An optimized method for ¹⁵N R₁ relaxation rate measurements in non-deuterated proteins

Margarida Gairí, Andrey Dyachenko, M. Teresa González, Miguel Feliz, Miquel Pons & Ernest Giralt

Journal of Biomolecular NMR

ISSN 0925-2738

J Biomol NMR DOI 10.1007/s10858-015-9937-4



Special Issue on NMR studies of membrane proteins Guest Editors: Rob Kaptein · Gerhard Wagner



Your article is published under the Creative Commons Attribution license which allows users to read, copy, distribute and make derivative works, as long as the author of the original work is cited. You may selfarchive this article on your own website, an institutional repository or funder's repository and make it publicly available immediately.



ARTICLE



An optimized method for ¹⁵N R₁ relaxation rate measurements in non-deuterated proteins

Margarida Gairí¹ · Andrey Dyachenko² · M. Teresa González¹ · Miguel Feliz¹ · Miquel Pons³ · Ernest Giralt²

Received: 29 October 2014/Accepted: 27 April 2015 © The Author(s) 2015. This article is published with open access at Springerlink.com

Abstract ¹⁵N longitudinal relaxation rates are extensively used for the characterization of protein dynamics; however, their accurate measurement is hindered by systematic errors. ¹⁵N CSA/¹H-¹⁵N dipolar cross-correlated relaxation (CC) and amide proton exchange saturation transfer from water protons are the two main sources of systematic errors in the determination of ${}^{15}N$ R₁ rates through ¹H-¹⁵N HSQC-based experiments. CC is usually suppressed through a train of 180° proton pulses applied during the variable ¹⁵N relaxation period (T), which can perturb water magnetization. Thus CC cancellation is required in such a way as to minimize water saturation effects. Here we examined the level of water saturation during the T period caused by various types of inversion proton pulses to suppress CC: (I) amide-selective IBURP-2; (II) cosine-modulated IBURP-2; (III) Watergate-like blocks; and (IV) non-selective hard. We additionally demonstrate the effect of uncontrolled saturation of

Electronic supplementary material The online version of this article (doi:10.1007/s10858-015-9937-4) contains supplementary material, which is available to authorized users.

Margarida Gairí mgairi@rmn.ub.edu

- Ernest Giralt ernest.giralt@irbbarcelona.org
- ¹ NMR Facility, Scientific and Technological Centers, University of Barcelona (CCiTUB), Baldiri Reixac 10, 08028 Barcelona, Spain
- ² Institute for Research in Biomedicine (IRB), Baldiri Reixac 10, 08028 Barcelona, Spain
- ³ Biomolecular NMR Laboratory and Organic Chemistry Department, University of Barcelona, Baldiri Reixac 10, 08028 Barcelona, Spain

aliphatic protons on ¹⁵N R₁ rates. In this paper we present an optimized pulse sequence that takes into account the crucial effect of controlling also the saturation of the aliphatic protons during ¹⁵N R₁ measurements in nondeuterated proteins. We show that using cosine-modulated IBURP-2 pulses spaced 40 ms to cancel CC in this optimized pulse program is the method of choice to minimize systematic errors coming from water and aliphatic protons saturation effects.

Introduction

The intrinsic dynamic properties of proteins play a key role in their function. Information about dynamics on several timescales can be studied through solution NMR spectroscopy using ¹⁵N spin relaxation experiments on ¹⁵Nlabeled protein samples. Longitudinal relaxation rate R₁, transverse relaxation rate R₂, and ¹⁵N–¹H steady-state NOE parameters of backbone ¹⁵N nuclei measured at various magnetic fields are commonly used to address global and local protein dynamics at ps–ns and μ s–ms timescales (Torchia 2011).

Quantitative analysis of backbone dynamics requires accurate and consistent relaxation measurements (Morin 2011). It has long been recognized that systematic errors in the measurements of ¹⁵N relaxation rates R_1 , R_2 and ¹⁵N–¹H NOE can be significant and far larger than random errors. Two main sources of these errors have been identified as amide proton exchange saturation transfer from water protons (Grzesiek and Bax 1993; Farrow et al. 1994; Ferrage et al. 2010; Chen and Tjandra 2011; Lakomek et al. 2012; Jurt and Zerbe 2012) on the one hand, and ¹⁵N CSA/¹H–¹⁵N dipolar cross-correlated (CC) relaxation (Boyd et al. 1990; Kay et al. 1992a; Palmer et al. 1992; Kumar et al. 2000; Gong and Ishima 2007; Ferrage et al.

2009; Ishima 2014) on the other. Water saturation has a strong effect on ¹⁵N R₁ measurements (Chen and Tjandra 2011). It can arise as a consequence of improper handling of water magnetization effects of RF pulses and/or pulsed field gradients (PFG) during the whole pulse program (Stonehouse et al. 1994). Saturation can be transferred through protonproton exchange or via NOE to exchangeable amide protons-directly detected in HSOC-based R1 measurements-during the inter-scan delay and it distorts proton amide signal intensities. In addition, varying degrees of water saturation during variable ¹⁵N magnetization decay periods can alter the intensity decay, thus causing systematic errors in R₁ relaxation rates (Lakomek et al. 2012). Cryogenic probeheads (Kovacs et al. 2005) exacerbate water suppression-related problems by strong radiation damping (RD) effects (Krishnan and Murali 2013) especially at high magnetic fields (Shishmarev and Otting 2011).

Another important source of systematic errors on R_1 measurements is longitudinal CC between ${}^{1}H^{-15}N$ DD and ${}^{15}N$ CSA. The CSA component increases also with magnetic fields. Insufficient cancellation of cross-correlation effects at very high fields is reported to generate significant deviations in ${}^{15}N$ R₁ rates for deuterated proteins (Ishima 2014). CC is typically suppressed by applying a series of proton 180° pulses during the variable ${}^{15}N$ relaxation period (Kay et al. 1992a). However, these pulses can perturb water magnetization if care is not taken to prevent it. In addition, they can cause unwanted saturation effects on protein protons. In particular, saturation of aliphatic protons in non-deuterated proteins may cause significant errors on ${}^{15}N$ R₁ measurements, a phenomenon that has not been properly addressed so far.

In this paper we present an optimized method to determine ¹⁵N R_1 rates on non-deuterated proteins, which can be used at high fields using cryogenic probes, in order to address the problems described above. The optimized sequence relies on the control of both water and aliphatic protons saturation while CC-suppressing pulses are applied during the variable ¹⁵N relaxation period.

