Structure, dielectric and optical properties of Nd³⁺ doped LiTaO₃ transparent ferroelectric glass-ceramic nanocomposites

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Abstract

Here, we present the structural, dielectric and optical properties of neodymium ion (Nd³⁺) doped novel transparent glass-ceramics containing LiTaO₃ nanocrystals in the Li₂O-Ta₂O₅-SiO₂-Al₂O₃ (LTSA) glass system prepared by the melt quenching technique. The precursor glasses were isothermally crystallized at 680°C for 3-100 h, following the differential thermal analysis (DTA) data, to obtain nanostructured glass-ceramics. They were characterized by X-ray diffraction (XRD), field emission scanning electron microscope (FESEM), transmission electron microscopy (TEM), Fourier transform infrared reflection spectra (FTIRRS), optical absorption and luminescence spectroscopy along with dielectric constant measurements. XRD, FESEM, TEM and FTIRRS confirm the nanocrystallization of LiTaO₃ (14-36 nm) in the LTSA glass matrix. A steep increase in dielectric constant (ϵ_r) of glass-ceramics with heat-treatment time is observed due to high dielectric constant ferroelectric LiTaO₃ formation. The measured NIR photoluminescence spectra have exhibited emission transitions of ${}^4F_{3/2} \rightarrow {}^4I_J$ (J = 9/2, 11/2 and 13/2) from Nd³⁺ ions upon excitation at 809 nm. It is observed that the photoluminescent intensity and excited state (${}^4F_{3/2}$) lifetime of Nd³⁺ ions decrease with increase in heat-treatment

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time due to concentration quenching effect. The absorption spectra and fluorescence measurements reveal that the incorporation of Nd³+ ions in the LiTaO₃ crystal lattice in the oxide glassy matrix is important for obtaining desirable fluorescence performance of the material.

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Photoluminescence; Lifetime; Concentration quenching; Dielectric constant

1. Introduction

Lithium tantalate (LiTaO₃, LT) single crystal is one of the most important lead-free ferroelectric materials in the $A^{1+}B^{5+}O_3$ type perovskite family rhombohedral crystal structure with crystal symmetry class R3c (unit cell dimensions: a = 5.1530 Å and c = 13.755 Å), having large nonlinear constant ($d_{33} = 13.6 \text{ pm/V}$ at 1064 nm) and second harmonic generation (SHG) coefficient ($d_{33}^{2w} = 40.0$ with respect to KDP at 1060 nm) [1-3]. Due to its above mentioned extraordinary optical properties, its single crystal exhibits unique piezoelectric, acousto-optic, electro-optic and non-linear optical (NLO) properties combined with good mechanical and chemical stability [4-6]. Thus, correlation of property alteration of LT single crystals, powders, thin films, glass-ceramics, etc. with processing parameters is an important area of exploration. Consequently, in recent times researchers have demonstrated the property monitoring based on preparation of LiTaO₃ powders [7] and thin films [8, 9] by different methods. Luminescence properties of Ho³⁺, Eu³⁺, Tb³⁺ etc. doped LiTaO₃ crystals, an another important area of exploration, which have also been investigated by various researchers [10-13]. Rare-earth (RE) doped transparent LiTaO₃ nanocrystallite containing glass ceramics, in which rare-earth ions selectively incorporated into the LiTaO₃ nanocrystals embedded in an oxide glassy matrix, can offer excellent luminescent properties due to the low phonon energy environment of LiTaO₃ nanocrystallites for luminescent ions, and good mechanical and chemical properties of oxide glassy matrix. This ability, combined with inherent nonlinear optical (NLO) properties of ferroelectric crystals, could offer a possibility to design self frequency doubling laser sources. Hence, this new material has attracted great attention in the continuous research for the development of novel optoelectronic devices [14-18]. Mukherjee and Varma have reported the crystallization and physical properties of LiTaO3 in a LiBO2-Ta2O5 reactive glass matrix,

however, they have not explored RE doped LiTaO₃ containing glass-ceramics [19]. As such, work performed on nanocrystalline LiTaO₃ containing aluminosilicate glass-matrix materials is very rare due to the difficulties in preparation of transparent precursor glass in general and glass-ceramics in particular which involves high temperature (about 1600°C) for its precursor glass melting [20]. Probably for this reason, the dielectric and fluorescence properties of Nd³⁺ ion doped transparent precursor glass and glass-ceramic composites of LiTaO₃ with heat-treatment time have not been reported so far. This fact has motivated us to explore this opportunity.

Glass materials are attractive hosts for rare earth ions (RE³⁺) because planar waveguides and optical fibers can be fabricated easily with them compared to crystalline materials. Nd³⁺ is one of the most widely studied luminescent ions, and Nd³⁺ doped crystalline, glass and ceramics have been applied in many fields. Nd³⁺ has been extensively used as laser active ion in several hosts due to high quantum efficiency of its some emission channels either by normal fluorescence in NIR region or by upconversion/frequency doubling in green-red wavelength [21-23]. Most of the Nd³⁺ lasers operate in NIR region on two lasing transitions ${}^4F_{3/2} \rightarrow {}^4I_{11/2,13/2}$ at around 1060 nm and 1349 nm wavelengths respectively based on four level systems. During the past few years the ${}^4F_{3/2} \rightarrow {}^4I_{9/2}$ laser transition near 900 nm of the neodymium ion has also attracted much attention for its applications especially in display technologies, as it opens the way for efficient generation of continuous blue laser radiation in the watt range by means of a compact, all-solidstate frequency doubling laser [24]. The main basis for such a laser is provided by novel highpower diode-laser sources that offer high beam brightness for longitudinal pumping because the three level nature of this laser transition requires high pump intensities for efficient laser output. Apart from its high quantum efficiency performance as laser active ion, Nd³⁺ has also been well exploited as a structural probe in studying the local field effects around the dopant ions [25] as it

possesses hypersensitive transitions. As we aware, there is no report on Nd³⁺ doped transparent LiTaO₃ glass-ceramics prepared by controlled crystallization of melt-quench precursor glass. Hence, it will be very interesting to investigate the evolution of spectroscopic along with structural and dielectric properties of doped Nd³⁺ ions in LiTaO₃ glass-ceramic nanocomposites.

