

# 1 Variations in oceanic plate bending along the Mariana trench

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## 9 Abstract

10 We quantify along-trench variations in plate flexural bending along the Mariana trench in the  
11 western Pacific Ocean. A 3-D interpreted flexural deformation surface of the subducting Pacific  
12 Plate was obtained by removing from the observed bathymetry the effects of sediment loading,  
13 isostatically-compensated topography based on gravity modeling, age-related lithospheric  
14 thermal subsidence, and residual short-wavelength features. We analyzed flexural bending of 75  
15 across-trench profile sections and calculated five best-fitting tectonic and plate parameters that  
16 control the flexural bending. Results of analysis revealed significant along-trench variations: the  
17 trench relief varies from 0.9 to 5.7 km, trench-axis vertical loading ( $-V_0$ ) from  $-0.73 \times 10^{12}$  to  
18  $3.17 \times 10^{12}$  N/m, and axial bending moment ( $-M_0$ ) from  $0.1 \times 10^{17}$  to  $2.7 \times 10^{17}$  N. The effective  
19 elastic plate thickness seaward of the outer-rise region ( $T_e^M$ ) ranges from 45 to 52 km, while that  
20 trench-ward of the outer-rise ( $T_e^m$ ) ranges from 19 to 40 km. This corresponds to a reduction in  
21  $T_e$  of 21-61%. The transition from  $T_e^M$  to  $T_e^m$  occurs at a breaking distance of 60-125 km from  
22 the trench axis, which is near the outer-rise and corresponds to the onset of observed pervasive  
23 normal faults. The Challenger Deep area is associated with the greatest trench relief and axial

24 vertical loading, while areas with seamounts at the trench axis are often associated with more  
25 subtle trench relief, smaller axial vertical loading, and greater topographic bulge at the outer-rise.

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27 *Keywords:* Mariana trench; Axial vertical force; Axial bending moment; Effective elastic  
28 thickness

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## 32 **1. Introduction**

33 The greatest flexural bending of Earth's oceanic lithosphere occurs at subduction zones.  
34 During subduction, the downgoing plate flexes in response to various types of tectonic forces,  
35 e.g., trench-axis vertical loading, axial bending moment, distributed sediment loading, and  
36 horizontal buckling. The flexural bending produces distinct seafloor sloping towards the trench  
37 axis, as well as gentle upward seafloor bulging near the outer-rise region seaward of the trench  
38 (e.g., [Hanks, 1971](#); [Bodine and Watts, 1979](#); [Harris and Chapman, 1994](#); [Bry and White, 2007](#)).  
39 Furthermore, as flexural deformation becomes significant, bending stresses could exceed the  
40 rock yield strength within the most deformed portions of the lithosphere (e.g., [McNutt and](#)  
41 [Menard, 1982](#); [Ranalli, 1994](#)), causing pervasive faulting and tensional earthquakes in the upper  
42 plate ([Christensen and Ruff, 1983](#); [Masson, 1991](#); [Ranero et al., 2005](#); [Naliboff et al., 2013](#)),  
43 local plastic yielding ([Turcotte et al., 1978](#); [Bodine and Watts, 1979](#); [McNutt, 1984](#); [McAdoo et](#)  
44 [al., 1985](#)), and a noticeable reduction in the effective elastic thickness of the plate, especially  
45 near the outer-rise region ([Judge and McNutt, 1991](#); [Levitt and Sandwell, 1995](#); [Watts, 2001](#);  
46 [Billen and Gurnis, 2005](#); [Contreras-Reyes and Osses, 2010](#)). Thus the observed spatial variations

47 in flexural bending of a specific subducting plate could provide important constraints on trench  
48 tectonic loading and lithospheric strength (e.g., [Mueller and Phillips, 1995](#); [Capitanio et al.,](#)  
49 [2009](#); [Capitanio and Morra, 2012](#)).

50 In this study, we investigated variations in flexural bending along the Mariana trench of the  
51 western Pacific Ocean (Fig. 1). We chose the Mariana trench as a study area for several reasons:  
52 (1) it exhibits significant along-trench changes in trench depth, slope, and outer-rise bulge (Fig.  
53 2); (2) it contains the greatest trench depth of the world, the Challenger Deep, with trench depth  
54 of about 10.9 km and trench relief of 5.7 km (Fig. 2); (3) high-resolution multi-beam bathymetric  
55 data are available for a significant portion of the trench axis and the outer-rise region, aiding the  
56 identification of detailed features; and (4) the 20-Ma difference in the crustal age of the  
57 subducting plate is relatively small comparing to the overall age of 140-160 Ma, thus facilitating  
58 analysis of factors unrelated to plate age.

59 A major challenge in flexural bending analysis is the identification of the “deformed shape”  
60 of a subducting plate from the complex seafloor topography that contains many other features  
61 unrelated to flexural bending, such as seamounts and volcanic ridges. Previous studies attempted  
62 to bypass this problem by choosing sparse topographic and free-air gravity profiles away from  
63 seamounts and ridges or designating these features as part of data uncertainties (e.g., [Bodine and](#)  
64 [Watts, 1979](#); [Judge and McNutt, 1991](#); [Levitt and Sandwell, 1995](#); [Billen and Gurnis, 2005](#); [Bry](#)  
65 [and White, 2007](#); [Contreras-Reyes and Osses, 2010](#)). However, in regions that contain abundant  
66 seamounts and ridges, such as near the Mariana trench (Fig. 2), these traditional approaches are  
67 inadequate for investigating the spatial variations in plate bending. In this paper, we present a  
68 new approach in identifying the deformed shape of a bending plate. Instead of using seafloor  
69 bathymetry directly, we conduct the analysis in two steps: First, we calculated “non-isostatic”

70 topography by removing from the observed bathymetry the effects of sediment loading,  
71 isostatically-compensated topography based on gravity modeling, and age-related lithospheric  
72 thermal subsidence. By removing these non-flexural effects, the resultant “non-isostatic”  
73 topography proves to be a much-improved approximation to the deformed shape of a bending  
74 plate. Second, we further removed short-wavelength features from the “non-isostatic”  
75 topography to obtain an interpreted 3-D surface of plate flexural bending.

76 We adopted a simplified model of a lithospheric plate of varying effective elastic thickness  
77 overlying an inviscid asthenosphere and analyzed flexural bending along 75 across-trench  
78 sections, each consists of ten profiles spanning over an along-trench distance of about  $0.2^\circ$ . Our  
79 analysis illustrated that these observed plate bending profiles could be explained by flexural  
80 deformation models assuming various forms of spatial variations in plate thickness. However,  
81 the vast majority of the observed plate bending profiles could be adequately reproduced by a  
82 simplified model, in which the deforming plate has only two characteristic values of effective  
83 elastic thickness, seaward ( $T_e^M$ ) and trench-ward ( $T_e^m$ ), respectively, of a breaking point near the  
84 outer-rise region. For each section, we then calculated five best-fitting parameters including axial  
85 vertical force ( $-V_0$ ), axial bending moment ( $-M_0$ ), maximum effective elastic thickness ( $T_e^M$ ),  
86 minimum elastic thickness ( $T_e^m$ ), and the breaking point distance ( $x_r$ ) between sections of the  
87 maximum and minimum elastic thickness. Modeling results revealed significant changes in  
88 tectonic loading and plate deformation along the Mariana trench.

89

## 90 **2. Identification of Plate Deformation Caused by Flexural Bending**

91 To better identify plate flexural bending near the trench axis, we first subtracted from the  
92 observed seafloor bathymetry the following predictable components that are not directly related

93 to trench-axis plate bending: (1) sediment loading; (2) isostatically-compensated topography,  
94 including features such as seamounts and volcanic ridges with crustal roots, calculated through  
95 gravity analysis; and (3) age-related lithospheric thermal subsidence (Muller et al., 2008)  
96 assuming 1-D vertical cooling of the lithosphere. The resultant “non-isostatically-compensated  
97 topography” should reflect primarily the topographic features that are dynamically supported by  
98 stresses in the lithospheric plate, including trench-related plate bending, as well as uncertainties  
99 in the above estimation of various loading features.

100 (1) *Seafloor Bathymetry*. We constructed a bathymetric database with grid spacing of 0.25’  
101 (Fig. 2), which combines two primary data sources: (1) high-resolution multi-beam bathymetric  
102 data from the database of the National Geophysical Data Center (NGDC, Lim et al., 2013); and  
103 (2) the GEBCO08 data with grid spacing of 0.5’ (<http://www.gebco.net>). Our study area covers a  
104 total along-trench distance of about 2,500 km, where multi-beam data are available for the  
105 distance of 100-2,000 km (Fig. 9). The trench depth (blue curve in Fig. 6a) is about 5-7 km near  
106 the Caroline Ridge (Region 1, at distance of 0-250 km, Fig. 2). At the Challenger Deep (at  
107 distance of about 400 km), the trench reaches a maximum depth of about 10.9 km. Another area  
108 of relatively deep trench is located at distance of about 650 km. The trench depth shows long-  
109 wavelength decrease northward. Prominent seamounts are located on the trench axis at distance  
110 of 1,350, 1,600, 1,800, and 2,300 km, respectively, reducing the local trench depth to only 5-6  
111 km (Figs. 2 and 6a).

112 The seafloor bathymetry also varies significantly seaward of the trench axis. The Caroline  
113 Ridge (Region 1, Fig. 2) is located near and off the trench axis at distance of about 0-250 km,  
114 while the Caroline Islands Chain (Region 2) intersects the trench axis at distance of about 600-  
115 700 km. A prominent trench-parallel belt of seamounts (Region 3), with trench-perpendicular

116 width of about 250 km, appears at distance of 800-1,350 km. Another prominent group of  
 117 seamounts, with relatively wide seamount bases and shallow apexes (Region 4), intersects the  
 118 trench axis at distance of 1,600-2,000 km. At a section of oblique subduction (at distance of  
 119 2,000-2,400 km), seamounts are absent seaward of the trench axis within 250 km. Starting at  
 120 distance of 2,400 km and northward, another prominent ridge (Region 5) approaches the trench  
 121 axis.

122 (2) *Sediment Loading*. We extracted sediment thickness data from the NGDC sediment  
 123 database (Divins, 2003) with grid spacing of 5' (Fig. 3a). Thick sediments appear in four regions:  
 124 the southwest corner of the Caroline Ridge region (up to 0.6 km of sediment thickness); the  
 125 eastern part of the study area at latitude 6°-15°N (up to 0.5 km); an off-axis region at latitude  
 126 17°-23°N (up to 0.6 km); and a narrow belt along the trench axis at latitude 12.5°-22°N (up to  
 127 0.25 km). For the rest of the study region, the interpolated sediment thickness is less than 0.1 km.  
 128 However, the interpolated sediment grids likely have under-sampled the true local sediment  
 129 thickness.

130 (3) *Isostatically-Compensated Topography*. For topographic features that are locally  
 131 compensated, e.g., seamounts and ridges with crustal roots, we calculated the isostatic  
 132 topography based on Airy-Heiskanen model. The isostatic topography is calculated as  
 133  $T_{iso} = (H_c - \overline{H}_c) \times (\rho_m - \rho_c) / (\rho_m - \rho_w)$ , where  $H_c$  is the gravity-derived crustal thickness,  $\overline{H}_c$  is  
 134 a reference crustal thickness, and  $\rho_w$ ,  $\rho_c$ , and  $\rho_m$  are densities of water, crust, and mantle,  
 135 respectively (Table 1). We used gravity-derived crustal thickness (Fig. 3c) calculated from  
 136 gravity inversion using methods similar to Parker (1973), Kuo and Forsyth (1988), and Wang et  
 137 al. (2011) and calibrated using available seismic data (see Appendix C in Supplementary  
 138 Materials).

139 Several regions are associated with relatively thick crust (Fig. 3c) and thus high values of  
140 calculated isostatic topography: a broad region close to the Caroline Ridge and Caroline Islands  
141 Chain at 6°-12°N (Regions 1 and 2, up to 27 km of crustal thickness and 5.5 km of isostatic  
142 topography); a trench-parallel belt at 12°-16.5°N (Region 3, up to 18 km crust and 3.2 km of  
143 isostatic topography); two E-W trending seamount groups at 17°-21.5°N (Region 4, up to 20 km  
144 crust and 3.7 km of isostatic topography); and an NW-SE trending group of ridges and  
145 seamounts at 25°-27.5°N (Region 5, also up to 20 km crust and 3.7 km of isostatic topography).  
146 For the rest of the study region, the calculated crustal thickness is about 3-6 km, corresponding to  
147 isostatic topography of -0.8 to 0 km.

148 (4) *Non-Isostatically-Compensated Topography*. We calculated non-isostatic topography ( $T_{n-iso}$ )  
149 by removing from the observed bathymetry (Fig. 2) the effects of sediment loading (Fig. 3a),  
150 thermal subsidence, and isostatically-compensated topography ( $T_{iso}$ ). On the map of non-isostatic  
151 topography (Fig. 3d), the most prominent feature is low topography along the Mariana trench  
152 with maximum depth near the Challenger Deep. The Caroline Ridge and Caroline Islands Chain  
153 are associated with much more subdued features in the map of non-isostatic topography (Fig. 3d)  
154 than in the map of observed bathymetry (Fig. 2). Similarly, some of the seamounts are more  
155 subdued in the non-isostatic topography than in the observed bathymetry. We suggest that the  
156 remaining short-wavelength non-isostatic topography of the ridges and seamounts are either  
157 supported by stresses in the lithospheric plate or due to inherent uncertainties in the estimation of  
158 crustal thickness using gravity analysis. Along the trench axis, the non-isostatic topography  
159 shows great long-wavelength variations with minimum values at along-trench distances of about  
160 400 km (Challenger Deep) and 650 km (black curve in Fig. 6b). The long-wavelength along-  
161 trench variations are greatly reduced on trench-parallel profiles taken at across-trench distances

162 of 100 km (the outer-rise region, red curve in Fig. 6b) and 550 km (far field, blue curve in Fig.  
163 6b). This suggests that the great along-trench axis variations in the non-isostatic topography  
164 (black curve in Fig. 6b) could reflect the significant along-trench variability in flexural bending  
165 of the subducting plate.

166 (5) *Flexural Bending of the Subducting Plate.* We extracted a total of 750 across-trench  
167 profiles, each of 600-km long, spanning at an interval of  $0.02^\circ$  (about 3.3 km) along the Mariana  
168 trench. Every ten profiles were stacked together to form a section (e.g., Figs. 4b-d), resulting in a  
169 total of 75 sections (Fig. 5; Figs. S1-8 of Appendix A in Supplementary Materials). For each  
170 section, we determined a flexural bending model (red dashed curves, Figs. 4 and 5; Figs. S1-8 in  
171 Supplementary Materials) that best captures the long-wavelength characteristics of the non-  
172 isostatic topography profile both across the trench strike (blue curves, Figs. 4 and 5; Figs. S1-8 in  
173 Supplementary Materials) and along the trench strike (Fig. 6b). On several across-trench sections  
174 (e.g., Sections 1-5, 56-58, 72-75), seamounts have covered up the trench axis or obscured a  
175 significant portion of the far-field reference seafloor depth (Supplementary Table 1). When  
176 estimating the overall shape of flexural bending, we ignored short-wavelength features of the  
177 seamounts and their periphery depression in the non-isostatic topography (Figs. 4 and 5; Figs.  
178 S1-8 in Supplementary Materials). The 75 sections were then interpolated to create a map of  
179 flexural bending (Fig. 3e). The greatest flexural bending occurs at along-trench distance of 350-  
180 650 km, including the Challenger Deep. The difference between the non-isostatic topography (Fig.  
181 3d) and the flexural bending model (Fig. 3e) is showing as a map of residuals (Fig. 3f), which  
182 appears to contain primarily local features such as seamounts and surrounding depression.

