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# Uncertainty Modelling and Stability Robustness Analysis of Nucleic Acid-Based Feedback Control Systems

Nuno M.G. Paulino<sup>†1</sup>, Mathias Foo<sup>2</sup>, Jongmin Kim<sup>3</sup> and Declan Bates<sup>†4</sup>

Abstract-Recent advances in nucleic acid-based chemistry have highlighted its potential for the implementation of biomolecular feedback circuits. Here, we focus on a proposed design framework, which is able to approximate the inputoutput behaviour of key linear operators used in feedback control circuits by combining three elementary chemical reactions. The implementation of such circuits using DNA strand displacement introduces non-linear internal dynamics due to annihilation reactions among different molecular species. In addition, experimental implementation of in silico designs introduces significant levels of uncertainty and variability in reaction rate constants and equilibrium concentrations. Previous work using this framework has overlooked the practical implications of these issues for the construction of nucleic acidbased feedback control circuits. Here, we analyse the impact of these nonlinearities and uncertainties on the stability of a biomolecular feedback loop. We show that a rigorous analysis of its nucleic acid-based implementation requires an investigation of the associated non-linear dynamics, to decide on realisable parameters and acceptable equilibrium concentrations. We also show how the level of experimental uncertainty that is tolerated by the feedback circuit can be quantified using the structured singular value. Our results constitute a first step towards the development of a rigorous robustness analysis framework for nucleic acid-based feedback control circuits.

#### I. INTRODUCTION

A primary aim of synthetic biology is the development of systematic design and analysis frameworks for biomolecular feedback control systems that can regulate concentrations of molecular species inside the cell [1]. The recently proposed use of chemical reaction networks (CRNs), implemented via nucleic acid-based chemistry, as a programming language for the design of such systems [2], has made it possible to conceptualise many biological circuits in line with classical feedback control ideas, [3], [4].

However, a direct translation of abstract mathematical operators (e.g. gain, summation) into biomolecular circuits via chemical reactions is often not straightforward. For example, implementation of the subtraction operator using CRNs generally results in a one-sided operation, (i.e. it can only compute a positive difference between two inputs), whereas in feedback control the error signal generated by

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the difference between the measured and the desired values can take both positive and negative values. This is a critical limitation, as a one-sided error computation can lead to poor control performance or even instability [5], [6].

To get around this problem, Oishi and Klavins [3] proposed a formalism that can perform a two-sided subtraction, by letting each signal in the circuit be the difference in the concentration of two chemical species,  $x^+$  and  $x^-$ . In this way, a systematic framework for implementing simple linear biomolecular feedback controllers can be developed using three elementary reactions, and the same formalism has recently been extended to include more complicated mathematical operators [7] and non-linear controllers [8], [9].

While previous studies have indicated satisfactory stability and performance properties for controllers designed using this framework, no formal robustness analysis of these circuits has so far been attempted. This is an important issue, because the implementation of CRNs using nucleic acids requires the matching of specified reaction rates for the required biomolecular reactions to realise a particular operator, which can be problematic in real experimental implementations.

The use of nucleic acids to implement CRNs is appealing because a simple kinetic model of toehold-mediated strand displacements could explain a large number of different strand displacement reactions from the DNA sequences [10]. However, predicting reaction rate constants from sequence information alone, i.e. without experimental inputs, is still challenging and can result in large deviations from predicted design values. Improved biophysical models [11] and computational tools [12] can help improve the precision of predicted kinetics of hybridization and toehold-mediated strand displacement. It is also reasonable to assume that the different reactions may be characterised and the rates measured. Finally, varying the concentrations of auxiliary species [2] can also be used to iteratively bring the reaction rates closer to their design values. Based on the above considerations, it is estimated that environmental factors, together with uncertainty of the characterisation, i.e., the difference between measured rates and the actual reaction rates in the implementation, could result in uncertainty in the values of rate constants of between 10% to 20%.

Thus, it is crucial to be able to formally quantify in advance the robustness of proposed CRN-based designs to specified levels of uncertainty. Although some formal definitions of robustness in systems biology have been proposed [13], the rigorous analysis of the effects of perturbations and uncertainties in synthetic circuits calls for more structured

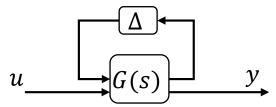


Fig. 1. Representation for  $\mu$ -analysis, where the nominal linear system G(s) is interconnected with the uncertainty structured in the system  $\Delta$ .

mathematical formulations for robustness analysis, estimation and model-based prediction [14].