Experimental section

Our study was carried out with a fully protonated ¹⁵N-labeled sample of the small and highly stable third Iggbinding domain of protein G (GB3). The NMR sample consisted of 1.5 mM ¹⁵N-labeled GB3 in 25 mM phosphate buffer, 25 mM NaCl, 0.02 % NaN₃ and 10 % D₂O, pH 6.5. ¹⁵N longitudinal relaxation experiments were performed on a 600 MHz Bruker Avance III and an 800 MHz Bruker DRX NMR spectrometer operating at 60.79 and 81.06 MHz ¹⁵N resonance frequencies, respectively. Both instruments were equipped with TCI cryogenic probes and a Z-gradient coil. Data were recorded at 298 K.

We used 1D versions of the pulse scheme shown in Fig. 1 to measure the water signal intensity—reporting the level of non-saturated water (M_{nsw}) —using a short angle readout pulse. Readout pulses were introduced just before the ¹⁵N relaxation period T (point *a*) and at the start of the regular ¹H acquisition time (point b). Experiments were carried out at 600 and 800 MHz. Simultaneous phase cycling (0°, 180°) of the receiver and the readout pulse allowed the cancelation of the protein signal. Readout pulses shorter than 1° were used in order to avoid artifacts caused by RD. The resulting water signal was integrated (I^{x}) in the region 6.5-2.9 ppm (spectral width of 3.65 ppm which corresponds to a range equal to 50 times the observed full peak width at half maximum). Water saturation is given with respect to the signal intensity measured following a single short angle proton pulse (I^{ref}). The water saturation level (M_{sw}) was calculated as:

$$M_{sw}^{x} = \left(1 - M_{nsw}^{x}\right) \times 100 = \left(1 - \frac{I^{x}}{I^{ref}}\right) \times 100$$

where I^x is the signal intensity at point x = a or b and I^{ref} is the intensity of the reference signal.

The relative water saturation at point b with respect to point a, as a percentage, is given by:

$$M_{sw}^{b/a} = \left(\frac{M_{nsw}^a - M_{nsw}^b}{M_{nsw}^a}\right) \times 100 = \left(\frac{I^a - I^b}{I^a}\right) \times 100$$

1D experiments were performed using fixed ¹⁵N relaxation periods of 80 and 800 ms. The first t₁ increment (10-12 µs) was measured in all cases. Saturated water was evaluated for several approaches to cancel CC, so the following 180° proton pulses (P) types were tested during T (Fig. 2): (I) a train of off-resonance amideselective IBURP-2 pulses (Geen and Freeman 1991) (1.9 ms length, offset 2400 Hz with respect to water frequency, at 600 MHz; 1.5 ms length, offset of 3000 Hz with respect to water frequency, at 800 MHz); (II) a train of cosine-modulated off-resonance IBURP-2 pulses which selectively invert two spectral ranges simultaneously, centered at +2400 and -2400 Hz (+3000 and -3000 Hz) with respect to the water frequency at 600 MHz (800 MHz), pulse lengths were 1.9 and 1.5 ms at 600 and 800 MHz, respectively, power levels were increased by 6 dB with respect to those used for IBURP-



Fig. 1 Pulse program for measuring ${}^{15}N$ R₁ relaxation rates in non-deuterated proteins, through a ${}^{1}H{}^{-15}N$ HSQC-based experiment. Narrow filled rectangles and open rectangles correspond to 90° and 180° flip angle pulses, respectively, with phase x unless indicated. The open bell in proton channel corresponds to a 90° sinc-shaped low-power water-selective pulse, a (sinx)/x function, with a duration of 1 ms. The *filled bells* (P) applied to proton during the variable ${}^{15}N$ relaxation delay T correspond to 180° pulses to cancel CC during this period. They are applied as a train of pulses spaced by intervals of duration δ , in such a way that the loop is repeated an even number of times. The number of 180° pulses (n) is varied to yield different ^{15}N relaxation delays (T = $2 \times n \times \delta$). Several types of 180° inversion elements (P) have been used to cancel CC in different R_1 experiments: (I) off-resonance selective amide proton IBURP-2 pulses (1.9 ms duration, offset = 2400 Hz at 600 MHz; 1.5 ms, offset = 3000 Hz at 800 MHz); (II) cosine modulated IBURP-2 pulses (1.9 ms duration, offset = ± 2400 Hz at 600 MHz; 1.5 ms, offset = ± 3000 Hz at 800 MHz); (III) Watergate-like block consisting of three pulses: water-selective low-power 90° pulse (1 ms,

2 pulses to account for the second frequency component; (III) a train of Watergate-like (Piotto et al. 1992) blocks, each one consisting of a 90° water-selective pulse (1 ms, square pulse) followed by a 180° hard pulse (20–24 μ s) and after by another 90° water-selective pulse (1 ms, square pulse) and (IV) a train of hard 180° pulses (20–24 μ s).

HSQC-based R₁ experiments were performed with the pulse sequence shown in Fig. 1 using the previously described CC-suppressing approaches. R1 rates at 600 and 800 MHz are shown in Fig. 3. The spectral widths for experiments carried out at 600 MHz were set to 8418 Hz (F2) and 2189 Hz (F1) for ¹H and ¹⁵N, respectively, with sampling durations of 107 ms (t_2) and 58 ms (t_1) . For experiments measured at 800 MHz, spectral widths of 11,160 Hz (F2) and 2838 Hz (F1) for ¹H and ¹⁵N, respectively, and sampling durations of 67 ms (t₂) and 53 ms (t_1) were used. Radio-frequency carriers were set to 4.7 ppm for ¹H and 118 ppm for ¹⁵N. For each relaxation time measurement, a series of eight 2D experiments with ¹⁵N relaxation delays T ranging from 0 to 800 ms (0, 80, 160, 240, 320, 400, 560, 800 ms) were collected in a randomized order. Eight scans per FID