183

### 184 **3. Modeling of Plate Flexural Bending**



185 In thin-plate approximation, the vertical deflection of a plate is governed by the balance  
 186 among various forces (Turcotte and Schubert, 2002):

$$187 \quad -\frac{d^2M}{dx^2} + \frac{d}{dx}\left(F\frac{dw}{dx}\right) + (\rho_m - \rho_w)gw = (\rho_s - \rho_w)gh_s(x), \quad (1)$$

188 where  $M$  is bending moment,  $F$  is horizontal buckling force,  $(\rho_m - \rho_w)gw$  represents hydrostatic  
 189 restoring force,  $(\rho_s - \rho_w)gh_s(x)$  is vertical sediment loading, and  $\rho_s$  and  $h_s$  are the sediment  
 190 density and thickness, respectively.

191 The bending moment is proportional to the vertical deflection by  $M = -D\frac{d^2w}{dx^2}$ , where flexural

192 rigidity  $D = \frac{ET_e^3}{12(1-\nu^2)}$ ,  $E$  is Young's modulus,  $\nu$  is Poisson's ratio, and  $T_e$  is effective elastic

193 plate thickness. The vertical force is related to both the bending moment and horizontal force by

$$194 \quad V = \frac{dM}{dx} - F\frac{dw}{dx}. \text{ As a first-order approximation, we ignored the horizontal buckling force}$$

195 similar to the analyses of Caldwell et al. (1976), Molnar and Atwater (1978), and Contreras-

196 Reyes and Osses (2010). Constant parameters assumed in the analysis are described in Table 1.

197 Boundary conditions of the vertical deflection include the following:  $w = 0$  and  $\frac{dw}{dx} = 0$  at  $x =$

$$198 \quad +\infty, \text{ while } D\frac{d^2w}{dx^2} = -M_0 \text{ and } \frac{dM}{dx} = -V_0 \text{ at } x = 0.$$

199 We used a simplified model of an elastic plate of two effective elastic thickness values. We

200 assumed that the effective elastic thickness changes from  $T_e^M$  (seaward of the outer-rise) to  $T_e^m$

201 (near the trench axis), in order to simultaneously replicate the observed steep slope trench-ward

202 of the outer-rise as well as the relatively long flexural wavelength seaward of the outer-rise

203 region (Turcotte et al. 1978; Judge and McNutt 1991). The transition occurs at a breaking

204 distance  $x_r$  near the outer-rise (Fig. 4a). The reduced effective elastic thickness is assumed to  
 205 reflect the onset of pervasive normal faulting within the upper plate near the outer-rise (Fig. 4a).

206 (1) *Calculation of Trench-Axis Vertical Force.* We found that the vertical force at the trench  
 207 axis is proportional to the total area of the vertical deflection integrated over the entire across-  
 208 trench profile. For the case of a constant plate thickness, the vertical deflection (Turcotte and

209 Schubert, 2002) is given by  $w(x) = \frac{\alpha^2 e^{-x/\alpha}}{2D} \left[ -M_0 \sin\left(\frac{x}{\alpha}\right) + (V_0 \alpha + M_0) \cos\left(\frac{x}{\alpha}\right) \right]$ , where the

210 flexural wavelength  $\alpha$  is defined by  $\alpha = \left[ \frac{4D}{(\rho_m - \rho_w)g} \right]^{1/4}$ . Direct integration of the above

211 equation yields

$$212 \quad -V_0 = (\rho_m - \rho_w)g \int_0^{+\infty} w(x) dx. \quad (2)$$

213 We conducted a series of tests for deflection of plates with variable thickness and found that the  
 214 above Eqn. 2 still holds for cases when the effective plate thickness varies horizontally. Thus by  
 215 integrating the interpreted vertical deflection of a given profile (Fig. 5), we can readily calculate  
 216 the trench-axis vertical force (Fig. 6c and Supplementary Table 1).

217 (2) *Inversion of Axial Bending Moment and Plate Thickness Variation.* We next inverted for  
 218 four best-fitting parameters,  $-M_0$ ,  $T_e^M$ ,  $T_e^m$ , and  $x_r$ , for each section. Using the finite-difference  
 219 method of Contreras-Reyes and Osses (2010), we discretized each profile section into a series of  
 220 nodes with a uniform spacing of 3 km (see Appendix B in Supplementary Materials). Sediment  
 221 loading was prescribed at each node point. For each section, we then inverted for a set of best-  
 222 fitting parameters that minimize the root mean square (RMS) misfit between the non-isostatic  
 223 topography ( $T_{n-iso}$ , blue curves in Figs. 4 and 5; Figs. S1-8 in Supplementary Materials) for areas

224 away from seamounts and the flexural bending model (red dashed curves in Figs. 4 and 5; Figs.  
225 S1-8 in Supplementary Materials).

226

## 227 **4. Results**

228 Results of analysis revealed that both the trench-axis loading and plate thickness vary  
229 significantly along the Mariana trench.

### 230 **4.1 Along-Trench Variations in Trench Relief and Axial Loading**

231 The trench relief, which was calculated from subtracting the trench-axis depth from a far-field  
232 reference depth, varies from 0.9 to 5.7 km along the Mariana trench (black curve in Fig. 6a and  
233 Supplementary Table 1). Within the first 230 km from the southwestern end of the trench, the  
234 trench relief ranges from 1.7 to 3.4 km. The greatest trench relief of 5.7 km is at the Challenger  
235 Deep. Another area of large trench relief of about 5.2 km is located east of the Challenger Deep  
236 at along-trench distance of about 650-670 km (Fig. 6a). In between the above two deep locations,  
237 the Caroline Islands Chain (Region 2, Fig. 2) has trench relief of about 4.0 km (Fig. 6a). From  
238 distance of 850 to 1,250 km, the trench relief gradually decreases from 3.9 to 1.8 km. From  
239 1,250 to 2,250 km, the trench relief ranges from 1.4 to 3.0 km with prominent trench-axis highs  
240 located at 1,250-1,300 km, 1,600-1,650 km, 1,950-2,050 km, respectively. The section of the  
241 trench at distance of 1,950-2,300 km, which is associated with relatively oblique convergence  
242 angles, has trench relief of 1.4 to 2.2 km.

243 The calculated axial vertical force (Fig. 6c) is in general proportional to the trench relief (Fig.  
244 6a). The two areas of great trench relief, at the Challenger Deep and its eastern section, were  
245 calculated to be subjected to large axial vertical loading of  $3.17 \times 10^{12}$  N/m and  $2.1 \times 10^{12}$  N/m,  
246 respectively (Supplementary Table1). Along the Mariana trench, several areas are associated

247 with upward, instead of downward, vertical force of small magnitude (red arrows in Fig. 6c and  
 248 black arrows in Fig. 3b; Supplementary Table 1); these sections account for more than 20%  
 249 length of the Mariana trench. The total trench-axis vertical force integrated over the 2,500-km-  
 250 long study area is about  $1.66 \times 10^{18}$  N; sections with trench relief greater than 3.0 km contribute to  
 251 more than 80% of the total vertical force. The axial vertical force averaged over the trench length  
 252 is about  $0.67 \times 10^{12}$  N/m.

253 The calculated trench-axis bending moment (Fig. 6d; Supplementary Table 1) also appears to  
 254 be correlated with trench relief (Figs. 6a). The axial bending moment is the smallest ( $-M_0 =$   
 255  $0.1 \times 10^{17}$  N) near the Challenger Deep (Fig. 6d). In contrast, the bending moment is in general  
 256 greater for sections of relatively small trench relief. The calculated bulge height at the outer-rise  
 257 ( $w_b$ ) ranges from 70 to 650 m (Supplementary Table 1). The bulge height is the smallest at the  
 258 Challenger Deep area ( $w_b = 70$  m), while large bulge height ( $w_b > 500$  m) is associated with  
 259 sections of relatively large axial bending moment ( $M_0 > 2.4 \times 10^{17}$  N) at distance of 920-1,000,  
 260 1,070-1,090, and 1,390-1,440 km (Supplementary Table 1). The across-trench distance of the  
 261 bulge height ( $x_b$ , blue curve in Fig. 6g) varies in the range of 69-180 km from the trench axis.

## 262 **4.2 Along-Trench Variations in Effective Elastic Thickness**

263 To replicate the far-field long-wavelength flexural bending, the effective elastic thickness of  
 264 the plate seaward of the outer-rise ( $T_e^M$ ) is calculated to range from 45 to 52 km (blue curve in  
 265 Fig. 6e; Supplementary Table 1). However, to replicate the observed steep seafloor slope towards  
 266 the trench axis, the effective elastic thickness trench-ward of the outer-rise ( $T_e^m$ ) is only 19 to 40  
 267 km (black curve in Fig. 6e). The transition from  $T_e^M$  to  $T_e^m$  occurs at a breaking distance  $x_r =$   
 268 60-125 km from the trench axis (black curve in Fig. 6g). The resultant reduction in the effective  
 269 elastic thickness, i.e.,  $1 - (T_e^m/T_e^M)$ , is in the range of 21-61% (Fig. 6f). The greatest reduction in

270  $T_e$  is about 61%, occurring near the Challenger Deep area, where the plate bends significantly  
 271 within a narrow distance of  $x_r = 75-85$  km. Reduction in  $T_e$  of greater than 50% also occurs at  
 272 four other areas at distance of 0-50, 1,180-1,230, 1,490-1,510, 1,760-1,860 km, respectively,  
 273 where the calculated breaking distance is relatively small ( $x_r < 90$  km, Fig. 6g; Supplementary  
 274 Table 1). In contrast, areas with smaller reduction in elastic thickness ( $< 30\%$ ), e.g., at distance  
 275 of 160-250, 1,320-1,350, 1,460-1,490, 1,560-1,610, and 1,980-2,140 km, are associated with  
 276 large breaking distance ( $x_r > 100$  km, Fig. 6g; Supplementary Table 1) or smaller trench relief ( $<$   
 277 2 km) (blue dots in Fig. 7b). Our results revealed that the reduction in  $T_e$  along the Mariana  
 278 trench does not exceed 61%, implying that an elastic core remains in the subducting plate despite  
 279 pervasive normal faulting caused by flexural bending near the trench axis (Fig. 4a).

280 For a plate of constant elastic thickness, the trench relief can be calculated as

281  $w_0 = \frac{\alpha^2(V_0\alpha + M_0)}{2D}$ , where flexural wavelength  $\alpha$  and flexural rigidity  $D$  are a function of  $T_e$ .

282 We compared the observed trench relief with the predicted values for the end-member cases of  
 283  $T_e = T_e^m$  and  $T_e^M$ , respectively (Fig. 7a). The  $w_0$  calculated assuming  $T_e = T_e^m$  is only 8 % greater  
 284 than the observed values with a regression coefficient of 0.99. This implies that the observed  
 285 trench relief is controlled primarily by  $T_e^m$ , and not by  $T_e^M$ .

286

## 287 **5. Discussion**

### 288 **5.1 Uncertainties in Data and Analysis**

289 Several aspects of the above analysis might be associated with uncertainties. (1) The northern  
 290 most 500-km of the trench-axis ( $\sim 21^\circ-25^\circ\text{N}$ ) lacks high-resolution multi-beam bathymetric data  
 291 (Fig. 9). While lacking multi-beam bathymetry is not likely to affect significantly our inverted  
 292 flexural bending parameters, it would prevent the identification of the onset location of normal

293 faults (Fig. 9). (2) The data coverage of sediment thickness might be highly non-uniform and  
294 sparse for much of the study region. However, our example test for Section 49, which has a  
295 maximum sediment thickness of 0.4 km, showed that the inverted flexural parameters change  
296 little with versus without considering sediment loading. Thus we infer that the lack of high-  
297 resolution data of sediment thickness might not change the overall pattern of the calculated  
298 flexural parameters. (3) There are inherent uncertainties associated with gravity-derived crustal  
299 thickness (e.g., [Wang et al., 2011](#)), leading to uncertainties in the calculated isostatic and non-  
300 isostatic topography. These uncertainties, however, are difficult to quantify without independent  
301 seismic constraints.

302 The subducting Pacific plate is concave along the Mariana trench. [Bonnardot et al. \(2008\)](#)  
303 numerically modeled the effects of trench curvature on the deformation of a subducting plate for  
304 different curvature radius values. We interpolated their modeling results for the estimated  
305 curvature values of the Mariana trench. The trench curvature appears to have greater effects on  
306 the modeling of axial vertical loading than on other parameters.

## 307 **5.2 Unique Characteristics of the “Seamount” Sections**

308 While the trench relief is most sensitive to the axial vertical force, the predicted topographic  
309 bulge height at the outer-rise is much more sensitive to the axial bending moment. We separated  
310 the 75 sections of the Mariana trench into four groups according to the average value of the  
311 calculated axial vertical force (Fig. 8a). The averaged value of the calculated trench relief is  
312 greater for sections of larger vertical force. In contrast, the averaged value of the bulge height at  
313 the outer-rise is greater for sections of larger axial bending moment (Fig. 8b).

314 Several areas of the Mariana trench are associated with small amplitudes of upward axial  
315 vertical force (red arrows in Fig. 6c). The averaged across-trench profiles of these “seamount”

316 sections (striped belt in Fig. 8c) are of relatively small vertical force. We further noted that these  
317 “seamount” sections are associated with relative large topographic bulge at the outer-rise  
318 (Supplementary Table 1). The averaged height of the topographic bulge for these “seamount”  
319 sections is 388 m, which is much greater than the averaged value of 288 m for the remaining  
320 “non-seamount” sections. Correspondingly, the calculated axial bending moment for individual  
321 “seamount” sections (Fig. 6d), as well as the averaged bending moment for all “seamount”  
322 sections (Fig. 8d), are greater than that of “non-seamount” sections. While the observed higher  
323 topographic bulge at the outer-rise could be caused by greater axial bending moment for the  
324 “seamount” sections, they might also be caused by significant horizontal buckling force in the  
325 plate, due to the resistance of the seamounts to subduction, which was not modeled in the present  
326 analysis.

### 327 **5.3 Causes of Reduction in Elastic Plate Thickness**

328 Results of analysis indicated a reduction in the effective plate thickness of 21-61% near the  
329 outer-rise region along the Mariana trench (Fig. 6f). Thus lateral changes in the plate property  
330 are likely to be significant, while the specific values of the plate thickness reduction depend on  
331 our specific model assumptions. Similar results were obtained from analysis of the central  
332 Mariana trench ([Oakley et al., 2008](#)).

333 We hypothesize that the calculated reduction in the effective elastic thickness reflects the  
334 effects of pervasive normal faulting in a bending plate (Fig. 4a). Under the trench-axis loading,  
335 the upper half of the bending plate would be in extension while the lower plate would be under  
336 compression. Computational geodynamic models showed that the loss of rock cohesion and  
337 strain weakening caused by slip on normal faults could significantly reduce the effective elastic  
338 strength of a lithospheric plate ([Rupke et al, 2004](#); [Faccenda et al., 2009](#)).

339 The development of normal faults is likely to be distributed over a broad region, and thus the  
340 reduced elastic thickness  $T_e^m$  and the breaking distance  $x_r$  are over simplifications. Nevertheless,  
341 the location of the calculated  $x_r$  in general appears to be consistent with the observed outer  
342 boundary of a zone of pervasive normal faults at sections with multi-beam bathymetry coverage  
343 (Fig. 9). The Challenger Deep area of the Mariana trench is associated with a relatively large  
344 reduction in the effective elastic thickness, which might reflect relatively extensive normal  
345 faulting in response to the large axial vertical force (Fig. 9b).

346

## 347 **6. Conclusions**

348 1. Results of analysis revealed significant variations in trench-axis loading and plate  
349 mechanical property along the Mariana trench. The trench relief varies from 0.9 to 5.7 km; the  
350 trench-axis vertical force varies from  $-0.73 \times 10^{12}$  to  $3.17 \times 10^{12}$  N/m; and the axial bending  
351 moment varies from  $0.1 \times 10^{17}$  to  $2.7 \times 10^{17}$  N.

