

Boron isotope analysis of silicate glass with very-low boron concentrations by secondary-ion mass spectrometry

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1 Abstract

2 Here we present an improved method for the determination of the boron isotopic composition of volcanic
3 glasses with boron concentrations of as low as $0.4 - 2.5 \mu\text{g/g}$, as is typical for mid-ocean ridge basalt
4 glasses. The analyses are completed by secondary-ion mass spectrometry using a Cameca 1280 large-radius
5 ion microprobe. Transmission and stability of the instrument and analytical protocol were optimised, which
6 led to an improvement of precision and reduction of surface contamination and analysis time compared to
7 earlier studies. Accuracy, reproducibility ($0.4 - 2.3 \text{‰}$, 2RSD), internal precision ($2\text{RSE} = 2.5 - 4.0 \text{‰}$ for
8 a single spot with $[\text{B}] = 1 \mu\text{g/g}$), matrix effects ($\ll 0.5 \text{‰}$ among komatiitic, dacitic and rhyolitic glass),
9 machine drift (no internal drift; long-term drift: $\sim 0.1 \text{‰/h}$), contamination ($\sim 3 - 8 \text{ ng/g}$) and machine
10 background (0.093 s^{-1}) were quantified and their influence on samples with low B concentrations was de-
11 termined.

12 The newly developed setup is capable of determining the B isotopic composition of basaltic glass with
13 $1 \mu\text{g/g}$ B with a precision and accuracy of $\pm 1.5 \text{‰}$ (2RSE) by completing 4 – 5 consecutive spot analyses
14 with a spatial resolution of $30 \times 30 \mu\text{m}^2$. Samples with slightly higher concentrations ($\geq 2.5 \mu\text{g/g}$) can be
15 analysed with a precision of better than $\pm 2 \text{‰}$ (internal 2RSE) with a single spot analysis, which takes
16 32 min.

17 Keywords: boron isotopes; MORB; low concentration; SIMS; ionprobe

18 Introduction

19 Boron is a moderately volatile, lithophile non-metal with a low atomic mass and two stable isotopes and a
20 $^{11}\text{B}/^{10}\text{B}$ variation of several tens of per mil in Earth's surface environments (Palmer & Swihart, 1996). The
21 strong enrichment of B in the crust and the significant difference in B isotopic compositions between con-
22 tinental crust, modern seawater and the depleted mantle make B a potentially powerful geochemical tracer
23 for the secular evolution of the ocean-crust-mantle system (e.g., Spivack & Edmond, 1987; Chaussidon &

24 Marty, 1995; Turner *et al.*, 2007). However, the very low abundance of B in mantle rocks and primitive
25 volcanic rocks confronts us with a major analytical hurdle. Boron isotope analyses of silicate materials at
26 the trace abundance level are highly challenging. The B isotopic compositions of the Earth's major reser-
27 voirs (i.e., continental crust, primitive and depleted mantle) are still poorly constrained, despite a several
28 decades-long history of research into B isotope geochemistry.

29 The pioneering work by Chaussidon and co-workers (Chaussidon & Jambon, 1994; Chaussidon & Marty,
30 1995; Chaussidon *et al.*, 1997) has demonstrated the capability of secondary-ion mass spectrometry (SIMS)
31 to determine the B isotopic composition of mantle-derived lavas, and exemplified the use of B isotopes as a
32 geochemical tool in these rocks. However, the small-radius mass spectrometers available in the 1980s and
33 '90s have been surpassed in terms of transmission and stability by the large-radius ion microprobes available
34 today, which provide a 3 to 5 times higher sensitivity for B isotope measurements in low-B samples (e.g.,
35 Chaussidon *et al.*, 2006; Gurenko & Kamenetsky, 2011; Shaw *et al.*, 2012, this study). Improvements
36 have also been made in controlling and quantifying surface contamination (Marschall & Ludwig, 2004).
37 Furthermore, the availability of new international reference materials (Gonfiantini *et al.*, 2003; Jochum
38 *et al.*, 2006) enables a fresh assessment of the capabilities and limits of in-situ B isotope analysis by SIMS.
39 The method presented in this paper provides the means for routine analysis of B isotope ratios in natural
40 volcanic glasses at concentration levels from depleted MORB to highly enriched rhyolites.

41 **Analytical setup**

42 Boron isotope ratios were determined by SIMS using a Cameca ims1280 ion microprobe at the North-
43 Eastern National Ionmicroprobe Facility (NENIMF) at the Woods Hole Oceanographic Institution between
44 October 2012 and September 2013. The parameters for B isotope analyses were: 25 – 40 nA, 22 keV $^{16}\text{O}^-$
45 primary ion beam; 10 kV secondary acceleration voltage; ± 40 eV energy window without offset; secondary
46 ion detection by a single electron multiplier (ETP, SGE Analytical Science) in counting mode (electronically
47 set $\tau = 28$ ns). The primary beam was tuned on an Al metal grid using a $200 \times 200 \mu\text{m}^2$ raster. The

48 energy distribution of the secondary beam was scanned periodically on each sample mount using the $^{28}\text{Si}^{2+}$
49 signal, and the energy window was centered on the maximum energy. On high-B samples (e.g., reference
50 glass B6; see below) the energy window was centered using the $^{11}\text{B}^+$ signal.

51 A $70 \times 50 \mu\text{m}$ raster was applied during presputtering to remove the gold layer and reduce surface con-
52 tamination. Presputtering lasted for 200 – 300s. Prior to each measurement, mass calibration of $^{10}\text{B}^+$ and
53 $^{11}\text{B}^+$ was performed, and fine tuning of the secondary column deflectors, stigmators, and lenses was done
54 manually in order to maximize the signal of $^{28}\text{Si}^{2+}$ on the electron multiplier, thereby maintaining maxi-
55 mum transmission throughout the session. Automatic beam centering using the $^{11}\text{B}^+$ signal was found to
56 be unreliable for samples with very low B concentrations. Gaussian illumination was used, which gave high
57 transmission of the rectangular rastered area through the square-shaped field aperture. A $4000 \times 4000 \mu\text{m}^2$
58 field aperture was used to exclude the edge of the secondary beam, again reducing the influence of sur-
59 face contamination (Marschall & Ludwig, 2004). The field of view of the sample through this aperture is
60 $30 \times 30 \mu\text{m}^2$, which defines the spatial resolution of the analysis and the ion-optical magnification ($130 \times$).
61 The largest contrast aperture ($400 \mu\text{m}$) was used to maximise transmission. We did not systematically test
62 other combinations of setups. No evidence for charging was found. Surface contamination is discussed
63 below in more detail. The instrument was operated in the X-Y mode, which produced a more ideal peak
64 shape and high transmission. No strong increase in transmission was observed when switching to circular
65 mode. However, this could be explored more rigorously in future to possibly further increase transmission.

66 The analyses were performed using a rastered area of $30 \times 30 \mu\text{m}^2$ in the centre of a larger pre-sputtered
67 area ($70 \times 50 \mu\text{m}^2$). It was found that performing the analyses in raster mode produces ablation pits with
68 a flatter bottom and less depth progression compared to spot analyses, as well as more constant secondary
69 ion signals over the duration of an analysis. The sputter rate on the basaltic glasses is $1.6 \pm 0.5 \text{ nm/s}$.
70 Dimensions of the sputtered areas were determined using an automatic-focus Zeiss AxioImager reflected
71 light microscope with the AxioVision imaging software including the topography package. Sputter pits
72 created by a regular 40-cycle analysis recommended for routine work are $\sim 3.1 \mu\text{m}$ deep, resulting in an

73 aspect ratio of the sputtered pits of ~ 0.1 . The amount of total sputtered material during an analysis is
74 $\sim 5.1 \times 10^3 \mu\text{m}^3$.

75 No dynamic transfer setting was used in order to avoid the inclusion of the contamination-afflicted edge
76 of the beam in the analyses. Forty cycles were measured per analysis, switching between masses $^{10}\text{B}^+$ and
77 $^{11}\text{B}^+$. The integration times per cycle were 20s for ^{10}B and 10s for ^{11}B . The mass resolution $m/\Delta m$ was
78 set to ~ 1250 at 10% intensity ratio, which is sufficient to remove possible molecular interferences (e.g.,
79 Ludwig *et al.*, 2011). The total time required for one analysis is 32 min. This analysis time could be reduced
80 to 22 min per spot, if a multi-collection system was used, as that would remove the settling time (a total of
81 5 s per cycle) and the integration time for the more abundant isotope (10 s per cycle).

82 The raw count rates were corrected for both the multiplier deadtime and for the slow changes in sec-
83 ondary ion intensity over the course of a single measurement (i.e., time interpolation). The latter was done
84 by averaging counts of ^{10}B from two subsequent cycles and calculate the $^{11}\text{B}/^{10}\text{B}$ ratio by dividing the
85 count rate of the intermediate ^{11}B measurement by that average. The extent of this intensity correction is
86 typically 0.1 – 1.0‰. Each analysis consisting of 40 analytical cycles, therefore, produced 39 isotope ra-
87 tios. The count rates were also corrected for the machine background of 0.093s^{-1} , which was determined
88 by analysing mass 9.7 on different glass samples over the course of one week for a total integration time
89 of 127 min. This correction is very small for natural volcanic glass, i.e. $< 0.25\text{‰}$ for MORB glasses with
90 $[\text{B}] = 0.4 \mu\text{g/g}$ and $< 0.10\text{‰}$ for MORB glasses with $[\text{B}] = 2.5 \mu\text{g/g}$. Finally, the 39 isotope ratios of each
91 analysis were filtered for statistical outliers $> 3\sigma$. These outliers are rare and are typically related to in-
92 tensity jumps in the primary beam, and any such cycles were rejected. On average, one in four analysis
93 has one cycle that is rejected. Isotope ratios calculated from averaging a number of ratios collected over
94 the course of a single analyses are positively biased, but this bias decreases with increasing counts N of the
95 minor isotope per measurement cycle, $\text{bias} = 1/N + 2/N^2$ (Ogliore *et al.*, 2011). The largest bias expected
96 for the MORB glasses with the lowest B concentration ($[\text{B}] = 0.4 \mu\text{g/g}$; ^{10}B : $N = 6500$ counts per cycle) is
97 $+0.15\text{‰}$. Results were not corrected for this bias.