In view of above, in the present work we focus our systematic analysis on the precursor glass preparation, crystallization, structure and optical properties of isothermally crystallized Nd³⁺ ion doped Li₂O-Ta₂O₅-SiO₂-Al₂O₃ transparent glasses. The crystallization process has been studied by differential thermal analysis (DTA), X-ray diffraction (XRD), field emission scanning electron microscopy (FESEM), transmission electron microscopy (TEM), Fourier transform infrared reflection spectra (FTIRRS), optical absorption, NIR-excited (809 nm) NIR fluorescence and excited state lifetime.

It is obvious that the ferroelectric crystal containing glass-ceramics would have relatively high dielectric constant than its precursor glass due to the high spontaneous polarization ($P_s = 0.50 \text{ C/m}^2$ for LiTaO₃) as well as induced polarization of ferroelectric crystal under applied electric field [1]. S. J. Kim et al. [26] reported the dielectric properties during phase transition of the LiTaO₃-SiO₂ glasses produced by the twin roller quenching method. Hence, a study has also been carried out on how the dielectric properties (dielectric constant) change with progression of isothermal nanocrystallization of precursor Nd³⁺ doped LTSA glasses.

2. Experimental procedures

The precursor glass having molar composition 25.53Li₂O-21.53Ta₂O₅-35.29SiO₂-17.65Al₂O₃ doped with Nd₂O₃ (0.5 wt% in excess) was prepared from high-purity chemicals such as Li₂CO₃ (GR, 99%, Loba Chemie), Ta₂O₅ (99.85%, Alfa Aesar), SiO₂ (99.8%, Sipur A1 Bremtheler

Quartzitwerk), Al₂O₃ (99.8%, CT 1200 SG, Almatis), and Nd₂O₃ (99.99%, Alfa Aesar) by conventional melt-quench technique. The well-mixed batch of about 250 g glass was melted in a platinum crucible in an electric furnace at 1600°C for 2h in air. The glass melt was poured onto a pre-heated iron mould. It was annealed at 600°C for 4 h to remove the internal stresses of the glass and then slowly cooled down to room temperature. The as-prepared glass block was cut into desired dimensions and optically polished for ceramization and to performing different measurements.

The density of precursor glass was measured using Archimedes principle using water as buoyancy liquid. The refractive index of precursor glass was measured by a Prism Coupler (Model: 2010/M, Metricon Corporation) at five different wavelengths (λ = 473, 532, 633, 1064 and 1552 nm). Differential thermal analysis (DTA) of precursor glass powder was carried out up to 1000°C from room temperature at the rate of 10°C/min with a SETARAM instrument (Model: TG/DTA 92, SETARAM Instrumentation) to establish the glass transition temperature (T_g) and the crystallization peak temperature (T_p). XRD data were recorded using an XPERT-PRO MPD diffractometer (PANalytical) with Ni-filtered CuK $_\alpha$ = 1.5406 Å radiation as the X-ray source to identify the developed crystalline phases. The 2 θ scan range was 10° to 80° with a step size of 0.05°.

A high resolution FESEM (Gemini Zeiss SupraTM 35 VP model of Carl Zeiss Microimaging GmbH) was used to observe the microstructure of freshly fractured surfaces of the heat-treated glass-ceramic nanocomposites after etching in 1% HF solution for 2 minutes, dried and then coated with a thin carbon film. The TEM images and selected area electron diffraction (SAED) of powdered glass-ceramic sample were obtained from FEI (Model: Tecnai G² 30ST, FEI Company) instrument. The FTIR reflectance spectra of all Nd³⁺-doped glass and glass-ceramics

were recorded using a FTIR spectrometer (Model: 1615 Series, Perkin-Elmer Corporation) in the wavenumber range 400-2000 cm⁻¹ with a spectral resolution of ± 2 cm⁻¹ and at 15° angle of incidence. Optical absorption spectra were recorded UV-Vis-NIR spectrophotometer (Lambda 20, Perkin-Elmer Corporation) at room temperature to monitor the changes of the environmental structure of the Nd³⁺ ions. The NIR excited (809 nm) emission spectra were measured on NIR enhanced continuous bench top modular spectrofluorimeter (QuantaMaster, Photon Technology International) attached with gated Hamamatsu NIR PMT (P1.7R) as detector and Xe arc lamp as excitation source. The excited state lifetime was measured with the same instrument using a Xe flash lamp of 75 W. The dielectric constant of glass and glass-ceramic nanocomposites was measured at room temperature using a LCR meter (Model: 3532-50 Hitester, Hioki) at 1 MHz frequency after coating the surfaces with a conductive silver paint followed by drying at 140°C for 1 h.

3. Results and discussion

3.1 Thermal, optical and other physical properties

The DTA curve of the precursor glass is shown in Fig. 1. It exhibits an inflection in the temperature range $680\text{-}715^{\circ}\text{C}$ followed by a very intense exothermic peak at 820°C (T_p) corresponding to the LiTaO₃ crystallization. The glass transition temperature (T_g) has been estimated to be 702°C from the point of intersection of the tangents drawn at the slope change as marked in Fig. 1 of the DTA curve.

