# LISA Pathfinder Micronewton Cold Gas Thrusters : In-flight Characterisation

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The LISA Pathfinder mission has demonstrated the ability to limit and measure the fluctuations in acceleration between two free falling test masses down to sub-femto-g levels. One of the key elements to achieve such a level of residual acceleration is the drag free control. In this scheme the spacecraft is used as a shield against any external disturbances by adjusting its relative position to a reference test mass. The actuators used to move the spacecraft are cold gas micro-propulsion thrusters. In this paper, we report in-flight characterisation of these thrusters in term of noise and artefacts during science operations using all the metrology capabilities of LISA Pathfinder. Using the LISA Pathfinder test masses as an inertial reference frame, an average thruster noise of  $\sim 0.17 \mu N/\sqrt{Hz}$ is observed and decomposed into a common (coherent) and an uncorrelated component. The very low noise and stability of the onboard metrology system associated with the quietness of the space environment allowed the measurement of the thruster noise down to  $\sim 20 \ \mu\text{Hz}$ , more than an order of magnitude below any ground measurement. Spectral lines were observed around  $\sim 1.5$  mHz and its harmonics and around 55 and 70 mHz. They are associated with the cold gas system itself and possibly to a clock synchronisation issue. The thruster noise-floor exhibits an excess of  $\sim 70\%$ compared to characterisation that have been made on ground on a single unit and without the feeding system. However this small excess has no impact on the LPF mission performance and is compatible with the noise budget for the upcoming LISA gravitational wave observatory. Over the whole mission, nominal and extension, the thrusters showed remarkable stability for both the science operations and the different manoeuvres necessary to maintain LPF on its orbit around L1. It is therefore concluded that a similar cold gas system would be a viable propulsion system for the future LISA mission.

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I. Introduction - LISA Pathfinder (LPF) was a European Space Agency mission [1] that demonstrated several technical milestones for the future gravitational wave observatory LISA [2]. In order to do so, two test masses (TM) of 1.92 kg and separated by 37 cm were placed in free fall at a level of differential acceleration of  $1.74 fm s^{-2}/\sqrt{Hz}$  down to a mHz [3]. One of the hardware components on which this performance relies is the cold gas micro-propulsion system [4],[5],[6]. LISA Pathfinder has a total of 12 cold gas micro-thrusters divided in two similar sets, prime and redundant, of 6 thrusters. Only one set can be used at the time.

In the low noise space environment, perturbations like solar wind, micro-meteorites or simply the self gravity of the satellite can strongly impact the performances. To prevent these perturbations from disturbing the test masses, the relative positions between the spacecraft (SC) and one of the free floating TMs is constantly monitored and maintained constant by adjusting the SC position using the thruster system. This strategy that forces the spacecraft to follow the test masses is called drag free and is implemented by the Drag Free and Attitude Control System (DFACS)[7]. The required range capability of the thrusters is directly dictated by the expected DC forces and torques on the spacecraft. During science operations for instance, the solar pressure plus the outgoing infrared radiation from the spacecraft, the sum estimated at around  $25\mu N$ , is the main contributor in term of DC force. This value was calculated along a direction perpendicular to the solar panel assuming a perfect pointing of

Table I. Thruster direction cosines

Thruster direction cosines					
Thruster	Х	Y	Z		
1	-0.078310	-0.864364	0.496732		
2	-0.787716	0.364364	0.496732		
3	0.787716	0.364364	0.496732		
4	0.078310	-0.864364	0.496732		
5	-0.709406	0.500000	0.496732		
6	0.709406	0.500000	0.496732		

the solar panels toward the sun. Six thrusters are used at the same time, consequently the mean thrusts required per thrusters by taking into account their orientation, see table I, is around 9  $\mu$ N.

This is for the static disturbances, but frequency dependent TM/SC motion have also to be nulled-out by the drag free loop up to 100 mHz. In any standard mission the noise coming from the micro-propulsion system will result in the jittering of the spacecraft. In LPF this jittering is strongly attenuated by the drag free loop as the commanded thrusts will mirror their own noise with a 180 degrees phase shift. Nevertheless some relative motion is still present due to the thruster noise. While the main measurement in LPF or in a gravitational wave observatory like LISA are accelerations between the freefalling test masses, motion of the surrounding spacecraft can introduce both measurement errors from cross-talk and dynamical coupling from force gradients [8]. In order to satisfy the expected performance of LPF, the requirement on the control accuracy of the S/C motion along the sensitive axis between 1mHz and 30mHz was  $2.5 \text{nm}/\sqrt{\text{Hz}}$ . That sets the level for each thruster noise at  $0.1\mu N/\sqrt{Hz}$  giving LPF's drag free controller design.

The cold gas propulsion system is also used for different manoeuvres that are not directly related to the science operations such as weekly station keeping to maintain the orbit around L1 (where TMs were kept in place electromagnetically) or spinning and de-spinning during the spacecraft/propulsion module separation (where TMs were still grabbed). These phases required to operate the thrusters at higher thrusts, up to  $500\mu N$  and with different requirements. The performance and characterisation during those phases are beyond the scope of this paper. Outside of these periods, the satellite was kept in *science operation* mode. Most of the results shown in this study come from special science periods when the satellite is in a *noise-only* state, i.e. when forces applied to the spacecraft or the test masses are only those necessary to keep it in drag free and attitude control mode. In the following, the data are referenced by the Day of the Year (DoY) measuring the number of days from the first of January 2016. LISA Pathfinder was launched on the 4th of December 2015 and was in commission mode

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during the month of February 2016. Science operations started on the 1st of March 2016. The mission was ended in July 2017.