352 2. Modeling of long-wavelength flexural bending seaward of the outer-rise region indicates  
353 that the effective elastic plate thickness of the incoming plate ( $T_e^M$ ) to be in the range of 45-52  
354 km. Trench-ward of the outer-rise, the observed steep seafloor slope indicates thinner effective  
355 elastic thickness ( $T_e^m$ ) of 19-40 km; the corresponding reduction in  $T_e$  is 21-61%. The transition  
356 from  $T_e^M$  to  $T_e^m$  occurs at a breaking distance of 60-125 km from the trench axis, corresponding  
357 to the onset of the observed zones of pervasive normal faulting.

358 3. The Challenger Deep area in the southwestern Mariana trench is associated with the  
359 greatest trench relief, axial vertical loading, and reduction in  $T_e$ . Several areas with seamounts at  
360 the trench axis are associated with shallower trench relief, smaller axial vertical force, and higher  
361 topographic bulge at the outer-rise.



362

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372

### 373 **References**

- 374 Billen, M.I., Gurnis, M., 2005. Constraints on subducting plate strength within the Kermadec  
375 trench, *J. Geophys. Res.*, 110, doi:10.1029/2004JB003308.
- 376 Bird, P., 2003. An updated digital model of plate boundaries, *Geochem., Geophys., Geosyst.*, 4,  
377 doi:10.1029/2001GC000252.
- 378 Bodine, J.H., Watts, A.B., 1979. On the lithospheric flexure seaward of the Bonin and Mariana  
379 Trenches, *Earth Planet. Sci. Lett.* 43, 132–148.
- 380 Bonnardot, M.A., Hassani, R., Tric, E., Ruellan, E., Regnier, M., 2008. Effect of margin  
381 curvature on plate deformation in a 3-D numerical model of subducton zones, *Geophys. J.*  
382 *Int.* 173, 1084-1094.
- 383 Bry, M., White, N., 2007. Reappraising elastic thickness variation at oceanic trenches, *J.*  
384 *Geophys. Res.* 112, doi:10.1029/2005JB004190.

- 385 Caldwell, J.G., Haxby, W.F., Karig, D.E., Turcotte, D.L., 1976. On the applicability of a  
386 universal elastic trench profile, *Earth Planet. Sci. Lett.* 31, 239–246.
- 387 Capitanio, F.A., Morra, G., 2012. The bending mechanics on a dynamic subduction system:  
388 Constraints from numerical modelling and global compilation analysis, *Tectonophys.*, 522-  
389 523, 224–234.
- 390 Capitanio, F.A., Morra, G., Goes, S., 2009. Dynamics of plate bending at the trench and slab-  
391 plate coupling, *Geochem., Geophys., Geosyst.*, 10,  
392 doi:10.1029/2008GC002348.
- 393 Christensen, D.H., Ruff, L.J., 1983. Outer-rise earthquakes and seismic coupling, *Geophys. Res.*  
394 *Lett.* 10, 697-700.
- 395 Contreras-Reyes, E., Osses, A., 2010. Lithospheric flexure modelling seaward of the  
396 Chile trench: Implications for oceanic plate weakening in the trench outer rise region,  
397 *Geophys. J. Int.* 182, 97-112.
- 398 Divins, D., 2003. Total sediment thickness of the world's oceans and marinal seas. NOAA  
399 National Geophysical Data Center, Boulder, CO.  
400 <http://www.ngdc.noaa.gov/mgg/sedthick/sedthick.html>
- 401 Faccenda, M., Gerya, T.V., Burlini, L., 2009. Deep slab hydration induced by bending-related  
402 variations in tectonic pressure, *Nat. Geosci.* 2, 790-793.
- 403 Hanks, T.C., 1971. The Kuril trench-Hokkaido Rise system: Large shallow earthquakes and  
404 simple models of deformation, *Geophys. J. R. Astr. Soc.*, 23, 173.
- 405 Harris, R., Chapman, D., 1994. A comparison of mechanical thickness estimates from trough and  
406 seamount loading in the southeastern Gulf of Alaska, *J. Geophys. Res.*, 95, 9297-9317.
- 407 Judge, A.V., McNutt, M.K., 1991. The relationship between plate curvature and elastic plate

- 408 thickness: A study of the Peru-Chile Trench, *J. Geophys. Res.* 96, 16625–16639.
- 409 Kuo, B-Y., Forsyth, D.W., 1988. Gravity anomalies of the ridge-transform system in the South  
410 Atlantic between 31° and 34.5°S: Upwelling centers and variations in crustal thickness,  
411 *Marine Geophys. Res.* 10, 205-232.
- 412 Levitt, D.A., Sandwell, D.T., 1995. Lithospheric bending at subduction zones based on depth  
413 soundings and satellite gravity, *J. Geophys. Res.* 100, 379–400.
- 414 Lim, E., Sutherland, M.G., Friday, D.Z., Eakins, B.W., and McLean, S.J., 2013. Bathymetric  
415 Digital Elevation Model of the Mariana Trench, NOAA National Geophysical Data Center,  
416 U.S. Dept. of Commerce, Boulder, CO, 4 pp.
- 417 Masson, D.G., 1991. Fault patterns at outer trench walls, *Mar. Geophys. Res.*, 13, 209-225.
- 418 McAdoo, D.C., Martin, C.F., Poulouse, S., 1985. Seasat observations of flexure: Evidence for a  
419 strong lithosphere, *Tectonophys.*, 116, 209–222.
- 420 McNutt, M.K., 1984. Lithospheric flexure and thermal anomalies, *J. Geophys. Res.*, 89, 11,180–  
421 11,194.
- 422 McNutt, M.K., Menard, H.W., 1982. Constraints on yield strength in the oceanic lithosphere  
423 derived from observations of flexure, *Geophys. J. R. astr. Soc.*, 71, 363-394.
- 424 Molnar, P., Atwater, T., 1978. Interarc spreading and cordilleran tectonics as alternates related  
425 the age of subducted oceanic lithosphere. *Earth planet. Sci. Lett.* 41, 330–340.
- 426 Müller, R.D., Sdrolias, M., Gaina, C., Roest, W.R., 2008. Age, spreading rates, and spreading  
427 asymmetry of the world's ocean crust, *Geochem., Geophys., Geosyst.*, 9,  
428 doi:10.1029/2007GC001743.
- 429 Mueller, S., Phillips, R.J., 1995. On the reliability of lithospheric constraints derived from  
430 models of outer-rise flexure, *Geophys. J. Int.*, 123, 887-902.

- 431 Naliboff, J.B., Billen, M.I., Gerya. T., Saunders. J., 2013. Dynamics of outer-rise faulting in  
432 oceanic-continental subduction systems, *Geochem., Geophys., Geosyst.*, 14,  
433 doi:10.1002/ggge.20155.
- 434 Oakley, A.J., Taylor, B., Moore, G.F., 2008. Pacific plate subduction beneath the central  
435 Mariana and Izu-Bonin fore-arcs: New insights from an old margin, *Geochem., Geophys.,*  
436 *Geosyst.*, 9, doi:10.1029/2007GC001820.
- 437 Parker, R.L., 1973. The rapid calculation of potential anomalies. *Geophys. J. Roy. Astron. Soc.*  
438 31, 447-455.
- 439 Ranalli, G., 1994. Nonlinear flexure and equivalent mechanical thickness of the lithosphere,  
440 *Tectonophys.* 240, 107-114.
- 441 Ranero, C.R., Villaseor, A., Phipps Morgan, J., and Wdwinrebe, W., 2005. Relationship between  
442 bending-faulting at trenches and intermediate-depth seismicity. *Geochem., Geophys.,*  
443 *Geosyst.*, 6, doi:10.1029/2005GC000997.
- 444 Rupke, L.H., Morgan, J.P., Hort, M., Connolly, J.D.L., 2004. Serpentine and the subduction zone  
445 water cycle. *Earth Planet. Sci. Lett.* 223, 17–34.
- 446 Turcotte, D.L., McAdoo, D.C., Caldwell, J.G., 1978. An elastic-perfectly plastic analysis of the  
447 bending of the lithosphere at a trench. *Tectonophys.*, 47, 193-208.
- 448 Turcotte, D., Schubert, G., 2002. *Geodynamics*, 2<sup>nd</sup> Edition, Cambridge University Press, New  
449 York, 456 pp.
- 450 Sandwell, D.T., and Smith, W.H.F., 2009. Global marine gravity from retracked Geosat and  
451 ERS-1 altimetry: Ridge segmentation versus spreading rate. *J. Geophys. Res.*, 114, B01411,  
452 doi:10.1029/2008JB006008.
- 453 Wang, T., Lin, J., Tucholke, B., Chen, Y.J., 2011. Crustal thickness anomalies in the North

454 Atlantic Ocean basin from gravity analysis. *Geochem., Geophys., Geosyst.*, 12,  
455 doi:10.1029/2010GC003402.

456 Watts, A.B., 2001. *Isostasy and flexure of the lithosphere*, Cambridge University Press, New  
457 York, 458 pp.

458

459 **Figure Captions**

460 **Figure 1.** Tectonic map of the western Pacific Ocean including the Mariana trench. The Pacific  
 461 Plate is subducting under the Mariana and Philippines Plates (Bird, 2003). Challenger Deep is  
 462 the deepest part of the Mariana trench and the world. Dashed lines mark the study region of Fig.  
 463 2.

464  
 465 **Figure 2.** Seafloor bathymetry of the Mariana trench and surrounding regions. Dashed lines  
 466 mark the study area as shown in Fig. 3. Along-trench distance is measured from the southwestern  
 467 end of the Mariana trench. Circled numbers indicate tectonic provinces discussed in the text.

468  
 469 **Figure 3.** Maps of the study area. (a) Sediment thickness. (b) Basement depth. Black arrows  
 470 indicate sections experiencing small upward instead of downward vertical force (along-trench  
 471 locations shown in Fig. 6c). (c) Crustal thickness calculated from gravity analysis. This map is  
 472 used to calculate isostatic topography due to crustal thickness variations. (d) Non-isostatic  
 473 topography. (e) Flexural bending model interpolated from results along across-trench sections.  
 474 Results from profile sections of relatively poor constraints due to significant seamount effects are  
 475 not shown. (f) Residuals showing the difference between non-isostatic topography (panel d) and  
 476 flexural bending model (panel e).

477  
 478 **Figure 4.** (a) Schematic model of plate flexural bending. The vertical force ( $-V_0$ ) and bending  
 479 moment ( $-M_0$ ) are applied at the trench axis. Distance  $x_r$  is where the effective elastic thickness is  
 480 reduced from  $T_e^M$  to  $T_e^m$ . Distance  $x_b$  is the location of maximum uplift at the outer-rise. Area  
 481 with stripes illustrates the approximate location of expected pervasive normal faulting failure in a

482 zone of tectonic extension within the upper plate; the effective elastic thickness of this section of  
 483 the plate is reduced to  $T_e^m$  due to mechanical weakening by normal faulting. (b), (c), and (d)  
 484 correspond to sections of the greatest curvature, greatest trench relief, and shallowest section  
 485 along the Mariana trench, respectively. Basement topography of every ten individual profiles  
 486 was stacked to form an averaged section for modeling: grey and green curves show areas with  
 487 and without multi-beam bathymetry, respectively. Blue curves are the calculated non-isostatic  
 488 topography. Red dashed curves show flexural bending models that best fit non-isostatic  
 489 topography away from seamounts.

490

491 **Figure 5.** Fourteen example sections of flexural bending at different locations along the Mariana  
 492 trench. Basement topography of every ten individual profiles was stacked to form an averaged  
 493 section for modeling: grey and green curves show areas with and without multi-beam bathymetry,  
 494 respectively. Blue curves are the calculated non-isostatic topography. Red dashed curves show  
 495 flexural bending models that best fit non-isostatic topography away from seamounts.

496

497 **Figure 6.** Tectonic variables and calculated parameters along the Mariana trench. In panels c-g,  
 498 results from profile sections of relatively poor constraints due to significant seamount effects are  
 499 not shown. (a) Blue curve is the observed trench depth. Black curve is trench relief (measured  
 500 from a far-field reference seafloor depth to the trench axis). (b) Calculated non-isostatic  
 501 topography on trench-parallel profiles along the trench axis (black curve), near the outer-rise  
 502 region (100 km away from the trench axis, red curve), and at the far-field (550 km from the  
 503 trench axis, blue curve). (c) Calculated trench-axis vertical loading ( $-V_0$ ). Red arrows indicate  
 504 sections that are subjected to small upward vertical loading. (d) Calculated axial bending

505 moment ( $-M_0$ ). (e) Calculated effective elastic thickness. Blue and black curves mark the  
 506 maximum and minimum elastic thickness  $T_e^M$  and  $T_e^m$ , respectively. (f) Calculated reduction  
 507 from  $T_e^M$  to  $T_e^m$ . (g) Blue and black curves mark the across-trench distances for locations of the  
 508 maximum bulge height near the outer-rise region ( $x_b$ ) and the transition from the maximum to  
 509 minimum effective elastic thickness ( $x_r$ ), respectively.

510

511 **Figure 7.** (a) Correlation of the observed trench relief with the calculated trench relief ( $w_0$ ) for a  
 512 constant plate thickness model assuming  $T_e = T_e^m$  (red dots) and  $T_e = T_e^M$  (blue dots). R is the  
 513 regression coefficient. (b)  $T_e$  reduction as a function of trench relief and breaking distance  $x_r$ .

514

515 **Figure 8.** (a) Averaged shapes of modeled best-fitting flexural bending profiles along the  
 516 Mariana trench for four ranges of axial vertical force ( $-V_0$ ). (b) Averaged shapes of modeled best-  
 517 fitting flexural bending profiles for four ranges of axial bending moment ( $-M_0$ ). (c) Range of  
 518 profiles with seamounts near the trench axis (stripe) in comparison to profiles with relatively  
 519 small (red) and large (green) axial vertical force. (d) Range of profiles with seamounts near the  
 520 trench axis (stripe) in comparison to profiles with relatively small (green) and large (red) axial  
 521 bending moment.

522

523 **Figure 9.** Map of shaded relief of the Mariana trench. Red curves mark the location of the trench  
 524 axis, while blue curves illustrate the calculated location of the transition from maximum to  
 525 minimum elastic thickness ( $x_r$ ). Areas lack of high-resolution multi-beam bathymetry data are  
 526 marked by light green shades. Inset maps (a) and (b) show enlarged areas near the southern



527 Mariana trench. Note the general good correlation between  $x_r$  (blue curves) and the seaward  
528 boundary of the observed pervasive trench-parallel normal faults.