square shaped pulse)-180° hard pulse (20-24 µs length at high power)-water-selective low-power 90° pulse (1 ms, square shaped pulse); (IV) non-selective hard pulses (20-24 µs at high power). Interspace delays $\delta = 5 \text{ ms}$ (pulsing rate $R_p = 200 \text{ s}^{-1}$) or $\delta = 40 \text{ ms}$ (pulsing rate $R_p = 25 \text{ s}^{-1}$) were used. $\tau = 2.78 \text{ ms}$. All PFG are along z-axis, smoothed-square shaped. Duration and strength are indicated in parenthesis: G1 (1 ms; 25 G/cm), G2 (1 ms; 40 G/cm), G3 (1 ms; 8.1 G/cm), G4 (1 ms; -25 G/cm at 600 MHz or 11.5 G/ cm at 800 MHz), G5 (1 ms; 4.5 G/cm), G6 (1 ms; 8.5 G/cm), G7 (0.5 ms; 15.5 G/cm), G8 (0.5 ms; 5.5 G/cm), G9 (0.5 ms; 6.5 G/cm), G10 (0.6 G/cm). Phase cycling: $\Phi_1 = 4(x)$, 4(-x); $\Phi_2 = y$, -y; $\Phi_3 = 2(x), 2(-x); \Phi_4 = 2(x), 2(-x); \Phi_5 = 2(y), 2(-y); \Phi_{rec} = x,$ -x, x, x, -x, x, x, -x. Quadrature detection in F1 is implemented using the gradient-enhanced echo/anti-echo scheme (Kay et al. 1992b) inverting the polarity of PFG G2 and also the phase Φ_5 for the second FID. 15N GARP decoupling sequences (Shaka and Keeler 1987) were applied during ¹H acquisition at a field strength of around 1 kHz

were recorded, and a recycle delay between scans of 1.7 s was used.

The aliphatic proton magnetization at the beginning of the variable ¹⁵N relaxation period in the pulse sequence of Fig. 1 is dephased in the xy-plane (XY-pulse program). Alternative sequences for measuring ¹⁵N R₁ rates have this magnetization on the z-axis (Z-pulse program). The 1D versions of the XY- and Z-sequences (Fig. 4) were used to evaluate the amount of amide proton (H^N) and aliphatic proton (H^R) polarization at several points of the pulse scheme. Longitudinal proton magnetization was measured with a 90° readout pulse followed by a Watergate block for water suppression (see box in Fig. 4) in the following points: (i) just before the ¹⁵N relaxation period (T); (ii) just after this period T, for T = 80 and T = 800 ms, and with amide-selective IBURP-2 pulses spaced 40 ms or cosine-modulated IBURP-2 pulses spaced 40 ms as inversion elements to cancel CC; (iii) at the final point of the pulse program after a delay (Δ) added to allow for proton magnetization to relax to equilibrium (+z) for a time period identical to the recycle delay in the pulse scheme used to measure ${}^{15}N$ R₁. In this last case, two values of Δ were evaluated: 1.7 and

Fig. 2 1 H spectrum of 15 N– GB3 and experimental excitation profiles at 800 MHz for several types of proton inversion elements used to cancel CC during the 15N relaxation period T: (I) IB-2: amide-selective IBURP-2 pulse (1.5 ms, offset 3000 Hz from water frequency); (II) cm IB-2: cosine-modulated IBURP-2 pulse (1.5 ms, offset \pm 3000 Hz from water frequency); (III) wg-like: Watergate-like pulses: waterselective 90°_{-x} pulse (1 ms, *square*)—hard 180°_{x} pulse (24 µs)—water-selective 90°_ pulse (1 ms, square); (IV) hard: non-selective hard 180° pulse (24 µs). Excitation profiles for I, II and IV were measured with a pulse program consisting of the selective 180° pulse, followed by a PFG and a final short readout pulse, while profile for scheme III was generated using a pulse program consisting of an initial 90° pulse, followed by a 90° sel-180° hard-90° sel block inserted in a PFG echo and a final shord readout pulse. Each spectrum was acquired with a single scan and the offsets vary along 12.000 Hz in 100 Hzsteps. A CuSO₄-doped water sample in D₂O was used



3.5 s. Thus, proton magnetization at point (iii) represents the amount of polarization at the beginning of the sequence under steady-state conditions. Experiments were performed at 800 MHz. The first t_1 increment (12 µs) was acquired. Eight scans were measured and a recycle delay d1 = 10 s was used. H^N and H^R polarization were evaluated by relative integration of the amide region between 9.55 and 7.25 ppm and relative integration of the aliphatic region between 2.25 and 0.25 ppm, respectively, in the corresponding GB3 ¹H spectra (see Fig. 5). Relative integrations were calculated respect to the amount of the equilibrium H^N or H^R polarization measured at point (iii) for $\Delta = 3.5$ s in the corresponding sequence.

 15 N R₁ rates measurements performed with Z-pulse program at 800 MHz were carried out in the same conditions as previously described for XY-pulse scheme (Fig. 1).



Fig. 3 ¹⁵N R₁ relaxation rates measured at 600 MHz (**a**) and 800 MHz (**b**) for non-deuterated ¹⁵N–GB3 using several CC-suppressing schemes during the variable ¹⁵N relaxation delay T. A proton pulsing rate R_p = 25 s⁻¹ (180° pulses *P* spaced δ = 40 ms, see Fig. 1) was used in all cases. **a**, **b** R₁ rates measured using IBURP-2 pulses (I), cosine-modulated IBURP-2 (II), Watergate-like

pulses (III) or hard pulses (IV) for suppressing CC; **c** correlation plot of R_1 rates measured at 600 MHz with approach (I) and with a deuterated GB3 sample (Lakomek et al. 2012); **d** correlation plot of R_1 rates measured at 600 MHz with approach (II) and with deuterated GB3; **e** correlation plot of R_1 rates measured at 600 MHz with approach (III) and with deuterated GB3

All spectra were processed with Topspin from Bruker, and relaxation time constants and fitting errors were extracted using the free software R programming language (www.r-project.org).

Results and discussion

Exploring the effect of water saturation caused by CC-suppressing pulses on 15 N R₁ measurements

Two sources of non-equilibrium water magnetization can be considered in experiments for measuring ¹⁵N R₁ relaxation in proteins: partially saturated water coming from water magnetization handling before the variable ¹⁵N relaxation period (Chen and Tjandra 2011) on one hand, and additional water saturation created as a side effect of the 180° proton pulses used to cancel CC effects during this period (Lakomek et al. 2012) on the other. Preventing water saturation while efficiently suppressing CC is crucial for accurately measuring R_1 rates. The pulse scheme depicted in Fig. 1 is designed for this purpose: PFG around 180° pulses during spin-echoes when water is transverse minimize RD; a water-flip-back pulse during the initial INEPT keeps water at +z axis; water magnetization is brought to +z at the start of ¹⁵N relaxation period T; bipolar PFG are used during the second part of t_1 evolution time where water remains at -z axis to reduce RD; and PFG are immediately applied to minimize RD every time water is placed at -z axis.