In section II, an overview of the thrusters hardware setup and in particular the gas feeding system is presented. Section III will describe the in-flights characterisation of the cold gas in term of noise and anomalies. In section IV potential consequences for LISA are presented and in section V some conclusions are given.

**II. Cold gas micronewton hardware -** Each thruster can be described with four main parts represented in figure 1 :

- A mass flow sensor : the flow is proportional to the level of thrust.
- A piezo valve that controls the flow according to the requested thrust.
- A micro-propulsion electronics unit (MPE) common to the six thrusters.
- The Nitrogen feeding system, common to the six thrusters.

The MPE runs an internal closed loop at 40 Hz that commands the aperture of the valve. This command is proportional to the error between the requested thrusts sent by the on-board computer at 10 Hz and the measured thrusts computed from the mass flow sensor. The valve allows thrusts ranging from 0 to 500  $\mu N$  for both science and manoeuvre operations.

The system is made fully redundant with two sets of 6 thrusters and two MPE. The two sets are located on three external panels, each one having a cluster of four thrusters, two from the prime sets or side-A and two from the redundant set or side-B.



Figure 1. Functional schematic of the full cold gas propulsion chain. For clarity reasons, only one micro thruster assembly (MTA) and one feed line are represented. The *DFACS* is the controller that commands the forces and torques applied to the spacecraft. The dispatching converts spacecraft forces and torques into requested thrusts. The MPE is the electronic box that controls the MTA : mass flow sensors and piezo valve according to the requested thrust.

The satellite is controlled with six thrusters used simultaneously. The +z direction is defined as the direction

normal to the solar panel. The attitude control during the science run forces the solar panel to face the sun so that the +z direction is in the sun pointing direction. All the thrusters have the same angle of  $\approx 60$  degree with respect to the sun and are all pointing opposite to it (see table I). Thus it is impossible with this thruster geometry to command a thrust in the -z direction. However, the solar radiation pressure (SRP) plus the outgoing infrared radiation exert a quasi constant force of the order of 25  $\mu$ N. that can be considered as a 7th thruster in the -z direction. Note that this virtual 7th thruster has a constant thrust and so doesn't provide additional thrusts combination to generate a given spacecraft motion. Nonetheless, it allows the *DFACS* control algorithm to move the spacecraft in the -z direction despite the thrusters geometry. Figure 2 shows a view of the LISA Pathfinder S/Cas seen from above and shows the orientation of the Z axis with respect to the solar panel.



Figure 2. A view of the LISA Pathfinder SC with the solar panel removed. The thrusters can be observed on the outside panels. The small view on the right side shows the orientation of the Z axis. The colloidal thrusters are on the top and bottom panels.

This has two consequences for science operations:

- As the thrusters of a given set are operated at almost the same amplitudes, close to  $\sim 9\mu N$ , the impact of cross-talk effects between the thrusters can be neglected.
- As the illumination of the spacecraft is quite constant, temperature variations are not expected to impact the gas system.

The feeding system is separated into a high pressure part where the  $N_2$  gas is loaded to a pressure of 292 bars at the start of the mission and a low pressure part at 1 bar from which the thrusters are fed (see figure 3). While the high pressure part keeps decreasing during the mission due to propellant consumption, the low pressure side



Figure 3. Schematic of a panel with a propulsion system attached. Each of these panels has one  $N_2$  feed line and four micro-thrusters, two from the prime A-side and two from the redundant B-side. One feed line can provide  $N_2$  either to the A-side or B-side thrusters set by opening or closing the low pressure valve. The low pressure part of the feed line is pressure regulated to maintain a constant calibrated thrust.

is maintained at a constant pressure by two pressure regulators. The stability of the low pressure is required for stable thruster operations. There are four high pressure tanks located at three extremities of the spacecraft (see figure 4), each extremity defining an independent feed line. Feed line one and two have their own high pressure tank while the third feed line has two tanks clustered together. A single feed branch – including both the high and low pressure plumbing illustrated in Fig 3 – is used by all 6 A (or B) thrusters, such that only a single tank is depleted at a time.

This propellant storage strategy has two objectives:

- Redundancy, in case of failure of one of the 3 independent high pressure feed lines
- Enough flexibility to have a dedicated tank emptying strategy to gravitationally balance the remaining fuel mass [9].

The propellant mass for the whole mission is about 10 kg which is ~ 2% of the total mass of the satellite. After 9 months of operations ~ 3 kg were spent, which means that the spacecraft is losing an average 10 g of mass per day. This mass loss directly modifies the gravity imbalance seen by the two test masses and its impact was calculated and included in the quasi DC drift of the differential acceleration ( $\Delta g$ ) between the two test masses. This effect is observed at first order below  $10^{-5}Hz$  so outside of the nominal measurement bandwidth, but still needs to be compensated by the tank emptying strategy to maintain a small differential gravitational field.

The functionality checks conducted during commissioning such as switching between feed lines or handling



Figure 4. Simplified sketch of the cold gas high pressure tank locations and feed branches of LISA Pathfinder. The two test masses are shown at the center of the spacecraft.

low and high simultaneous thrusts were successfully validated.

