

Dynamic and structural performances of offshore floating wind turbines in turbulent wind flow

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Abstract

A realistic turbulent wind field differs from a steady uniform one, in terms of the wind shear, the turbulence intensity and the coherence structure. Although it has been clear that an offshore floating wind turbine will behave differently in the turbulent wind, the individual effect of the above three items are not investigated sufficiently until now. The primary objective of the present research is to investigate in details how the wind shear, the turbulence intensity and the coherence influence the dynamic and structural responses of offshore floating wind turbines. Aero-hydro-servo-elastic coupled simulation of a semi-submersible floating wind turbine is run in time-domain. The wind shear has a limited effect on the global responses of the floating wind turbine although its influence on each individual blade is considerable. Comparatively, the floating wind turbine is quite sensitive to the turbulence intensity. In a wind field with high turbulence intensity, the platform motions become more violent and the structural loads are increased substantially. The proper orthogonal decomposition method is used to investigate the coherence quantitatively. A partial coherence structure helps to reduce the flow variation seen by the rotor and thereby beneficial to the safety of the floating wind turbine.

Keywords: turbulent wind; wind shear; turbulence intensity; coherence; offshore floating wind turbine; dynamic response; structural response

1. Introduction

Currently, great efforts are made around the world to pursue alternative energy sources, which are expected to be clean, sustainable and economic-efficient. Among various renewable energy resources, the application of wind energy has been proved successful, and the industry is trying to move to deep water zone to exploit the offshore wind energy. Since the proposal of the Hywind concept (Equinor, 2017a), the world's first full-scale offshore floating wind turbine, a set of floating wind turbine concepts have been proposed. Most recently, Hywind Scotland, the world's first floating wind farm, already starts to deliver electricity to the grid (Equinor, 2017b).

The model test has been accepted as a reliable approach to study the performances of offshore floating wind turbines. Duan et al. (2016) launched a model test program to investigate the vortex

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32 induced motion of a spar-type floating wind turbine. Li et al. (2018b) measured the dynamic response
33 of a semisubmersible floating wind turbine in experimental environment, and a free-rotation method
34 was proposed to correct the Reynolds number dissimilitude. Oguz et al. (2018) investigated the
35 dynamics of a TLP floating wind turbine with both numerical and experimental methods. Apart from
36 model test research, numerical simulation technology is also widely adopted by many researchers. Liu
37 et al. (2017) developed an OpenFOAM-based simulation tool for the fully coupled model of floating
38 wind turbines. The dynamic response and extreme structural response of an integrated floating turbine
39 were investigated numerically in (Li et al., 2018a; Li et al., 2018c). Their integrated concept was based
40 on the combination of a floating wind turbine, two tidal turbines and a wave energy converter.

41 So far, the uniform wind flow is commonly adopted in both numerical and experimental studies of
42 offshore wind turbines, which could simplify the aerodynamic modelling. Nevertheless, the wind field
43 in the natural world is turbulent rather than uniform. A realistic wind field varies with not only time,
44 but also space. Actually, the turbulence effect on the performance of the land-based wind turbine has
45 drawn the attention of researchers. Devinant et al. (2002) measured the aerodynamics of a fixed aerofoil
46 in high turbulence. They revealed the strong dependence of the aerodynamic properties on the
47 turbulence intensity, especially in the angle of attack range corresponding to aerofoil stall. Chamorro et
48 al. (2015) launched an experiment to study the unsteady behaviour of a full-scale 2.5 WM wind turbine
49 in turbulent inflow. A similar relationship was observed by Lee et al. (2018) in the field measurement
50 of a small vertical-axis wind turbine installed on the rooftop of a building. Barthelmie et al. (2007)
51 measured the power losses due to wake-induced turbulence at the Middelgrunden wind farm.
52 Approximate 10% energy losses were observed due to wakes. Based on the field measurement at the
53 Nysted wind farm, Barthelmie and Jensen (2010) also concluded that the energy absorption was
54 strongly dependent on the turbulence intensity. Recently, the offshore wind community begins to realize
55 the importance of inflow turbulence. Li et al. (2018) simulated the power production of a floating wind
56 turbine in full turbulent wind field.

57 In the realistic turbulent wind, the wind speed not only varies with the time (turbulence intensity),
58 but also with the space (coherence and wind shear). Although researches on the turbulence effect have
59 been documented, the above three factors are frequently investigated altogether and their individual
60 effect is not fully understood yet. This study aims to investigate the individual effect of wind shear,
61 turbulence intensity and coherence on the dynamic and structural performances of offshore floating
62 wind turbines. First, the full turbulent wind model will be interpreted in details. Afterwards, aero-hydro-
63 servo coupled analysis is conducted in time-domain to capture the performance of a semisubmersible
64 floating wind turbine with different levels of wind shear, turbulence intensity and coherence.

65 **2. Model Description**

66 A semisubmersible floating wind turbine, namely the OC4 DeepCwind semisubmersible concept
67 (Robertson et al., 2014), is considered here. As shown in Fig. 1, the DeepCwind concept mainly consists
68 of the wind turbine, the supporting structure, and the mooring line system.

69 The wind turbine is the NREL 5WM baseline wind turbine (Jonkman et al., 2009), which
70 incorporates a variable-speed torque controller and a blade pitch controller to regulate the power
71 generation based on the operational state. The diameter of the rotor is 126 m, and hub height is 90 m.
72 Please refer to (Jonkman et al., 2009) for more detailed parameters of the reference wind turbine.



73
74 Fig. 1. DeepCwind floating wind turbine system design (Coulling et al., 2013).

75 The supporting structure is a semisubmersible platform, made up of three main offset columns, one
76 central column, as well as a series of diagonal cross and horizontal bracing components. The main
77 scantlings of the semisubmersible are listed in Table 1.

