

1 Analysis and real-time prediction of the full-scale thrust for floating 2 wind turbine based on artificial intelligence

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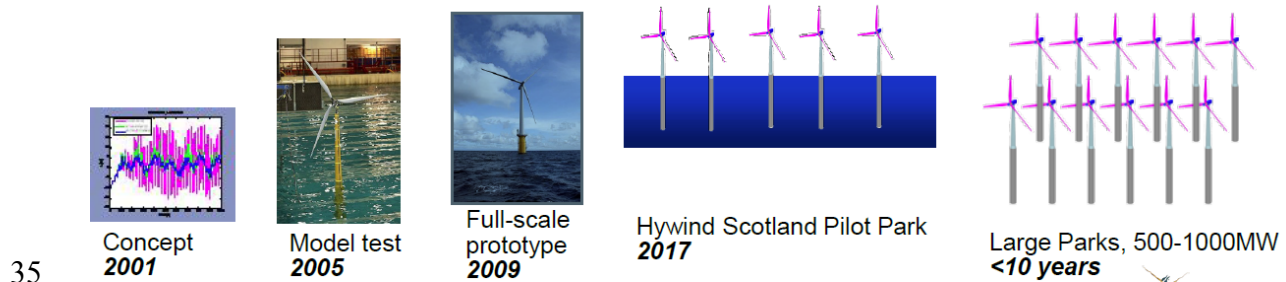
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7 Abstract: In this paper, numerous aero-hydro-servo-elastic coupled simulations are carried out
8 in time-domain to observe the performance of the real-time thrust acting on the rotor of the
9 OC3-Hywind offshore floating wind turbine. And the studying focuses on investigating the
10 correlation between inputs (surge motion, pitch motion, wind conditions, etc.) and the targeted
11 output (rotor thrust) in the time domain. Besides, artificial intelligence (AI) techniques are used
12 to estimate a prediction model of real-time thrust based on the data from simulations. To predict
13 the thrust, data for four comparative coupled environmental conditions are considered, by
14 which the effect of turbulence and wave spectrum on the thrust force is also investigated.
15 Moreover, a series of simulations of frequency-increasing regular wave conditions and speed-
16 increasing wind conditions are carried out to observe their effect on the real-time rotor thrust.
17 Additionally, the impact of the pitch and surge RAOs of the floating foundation and the wind
18 velocity are quantitatively studied. It reveals that the high-frequency response of thrust is
19 dominated by wave change, whereas low-frequency response is dominated by wind change.
20 Besides, one simulation model of the thrust acting on the rotor is estimated regarding high-
21 frequency and low-frequency response separately to account the dominating influence.

22 Keywords: Fully-coupled simulation; OC3-Hywind; real-time thrust prediction; Artificial
23 neural network; impact analyses on the real-time thrust

24 **1. Introduction**

25 During recent few years, the development of grid-connected floating wind turbines has
26 flourished. It is claimed that the more into the sea, there is more wind energy, since the increase
27 of wind velocity[1]. It is exciting that we have seen a few social achievements, especially the
28 deployment of the Hywind floating wind farm in Scotland in 2017, which is the first floating
29 wind farm in the world[2]. Indeed, this spar-type floating wind turbine has experienced
30 sufficient research as shown in Fig. 1. [3] deals with the modelling and parameter tuning OC3-
31 Hywind floating wind turbine with a tuned mass damper (TMD) installed in the nacelle. The
32 seakeeping performance of the OC3-Hywind platform is investigated in [4] by both numerical
33 solution and model test method. Moreover, a series of model test researches have been
34 performed at the same time[5-7].



36 **Fig. 1 The development of Hywind Scotland**

37 Different from fixed-bottom wind turbines, the floating platform naturally responds to wave
38 loadings which are Froude-scaled, while the turbine forces react to aerodynamic loads which
39 are Reynolds-scaled. Commonly Froude scaling is used in tank tests. However, using Froude
40 scaling also for the wind turbine rotor will lead to wrong aerodynamic loads compared to the
41 full-scale turbine. It is essential to model the aerodynamic loads for FOWTs correctly. One
42 possible method of shaping a floating wind turbine in the laboratory environment is to employ
43 a working rotor in a wind field generated by a series of fans. Experimental studies with a wind
44 field generated by a set of fans in a tank are usually used to examine the hydrodynamic or

45 aerodynamic performance of a floating wind model. For example, Martin et al. tried to account
46 for Reynolds effects by appropriate distortion of the rotor geometry[8], however, it is incredibly
47 challenging to construct and operate a working scale rotor due to the large size, not to mention
48 the very lightweight, and sophisticated control requirements, see references[9-12]. Moreover,
49 examples are given by Chujo et al. for a spar OWT in [13], by K Muller et al. for a semi-
50 submersible OWT in [14], and by Goupee et al. for the spar, semi-submersible and TLP types
51 in [15]. The experiments study are usually focused on a scaled model in the region of 1/50 –
52 1/100. However, both Froude similitude and Reynolds similitude cannot be satisfied
53 simultaneously in the small-scale model tests of ships and offshore structures.

54 And some common difficulties are concluded, such as lack of wind generation in a
55 hydrodynamic laboratory, issues of simulation of spread moorings, or issues of the similarity
56 between the model and full-scale prototype in [16]. Notably, the simulation of wind loading
57 for tank testing of floating wind turbines presents a variety of severe challenges, particularly
58 challenges concerning the wind generation. Therefore, many people turn to numerical tools,
59 for example, Liu et al.[17] develop a CFD simulation tool for the fully coupled model of
60 floating wind turbines. Godine Kok et al.[18] develop a hydrodynamic computer module to
61 evaluate the linear and nonlinear loads on floating wind turbines using a new fluid-impulse
62 formulation for coupling with the FAST program, for more details see reference[19]. However,
63 physical experiments are still needed, in spite of the difficulty in correctly modelling the
64 aerodynamic loads with the proper hydrodynamic influence taken into account.

65 A few recent articles review the approaches of simulating wind loading on floating wind
66 turbines in wave tanks. For example, Fowler et al. use a high-speed fan located on the model
67 in line with the rotor drivetrain to generate the thrust component of the wind load [20]. Bayati
68 et al. proposed a six-degree-of-freedom platform to simulate the motions of the platform of the
69 floating wind turbine[21]. Additionally, the six-degree-of-freedom platform is used to carry

70 out tests to measure the movements and the aero-dynamic thrust related to the instantaneous
71 position. Moreover, the velocity of the platform is calibrated in real time using a modified
72 version of the FAST aero-hydro-servo-elastic simulation code, which proves it possible to use
73 the thrust calculated from FAST to generate the equivalent thrust in the model test. More before,
74 Chujo et al. attach a lightweight line at the rotor hub and tension[22]. However, that approach
75 can only be justified for rough preliminary estimation of the maximum mooring offset.
76 Kraskowski et al.[23] use a fan rotating at a constant speed in stationary air generating a
77 predefined thrust to represent the mean rotor thrust, which and allows for easy adjustment of
78 the mean wind load and simulation of mean mooring offsets. Therefore, inspired by those
79 researches, the idea of “software in the loop” simulation(SIL) approach is formed.

80 Motivated by the problems mentioned above and research potentials, the authors aim at
81 proposing a new approach to utilise the idea of “software-in-the-loop” in which we use an
82 active control system inserted with AI model to handle the inputs from experiments and give
83 out a proper thrust command to drive a fan in real time in Fig. 2. By thus to enable the accurate
84 tank test of a floating wind turbine and overcomes the various limitations labs are facing.

