in N/K^{o}.$$

3. This action (2.12) preserves $\omega_{\tilde{U}} = \rho^* \omega_{\tilde{U}}$ and admits a (not necessarily equivariant) moment map $\phi: \tilde{U} \to \mathfrak{g}^*$ given by

$$\phi(\tilde{u}) = \tilde{u}(\check{c}),$$

where we fix henceforth an element $\check{c} \in \mathfrak{g}^*$ projecting to $c \in \mathfrak{n}^*$. Indeed, first we note that (2.13) is well-defined, i.e. all elements of a coset $\tilde{u} = n\mathrm{K}^{\mathrm{o}}$ have the same effect on \check{c} . This is because any two differ by an element of K^{o} , which acts trivially by (1.15) applied to the orbit $\mathrm{G}(\check{c})$ (which is clearly N-primary over U). Next, (2.13) is a moment map because we have, for $\tilde{u} = n\mathrm{K}^{\mathrm{o}}$ and all $(\nu, \gamma) \in \mathfrak{n} \times \mathfrak{g}$,

$$[\iota_{\gamma}\omega_{\tilde{\mathbf{U}}} + d\langle\phi(\cdot),\gamma\rangle](\nu(\tilde{u})) = \omega_{\tilde{\mathbf{U}}}(\gamma(\tilde{u}),\nu(\tilde{u})) + \frac{d}{dt}\langle\phi(\mathbf{e}^{t\nu}(\tilde{u})),\gamma\rangle\big|_{t=0}$$

$$= \omega_{\mathbf{U}}(\gamma(u),\nu(u)) + \frac{d}{dt}\langle\mathbf{e}^{t\nu}(n(\check{c})),\gamma\rangle\big|_{t=0}$$

$$= \langle n(c), [\nu,\gamma]\rangle + \langle n(\check{c}), [\gamma,\nu]\rangle$$

$$= 0.$$

If \tilde{G} is disconnected we still need to check that the action (2.12) preserves $\omega_{\tilde{U}}$. Since ρ in (2.11) intertwines it with the coadjoint action of G on U,

it suffices to show that $g^*\omega_{\rm U}=\omega_{\rm U}$, which is true because we have, for all $\nu,\nu'\in\mathfrak{n}$,

$$g^*\omega_{\mathrm{U}}(\nu(u), \nu'(u)) = \omega_{\mathrm{U}}(g_*(\nu(u)), g_*(\nu'(u)))$$

$$= \omega_{\mathrm{U}}((\mathrm{Ad}_g(\nu))(g(u)), (\mathrm{Ad}_g(\nu'))(g(u)))$$

$$= \langle g(u), [\mathrm{Ad}_g(\nu'), \mathrm{Ad}_g(\nu)] \rangle$$

$$= \langle u, [\nu', \nu] \rangle$$

$$= \omega_{\mathrm{U}}(\nu(u), \nu'(u)).$$

4. The given action of L on V induces an ω_V -preserving action of \tilde{G} on V, viz.

$$[n, l](v) = l(v).$$

Indeed, this is well-defined because [n, l] determines $l \in L$ up to multiplication by an element of K°, which we know acts trivially on V (by Cor. 1.15, applied this time to V as a K-primary L-space). Moreover it preserves ω_{V} , since the given action of L does.

5. This action (2.16) admits a (not necessarily equivariant) moment map $\psi: V \to \mathfrak{g}^*$ given by the following formula, where p^* denotes the injection $(\mathfrak{l}/\mathfrak{k})^* \xrightarrow{(2.11)} (\mathfrak{g}/\mathfrak{n})^* \xrightarrow{(2.6)} \mathfrak{g}^*$ viewed as a map $\operatorname{ann}_{\mathfrak{l}^*}(\mathfrak{k}) \to \operatorname{ann}_{\mathfrak{g}^*}(\mathfrak{n})$ (2.3):

(2.17)
$$\psi(v) = p^*(\Psi(v) - \check{c}_{|\mathfrak{l}}).$$

This is well-defined because our hypothesis that V is K-primary over $\{c_{|\mathfrak{k}}\}$ ensures $\Psi(v) - \check{c}_{|\mathfrak{l}} \in \operatorname{ann}_{\mathfrak{l}^*}(\mathfrak{k})$. To see that it is a moment map, we note that (2.8) implies that every $\gamma \in \mathfrak{g}$ can be written $\nu + \lambda$ for some $(\nu, \lambda) \in \mathfrak{n} \rtimes \mathfrak{l}$. Then,

$$\iota_{\gamma}\omega_{V} + d\langle\psi(\cdot),\gamma\rangle = \iota_{\lambda}\omega_{V} + d\langle\psi(\cdot),\gamma\rangle \qquad \text{by (2.16)}$$

$$= \iota_{\lambda}\omega_{V} + d\langle\psi(\cdot),\lambda\rangle \qquad \text{since range}(\psi) \subset \text{ann}_{\mathfrak{g}^{*}}(\mathfrak{n})$$

$$= \iota_{\lambda}\omega_{V} + d\langle\Psi(\cdot),\lambda\rangle \qquad \text{since } \langle\check{c},\lambda\rangle \text{ is constant}$$

$$= 0 \qquad \qquad \text{since } \Psi \text{ is a moment map.}$$

6. Putting the actions (2.12) and (2.16) together, we get an action of \tilde{G} on $\tilde{U} \times V$ which we claim is hamiltonian, i.e. its moment map $\phi + \psi$ is \tilde{G} -equivariant. In other words we will show that $\theta_{\tilde{U}} + \theta_{V} = 0$ where

