

1     **Thermal nature of mantle upwellings below the Ibero-western Maghreb**  
2                     **region inferred from teleseismic tomography**

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20    Key points:

- 21       • New high-resolution teleseismic *S*-wave tomographic model of the upper-mantle  
22       structure below the Ibero-western Maghreb region  
23       • Mantle upwellings below Canaries, Atlas and Gibraltar arc are sourced in the lower  
24       mantle and interact with the retreating Gibraltar slab  
25       • The main signature of the mantle upwellings is thermal in nature with temperature  
26       excesses of ~100-350°C

27

28

29 Abstract

30

31 Independent models of *P*- and *S*-wave velocity anomalies in the mantle derived from  
32 seismic tomography help to distinguish thermal signatures from those of partial melt,  
33 volatiles and compositional variations. Here, we use seismic data from *SW* Europe and *NW*  
34 Africa, spanning the region between the Pyrenees and the Canaries, in order to obtain a new  
35 *S-SKS* relative arrival-time tomographic model of the upper mantle below Iberia, Western  
36 Morocco and the Canaries. Similar to previous *P*-wave tomographic results, the *S*-wave  
37 model provides evidence for: (1) sub-vertical upper-mantle low-velocity structures below the  
38 Canaries, Atlas Ranges and Gibraltar Arc, which are interpreted as mantle upwellings fed by  
39 a common lower-mantle source below the Canaries; and (2) two low-velocity anomalies  
40 below the eastern Rif and Betics that we interpret as the result of the interaction between  
41 quasi-toroidal mantle flow induced by the Gibraltar slab and the mantle upwelling behind it.  
42 The analysis of teleseismic *P*- and *S*-wave arrival-time residuals and the conversion of the  
43 low-velocity anomalies to temperature variations suggest that the upwellings in the upper  
44 mantle below the Canaries, Atlas Ranges and Gibraltar Arc system may be solely thermal in  
45 nature, with temperature excesses in the range  $\sim 100\text{-}350^\circ\text{C}$ . Our results also indicate that  
46 local partial melting can be present at lithospheric depths, especially below the Atlas Ranges.  
47 The locations of thermal mantle upwellings are in good agreement with those of thinned  
48 lithosphere, moderate to high heat-flow measurements and recent magmatic activity at the  
49 surface.

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## 55 1. Introduction

56

57 Travel-time tomography exploits information contained within seismic datasets in order  
58 to constrain seismic velocity anomalies associated with variations in Earth's internal  
59 structure. Seismic velocity heterogeneities can result from variations in temperature,  
60 chemical composition, partial melt and volatile content, anisotropy and grain size  
61 (Cammarano et al., 2005; Faul & Jackson, 2002; Karato & Jung, 1998; Schmandt &  
62 Humphreys, 2010b). An ongoing challenge in imaging the Earth's internal structure is to  
63 infer the degree to which each of these mechanisms contributes to a velocity anomaly. Aided  
64 by continuous improvements in the quantity of seismic data and developments in inversion  
65 techniques, numerous studies have attempted to infer the nature of seismic velocity  
66 anomalies. The independent inversion of  $P$ - and  $S$ -wave velocity anomalies can be exploited  
67 to gain insight into the physical origin of anomalies. Importantly, the ratio of relative changes  
68 in shear to compressional wave velocities, defined as  $R_{S,P} = d\ln V_S / d\ln V_P = (dV_S/V_S) / (dV_P/V_P)$   
69 has been used as an indicator of the physical causes of the observed velocity variations  
70 (Cammarano et al., 2003; Masters et al., 2000; Resovsky & Trampert, 2003; Robertson &  
71 Woodhouse, 1996; Saltzer et al., 2001; Simmons et al., 2009). Geophysical and mineral-  
72 physics studies suggest that mantle temperature variations produce  $R_{S,P}$  of magnitudes in the  
73 range 1.1–2.2 (Cammarano et al., 2003; Goes et al., 2000; Koper et al., 1999). The effect of  
74 composition on  $V_P$  and  $V_S$  is small compared to that of temperature (for most plausible  
75 compositions) because of the very strong temperature sensitivity and therefore distinguishing  
76 between the two remains complicated (Goes et al., 2000). The presence of melt is known to

77 have the strongest effect on  $V_{P,S}$ , causing an increase of  $R_{S,P}$  up to 4.0, particularly for aligned  
78 film-like pores and melt fractions above 1% (Hammond & Humphreys, 2000; Takei, 2002).

79 A number of teleseismic studies have measured variations in  $V_P/V_S$  by jointly inverting  
80  $P$ - and  $S$ -arrival-time datasets (e.g., Hammond & Humphreys, 2000; Schmandt &  
81 Humphreys, 2010a). A recent work by Papaleo et al., (2018) estimated  $V_P/V_S$  from  
82 teleseismic data which constrain relative rather than absolute velocities; for this approach to  
83 be successful, accurate estimates of the background  $P$ - and  $S$ - wave velocity structure are  
84 required. However, both  $P$ - and  $S$ -wave velocity models are often not available for the same  
85 region and when they are, limitations arising from the resolving power of the tomography  
86 contaminate  $R_{S,P}$  estimates (e.g., Deschamps & Trampert, 2003). When one or both  $P$ - and  $S$ -  
87 wave velocity models lack good resolution, it is preferable to compare directly  $P$ - and  $S$ -wave  
88 relative arrival-time residuals, for common earthquake-station pairs (e.g., Bastow et al.,  
89 2005). This approach avoids difficulties related to the varying amplitude recovery of velocity  
90 anomalies (e.g. due to differing numbers of travel-time observations, different levels of noise  
91 and different regularization parameters) and other issues associated with the underdetermined  
92 nature of the tomographic inverse problem. Moreover, the ratio of the arrival-time residuals  
93 ( $a_{S,P}$ ) is proportional to the ratio of absolute  $dV_S$  and  $dV_P$  along the chosen station-event  
94 pairs; therefore the least-squares fit of  $a_{S,P}$  equals  $(V_P/V_S)R_{S,P}$  (Civiero et al., 2016). As this  
95 estimate is averaged along the whole path, from source to receiver, the spatial distribution of  
96 the anomaly may not be well known. However, if we measure it for several source-station  
97 pairs, this limitation can be mitigated.

98 The computation of  $R_{S,P}$  from the ratio of arrival-time residuals was first proposed by  
99 Hales and Doyle, (1967) who investigated variations in  $S$ - and  $P$ -wave residuals to infer the  
100 presence of melting beneath the Western United States. More recently, Rocha et al., (2011)  
101 focused on residuals to infer distinct temperature and compositional influences in Brazil.

102 Residual analysis performed for Ethiopia found a slope consistent with purely thermal  
103 variations, although in the shallow mantle partial melt is not excluded (Bastow et al., 2005;  
104 Civiero et al., 2016).

105 It has been demonstrated that in the upper mantle seismic velocities are primarily  
106 sensitive to temperature and less so to composition (Afonso et al., 2010; Cammarano et al.,  
107 2003; Goes et al., 2000; Sobolev et al., 1997). Down to depths of 200 km in the mantle,  
108 temperature variations induce strong relative  $dV_S$ -anomalies, whereas variations in  
109 composition generate weaker anomalies. For example, for a 100°C increase in temperature, a  
110 decrease of 0.7-4.5% in  $V_S$  is predicted, mostly due to the large effect of anelasticity at high  
111 temperature (Goes et al., 2000). The effects of realistic compositional variations instead  
112 produce velocity anomalies <1% and are therefore more difficult to retrieve (Goes et al.,  
113 2000). As such, while it may be reasonable to interpret  $dV_S$ -anomalies largely in terms of  
114 temperature distribution, it is challenging to make meaningful inferences about compositional  
115 variations (Forte & Perry, 2000). The amplitudes of tomographic velocity anomalies can thus  
116 be scaled to temperature anomalies (e.g., Cammarano et al., 2003; Goes et al., 2000; Sobolev  
117 et al., 1996; Yan et al., 1989) under the assumption that compositional variations are  
118 “second-order”. Using this reasoning, independent  $P$ - and  $S$ -wave tomographic models have  
119 been used to infer the thermal structure of the upper mantle below North America (Goes &  
120 van der Lee, 2002), Europe (Goes et al., 2000) and Australia (Goes et al., 2005). A similar  
121 approach was followed by Currie and Hyndman, (2006) for circum-Pacific back arcs and by  
122 Civiero et al., (2016, 2015) for Ethiopia, using different mantle composition assumptions.  
123 In this study we address the seismic and thermal structure of the upper-mantle below the  
124 Ibero-western Maghreb. This region, located between the western Mediterranean Sea and  
125 central-eastern Atlantic Ocean, has a complex tectonic history resulting from the convergence  
126 between the African and Eurasian plates (Lonergan and White, 1997). The western margin of

127 Iberia, covered by Paleozoic outcrops, forms the Iberian Massif which has been geologically  
128 stable for the last 300 Ma (Gibbons & Moreno, 2002). The Valencia Trough, a Neogene *SW*–  
129 *NE* oriented extensional basin, delimits the region offshore to the east-northeast (Fontboté et  
130 al., 1990). The Iberian-European plate boundary is represented by the Pyrenees, an orogen  
131 resulting from the collision of the two plates during the Cretaceous and Miocene  
132 (Choukroune, 1989). Southwards, the Gibraltar Arc System includes different structures such  
133 as the Alboran Basin and the Betic-Rif orogen (Gutscher et al., 2002). South of the Rif, the  
134 intra-continental Atlas Mountains consist of Paleozoic, Mesozoic and Cenozoic rocks with  
135 summits that exceed 4000 m elevation (e.g., Zeyen et al., 2005) (Fig. 1A).

136 Existing *S*-wave models for the region have been obtained by finite-frequency Rayleigh  
137 wave tomography (e.g., Palomeras et al., 2014; Peter et al., 2008; Schivardi and Morelli,  
138 2009). However, the depth extent of these models is more limited than that achieved by body-  
139 wave travel-time tomography. As a result, a comprehensive view of the upper mantle, as  
140 given by similarly resolved *P*- and *S*-wave velocity models, is not available. In this study, we  
141 present the first teleseismic travel-time tomography model which images *S*-wave velocities at  
142 high resolution down to the base of the mantle transition zone (*MTZ*). We further present a  
143 comparison with the *P*-wave velocity structure already imaged in a companion study (Civiero  
144 et al., 2018), which we will refer to as *IBEM-P18*. Finally, we investigate the nature of the  
145 low-velocity anomalies by analysing *P*- and *S*-wave arrival-time residuals and by converting  
146 the velocity anomalies to temperature estimates using a mineral-physics approach.

147

## 148 2. Data and Method

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### 150 2.1. Datasets

151

152 We used relative arrival-time residuals of teleseismic *S*- and *SKS*-phases observed at 410  
153 broadband seismic stations deployed in the Ibero-western Maghreb region to invert for 3*D* *S*-  
154 wave velocity perturbations in the mantle (Fig. 1). The inclusion of *SKS*-phases into the *S*-  
155 wave travel-time tomography provides additional, steeper, crossing paths down to the *MTZ*.  
156 The ensemble of stations used encompassed 203 stations from the IberArray deployment  
157 (Díaz et al., 2009), 91 from the PICASSO array (Platt et al., 2008), 20 from the WILAS  
158 project (Custódio et al., 2014) and 10 OBSs deployed by the NEAREST experiment (Carrara  
159 and NEAREST Team, 2008). An additional 86 temporary and permanent instruments from  
160 another 13 seismic networks were also used. In total, we used data from 17 international  
161 seismic experiments and permanent networks that operated over the course of six years, from  
162 2007 to 2013. Detailed information about the stations is available in the Supplementary Table  
163 S1. We selected 380 high-quality teleseismic earthquakes with magnitudes larger than 5.5  
164 and epicentral distances of 30°–135° (30°–95° for *S* phases and 95°–135° for *SKS* phases).

165 Relative arrival-time residuals were estimated using an adaptive stacking technique  
166 developed by Rawlinson and Kennett, (2004), which exploits the coherency in the arriving  
167 waveforms across the array. We first filtered seismograms by applying a 0.04–4 Hz band-pass  
168 Butterworth filter. Subsequently, the data were rotated into radial and tangential components  
169 and residuals were obtained for direct *S*-phases observed on the tangential component and for  
170 *SKS*-phases observed on the radial component. The automatic process was complemented by  
171 a visual inspection of the waveforms that ensured the selection of high-quality measurements  
172 only. After the rejection of poor-quality data, our dataset comprised 15619 *S*- and *SKS*-wave  
173 relative arrival-time residuals. The residuals range from approximately -5 to 5 s. Further  
174 details on the arrival-time residual estimation, the back-azimuthal variation of the residuals  
175 for a selection of stations located in different parts of the Ibero-western Maghreb region, and  
176 the epicentral distance distributions are shown in Supplementary Information S1.

177

## 178 2.2. Model parameterisation and inversion method

179

180 We adopt the same grid spacing used for *IBEM-P18* by Civiero et al., (2018), with 0.4°  
181 node spacing in latitude (from 26°N to 46°N) and longitude (from 19°W to 5°E) and ~35 km  
182 spacing in depth, extending from the Moho down to 800 km depth.