Here, we investigate the internal stability of the complete set of reactions resulting from an implementation of a linear feedback control circuit using nucleic acid-based chemistry. The assumption of high concentrations of intermediary species allows some simplification of the reactions, but still captures an important aspect missing from the inputoutput behaviour of the system: the circuit has non-linearities due to the annihilation reactions, which are not observable in the input-output dynamics. We show that the presence of these hidden internal dynamics has an impact on the equilibrium of the system, and also imposes lower bounds on the concentrations of the intermediary species. More importantly, the attraction of the equilibrium condition and its stabilising role also needs to be characterised. Secondly, we use the structured singular value (SSV) tool [15], [16] from robust control theory to perform a formal validation of the robustness of the proposed controller. An informal robustness and sensitivity analysis was carried out in [3], where the system was subjected to a square wave as reference input over 50 simulations in the presence of random  $\pm$  10% variability in the reaction rates. This led to mismatches in some of the reactions rates, causing significant steady state errors in the input tracking. In contrast to simulation-based approaches, the SSV analysis does not rely on statistics, making it a very strong certification method as long as the modelling of the underlying uncertainty is representative.

The use of SSV, or  $\mu$ -analysis, has been previously investigated in systems biology [17] in an attempt to quantify and validate robustness in natural biochemical networks [18], and it is now widely used in other domains such as aircraft [19], [20] and space engineering [21]. Once the uncertainty in the circuit is established, it can be aggregated in an uncertain system  $\Delta$  interconnected with the nominal system (see Fig. 1), representing an infinite family of systems accounting for all possible combinations of the uncertain parameters, where any variation of the parameters will cause a deviation from the nominal response. The SSV analysis searches for the  $\mu$  value (see e.g. [19]), which is a metric defining how far the system is from losing stability. If the domain of  $\Delta$  does not contain destabilising parameter combinations, then robust stability is guaranteed for the specified level of uncertainty.

#### II. UNCERTAINTY MODELLING

Fig. 2 shows the configuration of the feedback control system analysed throughout this work. It is a single-input/single-

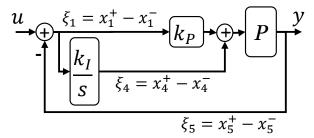


Fig. 2. The analysed feedback loop with proportional-integral controller from [3].

output (SISO) system with a single state  $\xi_4$  where the output of a constant plant  $P=1/(1+\varrho)$  tracks a reference input u by the action of a proportional and integral controller. The closed-loop feedback trajectories evolve accordingly to:

$$\dot{\xi}_4 = k_I \xi_1 = k_I (u - y) 
y = P(k_P \xi_1 + \xi_4)$$
(1)

# A. CRN Approximation of a Linear Feedback Control Circuit

Following the formalism introduced in [3], the configuration shown in Fig. 2 is approximated by a circuit composed of CRNs, where the negative and positive parts of a signal  $\xi$  are represented respectively by concentrations of two different species  $x^+$  and  $x^-$ . The final value of the signal is retrieved with  $\xi = x^+ - x^-$ .

The proposed CRNs for each of the operations in the circuit are based on three elementary reactions, namely catalysis, annihilation and degradation. Here, the nomenclature  $x^\pm$  represents compactly the two different species  $x^+$  and  $x^-$ . Similarly, the notation  $x^\pm \xrightarrow{\gamma^\pm} y^\pm$  represents compactly two simultaneous reactions

$$x^+ \xrightarrow{\gamma^+} y^+, \quad x^- \xrightarrow{\gamma^-} y^-$$
 (2)

The element  $\emptyset$  is used when the result of a reaction degrades, gets sequestered, or is removed, and no longer participates in the reactions.