98 Boron isotopes are reported in the delta notation relative to SRM 951 (U.S. National Institute of Stan-
99 dards and Technology, NIST; Catanzaro *et al.*, 1970): $\delta^{11}\text{B} = [({}^{11}\text{B}/{}^{10}\text{B}_{\text{sample}})/({}^{11}\text{B}/{}^{10}\text{B}_{\text{SRM 951}}) - 1] \cdot 1000$
100 (${}^{11}\text{B}/{}^{10}\text{B}_{\text{SRM 951}} = 4.04362$). Analytical uncertainties are discussed below.

101 Boron concentrations were determined using the Cameca ims1280 with the same setup for raster sizes
102 and aperture, mass resolution, 40 nA ${}^{16}\text{O}^-$ primary beam, and size of energy window with zero offset.
103 Pre-sputtering lasted for 2 min. Ten analytical cycles were analysed including masses ${}^{11}\text{B}^+$ and ${}^{28}\text{Si}^{2+}$.
104 Reference glass GOR132-G (Table 1; Jochum *et al.*, 2000) was used to determine ${}^{11}\text{B}^+ / {}^{28}\text{Si}^{2+}$ relative ion
105 yields¹, which were 490 ± 6 (2SD; $n = 4$) and 454 ± 20 (2SD; $n = 6$) for two different sessions. Reference
106 glass B6 (Gonfiantini *et al.*, 2003; Tonarini *et al.*, 2003) was also analysed in the latter session and resulted
107 in a relative ion yield of 435 ± 1 (2SD; $n = 2$), which is indistinguishable from the result on GOR132-G in
108 that session.

109 Reference materials and samples

110 Four different reference materials were used in this study, as given in Table 1. This includes natural glass
111 and glasses prepared from natural volcanic rocks. Glass compositions range from komatiitic to rhyolitic,
112 with boron concentrations of approximately 12 – 200 $\mu\text{g}/\text{g}$ (Table 1). GOR128-G and GOR132-G belong
113 to the set of MPI-DING glasses and were prepared from Gorgona (Columbia) komatiite samples GOR128
114 and GOR132, respectively (Jochum *et al.*, 2000). Boron concentrations and ${}^{11}\text{B}/{}^{10}\text{B}$ ratios in these glasses
115 are much higher than expected for mafic or ultramafic volcanic rocks and point to assimilation of seawater-
116 altered materials into the Gorgona magmas (Jochum *et al.*, 2006). The high abundance of B ($\sim 20 \mu\text{g}/\text{g}$;
117 Table 1) makes these glasses well suited as SIMS reference materials.

118 Reference material StHs6/80-G is another MPI-DING glass that was prepared by melting and quenching
119 a sample of dacitic ash from Mount St Helens (Washington, USA). It has a moderately high abundance of B
120 and a B isotopic composition that is similar to that of mid-ocean ridge basalts (Jochum *et al.*, 2000, 2006).

¹RIY = ${}^{11}\text{B}^+ / {}^{28}\text{Si}^{2+} \cdot [\text{Si}]/[\text{B}] \cdot (\bar{M}_{\text{B}}/I_{11\text{B}})/(\bar{M}_{\text{Si}}/I_{28\text{Si}})$; \bar{M} = mean atomic mass, I = isotopic abundance.

121 Reference material B6 is a natural obsidian glass from Lipari island (Aeolian archipelago, Italy) that
122 was characterised in a B isotope interlaboratory comparison study (Gonfiantini *et al.*, 2003; Tonarini *et al.*,
123 2003). The material is distributed by the International Atomic Energy Agency (IAEA). Its boron abundance
124 is the highest of all materials investigated here ($\sim 200 \mu\text{g/g}$; Table 1), which allows low-uncertainty SIMS
125 analysis of this reference material. The interlaboratory study resulted in $\delta^{11}\text{B}$ values determined by P-TIMS
126 of $-3.35 \pm 0.24 \text{‰}$, $-1.56 \pm 0.60 \text{‰}$, and $-0.45 \pm 0.60 \text{‰}$ (mean $\delta^{11}\text{B} = -1.79 \pm 2.93 \text{‰}$).

127 NIST glasses SRM 610 or SRM 612 were not used to determine the instrumental mass fractionation
128 (IMF), as they have been demonstrated to produce IMF values that are different from all silicate glasses
129 with natural compositions, and that that difference depends on machine type and setup (Rosner *et al.*, 2008;
130 Gurenko & Kamenetsky, 2011).

131 All glass samples analysed in this study were large fragments (0.5 – 2 mm diameter) that were mounted
132 in epoxy (Buehler Epothin) or pressed into indium mounts within 7 mm radius of the centre of the 12.7 mm
133 radius sample holder. Indium mounts generally produce a lower background during SIMS volatile analyses,
134 whereas it may be easier to produce a sample surface that is flat across the entire mount if the samples
135 are embedded in epoxy. It is neither expected nor observed that the epoxy should influence the machine
136 background for boron; however, the mounting materials were not systematically evaluated against each
137 other in the course of this study. The epoxy was annealed in a pressure chamber at 400 kPa to suppress the
138 formation of bubbles that would potentially disturb the surface and the stability of the vacuum. Polishing
139 was completed using a Buehler MiniMet 1000 polishing machine (1 μm diamond paste), which was set to
140 produce a flat and even surface throughout the epoxy and glass samples. Alumina polish (0.3 μm) was used
141 for the final polish. It was found that polishing new epoxy mounts within the first few weeks leads to a
142 relatively strong relief around the edges of the samples, while waiting several months before the final stage
143 of polishing leads to further hardening of the epoxy, and a very flat surface can be produced. Flat surfaces
144 without relief, pits or cracks are essential for accurate isotope measurements by SIMS (see discussion in
145 Kita *et al.*, 2009, for O isotopes).

146 No systematic investigation of surface flatness on the measured B isotope ratio was conducted here.
147 However, we analysed one glass sample that was in contact with two large exposed epoxy gas bubbles
148 (~ 0.5 mm diameter; see supplementary figure) at various distances to the gas bubbles to investigate the
149 possible influence of surface discontinuity. No systematic difference was found between analyses close to
150 a gas bubble ($< 200 \mu\text{m}$) compared to analyses at large distances from any gas bubble ($\sim 600 - 1000 \mu\text{m}$;
151 supplementary figure). This demonstrates that the B isotope measurements in our analytical setup are more
152 robust against surface discontinuities than O isotope analyses.

153 All analysis were completed at a distance of at least $100 \mu\text{m}$ from the edge of the samples. This includes
154 the reference materials and MORB glasses. Prior to gold coating, the grain mounts were cleaned using
155 96% ethanol followed by an ultrasound bath using distilled water from a Millipore ultrapure water system
156 ($18\text{M}\Omega$). The Millipore system typically reduces the B concentration in the water to $< 0.5 \mu\text{g/L}$ (from
157 $\sim 40 \mu\text{g/L}$ in tap water; Darbouret & Kano, 2000). Samples were always cleaned and coated immediately
158 before introducing them into the airlock of the mass spectrometer to reduce the possible deposition of
159 contamination on the sample surfaces.

160 Contamination

161 In general, thin sections and polished grain mounts are used for the analysis of B concentrations and B
162 isotopic compositions by SIMS. The surfaces of these samples are prone to the collection of contamination
163 with boron-bearing and other volatile or water-soluble compounds during sample preparation and during
164 storage (Shaw *et al.*, 1988; Chaussidon *et al.*, 1997; Hervig, 2002; Marschall & Ludwig, 2004). Surface
165 contamination has also been demonstrated to exist for Li, Na, K and Fe (Müller *et al.*, 2003; Marschall
166 & Ludwig, 2004), but is particularly recognised as a significant obstacle in the determination of accurate
167 B abundance and B isotope data in low-B concentration samples (Shaw *et al.*, 1988; Chaussidon *et al.*,
168 1997; Marschall & Ludwig, 2004). Levels of contamination may depend on the techniques of sample
169 preparation and cleaning procedures used, as well as analytical procedures during SIMS analysis. Additional

170 contamination may be produced during analysis by implantation of boron through the primary beam from
171 low-purity materials in the duoplasmatron (Chaussidon *et al.*, 1997). However, this is avoided by primary
172 beam mass filtering in modern SIMS instruments. Significant memory effects, as is potentially the case with
173 laser-ablation-ICP-MS (see discussion in le Roux *et al.*, 2004), have not been observed in SIMS, except after
174 sputtering boron-rich salts (Chaussidon *et al.*, 1997). Estimates for the equivalency of surface contamination
175 plus instrumental background contamination range from $< 10 \text{ ng/g}$ (Kent & Rossman, 2002; Marschall &
176 Ludwig, 2004) to $10 - 50 \text{ ng/g}$ (Chaussidon *et al.*, 1997) to $\geq 2 \text{ }\mu\text{g/g}$ (Domanik *et al.*, 1993; Marschall &
177 Ludwig, 2004).

178 A large range of minerals and glasses that are of geological interest show B abundances below $10 \text{ }\mu\text{g/g}$,
179 and many show abundances of approximately $1 \text{ }\mu\text{g/g}$ (e.g., Ottolini *et al.*, 2004; Marschall *et al.*, 2006a). For
180 example, mid-ocean ridge basalt (MORB) glasses commonly contain $0.4 - 2.5 \text{ }\mu\text{g/g}$ (Spivack & Edmond,
181 1987; Ryan & Langmuir, 1993; Chaussidon & Jambon, 1994; Leeman & Sisson, 2002; le Roux *et al.*,
182 2004). Boron surface contamination is, therefore, potentially significant for these type of samples, and its
183 suppression is imperative in order to reduce bias in B abundance and isotope measurements. Moreover, it is
184 critical to quantify the amount and B isotopic composition of the remaining surface contamination, so that
185 an estimate can be made of the bias introduced by the remaining contamination.