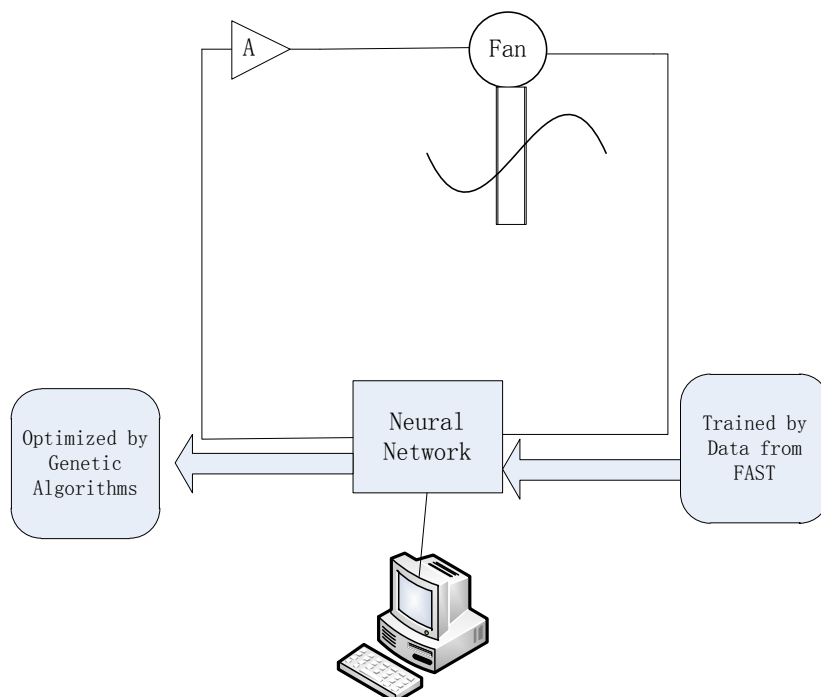


Fig. 2 The diagram of the software-in-the-loop application concept

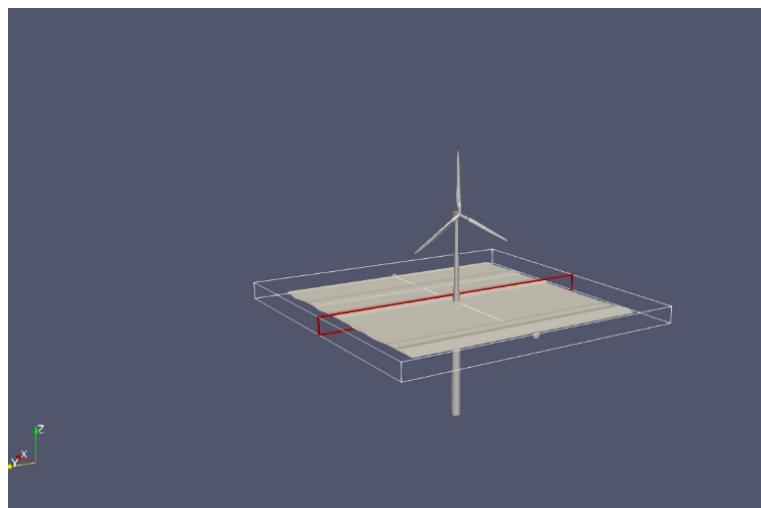
86
87 To simulate or predict actuated forces correctly acting on the model, AI models (simulation
88 model and prediction model) can simulate and predict the rotor thrust in real time with an
89 acceptable degree of accuracy are required. For this purpose, the fully-coupled calculation code
90 of floating wind turbines -FAST is used to run a number of simulations, and the data are
91 thoroughly analysed to obtain a reasonable control structure. A series of simulations of
92 frequency-increasing regular wave conditions and speed-increasing wind conditions are carried
93 out to observe their effect on the real-time rotor thrust. Additionally, the effect of the pitch and
94 surge RAOs of the floating foundation and the wind velocity are quantitatively studied. It
95 reveals that the high-frequency response of thrust is dominated by wave change, whereas low-
96 frequency response is dominated by wind change. Therefore, one simulation model of the thrust
97 acting on the rotor is estimated regarding high-frequency and low-frequency response
98 separately to account the dominating influences. To predict the thrust in more complexed
99 environmental conditions, data for four comparative coupled environmental conditions are
100 considered. They are respectively the irregular wave condition demonstrated by JONSWAP
101 spectrum combined with turbulent wind or the corresponding steady wind, the Pierson-
102 Moskowitz coupled with the same turbulent or steady wind conditions, by which the effect of
103 turbulence and wave spectrum on the thrust force is also investigated. The derived control
104 structures are trained by system identification or artificial neural network techniques to achieve
105 a simulation model and a prediction model separately.

106 The paper is structured as follows: In section 2, the classic model used to undertake the aero-
107 hydro-servo-elastic coupled simulations are described. Section 3 lists the environmental
108 conditions used to conduct the aero-hydro-servo-elastic coupled simulations. In Section 4, the
109 methodology regarding the mathematical expression of the fully-coupled six motions of the
110 platform, hydrodynamic calculation, aerodynamic calculation, mooring modelling, control and

111 electrical-drive are described in details. The analyses of the simulation result in various coupled
112 wave and wind conditions calculated by FAST, as well as the process of how to obtain the two
113 models are presented and discussed in Section 5. Section 6 summarises the conclusions from
114 the study.

115 2. Model description

116 The spar-buoy concept -“Hywind” floating wind turbine is initially developed by Statoil of
117 Norway. Then by NREL, aspects of the original data are adjusted slightly to make the platform
118 appropriate for supporting the NREL 5-MW baseline turbine. To differ from Statoil’s original
119 Hywind concept, it is referred to as “OC3-Hywind”. The OC3-Hywind floating wind turbine
120 concept is considered as the model in this study. We use the turbine specifications of the
121 National Renewable Energy Laboratory (NREL) offshore 5-MW baseline wind turbine, as
122 given in reference [7, 24]. Moreover, Fig.1 presents a model built from the data reflecting the
123 floating system properties documented in[7]. Seen in Fig. 3, a spar type platform is used to
124 carry the wind. Table 1 lists the specification details of the model



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126

127

Fig. 3 Building up the model for simulation

Table 1 General properties of the model

Item	Value
Rating	5 MW

Rotor configuration	Upwind, three blades
Cut-in, rated, cut-out wind speed	3m/s, 11.4m/s, 25m/s
Total draft below sea water level (SWL)	120m
Height to the base above SWL (top of platform)	10m
Height to Tower Top (Yaw Bearing) Above SWL	87.6m
The position of tower CM above SWL Along Tower centerline	43.4m
Tower Structural-Damping Ratio (All Modes)	1%
Nacelle dimension (length, width, height)	14.2m, 2.3m, 3.5m
Platform diameter above the taper	6.5m
Platform diameter below the taper	9.4m
Rotor nacelle assembly (RNA) mass	350000kg
Tower mass	249718kg
Platform mass, Including ballast	7466000kg
CM Location Below SWL Along Platform Centerline	89.9155 m
Platform Roll Inertia about CM	4,229,230,000 kg•m ²
Platform Pitch Inertia about CM	4,229,230,000 kg•m ²
Platform Yaw Inertia about Platform Centerline	164,230,000 kg•m ²
Number of mooring lines	3
Angle Between Adjacent Lines	120°
Depth to fairleads below SWL	70m
Depth to Anchors Below SWL	320m
Radius to Anchors from Platform Centerline	853.87m
Radius to Fairleads from Platform Centerline	5.2m
Unstretched Mooring Line Length	902.2m
Mooring Line Diameter	0.09 m
Equivalent Mooring Line Mass Density	77.7066 kg/m
Equivalent Mooring Line Weight in Water	698.094 N/m
Equivalent Mooring Line Extensional Stiffness	384,243,000 N
Additional Yaw Spring Stiffness	98,340,000 Nm/rad
Baseline control in Region 3	GSPI and constant torque

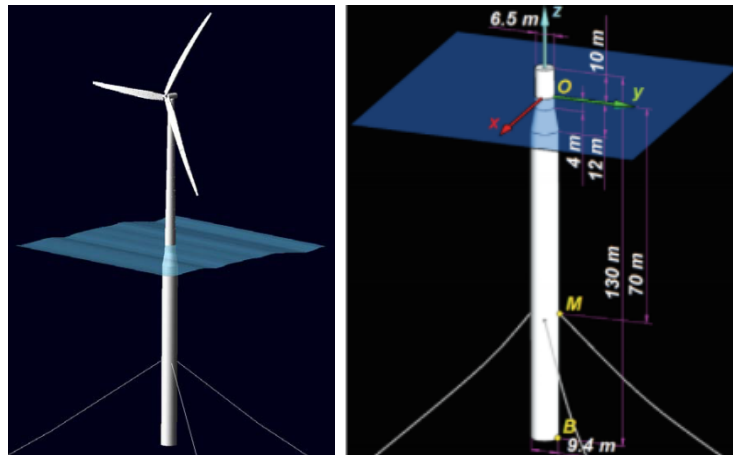


Fig. 4 Dimensions of the floating system

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131 As shown in Fig. 4, the tower base overlaps with the top of the platform and is located 10 m
 132 above the static water level (SWL). Although the top of the tower protrudes with the yaw
 133 bearing, it is situated at 87.6m above SWL. The distribution characteristics of the tower are
 134 based on the base diameter of 6.5m, which is matched with the top width of the platform.
 135 Moreover, the base thickness of the tower is 0.027m, top diameter 3.87m, thickness 0.019m.

136 The tower is 10 meters above the SWL and cantilever to the top of the floating platform, which
 137 is considered rigid. A 3-leg catenary moors the model and the mooring system is spread with
 138 a delta connection. However, in the calculation, the delta connection is eliminated, and the
 139 mooring system is modified with a yaw spring to meet the proper general yaw stiffness.

140

141 3. Load cases

142

143 Table 2 lists all the load cases used in the simulation. In order to observe the effect of the waves
 144 on thrust, LC 1 consists of 35 separate simulations in regular wave conditions of different wave
 145 frequency. Likewise, to investigate the effect of wind speed on thrust, LC 2 is selected. To
 146 further observe the performance of thrust in the sophisticated natural environment, LC 3 is
 147 considered. To investigated the influence the turbulence and wave spectrum has on real-time

148 thrust, LC4-LC are included as comparative load cases. For each simulation, it runs 600s to get
 149 an observation of the steady effect. However, only the data of 300s-500 s is selected for statistic
 150 calculation and analyses in section 5.1 and 5.2, to get rid of the transient effect arising in the
 151 initial simulation stage.

152 **Table 2 List of all the load cases used in the simulation**

Load Case	Enabled DOFs	Wave Conditions	Wind Conditions
LC1	Platform(floaters), tower, drivetrain, blades	Regular airy: H=2m, $\omega=0.1, 0.2, \dots, 3.5$ rad/s	Steady, no shear $u=8$ m/s
LC2	Platform, tower, drivetrain, blades	Regular airy: H=2m $\omega=0.5$ rad/s	Steady, no shear $u=2, 4, \dots, 18$ m/s
LC3	Platform, tower, drivetrain, blades	JONSWAP spectrum $H_s=4.55$ m, $T_p=9.00$ s, $\gamma=2.45$	IECKAI model $\bar{u}=11.40$, TI=20.45
LC4	Platform, tower, drivetrain, blades	JONSWAP spectrum $H_s=1.5$ m, $T_p=6.61$ s, $\gamma=1.00$	IECKAI model $\bar{u}=11.40$, TI=20.45
LC5	Platform, tower, drivetrain, blades	JONSWAP spectrum $H_s=4.55$ m, $T_p=9.00$ s, $\gamma=2.45$	Steady, no shear $u=11.40$ m/s
LC6	Platform, tower, drivetrain, blades	JONSWAP spectrum $H_s=1.5$ m, $T_p=6.61$ s, $\gamma=1.00$	Steady, no shear $u=11.40$ m/s

153

154 To simulate the natural environment, the sea states at the proposed site is characterised by the
 155 JONSWAP spectrum, which is described as follow in a lecture of MIT
 156 OpenCoursWare[25]:

157
$$S(\omega) = \frac{\alpha g^2}{\omega^5} \exp\left\{-1.25 \left(\frac{\omega_p}{\omega}\right)^4 \exp\left(-0.5 \left(\frac{\omega - \omega_p}{\sigma \omega_p}\right)^2\right)\right\} \gamma$$
 Equation 1

158 where, $\alpha = 5.061 \left(\frac{\omega_p}{2\pi}\right)^4 H_s^2 [1 - 0.287 \log \gamma]$; $\sigma = 0.07$ for $\omega < \omega_p$, and $\sigma = 0.09$ for $\omega \geq \omega_p$;

159 $\omega_p = 2\pi / T_p$. H_s is the significant wave frequency, ω_p is the peak frequency, γ is the peak
160 enhance coefficient.