183 Relative arrival-time residuals were inverted using the *FMTOMO* package (Rawlinson et  
184 al., 2006) to recover the mantle seismic structure. *FMTOMO* applies the Tau-P method  
185 (Kennett & Engdahl, 1991) to compute travel-times from the source to the edge of the *3D*  
186 model region. It then combines the Fast Marching Method as a forward solver (Rawlinson &  
187 Sambridge, 2004a, 2004b), which tracks the evolving wavefront from the edge of the local  
188 *3D* model to the receivers at the surface, and a subspace inversion technique, which adjusts  
189 model parameters to satisfy observations (Kennett et al., 1988).

190 The *3D* starting model for the crust and upper mantle was built by converting the  
191 *PRISM3D* *P*-wave model (Civiero et al., 2018) into an *S*-wave velocity model assuming a  
192  $V_P/V_S$  ratio calculated from the *ID ak135* *P*- and *S*-wavespeed models as a function of depth  
193 (Fig. S2a, b). The initial *S*-wave velocities in the topmost lower mantle (660-800 km depth)  
194 are those converted from the *LLNL* global *P*-wave tomographic model of Simmons et al.,  
195 (2012) (Fig. S2c). To reduce the effect of the unresolved crustal contributions to the arrival-  
196 time residuals we included in our starting model a realistic *a priori 3D* crust and Moho model  
197 taken from *PRISM3D*. We also performed two additional inversions using different starting  
198 models: the *ID ak135* *S*-wavespeed model (Kennett et al., 1995) and the *3D* *S*-wave global  
199 model (*SEMum2*) of French et al., (2013). Our final interpretation will be based on the model  
200 obtained using *PRISM3D* as a starting model, although the features we will focus on are  
201 clearly present in all three models.



202 Smoothing ( $\eta$ ) and damping ( $\varepsilon$ ) parameters of 5 and 5, respectively, are used and provide  
203 the optimum trade-off between data fit and model roughness (Fig. S3). However, different  
204 choices of  $\eta$ - $\varepsilon$  combinations in the neighbourhood of our preferred values do not significantly  
205 change the pattern of the anomalies. The final *S*-wave model solution reduces differential  
206 travel-time variance by 46.71%, from 2.56 s<sup>2</sup> to 1.36 s<sup>2</sup>.

207

### 208 3. Resolution tests

209

210 We carried out detailed resolution tests in order to assess whether the main features of our  
211 *S*-wave tomographic model are actually required by the data. First, we generated a set of  
212 synthetic arrival-time residuals by tracing rays through a known test structure. Then, we  
213 added Gaussian noise to the synthetic arrival-time residuals (0.4 s standard deviation) of the  
214 same order of magnitude as the data noise, which was estimated by the adaptive stacking  
215 approach. Finally, we inverted those residuals using the same algorithm and parameterization  
216 that is used for the inversion of the actual field dataset. The comparison between the synthetic  
217 input test model and the recovered output model provides a basis to assess the reliability of  
218 the features recovered in the inversion of the field dataset.

219 In this study, we performed two different types of resolution tests. In the first case, we  
220 introduced a checkerboard pattern of alternating positive and negative anomalies, with  
221 diameter of  $\pm \sim 200$  km and amplitude of  $\pm 0.50$  km/s, separated by a region of zero  
222 perturbation (Fig. 2). Slices at 250–500 km depth through the reconstructed *S*-wave model  
223 (Fig 2c,d) show a good agreement between the input and recovered models. However, the  
224 recovered anomalies do exhibit some vertical smearing along dominant ray paths, especially  
225 in the western (Atlantic Ocean) and eastern (east of the Alboran Sea) parts of the model.  
226 Bigger features, with diameters of  $\sim 300$  km, are still well resolved within the *MTZ* (Fig. S4).

227 The synthetic tests indicate that the  $S$ -wave data cannot resolve structures as clearly as the  $P$ -  
228 wave data (as demonstrated by a comparison with *IBEM-P18*, see Figures S5 and S6 in the  
229 Supplementary Information), especially below the *MTZ*. This is a fairly typical outcome,  
230 since the  $S$ -wave residuals tend to be more noisy than those of  $P$  waves, and the number of  
231 arrivals that are picked is lower (Bolton & Masters, 2001).

232 In the next test we used three vertical structures (amplitude  $\approx -0.30$  km/s, diameter  $\sim 100$   
233 km) located approximately below the Canaries (*A1*), Atlas Ranges (*A2*) and Gibraltar Arc  
234 (*A3*) as input (Fig. 3). The negative synthetic anomalies are positioned where upper-mantle  
235 upwellings have been imaged before (Civiero et al., 2018). We will use the resolving power  
236 of these simple structures to scale our tomographic model according to the estimated  
237 amplitude recovery, before converting to relative temperature anomalies (see Section 4.3). In  
238 spite of some smearing downward and along a number of oblique ray paths, all structures are  
239 well recovered. *A1* is the best resolved region, which may result from a good crossing-rays  
240 coverage likely due to the stations above *A1* being part of a permanent network, with a longer  
241 operating period compared to the temporary arrays which span *A2* and *A3*.

242

## 243 4. Results

244

### 245 4.1. $S$ -wave tomographic model

246

247 Figures 4 and 5 display depth slices and cross-sections through our  $S$ -wave velocity  
248 model, ranging from 70 km down to 730 km depth. The  $dV_S$ -structure overall exhibits a  
249 similar pattern of anomalies to  $dV_P$  in *IBEM-P18*, despite the decreased ray coverage.

250 The model shows a pronounced positive velocity anomaly corresponding to the location  
251 of the subducted Gibraltar slab ( $dV_S \approx 0.3$  km/s in its core), which dips near-vertically in the

252 upper mantle. The shape, dip and amplitude of the slab are consistent with the *P*-wave results  
253 (Civiero et al., 2018). We also observe a prominent high-velocity anomaly below Western  
254 Iberia and an *EW*-elongated high-velocity feature beneath the Southern Pyrenees, both of  
255 which extend through the upper mantle, similar to what was imaged in *IBEM-P18*. The most  
256 notable difference with the *IBEM-P18* is the absence of a strong high-velocity anomaly  
257 below the Gorringe Bank. This is likely due to the poorer resolution in this offshore area,  
258 which may result from the fact that we did not use TOPOMED OBS data in this study  
259 because the horizontal components lacked accurate orientations and had a poor signal-to-  
260 noise ratio.

261 In the oceanic domain, a prominent low-velocity anomaly is imaged in the upper mantle  
262 below *A1* ( $dV_S \approx -0.2$  km/s). *A2* and *A3* also exhibit moderate  $dV_S$  ranging from -0.3 km/s in  
263 the lithosphere to -0.1 km/s or less within the *MTZ*, with a similar geometry and vertical  
264 continuity to those seen in *IBEM-P18*. In cross-sections *AA'* and *CC'* (Fig. 5a,c) we observe  
265 that the low-velocity anomaly in *A3* is abruptly truncated to the east when it comes into  
266 contact with the high-velocity Gibraltar slab. Importantly, we observe a connection between  
267 these three slow anomalies at *MTZ*-depths, as also imaged in *IBEM-P18*.

268 Low-velocity anomalies below the Betic-Rif system ( $dV_S \approx -0.2$  km/s) are also imaged as  
269 previously found in *IBEM-P18*. They extend through the upper mantle and, together with the  
270 slow feature west the Gibraltar slab, they surround the high-velocity body in the Alboran Sea.  
271 Similar to what was imaged in *IBEM-P18*, the slow mantle structures below the Betic-Rif  
272 system appear disconnected from the lower mantle, although in some sections the structure  
273 seems to extend deeper.

274 The results of the inversions using the *ID ak135* and the *3D SEMum2* as starting models  
275 exhibit the same first-order characteristics as those discussed above (see Figs. S7 and S8).  
276 The main difference is the magnitude of the features imaged, which are weaker in the first

277 case (*ak135* reference model) and much stronger at lithospheric depths in the latter case (*3D*  
278 *SEMum2* model) (Fig. S8, a-c). The more pronounced anomalies of the latter model are not  
279 surprising, given that waveform inversions do not regularise amplitude as strongly as travel-  
280 time inversions do. In addition, damping back to a *3D* starting model instead of a *1D* model  
281 means that larger amplitudes will always be favoured in the former case (Rawlinson et al.,  
282 2010).

283

#### 284 4.2. Relative arrival-time residuals

285

286 To further explore the nature of the recovered perturbations,  $R_{S,P}$  is often calculated  
287 directly from the tomographic models (e.g., Cammarano et al., 2003; Karato & Karki, 2001;  
288 Saltzer et al., 2001; Takei, 2002). However, the strong differences in spatial resolution and  
289 smearing of our *P*- and *S*-wave models prevent us from using this method. Instead, we  
290 compare the *P*- and *S*-wave relative arrival-time residuals for common earthquake-station  
291 pairs, thus avoiding difficulties associated with the amplitude recovery of anomalies,  
292 potential smearing, spatial resolution dependence on station distribution and  
293 parameterization/regularization. The *P*-wave residuals are those of Civiero et al., (2018).  
294 Figure 6 shows the *S*-wave relative residuals plotted as a function of the *P*-wave relative  
295 residuals, for stations deployed in the three different areas where we image sub-vertical low-  
296 velocity anomalies extending within the *MTZ*: *A1* (Canaries), *A2* (Atlas Ranges) and *A3*  
297 (Gibraltar Arc). We examine *A3* by analyzing *P*- and *S*-wave residuals from stations near the  
298 Strait of Gibraltar that record phases that have mainly sampled the Gibraltar Arc. Only the  
299 highest-quality data are shown and outliers identified by visual inspection are removed prior  
300 to the analysis (see Table S2). The *P*- and *S*-wave data are calculated independently, but in  
301 general show a positive correlation, as expected. In order to fit a straight line through these

302 data, a standard least squares regression model should not be used, because it assumes that  
303 there are no errors in  $x$  (i.e. the  $P$ -wave residuals), which is clearly not the case (see Figure  
304 6). Instead, we apply the LinFitXY tool from MATLAB, which takes into account errors in  
305 both  $x$  ( $P$ -wave residuals) and  $y$  ( $S$ -wave residuals). This yields a slope  $a_{S,P}$ , which may be  
306 interpreted as a first order indication of the thermal/chemical origin of the velocity anomalies.

307 Previous studies using the ratio of  $P$ - to  $S$ -wave relative travel-time residuals found that  
308 regions with anomalies caused by thermal variations only have  $a_{S,P}$  ratios that vary between  
309 1.8 and 2.2 ( $R_{S,P} = 1.1$ - $1.3$ ) (Chung, 1971; Koper et al., 1999). Several other studies found  
310 that an  $a_{S,P}$  of around 2.9 may be on the high side to invoke only temperature perturbations  
311 (e.g., Gao et al., 2004) and values  $> 2.9$  likely require the presence of compositional  
312 heterogeneities (Rocha et al., 2011). A more recent residual analysis focused on the Afar used  
313 a much higher value, with  $a_{S,P}$  equal to 3.7, to invoke a thermal origin for the low-velocity  
314 anomaly that was imaged (Civiero et al., 2016). However, from our previous discussion of  
315 the expected range of  $R_{S,P}$  values,  $a_{S,P}$  up to  $\sim 3.8$  can reasonably be interpreted to correspond  
316 to an  $R_{S,P}$  within the thermal range.

317 The  $a_{S,P}$  in all our regions fall within the range 2.1- 3.8 ( $A1 = 2.5 \pm 0.1$ ;  $A2 = 3.8 \pm 0.2$ ;  $A3$   
318  $= 2.1 \pm 0.1$ ), suggesting that the low-velocity anomalies in the study region are likely due to  
319 temperature variations alone. We cannot, however, rule out the possibility that the presence  
320 of melt or compositional variations may contribute to the velocity anomalies we detect,  
321 especially for  $A2$ , which shows an  $a_{S,P}$  at the higher end of the thermal range.

322

### 323 4.3. Velocity-temperature conversion

324

325 We now convert the seismic-wave anomalies,  $dV_P$  and  $dV_S$ , of the three low-velocity  
326 regions,  $A1$ ,  $A2$  and  $A3$  at 250 km depth, to temperature anomalies following the method

327 described in Civiero et al., (2016). In order to convert the seismic velocities to temperature  
328 anomalies, one needs to account for the spatial variability of the seismic resolution, which  
329 has implications for the amplitude of the recovered anomalies. A practical way of doing so is  
330 to consider the amplitude scaling of the input and output synthetic anomalies. Thus, before  
331 proceeding with the velocity to temperature conversion, we scale the *P*- and *S*-wave  
332 tomographic velocity anomalies according to the amplitude recovery inferred from the center  
333 of the three idealized mantle upwelling models (Fig. 3, Fig. S9). The amplitude scaling is  
334 affected by some uncertainty, because the amplitude recovery depends on the shape of the  
335 anomalies and the idealized synthetic structures we use have a simple shape, which likely  
336 does not represent the geometries of the real anomalies. Furthermore, the level and  
337 distribution of data noise we impose, the choice of model parametrization, and the  
338 approximate forward theory we use to compute arrival-times will all contribute to uncertainty  
339 in the amplitude scaling relationship.

340 As mentioned in Section 1, the low-velocity anomalies in *A1*, *A2* and *A3* can have  
341 different origins, including temperature perturbations, compositional variations, the presence  
342 of melt and/or water, grain size and seismic anisotropy. However, if we assume that the effect  
343 of composition on  $V_P$  and  $V_S$  in the upper mantle is secondary to that of temperature and  
344 gains importance only with increasing depth (Cammarano et al., 2003; Goes et al., 2000),  
345 then we can interpret the seismic features solely in terms of temperature.