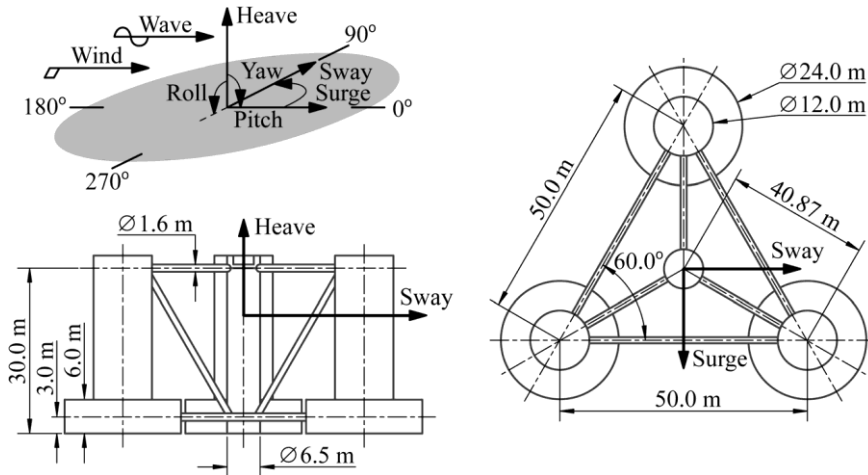


Fig. 2. Main dimensions of the submersible platform (Coulling et al., 2013).

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79

80 Table 1. Main scantlings of the platform.

Term	Value
Draft	20m
Elevation of platform top	10 m
Elevation of offset columns	12 m
Spacing between offset columns	50 m
Length of upper columns	26 m
Length of base columns	6 m
Depth to top of base columns	14 m
Diameter of main column	6.5 m
Diameter of offset (upper) columns	12 m
Diameter of base columns	24 m
Platform mass	13,473,000 kg
Displacement	13,986.8 m ³
Centre of mass	(0 m, 0 m, -13.5 m)
Platform roll inertia	6.827×10 ⁹ kg·m ²
Platform pitch inertia	6.827 ×10 ⁹ kg·m ²
Platform yaw inertia	1.226×10 ¹⁰ kg·m ²

81

82 The floating wind turbine is displaced at sea site with a water depth of 200 m. Three catenary lines
83 are used to hold the platform against sea waves and offshore wind. The three mooring lines are oriented
84 symmetrically at 60°, 180°, and 300° about the vertical axis. Fairleads are connected to the tops of
85 ballast tanks. The relevant properties of mooring lines are outlined in Table 2.

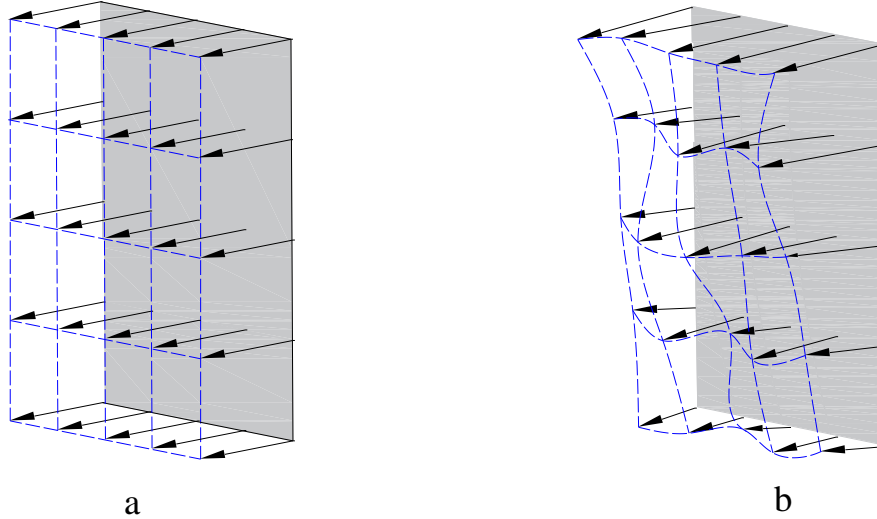
86 Table 2. Properties of mooring line.

Term	Value
Depth to anchor	200 m
Depth to fairlead	14 m
Radius to anchor	837.6 m
Radius to fairlead	40.868 m
Unstretched mooring line length	835.5 m
Mooring line diameter	0.0766 m
Equivalent line mass density	113.35 kg/m
Equivalent mooring line extensional stiffness	753.6 MN

87

88 3. Turbulent Wind Model

89 As shown in Fig. 3, two comparative wind fields are generated. The first one is a steady uniform
90 wind field, where the wind inflow is constant in both time and space scales. The second wind field is
91 turbulent, where the wind inflow varies with not only time but also with space. The second wind field
92 is turbulent, in terms of wind shear, time-scale inflow variation and space-scale inhomogeneity.



93 a
94 Fig. 3. Generated wind fields. (a) steady wind field; (b) turbulent wind field.

95 The wind shear describes how the near-ground wind varies vertically with the height. In the present
96 research, the power-law model is used to represent the wind profile,

97
$$u(z) = u \left(\frac{z}{90} \right)^\alpha \quad (1)$$

98 where α is the exponent parameter. Fig. 4 shows the wind profiles with different values of α . The wind
99 shear becomes more significant when α increases. When α is equal to 0, the wind shear is omitted and
100 the wind field reduces to the uniform one. In the present simulation, $\alpha = 0.15$ is used.

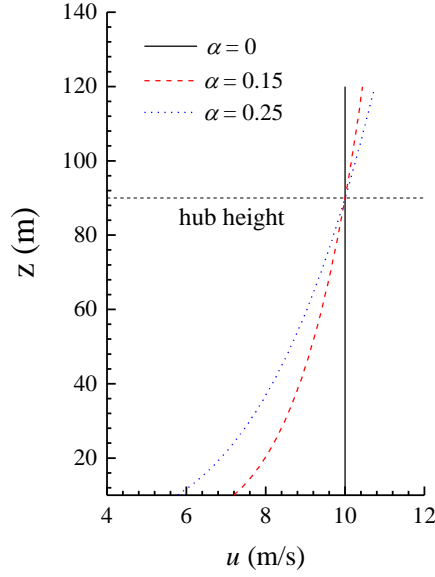


Fig. 4. Wind profile.