The pulse sequence in Fig. 1 can be divided into two blocks, A and B. Ideally, water magnetization should reach the +z axis just after each of these modules (points *a*, *b*) and retain equilibrium magnetization (no saturation). However, as a consequence of imperfections of RF pulses, water dephasing and rephasing caused by PFG, and strong RD effects present at high fields with cryoprobes, the path for water magnetization deviates from this ideal situation



Fig. 4 Pulse programs (XY and Z) to evaluate the amount of H^R and H^N polarization present at different points in the sequence used to determine ¹⁵N R₁ relaxation rates. A 90° readout pulse followed by a Watergate block (*boxed*) was applied in either point (i), (ii) or (iii). Δ corresponds to the recycle delay used in R₁ measurements. Therefore, point (iii) corresponds to the start of the scans under steady-state conditions. The first t₁ increment (12 µs) was acquired in the 1D versions of the pulse program. In the XY-pulse scheme the aliphatic

and is partially saturated. By comparing water saturation at points *a* and *b* in the pulse sequence, we can evaluate the additional water saturation introduced in block B, which includes mostly the effect of the CC-suppression scheme. Since CC-suppression is applied for a variable duration, water saturation can modify the intensity of amide proton decay differently for each value of T and introduce systematic errors in the measurement of ¹⁵N R₁ relaxation rates.

The most widely used approaches to suppress CC are: (i) amide-selective IBURP-2 pulses (Geen and Freeman 1991), inverting only magnetization of the amide protons (Lakomek et al. 2012; Chill et al. 2006); (ii) cosinemodulated selective pulses (Smallcombe 1993), exciting both the amide and the aliphatic region of the protein while performing zero excitation at the water resonance (Farrow et al. 1994; Zhu et al. 2000); and (iii) combination of selective and non-selective pulses in a Watergate-like (Piotto

protons magnetization remains in the transverse plane and is saturated by the PFG G1 applied before the T period. In the Z-pulse sequence the aliphatic protons magnetization is aligned along the z-axis at the start of this period. The same pulses, delays and PFG as those described in Fig. 1 were used. Phase cycling: $\Phi_1 = 4(x)$, 4(-x); $\Phi_2 = y, -y; \Phi_3 = 2(x), 2(-x); \Phi_4 = 2(x), 2(-x); \Phi_5 = 2(y), 2(-y);$ $\Phi_6 = x, -x; \Phi_7 = 2(x), 2(y), 2(-x), 2(-y); \Phi_8 = 2(-x), 2(-y),$ $2(x), 2(y); \Phi_{rec} = x, -x, x, x$

et al. 1992) way to selectively invert non-water protons (Chen and Tjandra 2011).

Here we evaluated the water saturation level caused by the three CC-suppressing schemes described above using a 1D version of the pulse program depicted in Fig. 1. Nonselective hard 180° proton pulses were also included in our study as a reference. Experiments were carried out at 800 and 600 MHz, using TCI cryoprobes with the following proton pulses: amide-selective 180° off-resonance IBURP-2 pulses (I); cosine-modulated 180° IBURP-2 pulses (II); Watergate scheme based on the block 90° water-selective pulse— 180° hard pulse— 90° water-selective pulse (III); and non-selective 180° hard pulses (IV). Figure 2 shows the corresponding experimental excitation profiles at 800 MHz. Experimental profiles at 600 MHz are presented in Fig. S1 SM. Selective excitation schemes should minimally affect water magnetization in contrast to the use of hard pulses. However, residual water saturation present



Fig. 5 a ¹⁵N R₁ relaxation rates measured at 800 MHz for nondeuterated ¹⁵N–GB3 using XY- and Z-pulse schemes (see Fig. 4), and CC-suppressing schemes based on amide-selective IBURP-2 pulses or cosine-modulated IBURP-2 pulses ($R_p = 25 \text{ s}^{-1}$), applied during the variable ¹⁵N relaxation delay (T). Deviations on average R₁ rates respect to those obtained with the XY-sequence and cosine-modulated IBURP-2 pulses, ($<R_1 > - <R_1 >_{XYcmIB2}$)/ $<R_1 >_{XYcmIB2}$ are shown inside the *box.* **b** ¹H spectra of GB3 at 800 MHz measured using the XY and Z-pulse schemes in Fig. 4. The spectra show the amount of non-saturated H^N and H^R magnetization at point (ii), measured

after the various schemes is significantly different, in practice.

Water saturation levels were measured after the application of the above mentioned CC-suppressing schemes for short (80 ms) and long (800 ms) 15 N relaxation periods T. In addition, the effect of pulsing rate (R_p) during

immediately after the application of a train of amide-selective IBURP-2 pulses or cosine-modulated IBURP-2 pulses (with 40 ms inter-pulse spacing) for relaxation periods T of 80 and 800 ms. The relative integral of the aliphatic region between 2.25 and 0.25 ppm is given in the *right* part of each spectrum. The relative integral of the amide region between 9.55 and 7.25 ppm measured at point (iii) with $\Delta = 1.7$ s is given inside the *boxes* at the *left* part of the spectra. Percentages are calculated respect to the amount of the equilibrium H^R or H^N polarization measured at point (iii) for $\Delta = 3.5$ s in the corresponding sequence

cancellation of CC was also studied using two rates: fast pulsing regime (200 s⁻¹, proton pulses spaced 5 ms) and slow pulsing regime (25 s⁻¹, proton pulses spaced 40 ms). Results are presented in Table 1.