### **III.** Thruster Characterisation -

#### A. Using commanded thrusts as a perturbation measurement

During the science operation, the thrusters are actuators in the multiple drag free control loops. As a consequence, their noise is not a directly measurable quantity. The first step is to estimate external perturbations on the spacecraft acceleration  $G_{SC_i}$ , the thrusters noise being one of them. For instance let's consider the sensitive axis x:

$$G_{SC_x} = -\frac{F_{cmd_x}}{M_{SC}} + \ddot{x_1} + \omega_1^2 \cdot x_1$$

Where  $x_1$  is the measurement of the sensor used to control the spacecraft in x,  $\omega_1^2$  is a stiffness term,  $\frac{F_{cmd_x}}{M_{SC}}$  is the total commanded thrusts in x divided by the spacecraft mass, and  $G_{SC_x}$  is any potential external disturbances along x.

The commanded forces sent by the drag free algorithm (DFACS) to the thrusters (ie :  $F_{cmd_x}$ ) can be used to measure the spacecraft disturbances under two conditions :

- The controller open loop gain is very high so that both  $x_1 \to 0$  and  $\ddot{x_1} \to 0$ .
- When the loop gain is very high, commanded forces will also null-out the noise of the sensor used in the loop. So sensing noises have to be low enough com-

pared to any other disturbances (i.e. :  $\ddot{x}_{1_{noise}} << G_{SC_x}$ ).

Both conditions are met for LISA Pathfinder within the bandwidth 0.1 - 10 mHz. The drag free loop counters the external noises effectively up to 100 mHz. The sensing noises have been estimated, approximately, at  $35 \cdot 10^{-15}$  mHz<sup>-1/2</sup> for the x interferometric position measurement, 2 nmHz<sup>-1/2</sup>, for y and z capacitive measurements and  $10^{-7}$  radHz<sup>-1/2</sup> for  $\theta_1$  capacitive measurement, see references [1] and [10]. Their impact on the spacecraft acceleration is well below the thrusters noise. Note that it also means that TMs are considered as perfect inertia reference frames for this study.

### B. Thruster noise and other external perturbations

Among the disturbances modelled in the  $\sim$  mHz range on the spacecraft, a thruster noise of 0.1  $\mu \text{NHz}^{-1/2}$  is the dominant one and is at least an order of magnitude above the solar wind protons and the solar radiation pressure fluctuations impact on the spacecraft motion. These quantities have been estimated from ACE([11])measurements and projected on LPF. Still, it is not possible to disentangle in the commanded thrusts, the thrusters noise contribution from a potential unknown noise source. So complementary measurements were performed to understand this level of noise. As one of NASA's contributions, LISA Pathfinder included a second set of thrusters, i.e. the *colloidal thrusters* [12], they have been used to remove this ambiguity. The commanded forces on the Z axis of the spacecraft with similar control scheme but with three different thrusters configurations (i.e. facing the same amplitude of the SRP) have been compared (see figure 5) :

- 1. The colloidal thrusters only are used for the drag free and solar radiation pressure (SRP) compensation (green line).
- 2. The colloidal thrusters are used for the drag free but the SRP is partially compensated by an open loop force commanded on four cold gas thrusters (red line). Cold gas thrusters number 1,2, 5 and 6 have been used for this, with thrusts of 18.4, 12.5, 5.3 and 11  $\mu N$ , respectively. The remaining 2 cold gas thrusters were set to provide 0 thrust, i.e. their noise contribution were not suppressed.
- 3. The cold gas thrusters only are used for drag free and SRP compensation (blue line).

When the colloidal thrusters only are used, the noise associated with the force compensation applied to the spacecraft (green line in the figure) is four times lower at 10 mHz than in the two other configurations. In the second configuration, an additional open loop thrust on four of the cold gas thrusters was sufficient to increase



Figure 5. Amplitude Spectral Density of commanded forces on the Z-axis of the spacecraft. The drag free is active in the three configurations. In green : Colloidal thrusters only are used. In red the colloidal thrusters are used for the drag free and an open loop thrusts is commanded on four cold gas thrusters. In blue, the cold gas thrusters only are used.

the commanded force noise up to  $0.4 \ \mu NHz^{-1/2}$ . This is exactly the noise level of the third configuration when the cold gas thrusters only are used for drag free operations. These measurements showed that cold gas thrusters noise are indeed the dominant source of external perturbation. As a consequence we can infer the level of noise of the cold gas thrusters by directly using the commanded thrusts.



Figure 6. Averaged commanded thrust Amplitude Spectral Density (ASD) during a "noise" run (DoY 95:104) below 10 mHz where the commanded thrust noise is dominant. Below 0.3 mHz, the impact of the rotation of the spacecraft is observed, see figure 9. Above, a white noise behaviour is observed with the presence of discreet lines.

#### C. Thruster noise general considerations

The six degrees of freedom of the spacecraft  $(x_{SC}, y_{SC}, z_{SC}, \theta_{SC}, \eta_{SC}, \Phi_{SC})$  are controlled with six thrusters. As a consequence there is only one set of commanded thrusts that can achieve a given spacecraft motion. Moreover all these axis are drag free controlled with closed loop transfert functions equal to one up to 10mHz [13]. Two major sources of correlations in the commanded thrusts unrelated to the thrusters were investigated :

i) To constantly point toward the sun and keep the antenna oriented toward earth, it is necessary to apply additional torques on the spacecraft at low frequencies (below 0.5 mHz) referred as the attitude control. These torques could create a correlated signal between commanded thrusts that doesn't relate to the thruster noise. They were measured using the forces applied on the y/z axis of the TMs and are orders of magnitude lower than the measured commanded torques in the 0.02 - 10 mHz frequency band.

ii) The scheme by which the DFACS commands forces and torques to counter the movement of the spacecraft is followed by a *dispatching algorithm* (i.e. a distribution of the thrusts among the thrusters in order to apply the requested forces and torques to the spacecraft) that can be suspected to introduce a level of correlation between the thrusters. A study of such correlations has been performed using ESA's LISA Pathfinder simulator which includes both the DFACS and the dispatching. The results show that no significant level of correlation is observed below 10 mHz and that, above, the correlations observed are explained by the impact of the DFACS.