161 Moreover, for the wind field with a turbulent inflow, the IEC Kaimal spectral turbulent model
162 is used to characterise the wind conditions. When the atmospheric stability is aero, the spectra
163 for the three wind components, $K = u, v, w$ (u is horizontal wind speed) are as follow[26]:

164
$$S(f) = \frac{4\sigma_K^2 L_K / \bar{u}_{hub}}{\left(1 + 6fL_K / \bar{u}_{hub}\right)^{5/3}}$$
 Equation 2

165 where f is the cyclic frequency and L_K is an integral scale parameter, which is defined as:

166
$$L_K = \begin{cases} 8.10\Lambda_U, & K = u \\ 2.70\Lambda_U, & K = v \\ 0.66\Lambda_U, & K = w \end{cases}$$

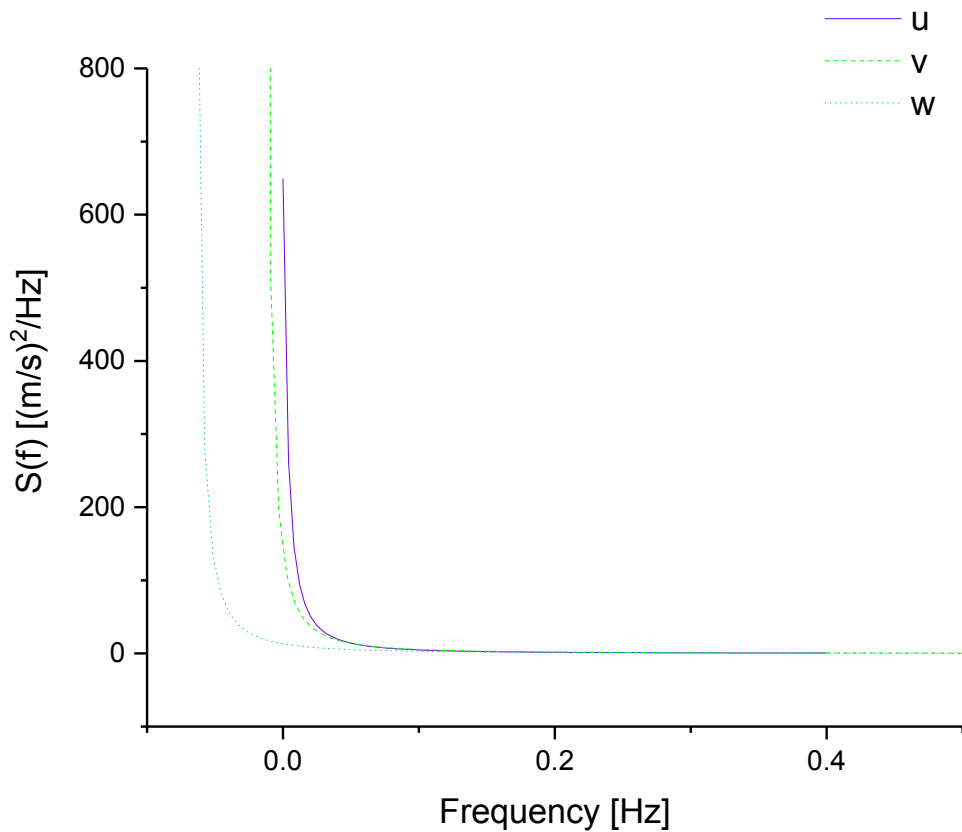
167 According to IEC 61400-3 code, the turbulence scale parameter $\Lambda_U = 0.7 \cdot \min(60m, H_{hub})$.
168

169 However, because $H_{hub} = 90m$, then $\Lambda_U = 42$. The relationships between the standard deviations

170 are established to be: $\sigma_v = 0.8\sigma_u$, $\sigma_w = 0.5\sigma_u$, $\sigma_u = \frac{TI}{100} \bar{u}_{hub}$, TI is the turbulence intensity, \bar{u}_{hub}

171 is the mean (total) wind speed at the reference height. Then the wind spectrum of LC 3 and LC

172 4 is as shown in Fig. 5.



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Fig. 5 Spectrum of the wind adapted in LC 3&4

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4. The methodology of the simulations

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To predict the sophisticated dynamics of the spar-type floating wind turbine, a coupled aero-

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servo-hydro-elastic model which integrates wind-inflow, the turbine control system (servo),

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hydrodynamic and structural-dynamic (elastic) models is adapted to run the simulations.

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4.1 Fully coupled six motions

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When the fully dynamic coupling between the floating foundation and the wind turbine is taken

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into consideration, the coupled motions of the floating foundation in the time domain is

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expressed as follow[27]:

183

$$M_{ij}(q,u,t)\ddot{q}_j = F_i(q,\dot{q},u,t)$$

Equation 3

184 where q is the input of the six DOFs (heave, sway, surge, roll, pitch and yaw); u is the control
 185 inputs; M_{ij} is the (??, ??) component of the inertia mass matrix; \ddot{q}_j is the acceleration of DOF ??;
 186 \dot{q} is the velocity input of the DOFs; F_i is the force applied on DOF i .

187 Because of the balance of forces for the system, the total force applied to the system should
 188 be as follow:

$$189 \quad *F_i + F_i = 0 \quad \text{Equation 4}$$

190 where $*F_i = *F_i^{Hub} + *F_i^{Nacelle} + *F_i^{Tower} + *F_i^{Floater} + *F_i^{Blades}$;

$$191 \quad F_i = F_i^{Aero} + F_i^{Hydro} + F_i^{Gravity} + F_i^{Elastic} + F_i^{Drivet} + F_i^{Mooring} .$$

192 4.2 Hydrodynamic methodology

193 In this study, both regular waves, and irregular (JONSWAP spectrum) waves are generated to
 194 form the coupled load conditions. Airy wave theory is adopted to model the wave kinematics,
 195 including both regular and irregular waves. For regular waves, the wave elevation (ζ) is
 196 represented as a sinusoid with a single amplitude (wave height) and frequency. Irregular is
 197 expressed as a summation or superposition of multiple wave components, as described by an
 198 appropriate wave spectrum[26]:

$$199 \quad \zeta(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} W(\omega) \sqrt{2\pi S_{\zeta}^{2-sided}(\omega)} e^{j\omega t} d\omega \quad \text{Equation 5}$$

200 where, $W(\omega)$ is the Fourier transform of a realisation of a white Gaussian noise time-series
 201 process with zero mean and unit variance. This equation represents the wave elevation as an
 202 inverse Fourier transform of the desired two-sided power spectral density, $S_{\zeta}^{2-sided}$ where j is
 203 an imaginary number, and ω is an individual wave frequency.

204 Potential-flow theory and strip theory (via Morison's equation) are used to model the
 205 hydrodynamic loads. The forces acting on the floater (floating foundation) consist of diffraction,

206 radiation and hydrostatic forces as well as the restoring forces from the mooring lines and the
 207 added inertia of the fully coupled turbine. Therefore, the total external load acting on the floating
 208 foundation $F_i^{Floator}$ is described as follow:

$$209 \quad F_i^{Floator} = -A_{ij}\ddot{q}_j + F_i^{Hydro} + F_i^{Mooring} \quad \text{Equation 6}$$

210 where, A_{ij} is the (??, ??) component of the added mass matrix, F_i^{Hydro} is the i th component of the
 211 applied hydrodynamic load on the floating foundation, $F_i^{Mooring}$ is the i th component of the
 212 load on the floating foundation from the mooring lines.

213 However, F_i^{Hydro} consists of three parts: radiation, diffraction and hydrostatic, hence:

$$214 \quad F_i^{Hydro} = F_i^{Excitation} + F_i^{Hydrostatic} - F_i^{Radiation}$$

215 where, $F_i^{Excitation}$ is the total excitation load from incident waves, which is the result of the
 216 undisturbed pressure field (Froude-Krylov) and wave scattering (diffraction loads), $F_i^{Hydrostatic}$ is
 217 the restoring forces of a freely moving body. $F_i^{Radiation}$ is steady-state hydrodynamic forces and
 218 moments. Thus we get

$$219 \quad F_i^{Hydrostatic} = \rho g V \delta_{i3} - C_{ij}^{Hydrostatic} q_j \quad \text{Equation 7}$$