The level of water saturation at point a was about 20 %, both at 800 and 600 MHz. The amount of saturated water

	(I) IB-2		(II) cm IB-2		(III) wg-like		(IV) hard	
	$25 \ s^{-1}$	$200 \ s^{-1}$	25 s^{-1}	$200 \ s^{-1}$	25 s^{-1}	$200 \ s^{-1}$	$25 \ s^{-1}$	$200 \ s^{-1}$
Degree of saturate	ed water (%) at	600 MHz						
$M^{b/a}_{sw(80ms)}$	10	14	10	13	11	25	25	90
M ^{b/a} _{sw(800ms)}	8	23	6	29	10	99	96	97
$\Delta M^{b/a}_{sw(800-80)}$	-2	9	-4	16	-1	74	71	7
Degree of saturate	ed water (%) at	800 MHz						
$M^{b/a}_{sw(80ms)}$	11	14	10	13	9	19	80	79
M ^{b/a} _{sw(800ms)}	8	23	7	26	10	63	97	99
$\Delta M^{b/a}_{sw(800-80)}$	-3	9	-3	13	1	44	17	20

Table 1 Level of saturated water generated between the start of 15 N relaxation delay (point *a*) and the start of proton acquisition period (point *b*) in the pulse program of Fig. 1

Water saturation (M_{sw}) is calculated, as described in "Experimental section", for different schemes to suppress CC: (I) IB-2: a train of amideselective IBURP-2 180° proton pulses; (II) cm IB-2: a train of cosine-modulated IBURP-2 180° proton pulses; (III) wg-like: a train of Watergatelike blocks consisting of 90° water-selective pulse—180° hard proton pulse—90° water-selective pulse; (IV) hard: a train of 180° hard proton pulses. Water saturation has been measured for two proton pulsing rates (R_p), 25 and 200 s⁻¹. For each one of the CC-suppressing schemes, two lengths of the ¹⁵N relaxation period (T), 80 and 800 ms, have been evaluated. $\Delta M_{sw(800-80)}^{b/a}$, calculated as the difference $M_{sw(800ms)}^{b/a} - M_{sw(80ms)}^{b/a}$, refers to the increment of water saturation occurred at long T periods. Results at 600 and 800 MHz are shown

generated between points *a* and *b*, $M_{sw}^{b/a}$, can be attributed mainly to the effect of proton pulses applied during ¹⁵N variable relaxation delay (T period) to cancel CC. Nonselective hard pulses caused extreme water saturation (80– 100 % in most cases), as expected. IBURP-2 pulses, cosine-modulated IBURP-2 pulses, and Watergate-like schemes gave comparable results at slow pulsing rates ($R_p = 25 \text{ s}^{-1}$), with levels of saturation around 10 %, and no significant differences ($\Delta M_{sw(800-80)}^{b/a}$) between short and long ¹⁵N relaxation periods.

Faster pulsing rates increase both the level of water saturation and the difference in saturation levels between short and long ¹⁵N relaxation delays, $\Delta M_{sw(800-80)}^{b/a}$. At pulsing rates of $R_p = 200 \text{ s}^{-1}$, water saturation was higher for the Watergate scheme than for IBURP-2-based sequences. In addition, a significant increase in the level of saturation was observed between the short and the large relaxation delay when Watergate-like blocks were used to suppress CC ($\Delta M_{sw(800-80)}^{b/a} = 74\%$ at 600 MHz and 44 % at 800 MHz, see Table 1). In contrast, when using IBURP-2-based sequences, water saturation differences between short and large T periods were less important, around 9–16 %.

These results indicate that the use of an inter-pulse delay of 40 ms ($R_p = 25 \text{ s}^{-1}$) allows the build-up of a steady-state polarization at the start of each scan, in the three approaches tested. Therefore the effect of water saturation is independent of the ¹⁵N relaxation delay and it should not affect the measured ¹⁵N R₁ rates. On the contrary, water

saturation increased along this relaxation period when faster pulsing rates were used ($R_p = 200 \text{ s}^{-1}$). Because this effect was stronger for Watergate-like than for IBURP-2-based sequences (IBURP-2 or cosine-modulated IBURP-2 pulses), the latter are expected to show better performance at cancelling CC with minimum water saturation on ¹⁵N R₁ measurements.

To study the impact of the previously CC-suppressing schemes on ¹⁵N longitudinal rates of non-deuterated proteins, eight protocols were compared for the measurement of ¹⁵N R₁ at 800 and 600 MHz. They included four proton inversion sequences (Fig. 2 and Fig. S1 SM) at two pulsing rates R_p (25 and 200 s⁻¹). The individual apparent R₁ rates obtained at R_p = 25 s⁻¹ are shown in Fig. 3. The average R₁, the average pairwise root mean square deviation (*rmsd*) and the correlation coefficients, R, are presented in Table 1 SM.

Measurements carried out at 600 MHz in minimal water saturation conditions (approaches I, II and III) at slow pulsing rates $R_p = 25 \text{ s}^{-1}$ resulted in identical R_1 rates for methods II (cosine-modulated IBURP-2) and III (Watergate-like pulses), while systematically shorter R_1 values were obtained for all residues in GB3 when using approach I (IBURP-2), see Fig. 3a. The same effect was observed at 800 MHz (Fig. 3b). The differences in R_1 values obtained by methods I and II were small (rmsd = 0.075 s⁻¹, 3.15 % at 600 MHz and 0.077 s⁻¹, 3.76 % at 800 MHz) but systematic and significant, and larger than the reproducibility of the individual experiments (rmsd = 0.030 s⁻¹, 1.2 % for duplicate measurements). R_1 rates measured at 600 MHz using the cosine-modulated IBURP-2 pulses at $R_p = 25 \text{ s}^{-1}$ (approach II) are in perfect agreement (rmsd = 0.022 s⁻¹, 0.9 %, Fig. 3d) with those reported for deuterated GB3¹ at the same field (Lakomek et al. 2012). The same happened for the Watergate-like scheme (Fig. 3e).

The use of hard 180° proton pulses (approach IV), causing strong water saturation, resulted in clear overestimation of R₁, at both 25 and 200 s⁻¹ pulsing rates (Table 1 SM). In addition, errors coming from water saturation effects varied along the protein sequence, being largest in the most solvent-exposed regions of the GB3 protein: residues 15–23 (A region) and 44–55 (B region), Fig. 3a, b. The effect of water saturation can be estimated by the clear correlation between the fractional increase in R₁ rates, [R₁(IV) – R₁(II)]/R₁(II), and the degree of saturation I_{sat}/I₀ measured in ¹H–¹⁵N HSQC spectra (Mori et al. 1995) acquired with and without water presaturation (Fig. S2 SM).