The wind inflow is also time-dependent and commonly a spectral method is used to model the time-scale inflow variation. The IEC Kaimal turbulence model (International Electrotechnical Commission, 2015) is used here

$$S(f) = \frac{4\sigma^2 L / u}{(1 + 6f \cdot L / u)^{5/3}} \quad (2)$$

where f is the cyclic frequency and L is an integral scale parameter dependent on the hub height. u is the mean wind speed at hub height. σ can be estimated by the turbulence intensity TI (%)

$$\sigma = \frac{TI}{100} u \quad (3)$$

The turbulence intensity represents the turbulence level, namely how strong the wind varies with time. In the present simulation, the turbulence intensity is set to 10%. Fig. 5 displays the spectra of rated wind ($u = 11.4$ m/s). As shown, the low-frequency oscillations dominate the wind.

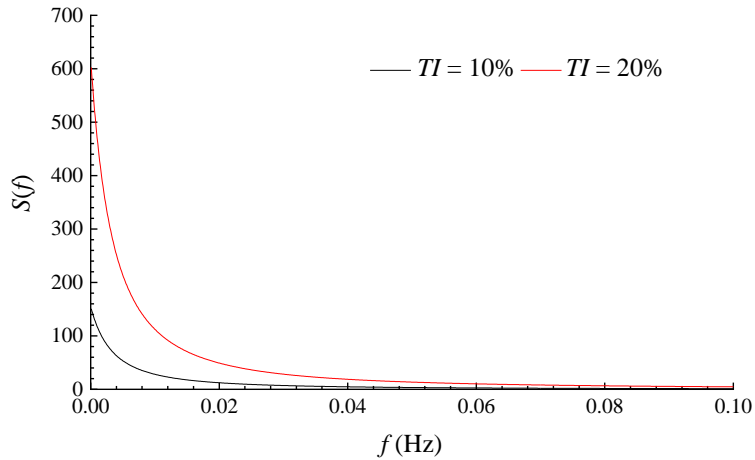


Fig. 5. Spectrum of IEC Kaimal turbulent wind.

115 Apart from the variation in time scale, the wind velocity is also inhomogeneous across the rotor
 116 plane at each time instant. This is due to the phase differences between two points in the wind field so
 117 that the time-scale phase difference leads to the space-scale inhomogeneity. The IEC coherence model
 118 is used to represent the correlation between two arbitrary points.

$$\begin{aligned}
 Coh_{i,j}(f) &= \frac{|S_{i,j}(f)|}{\sqrt{S_i(f)S_j(f)}} \\
 Coh_{i,j}(f) &= \exp \left[-12 \sqrt{\left(\frac{f \cdot r}{u} \right)^2 + \left(0.12 \frac{r}{L_c} \right)^2} \right] \\
 L_c &= \min(60, \text{hub height})
 \end{aligned} \tag{4}$$

120 where $S_{i,j}$ is the cross-spectra defining the correlation of the random wind speed at points i and j , r is the
 121 distance between the two points. L_c is the coherence scale parameter replying on the hub height.

122 4. Dynamic Analysis

123 4.1. Numerical model

124 The aero-hydro-servo-elastic coupled simulation code FAST (Jonkman and Buhl Jr, 2005)
 125 developed by the National Renewable Energy Laboratory (NREL) is used to simulate the dynamic
 126 performance of the DeepCwind floating wind turbine.

127 Assuming that the wave fluid is ideal, the wave-structure is addressed in the framework of potential
 128 flow theory. The wave radiation force is calculated with the convolution term to consider the free
 129 surface memory effect. Since the natural period of horizontal motion of the floating wind turbine is
 130 sufficiently long, second-order drift wave forces are also considered to capture the low-frequency
 131 responses of the floating wind turbine.

132 The blade element momentum (BEM) method is used to compute the wind force acting on the rotor.
 133 The blade is separated into a set of elements, and the interactions between neighbouring elements are
 134 neglected. By seeking the so-called induced velocity, the aerodynamic load on each element is
 135 determined using the lift and drag coefficients of the aerofoil. For an offshore floating wind turbine,
 136 both the platform motions and wind turbulence produce unsteadiness of the inflow seen by the rotor.
 137 The unsteady effect is accounted by the dynamic wake model developed by Minnema (1998), which
 138 can be regarded as a correction to the induced velocity determined by the BEM method.

139 A variable-speed torque controller and a blade pitch controller are incorporated to the wind turbine.
 140 The variable-speed torque controller is active in below-rated operational state. The control algorithm is
 141 to maximize the power output by adjusting the rotor speed while the blade pitch angle is fixed at zero.
 142 One the contrary, the blade-pitch controller works in over-rated state to regulate generator power by
 143 increasing the pitch angle of the blade.

144 The lumped-mass model is used for the dynamics of mooring lines connected to the floating platform.
 145 The mooring line is divided into a set of evenly-sized segments, which are represented by connected
 146 nodes and spring-damper systems. Each segment is divided into two components and the properties are
 147 assigned and lumped to the two nodes at each end of that segment, respectively. The connections
 148 between adjacent nodes are represented by damper-spring systems. Only the axial properties of the
 149 mooring lines are accounted whereas the torsional and bending properties are neglected.

150 4.2. Short-term extreme response

151 The extreme responses are estimated based on the mean up-crossing rate method. In an arbitrary
 152 time interval T , it can be assumed that the random number of up-crossing is approximated by the Poisson
 153 distribution on condition that the up-crossing is statistically independent. Once a level y is selected, the
 154 distribution of extreme value y_{max} for a random signal $y(t)$ is described as

$$155 \quad P(y_{max} \leq y) = \exp\left(-\int_0^T v^+(y, t) dt\right) \quad (5)$$

156 where $v^+(y, t)$ is the up-crossing rate corresponding to level y , which denotes the instantaneous
 157 frequency of the positive slop crossings of the defined level. In this circumstance, the probability of
 158 y_{max} exceeding a defined level y is given by

$$159 \quad \begin{aligned} P(y_{max} > y) &= 1 - \exp(-\hat{v}^+(y)T) \\ \hat{v}^+(y) &= \frac{1}{T} \int_0^T v^+(y, t) dt \end{aligned} \quad (6)$$