186 Here, we use the silica glass Herasil-102 (Heraeus Quarzglas GmbH, Germany) to characterise B surface
187 contamination. Herasil-102 was recommended as an appropriate material to quantify B surface contamina-
188 tion, as it is an ultrapure, homogeneous material that is available in large quantities, and its B concentration
189 is $\leq 1.1 \text{ ng/g}$ (Marschall & Ludwig, 2004). Any ion signals of $^{10}\text{B}^+$ and $^{11}\text{B}^+$ detected by the multiplier
190 during a regular measurement on Herasil-102 can be attributed to the combined contributions of surface
191 contamination, machine background, memory effects and B from the glass itself. Importantly, the abun-
192 dance of B in Herasil-102 is very low, so that the contributions from the other sources become dominant
193 and can be evaluated. A fragment of Herasil-102 was mounted in epoxy and polished together with the B
194 isotope reference materials and a number of MORB glasses and analysed for apparent B concentration and

195 B isotopic composition along with these samples using the same analytical setups.

196 Boron isotope measurements on Herasil-102 resulted in ion yields of $\sim 25 \text{ s}^{-1}$ for ^{11}B . In MORB glasses
197 with B concentrations of $0.4 - 2.5 \mu\text{g/g}$, the same count rate amounts to a contamination contribution of
198 8 ng/g or $0.3 - 2.0\%$ to the total counted signal. The B isotopic composition of the surface contamination
199 using Herasil-102 can only be determined with large uncertainty, due to the very low count rates. Four
200 analyses were completed and resulted in $\delta^{11}\text{B}$ values ranging from $-58 \pm 71 \text{ ‰}$ to $+2 \pm 45 \text{ ‰}$, with a mean
201 of $-36 \pm 27 \text{ ‰}$ (2SE). These values are used to estimate the systematic error of B surface contamination on
202 the B isotope analyses of the samples using a simple mass balance approach:

$$\delta^{11}\text{B}_d = X_s \cdot \delta^{11}\text{B}_s + X_c \cdot \delta^{11}\text{B}_c \quad (1)$$

203 where $\delta^{11}\text{B}_d$ is the determined B isotope value (measured value corrected for intensity and instrumental
204 mass fractionation), $\delta^{11}\text{B}_s$ is the the true value of the sample, and $\delta^{11}\text{B}_c$ is the B isotope value of the
205 contamination. X_s and X_c are the proportions of B from the sample and the contamination that contribute to
206 the signal ($X_s + X_c = 1$).

207 The bias for the determined B isotope value introduced from contamination, i.e. the difference between
208 $\delta^{11}\text{B}_d$ and $\delta^{11}\text{B}_s$ is displayed in Fig. 1a, assuming $\delta^{11}\text{B}_s = -7 \text{ ‰}$ for the sample and a contamination
209 $\delta^{11}\text{B}_c = -36 \text{ ‰}$ as discussed above. Note that the absolute values plotted in Fig. 1 depend on the difference
210 between $\delta^{11}\text{B}_s$ and $\delta^{11}\text{B}_c$.

211 MORB glass samples measurements with a $0.3 - 2.1\%$ signal contribution from contamination would
212 have to be corrected by $+0.09$ to $+0.23 \text{ ‰}$ to retrieve the uncontaminated isotopic composition of the sample
213 (Fig. 1). However, this systematic error has a large uncertainty, due to the large uncertainty of the isotopic
214 composition of the contamination component (Fig. 1) and, thus, no contamination correction of the $\delta^{11}\text{B}$
215 values was done during this study.

216 Furthermore, the systematic error introduced by contamination can be estimated as a function of B con-

217 centration of the analysed sample for the specific setup used in this study using Equation 1. In this case, X_c
 218 is derived from the apparent B concentration of the contamination $[B]_c = 8 \text{ ng/g}$ and $[B]_s$, the true B con-
 219 centration of the sample: $X_c = [B]_c / ([B]_s + [B]_c)$. The model demonstrates that for our analytical setup the
 220 potential systematic error introduced by contamination is only 0.23 ‰ for a sample with a B concentration
 221 of $1 \mu\text{g/g}$, or 0.58 ‰ ($\pm 2\text{SE} = 0.04 - 1.11 \text{ ‰}$) for $0.4 \mu\text{g/g}$ samples (Table 2; Fig. 1b). Larger biases would
 222 be expected for samples with much lower B contents or with exotic B isotopic compositions that diverge
 223 more from that estimated for the contamination.

224 Precision

225 The internal precision of a single spot analysis is defined here as the standard error of the mean ($\text{SE} =$
 226 SD/\sqrt{n}) of the $n = 39$ intensity-corrected $^{11}\text{B}/^{10}\text{B}$ ratios and are given here as two times the relative stan-
 227 dard error². In practice, the precision will depend on counting statistics, on the homogeneity of the analysed
 228 material on the scale of the measurement (micrometres in the case of SIMS), on surface contamination, on
 229 the stability of the mass spectrometer and the detector system, and, especially in case of a single-collector
 230 measurement, on the stability of the primary beam. At low concentrations, the standard error is dominated
 231 by counting statistics and can be predicted from Poisson statistics in a contamination-free measurement
 232 (e.g., Fitzsimons *et al.*, 2000):

$$\text{RSE}(\text{‰}) = 1000 \cdot \sqrt{\frac{1}{\sum N_{10}} + \frac{1}{\sum N_{11}}}, \quad (2)$$

233 where $\sum N_{10}$ and $\sum N_{11}$ are the total counts of ^{10}B and ^{11}B , respectively, over the course of the analysis.
 234 For example, in a sample containing $1 \mu\text{g/g}$ B, the count rates of ^{10}B and ^{11}B may be 3200 s^{-1} and 800 s^{-1} ,
 235 respectively, translating to 6.4×10^5 and 1.28×10^6 counts, respectively, for this analysis. The predicted
 236 precision is $2\text{RSE} = 3.06 \text{ ‰}$. Measurements that are significantly afflicted with surface contamination show

²Note that all errors and uncertainties discussed in this paper are 2 standard deviations or 2 standard errors, while previous publications in many cases report 1 RSD and 1 RSE uncertainties.

237 decreasing count rates over the course of the analysis and have a precision that is many times worse than
238 statistically predicted (Marschall & Ludwig, 2004).

239 The sensitivity of the instrument, i.e., the number of counts per second registered during an analysis of a
240 sample with a given concentration, depends on the setup of the instrument (mass resolution, energy filtering,
241 etc.), on the primary beam current (or more specifically on the product of beam density and analysed area)
242 and on the ion yield. The last is relatively constant among basaltic glasses, but may vary significantly (i.e., a
243 factor of 2) between basaltic and rhyolitic glasses. In our study, the sensitivity varied between approximately
244 50 and 80 counts/s/nA/($\mu\text{g/g}$) for ^{11}B on the GOR glasses. The resulting predicted precisions for primary
245 beam currents of 25 and 40 nA are displayed in Fig. 2. This figure also shows the precision of analyses
246 of the reference materials and a range of MORB glasses. Observed and predicted precisions agree well at
247 low concentrations, while counting statistics is not the limiting factor at concentrations $> 20 \mu\text{g/g}$ where
248 the internal precision reaches a practical lower limit of $\sim 0.5\%$ (2RSE). This practical limit is likely set
249 by the stability of the primary beam and of the ablation and ionisation conditions. These factors are much
250 less effective when both ions are counted simultaneously. The practical limit of precision may therefore be
251 lower on a multi-collector instrument. The majority of MORB glasses ($[\text{B}] = 0.4 - 1.2 \mu\text{g/g}$) were analysed
252 with a precision of between 2.5 and 4% (2RSE; Fig. 2).

253 The sensitivity of the instrument may also be expressed in terms of the useful ion yield (Hervig *et al.*,
254 2006), which depicts the number of counted ions of a particular isotope relative to the number of sput-
255 tered atoms of that isotope. For example, basaltic glass (assumed density = 2700 kg/m^3) with $[\text{B}] = 1 \mu\text{g/g}$
256 ($^{11}\text{B}/^{10}\text{B} = 4$) contains $1.2 \times 10^5 \text{ atoms}/\mu\text{m}^3$ of ^{11}B . Our analytical setup (sputter rate 1.65 nm/s ; sput-
257 tered area $30 \times 30 \mu\text{m}^2$) consumes $1.48 \mu\text{m}^3/\text{s}$ ($= 4 \text{ pg/s}$) of basalt glass. The amount of sputtered ^{11}B is,
258 therefore, $1.78 \times 10^5 \text{ atoms/s}$. The count rate on such a sample is typically 3200 s^{-1} . Hence, our useful ion
259 yield is 1.8%. This is approximately one order of magnitude higher than the useful ion yields reported for
260 boron for small-radius ion microprobes (Cameca 3f and 6f), which range from 0.14 to 0.31% (Hervig *et al.*,
261 2006). The total amount of boron consumed during a 32 min analysis of a $[\text{B}] = 1 \mu\text{g/g}$ glass including the

262 larger-area pre-sputtering is 14 fg from 14 ng of glass.

263 The precision of the measurement improves with the amount of boron consumed, which increases with
264 integration time and the number of analytical cycles for a given sample (Fig. 3). The analytical setup
265 used in this study permits MORB glasses with B concentrations of 1 $\mu\text{g/g}$ to be analysed with an observed
266 precision of $\sim 2.5 - 4.0\%$ (2RSE), which agrees with the Poisson statistical prediction of 3.1%. This
267 internal precision compares very favourably with previous studies using various SIMS and laser-ablation
268 inductively-coupled-plasma (LA-ICP-MS) instruments (Table 2).

269 Analysis of homogenous basaltic glass with 1 $\mu\text{g/g}$ B require ~ 5 analyses to achieve a 2RSE external
270 precision of 1.5% and ~ 10 analyses for 1.0% (2RSE). This would require 2.7 and 5.4 hours, respectively,
271 not including analyses of the reference materials. The practical results for multiple analyses on MORB
272 glasses demonstrate that the values are reproducible within the error given by internal precision, and produce
273 mean $\delta^{11}\text{B}$ values with 2RSE of 1 – 2% (Fig. 4).

274 **Analytical drift**

275 **Drift within individual measurements (internal drift)**

276 The internal precision of an isotope ratio measurement is influenced by the total integration time, as dis-
277 cussed above. It would, therefore, seem logical to extend analyses on low-concentration samples to very
278 long counting times to improve precision. However, such a practice reduces the number of samples that
279 could be analysed in a given time, which has practical and financial drawbacks. More importantly, though,
280 it also introduces additional potential sources of analytical bias. Analyses that sputter the sample for one
281 hour or longer create relatively deep pits potentially leading to a shift in instrumental fractionation over the
282 course of the analysis (Hervig *et al.*, 1992). Previous SIMS studies performing oxygen isotope and trace
283 element analyses recommend not to exceed aspect ratios of the sputter pit (depth divided by diameter) of 0.1
284 (Schuhmacher *et al.*, 1994) or 0.25 (Valley & Graham, 1991; Hervig *et al.*, 1992). The aspect ratio of the

285 pits produced by the 40-cycle analyses in our setup after 32 min is ~ 0.1 . Longer sputter times with more
286 cycles would increase the aspect ratio. For example, the pit created by a 400-cycle analysis is $\sim 25 \mu\text{m}$ deep
287 with an aspect ratio of ~ 0.8 . The shape of the pit is also asymmetrical due to the incidence angle of the
288 primary beam of $\sim 68^\circ$. The influence of sputter time or pit aspect ratio on the measured B isotope ratio was
289 not investigated systematically in this study, but it was found that a 400-cycle analysis on glass reference
290 material GOR132-G showed IMF values within 1‰ of the value determined from the first 40 analytical cy-
291 cles only for the first ~ 200 cycles. At that stage, the aspect ratio of the pit is ≥ 0.4 produced after a sputter
292 time of two hours. At an aspect ratio of ≥ 0.5 the IMF value is $\sim 2‰$ lower than the value determined from
293 the first 40 cycles. The B count rate decreased to approximately half of the initial value after 400 cycles.
294 These findings also demonstrate that determination of IMF values from analyses of references materials and
295 analyses of unknown samples need to be completed with the same analytical setup and the same number of
296 analytical cycles.

