# TWO-DIMENSIONAL OPTIMIZATION PROBLEM OF PLANT LOCATION

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ABSTRACT. In this paper, the following matters are presented: the adjoint problem of the two-dimensional matter propagation problem; the algorithm for determination of a domain in which a plant can be located so that the values of the pollution-level reflecting functional does not exceed a given value at considered sensitive areas; application of this algorithm for numerical experiments to a typical problem.

# 1. Equation of the suspended matter propagation and its adjoint equation (see [1])

The equation of the suspended matter propagation, i.e. the matter transport and diffusion equation in the horizontal 2D case has the following form:

$$\frac{\partial C}{\partial t} + u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} + \sigma C = f + \gamma \Delta C, \quad (x, y) \in G, \ 0 \le t \le T$$
 (1.1)

with the initial and boundary conditions:

$$C\big|_{t=0} = C^0, \quad C\big|_{\Gamma^-} = \varphi, \quad \frac{\partial C}{\partial n}\big|_{\Gamma^+} = 0,$$
 (1.2)

where x, y, t are the space and time variables; C is the matter concentration;  $\sigma$  is the decay coefficient; f is the source intensity;  $\gamma$  is the diffusion coefficient; u, v are respectively velocity components in the x and y directions, and satisfy the following equation:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0, (1.3)$$

 $\Gamma = \Gamma^+ + \Gamma^-$  with  $\Gamma^+$  is the boundary part, at which  $u_n \geq 0$ ;  $\Gamma^-$  is the boundary part, at which  $u_n < 0$ ;  $\Delta = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}$ -Laplace operator;  $u_n$  is the projection of the velocity on the external normal vector  $\vec{n}$ .

Using (1.3), the equation (1.1) can be rewritten as follows:

$$\frac{\partial C}{\partial t} + \frac{\partial uC}{\partial x} + \frac{\partial vC}{\partial y} + \sigma C = f + \gamma \Delta C. \tag{1.4}$$

Solution of the equation (1.1) may be determined under the form:  $C = C_1 + C_2$  where,  $C_1$  and  $C_2$  are the solutions of the following equation:

$$\frac{\partial C_1}{\partial t} + u \frac{\partial C_1}{\partial x} + v \frac{\partial C_1}{\partial y} + \sigma C_1 = \gamma \Delta C_1 \tag{1.5}$$

with the initial and boundary conditions:

$$C_1\big|_{t=0} = C^0, \quad C_1\big|_{\Gamma^-} = \varphi, \quad \frac{\partial C_1}{\partial n}\big|_{\Gamma^+} = 0$$

and

$$\frac{\partial C_2}{\partial t} + \frac{\partial uC_2}{\partial x} + \frac{\partial vC_2}{\partial y} + \sigma C_2 = \gamma \Delta C_2 + f \tag{1.6}$$

with the initial and boundary conditions:

$$C_2|_{t=0} = 0, \quad C_2|_{\Gamma^-} = 0, \quad \frac{\partial C_2}{\partial n}|_{\Gamma^+} = 0.$$
 (1.7)

We now establish the adjoint equation of the equation (1.6). By multiplying both sides of the equation (1.6) by a some function  $C_2^*$  and integrating the equation obtained on the area  $G \times [0, T]$ , we get:

$$\int_{0}^{T} dt \int_{G} C_{2}^{*} \frac{\partial C_{2}}{\partial t} dG + \int_{0}^{T} dt \int_{G} C_{2}^{*} \operatorname{div}(\vec{u}C_{2}) dG 
+ \int_{0}^{T} dt \int_{G} \sigma C_{2} C_{2}^{*} dG - \gamma \int_{0}^{T} dt \int_{G} C_{2}^{*} \Delta C_{2} dG = \int_{0}^{T} dt \int_{G} C_{2}^{*} f dG.$$
(1.8)

Let  $\gamma = \text{const}$ , using the partial integration technique, the Green formula and the condition (1.3), we have:

$$\int_{0}^{T} dt \int_{G} C_{2}^{*} \frac{\partial C_{2}}{\partial t} dG = \int_{G} C_{2}^{*} C_{2} \Big|_{0}^{T} dG - \int_{0}^{T} dt \int_{G} C_{2} \frac{\partial C_{2}^{*}}{\partial t} dG,$$

$$\int_{0}^{t} dt \int_{G} C_{2}^{*} \operatorname{div}(\vec{u}C_{2}) dG = \int_{0}^{T} dt \int_{\Gamma} u_{n} C_{2}^{*} C_{2} d\Gamma - \int_{0}^{T} dt \int_{G} C_{2} \operatorname{div}(\vec{u}C_{2}^{*}) dG,$$

$$\gamma \int_{0}^{T} dt \int_{G} C_{2}^{*} \Delta C_{2} dG = \gamma \int_{0}^{T} dt \int_{\Gamma} \left( C_{2}^{*} \frac{\partial C_{2}}{\partial n} - C_{2} \frac{\partial C_{2}^{*}}{\partial n} \right) d\Gamma + \gamma \int_{0}^{T} dt \int_{G} C_{2} \Delta C_{2}^{*} dG.$$