160 The mean up-crossing rate $\hat{v}^+(y)$ can be easily obtained from the time series of the signal that is
 161 going to be analysed. For example, if we have k independent realizations of the random process and let
 162 $n_j^+(y, T)$ denote the number of up-crossings in realization j , then the sample-based mean up-crossing
 163 rate is given by

$$164 \quad \begin{aligned} \hat{v}^+(y) &\approx \bar{v}^+(y) \\ \bar{v}^+(y) &= \frac{1}{kT} \sum_{j=1}^k n_j^+(y, T) \end{aligned} \quad (7)$$

165 Eq. (7) is the basic formula to approximate the mean up-crossing rate $\hat{v}^+(y)$ through numerical
 166 simulations. If the defined level y is not very high, then just a few simulation realizations of the random
 167 process will produce satisfactory approximation. Nevertheless, extensive simulations are required to
 168 evaluate the extreme values in the tail region. To save computation resources, the extrapolation method
 169 proposed by Naess and Gaidai (2009) is used in this study to extrapolate the mean up-crossing rate
 170 corresponding to high level y . The extrapolation method is based on the observation of marine structures
 171 so that it is applicable in this study. The mean up-crossing rate is approximated by

$$172 \quad \begin{aligned} \bar{v}^+(y) &\approx v_{fit}^+(y) \\ v_{fit}^+(y) &= q \cdot \exp\{-a(y-b)^c\}, y \geq y_0 \end{aligned} \quad (8)$$

173 where q , a , b and c are all constant values. y_0 is the lower limit of the sampled data used for the
 174 extrapolation. To ensure that the extrapolated rate is reliable, the 95% confidence interval (CI) of the
 175 raw rate is examined:

$$176 \quad CI_{\pm}(y) = \bar{v}^+(y) \pm \frac{1.96s(y)}{\sqrt{k}} \quad (9)$$

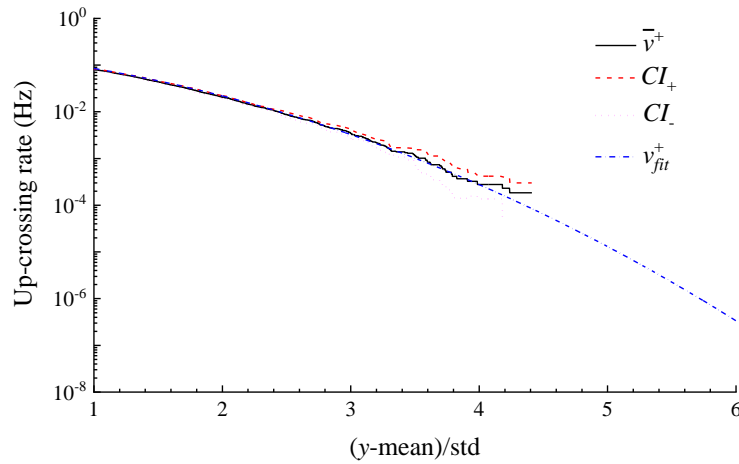
$$s^2(y) = \frac{1}{k-1} \sum_{j=1}^k \left(\frac{n_j^+(y, T)}{T} - \bar{v}^+(y) \right)^2$$

177 In the present research, the extrapolated up-crossing rate is based on 6 independent numerical
 178 realizations ($k = 6$). $y_0 = (\text{mean} + \text{std})$ is used, where ‘mean’ is the average mean response of the 6
 179 numerical realizations; ‘std’ is the average standard deviation of the 6 numerical realizations. Please
 180 note that y_0 varies with the load case so as to ensure the extrapolated rate is within the confidence interval.
 181 To put more emphasis on the more reliable sampled data, the weight factor proposed by Naess and
 182 Gaidai (2009) is used here

$$183 \quad \Theta = \sum_{j=1}^N w_j \left| \log(\bar{v}^+(y_j)) - \log(q) + a(y_j - b)^c \right|^2 \quad (10)$$

184 where Θ is the mean square error; $w_j = \left| \log(CI_+(y_j)) - \log(CI_-(y_j)) \right|^{-2}$ is the weight factor. The least
 185 square optimization method is used to get q , a , b and c by minimizing Θ .

186 Fig. 6 gives an example of the extrapolated up-crossing rate for the tower base bending moment
 187 under rated steady wind. As shown, the extrapolated rate is located within the 95% confidence interval
 188 indicating that the extrapolation is reliable. Hereinafter, the extreme value corresponding to $v_{fit}^+ = 10^{-5}$
 189 will be used to represent the short-term extreme response.



190
 191 Fig. 6. Extrapolation of the up-crossing rate of the tower base bending moment under rated steady wind based on 6
 192 numerical realizations.

193 4.3. Simulation setup

194 The aero-hydro-servo-elastic coupled dynamic analysis is conducted for a set of wind speeds,
 195 covering different operation states of the wind turbine. The wind speed refers to that measured at the

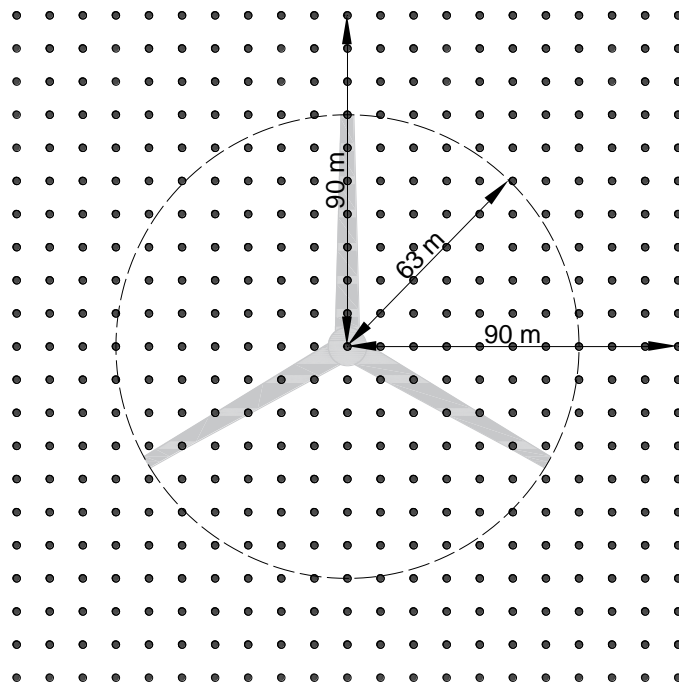
196 hub centre (90 m). The turbulent wind field is generated using the TurbSim (NWTC Information Portal,
 197 2016). Due to the platform motions, the wind turbine moves a lot during the simulation, and a wind grid
 198 with dimension 180 m \times 180 m is generated to cover the movement range of the rotor (see Fig. 7). 441
 199 points (21 \times 21) are uniformly distributed across the wind grid, at which the time-series of wind speed
 200 are generated. Table 3 summaries the environmental conditions considered in the present simulation.