529

530 Table 1. Constant Parameters

<i>Symbol</i>	<i>Description</i>	<i>Value</i>	<i>Unit</i>
$E$	Young's modulus	$7 \times 10^{10}$	Pa
$g$	Acceleration due to gravity	9.81	$\text{m s}^{-2}$
$\nu$	Poisson's ratio	0.25	
$\rho_m$	Mantle density	3,300	$\text{kg m}^{-3}$
$\rho_s$	Sediment density	2,000	$\text{kg m}^{-3}$
$\rho_c$	Crust density	2,700	$\text{kg m}^{-3}$
$\rho_w$	Water density	1,030	$\text{kg m}^{-3}$

531

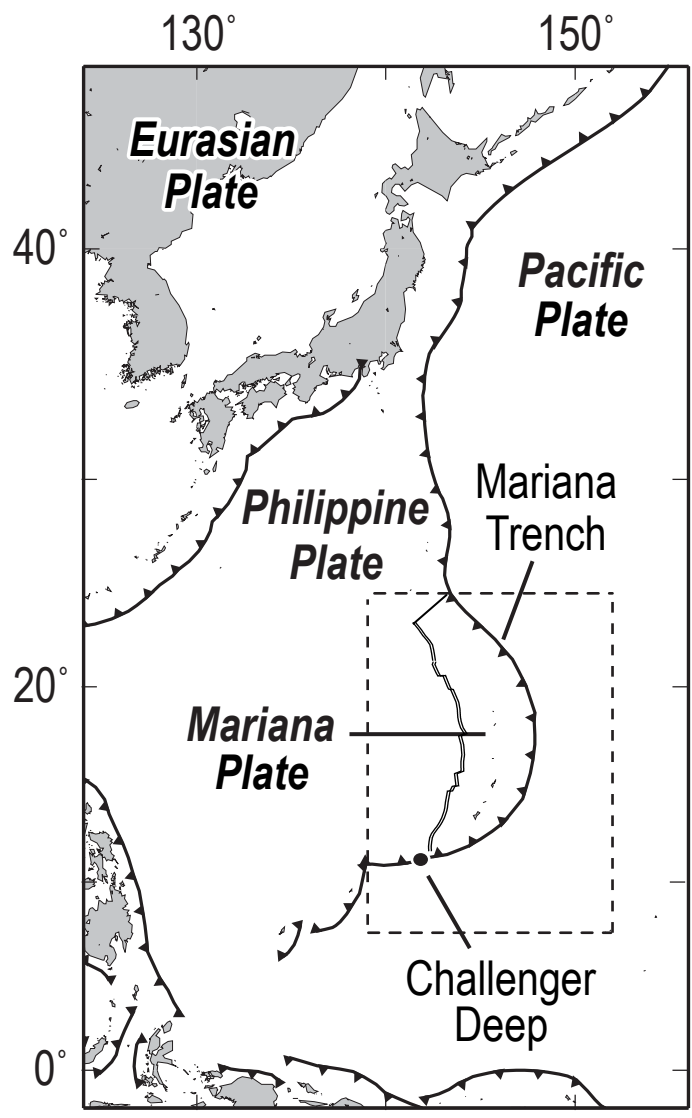


Fig. 1

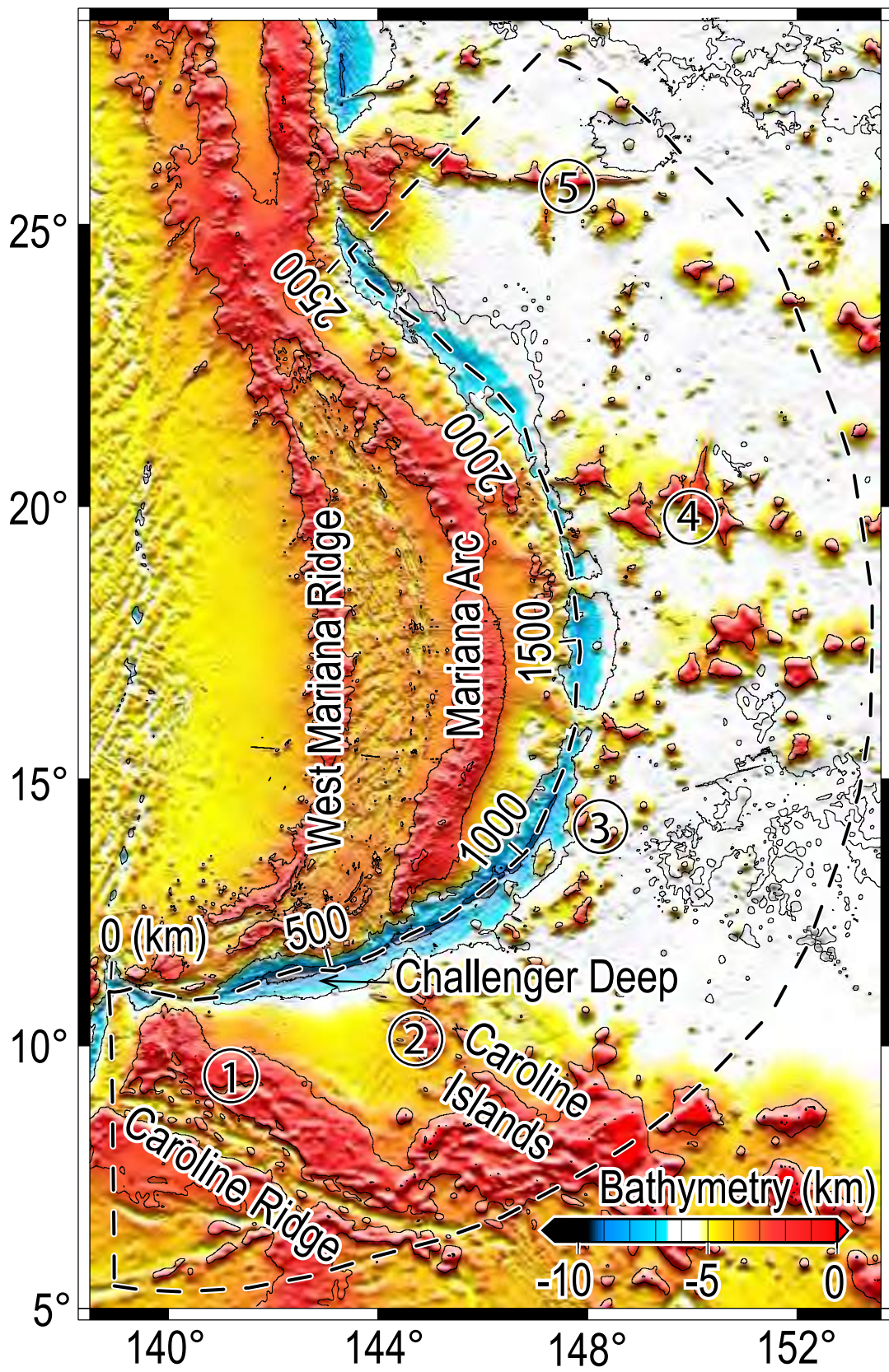


Fig. 2

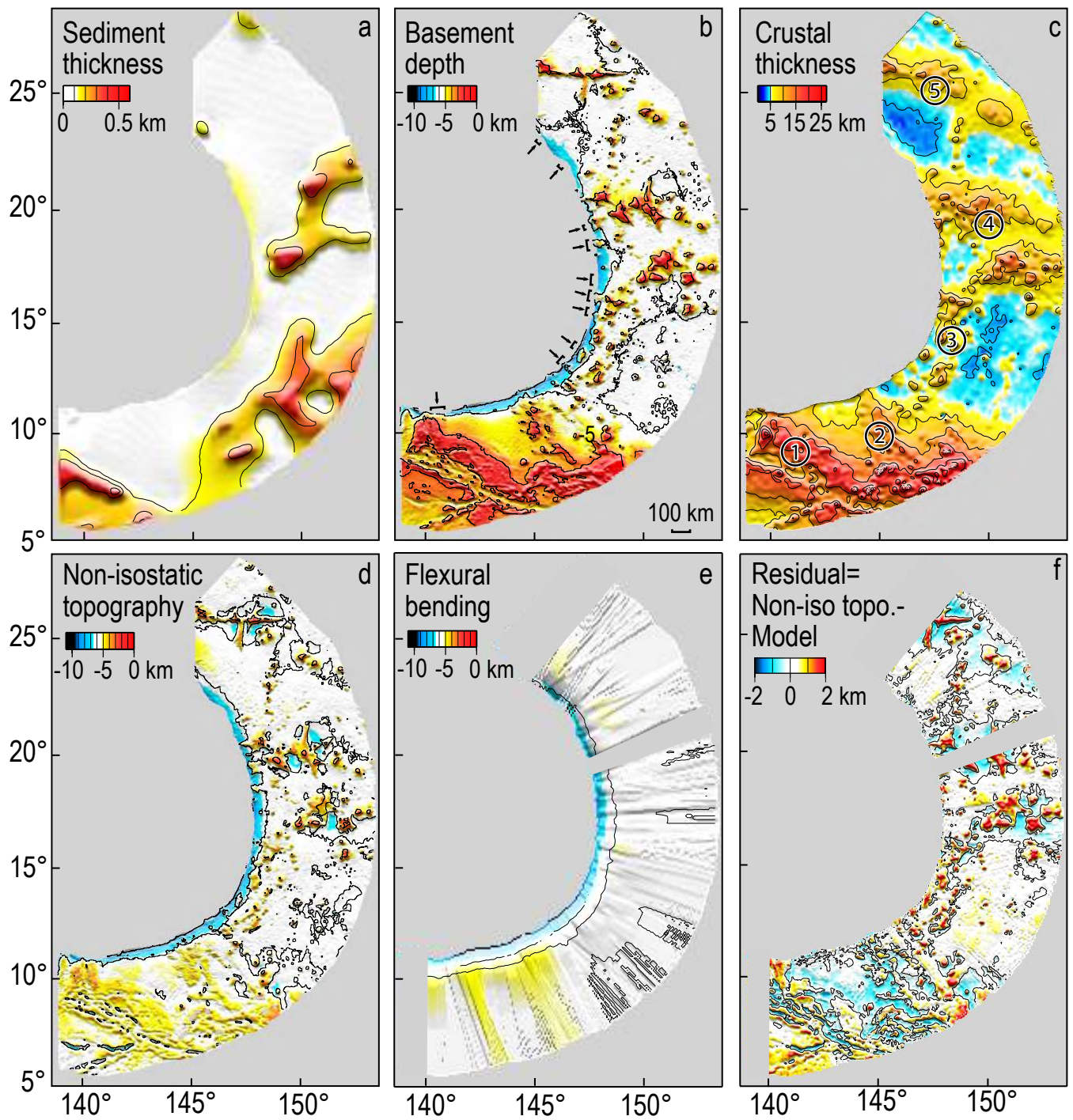


Fig. 3

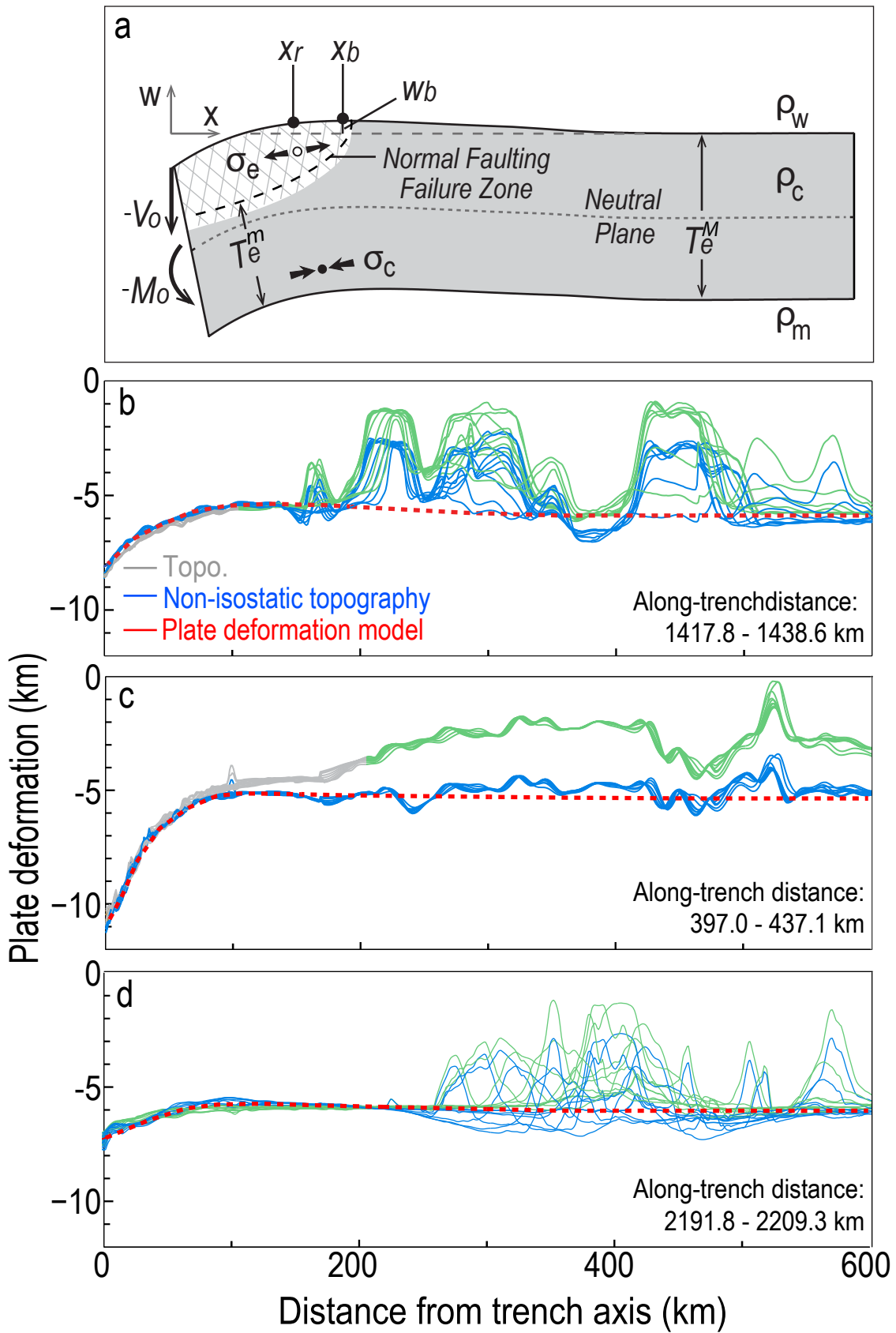


Fig. 4

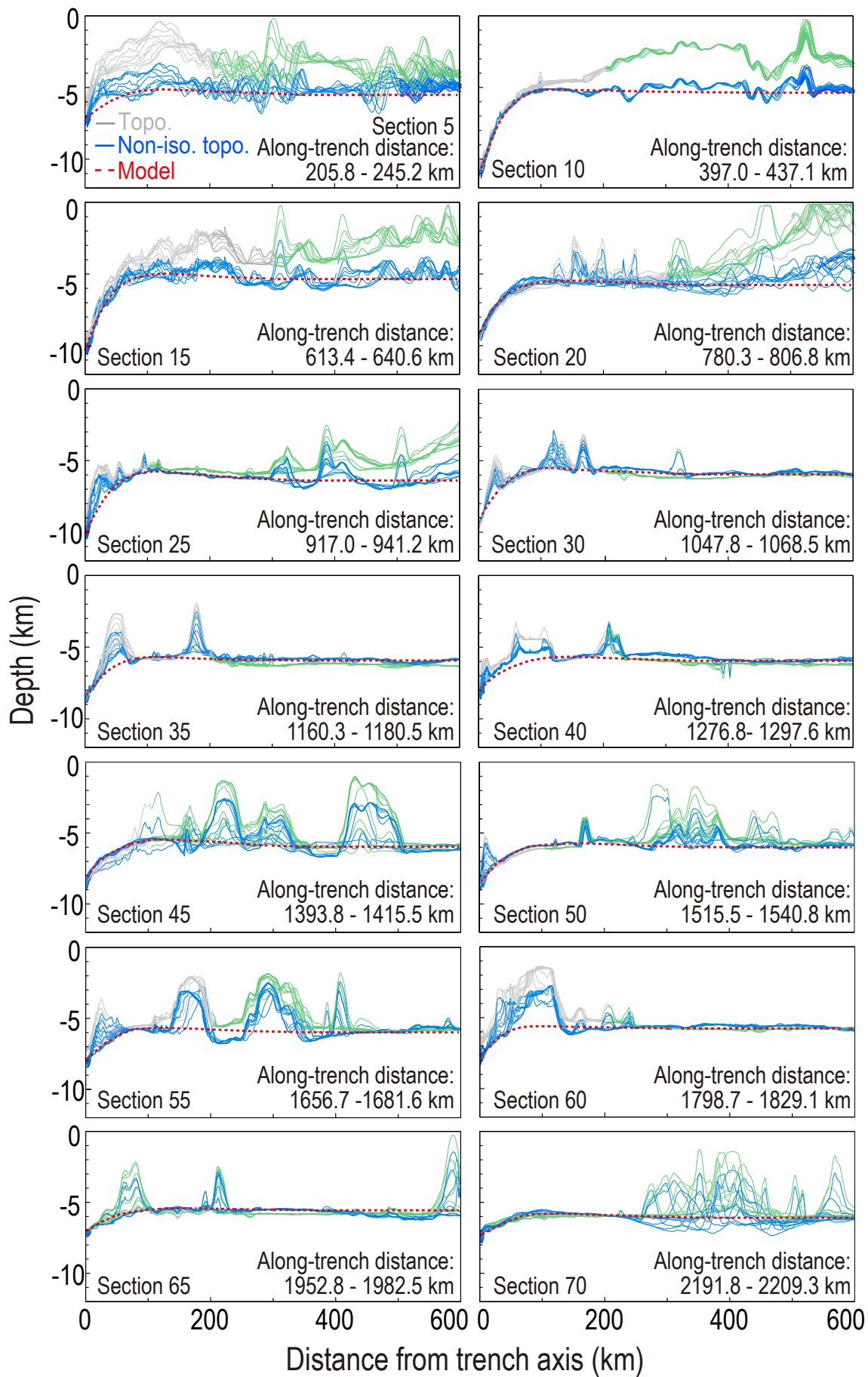


Fig. 5

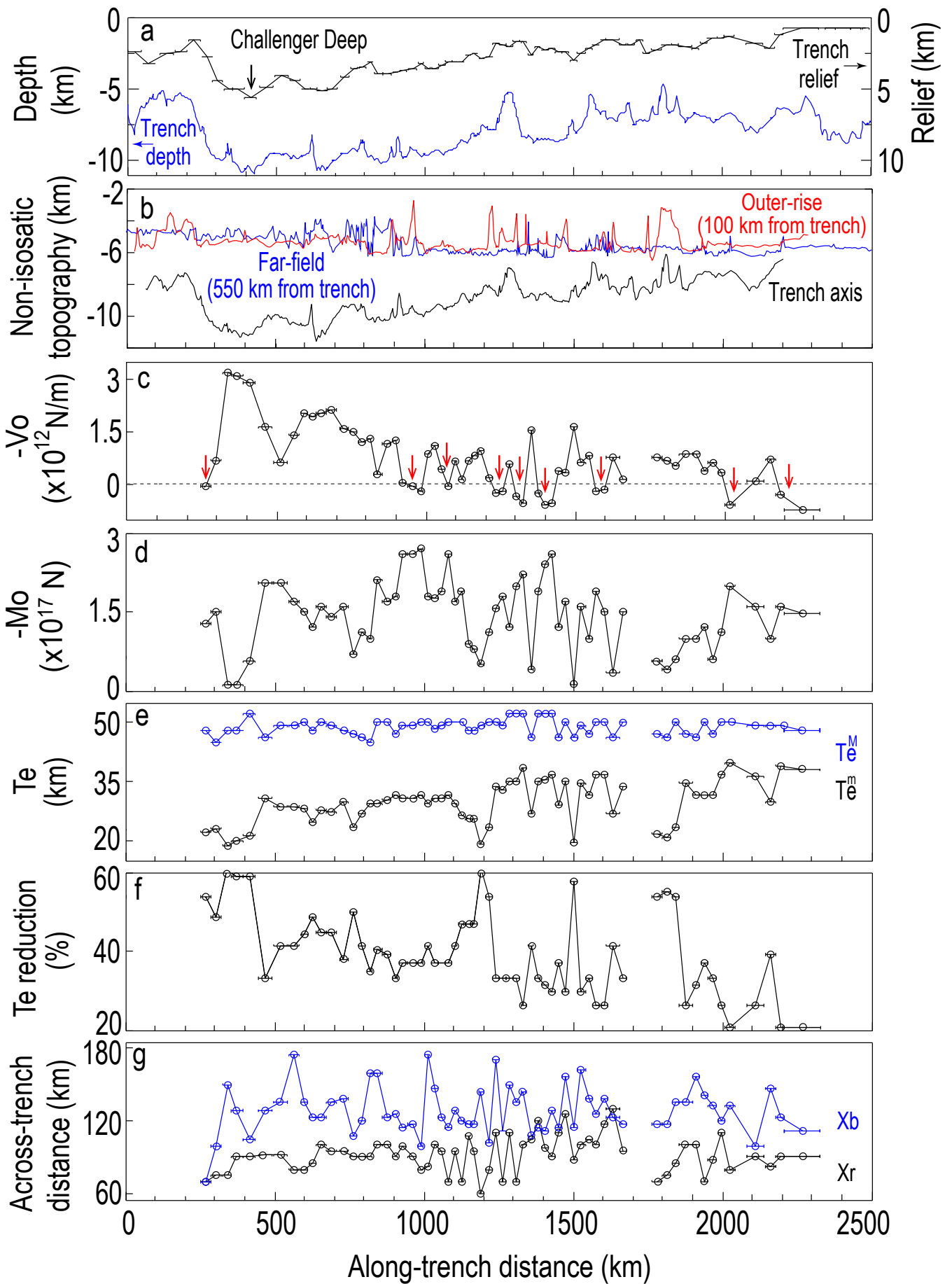


Fig. 6



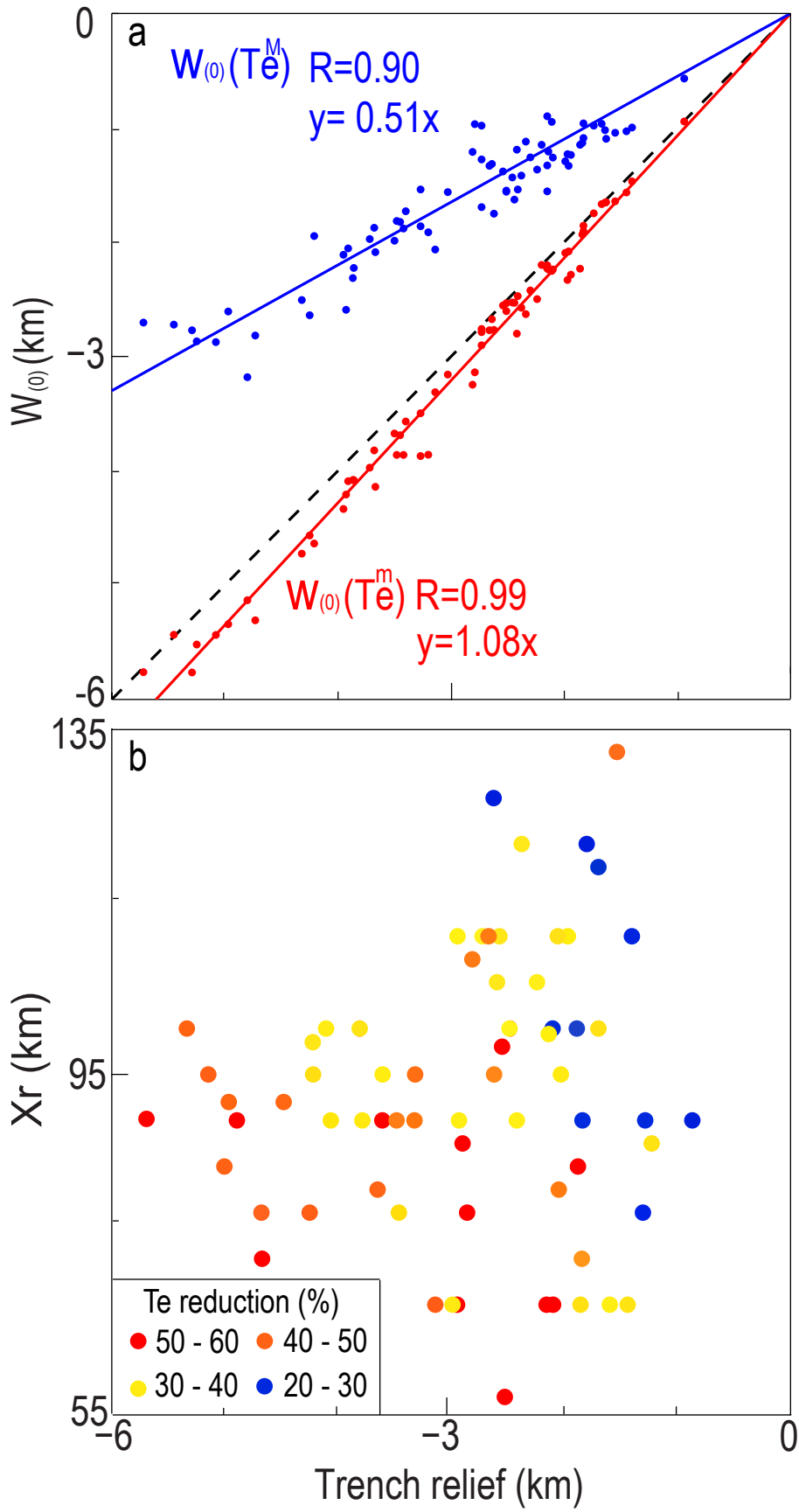


Fig. 7

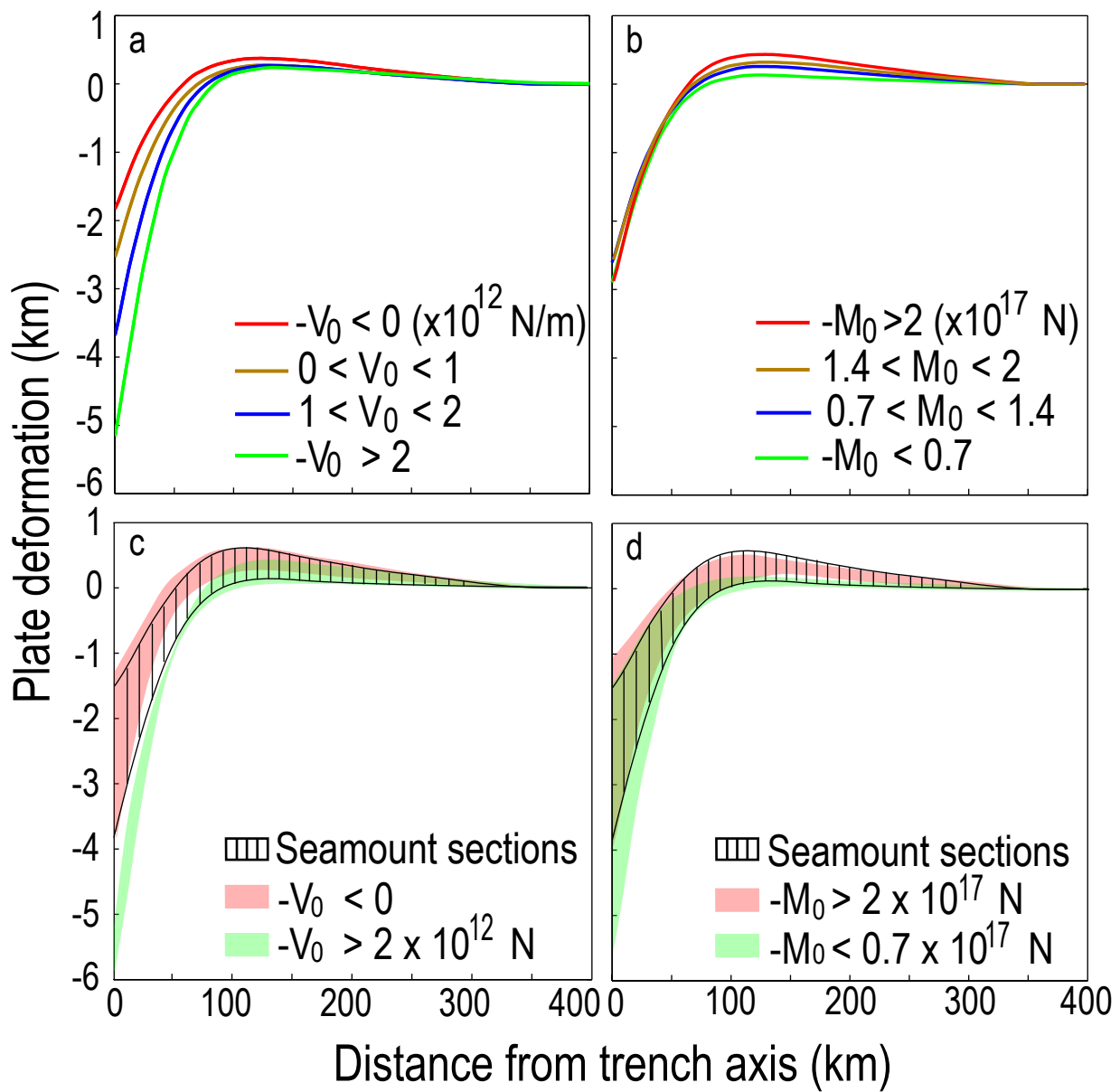


Fig. 8

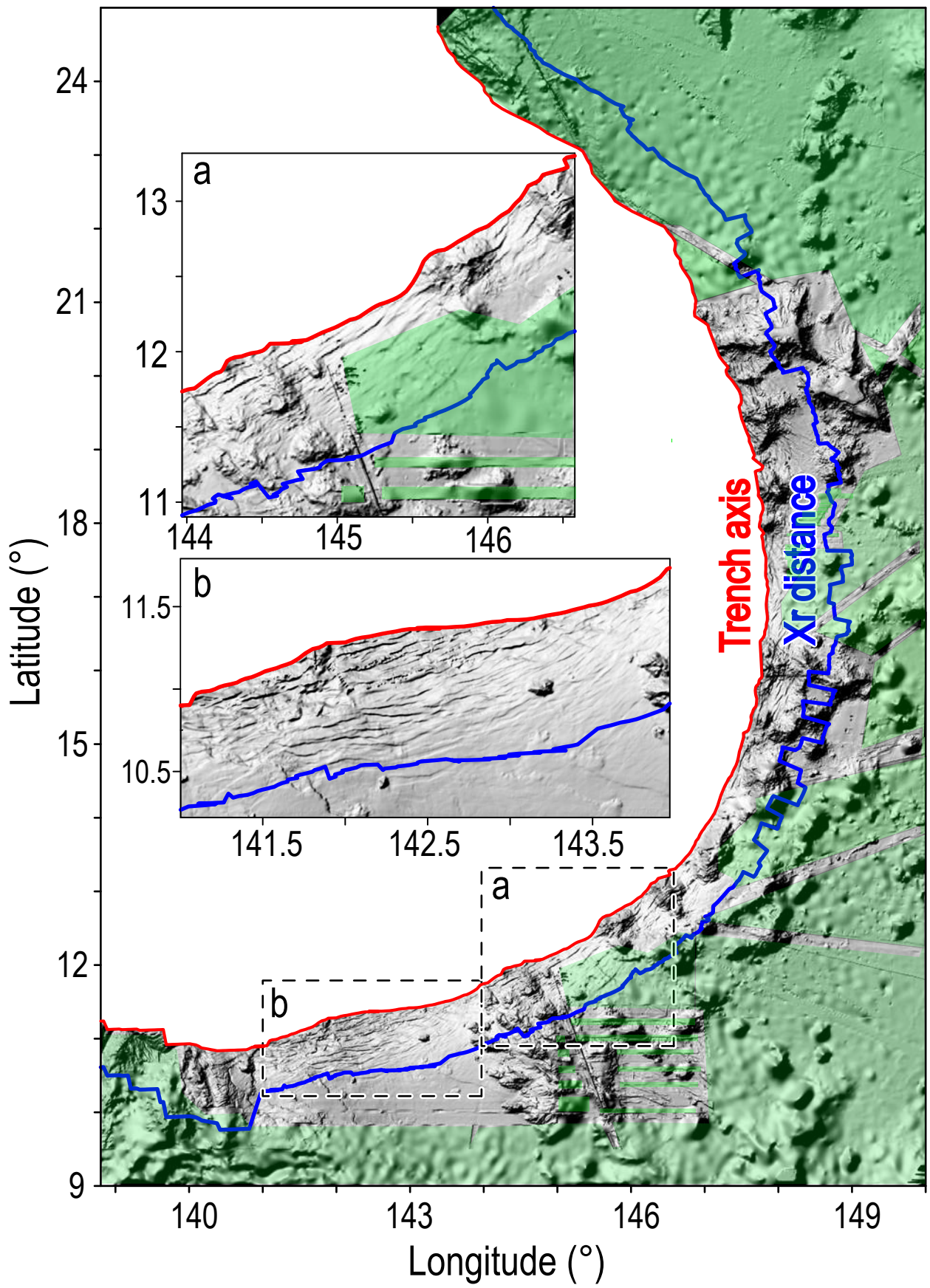


Fig. 9

# Supplementary Materials

**Supplementary Table 1. Best-Fitting Parameters for 75 Profile Sections**

Section	Along-trench distance (km)	Trench relief (km)	$-V_0$ ( $10^{12}$ N/m)	$-M_0$ ( $10^{17}$ N)	$T_e^m$ (km)	$T_e^M$ (km)	$x_r$ (km)	$w_b$ (m)	$T_e$ Reduction (%)
1*	0 - 48.8	-2.38	(0.84)	(0.63)	(22.7)	(49)	(70)	(142)	(53.6)
2*	53.1- 92.9	-3.36	(1.09)	(1.00)	(25.5)	(48)	(90)	(245)	(46.9)
3*	97.2 - 155.6	-2.59	(0.68)	(1.00)	(28.0)	(50)	(110)	(296)	(44.1)
4*	160.0 - 201.5	-2.47	(0.06)	(2.10)	(36.6)	(52)	(110)	(451)	(29.5)
5*	205.8 - 245.2	-1.69	(-0.30)	(1.80)	(38.1)	(54)	(120)	(385)	(29.5)
6	250.0 - 285.1	-2.74	-0.07	1.28	22.3	48	70	330	53.6
7	287.4 - 321.9	-4.16	0.69	1.50	23.1	45	75	417	48.7
8	326.0 - 357.2	-5.23	3.17	0.10	18.8	48	75	70	60.9
9	360.0 - 393.6	-5.40	3.11	0.10	18.8	48	85	141	60.9
10	397.0 - 437.1	-5.67	2.89	0.50	21.1	49	90	221	56.9
11	442.0 - 492.2	-4.76	1.62	2.05	30.8	46	92	311	33.1
12	498.1 - 541.8	-3.92	0.61	2.05	29.2	50	92	481	41.5
13	545.6 - 580.2	-4.28	1.41	1.85	29.2	50	82	350	41.5
14	583.4 - 611.1	-4.68	2.04	1.50	28.1	48	80	290	41.5
15	613.4 - 640.6	-4.94	1.93	1.20	24.6	48	85	352	48.7
16	643.4 - 670.5	-5.20	2.01	1.60	27.7	50	100	413	44.6
17	674.0 - 709.8	-5.05	2.13	1.40	27.1	49	95	368	44.6
18	717.4 - 746.7	-4.21	1.59	1.60	29.8	48	95	383	37.9