Putting these expressions into (1.8), one deduces:

$$\int_{0}^{T} dt \int_{G} C_{2} \left( -\frac{\partial C_{2}^{*}}{\partial t} - \operatorname{div}(\vec{u}C_{2}^{*}) + \sigma C_{2}^{*} - \gamma \Delta C_{2}^{*} \right) dG =$$

$$= \int_{0}^{T} dt \int_{G} C_{2}^{*} f dG - \int_{G} C_{2}^{*} C_{2}|_{t=T} dG + \int_{G} C_{2}^{*} C_{2}|_{t=0} dG - \int_{0}^{T} dt \int_{\Gamma} u_{n} C_{2} C_{2}^{*} d\Gamma$$

$$+ \gamma \int_{0}^{T} dt \int_{\Gamma} \left( C_{2}^{*} \frac{\partial C_{2}}{\partial n} - C_{2} \frac{\partial C_{2}^{*}}{\partial n} \right) d\Gamma. \tag{1.9}$$

Let the function  $C_2^*$  satisfy the following equation:

$$-\frac{\partial C_2^*}{\partial t} - \operatorname{div}(\vec{u}C_2^*) + \sigma C_2^* - \gamma \Delta C_2^* = p.$$
 (1.10)

From the initial and the boundary conditions (1.7), one yields;

$$\begin{split} &\int_{G} C_{2}C_{2}^{*}\big|_{t=0}dG = 0, \\ &\int_{G}^{T} dt \int_{\Gamma} u_{n}C_{2}C_{2}^{*}d\Gamma = \int_{0}^{T} dt \int_{\Gamma^{+}} u_{n}C_{2}C_{2}^{*}d\Gamma, \\ &\gamma \int_{0}^{T} dt \int_{\Gamma} \left(C_{2}^{*} \frac{\partial C_{2}}{\partial n} - C_{2} \frac{\partial C_{2}^{*}}{\partial n}\right) d\Gamma = \gamma \int_{0}^{T} dt \int_{\Gamma^{+}} \left( - C_{2} \frac{\partial C_{2}^{*}}{\partial n}\right) d\Gamma + \gamma \int_{0}^{T} dt \int_{\Gamma^{-}} C_{2}^{*} \frac{\partial C_{2}}{\partial n} d\Gamma. \end{split}$$

From the above expressions and (1.10), the equation (1.9) can be rewritten under the form:

$$\int_{0}^{T} dt \int_{G} pC_{2}dG = \int_{0}^{T} dt \int_{G} fC_{2}^{*}dG + \int_{G} C_{2}C_{2}^{*}\big|_{t=T}dG$$

$$+ \gamma \int_{0}^{T} dt \int_{\Gamma^{-}} C_{2}^{*} \frac{\partial C_{2}}{\partial n} d\Gamma - \int_{0}^{T} dt \int_{\Gamma^{+}} C_{2} \left(\gamma \frac{\partial C_{2}^{*}}{\partial n} + u_{n}C_{2}^{*}\right) d\Gamma. \quad (1.11)$$

Let the initial and boundary conditions of the equation (1.10) be chosen as follows:

$$C_2^*\big|_{t=T} = 0, \quad C_2^*\big|_{\Gamma^-} = 0, \quad \left(\gamma \frac{\partial C_2^*}{\partial n} + u_n C_2^*\right)\big|_{\Gamma^+} = 0.$$
 (1.12)

Then, from (1.11) and (1.12) we get the dual form:

$$\int_{0}^{T} dt \int_{G} pC_{2}dG = \int_{0}^{T} dt \int_{G} fC_{2}^{*}dG.$$
 (1.13)