297 The setup used in this study uses relatively short total sputter durations of 32 min for a single spot analysis
298 (Table 2). Ion count rates and the $^{11}\text{B}/^{10}\text{B}$ ratio are monitored over the course of each analysis (Fig. 5).
299 Potential drift in the $^{11}\text{B}/^{10}\text{B}$ ratio is evaluated through the slope of the linear regression of the isotope
300 ratio over the course of the 39 intensity-corrected measurement cycles (Fig. 5). These slopes are mostly not
301 significantly different from zero within 2 standard errors (Fig. 6). In addition, two important criteria must
302 be fulfilled for a set of analyses that do not show systematic analytical drift during a single analysis: (1) the
303 slopes for a set of analyses should follow a random distribution that is symmetrical around zero, and (2) the
304 magnitude of the slopes should decrease with increasing count rate and increase with increasing standard
305 deviation, i.e., with worse analytical precision. Both conditions are fulfilled for the set of 221 analyses
306 completed over the course of this study, demonstrating that no systematic drift occurs under the current
307 setup and analytical protocol (Fig. 6).

308 **Long-term drift**

309 Changes in instrumental mass fractionation in SIMS over the duration of several hours or days are observed
310 for O isotope analyses, and require frequent analysis of reference materials and a time-related correction
311 of the measured ratios for the samples (e.g., Valley *et al.*, 1998; Kita *et al.*, 2009). Long-term drift of IMF
312 has been observed to be negligible or within the precision of the analysed reference materials for B isotope
313 analyses in a number of previous studies (Chaussidon *et al.*, 1997; Marschall *et al.*, 2006b). In this study,
314 drift of the IMF was observed over the course of one day, with IMF values slowly changing by typically
315 + or -0.1‰ /h from early morning to late night. Note that this drift may be positive or negative (Fig. 7),
316 and that the total drift over an entire day never exceeded 1.8‰ . Observation and quantification of drift are
317 practically limited by the reproducibility of the analyses of the reference materials. Instrumental drift over
318 the course of one day can be corrected for through the repeated analysis of reference materials throughout
319 the analytical session. Alternatively, the drift-uncorrected IMF values for the entire session may be used,
320 which would increase the uncertainty on IMF from typically 1.5‰ to $\sim 2.4\text{‰}$ (Fig. 7).

321 **Matrix effects**

322 The matrix effect describes the dependency of the instrumental mass fractionation on the chemical compo-
323 sition or crystallographic structure of the analysed materials. Chemical matrix effects have been reported
324 for a number of trace element and isotope systems and require close chemical matching between standards
325 and samples or a good description of the matrix effect as a function of composition (Shimizu & Hart, 1982;
326 Eiler *et al.*, 1997; Page *et al.*, 2010). Matrix effects for B isotopes have been reported to be very small or
327 negligible for a large range of minerals and glasses with the exception of the NIST glass series SRM 61x
328 and one sample of a Li-rich tourmaline (Chaussidon & Albarède, 1992; Chaussidon *et al.*, 1997; Nakano
329 & Nakamura, 2001; Rosner *et al.*, 2008; Gurenko & Kamenetsky, 2011). Small matrix effects were also
330 reported between amphibole and rhyolitic glass B6 ($2.8 \pm 2.0\text{‰}$; 1σ), but were insignificant for the pair

331 phengite and B6 ($1.3 \pm 2.8 \text{‰}$; 1σ ; Pabst *et al.*, 2012). In this study, the possible effect of the chemical com-
332 position of the investigated material on instrumental fractionation of the two isotopes of B was investigated
333 by comparing IMF values determined for the four different reference materials. These four different glasses
334 vary in composition from komatiitic to rhyolitic with silica contents from ~ 45 to ~ 75 wt% (Table 1).

335 The recommended $\delta^{11}\text{B}$ values for these reference glasses are relatively well established and they were
336 repeatedly analysed in a number of laboratories by various methods. However, it should be noted that
337 there are still existing discrepancies between values reported by various laboratories that are larger than the
338 reported precisions (e.g., Gonfiantini *et al.*, 2003). Hence, although many TIMS and ICP-MS laboratories
339 routinely produce B isotope data on silicate minerals and glasses with reported analytical uncertainties of
340 $2\sigma < 0.5 \text{‰}$, reproducibilities including full sample dissolution and chemical B separation are more typical
341 in the range $0.5 - 1 \text{‰}$. Interlaboratory comparison reveals consistencies on the order of only $1.5 - 3.0 \text{‰}$
342 (Gonfiantini *et al.*, 2003; Tiepolo *et al.*, 2006; Hou *et al.*, 2010; Wei *et al.*, 2013). Consequently, there
343 are relatively large uncertainties on the reported $\delta^{11}\text{B}$ values, which limits the evaluation of possible SIMS
344 matrix effects. Note that if not taken into account, accumulation of all uncertainties on the reference values
345 can lead to a significant overestimation of SIMS matrix effects.

346 The results from this study confirm previous results that no significant matrix effects can be detected for
347 the large compositional range of reference glasses (Fig. 9, 10). The differences in IMF among the various
348 materials are well within the repeatability of the B isotope analyses on these samples; no correlation between
349 IMF and chemical composition was detected (Fig. 9b). The difference in IMF between komatiitic glass
350 GOR132-G and dacitic glass StHs6/80-G is $0.04 \pm 1.64 \text{‰}$ (2SD). The weighted mean of four sessions for
351 the differences in IMF between komatiitic glass GOR132-G and rhyolitic glass B6 is $0.18 \pm 1.36 \text{‰}$ (2SD).
352 The difference in IMF between the two komatiitic glasses GOR132-G and GOR128-G is $0.25 \pm 1.87 \text{‰}$
353 (2SD). All of these values are insignificant since the uncertainties are much larger than the values.

354 Note that these errors do not include uncertainties on the reference values, which, if propagated, would
355 make the observed differences even less significant. In conclusion, in the analytical setup used in this

356 study the potential matrix effect for B isotope analyses for the range of natural glasses with compositions
357 from komatiitic to rhyolitic is much smaller than the reproducibility of the reference glasses. A better
358 quantification of a possible small matrix effect in the low sub-permil range would require more accurate and
359 precise $\delta^{11}\text{B}$ values for the reference materials.

360 **Reproducibility and accuracy**

361 The primary B isotope standard, boric acid SRM 951 (NIST), cannot be used as a reference material in
362 SIMS to correct instrumental mass fractionation of measurements on silicate glass. Hence, accuracy of the
363 B isotope measurements can only be evaluated relative to secondary reference materials. Uncertainties on
364 the reported $\delta^{11}\text{B}$ values on those secondary reference materials, therefore, translate into a larger uncertainty
365 on the accuracy of the $\delta^{11}\text{B}$ values reported from the SIMS laboratory, if more than one reference material is
366 used for calibration. The use of an average IMF value determined from a set of reference materials reduces
367 the dependency on a single reference material and potentially reduces the inaccuracy of the reported $\delta^{11}\text{B}$
368 value.

369 The instrumental mass fractionation over a single session in this study based on all analysed reference ma-
370 terials showed values to be reproducible within ± 1.3 to ± 2.4 ‰ (2RSD) and ± 0.4 to ± 1.1 ‰ (2RSE) for all
371 sessions (Fig. 8). The reproducibility of individual reference materials are in the same range (Fig. 9a). These
372 values are also in the same range as the ones reported in previous studies using various SIMS instruments
373 or LA-ICP-MS (Table 2). For MPI-DING reference glasses GOR128-G, GOR132-G, and StHs6/80-G, the
374 reproducibility of 0.4 – 1.6 ‰ (2RSD; Fig. 9) is in a similar range as the internal precision on these materials
375 (0.5 – 1.9 ‰; both median and mean are 1.0 ‰).

376 Conclusions

377 The enhanced transmission and stability of the Cameca 1280 setup used at NENIMF (the Woods Hole SIMS
378 facility) in the course of this study led to an improvement of precision and reduced instrument drift, surface
379 contamination and analysis time compared to earlier studies. Accuracy, reproducibility, precision, matrix
380 effects, contamination and machine background were quantified and their influence on samples with low
381 B concentrations was determined. Single analyses were completed with a spatial resolution of $30 \times 30 \mu\text{m}$
382 within 32 min.

383 The accuracy of the SIMS analyses for multiple analyses of a homogenous material is determined by the
384 reproducibility of all analyses within a given analytical session ($\sim 2 - 5$ days), and was ± 0.4 to $\pm 2.3 \text{‰}$
385 (2RSD). Precision of a single B isotope analysis of basaltic glass with $1 \mu\text{g/g}$ B is determined by Poisson
386 statistics and is 3.1‰ (2RSD). Analysis of homogenous basaltic glass with $1 \mu\text{g/g}$ B require ~ 5 analyses to
387 achieve a 2RSE precision of better than $\pm 1.5 \text{‰}$. At concentrations exceeding $\sim 20 \mu\text{g/g}$, internal precision
388 reaches a lower limit of 0.5‰ (2RSE).

389 Chemical matrix effects are too small to be quantified, i.e. no significant differences in instrumental
390 mass fractionation were observed (even on the $1 - \sigma$ level) for the range of investigated glass compositions
391 ranging from komatiitic to rhyolitic. The analysis are demonstrated to show no systematic internal drift.
392 Long-term drift over the course of a day is limited ($< 1.8 \text{‰}$ throughout one day) and can be quantified
393 and corrected for through multiple analyses of reference glasses. Surface contamination can contribute
394 bias. However, this was negligible ($< 0.1 \pm 0.2 \text{‰}$) in samples exceeding $2.5 \mu\text{g/g}$ boron. At the low end
395 of concentrations found in MORB glasses ($0.4 \mu\text{g/g}$) the bias introduced by surface contamination is still
396 small ($0.6 \pm 1.1 \text{‰}$) compared to precision at that concentration ($4 - 5 \text{‰}$; 2RSE).