Thus, the subtraction block  $\xi_1 = u - \xi_5$  is approximated by the reactions:

$$u^{\pm} \xrightarrow{\gamma_{1}^{\pm}} u^{\pm} + x_{1}^{\pm} \quad , \quad x_{5}^{\pm} \xrightarrow{\gamma_{2}^{\pm}} x_{5}^{\pm} + x_{1}^{\mp}$$

$$x_{1}^{\pm} \xrightarrow{\gamma_{3}^{\pm}} \emptyset \quad , \quad x_{1}^{+} + x_{1}^{-} \xrightarrow{\eta} \emptyset$$
(3)

For the weighted integrator  $\dot{\xi}_4 = k_I \xi_1$ , the CRNs are:

$$x_1^{\pm} \xrightarrow{k_I^{\pm}} x_1^{\pm} + x_4^{\pm} \quad , \quad x_4^{+} + x_4^{-} \xrightarrow{\eta} \emptyset$$
 (4)

Finally, for the weighted summation  $y = \xi_5 = k_p \xi_1 + \xi_4$ , we have:

$$x_{1}^{\pm} \xrightarrow{\gamma_{7}^{\pm} k_{P}^{\pm}} x_{1}^{\pm} + x_{5}^{\pm} \quad , \quad x_{4}^{\pm} \xrightarrow{\gamma_{8}^{\pm}} x_{4}^{\pm} + x_{5}^{\pm}$$

$$x_{5}^{\pm} \xrightarrow{\gamma_{9}^{\pm} (1+\varrho)} \emptyset \quad , \quad x_{5}^{+} + x_{5}^{-} \xrightarrow{\eta} \emptyset$$
(5)

The parameters  $\gamma_i$ ,  $k_P$ ,  $k_I$ , and  $\eta$  represent the rates of the reactions, and  $\varrho$  is a constant (set to  $\varrho=2$  as in [3]). Using the law of mass action, where the rate of a chemical reaction

is proportional to the product of the concentrations of the reactants (see e.g. [22]), the dynamics of the CRNs in (3)-(5) can be modelled using the ordinary differential equations (ODE):

$$\dot{x}_1^+ = -\gamma_3^+ x_1^+ + \gamma_2^- x_5^- + \gamma_1^+ u^+ - \eta x_1^+ x_1^- \tag{6}$$

$$\dot{x}_1^- = -\gamma_3^- x_1^- + \gamma_2^+ x_5^+ + \gamma_1^- u^- - \eta x_1^+ x_1^- \tag{7}$$

$$\dot{x}_4^+ = k_I^+ x_1^+ - \eta x_4^+ x_4^- \tag{8}$$

$$\dot{x}_4^- = k_I^- x_1^- - \eta x_4^+ x_4^- \tag{9}$$

$$\dot{x}_5^+ = -\gamma_9^+ (1+\varrho) x_5^+ + \gamma_7^+ k_P^+ x_1^+ + \gamma_8^+ x_4^+ - \eta x_5^+ x_5^- \quad (10)$$

$$\dot{x}_{5}^{-} = -\gamma_{9}^{-}(1+\varrho)x_{5}^{-} + \gamma_{7}^{-}k_{P}^{-}x_{1}^{-} + \gamma_{8}^{-}x_{4}^{-} - \eta x_{5}^{+}x_{5}^{-}$$
(11)

The input-output systems have states computed by the difference in concentration of the species tracking the positive and negative components of the signal with  $\xi_i = x_i^+ - x_i^-$ . Considering the nominal case where  $\gamma_i^+ = \gamma_i^-$ , the dynamics of the CRNs are linear and are given by

$$\dot{\xi} = \begin{bmatrix} -\gamma_3 & 0 & -\gamma_2 \\ k_I & 0 & 0 \\ \gamma_7 k_P & \gamma_8 & -\gamma_9 (1+\varrho) \end{bmatrix} \xi + \begin{bmatrix} \gamma_1 \\ 0 \\ 0 \end{bmatrix} u \quad (12)$$

With the right choice of rates  $\gamma_i$ , and controller parameters  $k_I$  and  $k_P$ , the input-output behaviour of (12) mimics the input-output behaviour of the linear system shown in Fig. 2. In this sense, the abstract CRNs of (3)-(5) can implement a similar input-output response to the feedback system (1) with biomolecular reactions, albeit with the inclusion of extra dynamics.