The precursor glass is visually transparent, appearing blue-purple due to Nd³⁺ doping. The precursor glass samples were heat treated at 680°C near glass transition temperature for 0, 3, 5, 10, 20, 50 and 100 h after nucleating at 650°C for 2 h. The obtained samples were labeled as a, b,

c, d, e, f and g (Fig. 2) respectively for convenience. The transparency of the precursor glass persists in heat-treated samples although the samples have been appeared gradually translucent due to devitrification with progress of heat-treatment duration. Fig. 3 presents Cauchy fitting based on measured refractive indices at five different wavelengths (see experimental techniques) and shows the dependences of the refractive index on the wavelength for precursor glass (a) and the 10 h heat-treated glass-ceramics (d) sample. In general, refractive index decreases with increasing wavelength due to dispersion. This trend is observed in both the samples. In addition to this, the refractive index of the glass-ceramic sample (d) has increased in comparison with precursor glass (a) that can be seen in Fig. 3. The refractive indices n_F , n_D and n_C have been estimated at three standard wavelengths ($\lambda_F = 486.1$ nm, $\lambda_D = 589.2$ nm and $\lambda_C = 656.3$ nm respectively) from the dispersion curve (Fig. 3, curve a). From the measured glass density (p) and refractive index (n_D) at wavelength $\lambda_D = 589.2$ nm, other related optical properties have been determined using relevant expressions and the results are presented in Table 1. From Table 1, it is clear that the LTSA glass under study has high values of refractive index and density. The large refractive indices of this glass are due to high ionic refraction (23.4) of Ta⁵⁺ ions [27] having an empty or unfilled d-orbital (outer electronic configuration: 5d⁰6s⁰) which contributes strongly to the linear and nonlinear polarizability [28]. The high density of the glass has originated from the large packing effect of Ta in the glass matrix [29]. For the same reason, this glass possesses a high value of molar refractivity ($R_M = 13.39 \text{ cm}^3$) and electronic polarizability $(\alpha = 1.79 \text{ x } 10^{-21} \text{ cm}^3)$. Due to formation of high refractive index LiTaO₃ (RI = 2.1834 at 600 nm [30]), the heat-treated sample exhibit higher refractive indices as shown in Fig. 3, curve-d.

3.2 X-ray diffraction analysis

The X-ray diffractograms of precursor glass and cerammed glass-ceramics are shown in Fig. 4. The XRD pattern of the precursor glass exhibits broad humps characterizing its amorphous structure. The X-ray diffraction pattern of the glass-ceramics clearly shows the structural behavior expected after a thermal treatment of the precursor glass. With progression of heattreatment, several diffraction peaks have been appeared. From the analysis of these peaks it has been concluded that these peaks are attributed to rhombohedral LiTaO₃ (JCPDS Card File No. 29-0836) except a few diffraction peak around $2\theta = 22.98^{\circ}$, 25.21° , 44.46° and 47.02° which are due to the formation of β-spodumene (LiAlSi₂O₆) crystal phase (JCPDS Card File No. 35-0797) in minor quantity. Hsu and Speyer [31] have reported that the Ta₂O₅ acts as a nucleating agent in the Li₂O-Al₂O₃-SiO₂ glass system which exhibits β-quartz solid solution (ss) crystallization, where β-spodumene is one of them, peak in the temperature range 800-900°C particularly at higher concentration of Ta₂O₅. In the present case, the small crystallization peak (minor) of βspodumene in the DTA thermogram might have merged with the large crystallization peak (major) of the LiTaO₃ at 820°C. This could be the reason of exhibiting two crystal phases in the XRD patterns although DTA curve shows only single sharp crystallization peak around 820°C. It is clearly evidenced from the XRD analysis that the peak of LiAlSi₂O₆ ($2\theta = 25.21^{\circ}$) is more prominent in sample c with 5 h heat-treatment and it got diminished with respect to LiTaO₃ phase in samples d-g, indicating the stabilization of LiTaO₃ nanocrystallites with increase in heat-treatment duration. The XRD pattern of sample b differ due to its phase separated glassy nature and having different structure than the precursor glass (a) and 5 h heat-treated glassceramic nanocomposites (c). From the full width at half maximum (FWHM) of the most intense

diffraction peak (012) of LiTaO₃, the average crystallite size (diameter, d) is calculated by using the Scherrer's formula [32]

$$d = 0.9\lambda/\beta\cos\theta\tag{1}$$

where λ is the wavelength of X-ray radiation (CuK $_{\alpha}$ = 1.5406Å), β is the full width at half maximum (FWHM) of the peak at 20.

The diameter of the LiTaO₃ crystallites is plotted as a function of heat-treatment time in Fig. 5. The average crystallite size increases with heat-treatment duration and found to vary from 14 to 36 nm for samples b-g.

3.3 FESEM and TEM image analyses

The morphology and LiTaO₃ crystallite size in glass-ceramic nanocomposites have been examined by FESEM and TEM image analyses. FESEM images of the fractured surface of samples c and e have been presented in Figs. 6(a) and (b) respectively. From the FESEM micrographs, it is clearly observed that the glassy matrix of the heat-treated samples initially phase separated on nanometric scale followed by incipient precipitation of defined crystallites within the Li–Ta rich phase regions with increase in heat-treatment time. The droplets have irregular shapes and dispersed uniformly thought out the bulk glass matrix. The size of the droplets varies in the range 20-60 nm. The TEM bright field images and their selected area electron diffraction (SAED) patterns of the 3 and 50 h heat-treated samples (b and f) are shown in Figs. 7(a) and (b) respectively. From these images, it is observed that many spheriodal LiTaO₃ crystallites precipitated homogeneously from the glass matrix and remained homogeneously dispersed in the residual glass matrix. The crystallite size from TEM image of sample b is found to be around 15 nm and of sample f to be around 18 nm. The presence of fine spherical rings

around the central bright region in SAED pattern discloses the existence of LiTaO₃ nanocrystallites in the glassy matrix.