201 For each environmental condition, the simulation runs for a total length of 3800 seconds, and only
 202 the last 1-hour data will be collected to get rid of the transient effects in the early simulation stage. The
 203 simulation time increment is set to 0.0125 s.

204 Table 3. Environmental conditions

	U_w	H_s	T_p
Below-rated	8 m/s	4 m	6 s
Rated	11.4 m/s	5 m	8 s
Over-rated	14 m/s	6 m	10 s

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Fig. 7. Wind grid.

208 5. Simulation Results

209 5.1. Dynamic response in turbulent flow

210 First, the dynamic performance of the reference floating wind turbine in the uniform wind and the
 211 turbulent wind are compared.

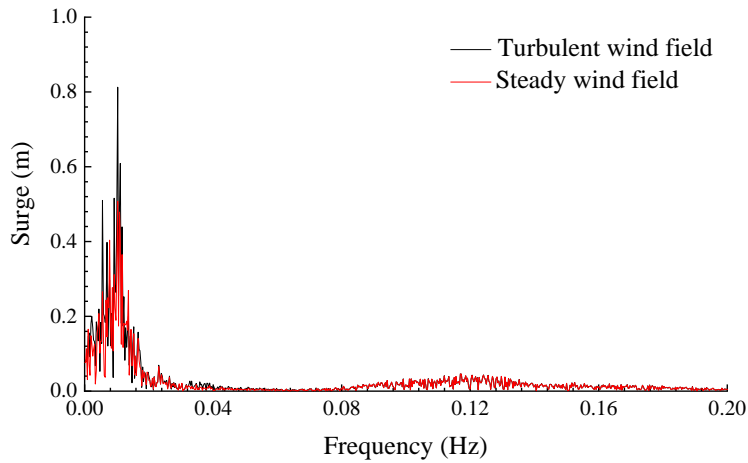
212 Table 4 summaries the standard deviations of the platform global motions under various operational
 213 states (below-rated, rated and over-rated) when the two comparative wind flow conditions are
 214 considered. It is clearly shown that the platform surge and pitch motions are increased substantially,
 215 regardless of the mean wind speed. Since the wind mainly induces horizontal loads, the turbulence has

216 a negligible effect on platform heave motion. Fig. 8 and Fig. 9 display the response feature of surge and
 217 pitch motions, respectively. In the steady wind, the platform motion is mainly induced at wave
 218 frequency range and resonant frequency. In the turbulent inflow, the platform motion is excited a lot at
 219 the resonant frequency whereas the wave frequency motion is hardly varied. It indicates that the
 220 aerodynamic turbulence effect is not effective at all on wave-induced response. Similar phenomenon
 221 has been documented previously in (Hu et al., 2016).

222 Table 4. Standard deviation of platform motions

	Below-rated		Rated		Over-rated	
	Turbulent wind	Steady wind	Turbulent wind	Steady wind	Turbulent wind	Steady wind
Surge	1.11	0.94	1.41	1.12	1.18	0.90
Heave	0.06	0.06	0.13	0.13	0.29	0.28
Pitch	0.32	0.21	0.57	0.40	0.72	0.54

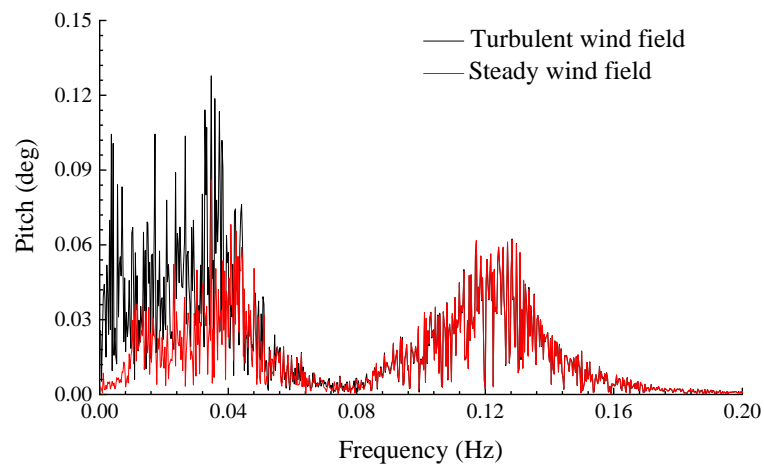
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Fig. 8. FFT analysis of platform surge motion.



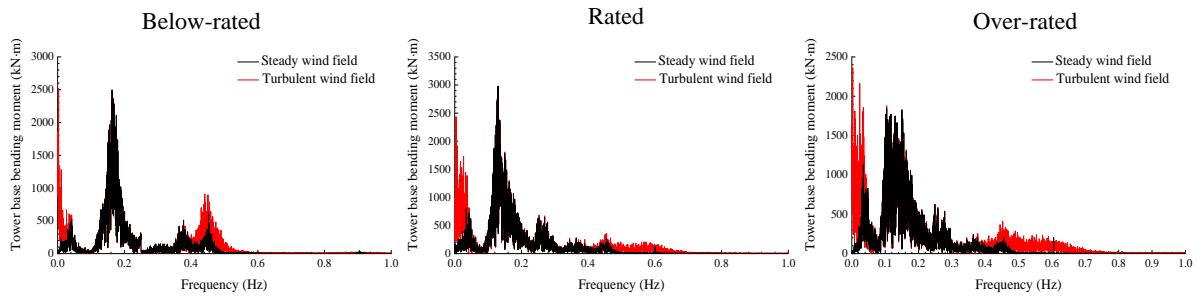
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Fig. 9. FFT analysis of platform pitch motion, rated operation state.

