## 3ES03 S<sub>1</sub> excited-state dynamics of OMpCA and its hydrogen-bonded complexes (Hiroshima Univ.<sup>1</sup> IMS<sup>2</sup>) <u>Y. Miyazaki<sup>1</sup></u>, Y. Inokuchi<sup>1</sup>, T. Ebata<sup>1</sup> and M. Ehara<sup>2</sup>

The photocycle of photoactive yellow protein (PYP) has drawn attention to understand the negative phototaxis behavior of *Halorhodospira halophila*. Upon photoexcitation at 446 nm ( $\lambda_{max}$ ), p-coumaric acid embedded in PYP initiates the photocycle by trans→cis isomerization<sup>1</sup>. Although various experimental and computational studies have been published, a nonradiative relaxation process that leads trans→cis isomerization is still a matter of discussion<sup>2-5</sup>. Recently, the S<sub>1</sub> lifetime of a jet-cooled Oxyester Methyl p-Coumaric Acid (OMpCA) and its hydrogen-bonded complexes was measured to elucidate the role of the hydrogen bond by employing picosecond time-resolved pump-probe spectroscopy<sup>5</sup>. It unveiled the following characteristics: 1) the lifetime of OMpCA at the S<sub>1</sub> band origin is 9 ps while that of OMpCA-H<sub>2</sub>O is 930 ps, and 2) OMpCA-H<sub>2</sub>O exhibits excess energy dependence of nonradiative decay that is best fitted with a biexponential decay curve of fast and slow components. In the present work, we adapted the kinetic argument that was used to describe the excited-state dynamics of trans-stilbene<sup>6</sup>, and estimated an energy barrier height on the S<sub>1</sub> potential surface by the RRK calculation.

A gas mixture of OMpCA and solvent diluted with Ne carrier gas was expanded into vacuum through the pulsed nozzle, and skimmed by a skimmer to generate a supersonic molecular beam. A picosecond tunable UV laser, obtained by second harmonic generation of the optical parametric generation/optical parametric amplifier (OPG/OPA) system of Nd:YAG laser, crossed with a molecular beam. Electronic spectrum was acquired by mass-resolved resonance 2-photon ionization (R2PI) method via  $S_1$  state. Pump-probe scheme was used to record the decay curve as a function of the time delay between pump and probe lasers. The  $S_1$ -lifetime was determined by the convolution fitting with a laser width of 12 ps.

The R2PI excitation spectra of OMpCA and its hydrogen-bonded complexes in the  $31000 - 33100 \text{ cm}^{-1}$ region are shown in Fig 1. Blue and red lines are the spectra recorded by nanosecond laser and picosecond laser, respectively. It is clearly seen that a spectral red-shift of hydrogen-bonded complexes increases i.e. 633 cm<sup>-1</sup>, 921 cm<sup>-1</sup>, 1396 cm<sup>-1</sup> as solvent

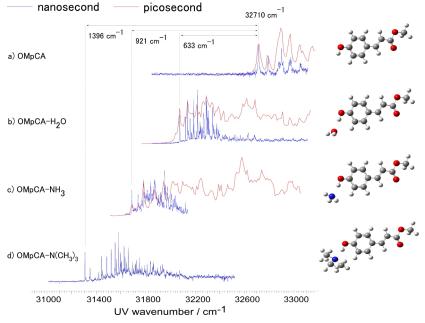


Fig 1. R2PI spectra of (a) OMpCA and (b)-(d) complexes

basicity increases from  $H_2O$ ,  $NH_3$  to  $N(CH_3)_3$ , respectively. The electronic excited-state is more stabilized than the ground state due to a slightly extended  $\pi$ -conjugation.

As mentioned, the excess energy dependence of the nonradiative decay is observed for  $OMpCA-H_2O$ . There exists potential energy barrier as a result of forming the hydrogen bond at the phenolic OH group. Fig 2. displays the time evolution excited at 310 cm<sup>-1</sup> above the S<sub>1</sub>-S<sub>0</sub> origin. The curve is fitted by a double exponential curve with a fast component (80 ps) and a slow component (220 ps). We applied a kinetic model involving IVR and nonradiative decay to describe the temporal behavior of the complex.

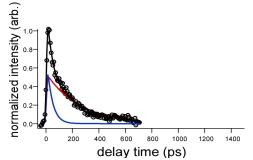


Fig 2. The observed decay curve of  $OMpCA-H_2O$  above 310 cm<sup>-1</sup> from the  $S_1-S_0$  band origin.

Fig 3 shows the schematic picture of  $S_1$ -state dynamics for OMpCA complexes, where  $\lambda$  is the wavelength,  $\sigma_{pump}$  is the absorption cross section,  $\sigma_{probe}$  is the ionization cross section,  $\beta$  is the dimentionless parameter that indicates the relative strength of the absorption from the  $S_1$  state. In short, an excess energy from the  $S_1$  band origin is first redistributed into the  $S_1$  vibrational manifold by IVR. A nonradiative process, e.g. isomerization and/or internal conversion, proceeds when an excess energy overcomes the barrier. Thus, a relative value of the two components could be observed in accord with an excess energy. In fact, the plot of decay rate constant as a function of excess energy for OMpCA-H<sub>2</sub>O in Fig 4. shows the validity of this kinetic argument, indicating that the fast component is IVR and the slow component is a nonradiative process. The extrapolated line is based on the RRK fitting with the barrier height of 205 cm<sup>-1</sup>. A detail of the analysis and dynamics associated with potential energy barrier for other OMpCA-complexes will be presented.

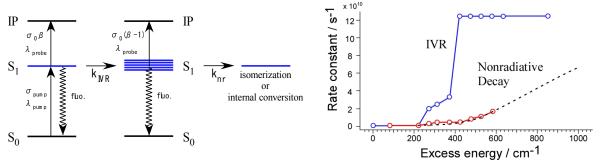


Fig 3. Schematic picture of the two-step kinetic model of the excited OMpCA complex

Fig 4. Plot of the observed decay rate constant of  $S_1$  state vs excess energy for OMpCA-H<sub>2</sub>O

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