So we expect the noise level of a given thruster, below 10 mHz, to be reflected as a noisy command of this particular thruster. Figure 6 shows the mean amplitude spectrum density, averaged over all 6 thrusters, of the commanded thrust during a science run ( $DoY \ 95:104$ ). A few remarks can be made:

- The thrusters were calibrated during *Station Keeping* by the ESA Mission Operation Center in Darmstadt. It is estimated to be precise to better than 10%. That sets our major source of error when estimating thrusters noise.
- Above 10 mHz the measured thruster noise increases due to the TM1  $\theta$  inertia sensor noise. This channel is used to control the  $\Theta_{SC}$  axis (rotation around the sensitive axis).
- Discrete lines are present at  $1.5 \ mHz$  and its harmonics. Groups of lines are also present around 55 and 70 mHz. This will be discussed later in this section.
- Below 0.3 *mHz*, the rise in amplitude is attributed to the rotation of the spacecraft, see figure 9.

• Discarding the observed lines, the noise between 0.3 mHz and 10 mHz is constant with an average value of 0.17  $\mu N/\sqrt{Hz}$ .



Figure 7. Evolution of the average Amplitude Spectral Density of each commanded thruster noises in the  $1.7 - 2.7 \ mHz$  frequency range as a function of time (*DoY*). The three panels show, from top to bottom, thrusters 1 - 2, 3 - 4 and 5 - 6. The dashed lines indicate the period during which NASA's colloidal thrusters were used.

Figure 7 shows the evolution, over most of the mission, of the noise for each individual thruster between 1 and 10 mHz, after removal of the spectral lines (see figure 6). Each data point corresponds to a 24h measurement and their statistical errors are smaller than the size of the data points. Up to DoY 272 set A thrusters were used and beyond there was a switch to set B. The black dotted lines, at DoY 179 and 339, define the period when NASA's colloidal thrusters were used and which induced a significant increase (a few degrees) in the temperature of the S/C. This is a possible explanation for the increase of the noise in this period. Otherwise, the noise levels display rather stable values up to DoY 179. Beyond DoY 272, when set B thrusters were used and large temperature variations were imposed on the S/C, slightly larger fluctuations are observed.

The thrusters noise exceeds the mission requirements by around 70% but it still fits within the noise margin and doesn't impact the performances of LISA Pathfinder or of LISA. It will be seen, in the next section, that part of this noise is a *common (coherent) noise* which will need further ground studies to elucidate.

Prior to the mission, ground characterisations were performed [14] on a thruster flight model but without the feeding system and only down to 2 mHz because of test bench limitations. These tests demonstrated a noise lower than  $0.1\mu \text{NHz}^{-1/2}$  for similar range of thrusts, thus below the values observed in flight. This seems to imply a

problem related with the feeding system common to both primary and redundant set of thrusters, even though, so far, no explicit investigation has demonstrated this.

#### D. Common (coherent) Noise Between Thrusters

Figure 8 shows the Amplitude Spectral Density (ASD) of the forces on the X, Y and Z axis during a noise run. One observes that the noise level on the Z axis around 1 mHz is significantly larger than those on the other axis. The dotted lines indicate the levels expected, around 1 mHz, if the thrusters noise measured by the individual thrusts commanded are projected on the X, Y and Z and considered as statistically independent noises. The mismatch between expected and measured values and that these lines are also above the measured X and Y forces, can be explained by the presence of a *common noise.* Because of the orientations of the thrusters (see table I), a common noise would add up on the Z axis whereas it would have a null impact on the X and Y axis where the sum of the direction cosines is quasi-null. This asymmetry is observed in figure 8. Figure 9 shows the Amplitude Spectral Density (ASD) of the torques around the X, Y and Z axis during the same period and the projection of measured thrusters noise as incoherent noises. The difference between these predictions and the measurements can, there again, be explained by the orientation of the thrusters (see table I) and the presence of a *common noise*. In order to evaluate the hypothesis of a common noise, a coherence analysis of the thruster noises has been performed.

Figure 10 shows the correlation values observed between thrusters 1 and 2 (red line) and between thrusters 5 and 6 (blue line), for DoY 95:104 in both cases. Below 10 mHz a strong degree of correlation is observed. At the frequencies corresponding to the observed lines in figure 8 the correlation is 100%, as discussed in the preceding section, these lines are common to all thrusters. Apart from these discrete cases, the correlation stays close to 90% for the first case and close to 50% for the second. This indicates that there exists a common component between thrusters. The same observation can be made for all other thruster pairs (the coherence levels varying between 90% and 50%) and for the duration of the mission although a significant decrease of coherence is observed for  $DoY \gtrsim 300$ .

The method to extract this common noise component is the following:

• Figure 5 shows that the common noise measured on the Z-axis is independent of the average commanded thrusts (between 0 and 18.4  $\mu$ N for the joint experiment with the colloidal or all thrusters at 9  $\mu$ N for coldgas only) but inherent to the coldgas micro-propulsion system. So the (strong) assumption is made that the common noise is equal on all thrusters, i.e. the level of common noise on 7

thruster<sub>i</sub> (in  $\mu N/\sqrt{Hz}$ ) is the same as on thruster<sub>i</sub>.