It is easy to verify that the problem (1.10), (1.12) is the adjoint problem of the (1.6), (1.7). Indeed, with the notation:

$$A = \frac{\partial}{\partial t} + \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \sigma - \gamma \Delta \quad A^* = -\frac{\partial}{\partial t} - \frac{\partial u}{\partial x} - \frac{\partial v}{\partial y} + \sigma - \gamma \Delta,$$

we have:

$$AC_2 = f, \quad A^*C_2^* = p,$$

$$(AC_2, C_2^*) = (f, C_2^*) = \int_0^T dt \int_G fC_2^* dG = \int_0^T dt \int_G pC_2 dG = (C_2, p) = (C_2, A^*C_2^*).$$

Use of the variable transformation  $t_1 = T - t$ , the equation (1.10) becomes:

$$\frac{\partial C_2^*}{\partial t_1} - \operatorname{div}(\vec{u}C_2^*) + \sigma C_2^* - \gamma \Delta C_2^* = p, 
C_2^*|_{t_1=0} = 0, \quad C_2^*|_{\Gamma^-} = 0, \quad \left(\gamma \frac{\partial C_2^*}{\partial r_1} + u_n C_2^*\right)|_{\Gamma^+} = 0.$$
(1.14)

For simplicity, by using (1.3), we obtain an another form of the adjoint equation (1.14):

$$\frac{\partial C_2^*}{\partial t_1} - u \frac{\partial C_2^*}{\partial x} - v \frac{\partial C_2^*}{\partial y} + \sigma C_2^* - \gamma \Delta C_2^* = p. \tag{1.15}$$

# 2. Pollution-level reflecting functionals (see [1])

Assume that the suspended matter concentration C is calculated from the equation (1.1). We consider the following functionals:

- a. The time-averaged amount of the matter concentration C on a sensitive area  $G_k \subset G$  for the period T:  $J_k^A = \frac{1}{T} \int\limits_0^T dt \int\limits_{G_k} CdG$ .
- b. The total amount of settling matter in the same area  $G_k \subset G$ :  $J_k^B = \int_0^1 dt \int_{G_k} aCdG$ , where, the constant a represents portion of matter which settles down, that are mainly the heavy matters and partly the suspended matters settling down by downward diffusion.

### c. Generalized functional:

$$J_k = \int_0^T dt \int_{G_k} pCdG \quad \text{where} \quad p = \begin{cases} \frac{1}{T} + a, & (x,y) \in G_k \\ 0, & (x,y) \notin G_k \end{cases}$$
 (2.1)

and p is a function referring to the economic, sanitary, ecological, health standards and so on.

### d. Global functional:

$$Y_p = \int_0^T dt \int_G pCdG \quad \text{where,} \quad p = \begin{cases} \frac{1}{T} + G_k, & (x,y) \in G_k, \ k = 1, 2, \dots, m \\ 0, & (x,y) \notin \bigcup_{k=1}^m G_k. \end{cases}$$

# 3. Optimization problem of plant location (see [1])

Let  $G_k$  (k = 1, 2, ..., m) be considered areas, recreation zones or other environmentally sensitive areas on the region G. Our problem is to determine the domain  $\Omega_k \subset G$  so that the pollution matter from a plant located in this domain  $\Omega_k$  satisfies the following condition for the sensitive area  $G_k$ :

$$Y_k = \int_0^T dt \int_{G_k} pCdG \le \overline{c}_k, \quad \text{where} \quad p = \begin{cases} \frac{1}{T} + a_k, & (x,y) \in G_k \\ 0, & (x,y) \notin G_k \end{cases}$$
(3.1)

and  $\overline{c}_k$  is a given figure.

If the determination of domain  $\Omega_k$  is impossible on the G, the reduction of rate of the pollution emission Q, will make the determination of the plant location possible.

Assume that on the region G there are m sensitive areas  $G_k$  (k = 1, ..., m) and the source of matter emission is located at a point  $r_0 = (x_0, y_0)$ . Then, the source intensity can be described by the function:  $f(x, y) = Q\delta(r - r_0)$ , Q = const

where,  $\delta(r) = \begin{cases} \infty, & r = r_0 \\ 0, & r \neq r_0 \end{cases}$  is Dirac function, and from (1.1), we get:

$$\frac{\partial C}{\partial t} + u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} + \sigma C = Q\delta(r - r_0) + \gamma \Delta C$$

with the conditions:  $C\big|_{t=0} = C^0$ ,  $C\big|_{\Gamma^-} = \varphi$ ,  $\frac{\partial C}{\partial n}\big|_{\Gamma^+} = 0$ .