346 In order to convert to temperature, we use a smooth  $dV_P/dT$  and  $dV_S/dT$  derivative for a  
347 pyrolytic composition along a  $1300^\circ$  adiabat, after the work of Styles et al., (2011) (Fig. 7).  
348 The full (metamorphic)  $dV/dT$  that includes the effects of phase transitions is also shown.  
349 However, we do not use the latter in our work because such small-scale shifts cannot be  
350 resolved with our dataset. The sensitivity of seismic velocities to temperature decreases very  
351 strongly with depth due to a combination of elastic and anelastic effects (Goes et al., 2004;

352 Styles et al., 2011; Xu et al., 2008). For example, the sensitivity of  $V_P$  ( $V_S$ ) to temperature  
353 along a 1300°C adiabat decreases from  $\sim -1.55\%$  ( $-2.5\%$ ) per 100°C at 100 km, to  $\sim -0.58\%$   
354 ( $-0.96\%$ ) per 100°C at 600 km. Errors in the isomorphic derivative lead to uncertainties in  
355 temperature anomalies of a few tens of degrees (Goes et al., 2000; Cammarano et al., 2003).  
356 In Figure 8 we present the resulting thermal structure for  $P$ - and  $S$ -wave velocity anomalies  
357 before and after applying the scaling correction at 250 km depth. As expected, the original  
358 temperature anomalies  $dT_P$  are higher than  $dT_S$  due to the poorer resolution of the  $S$ -wave  
359 model compared to *IBEM-P18* (higher noise levels, poorer data coverage, and hence greater  
360 regularisation). The scaled temperature anomalies  $dT_P$  and  $dT_S$  are to first order well  
361 correlated in all three regions *A1*, *A2* and *A3*. The highest positive thermal contrasts are  
362 observed below *A2* ( $dT_S \sim 200\text{--}350^\circ\text{C}$ ). Beneath *A1* and *A3*, we find weaker temperature  
363 anomalies, ranging from  $\sim 100$  to  $300^\circ\text{C}$ . In the *A2* case, the high temperatures together with  
364 the high  $a_{S,P}$  suggest that non-thermal effects (e.g., partial melt, volatiles) could likely have a  
365 contribution to explain the seismic anomaly.

366 We also plot the temperature estimates by applying the recovery amplitude from the same  
367 synthetic structures shown in Figure 3, but with two different maximum velocity amplitude:  
368 (i)  $dV_{S,P} = -0.1$  km/s (Fig. S10); (ii)  $dV_{S,P} = -0.6$  km/s (Fig. S11). These test show that if we  
369 change the maximum velocity perturbation we obtain slightly different results in terms of  
370 temperature anomaly range. However, the temperature anomalies in *A1* and *A3* after  
371 resolution correction are still in the range  $\sim 100\text{--}300^\circ\text{C}$ . The biggest differences appear in the  
372 *A2* region where in the  $S$ -wave model the thermal anomaly spans a larger range from  $\sim 100\text{--}$   
373  $300^\circ\text{C}$  (for  $dV_{S,P} = -0.1$  km/s, Fig. S10) to  $\sim 200\text{--}350^\circ\text{C}$  (Fig. 8) and much greater than  $400^\circ\text{C}$   
374 (for  $dV_{S,P} = -0.6$  km/s, Fig. S11). This uncertainty likely results from the poorer resolution of  
375 the  $S$ -wave model below *A2*.

376

## 377 5. Discussion

378

379 Our computation of independent teleseismic *P*- and *S*-wave velocity models allows us to  
380 address key questions concerning the mantle dynamics and physical state of the complex  
381 Iberia-NW Africa region. We will limit our interpretation to the upper mantle where our  
382 spatial resolution is higher.

383

## 384 5.1. Geodynamical context

385

386 The low-velocity anomalies imaged below the Atlas Ranges and the Betic-Rif system by  
387 previous tomographic studies have been attributed to passive mantle upwellings due to edge-  
388 driven convection/small-scale convection (e.g., Kaislaniemi & Van Hunen, 2014; Missenard  
389 & Cadoux, 2012; Ramdani, 1998) and lithospheric delamination (e.g., Bezada et al., 2014;  
390 Levander et al., 2014) or active mantle upwellings connected to the Canary plume (Civiero et  
391 al., 2018; Duggen et al., 2009; Miller et al., 2015). In particular, our previous *P*-wave model  
392 *IBEM-P18* allowed us to propose two main mechanisms for the generation of the low-  
393 velocity anomalies: (1) The mantle upwellings found below *A2* and *A3* rise from ~700 km  
394 depth and are connected to *A1* below the *MTZ*, with all three being fed by a broad lower-  
395 mantle plume (the Central Atlantic plume) imaged in global inversions (e.g., Simmons et al.,  
396 2012); and (2) Quasi-toroidal mantle flow, induced by the Gibraltar slab, drags the hot mantle  
397 material from the sub-slab upwelling (which is sourced below the *MTZ*) laterally and  
398 upwards around either lateral slab edge, to the eastern Rif and Betics. In this latter scenario,  
399 the subducted lithospheric slab, which brings cold lithospheric material down into the mantle  
400 while rolling back westward, has been associated with the prominent high-velocity body  
401 found below the western Alboran Sea. The pronounced low shear velocities that we find in



402 the upper mantle are to first order consistent with the *P*-wave low-velocity anomalies found  
403 in *IBEM-P18* (Civiero et al., 2018). From the base of the *MTZ* downwards, the resolving  
404 power of the *S*-wave model is inferior to that of *IBEM-P18*; therefore, any interpretation of  
405 the structure in the topmost lower mantle is avoided. However, although we cannot interpret  
406 the deeper roots of the anomalies, the structures seem compatible with the presence of  
407 multiple upper-mantle upwellings as suggested by *IBEM-P18*. In line with our previous  
408 geodynamical interpretation, we suggest that the *S*-wave low-velocity features imaged below  
409 *A1*, *A2* and *A3* are due to rising sub-vertical upwellings in the upper mantle originating from  
410 lower-mantle material located below the Canaries, which accumulates and spreads laterally  
411 below the *MTZ*. As indicated by several numerical and laboratory studies (e.g., Kumagai et  
412 al., 2007; Tosi & Yuen, 2011), the 660-km mantle discontinuity resists upward flow of  
413 mantle material, which thus ponds just below the discontinuity, heating the overlaying *MTZ*.  
414 In some locations the thermal boundary layer becomes unstable, developing secondary  
415 thinner instabilities in the upper mantle. The low-velocity features found beneath the Betic-  
416 Rif system may be due to the quasi-toroidal mantle flow created by the roll-back of the  
417 Gibraltar slab (Funiciello et al., 2006; Piromallo et al., 2006; Strak & Schellart, 2014), which  
418 drives the sub-slab hot mantle upwelling material from *A3* around the lateral slab edges. Due  
419 to the location of the low-velocity anomalies approximately below the basal lithospheric  
420 steps south and north of the Alboran Sea, edge-driven convection has been proposed as an  
421 alternative mechanism (e.g., Kaislaniemi and Van Hunen, 2014). However, because we  
422 would expect such anomalies to be of small scale and much more localised beneath the  
423 lithospheric steps, we favour our interpretation. In this framework, the two anomalies below  
424 the Betics and Rif are ultimately also fed by the material ponded below the *MTZ*. Seismic  
425 anisotropy measurements (Buontempo et al., 2008; Díaz et al., 2010; Miller et al., 2013) and  
426 modelling (e.g., Alpert et al., 2013) do support this hypothesis. Fast polarization directions

427 from *SKS* splitting analysis are remarkably parallel to the slow velocity anomalies beneath  
428 the Atlas Ranges, Betics and Rif. This correlates well with the general trend of slab rollback-  
429 induced toroidal mantle flow found around the Gibraltar slab. An alternative interpretation of  
430 the low-velocity feature below *A3*, based on numerical modeling of the anisotropic fabric in  
431 the mantle during subduction, is that it may reflect the presence of anisotropy in the upper  
432 mantle, which can potentially bias the tomographic inversion when isotropic structure is  
433 assumed (Bezada et al., 2016). Since *FMTOMO* cannot account for anisotropy we are unable  
434 to rule out this hypothesis and leave it as an open hypothesis for future work.

435 Next, we will address whether the physical mechanisms proposed for the observed  
436 seismic low-velocity anomalies *A1*, *A2* and *A3* are thermal or chemical in nature, and whether  
437 they may contain a significant proportion of partial melt and/or water.

438

## 439 5.2. Nature of the velocity anomalies

440

441 We compare our temperature model with those of previous studies in the region from  
442 thermobarometry calculations (Thurner et al., 2014), receiver functions (Morais et al., 2015),  
443 waveform analysis (Sun et al., 2014), and geophysical-petrological modelling (Fullea et al.,  
444 2010 and pers. communication).

445 The mantle structure below *A1* shows a  $T_{S,P}$  excess (after scaling) of  $\sim 100\text{-}300^\circ\text{C}$  (Fig. 8d,  
446 j). Our inferred temperatures in the mantle are a good match with the estimates determined by  
447 geodynamic, geophysical and petrological studies for hotspots ( $200^\circ\text{C}$  by Sleep, (1990) and  
448 McKenzie, (1984);  $215 \pm 35^\circ\text{C}$  by Shilling, (1991);  $200^\circ\text{C}$  by Mckenzie & O'nions, (1991);  
449  $200\text{-}300^\circ\text{C}$  by Zhao, (2001);  $162\text{-}235^\circ\text{C}$  by Putirka, (2005)). These values also fall within the  
450 large thermal range suggested by *PP/SS* precursor analyses in the region, which find a  
451 temperature anomaly of  $\sim 100\text{-}300^\circ\text{C}$  for the *MTZ*, strongly depending on the values of the

452 Clapeyron slope (e.g., for  $4.0 \text{ MPa K}^{-1}$  the range is  $84 \pm 47$ – $327 \pm 75^\circ\text{C}$ ) (Saki et al., 2015).  
453 Calculations from primary magma composition reveal that *AI* displays some of the lowest  
454 temperatures of all the ocean-island basalt (*OIB*) lavas,  $\sim 1400$ – $1500^\circ\text{C}$  ( $T$  excess: 100-  
455  $200^\circ\text{C}$ ) (Herzberg & Asimow, 2008), highly consistent with our results. For *OIBs*, there is no  
456 evidence of volatile-enrichment and source fertility. All are associated with thermal  
457 anomalies, which appear to be the only prerequisite for their formation. According to  
458 Korenaga, (2005), large mantle melting anomalies with small  $T$  excess are possible in hotspot  
459 settings such as the Canaries, due to the more fusible nature of the mantle. Several  
460 geophysical and geochemical studies (Day et al., 2010; Klügel et al., 2005; Lodge et al.,  
461 2012; Neumann et al., 2002) found shallow zones of partial melting in the western Canaries,  
462 consistent with depths of magma ponding deduced from petrological analysis and with  
463 seismic and volcanic activity (e.g., Almendros et al., 2007; Carracedo et al., 1999; De  
464 Gonzalez Vallejo et al., 2005; Soler et al., 1984). Our  $a_{S,P}$  for *AI* is within the expected  
465 thermal range ( $a_{S,P} \approx 2.5$ ,  $R_{S,P} \approx 1.5$ ) suggesting a purely thermal origin for the Canaries'  
466 upwelling. We also note that  $dT_S$  is similar to  $dT_P$  and even somewhat lower; thus, we tend  
467 to discard the hypothesis of high melt productivity at upper-mantle depths. Rather we suggest  
468 that local and shallow signatures of melt may exist, but they are not resolvable at the length-  
469 scale of our inversion.

470 The  $T_{S,P}$  excess beneath *A2* is within the range  $\sim 200$ – $350^\circ\text{C}$ , with some peaks of  $dT >$   
471  $400^\circ\text{C}$  below the most prominent basaltic centres of the High and Middle Atlas (Fig. 8e, k).  
472 Our findings are consistent with values derived from modelling of hotspot tracks beneath the  
473 continental lithosphere ( $\sim 300^\circ\text{C}$ , Yang & Leng, 2014). In addition, results from seismic  
474 waveform analysis combined with geodynamic modeling reveal a strong  $T$  excess of  $\sim 350 \pm$   
475  $90^\circ\text{C}$  in this region (for dry mantle), which is inferred to be solely of thermal origin, although  
476 a small quantity of partial melt and volatiles may be present (Sun et al., 2014).  $dT_S$  estimates

477 appear to be slightly higher than  $dT_P$  in the innermost parts of the low-velocity anomaly (see  
478 Fig. 8e, k). As our tomographic models are insensitive to the effects of local melt signatures  
479 and to rapid changes in velocity over short vertical distances, we do not exclude the  
480 possibility that other physical factors, such as small melt pockets and/or water-rich layers,  
481 may be present in small quantities and/or at small spatial scales in the upper mantle.  
482 Interestingly, the  $a_{S,P}$  ( $\approx 3.8$ ) and  $R_{S,P}$  ( $\approx 2.2$ ) values found for this area are on the higher side  
483 of the thermal range and may be attributed to other non-thermal effects (e.g., different  
484 compositions, volatiles, melt). A contribution of partial melt at shallow depths would also be  
485 consistent with magnetotelluric results (Anahnah et al., 2011) and geochemical studies  
486 (Duggen et al., 2009, 2005; El Azzouzi et al., 2010; Geldmacher et al., 2005; Hoernle et al.,  
487 1999, 1995; Lundstrom et al., 2003; Lustrino and Wilson, 2007). Receiver-function analyses  
488 from Morais et al., (2015) revealed a low regional  $S$ -velocity-layer atop the 410-km  
489 discontinuity. One hypothesis that explains these observations is that mantle upwelling across  
490 the 410-km discontinuity leads to water release and melting atop the 410-km discontinuity  
491 (Karato et al., 2006). The melt could be buoyant and flow upwards, or dense and remain  
492 stalled above the 410-km discontinuity.