3.4 Fourier transform infrared reflectance spectroscopy (FTIRRS)

The FTIR reflectance spectra of the as-prepared and heat-treated samples in the wavenumber range 400-2000 cm⁻¹ are shown in Fig. 8. It is seen from this figure that the precursor glass (curve-a) exhibits two broad reflection bands centered around 960 and 610 cm⁻¹ as a result of wider distribution of silicon and tantalate structural units respectively. This is an indication of the existence of structural disorder in the amorphous network with the presence of SiO₄ tetrahedra and TaO₆ octahedra having different number of non-bridging oxygen. In spite of the transparent nature of the heat-treated samples, their FTIR reflectance spectra (curves b and g) reveal narrowing of two main reflection bands with additional feature arising at 735 cm⁻¹ in comparison to the as prepared glass. In the FTIRR spectra, the stretching modes of the Si-O-Si bonds of the SiO₄ tetrahedra with nonbridging oxygen (NBO) atoms are active around 1000 cm⁻¹ (high energy side) and the stretching modes of the Ta-O bonds in the TaO₆ octahedra occur in the 600-650 cm⁻¹ ¹ range (low energy side) due to higher atomic weight of Ta than Si. The variation of Si-O (998 cm⁻¹) and Ta-O (602 cm⁻¹) stretching vibration bands intensities (here reflectivity) with heattreatment time is shown in the inset of Fig. 8. It is seen that with progression of heat-treatment the band intensities increase rapidly initially and then become almost saturated after a certain time of heat-treatment (10 h). The appearance of a low intensity band at 735 cm⁻¹ upon heattreatment related to the stretching mode of Al-O bond of AlO₄ tetrahedra of β-spodumene [33]. The prominent band occurred at 602 cm⁻¹ corresponds to the stretching mode of Ta-O bond of TaO₆ octahedral units of lithium tantalate [34, 35]. The reflection band centered at 602 cm⁻¹ is

assigned to LiTaO₃ crystal formation and the reflection band centered at 998 cm⁻¹ is assigned to Si-O stretching vibration of residual glass and β-spodumene crystal. Two bands generated at 998 and 602 cm⁻¹ in the FTIR spectrum (Fig. 8, curve-b) after 3 h heat-treatment at 680°C give the clear evidence of phase separation of precursor glass into Si-rich and Li-rich phases respectively, which exhibits very feeble sign of crystallization in the XRD pattern (Fig. 4, curve b). The gradual increase of relative intensity of band at 602 cm⁻¹ clearly indicates formation of LiTaO₃ crystal with the increase of heat treatment time. Thus from the investigations carried out on the measured FTIR reflectance spectra of Nd³⁺ doped Li₂O-Ta₂O₅-SiO₂-Al₂O₃ glass and glass-ceramics as described above provide the information of crystallization with initial phase separation followed by advancement of mainly LiTaO₃ crystal formation in the glass matrix. The results of the FTIRRS are in good agreement with that of XRD, FESEM and TEM studies. A similar observation has also been reported by Ito, et al [20].

3.5 Dielectric constant (ε_r)

Glass and glass-ceramics have certain advantages as dielectric materials because of their high dielectric strength. But the disadvantages of glass are a low permittivity ($\varepsilon_r = 4\text{-}15$) and a low thermal conductivity (about 1 W m⁻¹K⁻¹) [36]. In the present study, the as prepared Nd³⁺ doped Li₂O-Ta₂O₅-SiO₂-Al₂O₃ glass has exhibited relatively higher value (19.3) of dielectric constant (ε_r) than the common vitreous silica (3.8) or soda-lime silicate (7.2) or borosilicate glasses (4.1-4.9) [37] due to high ionic refraction of Ta⁵⁺ ions (23.4) [27]. This is due to its empty or unfilled d-orbital which contributes very strongly to its high polarizability [28, 38]. Its magnitudes show a sharp increase with increase in heat-treatment duration up to 5 h and thereafter it maintained saturation with a small decrease for any further heat treatment time as shown in Fig.9. This

suggests that, at the initial stages of heat treatment (3 h), separation of silica rich phase and Li—Ta enriched phases takes place and with the further heat-treatment, incipient precipitation of LiTaO₃ crystalline phase of high dielectric constant (ε_r = 52) [3] and spontaneous polarization (P_s = 0.50 C/m²) [1] occurs gradually which becomes well defined at 5 h and attains the maximum volume fraction of the crystalline phase. This is clearly evidenced from the observation of changes in XRD curves patterns as shown in the b and c of Fig. 4. Thus accumulation of Li⁺ ions in the phase-separated glass matrix initially could cause a slight increase of dielectric constant and with further heat treatment time due to formation of stable LiTaO₃ ferroelectric crystals remarkably increase the dielectric constant reaching the highest value for 5 h heat treated sample and then maintain almost same on further course of heat-treatment. The variation in the dielectric constant (ε_r) values among the heat-treated glass-ceramic nanocomposites are mostly due to volume fraction of crystal phases contained and also the distribution of the LiTaO₃ phase in the microstructure [39]. This result again ascertains the observations made from XRD and TEM analyses on the nanocrystallization of LiTaO₃ phase in the glass matrix.