In order to determine the domain  $\Omega$ , in which the plant can be located so that in all sensitive areas  $G_k$ , the generalized functional  $Y_k$  satisfies the condition (3.1), we do as follows:

**a.** Calculation of concentration C from the equation (1.5):

$$\frac{\partial C}{\partial t} + u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} + \sigma C = \gamma \Delta C$$

with the initial and boundary conditions  $C\big|_{t=0} = C^0$ ,  $C\big|_{\Gamma^-} = \varphi$ ,  $\frac{\partial C}{\partial n}\big|_{\Gamma^+} = 0$  and generalized functional (2.1):  $J_k = \int\limits_0^T dt \int\limits_{G_k} pCdG = \widetilde{c}_k$ .

**b.** Solving *m* adjoint equations:

$$\frac{\partial C_k^*}{\partial t_1} - u \frac{\partial C_k^*}{\partial x} - v \frac{\partial C_k^*}{\partial y} + \sigma C_k^* - \gamma \Delta C_k^* = p_k$$

where,  $p_k = \begin{cases} \frac{1}{T} + a_k, & (x, y) \in G_k \\ 0, & (x, y) \notin G_k \end{cases}$  with the conditions:

$$C_k^*\big|_{t_1=0} = 0, \quad C_k^*\big|_{\Gamma^-} = 0, \quad \left(v\frac{\partial C_k^*}{\partial n} + u_n C_k^*\right)\Big|_{\Gamma^+} = 0,$$

we obtain the solutions  $C_k^*$  (k = 1, 2, ..., m).

From the dual form (1.13), we get:

$$Y_{k}^{*} = \int_{0}^{T} dt \int_{G} p_{k} C dG = \int_{0}^{T} dt \int_{G} Q \delta(r - r_{0}) C_{k}^{*} dG$$
$$= \int_{0}^{T} Q C_{k}^{*}(r_{0}, t) dt = \int_{0}^{T} Q C_{k}^{*}(r_{0}, T - t_{1}) dt_{1}$$

which must satisfy the condition:  $Y_k^* \leq \overline{c}_k - \widetilde{c}_k = \overline{\overline{c}}_k$ .

Now we consider the function:  $Y_k^*(r) = Q \int_0^r C_k^*(r,t) dt$  and draw the iso-grams of  $Y_k^*(r) = \text{const.}$  Then,  $\Omega_k$  in which the functional  $Y_k^*(r) \leq \overline{c}_k$  are found out. If there is perchance no area  $\Omega_k$  inside G, it may be re-established anyway by reducing the discharge intensity Q.

c. Overlaying all the areas  $\Omega_k$   $(k=1,\ldots,m)$ , we obtain the domain  $\Omega$ ,  $(\Omega = \bigcap_{k=1}^m \Omega_k)$ .  $\Omega$  will be the domain in which the plant can be located so that pollution standards will be met in all the areas  $G_k \subset G$ ,  $(k=1,2,\ldots,m)$ .

## 4. Algorithm (see [2]-[4])

The equation (1.5) and the adjoint equation (1.15) may be rewritten in a common form:

$$\frac{\partial C}{\partial t} + \Lambda C = f \tag{4.1}$$

where, 
$$\Lambda = \Lambda_1 + \Lambda_2$$
,  $\Lambda_1 = \pm u \frac{\partial}{\partial x} - \gamma \frac{\partial^2}{\partial x^2} + \frac{\sigma}{2}$ ,  $\Lambda_2 = \pm v \frac{\partial}{\partial y} - \gamma \frac{\partial^2}{\partial y^2} + \frac{\sigma}{2}$ .

Equation (4.1) is solved by the method of the directional decomposition (splitting method):

$$\frac{C^{k+1} - C^k}{dt} + \Lambda \left[ \theta C^{k+1} + (1 - \theta) C^k \right] = f^{k+1}$$
or  $(I + dt\theta \Lambda) C^{k+1} = \left[ I - dt (1 - \theta) \Lambda \right] C^k + dt f^{k+1},$  (4.2)

where  $0 \le \theta \le 1$ , I is the unique operator.