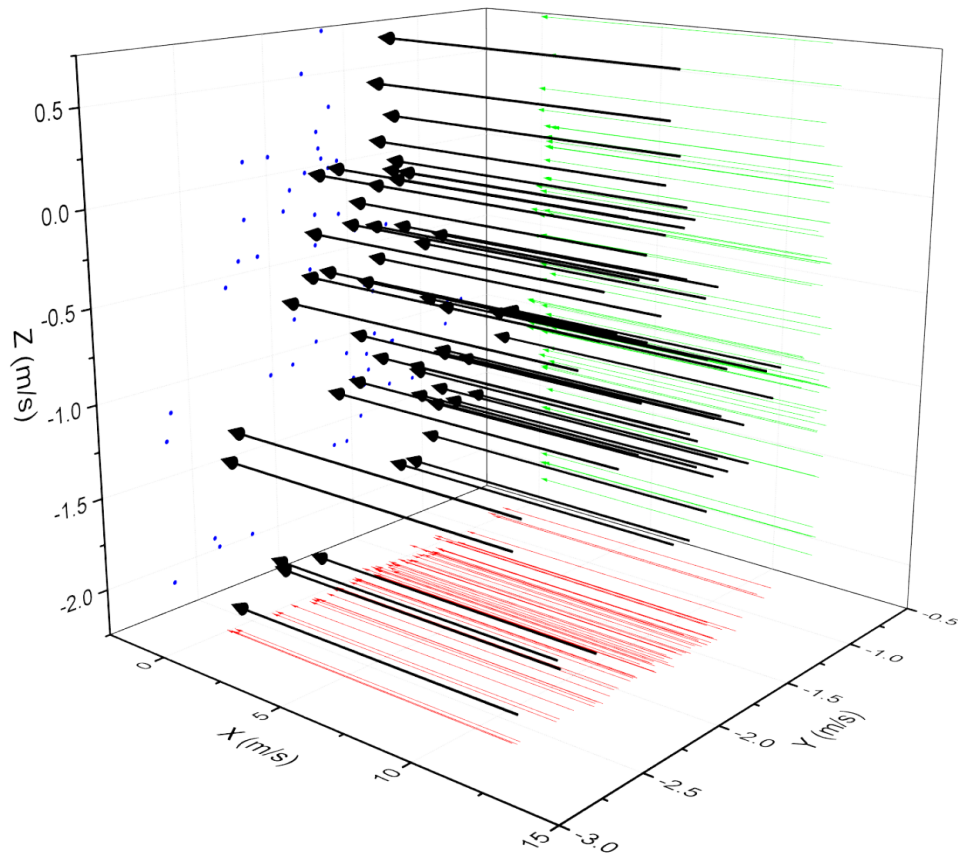
$$220 \quad F_i^{Radiation} = \int_0^t K_{ij}(t-\tau) \dot{q}_j(\tau) d\tau \quad \text{Equation 8}$$

$$221 \quad F_i^{Hydro} = F_i^{Excitation} + \rho g V \delta_{i3} - C_{ij}^{Hydrostatic} q_j - \int_0^t K_{ij}(t-\tau) \dot{q}_j(\tau) d\tau \quad \text{Equation 9}$$

222 where $\rho g V \delta_{i3}$ is the buoyancy force from Archimedes' principle; thereof, δ_{i3} is the (i,3)
 223 component of the Kronecker-Delta function; $-C_{ij}^{Hydrostatic} q_j$ is the hydrostatic and torque
 224 variations resulting from the effects of the water-plane area and the COB; thereof, $C_{ij}^{Hydrostatic}$ is
 225 the (??,??) component of the linear hydrostatic restoring matrix; K_{ij} is the wave radiation
 226 retardation kernel - the impulse-response function of the radiation problem; t is the Simulation
 227 time; τ is the user variable time.

228 4.3 Aerodynamic methodology

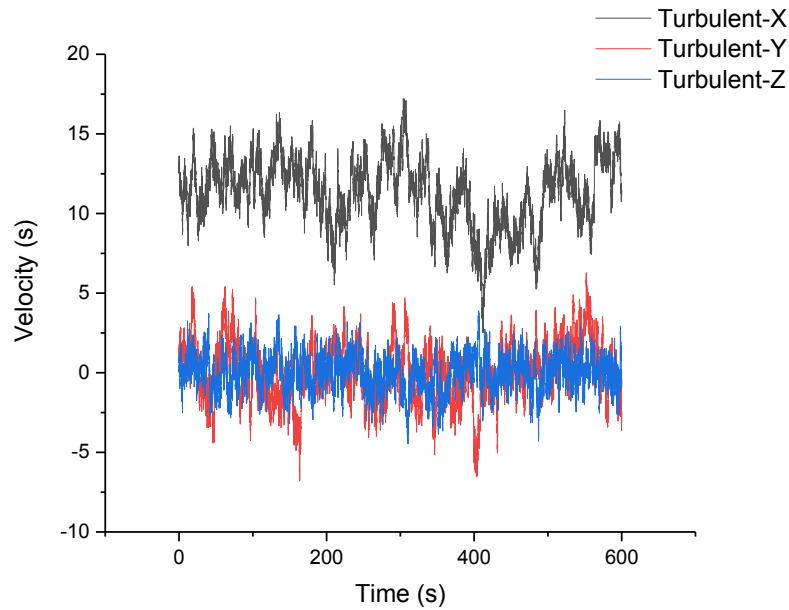
229 In this study, both steady wind fields and turbulent wind fields are generated to couple with
230 different wave conditions. For the convenience of describing the wind field, it is assumed that
231 the inflow wind direction is horizontal as shown in Fig. 6. Power-Law wind profile is used to
232 determine the mean u -component velocity at each height in the wind filed.



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234

Fig. 6 The global coordinate system concerning the turbulent wind



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Fig. 7 History of wind speed in the turbulent wind field

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As discussed in section 3, Fig. 5 displays the spectrum of turbulent wind. Moreover, Fig. 7

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plots the history of wind speed. As shown in Fig. 7, the slow slow-varying component forms

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the turbulence. However, in the natural world, not only the amplitude of the wind speed varies,

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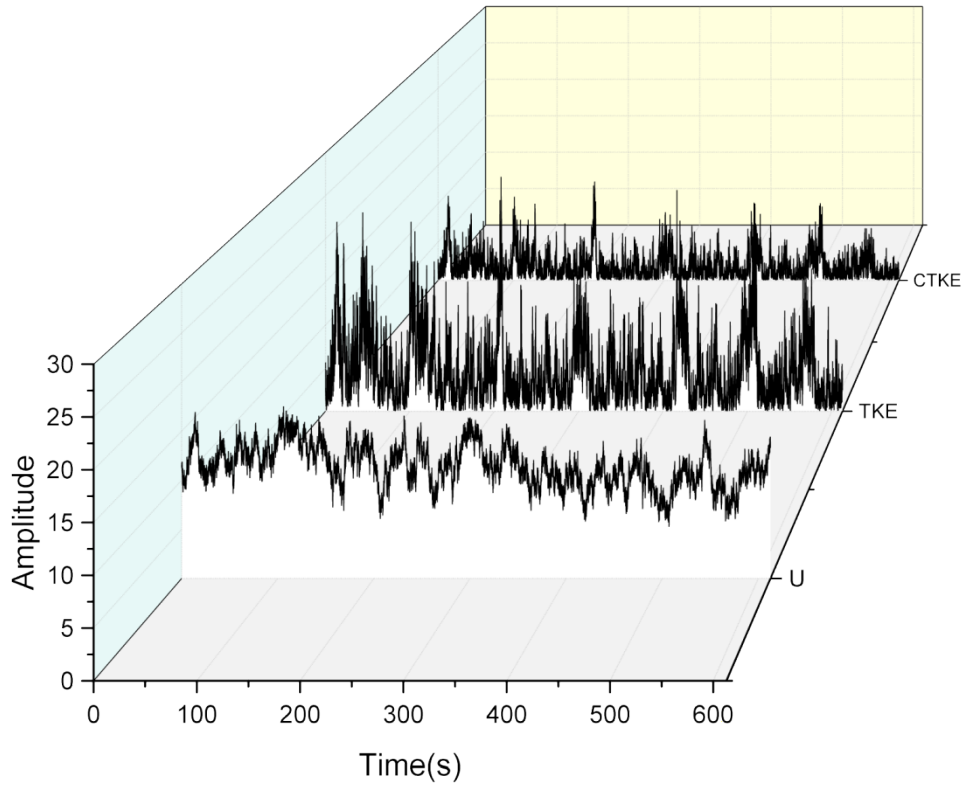
but also the speed direction is time-dependent. For analysis purposes, coherent structures have

241

been defined regarding coherent turbulent kinetic energy (CTKE), and the turbulent kinetic

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energy (TKE) is used to model the power spectrum as shown in Fig. 8.



243

244

Fig. 8 Energy distribution along the turbulent velocity

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4.4 Mooring modelling

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For mooring modelling, the closed-form solution[28] of the classic single line is reformulated

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into a piecewise, multi-segmented system as shown in Fig. 9 (where r_i is the node position

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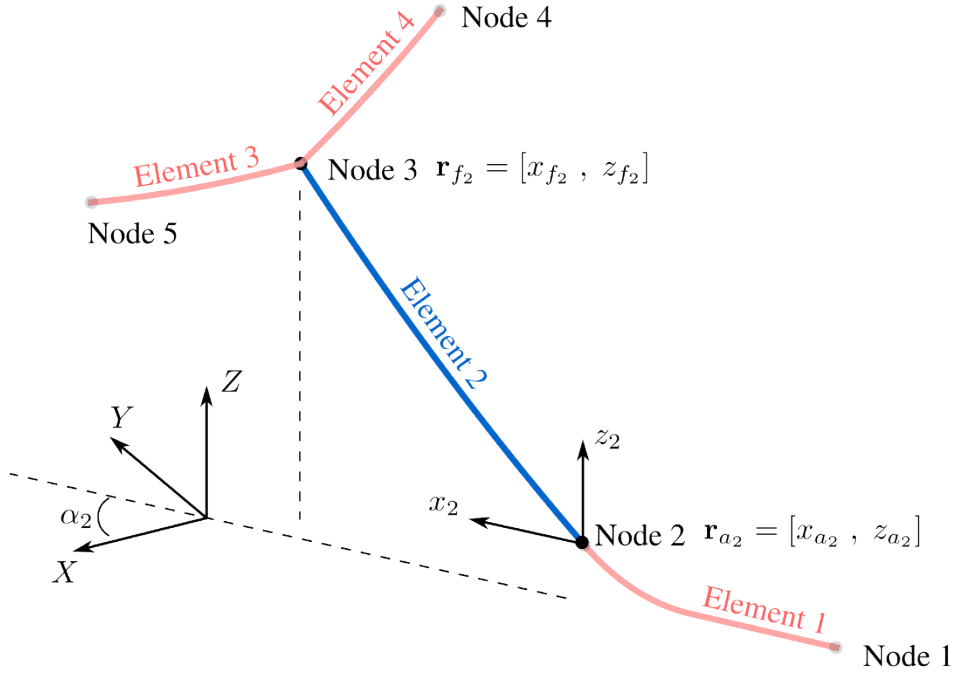
vector, $x_i z_i$ is a local frame, XYZ is the global coordinate system). This piecewise system is

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composed of a collection of nodes and elements. The mean forces in the mooring system,

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including elasticity, weight (in fluid), and geometric nonlinearities, are accounted.