397 In summary, the newly developed setup at NENIMF is capable of determining the B isotopic composition
398 of natural volcanic glasses, including basaltic glasses with very low B abundances such as depleted MORB
399 glass. Precision, accuracy and reproducibility of better than 1.5‰ (2RSE) is achieved, including possible

400 sources of error, such as surface contamination, drift and matrix effects. Samples with less than $\sim 4 \mu\text{g/g}$
401 require several analyses to achieve this precision, but analysis time is relatively short at ~ 30 min per spot.

402 The method presented here constitutes an improvement of analytical uncertainty by a factor of approxi-
403 mately two to four, while reducing the analysis time by a factor of three compared to previously reported ion
404 microprobe protocols. Analytical uncertainties are comparable to those of laser-ablation multiple-multiplier
405 inductively coupled plasma mass spectrometry, but with the analysed area reduced by two to three orders of
406 magnitude and the analysed sample volume reduced by three to four orders of magnitude. The method pre-
407 sented here, thus, not only provides the analytical capabilities to investigate MORB glasses at a geologically
408 meaningful level of uncertainty, but also to investigate spatially restricted samples, such as melt inclusions
409 and crystallite-rich glasses, which are not accessible by methods other than SIMS.

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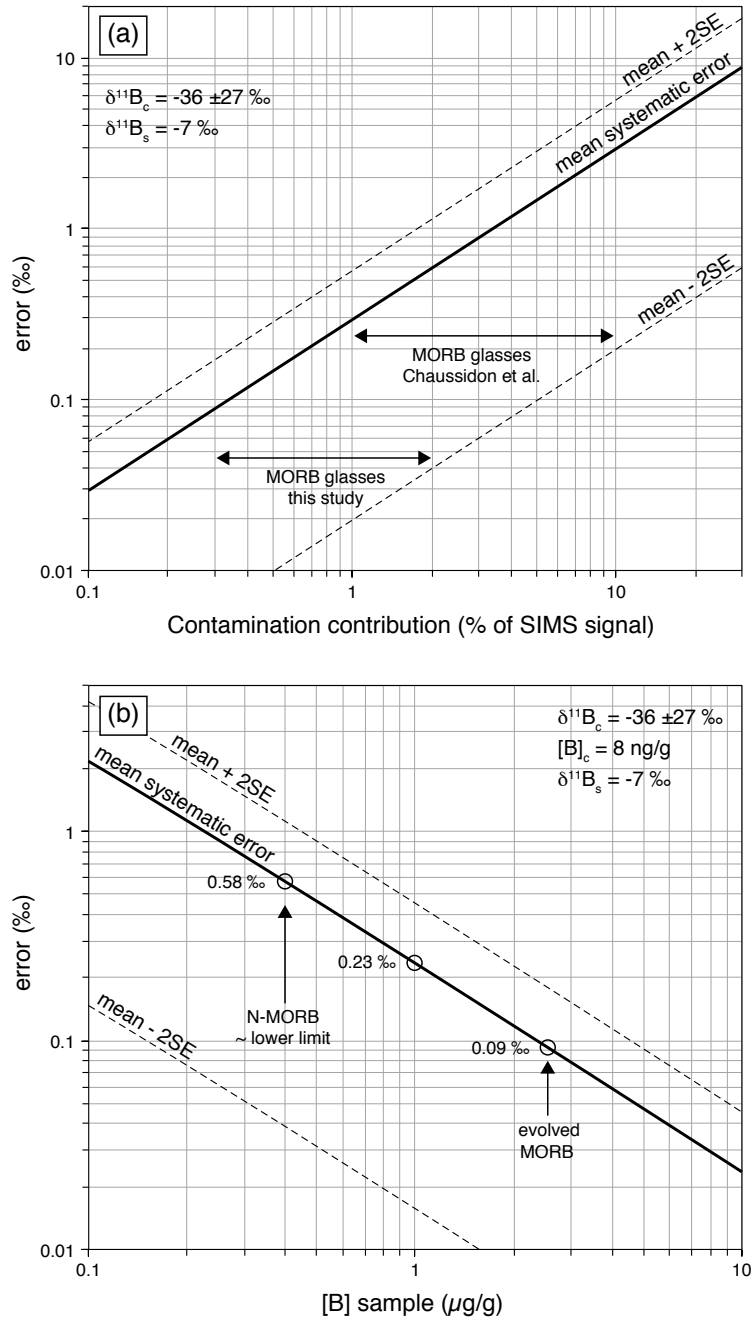


Fig. 1 (a) Systematic error introduced by surface and machine boron contamination as a function of the proportion of the contamination in the SIMS signal (sample $\delta^{11}\text{B}_s = -7.0 \text{‰}$; contamination $\delta^{11}\text{B}_c = -36 \pm 27 \text{‰}$). The continuous black line shows the systematic error, while the dashed lines show the uncertainty (2 standard error) on that systematic error. The contribution of surface contamination during MORB glass analysis in this study was determined to be 0.3 – 2.0%, while it was estimated at approximately five times higher in one previous study (Chaussidon *et al.*, 1997, no other studies have systematically quantified the contamination). **(b)** Systematic error introduced by boron contamination in this study as a function of sample boron concentration, assuming 8 ng/g contamination (B isotopic compositions as in (a)). MORB glass typically contains 0.4 – 2.5 $\mu\text{g/g}$ B. The systematic error expected from surface contamination at 1 $\mu\text{g/g}$ concentration in this study is +0.23‰ (2SE: +0.02 to +0.45‰). Note the large uncertainty on the error.

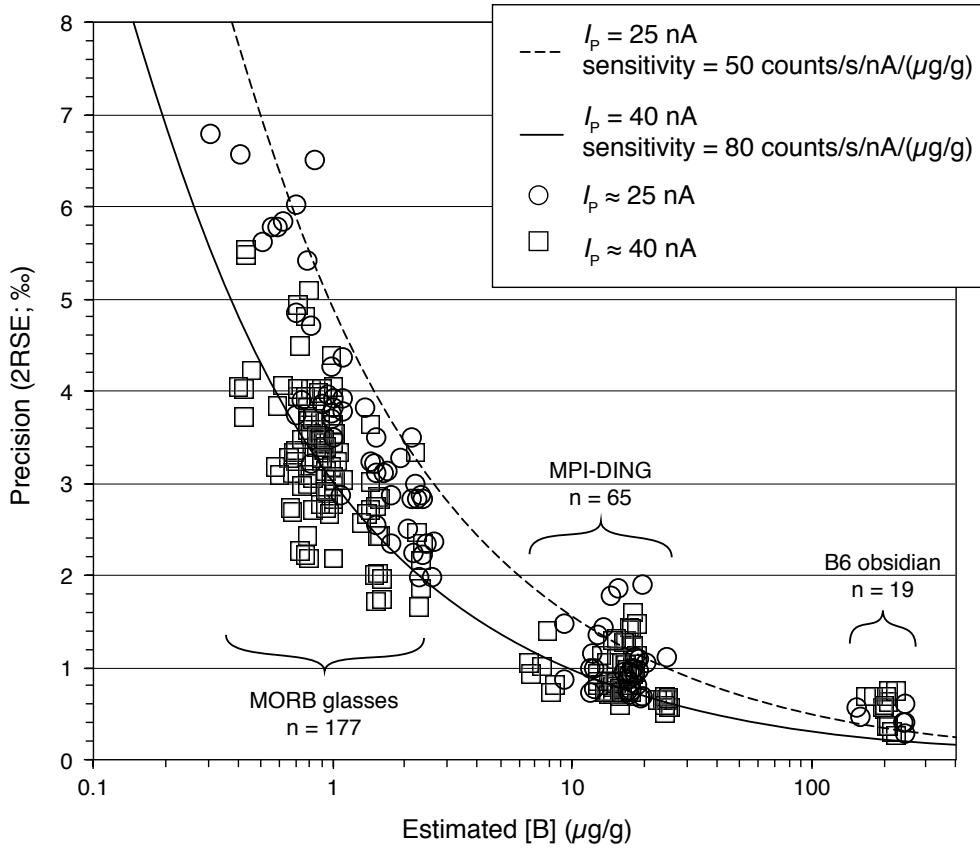


Fig. 2 Internal precision of individual B isotope analyses, expressed as two times the relative standard error, as a function of boron concentration (estimated from ^{11}B count rates) for reference materials and a range of MORB glasses. Two different primary beam currents were used, and sensitivity varied between sessions. Observed precision and precisions predicted from Poisson statistics (e.g., Fitzsimons *et al.*, 2000) agree well at low concentrations, while counting statistics is not the limiting factor at concentrations $> 20 \mu\text{g/g}$ where internal precision reaches a practical lower limit of $\sim 0.5\%$ (2RSE).