A common alternative to suppress the effect of nonuniform water saturation in R1 measurements performed with the classical CC-suppressing scheme employing hard 180° pulses (scheme IV, $R_p = 200 \text{ s}^{-1}$) consists of explicitly purging all proton polarization at the start of the inter-scan delay, using a 90 ¹H pulse followed by a pulse field gradient. Taking the R1 values obtained at 600 MHz using the cosine-modulated IBURP-2 ($R_p = 25 \text{ s}^{-1}$) scheme as reference (approach II) we compared the R₁ measured by approach IV with and without the purging element. Indeed, the use of the initial purge in scheme IV provides a better agreement with the reference experiment $(\text{rmsd} = 0.065 \text{ s}^{-1}, 2.6 \%, R = 0.931)$ than approach IV without purging (rmsd = 0.392 s^{-1} , 9.3 %, R = 0.195). However such an approach suffers from a significant loss of sensitivity (up to 35 % in the present experiment) and therefore scheme II should be clearly preferred.

Evaluating the effect of R_p on R_1 values measured with approaches II and III, showed that suppression of longitudinal CC at faster proton pulsing rate ($R_p = 200 \text{ s}^{-1}$) resulted in slightly higher R_1 values for approach II (cosinemodulated IBURP-2), and significantly larger R_1 rates for approach III (Watergate-like), with respect to those measured at slower pulsing rates ($R_p = 25 \text{ s}^{-1}$). Figures S3 SM and S4 SM show results at 600 and 800 MHz, respectively. R_1 differences along GB3 amino acid sequence between the two proton pulsing rates were not uniform. Higher deviations appeared for residues 15–23 (A region) and 44–55 (B region), coincident with regions more affected by water saturation effects (approach IV, hard pulses, Fig. 3). This observation is consistent with results shown in Table 1.

Sequence specific R_1 deviations depend also on the characteristics of the excitation profile of these selective pulses. With cosine-modulated IBURP-2 pulses, showing a wide zero-excitation region around water resonance (Fig. 2), R_1 differences lower than 3 % were observed in regions A and B. However, the use of Watergate-like pulses yielded R_1 rates around 6–10 % higher at faster pulsing rates (see Table 2 SM). Degradation of the performance of Watergate-like pulses at high proton pulsing rates may result from overuse of water-selective 90° pulses, continuously exciting water magnetization from +z-axis to the transverse plane and returning it back from +x/+y to +z. Whenever water magnetization is moved from equilibrium it remains in a vulnerable state, due to the strong RD effects, so it is more prone to be in a saturated state.

Our results showed that short pulsing rates resulted in minimal water saturation during the ¹⁵N relaxation period. It is clear that $R_p = 25 \text{ s}^{-1}$ is enough to cancel longitudinal CC relaxation in GB3 at 600 MHz. For larger non-deuterated proteins, longitudinal CC effects will be smaller because proton spin-flips, resulting from effective proton-proton cross-relaxation, induce exchange between the longitudinal doublet components at a rate faster than their R_1 difference. The larger the protein, the higher the amide proton spin-flip rate would be. Therefore a pulsing rate of 25 s⁻¹ should be high enough to cancel these undesired relaxation mechanisms when measuring ¹⁵N R_1 rates at 600 MHz or higher fields on non-deuterated proteins similar to GB3 in size or larger.

Insufficient cancellation of CC with the Watergate approach as a result of both off-resonance effects and pulse imperfection of the central 180° hard proton pulse in the Watergate block has been reported at very high fields (Ishima 2014). With a cosine-modulated IBURP-2 based scheme no 180° hard pulses are involved in CC cancellation. Moreover, larger dispersion in proton spectral widths at higher fields allows the application of these selective pulses at increased offsets from the water resonance, thereby minimizing residual effects of water saturation during the T period. The higher the field, the lower the residual effects would be. Consequently, the use of cosine-modulated IBURP-2 is preferable to Watergate-based schemes in order to obtain more accurate $^{15}N R_1$ rates.

Exploring the effect of saturation of aliphatic protons caused by CC-suppressing IBURP-2-based schemes on ¹⁵N R₁ measurements

While R_1 rates obtained at 600 MHz with the cosinemodulated IBURP-2 approach were very close (Fig. 3d) to

¹ The effect of deuteration of non-exchanging protons on ¹⁵N R₁ rates is expected to be insignificant for a protein the size of GB3. Thus ¹⁵N R₁ for a protonated sample of GB3 and for a deuterated one should be very similar. In addition, R₂ rates measured at 600 MHz with non-deuterated GB3 were also coincident with those obtained for deuterated GB3 (Lakomek et al. 2012); rmsd = 0.072 s⁻¹ (1.28 %).

those reported for deuterated GB3 at the same field (Lakomek et al. 2012), the amide-selective IBURP-2 scheme gave shorter R_1 values (Fig. 3c). Systematically shorter rates along the whole GB3 amino acid sequence could not derive from water-saturation effects, because saturated water is very similar for both schemes at $R_p = 25 \text{ s}^{-1}$ (Table 1). The origin of this discrepancy could be inefficient CC suppression or additional saturation effects caused, directly or indirectly, by CC-suppressing pulses, which may affect amide proton magnetization at the beginning of each scan. Differential saturation effects should be sensitive to the recycle delay while cross-correlation should not be affected by this delay.

¹⁵N R₁ rates were measured at 600 MHz using IBURP-2 and cosine modulated IBURP-2 with 1.7 and 3.5 s of recycle delay values. While R1 rates measured with the second approach were unaffected by recycle delay, R1 at long recycle delays using the IBURP-2 approached were larger than values at short delays (Fig. S5 SM). The same happened at 800 MHz. In addition, R1 rates at 600 MHz with IBURP-2 pulses and 3.5 s of recycle delay approached to rates reported for deuterated GB3 protein (Lakomek et al. 2012). These results suggested that the observed deviation originates from non-equilibrium magnetization of aliphatic protons caused by the effect of IBURP-2 or cosine modulated IBURP-2 pulses. While aliphatic protons (H^R) are affected by cosine modulated IBURP-2 pulses, IBURP-2 only excites the amide region $(\mathbf{H}^{\mathbf{N}}).$

The effect of the CC-suppressing train of pulses will depend on the aliphatic magnetization present at the beginning of the variable ¹⁵N relaxation period (T). In the pulse sequence shown in Fig. 1 the aliphatic protons magnetization remains on the transverse plane before the T period (XY-pulse program) and is saturated by the PFG G1 applied during the INEPT block. Other pulse sequences used to measure ¹⁵N R₁ rates (Lakomek et al. 2012) keep aliphatic protons magnetization aligned along the z-axis (Z-pulse programs) at the start of this period. In both XY- and Z-pulse programs, aliphatic protons magnetization will be affected by CC-suppressing pulses, but in a different way. Thus, depending on the pulse program and on the approach used to cancel CC a T-dependent aliphatic protons saturation level will develop during T.