493 The A3 region exhibits temperature anomalies from  $dV_P$  and  $dV_S$  on the order of 100-  
494 300°C (Fig. 8f, l). Few studies have tried to infer the temperature in the sub-slab domain. The  
495 petrologic results from Fullea et al., (2010) reveal a temperature of less than 1430°C, which  
496 is consistent with the lower bound of the  $T$  excess that we find. However, from our  $P$ - and  $S$ -  
497 wave velocity models we suggest that the hot mantle rising from below A3 is dragged around  
498 the retreating Gibraltar slab through quasi-toroidal mantle flow; therefore, the mantle  
499 material below A3 and below the Rif-Betics system should be in a similar thermal range. If  
500 so, our findings are in line with thermobarometry studies (Turner et al., 2014) which find  
501 mantle temperature estimates increasing from 1350°C to a maximum of 1430°C ( $\sim 50$ -130°C

502  $T$  excess), moving from the Rif to the Betics. This observation, together with the results from  
503 the residual analysis ( $a_{S,P} = 2.1$ ,  $R_{S,P} = 1.2$ ), thermal reconstruction and geodynamic  
504 interpretation, lead us to propose a weak thermal signature for the mantle upwellings in the  
505 A3 region and below the Rif and Betics.

506

### 507 5.3. Comparison with surface volcanism, heat flow and lithospheric thickness

508

509 Volcanism occurred across the Ibero-western Maghreb region from the late Cretaceous to  
510 the present (Lustrino & Wilson, 2007). According to Missenard et al., (2006), a linear *NE-SW*  
511 trend of anorogenic magmatism and significant seismicity, together with thinned lithosphere  
512 and uplifted topography, the so-called Moroccan Hot Line, crosses the A2 region, and cuts  
513 the eastern Rif. The magmatism continues towards southeastern Iberia and extends from the  
514 northern Moroccan passive margin, into the Alboran Sea and eastern Betics. The composition  
515 of these units changes in time and space from calc-alkaline in southern Iberia (comprising the  
516 Calatrava volcanic field) and the Alboran Sea during the Miocene, to alkaline in the Atlas  
517 Ranges during the Pliocene and Quaternary (Duggen et al., 2004; Hoernle et al., 1995). All  
518 these volcanic centres are often viewed within the wider context of *Na*-rich intra-plate  
519 magmatism, which characterizes the whole Euro-Mediterranean region and extends towards  
520 the eastern Atlantic Ocean (e.g., Lustrino & Wilson, 2007; Wilson & Downes, 1991, 2006).  
521 To the west of the Strait of Gibraltar, no clear evidence of volcanic activity has been found so  
522 far above the sub-slab low-velocity anomaly. A recent geophysical study on the Gulf of  
523 Cadiz domain (Neres et al., 2018) imaged a magmatic intrusion offshore below the  
524 Guadalquivir Bank, which may represent the southernmost expression of the Upper-  
525 Cretaceous Alkaline magmatic event that affected both onshore (Sintra-Sines-Monchique  
526 igneous complexes) and offshore West Iberia. In the Atlantic domain, A1 presents a long

527 volcanic record ranging from the oldest intra-plate manifestations occurring during the Late  
528 Cretaceous-Paleocene at just a few volcanic centers, to the recent magmatic activity in almost  
529 all the islands (e.g., Anguita & Hernán, 2000; Carracedo et al., 1999).

530 Surface heat flow (*SHF*) may play a role in generating these magmatic pockets and  
531 related temperature variations in the mantle and lithosphere can regionally change the heat  
532 input at the base of the crust (della Vedova et al., 2001). Rising magmas transport heat into  
533 the lithosphere by advection and can cause partial melting of lithospheric domains with lower  
534 than normal solidus temperatures in areas of active deformation. However, we should keep in  
535 mind that *SHF* measurements can be strongly scattered due to environmental factors, e.g.  
536 water circulation, paleo-climatic variations and can register large changes over short  
537 distances due to local geology and topography (Sclater et al., 1980).

538 Interestingly, the distribution of magmatic fields in the Gibraltar Arc System, Atlas  
539 Ranges and Canaries has a first order correlation with the low-velocity anomaly regions and  
540  $T$  excess found in our study; conversely, no recent volcanic activity occurs where cold high-  
541 velocity anomalies are imaged. Furthermore, correlating with extensive volcanic activity, a  
542 high averaged *SHF* value and a shallow lithosphere asthenosphere boundary (*LAB*),  
543 compared to the surrounding areas, are found. To illustrate this, we plot, on Figure 9, the  
544 magmatic centers which have been active in the last 75 Myr (from Lustrino & Wilson, 2007)  
545 together with the average of all available surface heat flow measurements (provided by J.  
546 Fullea, pers. comm.) for regions *A1*, *A2* and *A3*. We also include the averaged *SHF* values for  
547 eastern Rif and Betics, where the upward mantle flow induced by the slab rollback has been  
548 suggested to explain the vigorous Cenozoic volcanism in the Betic-Rif area. Below the top  
549 map in Figure 9, we show the *LAB* depth map (Fullea et al., 2015, 2010), our *S*-wave  
550 tomographic model and the thermal structure for regions *A1*, *A2* and *A3* derived from the *S*-  
551 wave tomography. Below the main volcanic fields, the *SHF* is high on average ( $> 60$

552  $\text{mWm}^{-2}$ ), although some important regional variations do occur. *A1* is characterized by *SHF*  
553 values of around  $73 \text{ mWm}^{-2}$  (see Canales & Dañobeitia, 1998), which are associated with  
554 massive volcanic outpourings (Fig. 9a) and moderate lithospheric thinning (*LAB* depth of 90–  
555 100 km) (Fig. 9b). In *A2*, *SHF* values (where available) are on average of  $63 \text{ mWm}^{-2}$  (Fig.  
556 9a), in line with the strong low-velocity and positive thermal anomalies found in our study,  
557 and correlate well with the lithospheric thinning observed (80-120 km depth) (Fig. 9b). A  
558 more recent study, however, found slightly higher *SHF* values of about  $80 \text{ mWm}^{-2}$  below the  
559 northeastern termination of the Middle Atlas domain (Chiozzi et al., 2017). Similarly, a local  
560 measurement of  $86 \text{ mWm}^{-2}$  has been found in the southern High Atlas by Ramdani, (1998)  
561 suggesting a lithosphere thickness of 50 km. When interpreted in light of these independent  
562 data, the low-velocity anomalies observed below *A2* may further indicate anomalously hot  
563 upper mantle with local areas of partial melting. The highest *SHF* values are observed in the  
564 eastern Rif (=  $\sim 89 \text{ mWm}^{-2}$ ) and eastern Betics ( $\sim 82 \text{ mWm}^{-2}$ , with peaks  $> 100 \text{ mWm}^{-2}$  in  
565 the eastern Alboran Basin), where the *LAB* is mapped as shallow as 50-150 km (Fig. 9a). This  
566 pattern of spatial coincidence between peaks in magmatic activity and strong low-velocity  
567 anomalies as imaged by our tomographic model is in good agreement with the existence of  
568 high mantle temperatures. Finally, below *A3* a few *SHF* values have been measured in the  
569 range  $50\text{-}60 \text{ mWm}^{-2}$ , which denote a moderate heat flow (Fig. 9a). According to the findings  
570 in Fullea et al., (2010) the *LAB* in *A3* is shallower than below the Alboran Sea, but much  
571 deeper ( $\sim 170\text{-}180$  km depth) compared to the *A2* region (Fig. 9b). However, a more recent  
572 geophysical study found contradictory results, which indicate a significant lithospheric  
573 thinning below *A3* (60-70 km depth), where low velocities are imaged (Palomeras et al.,  
574 2014). Further investigation of the *A3* domain is needed to conclusively understand why the  
575 upper-mantle upwelling that we find does not show a clear volcanic expression at the surface.

576 Overall, the spatial distribution of surface volcanism, *SHF* and depth of the *LAB* are in  
577 good agreement with the presence of upper-mantle thermal upwellings, which interact with  
578 the Gibraltar slab-induced mantle flow, facilitating decompression melting at sub-lithospheric  
579 depths and generating extensive magmatic provinces in Iberia and *NW* Morocco.

580

## 581 6. Conclusions

582

583 We combined *S*- and *SKS*-wave data from the same networks used in our recent *P*-wave  
584 tomography (Civiero et al., 2018) to compute a high-resolution shear-wave tomographic  
585 model of the Ibero-western Maghreb region. The *S*-velocity features are well resolved from  
586 the surface to the base of the transition zone (extending somewhat less deep than the *P*-wave  
587 resolution).

588 Our *S*-wave images show sub-vertical low-velocity anomalies below the Canaries, the  
589 Atlas Ranges and the Gibraltar Arc extending throughout the upper mantle. Although the  
590 resolution below the *MTZ* is poor, the *S*-wave anomaly trends strongly resemble those of the  
591 *P*-wave velocities, which suggest a lower-mantle origin. From our findings we interpret the  
592 structure as being mainly the signature of mantle upwellings rising from below the *MTZ* and  
593 sourced from hot material associated with the lower-mantle Central Atlantic plume. Other  
594 strong low-velocity anomalies are imaged below the eastern Rif and Betics and may represent  
595 the result of the interaction between the retreating Gibraltar slab and the mantle upwelling  
596 behind it, which pushes the hot mantle material around the lateral slab edges.

597 The conversion of the low-velocity anomalies below the Canaries, Atlas Ranges and  
598 Gibraltar Arc to temperature, together with analysis of arrival-time residuals, suggest that  
599 moderate temperature excesses of the order of 100 -350°C can explain the origin of these  
600 upwellings in the upper mantle. Surface heat flow observations, *LAB* depth estimates and



601 recent surface volcanism are consistent with the presence of mantle upwellings below the  
602 Canaries and Atlas Ranges. Current data availability does not allow for more definitive  
603 conclusions concerning the offshore Gibraltar Arc domain.

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 1117 **Figure 1.** A. Geological units of the Ibero-western Maghreb region (from  
 1118 <https://pubs.usgs.gov/dds/dds-060/>), Eurasia-Africa plate boundary (brown, Bird, 2003),  
 1119 potentially active faults from the SHARE database (black, Basili et al., 2013; Vilanova et al.,  
 1120 2014) and high-resolution fault traces (including debated faults) (blue, García-Mayordomo et  
 1121 al., 2012; Zitellini et al., 2009; Cabral, 2012). The geographic features cited in the text are  
 1122 indicated in black. GA: Gibraltar Arc; GB: Goringe Bank. B. Location of the seismic stations  
 1123 used in this study. Colours and symbols mark the different seismic networks. The six labelled  
 1124 stations are those for which residuals are shown in Figure S1. C. Distribution of the  
 1125 teleseismic events used (yellow dots).

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1128 **Figure 2.** Checkerboard resolution tests for our tomographic study, using alternating positive  
 1129 and negative velocity anomalies of ~200 km width and  $\pm 0.5$  km/s in amplitude separated by a  
 1130 narrow region of zero perturbation. Velocity perturbations are plotted relative to the 3D  
 1131 starting model. a, b) Input model at 250 km and 500 km depth respectively. c, d) Output S-

1132 velocity structure at 250 km and 500 km depth respectively. The raypaths and inversion  
 1133 parameters used are the same as in the inversion of actual data. Gaussian noise of 0.4s is  
 1134 added to the synthetic dataset to mimic that of the field dataset. Crustal structure is light grey-  
 1135 shaded. Regions with no piercing points are shaded darker grey. Black lines show coastlines.  
 1136 (e, f, i, j) Vertical cross-sections oriented east-west (e,i) and south-north (f, j), through the  
 1137 input model (orientations of the profiles are shown in depth slice a.). (g, h, k, l) Vertical  
 1138 cross-sections through the recovered model. These tests suggest a good resolution through the  
 1139 upper mantle for most of the region of interest (the oceanic domain north of the Canaries and  
 1140 the western African craton are excluded from our interpretation).

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1143 **Figure 3.** Structural resolution test, using synthetic vertical low-velocity structures below the  
 1144 Canaries (*A1*), the Atlas Ranges (*A2*) and the Gibraltar Arc (*A3*) (-0.3 km/s amplitude).  
 1145 Velocity perturbations are plotted relative to the 3D starting model. (a) Map view of the input  
 1146 model at 200 km depth. (b) Map view of the recovered model at 200 km depth. Boxes *A1*, *A2*  
 1147 and *A3* delimit the regions – Canaries, Atlas Ranges and Gibraltar Arc – for which we  
 1148 compute temperature conversions. (c) Input model through vertical cross-sections oriented  
 1149 west-east below *A2* (the same input structure is located below *A1* and *A3*). The profiles are  
 1150 shown in depth slice a.). (d, e, f) Vertical cross-sections through the recovered models. The  
 1151 raypaths and inversion parameters used are the same as for the inversion of actual  
 1152 observations. Gaussian noise with a standard deviation of 0.4s is added to the synthetic  
 1153 dataset to mimic the noise in the observations. Crustal structure is grey-shaded. The  
 1154 amplitude recovery of the vertical bodies below *A1* (d) and *A3* (f) in the depth range 0-200  
 1155 km is around 35%. The amplitude recovery of the vertical structure below *A2* (e) is around  
 1156 20% down to 450 km depth.