251

252

Fig. 9 the piece-wise, the multi-segmented system of a classic mooring line[29]

253

The solution process begins with the evaluation of the two continuous analysis catenary

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equations of each element based on the horizontal and vertical orifice offsets obtained through

255

the nodal displacement relationship. The element is defined as a component that connects two

256

adjacent nodes. Evaluating the force balance equation for each node as follow:

257

$$\{F\}_X^j = \sum_{i=1}^n [H_i \cos(\alpha_i) - F_{X_j}^{ext}] = 0$$

$$\{F\}_Y^j = \sum_{i=1}^n [H_i \sin(\alpha_i) - F_{Y_j}^{ext}] = 0$$

Equation 10

$$\{F\}_Z^j = \sum_{i=1}^n [V_i] - F_{Z_j}^{ext} + M_j g - \rho g B_j = 0$$

258

where, H is the horizontal fairlead force, V is the vertical fairlead force, $n = 1, 2, \dots$ is the

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elements at the node j , F^{ext} is the external force in the global coordinate system. Once the

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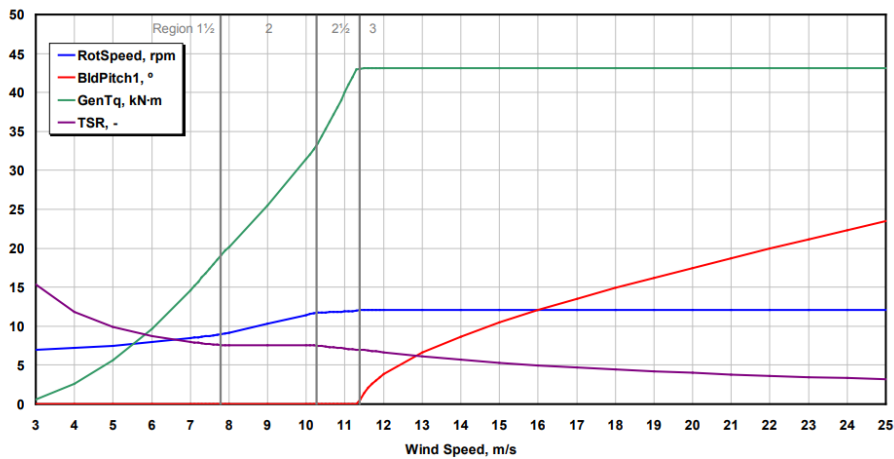
element fairlead (H, V) and anchor (H_a, V_a) values are known at the element level, their

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contribution at each element's anchor and fairlead are obtained by Equation 13.

262 4.5 Control and electrical-drive modelling

263 A conventional variable-speed, variable blade-pitch-to-feather control modelling strategy is
 264 usually used for onshore NREL 5-MW turbines, and an example of the steady-state behaviour
 265 of onshore NREL 5-MW turbine under operational control is given in Fig. 10 (where ‘RotSpeed’
 266 is for rotor speed, ‘BldPitch1’ is for the pitch angle for blade 1, ‘GenTq’ is for generator torque,
 267 and ‘TSR’ is for tip-speed ratio). However, the conventional pitch-to-feather control usually
 268 results in a reduction of the steady-state rotor thrust when increase wind speed above rated, when
 269 applied on a floating wind turbine as discussed by Nielsen et al. in [30, 31].



270

271 **Fig. 10 Steady-State behaviour of onshore NREL 5-MW turbine under operational control[32]**

272 Therefore, in this research, two modifications are adapted as listed in Table 3, and a time step is
 273 added for the Bladed-style DLL controllers, which is independent of the ServoDyn time step.
 274 Moreover, a linear ramp and first-order low-pass filter are applied to the blade-pitch command
 275 from the Bladed-style DLL.

276

Table 3 Control System Property Modifications[32]

Proportional Gain at Minimum Blade-Pitch Setting	0.006275604 s
Integral Gain at Minimum Blade-Pitch Setting	0.0008965149
Constant (Rated) Generator Torque in Region 3	43,093.55 $N \cdot m$

277

278 5. Results

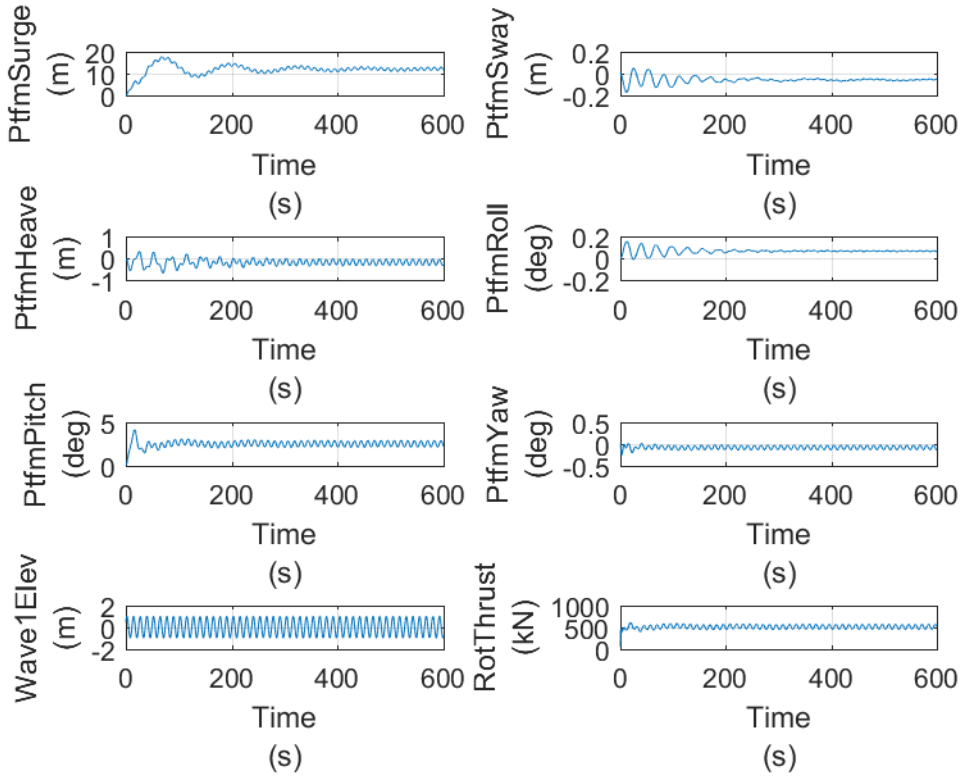
279 This section provides the analyses of data from simulations and proposes two models. The
280 simulation model of the thrust acting on the rotor is obtained by the data from LC1 and LC 2;
281 the simulation model is estimated by system identification techniques regarding high frequency
282 and low-frequency response separately to account the dominating influence. The prediction
283 model is obtained by artificial neural network method to train a black box by the data from
284 simulations of LC 3-LC 6.

285 5.1 Rotor thrust prediction under LC 1

286 In this section, the wind speed of uniform wind field is 8m/s, and 35 frequencies of 0.1 -3.5
287 rad/s of regular waves with a constant wave height of 2m, are used to run the simulations. The
288 outputs of one standard simulation are shown in Fig. 11. As pointed out in [33, 34]], the same
289 phenomenon is observed from the series simulation, that for numerous environmental condition,
290 the significant response is the surge and pitch motions. Moreover, thrust is fundamentally due
291 to the relative wind speed experienced by the turbine blades. For example, if the platform is
292 fixed, then the relative wind speed experienced by the edges will be the same as the actual wind
293 speed.

294 However, for a floating wind turbine, the platform is moving back and forward relative to the
295 wind, and in this case, the turbine blades experience the relative wind speed. Since we are only
296 looking at wind and wave direction in the x-axis, then the platform motions are predominately
297 pitched and surge, i.e. roll, yaw, sway is practically zero as shown in Fig. 11. When the platform
298 pitches or surges this adds a component to the relative wind speed. Hence when predicting the
299 thrust for a floating wind turbine, the authors would like to use surge, pitch (which provide the
300 corresponding component of wind) and the actual wind speed. Therefore, the surge and pitch

301 motions of the platform are extracted from the outputs to observe their correlation with the
 302 aerodynamic thrust in this section.



303

304 **Fig. 11 Some outputs from one standard simulation under $\omega=0.5$ rad/s, $u=8$ m/s**

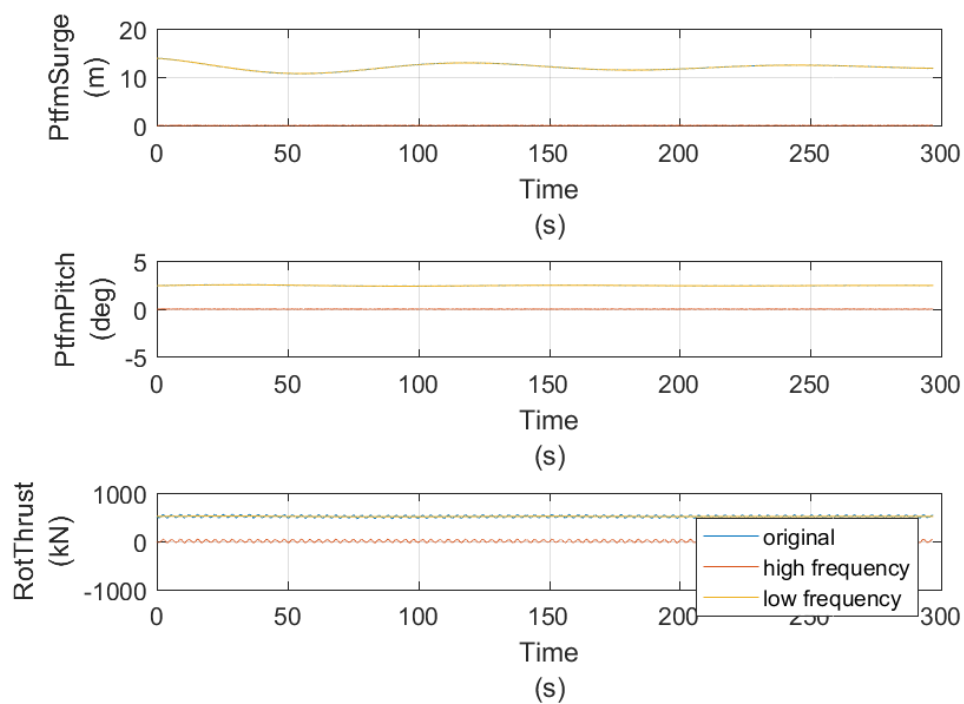
305 To observe the correlation between Surge, pitch, and thrust, two 4th order Butterworth filters
 306 are designed to separate the thrust, surge, and pitch signals into two independent components
 307 (high-frequency component and low-frequency component), respectively. The cut-off
 308 frequency for both high pass filter and low pass filter is 0.01 Hz, to include all the data of the
 309 three signals (surge, pitch, and thrust). Also, the transfer function is expressed regarding b and
 310 a as:

$$311 \quad H(z) = \frac{B(z)}{A(z)} = \frac{b(1) + b(2)z^{-1} + b(3)z^{-2} + b(4)z^{-3} + b(5)z^{-4}}{a(1) + a(2)z^{-1} + a(3)z^{-2} + a(4)z^{-3} + a(5)z^{-4}} \quad \text{Equation 11}$$