Of course, the state of non-amide protons magnetization has a significant effect on effective longitudinal relaxation time of amide protons (Pervushin et al. 2002; Schanda 2009), the protons excited at the start of HSQC-based pulse programs to measure ¹⁵N R₁ rates. Fast H^N relaxation relies on keeping the non-amide magnetization fully aligned along +z, in a non-perturbed state, so saturation of nonamide spins reduce the effective longitudinal H^N rates (Schanda 2009). In consequence, slower amide relaxation will take place at high levels of aliphatic spins saturation respect to low saturation states. Slow H^N recovery rates will result in a smaller amount of H^N magnetization at the start of the pulse program for a fixed recycle delay versus high H^N recovery rates. Thus, an increase in H^R magnetization saturation at long versus short T periods is expected to produce artificially higher ¹⁵N R₁ rates, respect to a T-independent aliphatic protons saturation level. On the contrary, shorter ¹⁵N R₁ rates would result from a decreased saturation of aliphatic protons at long respect to short T periods.

We used 1D versions of the pulse programs (XY- and Z) shown in Fig. 4 to measure H^N and H^R magnetization after a short (80 ms) and a long (800 ms) ¹⁵N relaxation period, in which IBURP-2 or cosine modulated IBURP-2 sequences were applied to cancel CC. Measurements were performed at 800 MHz. Figure 5 and Table 3 SM show a comparison of the available longitudinal magnetization (non-saturated) after the ¹⁵N relaxation period (point ii) as a function of the CC-suppression scheme. Also, magnetization was evaluated at point (iii), after a Δ period equivalent to the recycle delay used on the R₁ measurement experiments, representing the available proton magnetization at the start of the scan. The amount of H^R and H^N magnetization present after $\Delta = 1.7$ s (recycle delay used on ${}^{15}N$ R₁ measurements shown in Table 1 SM) was calculated respect to the magnetization measured with $\Delta = 3.5$ s, and the percentage of H^N magnetization is represented in the box at the left over the corresponding 1D spectra (Fig. 5b). In addition, ¹⁵N R₁ rates were also measured at 800 MHz with the Z-pulse program using the two CC-suppressing schemes (IBURP-2 and cosine modulated IBURP-2). Average R₁ rates and the individual apparent R₁ rates for both XY and Z-pulse programs are presented in Fig. 5a.

With the XY pulse program no signal from aliphatic protons was observed just before the ¹⁵N relaxation period (point *i* in Fig. 4, Table 3 SM). The longitudinal aliphatic magnetization grows during the ¹⁵N relaxation time when the aliphatic protons are not excited, this is when CC is suppressed by means of amide-selective IBURP-2 pulses. Because of this, a higher amount of longitudinal H^R magnetization has developed at long T periods (800 ms), resulting in a faster recovery of the H^N magnetization respect to short T periods (80 ms), see Fig. 5b. A high degree of H^N at long periods, at the start of each scan, explains the shorter ¹⁵N R₁ rates measured with IBURP-2 (Fig. 5). On the contrary, the train of cosine-modulated IBURP-2 pulses maintains a full saturation of aliphatic protons constant during the T period so the recovery rate of the H^N magnetization is unaffected. Thus, ¹⁵N R₁ relaxation rates measured with this scheme are not subject to saturation effects.

On the other hand, when the aliphatic protons magnetization is placed on the z-axis at the beginning of the ¹⁵N relaxation T period (Z pulse program), cosine-modulated IBURP-2 pulses cause T period-dependent variable degrees of saturation, the larger the relaxation period the larger the saturation level (Fig. 5b). As a consequence, the amount of H^N polarization at the beginning of each scan is reduced for long T values, resulting in artificial and significant increased ¹⁵N R₁ rates (Fig S6 SM). The use of amideselective IBURP-2 pulses to cancel CC, in principle, should not affect the aliphatic protons as their magnetization is not directly perturbed. However, amide protons inversion during the train of IBURP-2 pulses perturbs aliphatic protons magnetization by cross-relaxation and the non-equilibrium aliphatic magnetization eventually feedsback to the amide protons in a way that depends on the ¹⁵N relaxation delay. Actually, a slightly higher degree of aliphatic protons saturation is observed at long T periods resulting in a reduced amount of H^N magnetization at the start of the scan, versus that observed at short T periods. This is consistent with an erroneous increase of measured 15 N R₁ rates with this scheme.

Concluding remarks

Systematic errors in ¹⁵N R_1 measurements arise from unwanted saturation of water and aliphatic protons by the train of 180° proton pulses used to eliminate cross-correlations effects. Errors coming from water saturation show a high degree of variability, being more important for residues in solvent-exposed regions of the protein, whose amide protons show fast exchange with water.

We have shown that saturation of aliphatic protons is another important source of systematic errors on $^{15}N R_1$ rates measurements in non-deuterated proteins. Errors show less variability along the amino acid sequence, although they are probably sensitive to local variations in the relaxation rates of relevant protons.

Maintaining the transverse magnetization of aliphatic protons in a dephased state (saturated state) during the variable ¹⁵N relaxation period T of the XY-pulse scheme by using the cosine modulated IBURP-2 scheme ensures that the amount of H^N polarization at the start of the pulse sequence does not depend on the length of the variable relaxation delay T.

In conclusion, the optimized pulse scheme shown in Fig. 1 (XY-sequence) using cosine-modulated IBURP-2 pulses spaced 40 ms minimize systematic errors on $^{15}N R_1$ rates measurements in non-deuterated proteins. This sequence efficiently controls both water and aliphatic protons saturation while cancelling CC effects during the variable ^{15}N relaxation period. The optimized pulse program

applied to protonated GB3 gives identical $^{15}N R_1$ values to those obtained with a deuterated sample.