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1159 **Figure 4.** Depth slices through the tomographic *S*-wave model at depths between 70 and 730  
 1160 km. Velocities are plotted relative to a 1D laterally averaged depth-dependent version of the  
 1161 starting model. Regions with no piercing points are shaded grey. Black lines show coastlines.  
 1162 These maps reveal velocity anomalies similar to those imaged in *IBEM-P18* (see Figure S5  
 1163 for comparison). The anomalies discussed in the text are the prominent low-velocity regions  
 1164 below *A1*, *A2* and *A3* (boxes delimiting the regions are indicated in black in a), which extend  
 1165 through the *MTZ*, and the two low-velocity anomalies below eastern Rif and Betics that  
 1166 surround the high-velocity body below the Alboran Sea.

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1168 **Figure 5.** Vertical cross-sections through our *S*-wave model. The orientation of the profiles is  
 1169 indicated with black lines in the 250 km depth slice. Topography profiles (from Smith and  
 1170 Sandwell, 1997) and geographic names are shown above each cross-section. Velocities are  
 1171 plotted relative to a 1D laterally averaged depth-dependent version of the starting model.  
 1172 Regions with no piercing points are shaded grey. The most prominent high-velocity feature is  
 1173 that below the Alboran Sea (profiles a, b and c). The most relevant low-velocity anomalies,  
 1174 which are discussed in the text, are located below *A3* (profile a), below *A2* (profiles d and g)  
 1175 and below *A1* (profiles e, d and g). These features extend through the *MTZ* and appear  
 1176 connected at lower-mantle depths.

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1178 **Figure 6.** *P*-wave (from Civiero et al., 2018) versus *S*-wave relative travel-time residuals  
 1179 (with associated errors) for common earthquakes and stations, for stations located in regions  
 1180 *A1*, *A2* and *A3* (Figure 3). The values of the slopes,  $a_{S,P}$ , fall within the thermal range for all  
 1181 the selected regions (*A1* =  $\sim 2.5 \pm 0.1$ ; *A2* =  $\sim 3.8 \pm 0.2$ ; *A3* =  $\sim 2.1 \pm 0.1$ ).

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**Figure 7.** Profiles of the  $dV_P/dT$  (blue line) and  $dV_S/dT$  (red line) derivatives that we use to convert the recovered velocity to temperature anomalies. The thick profiles corresponds to the smoothed isomorphic  $dV_P/dT$  and  $dV_S/dT$  derivatives, which do not account for effects of phase-boundary topography. The dashed blue and red profiles are respectively the metamorphic  $dV_P/dT$  and  $dV_S/dT$  derivatives, which include the effects of phase transitions. The derivatives were computed along a  $1300^\circ\text{C}$  adiabat for a pyrolite composition using mineral parameters from the database stx08 (Xu et al., 2008), with composite attenuation model Qg (above 400 km) (Van Wijk et al., 2008) and Q6 (below 400 km) (Goes et al., 2004). For our conversion, we use the isomorphic derivatives because the teleseismic body-wave tomography cannot resolve localized phase-boundary anomalies.

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**Figure 8.** Horizontal slices at 250 km depth showing the thermal anomalies obtained from the conversion of  $P$ - and  $S$ -wave velocities using the  $dV/dT$  curves in Figure 7 below the upwellings within  $A1$ ,  $A2$  and  $A3$ . a, b, c) Original  $T$  excesses derived from the  $P$ -wave velocities in *IBEM-P18* (without accounting for the estimated amplitude recovery). d, e, f) Scaled  $T$  excesses derived from the  $P$ -wave velocities in *IBEM-P18* using the estimated amplitude recovery from the resolution test in Figure S9. g, h, i) Original  $T$  excesses derived from the  $S$ -wave velocities (without accounting for the amplitude recovery). j, k, l) Scaled  $T$  excesses derived from the  $S$ -wave velocities using the estimated amplitude recovery from the resolution test in Figure 3. The scaled  $dT_P$  and  $dT_S$  (d, e, f and j, k, l) are to first order well correlated in all three regions  $A1$ ,  $A2$  and  $A3$ . Regions with no piercing points are shaded white. Grey lines show coastlines. The spacing between the contours is  $100^\circ\text{C}$ .

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**Figure 9.** a) Topography map of the Ibero-western Maghreb region with yellow triangles indicating volcanic centres from Late-Cretaceous to present (modified from Lustrino & Wilson, 2007) and the averaged  $SHF$  values for  $A1$ ,  $A2$ ,  $A3$ , the Rif and Betics written in black (from J. Fullea, pers. comm.). b)  $LAB$  depth map (Fullea et al., 2010, 2015). c) Our  $dV_S$  model ( $dV_S$  in %) at 150 km depth. d) Our scaled thermal structure for regions  $A1$ ,  $A2$  and  $A3$  (delimitated in Figure 4a) from the  $S$ -wave velocities at 150 km depth. The regions where no data are available are masked out in grey. Regions with no piercing points are shaded white. The positive temperature anomalies correlate well with the low-velocity features  $A1$  and  $A2$ , with the thinning of the lithosphere, and at the surface with the most recent magmatic centres and highest  $SHF$  values observed. The Rif-Betics system also shows a good spatial coincidence with shallow  $LAB$  depth,  $SHF$  values and volcanic strips. No recent volcanism and a moderate  $SHF$  are found for  $A3$ , which may result from the presence of the subducted slab.

Figure 1.



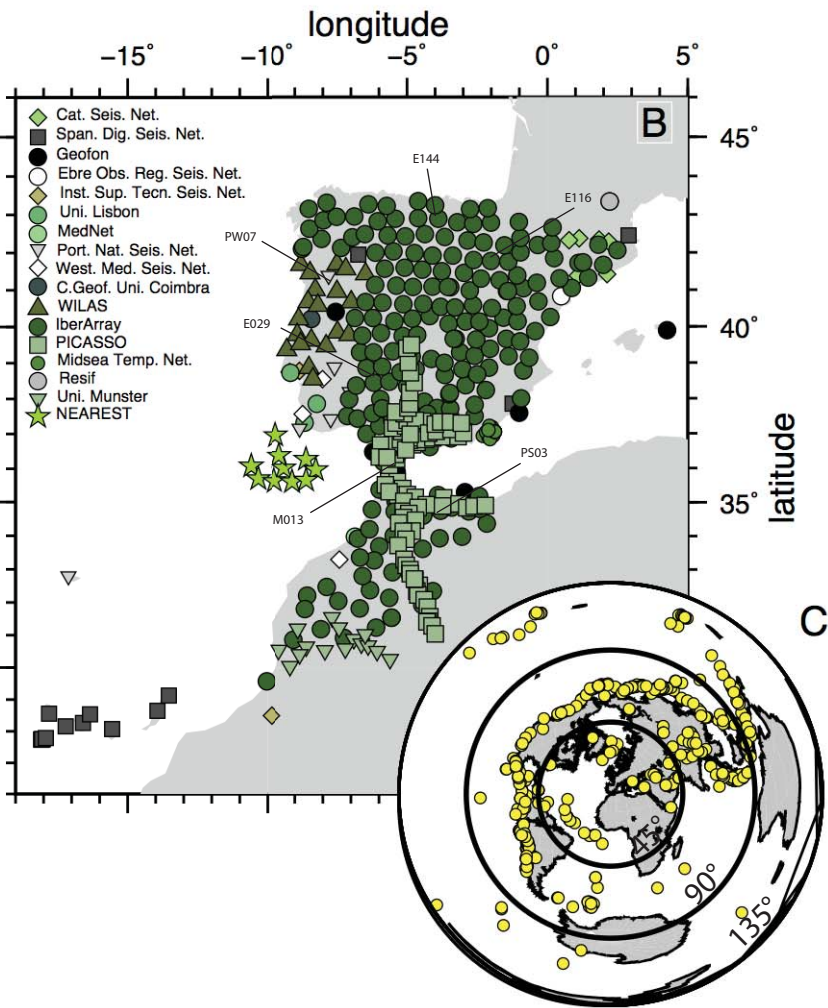
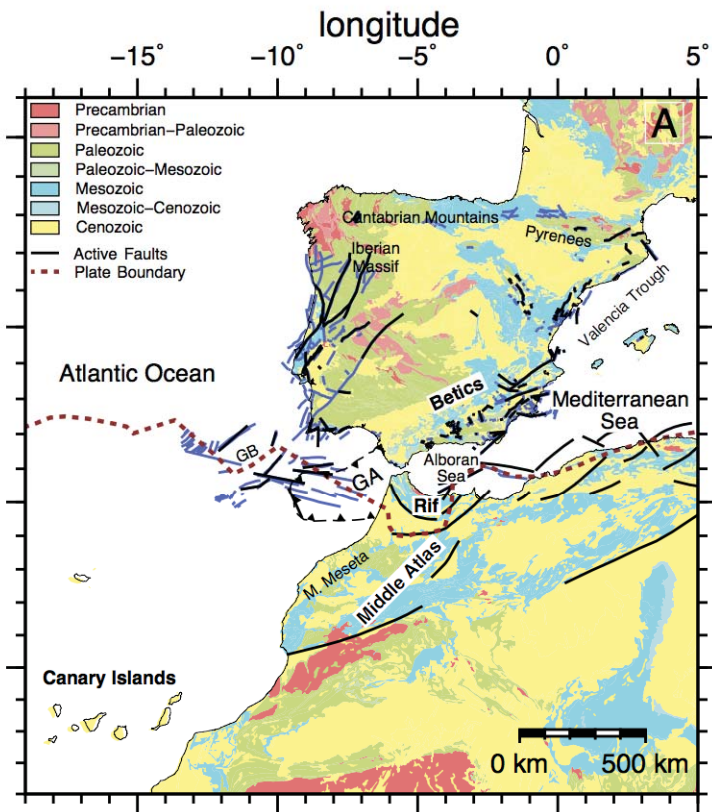


Figure 2.