- Using a coherence algorithm for each pair of thrusters that estimates the cross power spectrum between two time series, the level of the common component is extracted. There are therefore 15 such values from as many pairs.
- The average of these values is then used to estimate the common component. Taking the minimum of these values would diminish this estimate by only 10%.

Figure 11 displays, for each thruster, the mean noise level (in  $\mu N\sqrt{\text{Hz}}$ ) during the whole mission after the impact of the spectral lines has been removed. These values are quite stable although some evolution is seen beyond *DoY 300*. The black points and dotted line shows the average *common (coherent) noise* over all thruster pairs and for each measurement using the method outlined above.

Figure 12 displays the level of uncorrelated thrust (in  $\mu N\sqrt{Hz}$ ) for each thruster. The uncorrelated noise level is calculated by subtracting, in power, the common noise level from the total power. From this study, it is concluded that the uncorrelated (statistically independent) thruster noise of the thrusters can be estimated, in the mHz range, as a white noise with levels between 0.5 and 0.15  $\mu N/\sqrt{Hz}$ . The value for thruster 1 at DoY90 appearing as an outlier. The uncorrelated noise levels appear to be stable in time with an average level of ~ 0.10  $\mu N/\sqrt{Hz}$ , suggesting that the evolution of the total noise level beyond DoY 300 (see figure 11) could be due to the evolution of the common noise contribution.



Figure 8. Amplitude Spectral Density (ASD) of commanded forces for DoY 95:104. The dotted lines indicates the expected levels if the individual measured thruster noises were statistically incoherent.





Figure 9. Amplitude Spectral Density (ASD) of commanded torques for DoY 95:104. The dotted lines indicates the expected levels if the individual measured thruster noises were statistically incoherent.



Figure 10. Correlation (absolute value and phase difference) between thrusters 1 and 2 (red lines) and between thrusters 5 and 6 (blue lines) as measured during DoY 95:104. Above 10 mHz, the correlations are due to the DFACS system.

## E. Thrusters lines

As can be seen on figure 6, a number of discrete lines are observed at 1.5 mHz and its multiples. Other sets of lines are seen around 55 and 70 mHz. In fact all these lines appear to originate from the same source and, as shown above, are a feature of the cold gas propulsion system and 100% correlated between individual thrusts. The lines observed at 55 and 70 mHz are, as shown by figures 13 and 14 , multiple lines separated by 1.5 mHz indicating that they probably have a common origin to the lines observed at lower frequencies.

The observation that these lines are extremely narrow seems to point to an electronic effect. As an example, a quality factor (QF =  $f/\Delta f(@3db)$ ) of ~ 2000 is measured for the 3 mHz line during DoY 95:104.



Figure 11. Mean Thruster noise as a function of DoY. The black dots indicate the average value of the common (coherent) part of the thrust, calculated between all possible pairs.



Figure 12. Extracted uncorrelated thruster noises as a function of DoY. See text for further details

It should be noted that they produce a real motion of the spacecraft since they can be seen by measurements of the distance between  $TM_1$  and the spacecraft.

We are currently investigating a clock synchronisation problem between the thruster electronic that runs a loop at 40 Hz and the DFACS that request commanded thrusts at 10 Hz based on a separate clock. This seems to be confirmed by a study of their behaviour during the mission that shows a correlation between the frequency of the lines and the average temperature of the spacecraft (see figure 15) with a coupling coefficient of 0.02 mHz/Kfor the line at 1.5 mHz. Further studies should be performed on a dedicated ground based testbed in order to track and eventually confirm their origin. It may be interesting to note that the line seen at 70 mHz shows a similar dependence upon temperature (but with a slope of 0.93 mHz/K) whereas the 55 mHz line has the same slope as the 70 mHz line but with a negative value (see figure 15). It is worth noting that the temperature dependence df/dT of the 70 mHz line is equal to 70/1.5 times the df/dT of the 1.5 mHz line.



Figure 13. Observed lines around 70 mHz. The lines are separated by 1.52 mHz. The lowest frequency line is observed at 1.54 mHz. Data for DoY 95:104.



Figure 14. Observed lines around 50 mHz. The lines are separated by 1.55 mHz The lowest frequency line is observed at 1.54 mHz. Data for DoY 95:104.

**IV.** Projection for LISA - Even though greatly reduced by the drag free controller, the thrusters noise is an important source of spacecraft motion with respect to the test masses above  $\sim mHz$  [13]. Measurements of the thrusters induced motion are given on LPF by in-loop drag free controlled channels :  $x_1$ ,  $y_1 + y_2$ ,  $z_1 + z_2$ ,  $y_1 - y_2$ ,  $z_1$ - $z_2$  and  $\theta_1$ . The motion itself does not directly couple with the test mass to test mass measurement error. In LPF, the optical bench or spacecraft motion is correlated between the two TMs so that it is eliminated when the differential acceleration  $(\Delta g)$  is computed. For LISA, the optical bench displacement noise will also be completely suppressed after reconstruction of the test masstest mass measurement between two satellites with the Time Delay Interferometry algorithm. However, thruster noises, through spacecraft motion, will impact LISA via



Figure 15. Frequency evolution of 1.5, 55 and 70 mHz thruster lines during the mission and the average temperature of the satellite as a function of DoY. A strong correlation is seen with a coefficient of 0.02 mHz/K for the 1.5 mHz line.

two main physical effects:

- Optical imperfections (misalignments and wavefront error) will couple spacecraft motion into optical path length error. This source of noise is usually referred to as tilt to length coupling (TTL) for the spacecraft angular jitter coupling.
- A moving spacecraft will couple the TM-spacecraft relative position with any gradient field such as the spacecraft self-gravity or the electrostatic field. This coupling generates a parasitic force on the TM proportional at the first order with the TM position (ie : a stiffness term).