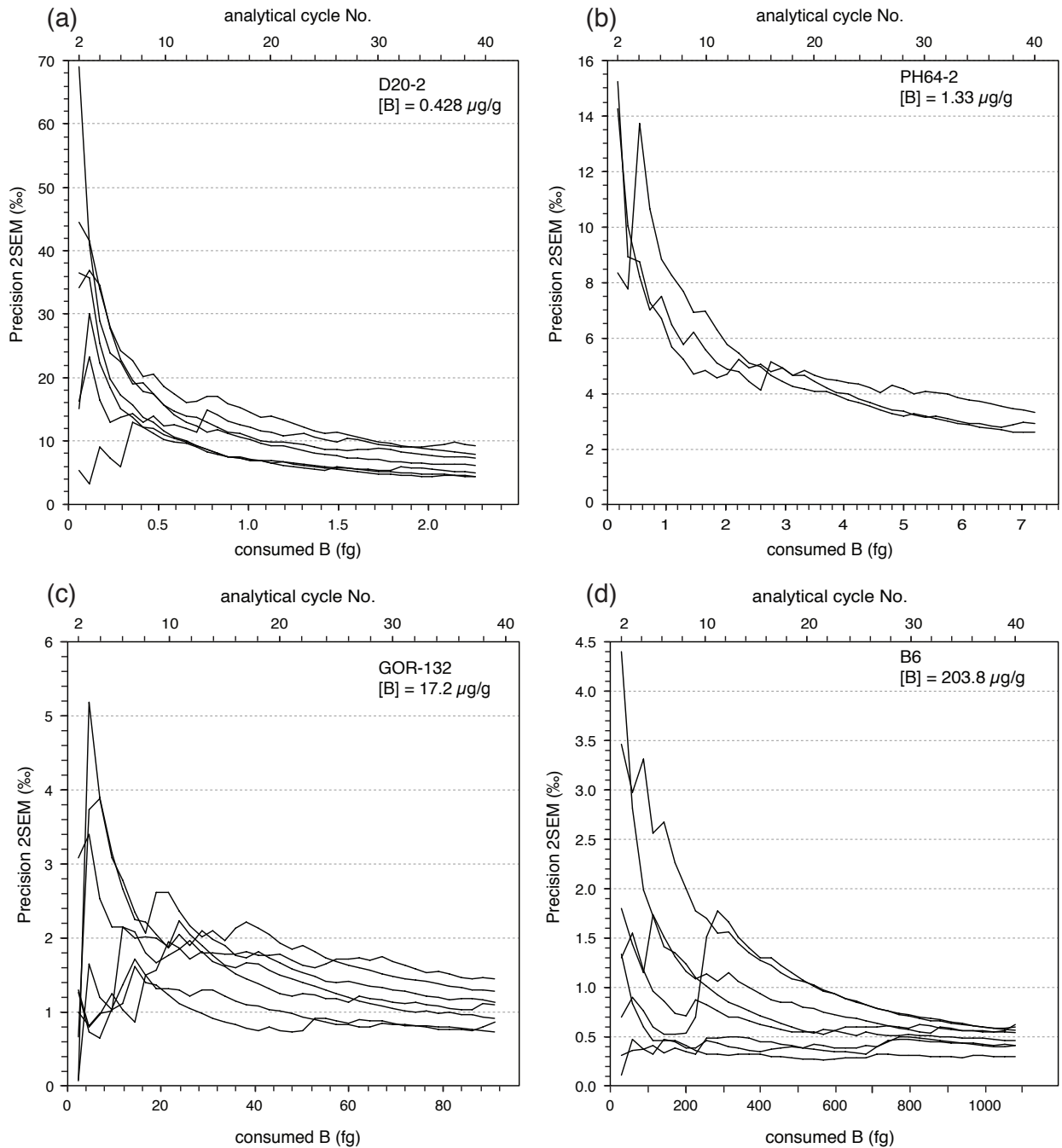


Fig. 3 Internal precision of individual B isotope analyses, expressed as two times the relative standard error, as a function of boron consumed during the analysis ($1 \text{ fg} = 10^{-15} \text{ g}$) for reference materials and two different MORB glasses. The amount of boron consumed during presputtering is not included in order to facilitate comparison among the samples. The number of analytical cycles are given on the secondary x-axis on top. Note the different scales on x and y axes among figures (a) to (d).

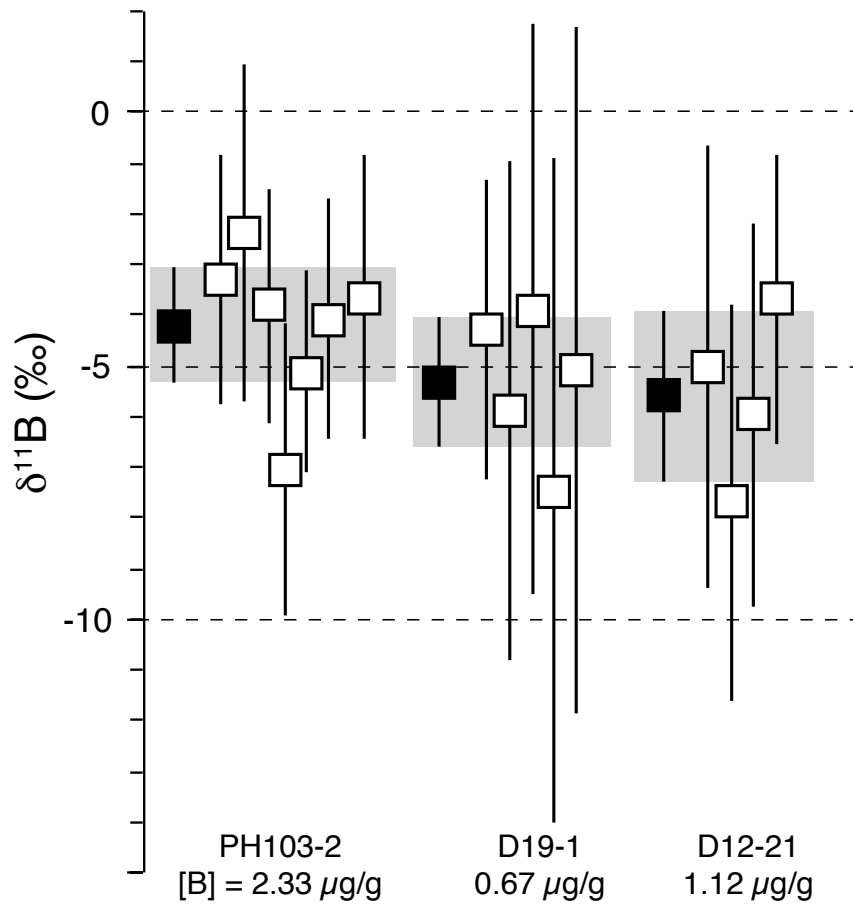


Fig. 4 Repeated measurements of B isotopic composition of three different MORB glass samples with B concentrations between 0.67 and 2.33 $\mu\text{g/g}$. White squares mark individual measurements with internal 2SD error bars; solid squares mark sample means with 2SE error bars also highlighted by grey fields.

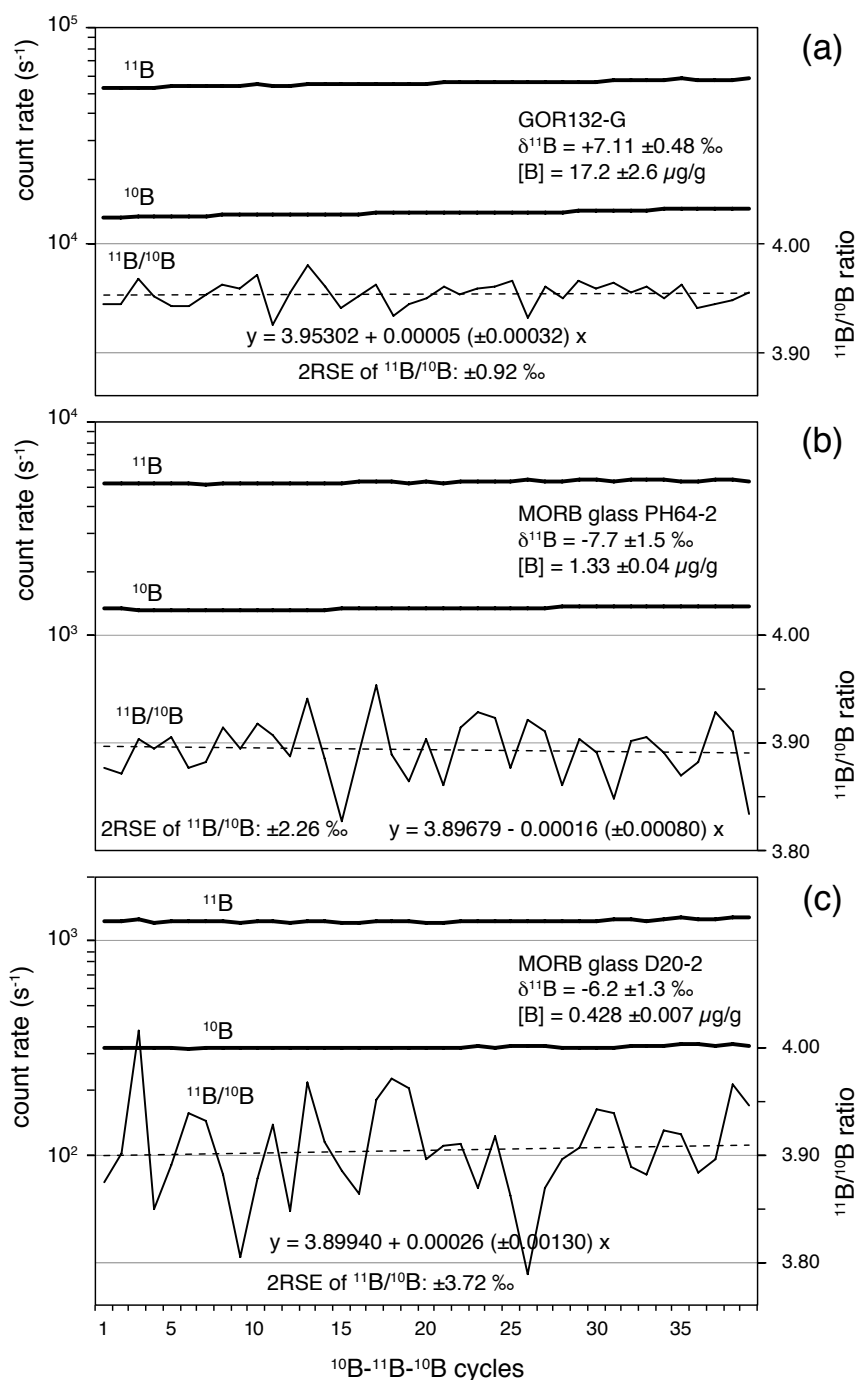


Fig. 5 Three examples of typical ^{11}B and ^{10}B signals and $^{11}\text{B}/^{10}\text{B}$ ratios over a 40-cycles analysis (intensity-corrected to 39 isotope ratios): **(a)** Reference material MPI-DING glass GOR132-G, **(b)** MORB glass PH64-2 with $[\text{B}] = 1.33 \mu\text{g/g}$, and **(c)** MORB glass D20-2 with $[\text{B}] = 0.428 \mu\text{g/g}$. The dashed lines and the linear equations in each diagram show the linear regression of the isotope ratios over the analyses, including 2SE on the slope of this regression. Note that the slopes of the regression lines are not significantly different from zero and are unsystematic with positive and negative values.