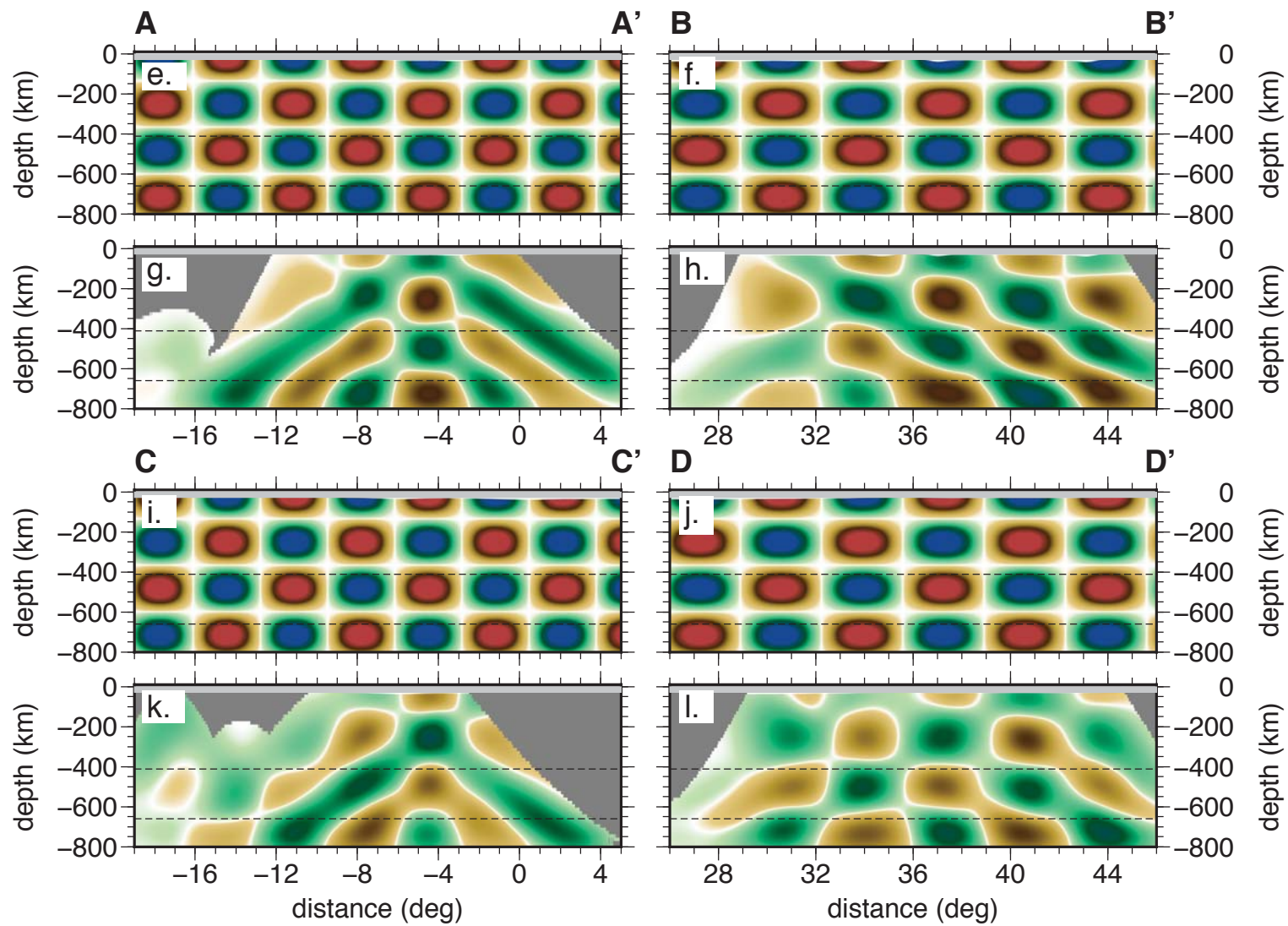
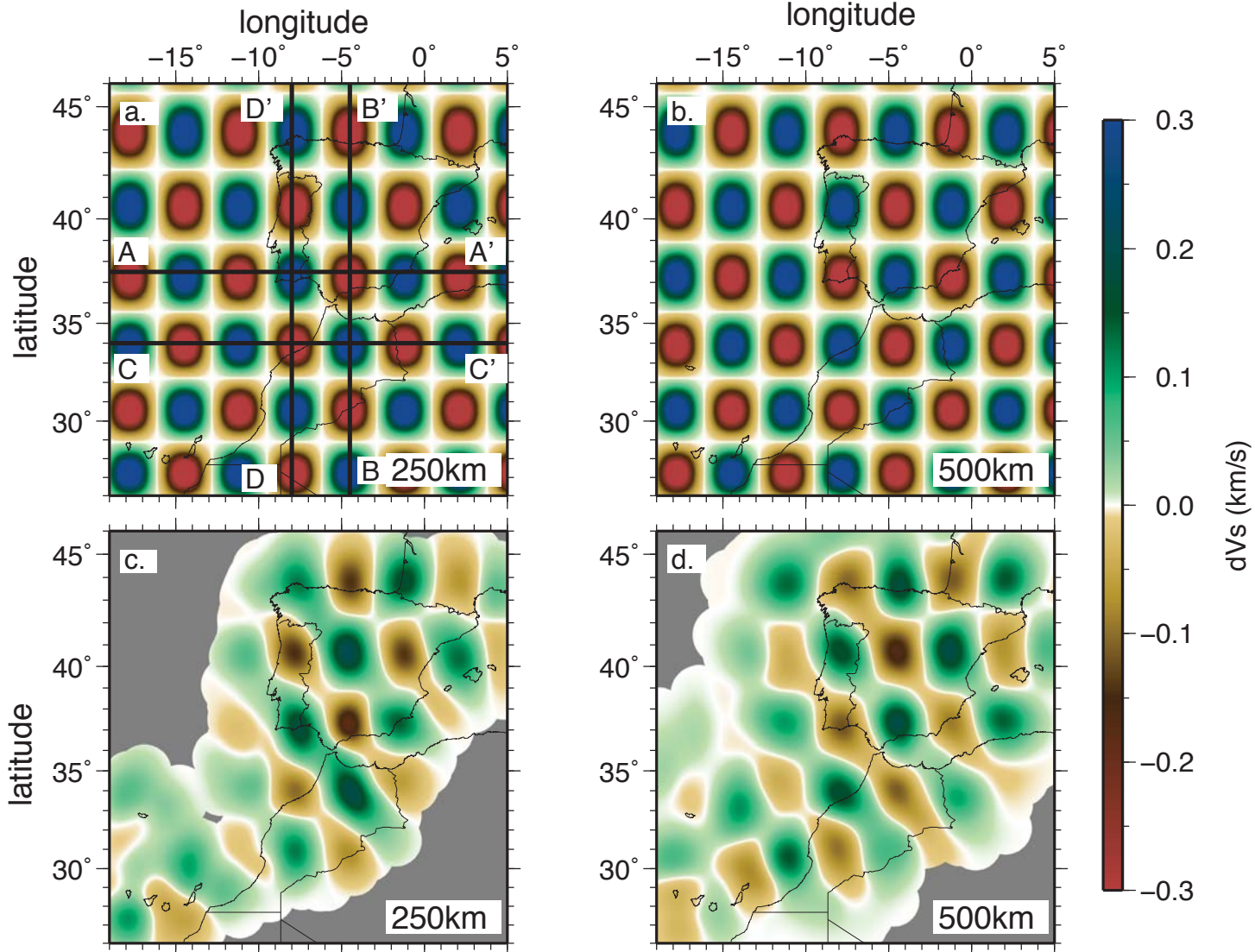
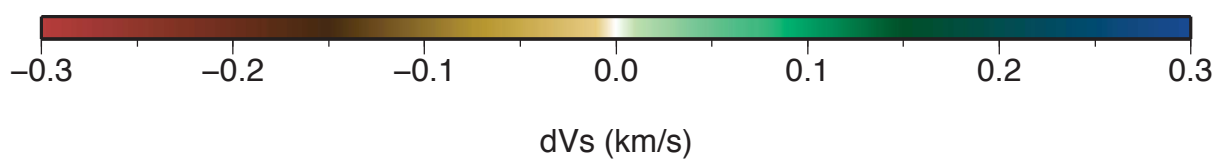
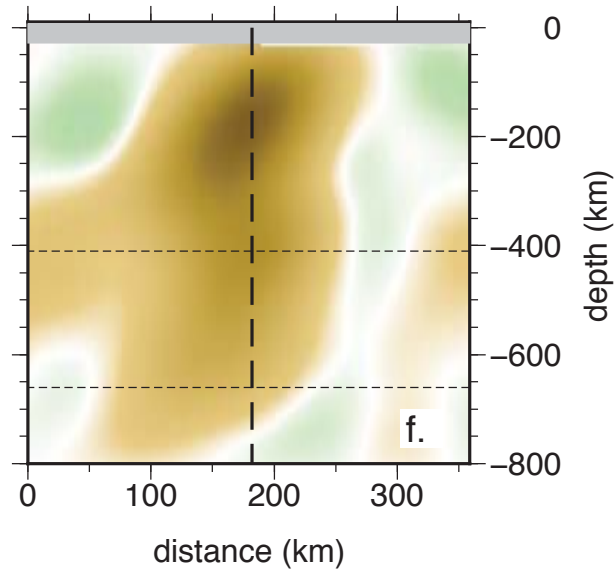
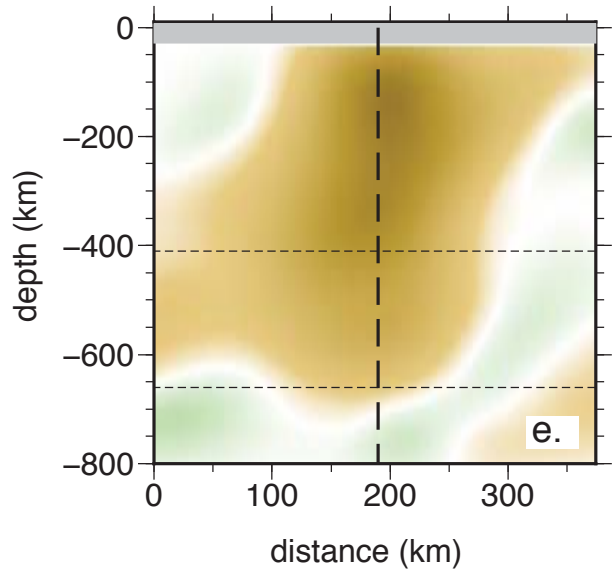
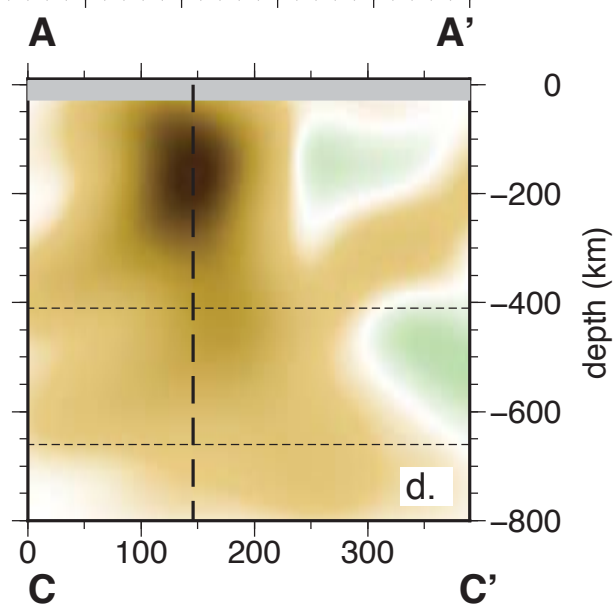
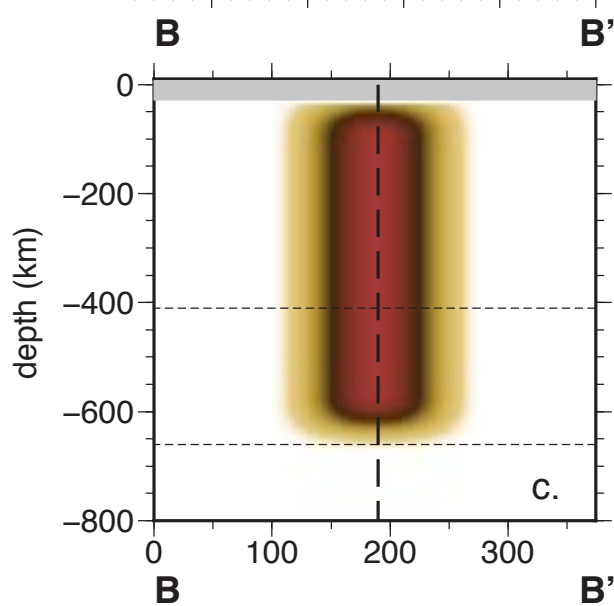
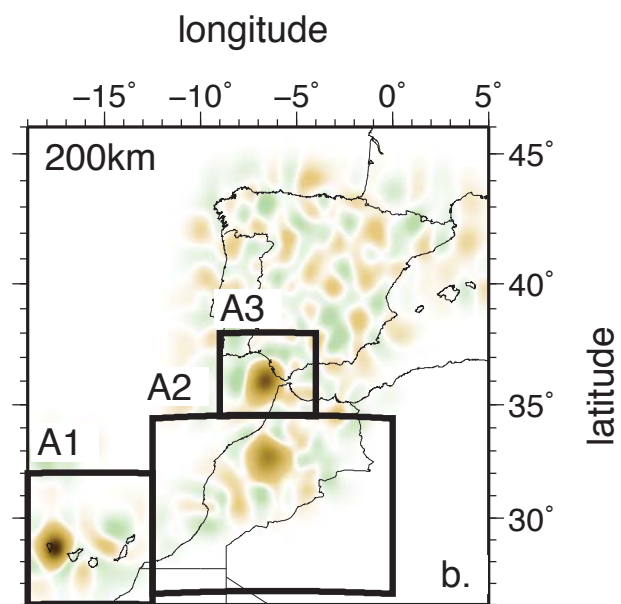
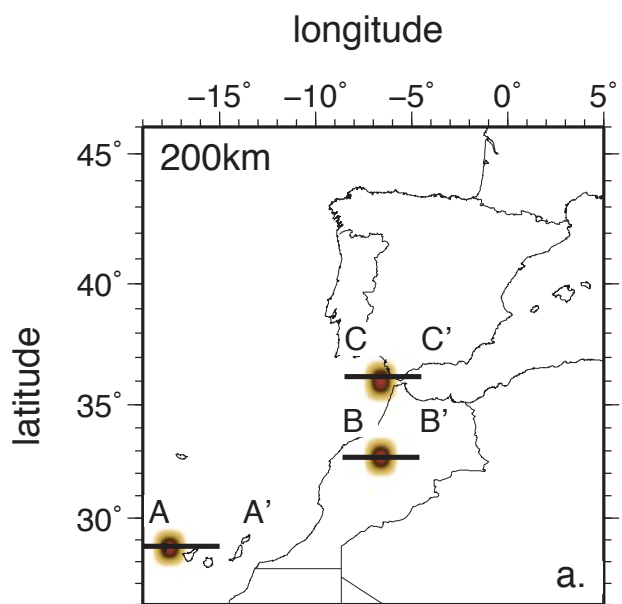


Figure 3.