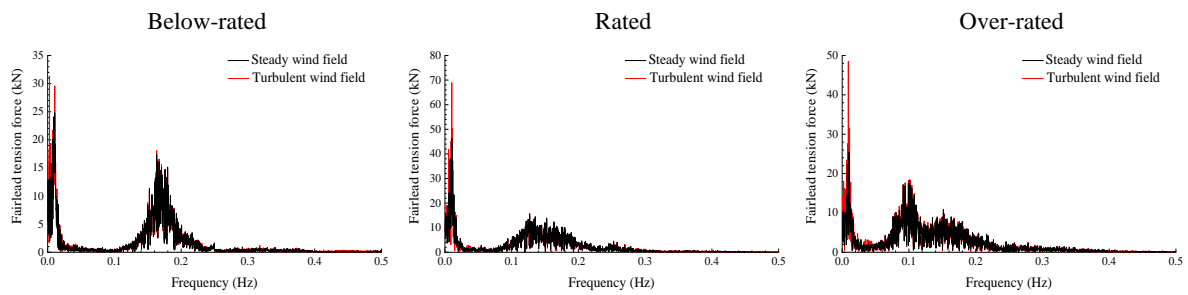
228 Apart from platform motions, the effect of turbulent inflow on structural loads is also considerable.
 229 Fig. 10 plots the FFT analysis result of the tower base fore-aft bending moment in the three operation
 230 states. Although the wave-frequency response is generally independent from the turbulent inflow, the
 231 low-frequency and the high-frequency are quite sensitive to the turbulent inflow. In all three operation

232 scenarios, the bending moment response at low-frequency range is excited the most. It is attributed to
 233 the dominating slow-varying inflow in the turbulent wind field (see Fig. 5). In the meanwhile, the blades
 234 experience wind speed variation over the rotation process due to the spatial inhomogeneity, leading to
 235 the high-frequency range response. Similar conclusions can be drawn from the fairlead tension force,
 236 which is displayed in Fig. 11. The aerodynamic loads are not applied to the fairlead directly but through
 237 the platform movement. Therefore, high-frequency fairlead tension force response is not observed since
 238 the platform's natural frequency is sufficiently low.



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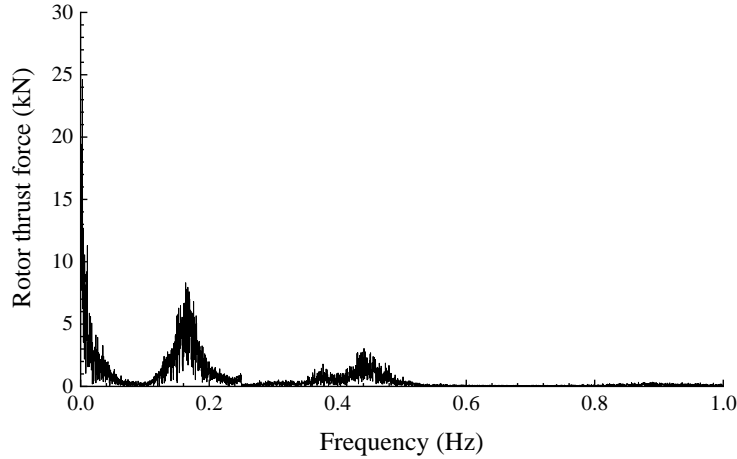
Fig. 10. FFT analysis of tower base fore-aft bending moment.



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Fig. 11. FFT analysis of fairlead tension force.

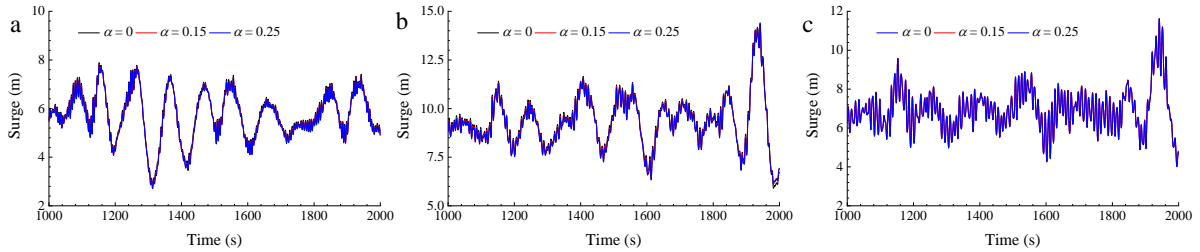
243 Since the floating wind turbine is subject to identical wave excitations in the steady wind and the
 244 turbulent wind, the amplitude of platform motions and structural loads could be purely attributed to the
 245 wind force. Fig. 12 illustrates the response character of rotor thrust force. The majority of response
 246 energy is located within the low-frequency range, mainly induced by the turbulence intensity. Besides,
 247 the response is also observed around 0.44 Hz, namely the 3P frequency of the rotor speed. The 3P
 248 frequency response is induced by the spatial inhomogeneity of the wind field since the wind speed seen
 249 by the blade experiences variation during the rotation process. Two aspects contribute to the spatial
 250 inhomogeneity. The first one is the wind shear, representing how the wind inflow varies vertically with
 251 the height. Secondly, the phase lag between two points in the rotor plane also leads to the space-scale
 252 inflow variation, and it is represented by the coherence model. In the following part, the individual
 253 effect of wind shear, turbulence intensity and coherence will be clarified.



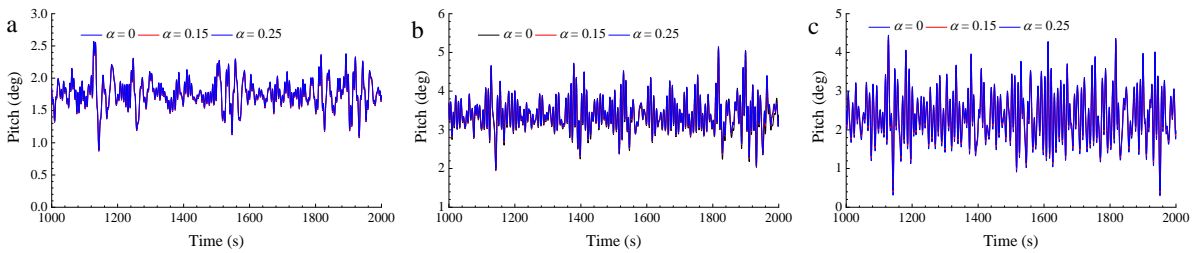
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255 Fig. 12. FFT analysis result of thrust force in turbulent wind field, below-rated operation state.