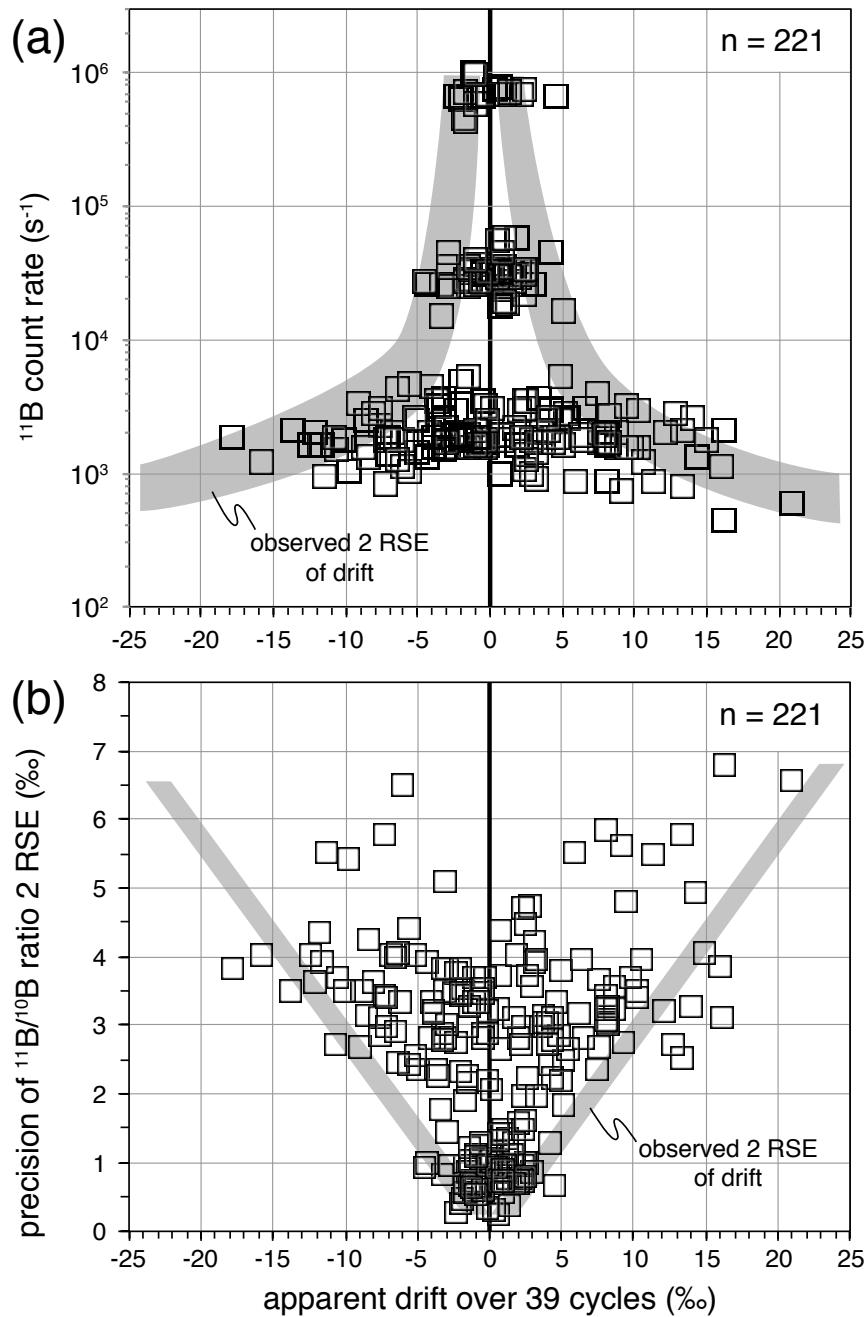


Fig. 6 Apparent drift during single spot analyses as calculated from the slope of the linear regression through the $^{11}\text{B}/^{10}\text{B}$ ratios over the 40-cycles analyses intensity-corrected to 39 isotope ratios (see Fig. 5). The observed drift is not systematic and shows a symmetric distribution between positive and negative slopes centered around 0, and the magnitudes of the slopes decrease with (a) ^{11}B count rate and (b) analytical precision. The grey fields mark the 2SE of the slope of the linear regression lines as observed from the same analyses. Only analyses that plot outside the grey envelopes show $> 2\text{SE}$ significant drift.

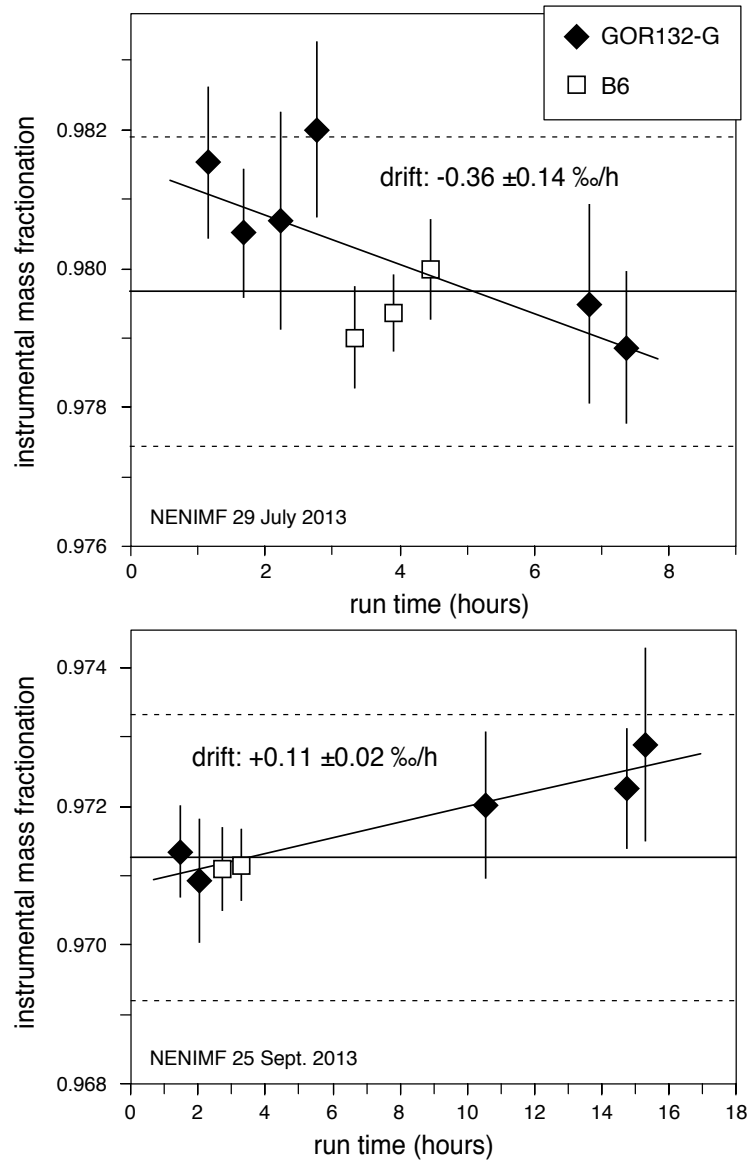


Fig. 7 Two examples of drift of instrumental mass fractionation over the course of an analytical day (run time given in hours). The IMF drift over one day was typically between -0.4 and $+0.3 \text{ ‰/h}$, and never exceeded 1.8 ‰ in total over the course of one day. Without any drift correction, all analysis of one session (up to 6 days) vary around one mean value with a variation of less than $\pm 2.4 \text{ ‰}$ (2SD) during an entire session, as indicated by the solid and dashed horizontal lines, respectively.

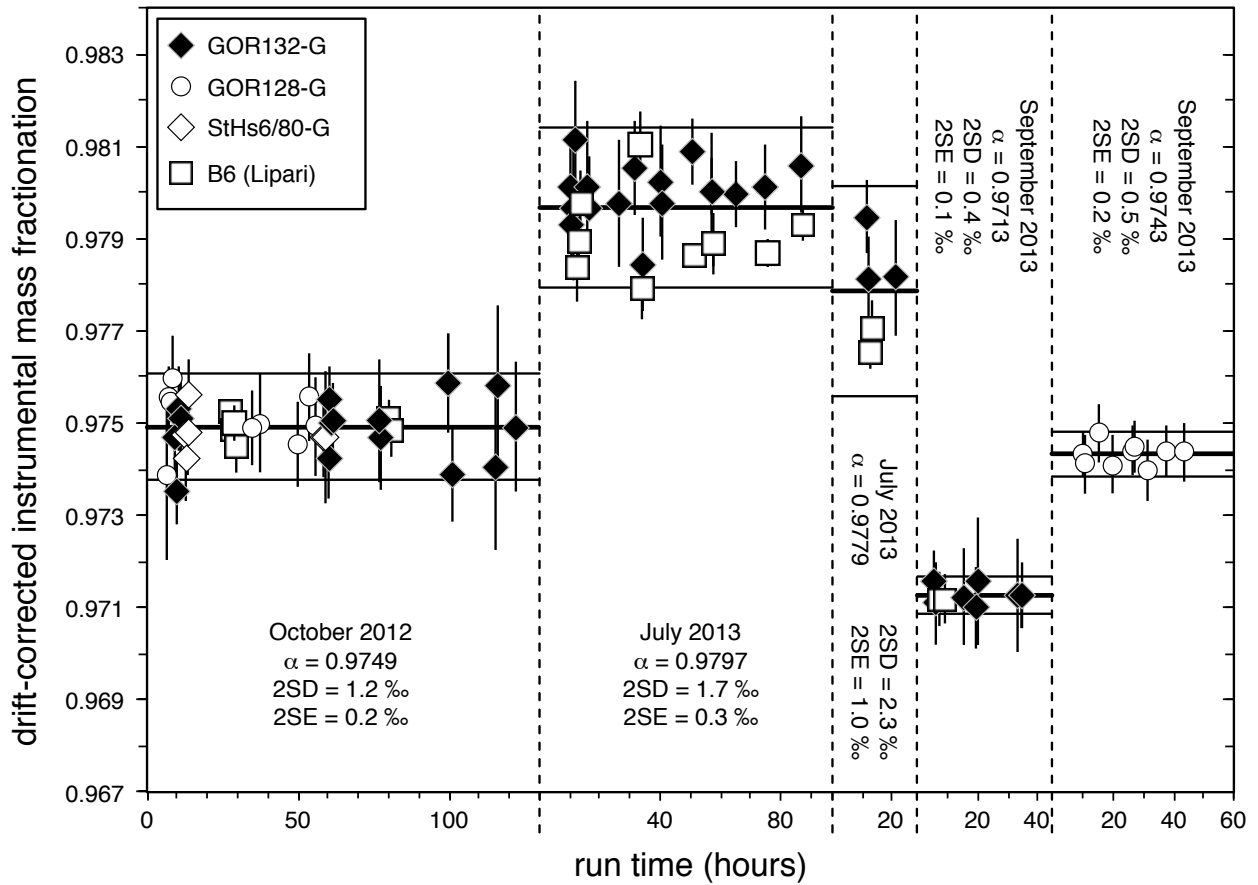


Fig. 8 Instrumental mass fractionation (corrected for intra-day instrumental drift) during five different analytical sessions over the course of one year, plotted against run time of each session. The longest session lasted one week. Reproducibility of instrumental mass fractionation was accurate within ± 0.4 to ± 2.3 ‰ (2SD) as demonstrated through repeated measurements of a number of reference materials.