3.6 UV-Visible-NIR absorption spectra

The room temperature measured absorption spectra of the Nd³⁺-doped precursor glass (a) and 5 and 100 h heat-treated glass-ceramic samples (c and g respectively) in the visible-NIR range have been presented in Fig. 10. The spectra reveal absorption peaks due to the $4f^3$ - $4f^3$ forced electric dipole transitions from the ground ${}^4I_{9/2}$ state to different excited states of Nd³⁺ ion in $4f^3$ configuration. All the peaks ${}^4I_{9/2} \rightarrow {}^4G_{9/2}$ (512 nm), ${}^2K_{13/2} + {}^4G_{7/2} + {}^4G_{9/2}$ (526 nm), ${}^4G_{5/2} + {}^2G_{7/2}$ (583 nm), ${}^2H_{11/2}$ (626 nm), ${}^4F_{9/2}$ (679 nm), ${}^4F_{7/2} + {}^4S_{3/2}$ (739 nm), ${}^4F_{5/2} + {}^2H_{9/2}$ (806 nm) and ${}^4F_{3/2}$ (880 nm) are assigned in accordance with Carnall's convention [40, 41]. From this figure it is noticed

that the base line of absorption spectra of heat-treated samples (c and g) has been elevated significantly with the diminishing intensities of the absorption peaks. This uplifting can be attributed to scattering of short wavelength light by the crystals [42, 43] or may be due to the difference in refractive index of crystalline phase (RI of LiTaO₃ is 2.1834 at 600 nm [30] with that of residual glassy matrix (RI = 1.7821 at 656.3 nm, see Table 1). Since the crystallites (14-36 nm) are smaller than the visible wavelength, a Rayleigh scattering model should be applicable [44]. According to this model, the scattering loss, τ is given by

$$\tau = \frac{32\pi^4 d^3 (n\Delta n)^2}{3\lambda^4} NV$$
 (2)

where d is the particle diameter, λ the wavelength of light, n the refractive index, N the number density of particles, and V the volume of the particle. With progression of heat-treatment, the number and sizes of nanocrystallites developed in the glassy matrix increase and hence the scattering centre and scattering loss increase that corresponds to a decrease in the visible transparency of the glass-ceramic nanocomposites. There is no significant difference in the shapes of absorption bands for glass and crystallized sample. During heat-treatment of the glasses, the hypersensitive transition ${}^4I_{9/2} \rightarrow {}^4G_{5/2}$, ${}^2G_{7/2}$ (obeying selection rule: $\Delta J \le 2$, $\Delta L \le 2$ and $\Delta S = 0$ with large value, 0.898 of squared reduced matrix element, $||U^{(2)}||^2$) of Nd³⁺ ion lying around 583 nm has been selected to understand the environment change closely around the rare-earth ion which is shown in Fig. 10 (inset). The spectrum is progressively sharpened in 100 h heat-treated sample (around 70 cm⁻¹) due to the line narrowing associated to the glass-crystal transformation around the rare-earth ion [45]. The band shift (around 70 cm⁻¹ towards lower wavelength) is the sign for the transformation of environmental structure of Nd³⁺ site from

amorphous to ordered crystalline, which confirms the incorporation of Nd³⁺ ions into LiTaO₃ crystallites during crystallization. The reason of band shifting towards lower wavelength can be well explained on the basis of covalent characteristics of hosts (glass and LiTaO₃ nanocrystals). The degree of covalent character of a host is estimated approximately using the formula [46]:

Covalent character (%) = exp
$$[-0.25(\Delta \chi)^2]$$
 x 100 (3)

where $\Delta \chi$ is the electronegativity of the host, i.e. the electronegativity difference $(\chi_A - \chi_C)$ of the anions and cations. The average electronegativity of anions (χ_A) or cations (χ_C) in the host is evaluated using simple additive relation [47]:

$$\chi_A \text{ or } \chi_C = \sum N_i \chi_i / \sum N_i \tag{4}$$

where N_i and χ_i are the number of individual constituent atom per mole and its electronegativity, respectively. The calculated covalent character of the precursor glass (LTSA) composition in this study is about 34% and LiTaO₃ nanocrystal is about 28%. As the degree of covalency is less in the LiTaO₃ nanocrystal host, so the absorption band shifts towards lower wavelength. This phenomenon is popularly known as the "nephelauxetic effect" after Jørgensen [48]. Similar correlation between the rare earth band position shift and the degree of covalent character of various hosts has been established by Karmakar, et al. [47, 49, 50].

3.7 NIR-excited NIR fluorescence and lifetime

The infrared fluorescence spectra (λ_{ex} = 809 nm) of the samples around 1069 nm are shown in Fig. 11. The emission band intensity around 1069 nm decreases with progression of heat-treatment. This decrease in emission intensity is due to the clustering of Nd³⁺ ions which is extremely sensitive to concentration quenching [51]. Dejneka [52] has demonstrated in fluoride glasses that clustering thereby quenching occurs when the Eu³⁺-Eu³⁺ ionic separation is around

40 Å. In the present case, the Nd^{3+} - Nd^{3+} ionic separation (R_i) in the precursor glass is found to be about 26 Å which was calculated using the relation [53]:

$$R_i(A) = (1/N_{M^{3+}})^{1/3}$$
(5)

where N_{Nd}^{3+} is the Nd^{3+} ion concentration as already provided in Table 1. It is, therefore, seen that the Nd^{3+} - Nd^{3+} ionic separation (R_i) is in the quenching region. Theoretically, the rate of relaxation due to concentration quenching varies as R_i^{-6} [53-55]. With the progression of heat-treatment, the LiTaO₃ crystal phase has been formed and the Nd^{3+} ions partitioned into the residual glassy phase by reducing the inter-ionic separation less than 26 Å of precursor glasses. This fact results in reduction in fluorescence intensity (see curve-b, c, d and g, Fig. 11) due to concentration quenching. The emission bands become sharper and take shapes as in crystalline host with progress of heat-treatment duration. All these observations indicate that the Nd^{3+} ions enter into the LiTaO₃ crystalline phase and therefore, environment around Nd^{3+} ions is changed with progression of heat-treatment (see section 3.6).