(2.19)
$$\begin{aligned} \theta_{\tilde{\mathbf{U}}}(\tilde{g}) &:= \phi(\tilde{g}(\tilde{u})) - g(\phi(\tilde{u})), \\ \theta_{\mathbf{V}}(\tilde{g}) &:= \psi(\tilde{g}(v)) - g(\psi(v)). \end{aligned}$$

To work out these 'non-equivariance cocycles' (which as the notation anticipates, will be independent of \tilde{u} and v; note that for disconnected X this is emphatically *not* guaranteed by the general theory [S70, 11.17; O04,

4.5.21), we pick $\tilde{g} = [n, l]$ projecting to g = nl and compute, on the \tilde{U} side,

$$\theta_{\tilde{\mathbf{U}}}(\tilde{g}) = \phi(n\underline{l}(\tilde{u})) - (nl)(\phi(\tilde{u})) \qquad \text{by (2.12)}$$

$$= (n\underline{l}(\tilde{u}) - nl\tilde{u})(\check{c}) \qquad \text{by (2.13)}$$

$$= n\underline{l}(\tilde{u})(\check{c} - l(\check{c}))$$

$$= \check{c} - l(\check{c}) \qquad \text{by Lemma (2.3ii),}$$

which lemma applies because $l \in G_c$ implies $\langle \check{c} - l(\check{c}), \mathfrak{n} \rangle = \langle c - l(c), \mathfrak{n} \rangle = 0$, i.e. we have

$$(2.21) \check{c} - l(\check{c}) \in \operatorname{ann}_{\mathfrak{a}^*}(\mathfrak{n}).$$

Meanwhile on the V side we find, for all $\gamma = \nu + \lambda$ in $\mathfrak{g} = \mathfrak{n} + \mathfrak{l}$,

$$\langle \theta_{\mathrm{V}}(\tilde{g}), \gamma \rangle = \langle \psi(l(v)) - (nl)(\psi(v)), \gamma \rangle \qquad \text{by (2.16)}$$

$$= \langle \psi(l(v)) - l(\psi(v)), \gamma \rangle \qquad \text{by (2.3i, ii)}$$

$$= \langle \psi(l(v)) - l(\psi(v)), \lambda \rangle \qquad \text{since range}(\psi) \subset \operatorname{ann}_{\mathfrak{g}^*}(\mathfrak{n})$$

$$(2.22) \qquad = \langle \Psi(l(v)) - \check{c}_{\mid \mathbf{l}} - l(\Psi(v) - \check{c}_{\mid \mathbf{l}}), \lambda \rangle \qquad \text{by (2.17)}$$

$$= \langle l(\check{c}) - \check{c}, \lambda \rangle \qquad \text{since } \Psi \text{ is L-equivariant}$$

$$= \langle l(\check{c}) - \check{c}, \gamma \rangle \qquad \text{by (2.21)}.$$

Thus we have $\theta_{\tilde{U}} + \theta_{V} = 0$, which proves that $\phi + \psi$ is \tilde{G} -equivariant.

7. Next we consider the subgroup $\Gamma = \{ \llbracket k^{-1}, k \rrbracket : k \in K \}$ in (2.8) and observe that its orbits in $\tilde{\mathbb{U}} \times \mathbb{V}$, under the action just considered, are exactly the points of $X = \tilde{\mathbb{U}} \times_{\Gamma} \mathbb{V}$. Since $\tilde{\mathbb{G}}$ normalizes Γ , it follows that its action takes Γ -orbit to Γ -orbit and hence descends to an action of $G = \tilde{\mathbb{G}}/\Gamma$ on X, viz.

(2.23)
$$nl([\tilde{u}, v]) = [n\underline{l}(\tilde{u}), l(v)].$$

Moreover the Γ -equivariance of $\phi + \psi$ means precisely that this map is constant on Γ -orbits, hence descends to an equivariant moment map $\Phi : X \to \mathfrak{g}^*$. Tracing through its construction yields explicitly, with notation as above,

$$(2.24) \qquad \langle \Phi([\tilde{u}, v]), \nu + \lambda \rangle = \langle \rho(\tilde{u}), \nu \rangle + \langle \tilde{u}(\check{c}) - \check{c}, \lambda \rangle + \langle \Psi(v), \lambda \rangle.$$

Finally we check that Φ does not depend on the choice of \check{c} . Indeed if $\check{c}' \in \mathfrak{g}^*$ is another choice projecting to $c \in \mathfrak{n}^*$, then the difference $z = \check{c}' - \check{c}$ lies in $\operatorname{ann}_{\mathfrak{g}^*}(\mathfrak{n})$. So (2.3ii) implies $\tilde{u}(z) - z = 0$, which shows that the change leaves (2.24) unaltered. This completes the proof.

As in (1.9), X and V determine each other to within isomorphism. More precisely we have the following proposition, whose proof is straightforward and left to the reader.