Using approximation:

$$[I + dt\theta(\Lambda_1 + \Lambda_2)] = (I + dt\theta\Lambda_1)(I + dt\theta\Lambda_2) + 0(dt^2)$$

from (4.2), one deduces:

$$(I + dt\theta\Lambda_1)(I + dt\theta\Lambda_2)C^{k+1} = dtf^{k+1} + [I - dt(1 - \theta)\Lambda]C^k.$$

The computational process contains two steps:

$$(I + dt\theta \Lambda_1)C^{k+1/2} = [I - dt(1-\theta)\Lambda]C^k + dtf^{k+1}$$
(4.3)

$$(I + dt\theta \Lambda_2)C^{k+1} = C^{k+1/2}. (4.4)$$

a. Discretizing the equation (4.3) by an implicit finite difference scheme:

$$\begin{split} &\Lambda_1 C^{k+1/2} = \frac{\left(\pm u + |u|\right)_{m,n}^{k+1/2}}{2} \frac{\left(C_{m,n}^{k+1/2} - C_{m-1,n}^{k+1/2}\right)}{\Delta x} \\ &+ \frac{\left(\pm u - |u|\right)_{m,n}^{k+1/2}}{2} \frac{\left(C_{m+1,n}^{k+1/2} - C_{m,n}^{k+1/2}\right)}{\Delta x} - \gamma \frac{\left(C_{m+1,n}^{k+1/2} - 2C_{m,n}^{k+1/2} + C_{m-1,n}^{k+1/2}\right)}{\Delta x^2} + \frac{\sigma}{2} \,, \\ &\Lambda C^k = \pm u_{m,n}^{k+1/2} \frac{\left(C_{m+1,n}^k - C_{m-1,n}^k\right)}{2\Delta x} - \gamma \frac{C_{m+1,n}^k - 2C_{m,n}^k + C_{m-1,n}^k}{\Delta x^2} \\ &\pm v_{m,n}^{k+1/2} \frac{\left(C_{m,n+1}^k - C_{m,n-1}^k\right)}{2\Delta y} - \gamma \frac{C_{m,n+1}^k - 2C_{m,n}^k + C_{m,n-1}^k}{\Delta y^2} + \sigma, \end{split}$$

we obtain:

$$a_m C_{m+1,n}^{k+1/2} + b_m C_{m,n}^{k+1/2} + c_m C_{m-1,n}^{k+1/2} = d_m, (4.5)$$

where,

$$a_{m} = \frac{\left(\pm u - |u|\right)_{m,n}^{k+1/2} \theta dt}{2\Delta x} - \frac{\gamma \theta dt}{(\Delta x)^{2}}, \quad b_{m} = 1 + \frac{\theta |u|_{m,n}^{k+1/2} dt}{\Delta x} + 2\frac{\gamma \theta dt}{(\Delta x)^{2}} + \frac{\sigma dt}{2},$$

$$c_{m} = -\frac{\left(\pm u + |u|\right)_{m,n}^{k+1/2} \theta dt}{2\Delta x} - \frac{\gamma \theta dt}{(\Delta x)^{2}}, \quad d_{m} = dt f_{m,n}^{k+1} + [I - dt(1 - \theta)\Lambda] C_{m,n}^{k}.$$

It is easy to verify that:  $b_m > 0$ ,  $a_m < 0$ ,  $c_m < 0$  and  $|b_m| \ge |a_m| + |c_m| + \delta$ ,  $\delta > 0$ .

So, the linear equation system (4.5) has the unique solution and the computational error of the following double sweep method

$$C_{m,n}^{k+1} = L_m C_{m+1,n}^{k+1} + K_m, (4.6)$$

where, 
$$L_m = \frac{-a_m}{b_m + c_m L_{m-1}}$$
,  $K_m = \frac{d_m - c_m K_{m-1}}{b_m + c_m L_{m-1}}$ , is not accumulated (see [5]).

**b.** Discretizing the equation (4.4) by a difference scheme:

$$\Lambda_2 C^{k+1} = \frac{\left(\pm v + |v|\right)_{m,n}^{k+1} \left(C_{m,n}^{k+1} - C_{m,n-1}^{k+1}\right)}{2} + \frac{\left(\pm v - |v|\right)_{m,n}^{k+1} \left(C_{m,n+1}^{k+1} - C_{m,n}^{k+1}\right)}{2} - \gamma \frac{\left(C_{m,n+1}^{k+1} - 2C_{m,n}^{k+1} + C_{m,n-1}^{k+1}\right)}{\Delta y^2} + \frac{\sigma}{2} C_{m,n}^{k+1}$$

we also get:

$$\tilde{a}_n C_{m,n+1}^{k+1} + \tilde{b}_n C_{m,n}^{k+1} + \tilde{c}_n C_{m,n-1}^{k+1} = \tilde{d}_n, \tag{4.7}$$