**Figure 4.**

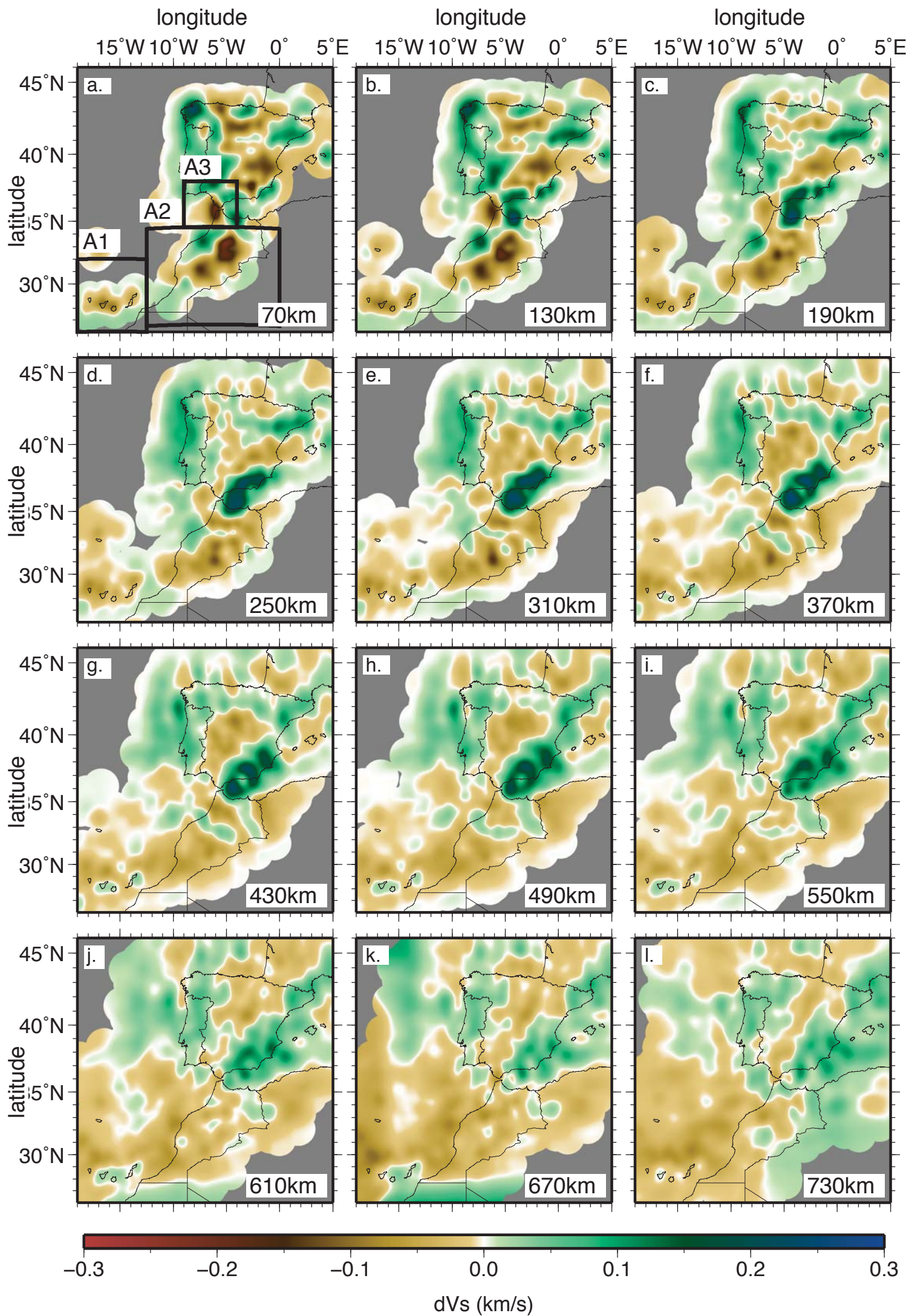


Figure 5.



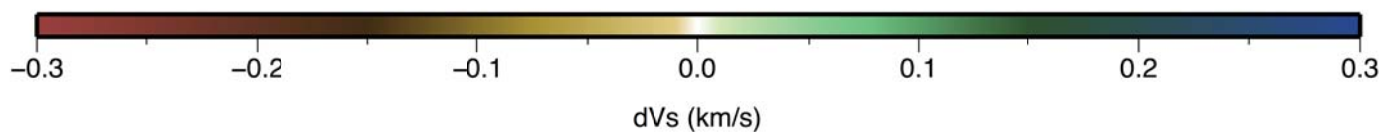
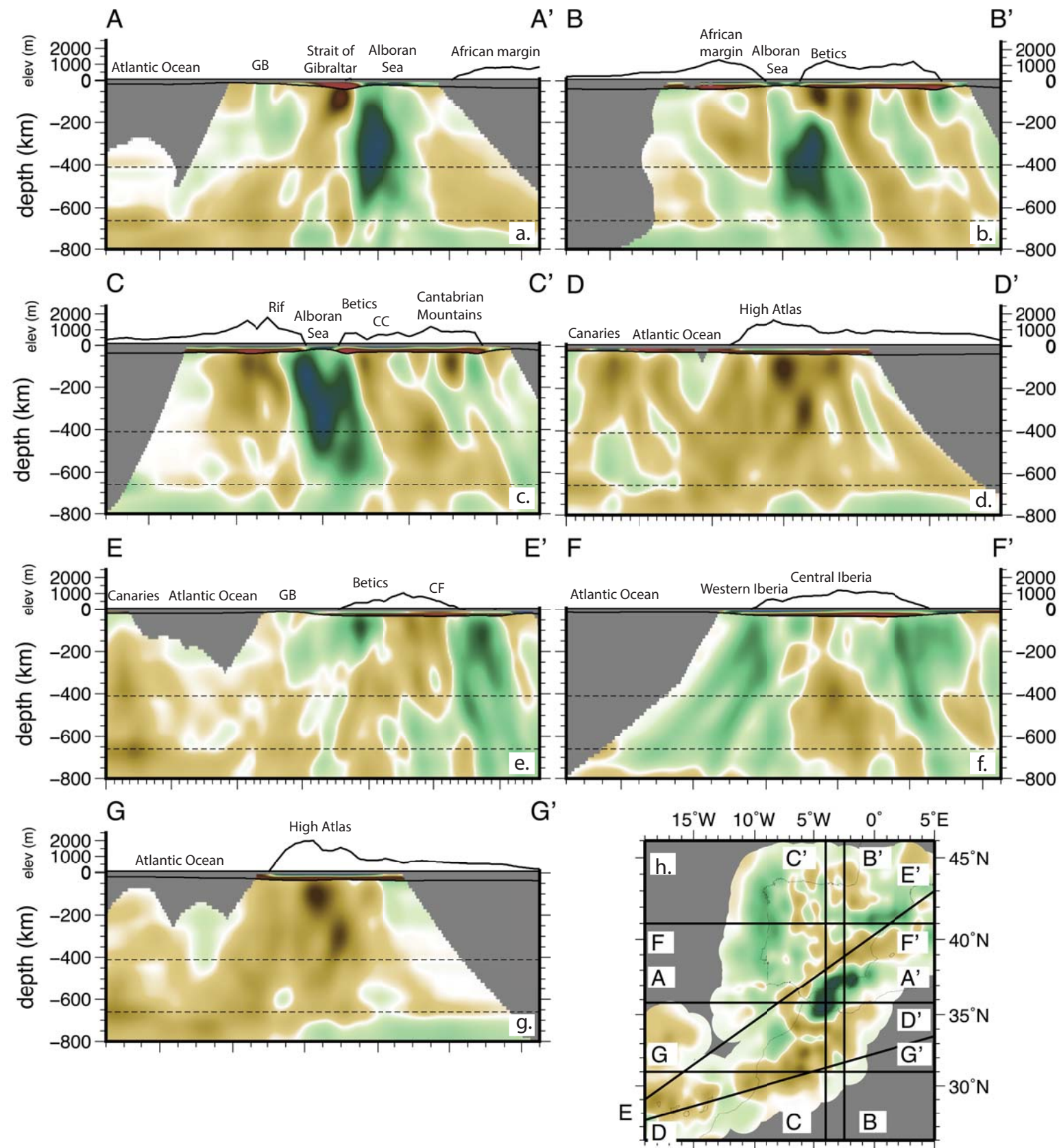
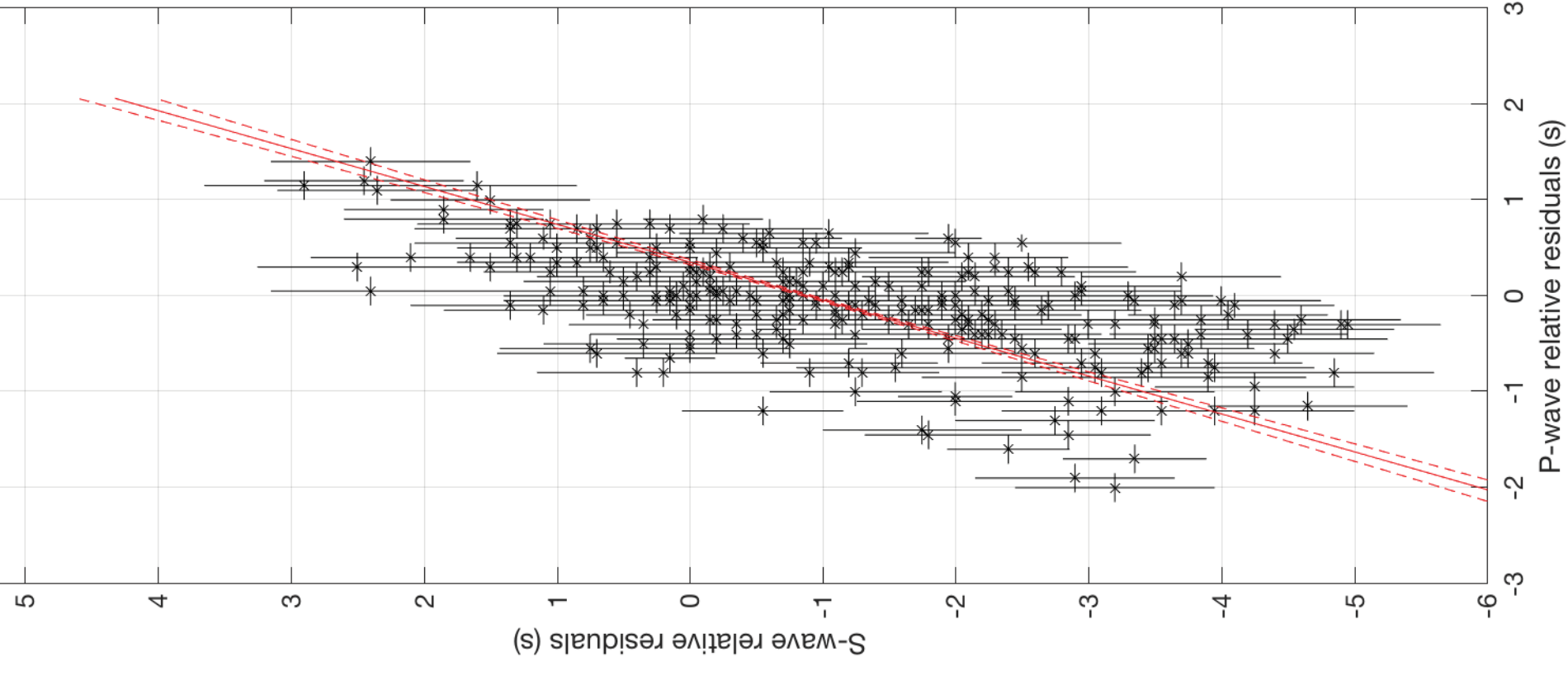


Figure 6.

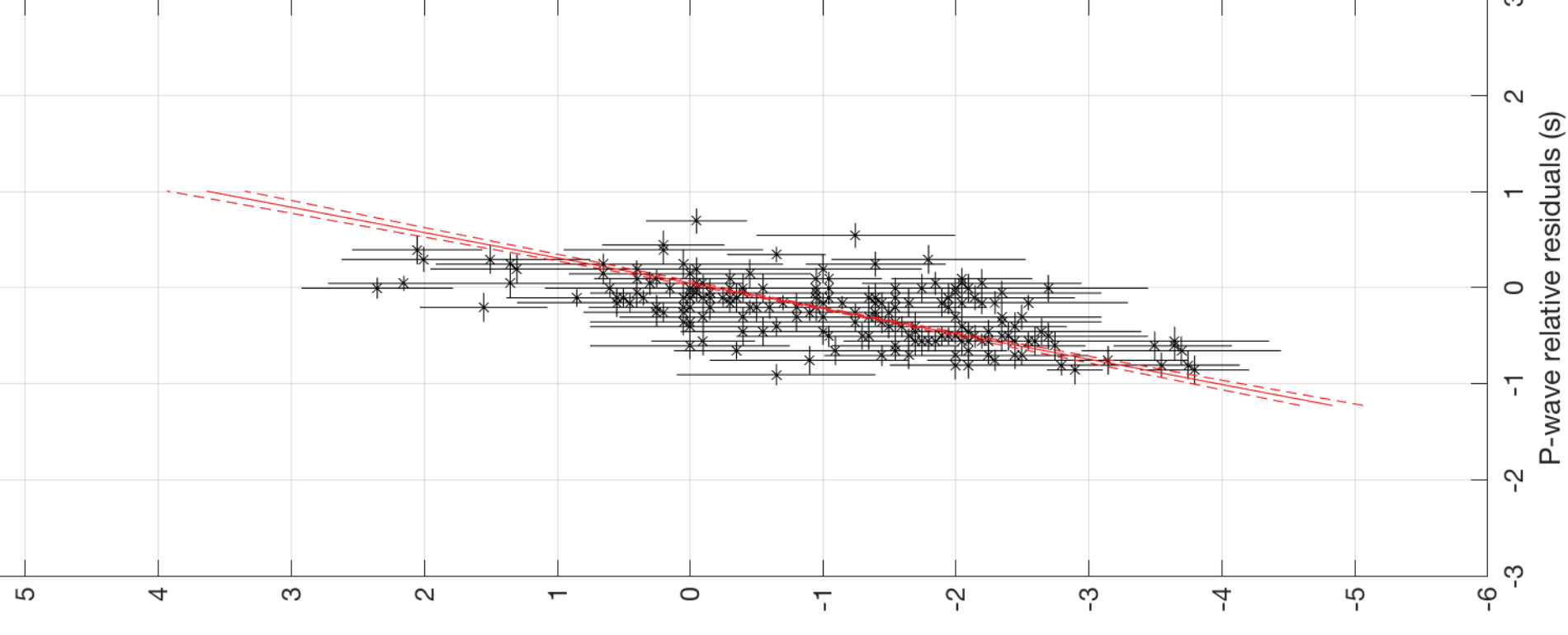
$$Y_6 = (2.5294 \pm 0.13751) * X + (-0.88779 \pm 0.040636)$$