(2.25) Proposition. Under the correspondence of Theorem (2.2),

- (i) Two N-primary hamiltonian G-spaces X_1, X_2 over U are isomorphic iff the corresponding $V_i = \Pi_i^{-1}(c)$ are isomorphic as hamiltonian L-spaces.
- (ii) X is G-homogeneous iff V is L-homogeneous, and X is the coadjoint orbit $G(\check{c})$ (where $\check{c}_{|\mathfrak{n}}=c$) iff V is the coadjoint orbit $L(\check{c}_{|\mathfrak{l}})$.

3. Kœnig's theorem and the Mackey obstruction

In this section we spell out three corollaries of Theorem (2.2) and its proof. We keep the notations K, Γ , \tilde{U} , \tilde{G} , \tilde{N} , \tilde{L} established in (1.6–1.7) and (2.8–2.10). The first corollary merely records how we reconstructed the moment map Φ from two independent pieces:

(3.1) Corollary 1 (Generalized Kœnig Theorem). Let (X, ω, Φ) be an N-primary hamiltonian G-space over U with fiber V. Then not only does G act on $X = \tilde{U} \times_{\Gamma} V$, but the product $\tilde{G} \times \tilde{G}$ acts on $\tilde{U} \times V$ with moment map (ϕ, ψ) such that

(3.2)
$$\Phi([\tilde{u}, v]) = \phi(\tilde{u}) + \psi(v).$$

The second action is really an action of \tilde{G}/\tilde{N} with moment map $\psi: V \to \operatorname{ann}_{\mathfrak{g}^*}(\mathfrak{n})$.

Unlike Φ , the maps ϕ , ψ , $\theta_{\tilde{U}}$, θ_{V} do depend on the choice of \check{c} such that $\check{c}_{|\mathfrak{n}} = c$. The next corollary extracts the intrinsic part of this construction:

(3.3) Corollary 2. Attached to each $U = N(c) \in (\mathfrak{n}^*/N)^G$ is a well-defined cohomology class $[\theta_{\tilde{U}}] \in H^1(\tilde{G}/\tilde{N}, (\mathfrak{g}/\mathfrak{n})^*)$ which measures the obstruction to making U a hamiltonian G-space, and vanishes if $c_{|\mathfrak{k}} = 0$. If $c_{|\mathfrak{k}} \neq 0$, the derived class $[\langle D\theta_{\tilde{U}}(e)(\cdot), \cdot \rangle] \in H^2(\mathfrak{g}/\mathfrak{n}, \mathbf{R})$ is that of the central extension

$$(3.4) 0 \longrightarrow \mathfrak{k}/\mathfrak{j} \longrightarrow \mathfrak{l}/\mathfrak{j} \longrightarrow \mathfrak{l}/\mathfrak{k} \longrightarrow 0$$

where $\mathfrak{j} = \ker(c_{|\mathfrak{k}})$.

Proof. These are all elementary properties of the function $\theta_{\tilde{\mathbf{U}}}$ in (2.20). First, by an easy application of (2.3ii), it satisfies the cocycle identity $\theta_{\tilde{\mathbf{U}}}(\tilde{g}\tilde{g}') = \theta_{\tilde{\mathbf{U}}}(\tilde{g}) + g(\theta_{\tilde{\mathbf{U}}}(\tilde{g}'))$ which defines $\mathbf{Z}^1(\tilde{\mathbf{G}}, \mathfrak{g}^*)$ [B80, p. 112]. As it takes its values in $\mathrm{ann}_{\mathfrak{g}^*}(\mathfrak{n})$ (2.21) and clearly only depends on $\tilde{g} = [n, l]$ via its class [N, l], we may equally regard it as a member of $\mathbf{Z}^1(\tilde{\mathbf{G}}/\tilde{\mathbf{N}}, (\mathfrak{g}/\mathfrak{n})^*)$. Moreover, replacing \tilde{c} by $\tilde{c} + z$, also projecting on $c \in \mathfrak{n}^*$, alters the cocycle by a term z - l(z) which is the coboundary of $-z \in \mathrm{ann}_{\mathfrak{g}^*}(\mathfrak{n})$. So the resulting cohomology class $[\theta_{\tilde{\mathbf{U}}}]$ does not depend on our choice of \tilde{c} . Next, if $c_{|\mathfrak{k}} = 0$ then we may take \tilde{c} to be c on \mathfrak{n} and 0 on \mathfrak{l} , in which case the cocycle vanishes: $\langle \tilde{c} - l(\tilde{c}), \mathfrak{n} + \mathfrak{l} \rangle = \langle c - l(c), \mathfrak{n} \rangle + \langle 0 - l(0), \mathfrak{l} \rangle = 0$ since l stabilizes c. Finally

we consider the derivative $f_{\tilde{\mathbf{U}}} = \langle \mathrm{D}\theta_{\tilde{\mathbf{U}}}(e)(\cdot), \cdot \rangle$:

$$\begin{split} f_{\tilde{\mathbf{U}}}(\nu+\lambda,\nu'+\lambda') &= \frac{d}{dt} \left\langle \theta_{\tilde{\mathbf{U}}}(\llbracket \mathbf{e}^{t\nu},\mathbf{e}^{t\lambda} \rrbracket),\nu'+\lambda' \right\rangle \Big|_{t=0} \\ &= \frac{d}{dt} \left\langle \check{c} - \mathbf{e}^{t\lambda}(\check{c}),\nu'+\lambda' \right\rangle \Big|_{t=0} \\ &= \left\langle -\lambda(\check{c}),\lambda' \right\rangle \qquad \text{since } \left\langle \mathfrak{l}(c),\mathfrak{n} \right\rangle = 0 \\ (3.5) &= \left\langle \check{c},[\lambda,\lambda'] \right\rangle. \end{split}$$