Both these effects were characterised on LPF but will impact LISA differently. In LPF the test masses are affected by the same spacecraft motion, so a significant part of the stiffness noises contribution appeared as correlated noise and vanished when computing  $\Delta g$ . The stiffness coefficient was measured via dedicated experiments at  $\omega_1^2 = (-525 \pm 30) \times 10^{-9}/s^2$  while differential stiffness  $(\Delta \omega_{12}^2)$  was consistent with zero [1]. For LISA, spacecraft motion will be uncorrelated between two distant spacecraft so that contributions for a test mass to test mass measurement will be incoherent and sum up quadratically in the noise budget. If we assume similar levels of stiffness and design of drag free controller, the expression  $\sqrt{2} \cdot x1 \cdot \omega_1^2$  gives a rough estimation of the thrusters impact on stiffness noise for LISA. The additional factor of  $\sqrt{2}$  is applied to take into account the effect on both test masses. This contribution has a maximum of  $6.7 \times 10^{-15} \text{ms}^{-2}$  at 25 mHz which is well below the LISA noise curve requirement in terms of differential acceleration as shown in [1].

In LPF spacecraft pickup into  $\Delta g$  was mainly due to misalignment between the TMs, the GRS and the spacecraft. The resulting contribution appeared as a *bump* between 20mHz and 100mHz. It was subtracted by fitting a simple model using other TM's degrees of Freedom measurements [8]. The optical layout in LISA to do a test mass to test mass measurement is much more complex as it involves multiple interferometers and imaging systems. More sources of misalignment as well as the wavefront error of the transmitted beam will greatly increase the level of coupling between angular or lateral jitter of the spacecraft and the main measurement. Postprocessing subtraction of this noise will be mandatory to achieve the required performance for LISA. The subtraction method won't be as straightforward as for LPF since it will rely on the Time Delay Interferometry algorithm first to reconstruct laser noise free signals. So, minimizing the thrusters noise impact on the spacecraft jitter by DFACS design or by a better handling of the common noise and lines will be crucial for this source of noise.

**V. Conclusion** - The full cold gas propulsion system (i.e. the thrusters with their feeding system) was characterised with a flat white noise component of ~  $0.17\mu$ N/ $\sqrt{\text{Hz}}$  down to 0.02 mHz. This noise has been decomposed into a common (coherent) and incoherent part, both of comparable amplitudes. The presence of lines and the common (coherent) part of the noise were identified as a cold gas thruster system contribution by joint use of the NASA provided colloidal propulsion system with the cold gas system.

As there is no clear evidence on the origin of these artefacts, the observations that have been made in-flight (correlation of the lines with the temperature, a noise insensitive to the thrusters used) should lead to specific ground experiments in order to elucidate and remedy these effects.

The gain of the thrusters were calibrated by ESA's Mission Operation Center during station keeping and the values obtained seem reliable to better than 10%.

For LISA, the gain fluctuation have no consequences on its predicted performances. However, the observed spectral lines and the excess coherent noise will induce spacecraft motion that will couple with the test masses sensitive axis motion through stiffness and optical imperfections. Their expected level could be non-negligible. Dedicated studies or design strategy should therefore be devoted to explain or suppress their presence.