19	749.7 - 777.2	-3.63	1.50	0.70	23.5	47	90	232	50.0
20	780.3 - 806.8	-3.44	1.19	1.10	26.9	46	90	287	31.5
21*	809.1 - 829.4	-3.08	1.32	1.00	29.3	45	90	260	35.0
22*	831.9 - 859.9	-3.43	0.27	2.00	29.7	50	100	371	40.6
23	862.9 - 889.3	-3.82	1.14	1.70	30.4	50	100	343	39.2
24	891.8 - 914.7	-3.89	1.27	1.80	31.5	47	90	374	33.1
25*	917.0 - 941.2	-3.85	0.03	2.60	30.9	49	99	650	37.0
26*	945.1 - 975.7	-3.67	-0.07	2.60	30.9	49	90	603	37.0
27*	978.1- 999.9	-3.38	-0.20	2.70	31.5	50	80	528	37.0
28*	1002.3 - 1022.7	-3.64	0.88	1.80	29.2	50	83	287	41.5
29*	1024.9 -1045.5	-3.82	1.10	1.75	30.6	48	100	394	37.0
30*	1047.8 -1068.5	-3.22	0.45	1.90	30.9	49	95	448	37.0
31	1070.8 - 1090.9	-3.15	-0.07	2.60	31.5	50	70	517	37.0
32	1093.1 - 1111.1	-3.40	0.65	1.70	29.2	50	95	373	41.5
33	1115.3 - 1135.7	-3.21	0.12	1.90	26.6	50	70	408	46.9
34*	1138.1 - 1158.1	-2.68	0.65	0.90	25.5	48	108	326	46.9
35*	1160.3 - 1180.5	-2.60	0.81	0.80	25.5	48	95	267	46.9
36	1182.7 - 1202.9	-2.75	0.98	0.50	19.2	49	60	164	60.9
37*	1205.3 - 1227.3	-2.68	0.18	1.10	23.2	50	80	316	53.6
38*	1229.6 - 1250.2	-1.77	-0.27	1.55	33.5	50	110	307	33.1
39*	1252.6 - 1274.3	-1.92	-0.21	1.80	32.8	49	70	314	33.1
40	1276.8 - 1297.6	-2.09	0.59	1.20	34.8	52	110	288	33.1
41*	1300.0 - 1321.5	-1.80	-0.35	2.00	34.8	52	70	385	33.1
42*	1323.9 - 1345.5	-1.79	-0.55	2.20	38.3	52	100	366	26.3
43	1347.8 - 1368.7	-2.67	1.54	0.40	26.9	46	105	131	41.5
44*	1371.1 - 1391.4	-2.14	-0.26	1.90	34.8	52	120	450	33.1
45	1393.8 - 1415.5	-2.24	-0.57	2.40	35.6	52	98	514	31.6
46*	1417.8 - 1438.6	-2.33	-0.55	2.60	36.6	52	90	509	29.5
47	1441.3 - 1463.6	-2.48	0.40	1.20	29.0	46	110	421	37.0
48*	1465.9 - 1487.3	-2.46	0.33	1.70	35.2	50	125	368	29.5
49	1490.0 - 1512.8	-2.99	1.66	0.13	19.9	47	90	82	57.6
50*	1515.5 - 1540.8	-2.58	0.60	1.60	34.5	49	100	303	29.5
51	1543.0 - 1563.4	-2.35	0.82	1.00	31.5	47	105	233	33.1
52*	1565.8 - 1590.8	-1.95	-0.18	1.90	36.8	50	100	328	26.3
53*	1599.3 - 1610.0	-1.60	-0.14	1.50	36.8	50	117	332	26.3
54	1618.5 - 1654.4	-1.62	0.76	0.50	30.6	50	110	186	38.7
55*	1656.7 - 1681.6	-2.05	0.15	1.50	33.5	50	95	325	33.1
56*	1683.9 - 1709.2	-2.09	(-0.10)	(1.70)	(33.5)	(50)	(100)	(384)	(33.1)
57*	1712.1 - 1736.0	-1.79	(0.22)	(1.20)	(33.5)	(46)	(100)	(280)	(33.1)