#### B. Implementation with Nucleic Acids

The abstract CRNs given above can be implemented experimentally using DNA strand displacement (DSD) reactions [2],[3],[11]. The parameters  $\gamma_i^\pm$ ,  $k_I^\pm$ , and  $k_P^\pm$  in the CRNs are set as functions of nominal strand displacement rates  $q_i^\pm$ . The individual dependencies of  $\gamma_i^\pm$  on  $q_i^\pm$  are inferred from the DSD implementation in [3], and set as

$$\gamma_i^{\pm} = \frac{1}{2} q_i^{\pm} C_{max}, \ i \in \{1, 2, 3, 7, 8, 9\}$$
 (13)

$$k_I^{\pm} = \frac{1}{2} q_5^{\pm} C_{max}, \quad k_P^{\pm} = \frac{q_7^{\pm}}{q_8^{\pm}}$$
 (14)

The scaling parameter  $C_{max}$  results from the DSD implementation of the feedback system in [3], which relies on bimolecular reactions mediated by intermediary auxiliary species with initial concentrations given by  $C_{max}$ . If  $C_{max} \gg x_i^\pm(t)$ , the concentration of the auxiliary species can be considered almost constant, and the set of reactions of the DSD network can be approximated by the unimolecular reactions of the CRNs in (3)-(5). Thus it is important to consider a set of parameters where the concentrations remain small with respect to  $C_{max}$ .

The annihilation reaction  $\eta$  depends on a maximum displacement rate  $q_{max}$  [2], and the DSD implementation proposed in [3] results in the equivalency

$$\eta = \frac{1}{2}q_{max} \tag{15}$$

If  $\eta \gg \gamma_i$ , the annihilation reaction keeps the concentrations  $x_i^+$  and  $x_i^-$  minimal without changing the transmission of the signal  $\xi_i$  [3]. Often this entails that for steady state, if  $\xi_i > 0$  then  $x_i^+ \approx \xi_i$  and  $x_i^- \approx 0$ . For  $\xi_i < 0$  we have the inverse with  $x_i^- \approx \xi_i$  and  $x_i^+ \approx 0$ .

### C. Implications of Non-Linear Dynamics

While the CRN model provides a framework to approximate the dynamics in (1) using concentrations of chemical species, when dealing with the DSD implementation it is important to note that the dynamics of  $\xi_i$  do not provide a complete description of the system. For example, the requirement that  $x_i^\pm(t) \ll C_{max}$  cannot be verified from the input-output response alone, since the concentrations of  $x_i$  are not observable in the states  $\xi_i$ , or y.

In addition, the stability and steady states of the system cannot be inferred from (12) alone, due to the absence of the non-linear annihilation reactions. Consider the non-negative state vector of species concentrations  $x=\begin{bmatrix}x_1^+,x_1^-,x_4^+,x_4^-,x_5^+,x_5^-\end{bmatrix}^T\in\mathbb{R}_0^+$ . The system of ODE's in (6)-(11) can be represented compactly as

$$\dot{x} = Ax + Bu - \eta g(x) \tag{16}$$

The input vector  $u = [u^+, u^-]^T$  contains both positive and negative components of the reference input, and the linear part of the model is given by

$$A = \left[ \begin{array}{cccccc} -\gamma_3^+ & 0 & 0 & 0 & 0 & \gamma_2^- \\ 0 & -\gamma_3^- & 0 & 0 & \gamma_2^+ & 0 \\ k_I^+ & 0 & 0 & 0 & 0 & 0 \\ 0 & k_I^- & 0 & 0 & 0 & 0 \\ \gamma_7^+ k_P^+ & 0 & \gamma_8^+ & 0 & -\gamma_9^+ (1+\varrho) & 0 \\ 0 & \gamma_7^- k_P^- & 0 & \gamma_8^- & 0 & -\gamma_9^- (1+\varrho) \end{array} \right]$$

$$B = \begin{bmatrix} \gamma_1^+ & 0 \\ 0 & \gamma_1^- \\ 0 & 0 \\ 0 & 0 \\ 0 & 0 \end{bmatrix}, C = \begin{bmatrix} 0 & 0 & 0 & 0 & 1 & -1 \end{bmatrix}$$

A non-linear term g(x) is now present containing the multiplications due to bimolecular annihilation reactions:

$$g(x) = \left[ \begin{array}{c} x_1^+ x_1^- \\ x_1^+ x_1^- \\ x_4^+ x_1^- \\ x_4^+ x_4^- \\ x_5^+ x_5^- \\ x_5^+ x_5^- \end{array} \right]$$

As mentioned above, the states  $\xi_i \in \mathbb{R}$  are obtained from the six non-negative components  $x_i^\pm \in \mathbb{R}_0^+$ . This subtraction results in the projection of the trajectories  $x_i^\pm$  in a  $\mathbb{R}^3$  manifold, defined by the projection matrix  $\Psi: \mathbb{R}^6 \to \mathbb{R}^3$ 

$$\Psi = \begin{bmatrix} 1 & -1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & -1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & -1 \end{bmatrix}$$
 (17)