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- [1] M. Armano, H. Audley, G. Auger, J. T. Baird, M. Bassan, P. Binetruy, M. Born, D. Bortoluzzi, N. Brandt, M. Caleno, L. Carbone, A. Cavalleri, A. Cesarini, G. Ciani, G. Congedo, A. M. Cruise, K. Danzmann, M. de Deus Silva, R. De Rosa, M. Diaz-Aguiló, L. Di Fiore, I. Diepholz, G. Dixon, R. Dolesi, N. Dunbar, L. Ferraioli, V. Ferroni, W. Fichter, E. D. Fitzsimons, R. Flatscher, M. Freschi, A. F. García Marín, C. García Marirrodriga, R. Gerndt, L. Gesa, F. Gibert, D. Giardini, R. Giusteri, F. Guzmán, A. Grado, C. Grimani, A. Grynagier, J. Grzymisch, I. Harrison, G. Heinzel, M. Hewitson, D. Hollington, D. Hoyland, M. Hueller, H. Inchauspé, O. Jennrich, P. Jetzer, U. Johann, B. Johlander, N. Karnesis, B. Kaune, N. Korsakova, C. J. Killow, J. A. Lobo, I. Lloro, L. Liu, J. P. López-Zaragoza, R. Maarschalkerweerd, D. Mance, V. Martín, L. Martin-Polo, J. Martino, F. Martin-Porqueras, S. Madden, I. Mateos, P. W. Mc-Namara, J. Mendes, L. Mendes, A. Monsky, D. Nicolodi, M. Nofrarias, S. Paczkowski, M. Perreur-Lloyd, A. Petiteau, P. Pivato, E. Plagnol, P. Prat, U. Ragnit, B. Raïs, J. Ramos-Castro, J. Reiche, D. I. Robertson, H. Rozemeijer, F. Rivas, G. Russano, J. Sanjuán, P. Sarra, A. Schleicher, D. Shaul, J. Slutsky, C. F. Sopuerta, R. Stanga, F. Steier, T. Sumner, D. Texier, J. I. Thorpe, C. Trenkel, M. Tröbs, H. B. Tu, D. Vetrugno, S. Vitale, V. Wand, G. Wanner, H. Ward, C. Warren, P. J. Wass, D. Wealthy, W. J. Weber, L. Wissel, A. Wittchen, A. Zambotti, C. Zanoni, T. Ziegler, and P. Zweifel, Phys. Rev. Lett. **116**, 231101 (2016).
- [2] The eLISA Consortium, :, P. A. Seoane, S. Aoudia, H. Audley, G. Auger, S. Babak, J. Baker, E. Barausse, S. Barke, M. Bassan, V. Beckmann, M. Benacquista, P. L. Bender, E. Berti, P. Binétruy, J. Bogenstahl, C. Bonvin, D. Bortoluzzi, N. C. Brause, J. Brossard, S. Buchman, I. Bykov, J. Camp, C. Caprini, A. Cavalleri, M. Cerdonio, G. Ciani, M. Colpi, G. Congedo, J. Conklin, N. Cornish, K. Danzmann, G. de Vine, D. De-Bra, M. Dewi Freitag, L. Di Fiore, M. Diaz Aguilo, I. Diepholz, R. Dolesi, M. Dotti, G. Fernández Barranco, L. Ferraioli, V. Ferroni, N. Finetti, E. Fitzsimons, J. Gair, F. Galeazzi, A. Garcia, O. Gerberding, L. Gesa, D. Giardini, F. Gibert, C. Grimani, P. Groot, F. Guzman Cervantes, Z. Haiman, H. Halloin, G. Heinzel, M. Hewitson, C. Hogan, D. Holz, A. Hornstrup, D. Hoyland, C. D. Hoyle, M. Hueller, S. Hughes, P. Jetzer, V. Kalogera, N. Karnesis, M. Kilic, C. Killow, W. Klipstein, E. Kochkina, N. Korsakova, A. Krolak, S. Larson, M. Lieser, T. Littenberg, J. Livas, I. Lloro, D. Mance, P. Madau, P. Maghami, C. Mahrdt, T. Marsh, I. Mateos, L. Mayer, D. McClelland, K. McKenzie, S. McWilliams, S. Merkowitz, C. Miller, S. Mitryk, J. Moerschell, S. Mohanty, A. Monsky, G. Mueller, V. Müller, G. Nelemans, D. Nicolodi, S. Nissanke, M. Nofrarias, K. Numata, F. Ohme, M. Otto, M. Perreur-Lloyd, A. Petiteau, E. S. Phinney, E. Plagnol, S. Pollack, E. Porter, P. Prat, A. Preston, T. Prince, J. Reiche, D. Richstone, D. Robertson, E. M. Rossi, S. Rosswog, L. Rubbo, A. Ruiter, J. Sanjuan, B. S. Sathyaprakash, S. Schlamminger, B. Schutz, D. Schütze, A. Sesana, D. Shaddock, S. Shah, B. Sheard, C. F. Sopuerta, A. Spector,

R. Spero, R. Stanga, R. Stebbins, G. Stede, F. Steier,
T. Sumner, K.-X. Sun, A. Sutton, T. Tanaka, D. Tanner,
I. Thorpe, M. Tröbs, M. Tinto, H.-B. Tu, M. Vallisneri,
D. Vetrugno, S. Vitale, M. Volonteri, V. Wand, Y. Wang,
G. Wanner, H. Ward, B. Ware, P. Wass, W. J. Weber,
Y. Yu, N. Yunes, and P. Zweifel, ArXiv e-prints (2013),
arXiv:1305.5720 [astro-ph.CO].