312 where, b and a are the transfer function coefficients of the 4th-order filters. And

$$313 \quad b = [0.9990, -3.9959, 5.9938, -3.9959, 0.9990], \quad a = [1.0000, -3.9979, 5.9938, -3.9938, 0.9979]$$

314 for the high pass filter; $b = [2.3759, 9.5035, 14.2550, 9.5035, 2.3759] \times 10^{-14}$,
 315 $a = [1.0000, -3.9979, 5.9938, -3.9938, 0.9979]$ for the low pass filter. Once the filters designed,
 316 the three signals are separated. One example when $\omega = 1.9$ rad/s, $u = 8$ m/s is given in Fig.9(
 317 see separation analyses for other cases in the supplementary dataset). Fig. 12 illustrates that
 318 the filters are effective to separate the properties of the original signal into two independent
 319 parts.



320

321

Fig. 12 Separation of the three signals at $\omega = 1.9$ rad/s, $u = 8$ m/s

wind		wave			surge_high		pitch_high		thrust_high		surge_lo	pitch_low	thrust_low
wind(m/s)	w(rad/s)	f(Hz)	Hs(m)	H(m)	f(Hz)	A(deg)	f(hz)	T(KN)	f(hz)	H(m)	A(deg)	T(KN)	
8	0.2	0.032	2	1.49	0.033	0.133	0.031	11.969	0.033	11.817	2.422	504.803	
8	0.5	0.08	2	1.504	0.080	0.760	0.080	116.955	0.086	11.963	2.455	506.259	
8	0.7	0.111	2	0.859	0.111	0.098	0.111	37.112	0.112	11.951	2.449	505.373	
8	0.9	0.143	2	0.509	0.142	0.076	0.142	42.554	0.143	11.983	2.449	505.133	
8	1.2	0.191	2	0.249	0.188	0.049	0.188	49.369	0.195	11.936	2.445	504.777	
8	1.5	0.239	2	0.119	0.236	0.029	0.235	51.131	0.243	11.962	2.446	504.705	
8	1.6	0.255	2	0.090	0.254	0.024	0.250	50.143	0.263	11.957	2.446	504.722	
8	1.9	0.302	2	0.040	0.298	0.012	0.295	46.066	0.302	11.958	2.446	504.711	
8	2.2	0.35	2	0.018	0.334	0.006	0.342	43.475	0.350	11.966	2.447	504.744	
8	2.5	0.398	2	0.012	0.556	0.005	0.386	44.313	0.397	11.950	2.445	504.512	
8	2.9	0.462	2	0.020	0.455	0.011	0.453	63.071	0.461	11.958	2.446	504.636	
8	3.5	0.557	2	0.014	0.548	0.009	0.547	24.485	0.558	11.961	2.446	504.691	

322

323

324 **Fig. 13 the separated components of the surge, pitch, and thrust for a range of ω at the wind speed of**
 325 **8m/s**

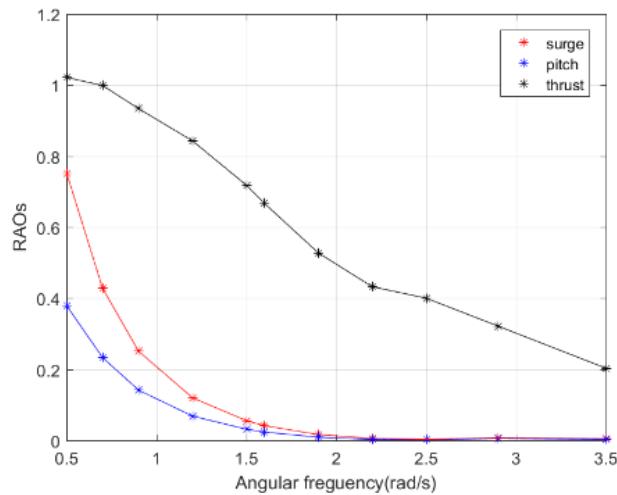
326 The static calculation results of the separated components for all simulations included in load
 327 cases 1 are shown in Fig. 13. Looking at results, it is noticed that low-frequency components
 328 of the three signals (surge_low, pitch_low, thrust_low) tend to be independent on wave
 329 frequency. However, the high-frequency components (surge_high, pitch_high, thrust_high)
 330 tend to be dependent on the wave frequency. Therefore, if surge_RAO, pitch_RAO and
 331 thrust_RAO are defined as follow :

332
$$surge_RAO = \frac{surge}{Hs} \quad \text{Equation 12}$$

333
$$pitch_RAO = \frac{pitch}{Hs} \quad \text{Equation 13}$$

334
$$thrust_RAO = \frac{thrust}{120} \quad \text{Equation 14}$$

335 where, Hs is the significant wave height, 120 is the introduce scaling coefficient. Thus, the
 336 RAOs are presented in Fig. 14.



337

338 **Fig. 14 the RAOs of the surge, pitch, and thrust at a range of ω at the wind speed of 8m/s**

339 It is noticed in Fig. 14, there might be a linear correlation between thrust_high, surge_high
 340 and pitch_high when the wave frequency increase. Then the structure of the mathematical

341 model to predict the high-frequency component of the aerodynamic thrust is assumed to be
 342 as:

$$343 \quad RAO_thrust_high = k_1 \times RAO_pitch_high + k_2 \times RAO_surge_high + c(\omega) \quad \text{Equation 15}$$

344 where k_1 , k_2 are model parameters, and $c(\omega)$ is a constant modification relative to wind
 345 speed(see section 5.1.2), which could be identified by using the data of inputs(
 346 RAO_surge_high , RAO_pitch_high) and outputs($thrust_RAO_high$).

348 5.2 Rotor thrust under LC 2

349 As learned from above, the high-frequency component of the three observed signals are
 350 independent of wave frequencies, in this section, the frequency of regular waves (airy) is 0.5
 351 rad/s, and velocities of $u=2, 4, \dots, 18$ m/s of steady wind, are used to run the simulations, for
 352 observing the effect of wind speed on aerodynamic thrust. One example of the simulation at ω
 353 $=0.5$ rad/s, $u=12$ m/s after applying the data through the high pass and low pass filters is shown
 354 in Fig. 15

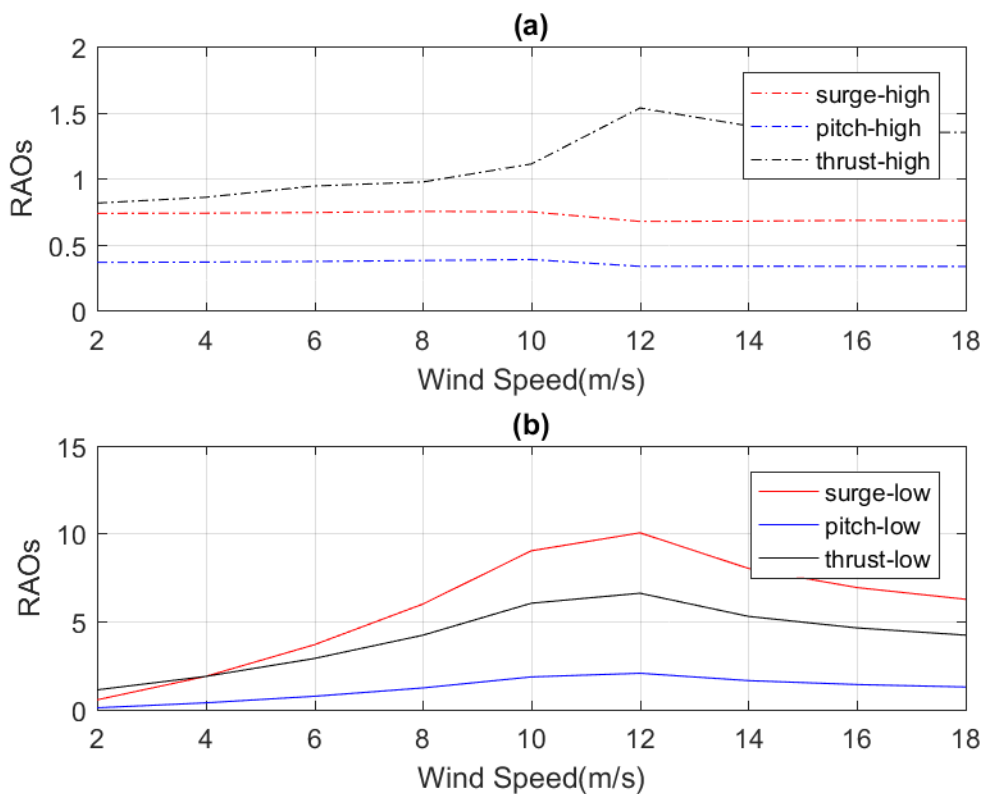
wind		wave			surge_high		pitch_high		thrust_high		surge_low	pitch_low	thrust_low
wind(m/s)	w(rad/s)	f(Hz)	Hs(m)	H(m)	f(Hz)	A(deg)	f(hz)	T(KN)	f(hz)	H(m)	A(deg)	T(KN)	
2	0.5	0.08	2	1.474	0.080	0.732	0.080	97.768	0.080	1.102	0.196	134.074	
4	0.5	0.08	2	1.477	0.080	0.736	0.080	103.036	0.080	3.763	0.761	225.799	
6	0.5	0.08	2	1.487	0.080	0.746	0.080	113.295	0.080	7.355	1.508	347.780	
8	0.5	0.08	2	1.504	0.080	0.760	0.080	116.955	0.086	11.973	2.456	506.397	
10	0.5	0.08	2	1.497	0.081	0.775	0.079	133.277	0.087	18.030	3.707	724.560	
12	0.5	0.08	2	1.352	0.081	0.672	0.079	184.146	0.083	20.070	4.121	792.488	
14	0.5	0.08	2	1.356	0.081	0.673	0.079	167.927	0.083	16.023	3.286	634.720	
16	0.5	0.08	2	1.368	0.079	0.672	0.078	161.980	0.087	13.850	2.840	556.172	
18	0.5	0.08	2	1.362	0.079	0.670	0.078	162.021	0.087	12.524	2.563	507.111	