256 **5.2. Wind shear effect**

257 Fig. 13 and Fig. 14 plot the platform surge and pitch motions, respectively. It appears that the
258 platform motions are not sensitive to the wind shear at all. Despite that the wind shear exponent
259 increases from 0 to 0.25, the responses of surge and pitch motions are hardly varied. When a blade is
260 experiencing the high wind velocity region (up half of the rotor plane), the other two blades are within
261 low wind velocity region (down half of the rotor plane). In this case, the resultant thrust force induced
262 by the three blades remains relatively stable. Consequently, the wind shear has a negligible influence
263 on the platform motions.

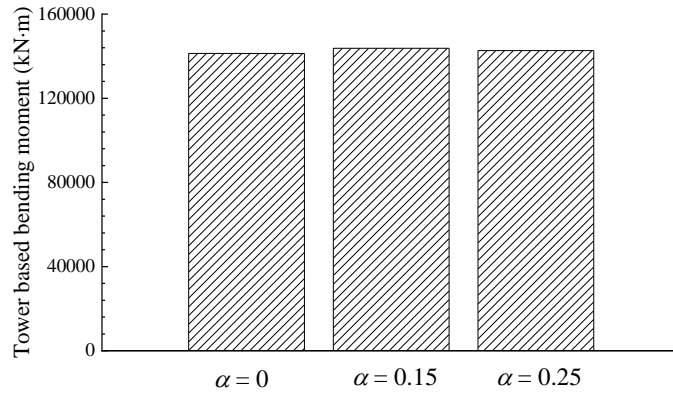


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265 Fig. 13. Time series of platform surge motions. (a) below-rated condition; (b) Rated condition; (c) Over-rated condition.



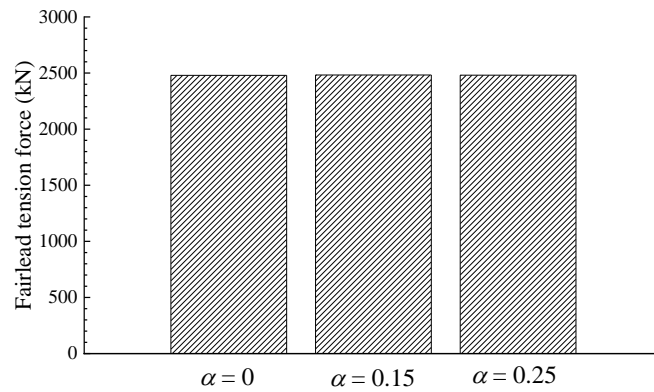
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267 Fig. 14. Time series of platform pitch motions. (a) below-rated condition; (b) Rated condition; (c) Over-rated condition.

268 Fig. 15 and Fig. 16 demonstrate the extreme tower base bending moment and the extreme fairlead
269 tension in the presence of wind shear, respectively. As discussed above, the rotor thrust varies hardly
270 with the wind shear, and thereby the extreme tower base fore-aft bending moment and the extreme
271 fairlead tension remains nearly unchanged regardless of the wind shear.



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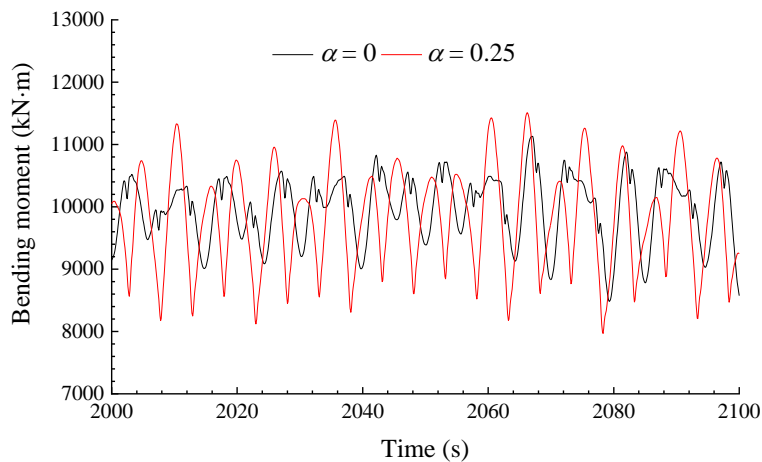
Fig. 15. Extreme tower base fore-aft bending moment, rated operation condition.



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Fig. 16. Extreme fairlead tension, rated operation condition.

276 Although the platform motions and the structural loads at tower base and fairlead are not sensitive
 277 to the wind shear, the local loads at blade root depend strongly on the wind shear. According to the time
 278 series plotted in Fig. 17, the out-of-plane bending moment at blade root becomes quite unstable in the
 279 presence of large wind shear. As explained before, the blade will experience high-speed and low-speed
 280 region alternately due to the wind shear. Although the resultant force of the three blades remains stable,
 281 the load applied on each blade varies violently.



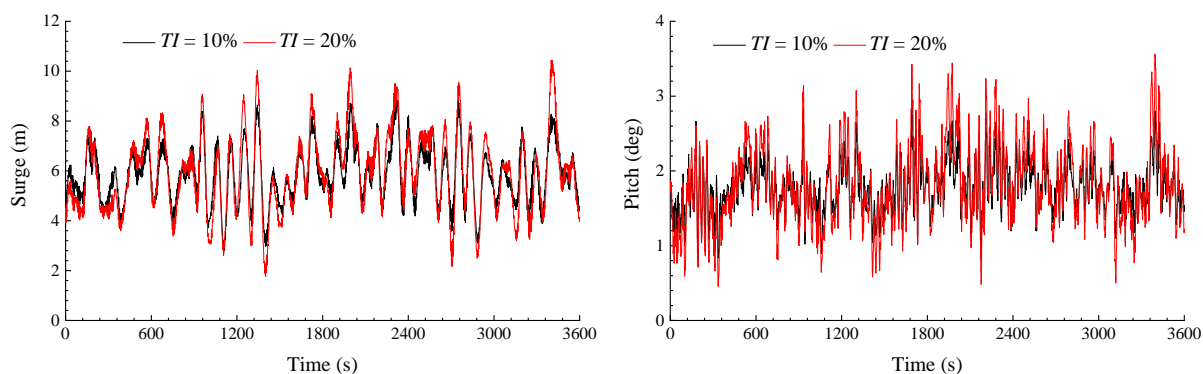
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Fig. 17. Times series of blade root bending moment.