where,

$$\tilde{a}_{m} = \frac{\left(\pm v - |v|\right)_{m,n}^{k+1} \theta dt}{2\Delta y} - \frac{\gamma \theta dt}{(\Delta y)^{2}}, \quad \tilde{b}_{m} = 1 + \frac{\theta |v|_{m,n}^{k+1} dt}{\Delta y} + 2\frac{\gamma \theta dt}{(\Delta y)^{2}} + \frac{dt\sigma}{2},$$

$$\tilde{c}_{m} = -\frac{\left(\pm v + |v|\right)_{m,n}^{k+1} dt}{2\Delta y} - \frac{\gamma \theta dt}{(\Delta y)^{2}}, \quad \tilde{d}_{m} = C_{m,n}^{k+1/2}.$$

Obviously:  $\tilde{b}_m > 0$ ,  $\tilde{a}_m < 0$ ,  $\tilde{c}_m < 0$  and  $|\tilde{b}_m| \ge |\tilde{a}_m| + |\tilde{c}_m| + \delta$ ,  $\delta > 0$ .

Also, the equation system (4.7) has the unique solution and the double sweep method (4.6) does not produce an accumulated computational error.

## 5. Numerical experiments

The mentioned-above algorithm is applied to solve the following optimization problem of plant location:

- The computed rectangular region  $G = 1000 \, m \times 1000 \, m$  is covered by a uniform grid  $51 \times 51$  with spacing steps:  $dx = 20 \, m$ ,  $dy = 20 \, m$ .
  - A constant velocity field (u, v):  $u = 0.5 \,\mathrm{m/s}, v = -0.5 \,\mathrm{m/s}.$
  - Diffusion coefficient :  $\gamma = 0.5 \,\mathrm{m}^2/\mathrm{s}$ .
  - Decay coefficient:  $\sigma = 0.0005 \,\mathrm{s}^{-1}$ .
  - Time step: dt = 5 s.
  - Time simulation:  $T = 20000 \,\mathrm{s}$ .
- 3 considered sensitive rectangular areas  $G_k$  inside G (k = 1, 2, 3) with the left-bottom corner coordinates and the right-top corner coordinates are as follows:
  - $+G_1 = [(24.5, 8.5), (25.5, 9.5)], +G_2 = [(37.5, 12.5), (39.5, 14.5)],$
  - $+G_3 = [(29.5, 33.5), (30.5, 34.5)].$
  - Standard concentration:  $\bar{c}_k = 10 \, mg/l \ (k = 1, 2, 3)$ .

The numerical results are illustrated in Fig. 1. In this figure, the figure on the

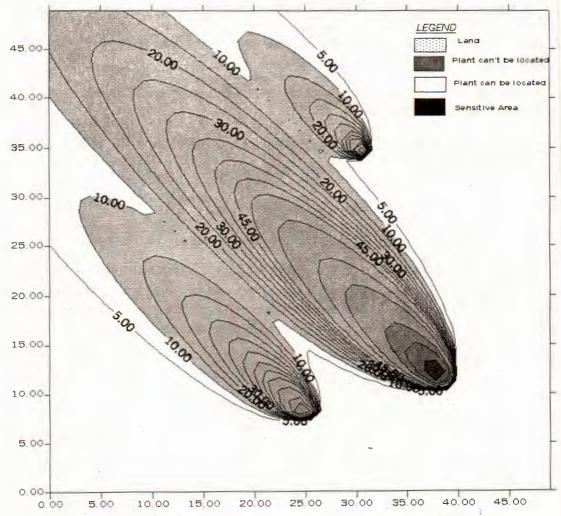


Fig. 1. Distribution of value of the pollution level-reflecting functionals  $Y_k^*$  for problem 1

contour lines indicates value of the pollution-level reflecting functionals  $Y_k^*$ . As a result, the domain  $\Omega$  where the plant can be located so that the sanitary condition in the all areas  $G_k$  are satisfied (that means  $Y_k^* \leq \overline{c}_k$ ) is in white.

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## BÀI TOÁN 2 CHIỀU TỐI ƯU XÁC ĐỊNH VỊ TRÍ NGUỒN THẢI

Bài báo trình bày các vấn đề sau: Bài toán liên hợp với bài toán lan truyền vật chất 2 chiều. Thuật toán xác định miền có thể đặt xí nghiệp sao cho phiếm hàm biểu thị mức độ ô nhiễm không vượt quá mức độ cho phép ở các vùng nhạy cảm quan tâm. Đã áp dụng thuật toán này để tính toán cho một bài toán mẫu.