58*	1738.2 - 1760.4	-2.40	(1.05)	(0.60)	(26.9)	(46)	(100)	(179)	(31.5)

59*	1763.6 - 1796.3	-2.28	0.77	0.55	21.8	47	70	164	53.6
60*	1798.7 - 1829.1	-2.05	0.69	0.40	20.6	46	75	164	55.2
61*	1832.0 - 1857.7	-2.10	0.54	0.60	23.2	50	85	160	53.6
62	1860.5 - 1899.9	-2.10	0.84	1.00	34.6	46	100	242	31.2
63	1902.3 - 1923.1	-2.38	0.88	1.00	31.6	46	100	265	31.2
64*	1925.7 - 1949.9	-1.87	0.40	1.20	31.5	50	70	246	37.0
65*	1952.8 - 1982.5	-1.49	0.63	0.60	31.5	47	87	140	33.1
66	1984.9 - 2010.9	-1.57	0.32	1.10	36.8	50	110	234	26.3
67	2013.2 - 2044.7	-1.39	-0.59	2.00	39.7	50	80	349	20.6
68	2061.7 - 2138.6	-1.91	0.10	1.60	36.1	49	90	276	26.3
69	2142.6 - 2179.2	-2.20	0.70	1.00	29.8	49	83	176	39.2
70	2191.8 - 2209.3	-1.53	-0.30	1.60	36.1	49	90	306	26.3
71	2212.5 - 2328.9	-0.87	-0.73	1.48	38.1	48	90	326	20.6
72	2341.7 - 2382.3	-	-	-	-	-	-	-	-
73	2384.6 - 2432.5	-	-	-	-	-	-	-	-
74	2436.9 - 2469.5	-	-	-	-	-	-	-	-
75	2472.5 - 2495.5	-	-	-	-	-	-	-	-

Notes:

- The star symbol (\*) marks sections with seamounts near the trench axis.
- Parameters bracketed with parentheses, i.e., Sections 1-5 and 56-58, are relatively poorly constrained due to significant effects of seamounts near the profile sections.
- For Sections 72-75, sediment thickness data are not available and thus several parameters were not calculated.

## Appendix A. Topography and Flexural Bending of 75 Profile Sections of the Mariana Trench

We analyzed flexural bending along 75 across-trench sections, each consists of ten profiles spanning over an along-trench distance of about  $0.2^\circ$  (Figs. S1-8). For each profile section, we calculated a flexural bending model that best matches non-isostatic topography (blue curves) of areas away from seamounts.

### Figure Captions

**Figure S1-8.** A total of 75 across-trench profile sections of the Mariana trench. Basement topography of every ten individual profiles was stacked to form an averaged section for modeling: grey and green curves show areas with and without multi-beam bathymetry, respectively. Blue curves are the calculated non-isostatic topography. Red dashed curves show flexural bending models that best fit non-isostatic topography away from seamounts.