As above one verifies that this defines a cocycle in $Z^2(\mathfrak{g/n}, \mathbf{R})$ [B80, p. 196] whose class in $H^2(\mathfrak{g/n}, \mathbf{R}) = H^2(\mathfrak{l/k}, \mathbf{R})$ (2.11) does not depend on the choice of \check{c} . To see that it coincides with the cohomology class of the extension (3.4), we recall that the latter is that of the cocycle [B80, p. 198]

(3.6)
$$f(\lambda + \mathfrak{k}, \lambda' + \mathfrak{k}) = \langle c, [s(\lambda + \mathfrak{k}), s(\lambda' + \mathfrak{k})] - s([\lambda + \mathfrak{k}, \lambda' + \mathfrak{k}]) \rangle$$

where s is any linear section of the projection $\mathfrak{l}/\mathfrak{j} \to \mathfrak{l}/\mathfrak{k}$, and where we have identified $\mathfrak{k}/\mathfrak{j}$ with \mathbf{R} by sending $\varkappa + \mathfrak{j}$ to $\langle c, \varkappa \rangle$. Choosing $s(\lambda + \mathfrak{k}) = q(\lambda) + \mathfrak{j}$ where p (resp. q) is the projection onto \mathfrak{k} (resp. \mathfrak{v}) associated to a vector space decomposition $\mathfrak{l} = \mathfrak{k} \oplus \mathfrak{v}$, this is

$$(3.7) = \langle c, [q(\lambda), q(\lambda')] - q([\lambda, \lambda']) \rangle$$

$$= \langle c, [\lambda - p(\lambda), \lambda' - p(\lambda')] - [\lambda, \lambda'] + p([\lambda, \lambda']) \rangle$$

$$= \langle c \circ p, [\lambda, \lambda'] \rangle$$

since $\langle \mathfrak{l}(c), \mathfrak{n} \rangle = 0$. Now $c \circ p$ defines an extension of c to \mathfrak{l} and thereby to $\mathfrak{g} = \mathfrak{n} + \mathfrak{l}$, which we can assume coincides with \check{c} (by changing the latter if necessary). So (3.5) and (3.7) belong to the same cohomology class, as was to be shown.

(3.8) **Definition.** We shall refer to the class $[\theta_{\tilde{U}}]$ in (3.3) as the *symplectic Mackey obstruction* of U, and to the class $[\langle D\theta_{\tilde{U}}(e)(\cdot), \cdot \rangle]$, or the central extension (3.4) it represents, as the *infinitesimal Mackey obstruction* of U.

Given $\Theta \in H^1(G, \mathfrak{g}^*)$ let us also agree to call **hamiltonian** (G, Θ) -space a symplectic manifold (X, ω) with an ω -preserving action of G which admits a moment map ϕ satisfying $\phi(g(x)) = g(\phi(x)) + \theta(g)$ identically for some $\theta \in \Theta$. (We emphasize that ϕ is not part of the structure: only its existence is.) With this in hand we can rephrase our solution (2.2) of the extension problem in a form more closely parallel to the representation-theoretic version (3), namely:

(3.9) Corollary 3. Let N be normal in G and let $U \in (\mathfrak{n}^*/N)^G$ have symplectic Mackey obstruction Θ . Then the correspondence $X \rightleftarrows V$ of Theorem 2.2 can be regarded as a correspondence

$$(3.10) \qquad \begin{array}{c} \text{N-primary hamiltonian} \\ \text{G-spaces X over U} \end{array} \quad \rightleftarrows \quad \begin{array}{c} \text{arbitrary hamiltonian} \\ (\tilde{G}/\tilde{N}, -\Theta)\text{-spaces V}. \end{array}$$

Proof. Let (X, ω, Φ) be N-primary over U, and consider the corresponding K-primary L-space (V, ω_V, Ψ) given by Theorem (2.2). We know (1.15) that K° acts trivially so that V is in fact a $L/K^{\circ} = \tilde{G}/\tilde{N}$ -space (2.11). Moreover we have seen that this action admits a moment map $\psi : V \to (\mathfrak{g}/\mathfrak{n})^*$ (2.17) whose non-equivariance cocycle (2.19) has cohomology class $-\Theta$.

Conversely, let (V, ω_V) be any hamiltonian $(\tilde{G}/\tilde{N}, -\Theta)$ -space. By (2.11) we can regard it as a hamiltonian $(L/K^o, -\Theta)$ -space, $\Theta \in H^1(L/K^o, (I/\mathfrak{k})^*)$ being the class of the cocycle $\theta_{\tilde{U}}(lK^o) = (\check{c} - l(\check{c}))_{|I} \in \operatorname{ann}_{I^*}(\mathfrak{k})$ (2.3iii, 2.20, 2.21). By definition this means that we can choose a moment map $\psi : V \to \operatorname{ann}_{I^*}(\mathfrak{k})$ so that

(3.11)
$$\psi(l(v)) = l(\psi(v)) + (l(\check{c}) - \check{c})_{|i|}$$

(equality in $\operatorname{ann}_{\mathfrak{l}^*}(\mathfrak{k}) = (\mathfrak{l}/\mathfrak{k})^*$). Now, lifting the action from L/K° to L enables us to add $\check{c}_{|\mathfrak{l}}$ to both sides to obtain the L-equivariance of $\Psi := \psi(\cdot) + \check{c}_{|\mathfrak{l}}$. Thus (V, ω_V, Ψ) is a K-primary hamiltonian L-space over $\check{c}_{|\mathfrak{k}}$. By (2.2) this corresponds to an N-primary hamiltonian G-space over U, as desired.