The room temperature fluorescence decay curves of the emission transition (${}^4F_{3/2} \rightarrow {}^4I_{11/2}$) at 1069 nm with an excitation at 809 nm for Nd³⁺ ions in as-prepared glass and glass-ceramic nanocomposites have been depicted in Fig. 12. The measured curves demonstrate a single exponential decay. The excited state lifetime (τ) for all has been estimated from these decay curves and the results of sample a, b, c, d and f are shown in the inset of Fig. 12. It is seen that the excited state (${}^4F_{3/2}$) lifetime (τ) decreases with increase in heat-treatment duration. This result indicates that Nd³⁺ ions undergoing clustering upon formation of glass-ceramics. This change in the heat treatment conditions leads to changes in the volume of crystallinity in the glass-ceramics and in the lifetime of the ${}^4F_{3/2}$ state due to the concentration quenching [45, 52]. Further

investigations are in progress to examine the concentration effects and also the limiting concentration of the cluster formation in this host.

4. Conclusions

Transparent glass-ceramics containing LiTaO₃ nanocrystals in the aluminosilicate glass matrix were prepared and the formation of LiTaO₃ nanocrystals has been confirmed by XRD, FTIRR spectra, FESEM and TEM images. The nanocrystallite size of LiTaO₃ has been evaluated from XRD and found to vary in the range 14-36 nm. This evaluation correlates well with those obtained from TEM images. The increase of dielectric constant in comparison to precursor glass confirms the formation of high dielectric constant ferroelectric LiTaO₃ (ε_r = 52) in glassy matrix. The absorption spectra and NIR fluorescence spectra of glass-ceramic nanocomposites demonstrate that the Nd³⁺ ion has entered into the LiTaO₃ crystalline phase and form Nd³⁺:LiTaO₃ nanocrystals. The decrease in NIR excited emission intensity and fluorescence lifetimes with progression of heat-treatment time have been attributed to the clustering of Nd³⁺ ions which resulted in concentration quenching.

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References

- [1] W.P. Risk, T.R. Gosnell, A.V. Nurmikko, Compact Blue-Green Lasers, Cambridge University Press, Cambridge, 2003.
- [2] International Centre for Diffraction Data, JCPDS card file no. 29-0836.
- [3] A.J. Moses, The Practicing Scientist's Handbook, Van Nostrand Reinhold Company, New York, 1978, p. 558.
- [4] K.S. Abedin, T. Tsuritani, M. Sato, H. Ito, Appl. Phys. Lett. 70 (1997) 10.
- [5] S. Zhu, Y. Zhu, Z. Yang, H. Wang, Z. Zhang, J. Hong, C. Ge, N. Ming, Appl. Phys. Lett. 67 (1995) 320.
- [6] K. Mizuuchi, K. Yamamoto, Appl. Phys. Lett. 66 (1995) 2943.
- [7] F. Zheng, H. Liu, D. Liu, S. Yao, T. Yan, J. Wang, J. Alloys Compd. 477 (2009) 688.
- [8] Z.X. Cheng, H. Kimura, K. Ozawa, A. Miyazaki, C.V. Kanna, J. Alloys Compd. 402 (2005) 208.
- [9] S. Youssef, R. Al Asmar, J. Podlecki, B. Sorli, A. Foucaran, Eur. Phys. J. Appl. Phys. 43 (2008) 65.
- [10] I. Sokólska, J. Alloys Compd. 341 (2002) 288.
- [11] I. Sokólska, S. Kück, G.D. Dzik, M. Baba, J. Alloys Compd. 323-324 (2001) 273.
- [12] G. Gasparotto, S.A.M. Lima, M.R. Davolos, J.A. Varela, E. Longo, M.A. Zaghete, J. Lumin. 128 (2008) 1606.
- [13] J.B. Gruber, T.H. Allik, D.K. Sardar, R.M. Yow, M. Scripsick, B. Wechsler, J. Lumin. 117 (2006) 233.
- [14] H. Jain, Ferroelectrics, 306 (2004) 111.

- [15] W.R. Romanowski, I. Sokólska, G.D. Dzik, S. Gołab, J. Alloys Compd. 300-301 (2000)152.
- [16] H. Hase, H. Nasu, A. Mito, T. Hashimoto, J. Matsuoka, K. Kamiya, Jpn. J. Appl. Phys. 35 (1996) 5355.
- [17] A. Tarafder, K. Annapurna, R.S. Chaliha, V.S. Tiwari, P.K. Gupta, B. Karmakar, J. Am. Ceram. Soc. 92 (2009) 1934.
- [18] A. Tarafder, K. Annapurna, R.S. Chaliha, V.S. Tiwari, P.K. Gupta, B. Karmakar, J. Mater. Sci. 44 (2009) 4495.
- [19] P. Mukherjee, K.B.R. Varma, Ferroelectrics 306 (2004) 129.
- [20] S. Ito, T. Kokubo, M. Tashiro, J. Mater. Sci. 13 (1978) 930.
- [21] P. Zeller, P. Peuser, Opt. Lett. 25 (2000) 34.
- [22] F. Lahoz, I.R. Martín, U.R. Roríguez-Mendoza, I. Iparraguirre, J. Azkargorta, A. Mendioroz, R. Balda, J. Fernández, V. Lavín, Opt. Mater. 27 (2005) 1762.
- [23] M.C. Silva, F.H. Cristovan, C.M. Nascimento, M.J.V. Bell, W.O. Cruz, A. Marletta, J. Non-Cryst. Solids 352 (2006) 5296.
- [24] V. Seznec, H.L. Ma, X.H. Zhang, V. Nazabal, J-L. Adam, X.S. Qiao, X.P. Fan, Opt. Mater. 29 (2006) 371.
- [25] S. Mohan, K.S. Thind, G. Sharma, L. Gerward, Spectrochim. Acta A 70 (2008) 1173.
- [26] S.J. Kim, J.E Kim., Y.S. Yang, Ferroelectrics 272 (2002) 321.
- [27] M.B. Volf, Chemical Approach to Glass, Elsevier, Amsterdam, 1984, p.125.
- [28] M. Yamane, Y. Asahara, Glasses for Photonics, Cambridge University Press, Cambridge, UK, 2000, p.173.
- [29] C. Hirayama, D. Berg, J. Am. Ceram. Soc. 46 (1963) 85.