Applying this projection to the non-linear system in (16) we have that  $\Psi \cdot g(x) = 0$ , and the trajectories of  $\xi_i$  result from the linear dynamics

$$\dot{\xi} = \Psi \left( Ax + Bu - \eta g(x) \right) = \frac{1}{2} \Psi A \Psi^T \xi + \Psi Bu \qquad (18)$$

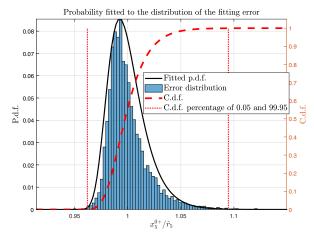


Fig. 3. Probability density function fitted to the distribution of the equilibrium estimation  $(1+e_6)=x_5^{0+}/\hat{r}_5$  (solid black), and the limits (dotted red) containing 99.9% of the respective cumulative distribution function (dashed red).

An immediate consequence is that the steady state of (18) tells us nothing about the steady state of the DNA strand concentrations, which are the solution to (16).

The other significant consequence is that, in our example, while  $\frac{1}{2}\Psi A\Psi^T$  is Hurwitz, the matrix A is not. The stability of the experimental implementation using DSD depends also on  $\eta g(x)$  and the equilibrium solution. For example, with initial conditions x(0)=0, at t=0 the system in (16) is unstable (g(x)=0) and A is not Hurwitz). The nonlinear term  $\eta g(x)$  is an adaptive stabilising term, where the state x needs to converge to an equilibrium where the system becomes stable. By collapsing the trajectories with the projection (17), this information about the internal stability is hidden.

#### D. Equilibria Change with Uncertainty on the Parameters

Our objective is to measure the robustness of the system to uncertainty. The strand displacement rates  $q_i^\pm$  depend on the designed DNA toehold sequences [10], which can be hard to predict. Even after experimental characterisation, there will always be an error between the measured and real value of the reaction rates. Here, we consider that the reaction rates are never perfectly known and are therefore uncertainties in the system.

Considering the start of the time response to be when  $u^+ = u^- = 0$ , the equilibrium condition  $x^0$  is the non-trivial solution to the non-linear equation

$$\eta \cdot g(x^0) = Ax^0 \tag{19}$$

Thus the solution depends on the parametrisation of A and the strand displacement rates  $q_i^{\pm}$ . Since it depends directly on the uncertain parameters, the solution is a moving equilibrium, [23], that changes with the values of the rates.

However, through numerical simulation of the dynamics in (16) it is possible to produce an arbitrary number of pairs associating an equilibrium condition  $x^0$  with the corresponding values of  $q_i^{\pm}$ , opening the way for estimation techniques. The SSV analysis framework requires rational representations of

uncertainty, and therefore each of the equilibrium components must be approximated by a polynomial function of the rates. To do this, the 14 uncertain rates  $q_i^\pm$  are organised in a column vector  $\bar{q}$ , and an estimator is obtained through the linear function

$$x^{0} = r(\bar{q}) + \varepsilon = V\bar{q} + \varepsilon \tag{20}$$

The samples of  $x^0$  and  $\bar{q}$  are arranged respectively as columns of R and Q, and  $V = RQ^T \left(QQ^T\right)^{-1}$  is the outcome of a Least Squares (LS) fitting. The fit of each of the components is further improved by expressing the LS fitting error  $\varepsilon_i$ ,  $i \in \{1, \ldots, 6\}$  as a quadratic function of  $r_i(\bar{q})$ :

$$\hat{r}_i(\bar{q}) = r_i(\bar{q}) + \left(\alpha_i + \beta_i r_i(\bar{q}) + \rho_i \left(r_i(\bar{q})\right)^2\right) + \vartheta_i \quad (21)$$

This new function reduces quadratic distortions of the errors and  $\varepsilon_i < \vartheta_i$ . The fittings of the matrix V and the three column vectors  $\alpha$ ,  $\beta$ , and  $\gamma$  in (21) were done using 10000 random samples of  $q_i^{\pm}$  in the intervals in Table I, to arrive at the estimation function and error

$$x^{0}(\bar{q}) = \hat{r}(\bar{q})(1+e) \tag{22}$$

Despite the presence of a small estimation error, we have now computed a rational function that relates the uncertain parameters and the equilibrium state.