Acknowledgments The clones used for the expression of GB3 were derived from those kindly provided by Dr. Donghan Lee (Max Planck Institute for Biophysical Chemistry). We thank Dr. Jesús García for providing us with the GB3 protein sample, and Dr. Oscar Millet for useful discussions and critical reading of the manuscript. Spectrometers at the NMR Facility of the Scientific and Technological Centers of the University of Barcelona were used in this work. This study was supported by MINECO-FEDER (grants BIO2013-40716-R and BIO2013-45793-R) and Generalitat de Catalunya (XRB and Grant 2014-SGR-1251).

Open Access This article is distributed under the terms of the Creative Commons Attribution 4.0 International License (http:// creativecommons.org/licenses/by/4.0/), which permits unrestricted use, distribution, and reproduction in any medium, provided you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons license, and indicate if changes were made.

References

- Boyd J, Hommel U, Campbell ID (1990) Influence of crosscorrelation between dipolar and anisotropic chemical shift relaxation mechanisms upon longitudinal relaxation rates of ¹⁵N in macromolecules. Chem Phys Lett 175:477–482
- Chen K, Tjandra N (2011) Water proton spin saturation affects measured protein backbone ¹⁵N spin relaxation rates. J Magn Reson 213:151–157
- Chill JH, Louis JM, Baber JL, Bax A (2006) Measurement of ¹⁵N relaxation in the detergent-solubilized tetrameric KcsA potassium channel. J Biomol NMR 36:123–136
- Farrow NA, Muhandiram R, Singer AU, Pascal SM, Kay CM, Gish G, Shoelson SE, Pawson T, Forman-Kay JD, Kay LE (1994) Backbone dynamics of a free and a phosphopeptide-compexed Src homology 2 domain studied by ¹⁵N NMR relaxation. Biochemistry 33:5984–6003
- Ferrage F, Cowburn D, Ghose R (2009) Accurate sampling of high-frequency motions in proteins by steady-state ¹⁵N-{¹H} nuclear Overhauser effect measurements in the presence of cross-correlated relaxation. J Am Chem Soc 131:6048–6049
- Ferrage F, Reichel A, Battacharya S, Cowburn D, Ghose R (2010) On the measurement of ¹⁵N-{¹H} nuclear Overhauser effects. 2. Effects of the saturation scheme and water signal suppression. J Magn Reson 207:294–303
- Geen H, Freeman R (1991) Band-selective radiofrequency pulses. J Magn Reson 93(1):93–141
- Gong Q, Ishima R (2007) ¹⁵N-{¹H} NOE experiment at high magnetic field strengths. J Biomol NMR 37:147–157
- Grzesiek S, Bax A (1993) The importance of not saturating H₂O in protein NMR. Application to sensitivity enhancement and NOE measurements. J Am Chem Soc 115:12593–12594
- Ishima R (2014) A probe to monitor performance of ¹⁵N longitudinal relaxation experiments for proteins in solution. J Biomol NMR 58:113–122
- Jurt S, Zerbe O (2012) A study on the influence of fast amide exchange on the accuracy of 15N relaxation rate constants. J Biomol NMR 54:389–400
- Kay LE, Nicholson LK, Delaglio F, Bax A, Torchia DA (1992a) Pulse sequences for removal of the effects of cross correlation between dipolar and chemical-shift anisotropy relaxation mechanisms on

the measurement of heteronuclear T_1 and T_2 values in proteins. J Magn Reson 97:359–375

- Kay LE, Keifer P, Saarinen T (1992b) Pure absorption gradient enhanced heteronuclear single quantum correlation spectroscopy with improved sensitivity. J Am Chem Soc 114:10663–10665
- Kovacs H, Moskau D, Spraul M (2005) Cryogenically cooled probes. A leap in NMR technology. Prog Nucl Magn Reson Spectrosc 46:131–155
- Krishnan VV, Murali N (2013) Radiation damping in modern experiments: progress and challenges. Prog Nucl Magn Reson Spectrosc 68:41–57
- Kumar A, Grace RCR, Madhu PK (2000) Cross-correlations in NMR. Prog Nucl Magn Reson Spectrosc 37:191–319
- Lakomek NA, Ying J, Bax A (2012) Measurement of ¹⁵N relaxation rates in perdeuterates proteins by TROSY-based methods. J Biomol NMR 53:209–221
- Mori S, Abeygunawardana C, Johnson MO, Vanzijl PCM (1995) Improved sensitivity of HSQC spectra of exchanging protons at short interscan delays using a new fast HSQC (FHSQC) detection scheme that avoids water saturation. J Magn Reson B 108:94–98
- Morin S (2011) A practical guide to protein dynamics from ¹⁵N spin relaxation in solution. Prog Nucl Magn Reson Spectrosc 59:245–262
- Palmer AG, Skelton NJ, Chazin WJ, Wright PE, Rance M (1992) Suppression of the effects of cross-correlation between dipolar

and anisotropic chemical-shift relaxation mechanisms in the measurement of spin-spin relaxation rates. Mol Phys 75:688-711

- Pervushin K, Vögeli B, Eletsky A (2002) Longitudinal ¹H relaxation optimized in TROSY NMR spectroscopy. J Am Chem Soc 124:12898–12902
- Piotto M, Saudek V, Sklenar V (1992) Gradient-tailored excitation for single-quantum NMR spectroscopy of aqueous solutions. J Biomol NMR 2:661–665
- Schanda P (2009) Fast-pulsing longitudinal relaxation optimized techniques: enriching the toolbox of fast biomolecular NMR spectroscopy. Prog Nucl Magn Reson Spectrosc 55:238–265
- Shaka AJ, Keeler J (1987) Broadband spin decoupling in isotropic liquids. Prog Nucl Magn Reson Spectrosc 19:47–129
- Shishmarev D, Otting G (2011) Radiation damping on cryoprobes. J Magn Reson 213:76–81
- Smallcombe SH (1993) Solvent suppression with symmetricallyshifted pulses. J Am Chem Soc 115:4776–4785
- Stonehouse J, Shaw GL, Keeler J, Laue ED (1994) Minimizing sensitivity losses in gradient-selected 15N–1H HSQC spectra of proteins. J Magn Reson A 107:178–184
- Torchia DA (2011) A practical guide to protein dynamics from 15N spin relaxation in solution. J Magn Reson 212:1-10
- Zhu G, Xia Y, Nicholson LK, Sze KH (2000) Protein dynamics measurements by TROSY-based NMR experiments. J Magn Reson 143:423–426