(3.12) Remark. Another justification for Definition (3.8) is that it appears to correctly predict, at least at the Lie algebra level, the Mackey obstruction (3) of the representation(s) obtained by "quantizing" U. To be specific, let us first note that (3.4) does not always integrate to a Lie group extension. (Indeed it is not hard to construct examples where the subgroup J° generated by j in (3.4) is not even closed.) However, when U is *integral* in the sense that K admits a character χ with differential $ic_{|\mathfrak{k}}$, then $J = \ker(\chi)$ is closed and the extension

$$(3.13) 1 \longrightarrow K/J \longrightarrow L/J \longrightarrow L/K \longrightarrow 1$$

integrates (3.4) and is precisely the Mackey obstruction found in [A71, Thm IV.4.1; D70, §6], in the case where N is nilpotent and G is solvable.

4. Examples

(4.1) Souriau's theorems. We consider here the group G of all block upper triangular matrices of the form

(4.2)
$$g = \begin{pmatrix} 1 & \overline{\boldsymbol{b}} \mathbf{A} & \frac{1}{2} \|\boldsymbol{b}\|^2 & f \\ \mathbf{A} & \boldsymbol{b} & \boldsymbol{c} \\ & 1 & e \\ & & 1 \end{pmatrix}$$

where $A \in SO(3)$, $b, c \in \mathbb{R}^3$, $e, f \in \mathbb{R}$ and the bar means transpose. The quotient obtained by forgetting the first row and column is the Galilei group and G is (up to a \mathbb{Z}_2 cover which need not concern us here) its universal central extension, where f is the central parameter. The axioms of classical

mechanics as spelled out in [S70, 13.1, 13.7], state: "A free dynamical system is represented by a connected symplectic manifold (X, ω) together with a ω -preserving action of the Galilei group admitting a moment map". As one knows [K68, p. 66], the resulting action of G always admits an equivariant moment map, so we can rephrase this as: A free dynamical system is represented by a connected hamiltonian G-space where the center acts trivially. Identifying \mathfrak{g}^* with \mathbf{R}^{11} by writing $(\mathbf{L}, \mathbf{G}, \mathbf{P}, \mathbf{E}, \mathbf{M})$ for the value of the 1-form

$$(4.3) \qquad \frac{1}{2}\operatorname{Tr}\left(\begin{pmatrix} 0 & L_3 & -L_2 \\ -L_3 & 0 & L_1 \\ L_2 & -L_1 & 0 \end{pmatrix}d\mathbf{A}\right) - \langle \mathbf{G}, d\mathbf{b} \rangle + \langle \mathbf{P}, d\mathbf{c} \rangle - \mathbf{E}de - \mathbf{M}df$$

at the identity, we observe that this implies, first, that X is primary over a point $\{M\}$ (known as the *total mass* of the system), and secondly, when $M \neq 0$, that X is also primary over a coadjoint orbit \mathbf{R}^6 of the normal Heisenberg subgroup $N \subset G$ defined by $A = \mathrm{id}$, e = 0. Applying the results (1.9), (1.21iii) and (3.1), we obtain after computing the moment map ϕ :

1) X is symplectomorphic to the product of $U = (\mathbf{R}^6, Md\mathbf{V} \wedge d\mathbf{R})$ by a symplectic manifold V on which N acts trivially. This is the original barycentric decomposition theorem, first proved in this generality in [S70, Thm 13.15], but which in essence goes back at least to Euler [E50, p. 187]:

"Quelque composé que soit le mouvement [x] d'un corps solide, on le peut toujours décomposer en un mouvement progressif [u] et en un mouvement de rotation [v]. Le premier s'estime par le mouvement du centre de gravité du corps, et il est toujours permis de considèrer ce mouvement séparément et indépendamment de l'autre mouvement."

2) The group $G \times SO(3) \times \mathbf{R}$ acts symplectically on $X = U \times V$ and the moment map $\Phi : X \to \mathfrak{g}^*$ splits as $\Phi(u,v) = M(\mathbf{R} \times \mathbf{V}, \mathbf{R}, \mathbf{V}, \frac{1}{2} ||\mathbf{V}||^2, 1) + (\mathbf{L}', 0, 0, \mathbf{E}', 0)$ where \mathbf{L}' and \mathbf{E}' , known as the *proper* angular momentum and energy, depend only on v. This is the classical Kœnig theorem [S70, 13.35], thus named by Painlevé [P30] but going back to Laplace for \mathbf{L} [L99, p. 70] and to Kœnig for \mathbf{E} [K51, p. 173]:

"Vis viva [2E], quæ inest corporibus A et B, æqualis est Vi viva communi [M||V||²], una cum Viribus vivis propriis amborum corporum [2E']."

(4.4) Remark. [S71] gives another application of the case where N is Heisenberg, to the description of a particle in a constant electromagnetic field.