- [3] M. Armano et al., Phys. Rev. Lett. 120, 061101 (2018).
- [4] G. Morris, N. Dunbar, P. Bianco, and C. Edwards, in 49th AIAA/ASME/SAE/ASEE Joint Propulsion Conference, San Jose, California, USA : Design of a Cold-Gas Micropropulsion system for LISA Pathfinder (July 2013).
- [5] G. Noci, G. Matticari, P. Siciliano, L. Fallerini, L. Boschini, and V. Vettorello, in 45th AIAA/ASME/SAE/ASEE Joint Propulsion Conference and Exhibit, Denver, Colorado, USA : Cold Gas Micro Propulsion System for Scientific Satellite Fine Pointing: Review of Development and Qualification Activities at Thales Alenia Space Italia (July 2009).
- [6] A. Polli, in 65th International Astronautical Congress 2014, Space Propulsion Symposium (C4), Propulsion Technology (1) (3), Toronto, Canada : European Cold Gas Micro Propulsions System reached TRL-9 on board GAIA S/C. Discussion on in flight performance and system flexibility and capabilities for being implemented in other mission requiring ultra-fins satellite positioning and attitude (September 2014).
- [7] A.Schleicher, T.Ziegler, R.Schubert, N.Brandt, P.Bergner, U.Johann, W.Fichter, and J.Grzymisch, CEAS Space Journal 10, 471 (2010).
- [8] M. Armano, H. Audley, J. Baird, P. Binetruy, M. Born, D. Bortoluzzi, E. Castelli, A. Cavalleri, A. Cesarini, A. M. Cruise, K. Danzmann, M. de Deus Silva, I. Diepholz, G. Dixon, R. Dolesi, L. Ferraioli, V. Ferroni, E. D. Fitzsimons, M. Freschi, L. Gesa, F. Gibert, D. Giardini, R. Giusteri, C. Grimani, J. Grzymisch, I. Harrison, G. Heinzel, M. Hewitson, D. Hollington, D. Hoyland, M. Hueller, H. Inchauspé, O. Jennrich, P. Jetzer, N. Karnesis, B. Kaune, N. Korsakova, C. J. Killow, J. A. Lobo, I. Lloro, L. Liu, J. P. López-Zaragoza, R. Maarschalkerweerd, D. Mance, N. Meshksar, V. Martín, L. Martin-Polo, J. Martino, F. Martin-Porqueras, I. Mateos, P. W. McNamara, J. Mendes, L. Mendes, M. Nofrarias, S. Paczkowski, M. Perreur-Lloyd, A. Petiteau, P. Pivato, E. Plagnol, J. Ramos-Castro, J. Reiche, D. I. Robertson, F. Rivas, G. Russano, J. Slutsky, C. F. Sopuerta, T. Sumner, D. Texier, J. I. Thorpe, D. Vetrugno, S. Vitale, G. Wanner, H. Ward, P. Wass, W. J. Weber, L. Wissel, A. Wittchen, and P. Zweifel, Phys. Rev. D 97, 122002 (2018), arXiv:1806.08581 [astro-ph.IM].
- [9] M. Armano, H. Audley, G. Auger, J. Baird, P. Binetruy, M. Born, D. Bortoluzzi, N. Brandt, A. Bursi, M. Caleno, A. Cavalleri, A. Cesarini, M. Cruise, K. Danzmann, M. de Deus Silva, D. Desiderio, E. Piersanti, I. Diepholz, R. Dolesi, N. Dunbar, L. Ferraioli, V. Ferroni, E. Fitzsimons, R. Flatscher, M. Freschi, J. Gallegos, C. G. Marirrodriga, R. Gerndt, L. Gesa, F. Gibert, D. Giardini, R. Giusteri, C. Grimani, J. Grzymisch, I. Harrison, G. Heinzel, M. Hewitson, D. Hollington, M. Hueller,

J. Huesler, H. Inchauspé, O. Jennrich, P. Jetzer, B. Johlander, N. Karnesis, B. Kaune, N. Korsakova, C. Killow, I. Lloro, L. Liu, J. P. López-Zaragoza, R. Maarschalkerweerd, S. Madden, D. Mance, V. Martín, L. Martin-Polo, J. Martino, F. Martin-Porqueras, I. Mateos, P. W. McNamara, J. Mendes, L. Mendes, A. Moroni, M. Nofrarias, S. Paczkowski, M. Perreur-Lloyd, A. Petiteau, P. Pivato, E. Plagnol, P. Prat, U. Ragnit, J. Ramos-Castro, J. Reiche, J. A. R. Perez, D. Robertson, H. Rozemeijer, F. Rivas, G. Russano, P. Sarra, A. Schleicher, J. Slutsky, C. F. Sopuerta, T. Sumner, D. Texier, J. I. Thorpe, R. Tomlinson, C. Trenkel, D. Vetrugno, S. Vitale, G. Wanner, H. Ward, C. Warren, P. J. Wass, D. Wealthy, W. J. Weber, A. Wittchen, C. Zanoni, T. Ziegler, and P. Zweifel, Classical and Quantum Gravity **33**, 235015 (2016).

- [10] M. Armano *et al.* (LISA Pathfinder), Phys. Rev. D96, 062004 (2017).
- [11] D. McComas, S. Bame, P. Barker, W. Feldman, J. Phillips, P. Riley, and J. Griffee, Space Sci. Rev. 86, 563 (1998).
- [12] G. Anderson, J. Anderson, M. Anderson, G. Aveni, D. Bame, P. Barela, K. Blackman, A. Carmain, L. Chen, M. Cherng, et al., Physical Review D 98, 102005 (2018).

- [13] M. Armano, H. Audley, J. Baird, P. Binetruy, M. Born, D. Bortoluzzi, E. Castelli, A. Cavalleri, A. Cesarini, M. Cruise, K. Danzmann, M. de Deus Silva, I. Diepholz, G. Dixon, R. Dolesi, L. Ferraioli, V. Ferroni, E. Fitzsimons, M. Freschi, L. Gesa, F. Gibert, D. Giardini, R. Giusteri, C. Grimani, J. Grzymisch, I. Harrison, G. Heinzel, M. Hewitson, D. Hollington, D. Hoyland, M. Hueller, H. Inchauspe, O. Jennrich, P. Jetzer, N. Karnesis, B. Kaune, N. Korsakova, C. J. Killow, A. Lobo, I. Lloro, L. Liu, J.-P. P. Lopez-Zaragoza, R. Maarschalkerweerd, D. Mance, N. Meshksar, V. Martin, L. Martin-Polo, J. Martino, F. Martin-Porqueras, I. Mateos, P. McNamara, J. Mendes, L. Mendes, M. Nofrarias, S. Paczkowski, M. Perreur-Lloyd, A. Petiteau, P. Pivato, E. Plagnol, J. Ramos-Castro, J. Reiche, D. Robertson, F. Rivas, G. Russano, J. Slutsky, C. Sopuerta, T. Sumner, D. Texier, I. Thorpe, D. Vetrugno, S. Vitale, G. Wanner, H. Ward, P. Wass, W. Weber, L. Wissel, A. Wittchen, and P. Zweifel, arXiv e-prints (2018), arXiv:1812.05491 [astro-ph.IM].
- [14] J. Jarrige *et al.*, Journal of Propulsion and Power **30**, 934 (2014).