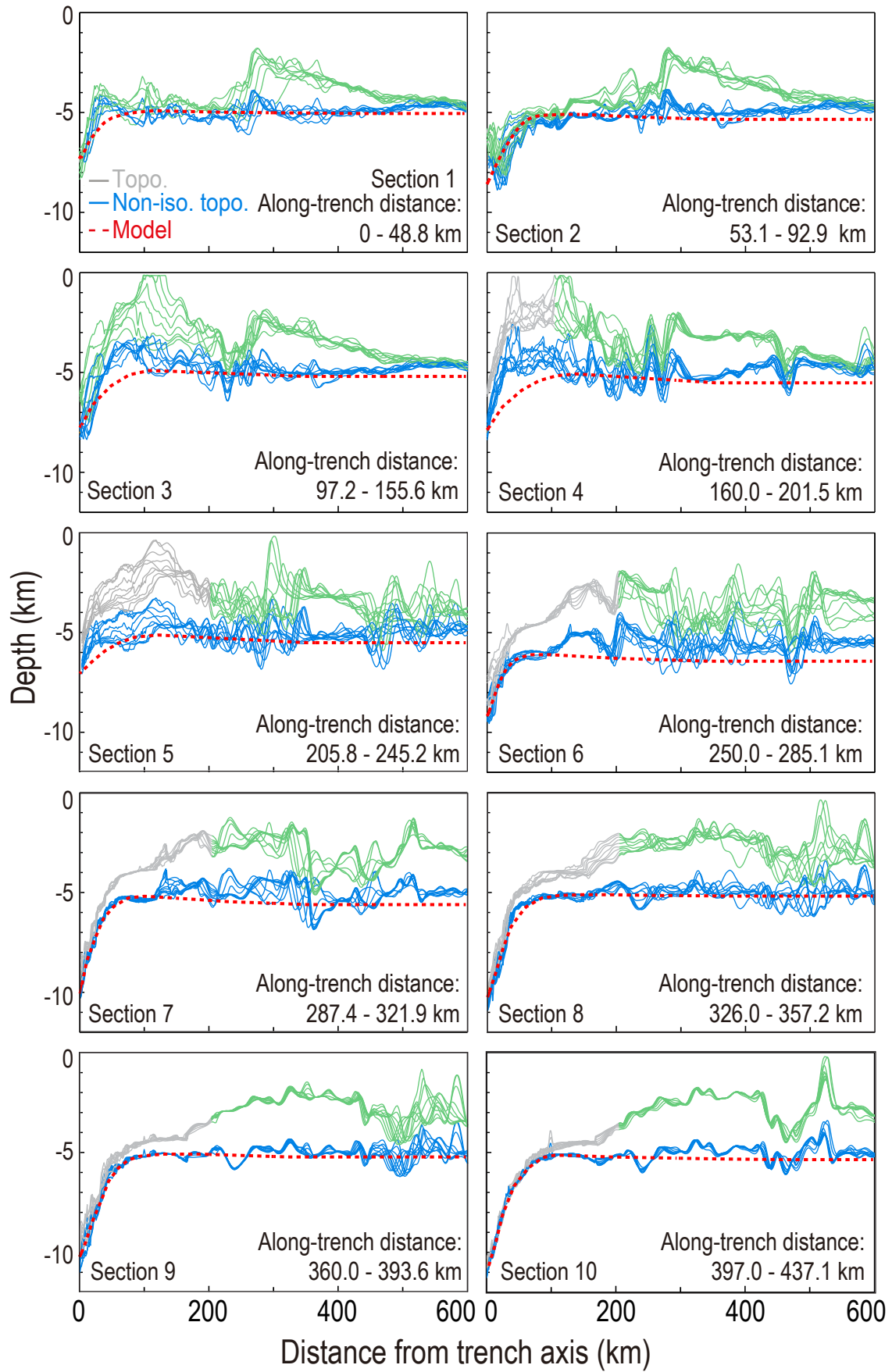


Fig. S1

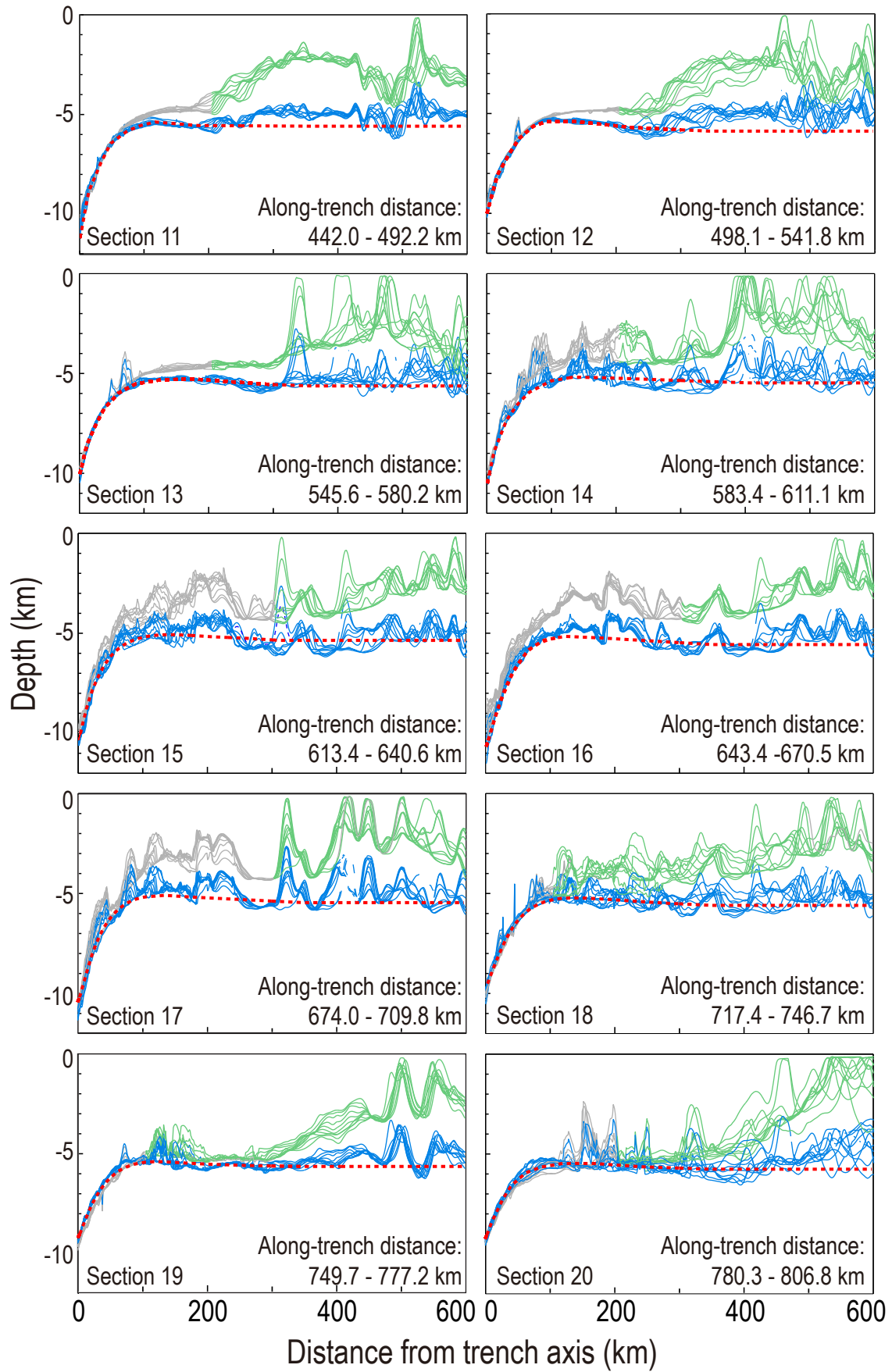


Fig. S2

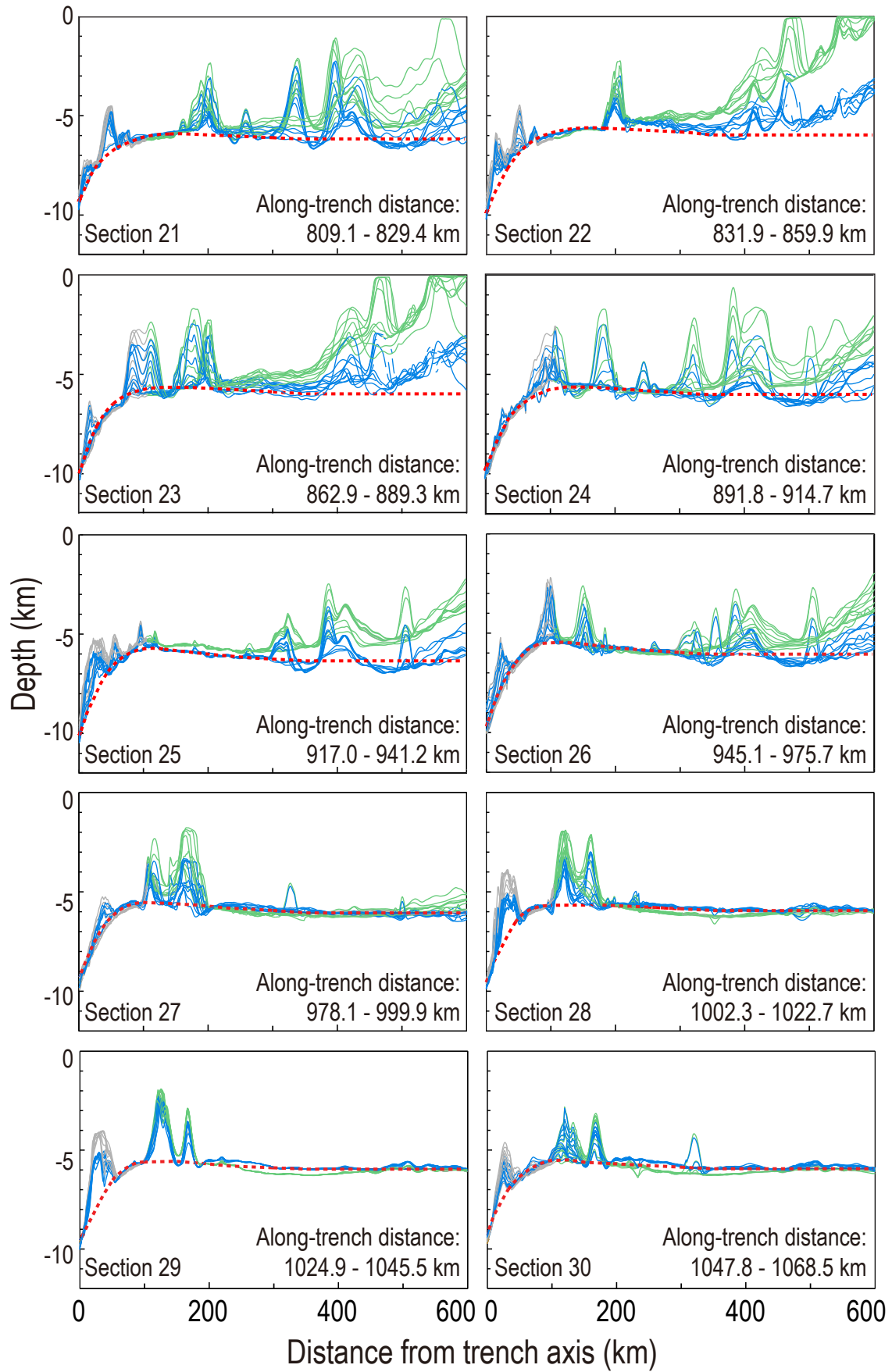


Fig. S3



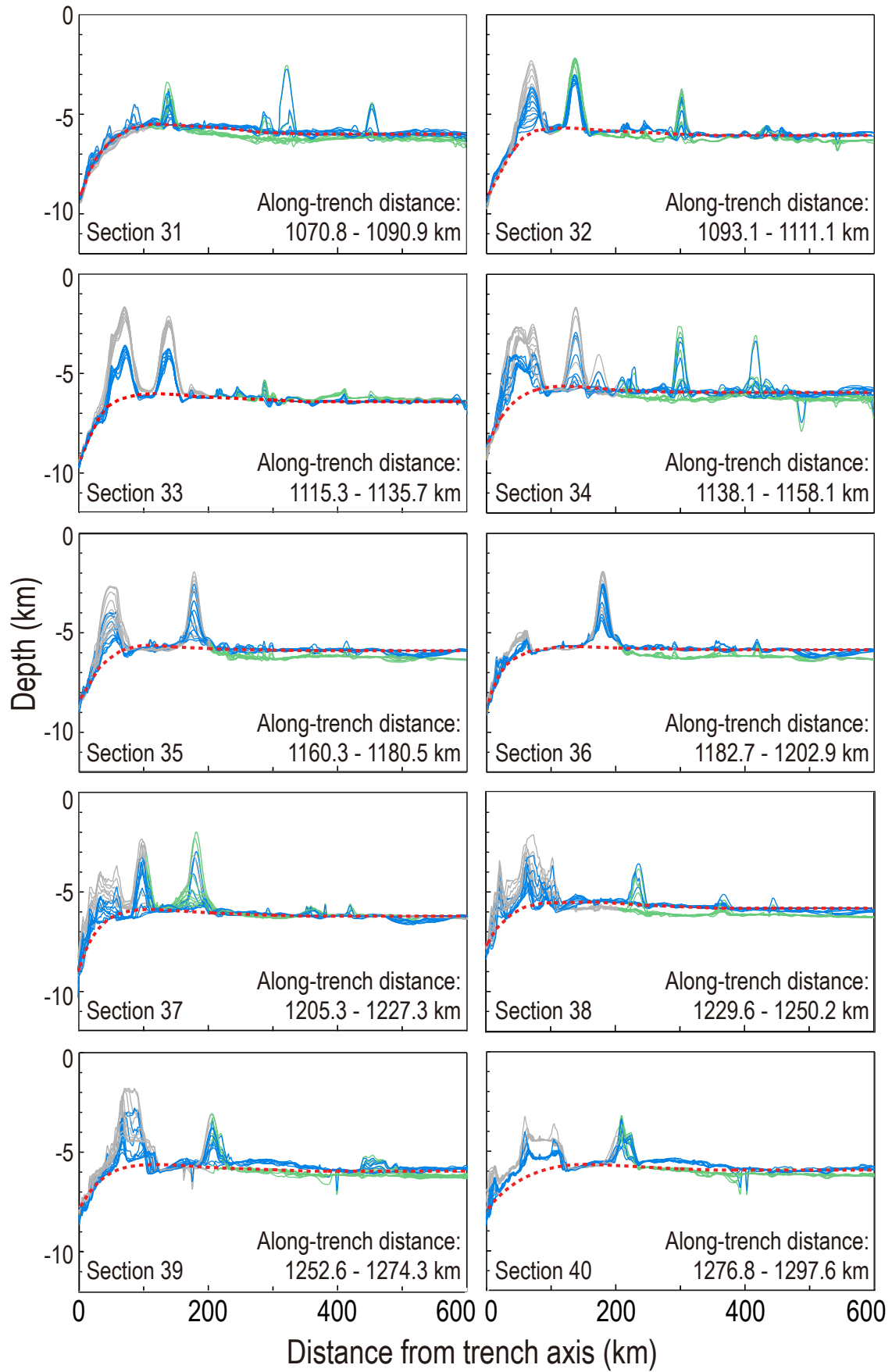


Fig. S4

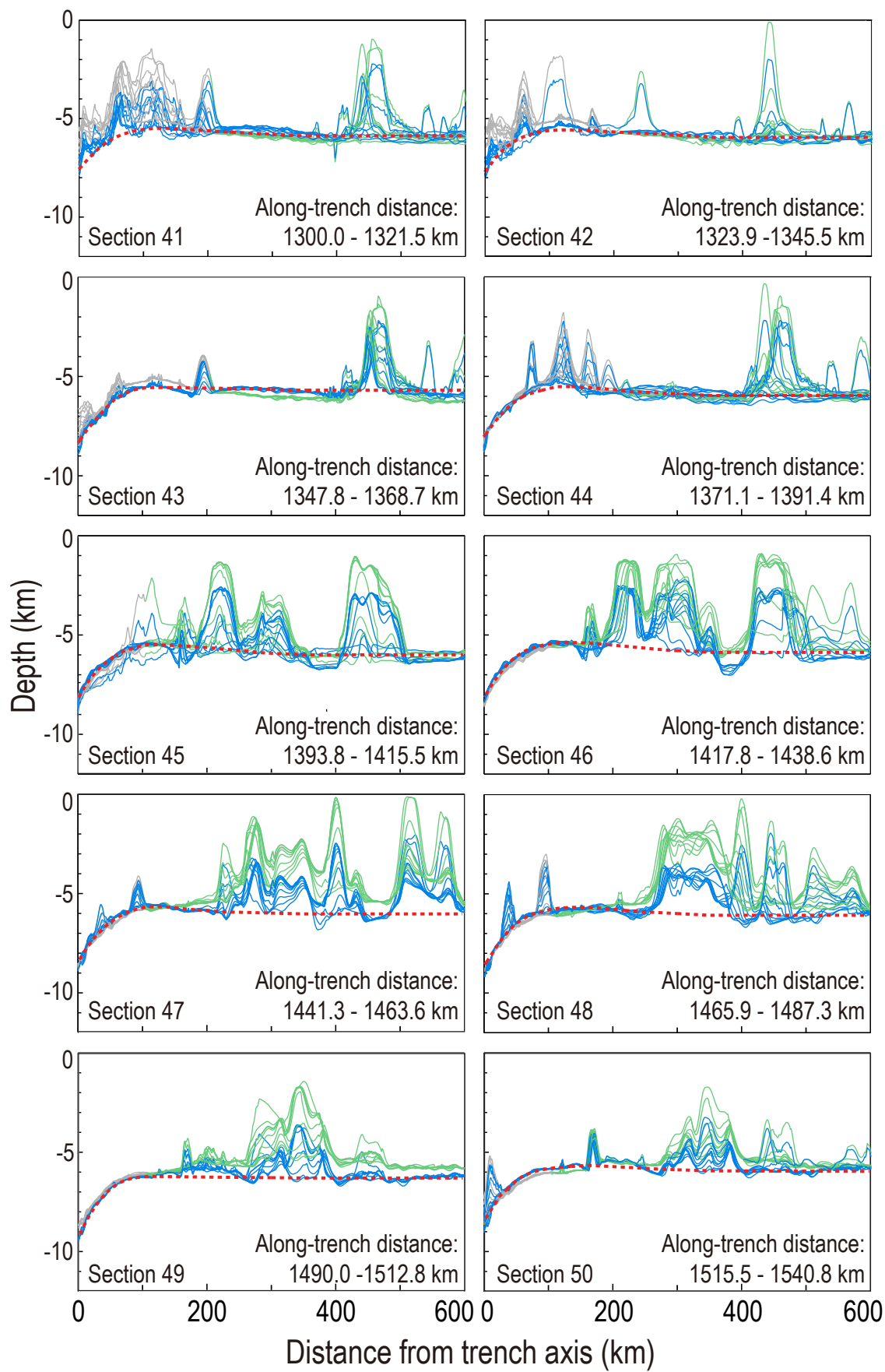


Fig. S5

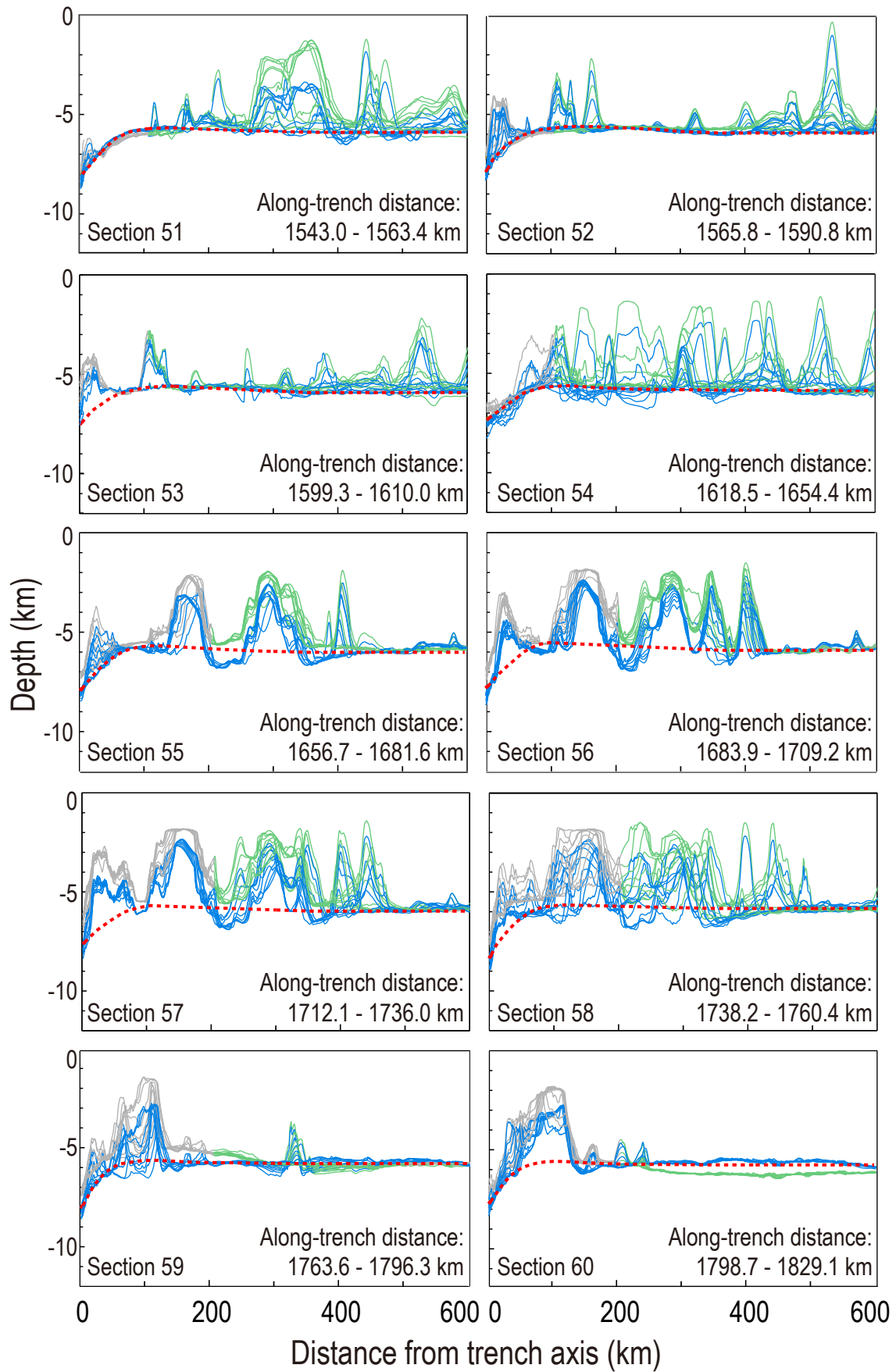


Fig. S6

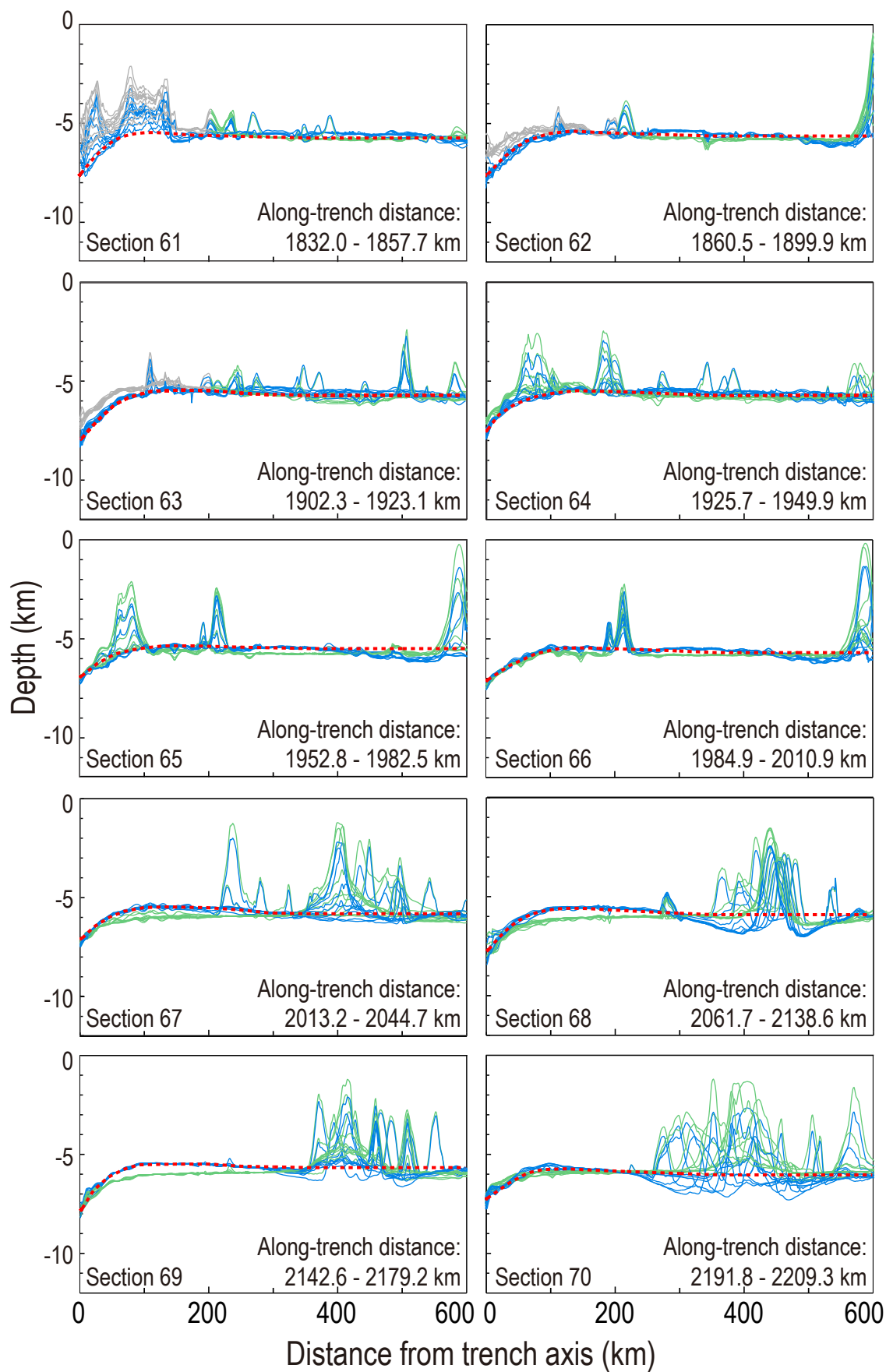


Fig. S7

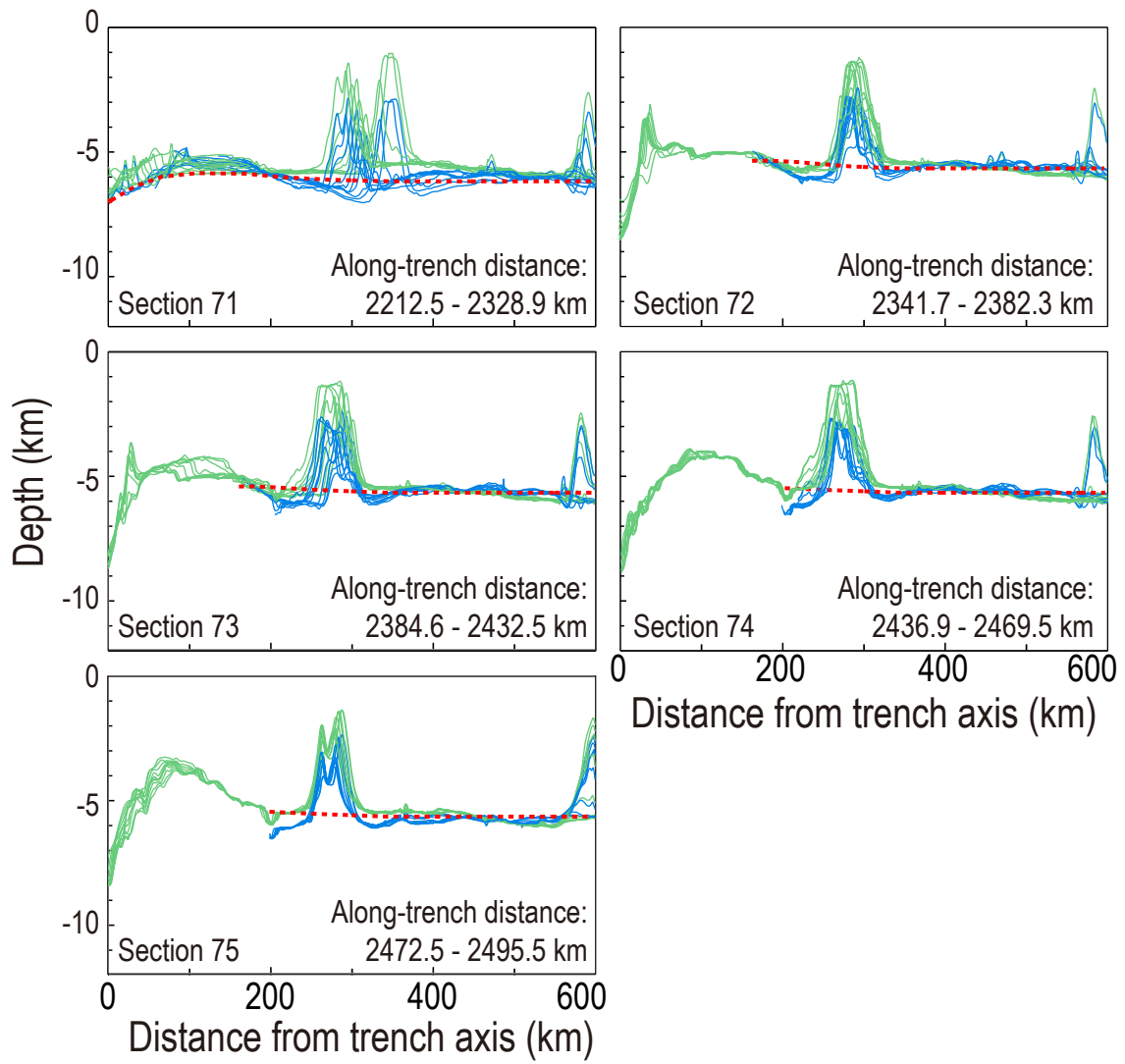


Fig. S8

## Appendix B. Numerical Solutions of Flexural Bending of a Plate with Variable Elastic Thickness

The flexural bending of a thin elastic plate is described by the following equation (Turcotte and Schubert, 2002):

$$-\frac{d^2M}{dx^2} + \frac{d}{dx}\left(F\frac{dw}{dx}\right) + \Delta\rho gw = q, \quad (\text{S1})$$

where  $w(x)$  is vertical deflection of the plate,  $M(x)$  is bending moment,  $F(x)$  is horizontal buckling force,  $\Delta\rho gw = (\rho_m - \rho_w)gw$  represents hydrostatic restoring force,  $q(x) = (\rho_s - \rho_w)gh_s(x)$  is vertical sediment loading,  $h_s(x)$  is sediment thickness, and  $\rho_w$ ,  $\rho_s$ , and  $\rho_m$  are densities of water, sediment, and mantle, respectively.

Assuming that the spatial variation in the buckling force,  $\frac{dF(x)}{dx}$ , is relatively small,

Eqn. S1 can be re-written as the following set of second-order differential equations (Contreras-Reyes and Osses, 2010):

$$\frac{d^2w}{dx^2} = -\frac{M}{D}; \quad (\text{S2})$$

$$\frac{d^2M}{dx^2} = -F\frac{M}{D} + \Delta\rho gw - q, \quad (\text{S3})$$

where flexural rigidity is  $D(x) = \frac{ET_e(x)^3}{12(1-\nu^2)}$ , while  $T_e$ ,  $E$ , and  $\nu$  are the effective elastic

thickness of the plate, Young's modulus, and Poisson's ratio, respectively.

To solve for the four variables,  $w$ ,  $\frac{dw}{dx}$ ,  $M$ , and  $\frac{dM}{dx}$ , the above Eqns. S2 and S3 can be

further re-written as a set of first-order differential equations in a matrix form:

$$\frac{d}{dx} \begin{pmatrix} w \\ \frac{dw}{dx} \\ M \\ \frac{dM}{dx} \end{pmatrix} + \begin{pmatrix} 0 & -1 & 0 & 0 \\ 0 & 0 & \frac{1}{D(x)} & 0 \\ 0 & 0 & 0 & -1 \\ -\Delta\rho g & 0 & \frac{F(x)}{D(x)} & 0 \end{pmatrix} \times \begin{pmatrix} w \\ \frac{dw}{dx} \\ M \\ \frac{dM}{dx} \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ -q(x) \end{pmatrix} \quad (\text{S4})$$

The above Eqn. S4 is then solved using the finite-difference method.

In finite-difference operation, the  $x$  domain of length  $L$  is discretized into  $N$  grid points of equal distance, i.e.,  $x(i) = (i-1)\Delta x$ , for  $i=1:N$ , where  $\Delta x = L/(N-1)$ . For grid points away from the boundary, i.e., for  $i=2:N-1$ , the four rows of the first-order differential equations in Eqn. S4 are described by the following:

$$\frac{dw(i)}{dx} = \frac{w(i+1) - w(i-1)}{2\Delta x}; \quad (\text{S5})$$

$$\frac{dw(i+1)/dx - dw(i-1)/dx}{dx} + \frac{M(i)}{D(i)} = 0; \quad (\text{S6})$$

$$\frac{dM(i)}{dx} = \frac{M(i+1) - M(i-1)}{2\Delta x}; \quad (\text{S7})$$

$$\frac{dM(i+1)/dx - dM(i-1)/dx}{2\Delta x} - \Delta\rho g w(i) + \frac{F(i)M(i)}{D(i)} = -q(i). \quad (\text{S8})$$

When applied to the grid points of  $i=2:N-1$ , Eqns. S5-S8 yield a total of  $4N-8$  constraints.

The boundary conditions at  $x=0$  and  $x=L$  are prescribed as the following:

$$M=M_0, \quad \frac{d^2 M}{dx^2} = V_0 + F_0 \frac{dw}{dx}, \quad \text{at } x=0, \quad (\text{S9})$$