To account for the estimation error we add corresponding uncertainties, as suggested in [23], obtained from the distributions of the fitting errors for each component  $e_i$ . These were fitted with a probability density function (p.d.f.). Limits of variability for  $e_i$  were defined containing 99.9% of the cumulative density function to define six additional uncertainties, which cover the error in tracking the moving equilibria  $x_i^0 = \hat{r}_i(\bar{q}) \cdot \delta_{ei}$ . The resulting limits of all uncertainties  $\delta_{ei}$  in the presence of 13% of uncertainty on  $q_i^\pm$  are listed in Table I. See the case of error  $(1+e_5)$  in Fig. 3 when estimating  $x_5^{0+}$ .

Linearising (16) around the equilibrium point  $(x, u) = (x^0, 0)$  then allows for a frequency representation of the system's response around its equilibrium, which is suitable for analysis within the SSV framework. The error dynamics

TABLE I  $\label{eq:table_eq} \text{Uncertain set } \Delta \text{ comprised by the reaction rates } q_i \text{ and the fitting error of the equilibrium } \delta_{ei}.$ 

Name	Nominal	Variability
$q_i^{\pm}, i \in \{1, 2, 3, 5, 7, 8, 9\}$	800	[13, 13] %
$\delta_{e1}$	1	$[0.967081 \ 1.09241]$
$\delta_{e2}$	1	[0.967098  1.09217]
$\delta_{e3}$	1	[0.9624631.06614]
$\delta_{e4}$	1	[0.962197  1.06822]
$\delta_{e5}$	1	[0.9619241.09484]
$\delta_{e6}$	1	[0.96421  1.10709]

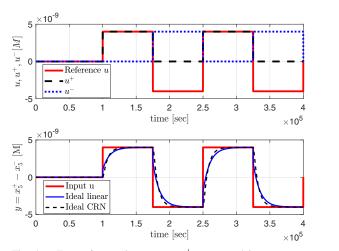


Fig. 4. Top: reference input  $u=u^+-u^-$  and its two components. Bottom: output of the projected CRN in (18) (dashed black) approximates the response of the ideal linear model in (1) (blue).

around equilibrium  $x^0$  are

$$\dot{x}_e = \frac{\partial f}{\partial x} \Big|_{x^0} x_e + \frac{\partial f}{\partial u} \Big|_{x^0} u_e$$

$$= (A - \eta J(x^0)) x_e + B u_e \qquad (23)$$

$$y_e = x_{5e}^+ - x_{5e}^- \qquad (24)$$

where  $x_e = x - x^0$ ,  $u_e = u - u^0$ , and  $y = y_e - y^0$ . The Jacobian matrix  $J(x^0)$  is given by

$$J(x^{0}) = \frac{\partial g(x)}{\partial x} \Big|_{x^{0},0}$$

$$= \begin{bmatrix} x_{1}^{0-} & x_{1}^{0+} & 0 & 0 & 0 & 0\\ x_{1}^{0-} & x_{1}^{0+} & 0 & 0 & 0 & 0\\ 0 & 0 & x_{4}^{0-} & x_{4}^{0+} & 0 & 0\\ 0 & 0 & x_{4}^{0-} & x_{4}^{0+} & 0 & 0\\ 0 & 0 & 0 & 0 & x_{5}^{0-} & x_{5}^{0+}\\ 0 & 0 & 0 & 0 & x_{5}^{0-} & x_{5}^{0+} \end{bmatrix}$$
(25)

The function in (22) is used to include  $\Delta$  in the Jacobian through  $J(\hat{r}_i(\bar{q}) \cdot \delta_{ei})$  to account for the moving equilibrium conditions. In this way, the equilibrium becomes correlated with the uncertainty, as in the original non-linear system. This step is important in order to remove the conservatism that would be introduced by assuming an independent variability of the equilibrium conditions.

# III. ANALYSIS RESULTS

The model in (16) for the non-linear system was implemented and simulated in Matlab/Simulink<sup>®</sup>, and analysed with the Robust Control Toolbox from Mathworks [24].