- [30] CRC Handbook of Materials Science, Vol. III, edited by C.T. Lynch, CRC Press, Cleveland, Ohio, 1975, p.170.
- [31] J.Y. Hsu, R.F. Speyer, J. Am. Ceram. Soc. 72 (1989) 2334.
- [32] B.D. Cullity, Elements of X-Ray Diffraction, Addison-Wesley Publishing Co., London, 2nd Edition, 1978, p. 101.
- [33] V.L. Burdick, D.E. Day, J. Am. Ceram. Soc. 50 (1967) 97.
- [34] H. Ono, Y. Hosokawa, K. Shinoda, K. Koyanagi, H. Yamaguchi, Thin Solid Films 381 (2001) 57.
- [35] J.Y. Zhang, I.W. Boyd, V. Dusastre, D.E. Williams, J. Phys. D: Appl. Phys. 32 (1999)L91-95.
- [36] A.J. Moulson, and J.M. Herbert, Electroceramics: Materials. Properties. Applications, Chapman & Hall, London, 1990, p.222.
- I.A. Blech, Properties of Materials in Electronics Engineering Handbook, 2nd Edition,
 D.G. Fink, D. Christiansen Ed., McGraw-Hill, New York, 1986, p. 6-30.
- [38] W.P. Risk, T.R. Gosnell, A.V. Nurmikko, Compact Blue-Green Lasers, Cambridge University Press, Cambridge, 2003.
- [39] D.E. Vernacotola, Key Engg. Mater. 94-95 (1994) 379-408.
- [40] W.T. Carnall, P.R. Fields, K. Rajnak, J. Chem. Phys. 49 (1968) 4424.
- [41] Y Chen, Y. Huang, M. Huang, R. Chen, Z. Luo, J. Am. Ceram. Soc. 88 (2005) 19.
- [42] G.H. Beall, D.A. Duke, Glass-Ceramic Technology in Glass Science and Technology, Vol. 1, D.R. Uhlmann, N.J. Kreidl (Eds.), Academic Press, New York, 1983, p. 403.
- [43] G.H. Beall, D.A. Duke, J. Mater. Sci. 4 (1969) 340.
- [44] H.C. Van De Hulst, Light Scattering by Small Particles, Wiley, New York, 1957.

- [45] M. Mortier, A. Monteville, G. Patriarche, G. Mazé, F. Auzel, Opt. Mater. 16 (2001) 255.
- [46] L. Pauling, The Nature of Chemical Bond, 3rd ed. Cornell University Press, New York, 1970.
- [47] B. Karmakar, J. Solid State Chem. 178 (2005) 2663.
- [48] C.K. Jørgensen, Modern Aspects of Ligand Field Theory, North-Holland Publishing Co., Amsterdam, 1971.
- [49] B. Karmakar, R. N. Dwivedi, J. Non-Cryst. Solids 342 (2004) 132.
- [50] T. Som, B. Karmakar, J. Alloys Compd. 476 (2009) 383.
- [51] P. Riello, S. Bucella, L. Zamengo, U. Anselmi-Tamburini, R. Francini, S. Pietrantoni, Z.A. Munir, J. Eur. Ceram. Soc. 26 (2006) 3301.
- [52] M.J. Dejneka, J. Non-Cryst. Solids 239 (1998) 149.
- [53] K. Pátek, Glass Lasers, Butterworth & Co (Publishers) Ltd., London, 1970, p.95.
- [54] J.H. Campbell, T.I. Suratwala, J. Non-Cryst. Solids 263-264 (2000) 318.
- [55] Uk Kang, A.A. Zhilin, D.P. Logvinov, A.A. Onushchenko, V.A. Savost'yanov, T.I. Chuvaeva, A.V. Shashkin, Glass Phys. Chem. 27 (2001) 344.

FIGURE CAPTIONS

- Fig. 1. DTA curve of Nd³⁺:LiTaO₃ precursor glass powder.
- Fig. 2. (Color online) Photograph of Nd³⁺:LiTaO₃ precursor glass and glass-ceramic nanocomposites (thickness: 2 mm) laid over the writing to show their transparency.
- Fig. 3. Variation of refractive indices (Cauchy fitted) of Nd³⁺:LiTaO₃ (a) precursor glass and (d) 10 h heat-treated glass-ceramic nanocomposites as a function of wavelength.
- Fig. 4. XRD patterns of the samples a-g.
- Fig. 5. Variation of crystal size as a function of heat-treatment time.
- Fig. 6. FESEM images of samples (a) c and (b) e.
- Fig. 7. (a) TEM images and SAEDs (Inset) of samples (a) b and (b) f.
- Fig. 8. FTIRR spectra of (a) precursor glass, (b) 3 h and (g) 100 h heat-treated samples.
- Fig. 9. Variation of dielectric constant of Nd³⁺:LiTaO₃ precursor glass (a) and glass-ceramic nanocomposites (b-g) as a function of heat-treatment time.

Fig. 10. Absorption spectra of samples a, c and g (thickness: 2 mm). Inset shows the peak position shifting of the hypersensitive transition, ${}^4I_{9/2} \rightarrow ({}^4G_{5/2}, {}^2G_{7/2})$, of Nd³⁺ ions in these samples in the wavelength range 575-595 nm.

Fig. 11. (Color online) Near infrared (NIR) fluorescence spectra (λ_{ex} = 809 nm) of Nd³⁺:LiTaO₃ precursor glass and glass-ceramic nanocomposites obtained after heat-treatment for various duration (a, b, c, d and g indicate the sample identity).