(4.5) A primary orbit which doesn't split. Let G be the solvable group of all upper triangular matrices of the form

(4.6)
$$g = \begin{pmatrix} 1 & c & 0 & e & f \\ 1 & 0 & 0 & e \\ & e^{ja} & 0 & b \\ & & 1 & a \\ & & & 1 \end{pmatrix} \qquad a, c, e, f \in \mathbf{R}$$
$$b \in \mathbf{C}$$
$$j = 2\pi\sqrt{-1},$$

and let N be the normal subgroup in which c = e = 0. We identify \mathfrak{g}^* with $\mathbf{R} \times \mathbf{C} \times \mathbf{R}^3$ by writing (p, z, r, s, t) for the value of the 1-form

$$(4.7) pda + \operatorname{Re}(\bar{z}db) - rdc - sde - tdf$$

at the identity; likewise we identify \mathfrak{n}^* with triples (p, z, t) so that the projection $\mathfrak{g}^* \to \mathfrak{n}^*$ writes $(p, z, r, s, t) \mapsto (p, z, t)$. The coadjoint action of G preserves the hyperplane t = 1 and writes there

(4.8)
$$g \begin{pmatrix} p \\ z \\ r \\ s \\ 1 \end{pmatrix} = \begin{pmatrix} p + e + \operatorname{Re}(\overline{jb} e^{ja} z) \\ e^{ja} z \\ r + e \\ s + a - c \\ 1 \end{pmatrix}$$

in \mathfrak{g}^* , and likewise in \mathfrak{n}^* with the rows for r, s erased. Taking $\check{c} = (0, 1, 0, 0, 1)$, it is clear from this that the coadjoint orbit $X = G(\check{c})$ is N-primary over U with fiber V, where

(4.9)
$$X = \{ (p, e^{jq}, r, s, 1) : p, q, r, s \in \mathbf{R} \},\$$

(4.10)
$$U = \{(p, e^{jq}, 1) : p, q \in \mathbf{R}\}, \qquad \omega_U = dp \wedge dq,$$

(4.11)
$$V = \{ (0, 1, r, s, 1) : r, s \in \mathbf{R} \}, \qquad \omega_{V} = dr \wedge ds.$$

We claim that despite first appearances, X splits neither as $U \times V$ nor in fact in any other way. A first hint of this is the observation that its 2-form works out as $\omega_X = dp \wedge dq + dq \wedge dr + dr \wedge ds$ which is not the sum of ω_U and ω_V . For an actual proof we note that since the fiber V is connected it is enough to see that Γ , or equivalently the stabilizer K, acts nontrivially on V (1.23). But this is clear since one finds

(4.12)
$$K = \left\{ k = \begin{pmatrix} 1 & 0 & 0 & 0 & f \\ 1 & 0 & 0 & 0 & 0 \\ 1 & 1 & 0 & b \\ 1 & 1 & 1 \end{pmatrix} : \quad a \in \mathbf{Z} \\ b, f \in \mathbf{R} \right\}$$

(hence $\Gamma = \mathbf{Z}$) and, from (4.8), k(0, 1, r, s, 1) = (0, 1, r, s + a, 1).

(4.13) Remarks. (i) As the reader may verify, this orbit also provides an example of nonvanishing Mackey obstruction (3.8), viz. the cohomology class of

(4.14)
$$\theta_{\tilde{\mathbf{U}}}(g, l_{45}) = (0, 0, -e, c - l_{45}, 0).$$

Here we have identified the group \tilde{G} (2.8) with the direct product $G \times \mathbf{Z}$ by sending [n, l] to the pair $(nl, \text{entry } l_{45} \text{ of } l)$.

(ii) It is of interest to note that the representations attached to X are not primary when restricted to N. In fact, [**Z96**, Prop. 3.8] applied to the normal subgroup P^o in which a = e = 0 shows that we have $X = \operatorname{Ind}_{P}^{G}\{\check{c}_{|\mathfrak{p}}\}$ where P is the subgroup in which $a \in \mathbf{Z}$, e = 0. Accordingly one finds that the representations attached to X by [A71] are the $\operatorname{Ind}_{P}^{G}\chi_{w}$ where

(4.15)
$$\chi_w \begin{pmatrix} 1 & c & 0 & 0 & f \\ 1 & 0 & 0 & 0 \\ 1 & 0 & b \\ 1 & a \\ 1 & 1 \end{pmatrix} = w^a e^{i \operatorname{Re}(b)} e^{-if}, \qquad w = e^{i\theta} \in S^1,$$

is the most general character of P with differential $i\check{c}_{|\mathfrak{p}}$. These representations can be realized in $L^2(G/P) = L^2(S^1 \times \mathbf{R})$ by

(4.16)
$$(gE)(z,r) = e^{i(\theta+r)a}e^{i\operatorname{Re}(\bar{z}b)}e^{-ic(r-e)}e^{-if}E(ze^{-ja},r-e).$$

Likewise one finds that the representations attached to U are the $\operatorname{Ind}_{N\cap P}^{N} \eta_w$ where η_w is the restriction of χ_w , which we can realize in $L^2(S^1)$ by

(4.17)
$$(nF)(z) = e^{i\theta a} e^{i\operatorname{Re}(\bar{z}b)} e^{-if} F(ze^{-ja}).$$

Now it is clear that on restricting (4.16) to $N = \{g : c = e = 0\}$ one gets not one of the representations (4.17) but a direct integral of all of them.