$$w=0, \quad \frac{dw}{dx} = 0, \quad \text{at } x=L, \quad (\text{S10})$$

where  $V_0$ ,  $M_0$ , and  $F_0$  are, respectively, the given vertical force, bending moment, and horizontal force at  $x=0$ .

The boundary conditions for Eqn. S4 can be further expressed in 8 sets of finite-difference operations:

$$\frac{dw(1)}{dx} = \frac{w(2) - w(1)}{\Delta x}; \quad (\text{S11})$$

$$w(N)=0; \quad (\text{S12})$$

$$\frac{dw(2)/dx - dw(1)/dx}{\Delta x} + \frac{M_0}{D_0} = 0; \quad (\text{S13})$$

$$\frac{-dw(N-1)/dx}{dx} + \frac{M(N)}{D(N)} = 0; \quad (\text{S14})$$

$$\frac{dM(1)}{dx} = \frac{M(2) - M_0}{\Delta x}; \quad (\text{S15})$$

$$\frac{dM(N)}{dx} = \frac{M(N) - M(N-1)}{\Delta x}; \quad (\text{S16})$$

$$\frac{dM(2)/dx - dM(1)/dx}{\Delta x} - \Delta\rho gw(1) + \frac{F_0 M_0}{D_0} = -q(1); \quad (\text{S17})$$

$$\frac{dM(N)/dx - dM(N-1)/dx}{\Delta x} - \Delta\rho gw(N) + \frac{F(N)M(N)}{D(N)} = -q(N); \quad (\text{S18})$$

Eqns. S11-S18 yield a total of 8 constraints. In sum, Eqns. S5-S8 together with Eqns.

S11-S18 provide a total of 4N constraints, which are used through matrix inversion to

solve for the 4N unknowns of  $w(i)$ ,  $\frac{dw(i)}{dx}$ ,  $M(i)$ , and  $\frac{dM(i)}{dx}$ , for  $i=1:N$ .

### Appendix C. Gravity-Derived Crustal Thickness

We used gravity-derived crustal thickness (Fig. 3c) to calculate isostatic topography.

The crustal thickness for the study region was calculated by following steps:

(1) *Data and Analysis*. Bathymetric data were obtained from the National Geophysical Data Center (NGDC, <http://www.ngdc.noaa.gov>) and GEBCO (<http://www.gebco.net>); free-air gravity anomaly (FAA) data were from Sandwell and Smith (2009); sediment thickness data were from the NGDC database (Divins, 2003); and the crustal age data were from Müller (2008).



(2) *Thermal Correction.* The gravity effects due to age-dependent cooling of the oceanic lithosphere were calculated from a 1-D plate cooling model, assuming the top and bottom temperatures of a 100-km-thick mantle layer with  $T_s=0^\circ\text{C}$  and  $T_m=1,350^\circ\text{C}$ , respectively (Turcotte and Schubert, 2002). The calculated thermal structure was then converted into a 3-D mantle density grid, in which the density anomaly  $\Delta\rho = -\alpha\rho_0(T - T_0)$ , where  $T$  is mantle temperature,  $T_0 = 1,350^\circ\text{C}$  and  $\rho_0 = 3.3*10^3 \text{ kg/m}^3$  are reference temperature and density, respectively, and  $\alpha = 3*10^{-5} \text{ }^\circ\text{C}^{-1}$  is the volumetric coefficient of thermal expansion.

(3) *Residual Mantle Bouguer Anomaly.* Assuming a reference crustal thickness of 6 km, the gravitational effects of the water/crust and crust/mantle density interfaces were removed from the FAA using the Parker spectrum method (Parker, 1973) to obtain the mantle Bouguer anomaly (MBA). The gravitational effects of lithospheric cooling, as calculated in the above Step 2, were then removed from the MBA to obtain residual mantle Bouguer anomaly (RMBA).

(4) *Crustal Thickness.* The crustal thickness was calculated from downward continuation of the RMBA signal to a reference depth using the methods of Parker (1973) and Kuo and Forsyth (1988). The best-fitting parameters of mantle and crustal densities used in the calculations were obtained from calibration of the gravity-derived models with constraints from available seismic refraction profiles.