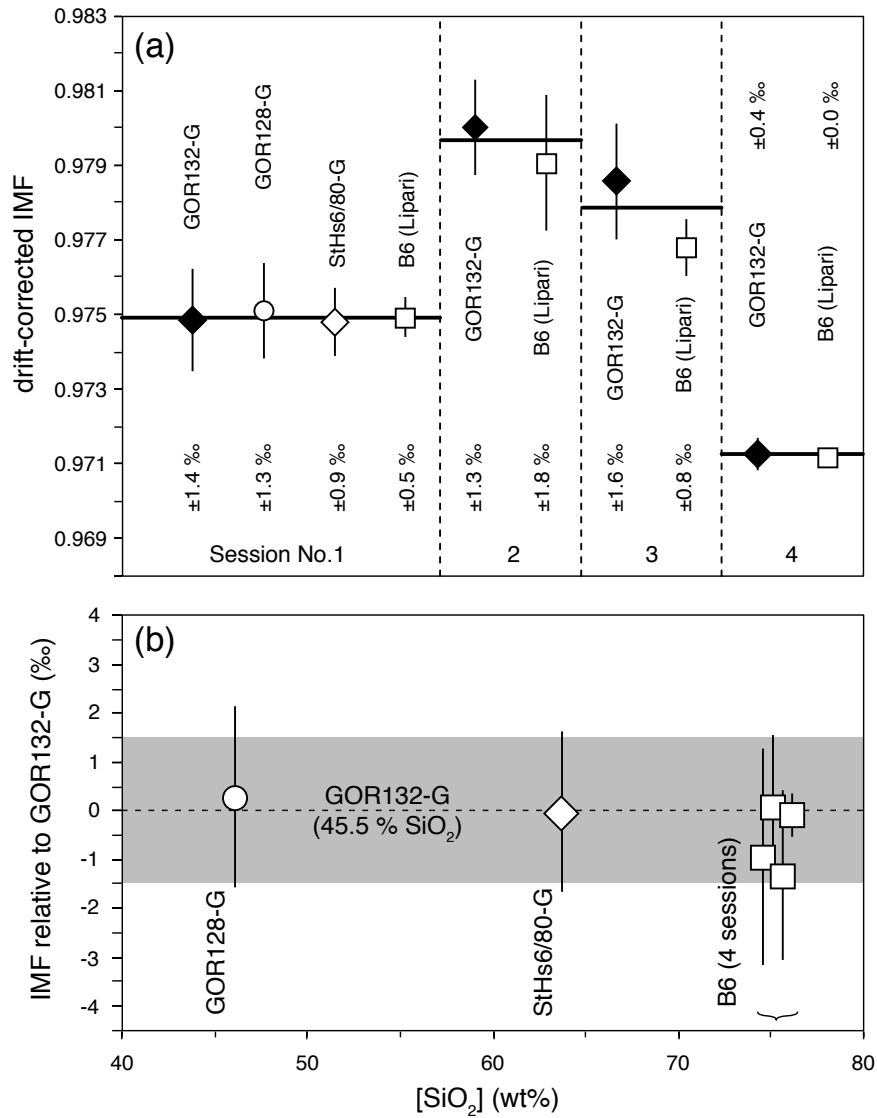


Fig. 9 (a) Drift-corrected instrumental mass fractionation (IMF) in four different analytical sessions for four different reference materials. IMF may vary between sessions, but no significant differences in IMF between different materials were observed. The values given depict reproducibilities of the materials throughout the session (2SD). (b) Drift-corrected IMF relative to reference material MPI-DING glass GOR132-G plotted versus silica content. No significant difference was observed in the IMF between reference materials, despite a large compositional range from komatiitic to rhyolitic composition. Hence, no compositional matrix effect can be detected for B isotope analyses in our SIMS lab for glasses of natural composition ranging from komatiite to rhyolite. Note that the propagated errors do not include uncertainties (precision or accuracy) on the reference values. The grey bar represents typical reproducibility (2SD) of GOR132-G.

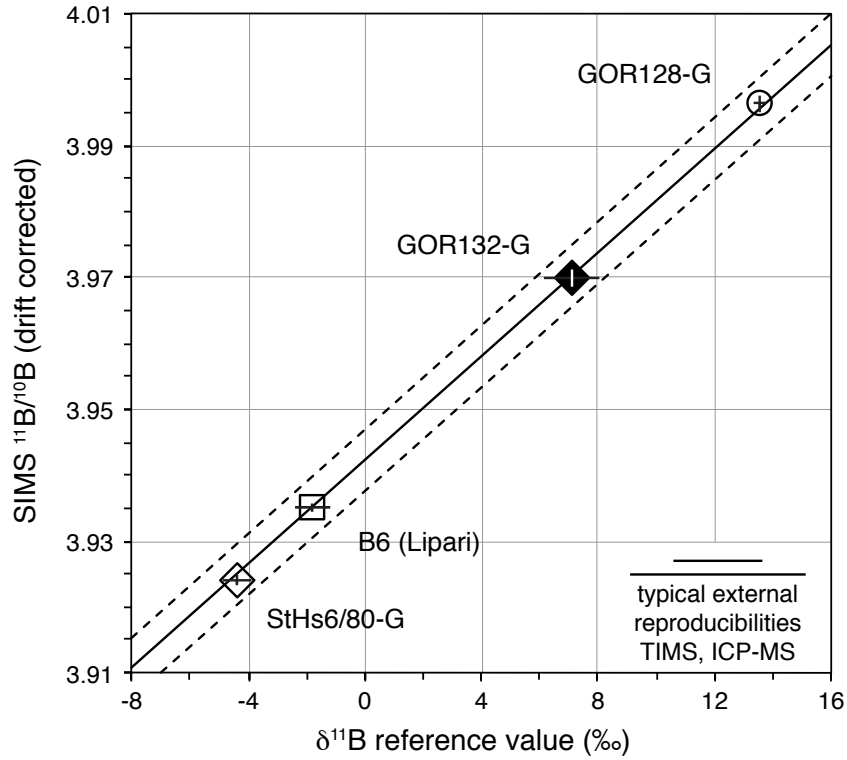


Fig. 10 Drift-corrected ¹¹B/¹⁰B ratio versus recommended δ¹¹B values for four different glass reference materials. The reference values have a typical precision of ±0.5permil. However, interlab comparison studies show that the accuracy of the TIMS and ICP-MS analyses are more typically in the range of ±1.5 – 3permil, as indicated by the bars in the lower right corner. The diagonal black line indicates constant instrumental mass fractionation ($\alpha = 0.9749 \pm 0.0013$). Note that the reference materials range from komatiitic to rhyolitic in composition, and no matrix effect is detectable among these materials. This plot also shows that the IMF is independent of the absolute ¹¹B/¹⁰B ratio.

Table 1 Reference materials used in this study

Name	Locality	Composition	[SiO ₂] (wt%)	[B] ($\mu\text{g/g}$)	2SD	$\delta^{11}\text{B}$ (‰)	2SD	References
MPI-DING glasses								
GOR128-G	Gorgona	komatiite	46.1	23.5	2.8	+13.55	0.21	[1,2,3]
GOR132-G	Gorgona	komatiite	45.5	17.2	2.6	+7.11	0.97	[1,2,3]
StHs6/80-G	Mt St. Helens	dacite	63.7	11.8	1.3	-4.48	0.29	[1,2,3,4]
IAEA reference glass								
B6	Lipari	rhyolite	75.2	203.8	8.9	-1.79	0.6	[4,5,6,7]

All listed boron isotope values are values determined by P-TIMS. Uncertainties represent 2SD on the reported results and do not include uncertainties from full repeats including sample dissolution or from interlaboratory differences, which are typically on the order of $\pm 1.5 - 3$ ‰. References are: [1] Jochum *et al.* (2006), [2] Rosner & Meixner (2004), [3] Tiepolo *et al.* (2006), [4] Rosner *et al.* (2008), [5] Gonfiantini *et al.* (2003), [6] Wei *et al.* (2013), [7] Hou *et al.* (2010).

Table 2 Uncertainties of B isotope analysis in this study and published studies

Study	Instrument	Detector(s)	Precision of single analysis at [B] = 1 µg/g	Reproducibility of reference materials per session (2SD)	Contamination and background (apparent concentr.)	Spatial resolution (µm)	Analysis time per spot (min)
Secondary-ion mass spectrometry							
this study	Cameca 1280	single multiplier	±3.1‰	0.4 – 2.3‰	3 – 8 ng/g	30 × 30	32
G&K11	Cameca 1280	single multiplier	±5.5‰	1.5 – 3.0‰	not determined	30 × 30	90
CH06	Cameca 1270	single multiplier	~ ±5‰	1.9 – 3.0‰	not determined	~ 20 (diameter)	105
CH97	Cameca 3f	single multiplier	±5.7‰	0.9 – 4.4‰	10 – 50 ng/g	100 (diameter)	90
H01	Cameca 3f	single multiplier	±12.9‰	~ 3.0‰	not determined	30 (diameter)	90
S01	Cameca 6f	single multiplier	±6.6‰	4.7‰	not determined	~ 20 (diameter)	40
Laser-Ablation-Multiple-Multiplier-ICP-MS							
leR04	VG Elemental + Nd-YAG laser	two multipliers	±3.6‰	0.8 – 1.0‰	not determined	1650 × 500	6

Listed studies are: G&K11 Gurenko & Kamenetsky (2011) (WHOI prior to 2010); CH06 Chaussidon *et al.* (2006) (CRPG, Nancy); CH97 Chaussidon & Jambon (1994), Chaussidon *et al.* (1997) (CRPG, Nancy); H01 Hoppe *et al.* (2001) (MPI, Mainz); S01 Sugiura *et al.* (2001) (Univ. Tokyo); leR04 le Roux *et al.* (2004) (DTM, Washington D.C.). Note that some of the values listed here are not listed in the original papers, but had to be reconstructed from the given values in the analytical sections and from error bars given in diagrams and converted to 2 standard deviation and 2 standard error, respectively, for this comparison. Values for single analysis precision are average values for the different methods and laboratories.