Fig. 12. (Color online) Decay curves for the ${}^4F_{3/2} \rightarrow {}^4I_{11/2}$ transition of Nd³⁺ ion at1069 nm under excitation at 809 nm of samples a, b, c, d and f.

 $\label{eq:continuous_continuous$

Properties	Corresponding
	value
Average molecular weight, Mav	142.37
Density, ρ (g.cm ⁻³)	4.50
Refractive indices:	
n _F (at 486.1 nm)	1.8053
n _D (at 589.2 nm)	1.7894
n _C (at 656.3 nm)	1.7821
Abbe number (v_D)	34
Dispersive power $(1/v_D)$	0.03
Molar refractivity, $R_{\rm M}$ (cm ³)	13.39
Electronic polarizability, α (cm ³)	1.79 x 10 ⁻²¹
Nd ³⁺ ion concentration, N _{Nd³⁺} (ions/cm ³)	5.66×10^{19}
Nd ³⁺ -Nd ³⁺ inter ionic distance, R _i (Å)	26
Glass transition temperature, T _g (°C)	702
Crystallization peak, T _p (°C)	820

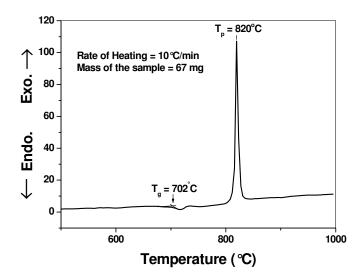


Fig. 1. DTA curve of Nd³⁺:LiTaO₃ precursor glass powder.

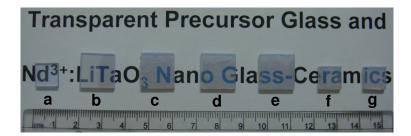


Fig. 2. (Color online) Photograph of Nd³⁺:LiTaO₃ precursor glass and glass-ceramic nanocomposites (thickness: 2 mm) laid over the writing to show their transparency.

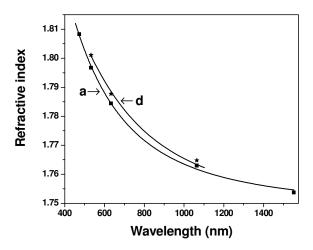


Fig. 3. Variation of refractive indices (Cauchy fitted) of Nd³⁺:LiTaO₃ (a) precursor glass and (d) 10 h heat-treated glass-ceramic nanocomposites as a function of wavelength.

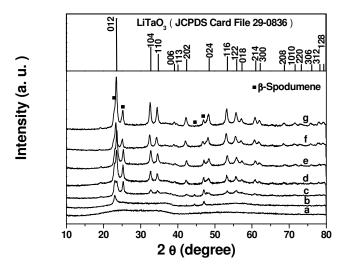


Fig. 4. XRD patterns of the samples a-g.

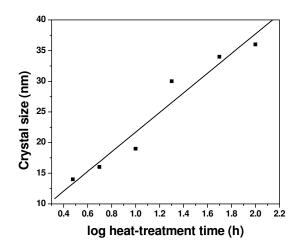


Fig. 5. Variation of crystal size as a function of heat-treatment time.

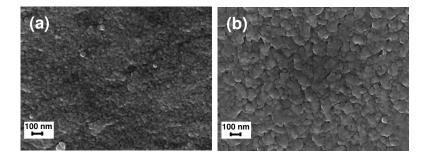


Fig. 6. FESEM image of samples (a) c and (b) e.

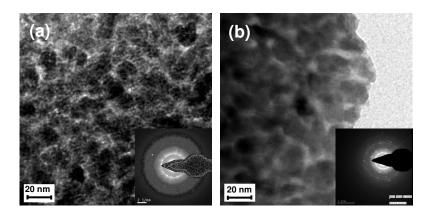


Fig. 7. (a) TEM images and SAEDs (Inset) of samples (a) b and (b) f.

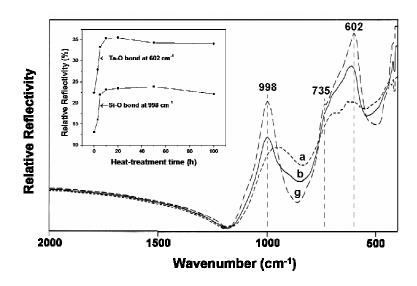


Fig. 8. FTIRR spectra of (a) precursor glass, (b) 3 h and (g) 100 h heat-treated samples.

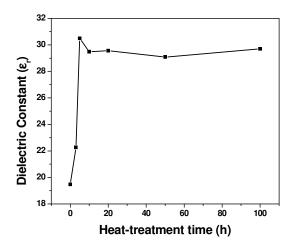


Fig. 9. Variation of dielectric constant of Nd³⁺:LiTaO₃ precursor glass (a) and glass-ceramic nanocomposites (b-g) as a function of heat-treatment time.

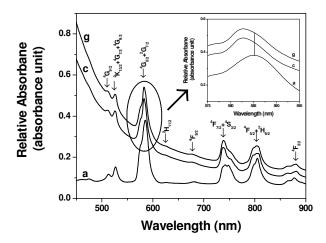


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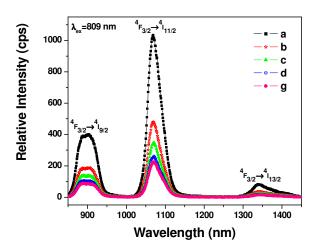


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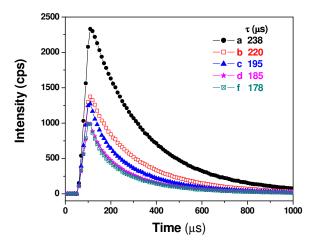


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