355 **Fig. 15 the separated components of the surge, pitch, and thrust for a range of u at the $\omega=0.5$ rad/s**
 356 Similarly, the static calculation of the separated components for all simulations included in
 357 load cases 2 is carried out, and the results are listed in Fig. 15. Different from being independent
 358 of wave frequencies, the low-frequency components of the three signals vary dramatically as
 359 the wind speed changes. As the RAOs of the three signals are defined above, the analyses
 360 carried out all based on RAOs in the section, as shown in Fig. 16. Learn from Fig. 16(a), the
 361

362 high-frequency components of surge and pitch tend to be steady when the wind speed increase,
 363 while the high- requency thrust component increses slightly, which is primarily described by
 364 the $c(u)$ in the structure of the mathematical model. So when Equation17 modified with an
 365 offset the structure represented by equation 17 can predict the high-frequency component when
 366 wind speed changes. Therefore, the fabric used to predict the thrust-high is modified as:

$$367 \quad RAO_thrust_high = k_1 \times RAO_pitch_high + k_2 \times RAO_surge_high + c(\omega) + c_1$$

368 Equation 16
 369 where, c_1 is a constant offset related to wind speed.



370
 371 **Fig. 16 Behaviours of the high or low components of the surge, pitch, and thrust at a range of wind speed**
 372 **at $\omega = 0.5$ rad/s**

373 Additionally, it is observed in Fig. 16(b) that there is an apparent linear correlation among
 374 surge_low, pitch_low, thrust_low when the wind speed increase from 2m/s to 8m/s, and they
 375 all reach the maximum when the wind speed is 12m/s. Therefore the structure of the

376 mathematical model to predict the low-frequency component of aerodynamic thrust is assumed
377 to be as:

$$378 \quad RAO_thrust_low = k_3 \times RAO_pitch_low + k_4 \times RAO_surge_low + k_5 \times wind_speed + c_2$$

379 **Equation 17**
380 where k_3 , k_4 , k_5 are also model parameters, but different from k_1 , k_2 , and c_2 is an allowable
381 constant offset to make up the uncertainty might result from the linear structure. And it will
382 be identified by using the data of inputs(RAO_surge_high , RAO_thrust_high) and
383 outputs($thrust_RAO_high$) as discussed below.

384 5.2 Training of the simulation model of rotor thrust

385 As discussed above, two structures (Equation 18&19) are obtained to predict the aerodynamic
386 thrust:

$$387 \quad RAO_thrust_high = k_1 \times RAO_pitch_high + k_2 \times RAO_surge_high + c(\omega) + c_1$$

$$388 \quad RAO_thrust_low = k_3 \times RAO_pitch_low + k_4 \times RAO_surge_low + k_5 \times wind_speed + c_2$$

389 For the simplicity of scientific discussion, RAO_pitch_high , RAO_surge_high are
390 represented as inputs u_1, u_2 , and RAO_pitch_low , RAO_surge_low as inputs $*u_1, *u_2$ to
391 differ from, RAO_pitch_high RAO_surge_high . Similarly, the output
392 $thrust_RAO_high$ will be discussed as Y_{low} , and Y_{high} . As the model structures and some input
393 and output signals are obtained, the values of adjustable parameters in the given model structure
394 could be estimated by minimising the error between the model output and the measured
395 response. The output y_{model} of the linear model is provided by:

$$396 \quad y_{model}(t) = Gu(t)$$

397 where G is the transfer function, $y_{model}(t)$ is the simulated response of the model for a given
 398 input $u(t)$ (such as $u_1, u_2, *u_1, *u_2$). The principal of determining G is minimising the difference
 399 between the model output $y_{model}(t)$ and the measured output $y_{meas}(t)$ (such as Y_{low}, Y_{high}). The
 400 minimisation criterion is a weighted norm of the error, $v(t)$, where: $v(t) = y_{meas}(t) - y_{model}(t)$.

401 Therefore for Equation 18, the output $y_{model}(t)$ is related to four inputs in total, which are
 402 $RAO_pitch_high, RAO_surge_high, \omega, Y_{high}$ (discussed as u_1, u_2, u_3, u_4). Then the transfer
 403 functions are obtained as follow:

404 From u_1 to y_{model} : $G_1(s) = \frac{864.9s - 107.4}{s^2 + 1.853s + 2.271}$

405 From u_2 to y_{model} : $G_2(s) = \frac{-266.6s + 90.01}{s^2 + 2.105s + 1.861}$

406 From u_3 to y_{model} : $G_3(s) = \frac{52.48s - 9.524}{s^2 + 5.725e^{-10}s + 1.374}$

407 From u_4 to y_{model} : $G_4(s) = \frac{-0.2402s + 1.105}{s^2 + 1.379s + 1.317}$

408 Fit to estimation data reaches 99.12%, the mean square error is 0.0206 when compared to the
 409 original signal.

410 For Equation 19, the output $y_{model}(t)$ is related to four inputs in total, which are
 411 $RAO_pitch_low, RAO_surge_low, wind_speed, Y_{low}$ (discussed as $*u_1, *u_2, *u_3, *u_4$). Then
 412 the transfer functions are obtained as follow:

413 From $*u_1$ to y_{model} : $G_1(s) = \frac{-19.72s + 2.124}{s^2 + 3.86s + 0.1255}$

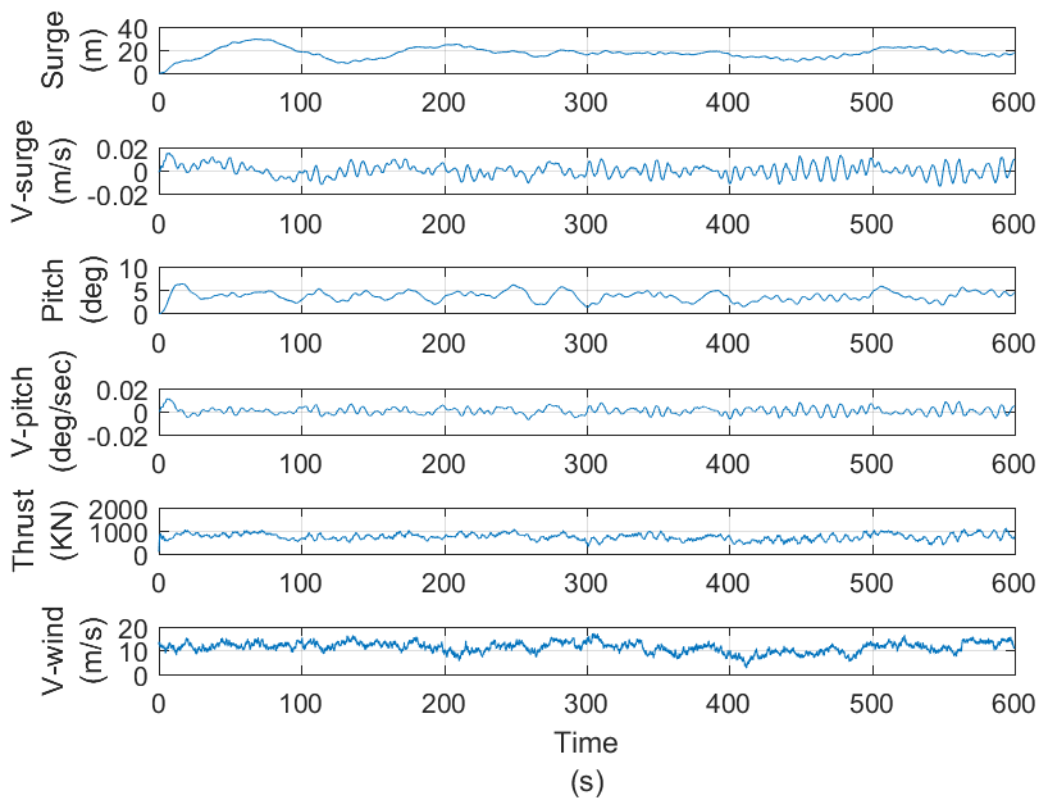
414 From $*u_2$ to y_{model} : $G_2(s) = \frac{-4.211s + 2.433}{s^2 + 4.208s + 1.718}$