(4.18) A primary orbit which splits nontrivially. Let H be the subgroup of (4.6) in which c = 0. Its coadjoint action is obtained by making c = 0 and erasing the row for r in (4.8). In particular we see that the orbit

(4.19)
$$Y = H(0, 1, 0, 1) = \{(p, e^{jq}, q, 1) : p, q \in \mathbf{R}\}$$

is N-primary over the same orbit (4.10) of the same normal subgroup N, and is its universal covering.

- (4.20) Remark. As in (4.13ii) one can verify that the representations attached to Y are not primary when restricted to N. This raises the question of a possible correlation between an irreducible representation being N-primary and the corresponding orbit being *trivially split* N-primary (1.4).
- (4.21) A primary space which doesn't split even topologically. All our examples so far were at least homeomorphic to the product of a covering of U with another manifold, so it is natural to ask how far from such a product a primary space can get. In this section we exploit the close similarity of (1.8) with the Kodaira-Thurston construction [T76; W77, p. 10] to show that, at least as long as we don't insist on homogeneity under an ambient Lie group G acting in hamiltonian fashion, examples exist with no such homeomorphism.

To this end, we let N and U be same group and cylinder orbit as in (4.10), so that $\tilde{\mathbf{U}}$ and Γ are respectively the plane \mathbf{R}^2 and the integers \mathbf{Z} . We take for V the flat torus $\mathbf{T}^2 = \{(\mathbf{e}^{jr}, \mathbf{e}^{js}) : r, s \in \mathbf{R}\}$ where the integers act by powers of a Dehn twist: $k(\mathbf{R}, \mathbf{S}) = (\mathbf{R}\mathbf{S}^k, \mathbf{S})$, and we form the associated

bundle $X = \tilde{U} \times_{\Gamma} V$. Thus X is the quotient of $(\mathbf{R}^2 \times \mathbf{T}^2, dp \wedge dq + dr \wedge ds)$ by the **Z**-action

(4.22)
$$k(p, q, e^{jr}, e^{js}) = (p, q + k, e^{j(r+ks)}, e^{js}).$$

Our claim is that X is not homeomorphic to the product of any covering of U by any surface. To see this, we observe that $\mathbf{R}^2 \times \mathbf{T}^2$ can be regarded as the right quotient of the group M of real matrices

$$(4.23) \quad m = \begin{pmatrix} 1 & 0 & 0 & p \\ & 1 & r & s \\ & & 1 & q \\ & & & 1 \end{pmatrix} \text{ by the subgroup } \mathbf{Z} \oplus \mathbf{Z} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ & 1 & \mathbf{Z} & \mathbf{Z} \\ & & 1 & 0 \\ & & & 1 \end{pmatrix}.$$

Moreover (4.22) coincides with the residual right action on $M/(\mathbf{Z} \oplus \mathbf{Z})$ of the Heisenberg group

(4.24)
$$\mathbf{H}_{\mathbf{Z}} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ & 1 & \mathbf{Z} & \mathbf{Z} \\ & & 1 & \mathbf{Z} \\ & & & 1 \end{pmatrix}$$

which normalizes $\mathbf{Z} \oplus \mathbf{Z}$. Thus we have $X = M/H_{\mathbf{Z}}$ and therefore $\pi_1(X) = H_{\mathbf{Z}}$. Now since $H_{\mathbf{Z}}$ does not admit \mathbf{Z} as a direct factor, it follows that X cannot be the product of the cylinder U (or any finite covering) by anything. Nor is it the product of the plane \tilde{U} by any surface, because no surface has fundamental group $H_{\mathbf{Z}}$ [M77, pp. 135, 143].

References

- [A71] Louis Auslander and Bertram Kostant, *Polarization and unitary representations of solvable Lie groups*. Invent. Math. **14** (1971) 255–354.
- [B72] Pierre Bernat et al., Représentations des groupes de Lie résolubles. Dunod, Paris, 1972.
- [B87] Arthur L. Besse, *Einstein manifolds*. Springer-Verlag, Berlin, 1987.
- [B67] Nicolas Bourbaki, Variétés différentielles et analytiques. Fascicule de résultats. Paragraphes 1 à 7. Hermann, Paris, 1967.
- [B80] Nicolas Bourbaki, Algèbre. Chapitre 10. Algèbre homologique. Masson, Paris, 1980.
- [C99] Ana Cannas da Silva and Alan Weinstein, Geometric models for noncommutative algebras. Berkeley Mathematics Lecture Notes, vol. 10. American Mathematical Society, Providence, RI, 1999.
- [C37] Alfred H. Clifford, Representations induced in an invariant subgroup. Ann. of Math. (2) 38 (1937) 533–550.