In our example,  $C_{max}=1\mu M$  and the displacement rates  $q_i=800M^{-1}s^{-1}$  (as in [3]). The maximum rate was set to  $q_{max}=10^6M^{-1}s^{-1}$ . The reference input u(t) is given by a square wave oscillating between  $\pm 4\times 10^{-9}M$ . The desired alternating step input is composed of the difference between two concentrations in the presence of the annihilation reaction  $u^++u^- \stackrel{\eta}{\to} \emptyset$  (see Fig. 4). The trajectories  $x_i^\pm$  for the complete nominal system in (16) are shown in Fig. 5

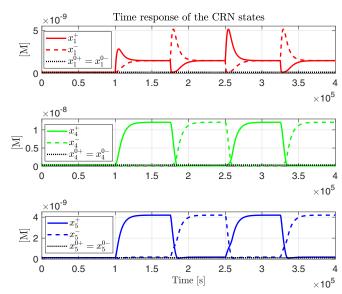


Fig. 5. Trajectories of the states for the non-linear system in (16), with initial conditions around the equilibrium solution of (19).

when the initial conditions are set to  $x^0$ . The tracking error  $\xi_1 = x_1^+ - x_1^-$  can assume positive and negative values and is null at steady state. The trajectories of  $x_1^\pm$  are not inferable from  $\xi_1$ , where  $x_1^+$  and  $x_1^-$  are always positive, and alternate their role with each new step on the reference.

For the robustness analysis we consider that variability is present for all 14 strand displacement reaction rates independently, which can allow for  $q_i^+ \neq q_i^-$ . The nominal and uncertain values for  $q_i^\pm$  and the fitting errors  $\delta_{ei}$  are summarised in Table I. The final system has a high repetition of uncertainties, and for strictly real uncertainties it can be computationally hard to obtain good lower bounds on the SSV (=  $1/\mu$ ). However, it is sufficient to demonstrate that the upper-bound of the SSV is less than 1 for all frequencies to guarantee that the system is stable for all possible parameterisations (see e.g. [19]).

As shown in Fig. 6, for uncertainty in the reaction rates of up to 13%, the SSV analysis guarantees that the closedloop system is stable for all uncertain parameters shown in Table I, since the upper bound on  $1/\mu$  remains below 1 for all frequencies. For levels of 14% or more the analysis is not conclusive. Since the fitting errors  $\delta_{ei}$  introduce conservatism, which does not exist in the non-linear system where the equilibrium is always the true value, a better description of  $\hat{r}(\bar{q})$  would improve the accuracy of the bound. Even with this extra artificial conservatism, however, our analysis shows that the system is guaranteed to be robustly stable for a 13% variation in the reaction rates, which falls within the expected 10 - 20% level of experimental uncertainty. This represents crucial information to help experimentalists estimate the accuracy with which reaction rates need to be implemented for successful construction of these circuits.

#### IV. CONCLUSIONS

In this paper we investigated the problem of designing linear feedback control circuits that can be practically im-

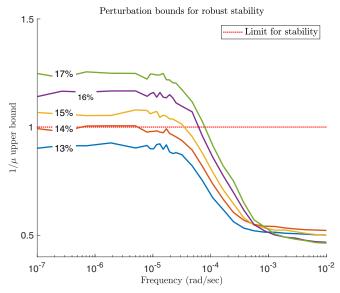


Fig. 6. Upper bounds on destabilising  $1/\mu$ , for different levels of uncertainty. The upper bounds are lower than 1 across all frequencies for uncertainty levels of 13% or lower.

plemented using nucleic acid-based chemistry. We showed how the implementation of even simple linear circuits with distinct chemical species depends on non-linear terms that drive important dynamics which are not observable in the system's input-output response. Rigorous analysis of the robustness of such circuits requires that we focus on the dynamics of the complete state vector and analyse the internal stability of the system. This provides information about closed loop stability and the system's steady state that is highly relevant when choosing experimental parameters.

The use of the SSV framework provides a method for the rigorous quantification of how much the circuit implementation deviates from the ideal response in the presence of experimental uncertainty, and allows a certification of the stability of the system around its equilibrium conditions. To achieve a linear representation, it was necessary to express the equilibrium conditions as a function of the uncertainty, to account for shifts due to parameter variability. We found the system to be robustly stable around its equilibrium for experimentally feasible uncertainty levels in the reaction rates. Further work is ongoing to investigate the significance of the non-linear terms in the dynamics of more complex feedback circuits, to improve the estimation of the moving equilibrium, to include additional sources of uncertainty, to quantify the effects of uncertainties on performance, and to obtain improved estimates of likely ranges of the uncertainties based on experimental data.

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