A1



$$Y_6 = (3.7965 \pm 0.23144) * X + (-0.20335 \pm 0.071037)$$

A2



$$Y_6 = (2.1528 \pm 0.10139) * X + (-0.59173 \pm 0.062726)$$

A3

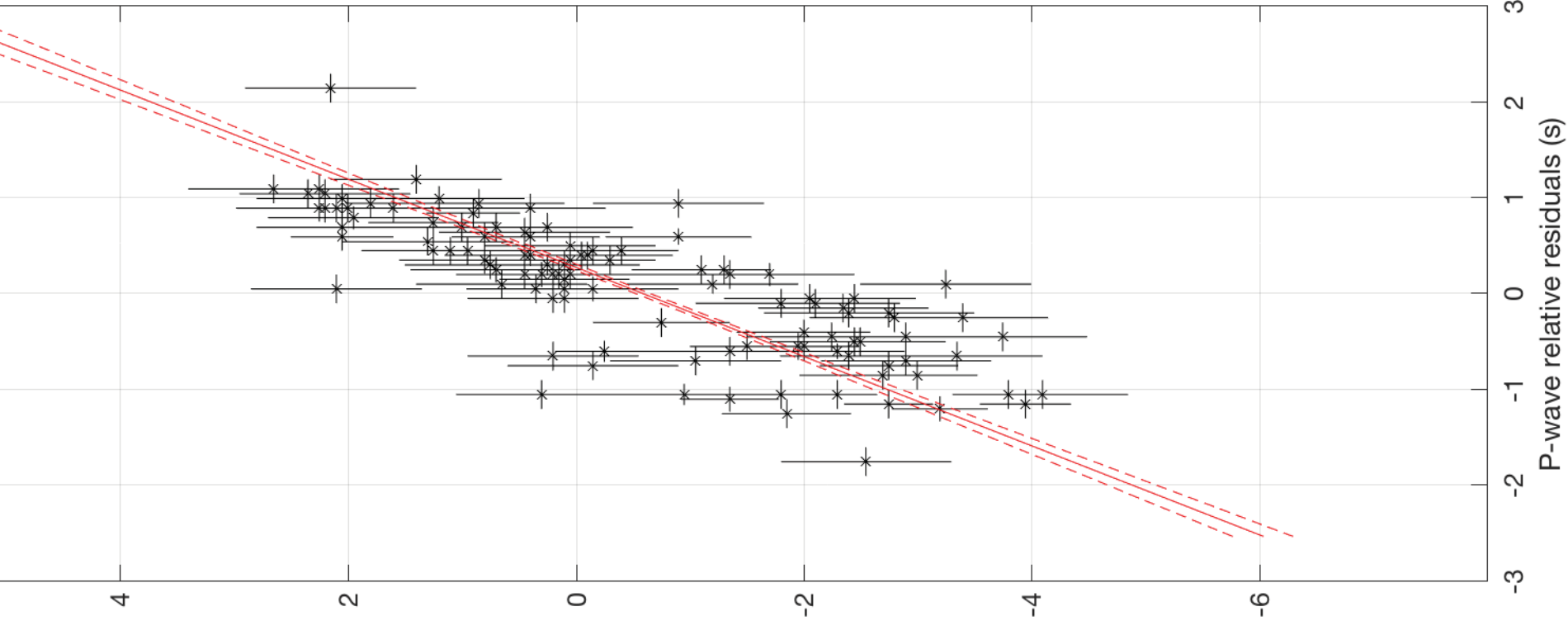


Figure 7.

$d \ln V / dT$  (%/100K)

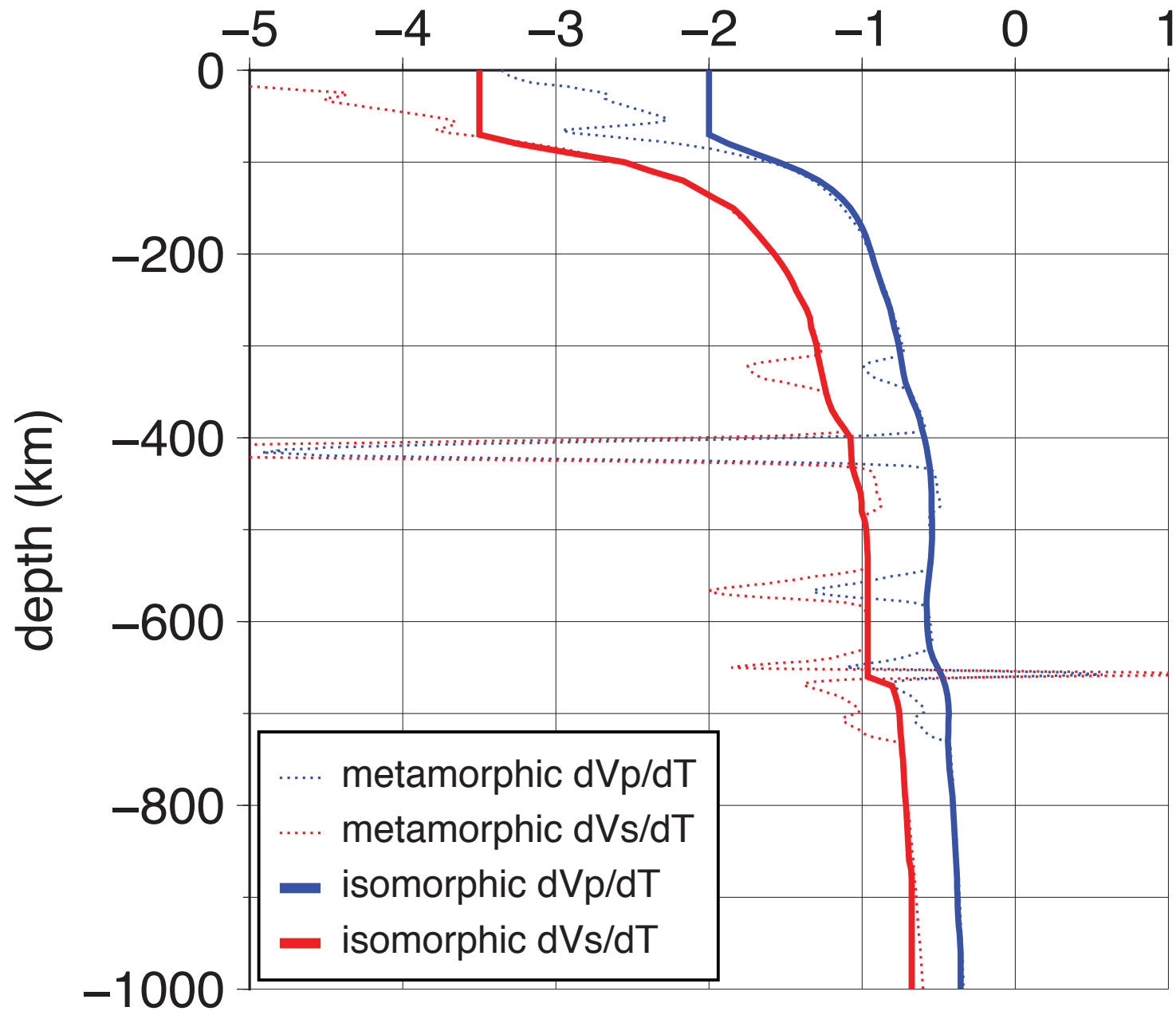
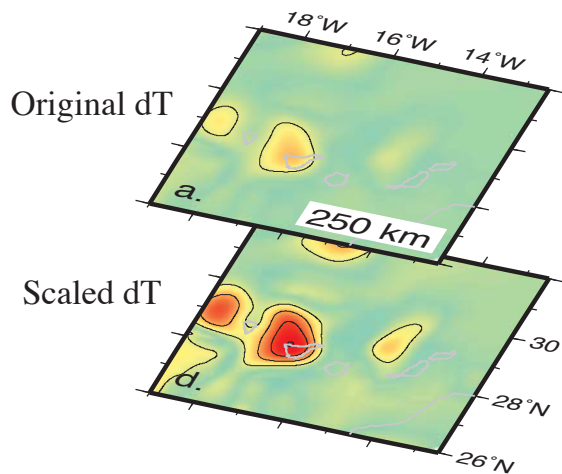


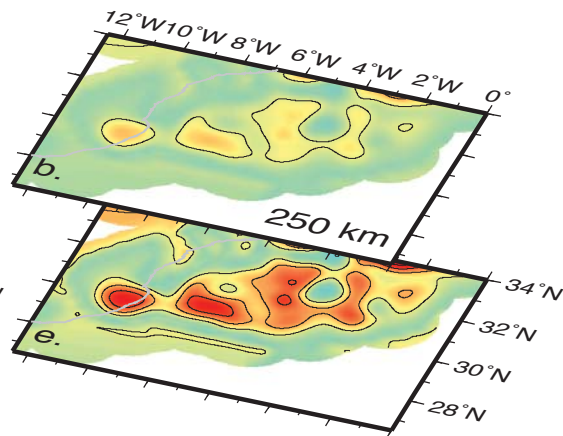
Figure 8.

*P model*

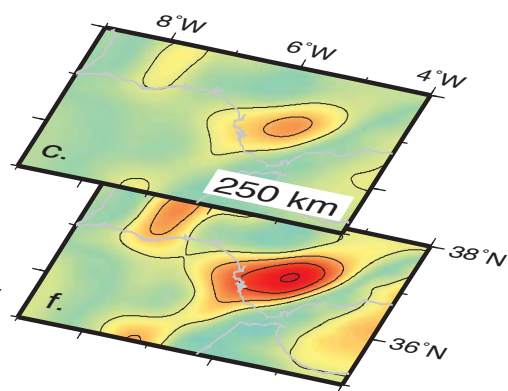
**A1**



**A2**

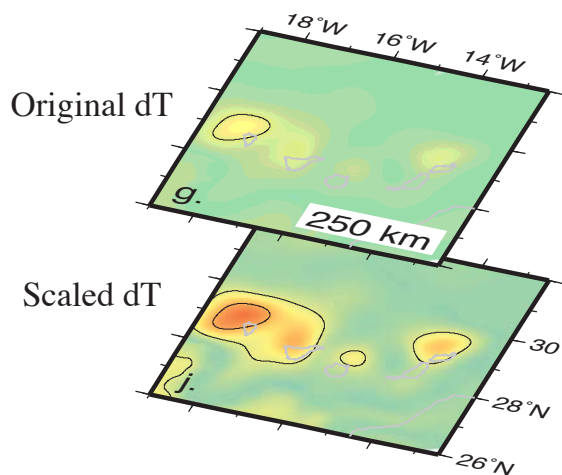


**A3**

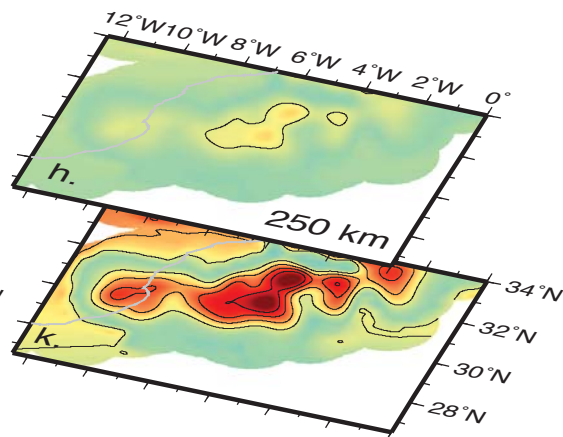


*S model*

**A1**



**A2**



**A3**

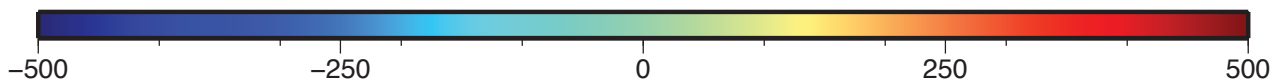
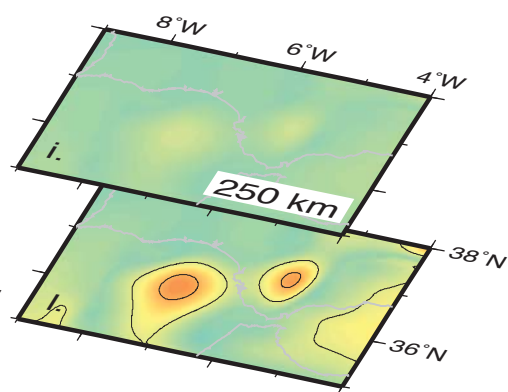


Figure 9.