415 From $*u_3$ to y_{model} : $G_3(s) = \frac{3.712s - 0.2842}{s^2 + 1.407s + 1.009}$

416 From $*u_4$ to y_{model} : $G_4(s) = \frac{0.1383s - 0.1432}{s^2 + 0.04644s + 0.6857}$

417 **5.3 Rotor thrust under LC 3-6 and the prediction model of rotor thrust**

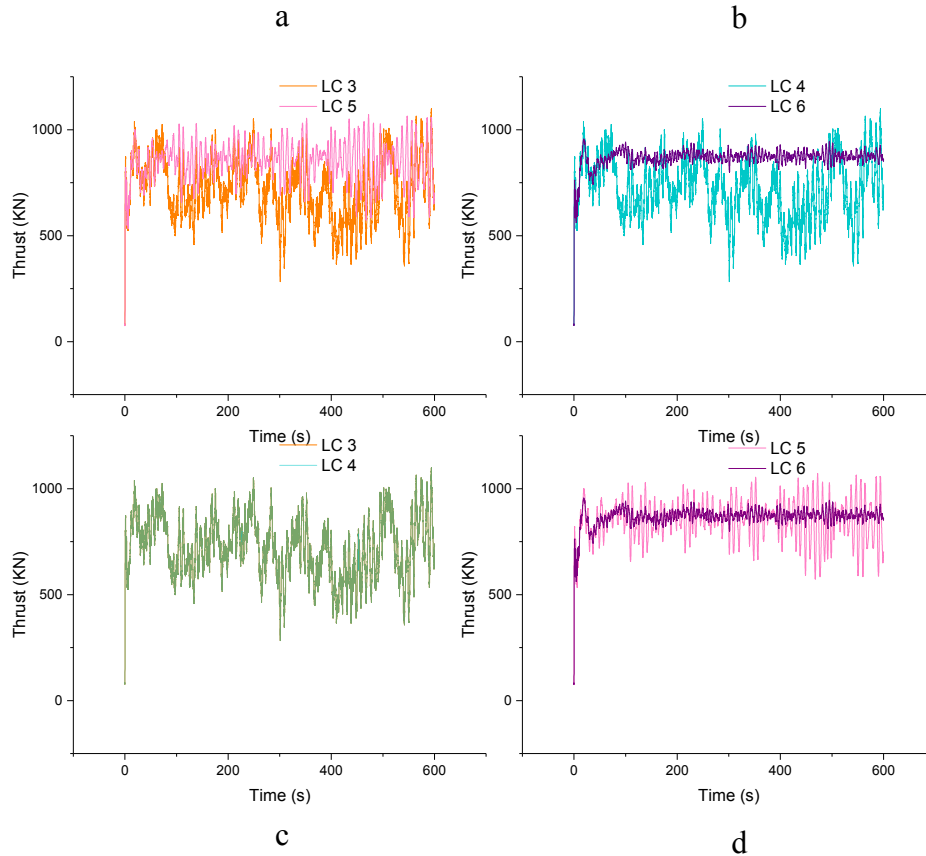
418 Real sea states are more sophisticated than regular-wave and steady-wind conditions. For LC3-
 419 LC4, JONSWAP spectrum and IECKAI turbulent model are adapted to characterise the
 420 environmental conditions, and corresponding simulations are undertaken. To observe the effect
 421 of turbulence inflow on the aerodynamic thrust, the identical steady wind cases (load cases
 422 4.X) are simulated as well. The output of some significant response of LC3 is shown in Fig. 17



423

424

Fig. 17 Significant outputs from a simulation of LC 3



425

426

Fig. 18 Thrust comparison of LC 3 -LC 6

427

It is learned from Fig. 18(a & b) that the average thrust force is just slightly varied by the

428

turbulence. Although the inflow is turbulent, the mean wind speed of LC 3, LC 4 is identical

429

to that of the uniform wind field LC 5, LC 6. Therefore, the mean thrust force acting on the

430

rotor is nearly the same. Nevertheless, the thrust force becomes very unstable in the presence

431

of turbulence. Fig. 18(c) compares the thrust forces of identical turbulent inflows but different

432

wave spectrum parameters, and it is learned that in the presence of turbulence the influence of

433

wave changes on aerodynamic thrust are ignorable. However, as shown in Fig. 18(d), which

434

compares the thrust forces of identical steady wind inflow but different wave conditions, the

435

changes in parameters of wave spectrum have impacts on the behaviour of the thrust, though

436

the mean thrust force is almost the same. Therefore, data of five inputs are included in training

437

a model to predict the thrust force acting on the rotor for sophisticated environmental

438 conditions like LC 3-LC 6. If the model is written in the term of $A(z)y_{model}(t) = B(z)u(t) + e(t)$,
439 the inputs

440 $u(t)$ includes the pitch($u_1(t)$), the velocity of pitch ($u_2(t)$), the surge($u_3(t)$), the velocity of
441 surge, and the velocity of wind($u_5(t)$). Of course, $y_{model}(t)$ is the real-time thrust, $A(z)$ and
442 $B(z)$ are the estimable parameters of the model. Thus one thrust force prediction model is
443 obtained, $e(t)$ is the noise source. The estimable parameters are as follow:

444
$$A(z) = 1 - 2.191z^{-1} + 1.884z^{-2} - 1.069z^{-3} + 0.5424z^{-4} - 0.3577z^{-5} + 0.2575z^{-6} - 0.08479z^{-7}$$
$$- 0.001325z^{-8} + 0.04657z^{-9} - 0.02854z^{-10}$$

445
$$B_1(z) = 59.61z^{-1} - 330.3z^{-5} + 634.7z^{-6} - 364z^{-10}$$

446
$$B_2(z) = 99.66z^{-1} - 158.8z^{-3} - 119.6z^{-4} - 275.7z^{-5} - 584.1z^{-7} - 351.5z^{-8} - 29.44z^{-9}$$

447
$$B_3(z) = -51.46z^{-1} + 31.95z^{-4} + 30.17z^{-5} - 10.65z^{-10}$$

448
$$B_4(z) = -43.35z^{-1} + 24.39z^{-2} + 27.39z^{-3} - 4.442z^{-7} + 19.42z^{-8} + 50.99z^{-9} + 51.06z^{-10}$$

449
$$B_5(z) = -0.1884z^{-1} + 0.3314z^{-2} - 0.214z^{-3} + 0.3012z^{-4} - 0.4324z^{-5} + 0.4156z^{-6} - 0.2558z^{-7}$$
$$+ 0.115z^{-8} - 0.2601z^{-9} + 0.2063z^{-10}$$

450 Generally, the bigger the dataset is, the more accurate model will be obtained, and the model
451 will slightly increase when the database reaches a certain level. For the model proposed in this
452 paper, the size of the dataset is 192000. And the accuracy of fitting to estimate data is [98.88
453 98.73 99.84 99.69] %, and the mean square error is [2.396 2.336 0.02436 0.01664] respectively
454 for LC3-LC6. The final prediction error for the above model is 1.194.

455 5 Conclusions

456 In this paper, numerous simulations based on the classic model of ‘OC3-hywind’ floating wind
457 turbine have been carried out, and the data from simulations has been analysed in great details,
458 Moreover, artificial neural network techniques are used to obtain the structure of the thrust
459 force simulation and prediction models of the real-time thrust acting on the rotor of FOWTs. .
460 It could be used as a software-in-the-loop application in the experiments to provide real-time
461 thrust varying with the input signals(like surge, surge velocity, pitch, pitch_velocity, wind
462 speed, wave condition ed.).

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588 **List of equations**

589 $S(\omega) = \frac{\alpha g^2}{\omega^5} \exp\{-1.25 \left(\frac{\omega_p}{\omega}\right)^4 \exp(-0.5(\frac{\omega-\omega_p}{\sigma\omega_p})^2)\} \gamma$ Equation 18

590 $S(f) = \frac{4\sigma_K^2 L_K / \bar{u}_{hub}}{\left(1 + 6fL_K / \bar{u}_{hub}\right)^{5/3}}$ Equation 29

591 $M_{ij}(q, u, t)\ddot{q}_j = F_i(q, \dot{q}, u, t)$ Equation 310

592 $*F_i + F_i = 0$ Equation 411

593 $\zeta(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} W(\omega) \sqrt{2\pi S_{\zeta}^{2-sided}(\omega)} e^{j\omega t} d\omega$ Equation 511

594 $F_i^{Floator} = -A_{ij}\ddot{q}_j + F_i^{Hydro} + F_i^{Mooring}$ Equation 612

595 $F_i^{Hydrostatic} = \rho g V \delta_{i3} - C_{ij}^{Hydrostatic} q$ Equation 712

596 $F_i^{Radiation} = \int_0^t K_{ij}(t-\tau)\dot{q}_j(\tau)d\tau$ Equation 812

597 $F_i^{Hydro} = F_i^{Excitation} + \rho g V \delta_{i3} - C_{ij}^{Hydrostatic} q_j - \int_0^t K_{ij}(t-\tau)\dot{q}_j(\tau)d\tau$ Equation 912

598 $u(z) = \bar{u}_{hub} \left(\frac{z}{H_{hub}}\right)^p$ Equation 1014

599 $Coh_{i,j} = \exp\left(-a \sqrt{\left(\frac{fr}{\bar{u}_{hub}}\right)^2 + \left(0.12 \frac{r}{L_c}\right)^2}\right)$ Equation 1115

600 $\{F\}_X^j = \sum_{i=1}^n [H_i \cos(\alpha_i) - F_{X_j}^{ext}] = 0$ Equation 1216

601 $\{F\}_Y^j = \sum_{i=1}^n [H_i \sin(\alpha_i) - F_{Y_j}^{ext}] = 0$ Equation 1319

602 $\{F\}_Z^j = \sum_{i=1}^n [V_i] - F_{Z_j}^{ext} + M_j g - \rho g B_j = 0$ Equation 1421

603 $surge_RAO = \frac{surge}{Hs}$ Equation 1521

604 $pitch_RAO = \frac{pitch}{Hs}$ Equation 1621

604 $thrust_RAO = \frac{thrust}{120}$ Equation 1621

605	$RAO_thrust_high = k_1 \times RAO_pitch_high + k_2 \times RAO_surge_high + c(\omega)$	
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607	Equation 18	23
608	$RAO_thrust_low = k_3 \times RAO_pitch_low + k_4 \times RAO_surge_low + k_5 \times wind_speed + c_2$	
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