- [C93] David H. Collingwood and William M. McGovern, *Nilpotent orbits in semisimple Lie algebras*. Van Nostrand Reinhold Co., New York, 1993.
- [D70] Michel Duflo, Sur les représentations unitaires des groupes de Lie contenant un sous-groupe invariant nilpotent. C. R. Acad. Sci. Paris Sér. A 270 (1970) 578–581.
- [E50] Leonhard Euler, Découverte d'un nouveau principe de Mecanique. Hist. Acad. Roy. Berlin 6 (1750) 185–217.
- [F88] J. M. G. Fell and Robert S. Doran, Representations of *-algebras, locally compact groups, and Banach *-algebraic bundles. Academic Press, Boston, 1988.
- [G82] Victor Guillemin and Shlomo Sternberg, Geometric quantization and multiplicities of group representations. Invent. Math. 67 (1982) 515–538.
- [G84a] Victor Guillemin and Shlomo Sternberg, *Multiplicity-free spaces*. J. Differential Geom. **19** (1984) 31–56.
- [G84b] Victor Guillemin and Shlomo Sternberg, Symplectic techniques in physics. Cambridge University Press, Cambridge, 1984.
- [H63] Edwin Hewitt and Kenneth A. Ross, *Abstract harmonic analysis*, vol. 1. Springer-Verlag, Berlin, 1963.
- [184] Patrick Iglesias, Un théorème de décomposition, relatif à une classe de systèmes dynamiques isolés. J. Geom. Phys. 1 (1984) 37–44.
- [I91] Patrick Iglesias, Les SO(3)-variétés symplectiques et leur classification en dimension 4. Bull. Soc. Math. France 119 (1991) 371–396.
- [K78] David A. Kazhdan, Bertram Kostant, and Shlomo Sternberg, Hamiltonian group actions and dynamical systems of Calogero type. Comm. Pure Appl. Math. 31 (1978) 481–507.
- [K92] Donald R. King, The component groups of nilpotents in exceptional simple real Lie algebras. Comm. Algebra 20 (1992) 219–284.
- [K62] Aleksandr A. Kirillov, Unitary representations of nilpotent Lie groups. Uspehi Mat. Nauk 17 (1962) 57–110.
- [K51] Johann Samuel Kœnig, De universali principio æquilibrii et motus, in vi viva reperto, deque nexu inter vim vivam et actionem, utriusque minimo, dissertatio. Nova Acta Eruditorum (1751) 125–135, 162–176.
- [K68] Bertram Kostant, Symplectic manifolds and unitary representations. MIT Lecture Notes, 1968.
- [K70] Bertram Kostant, Quantization and unitary representations. I. Prequantization. Lecture Notes in Math. 170 (1970) 87–208.
- [L06] Nicolaas P. Landsman, Lie groupoids and Lie algebroids in physics and noncommutative geometry. J. Geom. Phys. **56** (2006) 24–54.

- [L99] Pierre-Simon de Laplace, *Traité de mécanique céleste*, vol. 1 (Books 1 and 2). Crapelet, Paris, 1799.
- [L92] Wojciech Lisiecki, A symplectic analog of Mackey's theory of representations of group extensions. In: Tomasz Mazur (ed.) Classical analysis (Kazimierz Dolny, 1991), pp. 124–131. World Scientific, River Edge, NJ, 1992.
- [L94] Zongyi Li, The Mackey obstruction and the coadjoint orbits. Trans. Amer. Math. Soc. 346 (1994) 693–705.
- [M58] George W. Mackey, Unitary representations of group extensions. I. Acta Math. 99 (1958) 265–311.
- [M76] Charles-Michel Marle, Symplectic manifolds, dynamical groups, and Hamiltonian mechanics. In: Michel Cahen and Moshé Flato (eds.) Differential geometry and relativity, pp. 249–269. Reidel, Dordrecht, 1976.
- [M74] Jerrold E. Marsden and Alan Weinstein, Reduction of symplectic manifolds with symmetry. Rep. Math. Phys. 5 (1974) 121–130.
- [M77] William S. Massey, Algebraic topology: An introduction. Springer-Verlag, New York, 1977.
- [O84] Anatol Odzijewicz, The reduction of symplectic structure and Sternberg construction. Ann. Inst. H. Poincaré Phys. Théor. 40 (1984) 361–370.
- [O04] Juan-Pablo Ortega and Tudor S. Raţiu, Momentum maps and Hamiltonian reduction. Birkhäuser, Boston, 2004.
- [P30] Paul Painlevé, Cours de mécanique. Professé à l'École polytechnique, vol. 1. Gauthier-Villars, Paris, 1930.
- [P78] Lajos Pukánszky, Unitary representations of Lie groups with cocompact radical and applications. Trans. Amer. Math. Soc. 236 (1978) 1–49.
- [S70] Jean-Marie Souriau, Structure des systèmes dynamiques. Dunod, Paris, 1970. (Reprint: Éditions Jacques Gabay, Sceaux, 2008. Translation: Structure of dynamical systems: A symplectic view of physics. Birkhäuser, Boston, 1997.)
- [S71] Jean-Marie Souriau, Variétés symplectiques et cohomologie en mécanique.
 Publ. Dép. Math. Univ. Lyon 1 8 (Suppl. 4) (1971) 1–11.
- [T54] Robert L. Taylor, Covering groups of nonconnected topological groups. Proc. Amer. Math. Soc. 5 (1954) 753–768.
- [T76] William P. Thurston, Some simple examples of symplectic manifolds. Proc. Amer. Math. Soc. **55** (1976) 467–468.
- [W77] Alan Weinstein, Lectures on symplectic manifolds. Regional Conference Series in Mathematics, vol. 29. Amer. Math. Soc., Providence, RI, 1977.

[Z96] François Ziegler, *Théorie de Mackey symplectique*. In: Méthode des Orbites et Représentations Quantiques, pp. 69–96. Ph. D. Thesis, Université de Provence, 1996. arXiv:1011.5056.

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