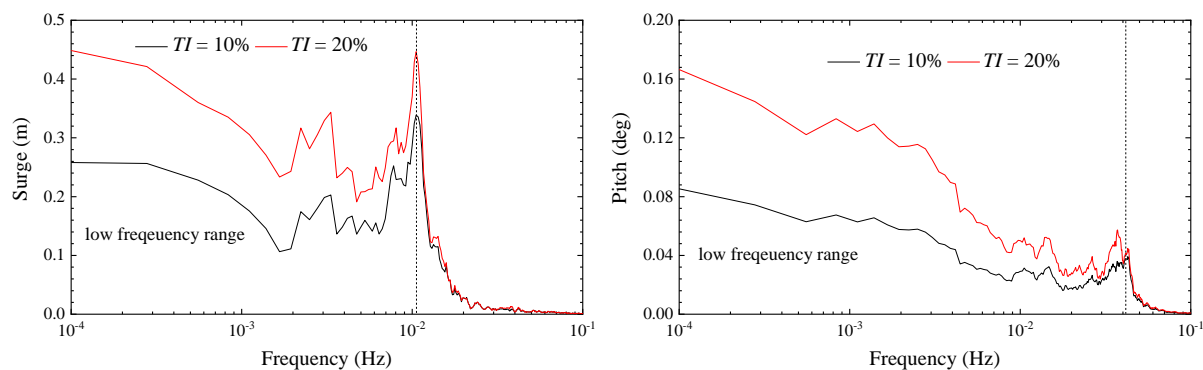
284 5.3. Turbulence intensity effect

285 The performances of the floating wind turbine with different turbulence intensities (10% and 20%)
 286 are examined in this sub-section. The law exponent parameter α is set to zero to eliminate the wind
 287 shear effect.

288 Fig. 18 plots the time series of platform surge and pitch motions under below-rated wind. It is easy
 289 to identify that the platform motions become increasingly violent when the turbulence intensity
 290 increases to 20%. To interpret the turbulence intensity effect more clearly, the time series of platform
 291 motions are analyzed with the FFT method and the results are presented in Fig. 19. The turbulence
 292 intensity effect is only observed within the low-frequency region (lower than the resonant frequency).
 293 In the right side of the vertical dash line (representing the resonant frequency), the two curves match
 294 well. When the turbulence intensity increases, the resonant response is somewhat amplified. Moreover,
 295 the quasi-static response at very low frequency range (10^{-4} Hz \sim 10^{-3} Hz) is further induced.



296
 297 Fig. 18. Time series of platform motions, below-rated.



298
 299 Fig. 19. FFT analysis results of platform motions, below-rated.

300 In addition to the platform motions, the turbulence intensity also has observable consequences for
 301 structural loads. According to Table 5, the extreme responses of the tower base bending moment and
 302 the fairlead tension increase considerably with the turbulence intensity.

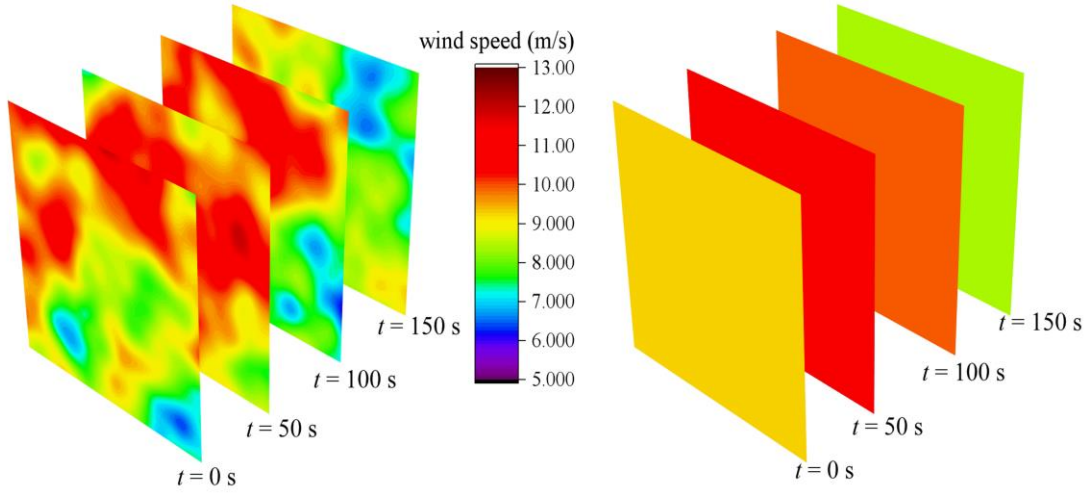
303 Table 5. Short-term extreme responses

	$TI = 10\%$	$TI = 20\%$
Tower base bending moment (kN·m)	1.09×10^5	1.19×10^5
Fairlead tension force (kN)	2159	2462

304

305 5.4. Coherence structure effect

306 In addition to the wind shear and the turbulence intensity, the blade also experiences inflow variation
 307 due to the inhomogeneity of the wind field and the inhomogeneity is caused by the coherence (phase
 308 difference). We adopt the hub centre as the reference point. An unsteady uniform wind field, in which
 309 the phase difference between any point and the reference point is zero, is generated to illustrate the
 310 coherence structure effect. As shown in Fig. 20, the wind speed is not uniformly distributed in the
 311 spatially coherent wind field at each time instant due to the phase lag. On the contrary, a completely
 312 coherent air inflow is uniformly distributed across the space at each time instant since the phase
 313 difference between any two points are zero. Of course, the wind speed also varies with time since the
 314 time-scale turbulence is not eliminated. Please note that, at each time instant, the reference point wind
 315 speeds of the two comparative wind fields are equal.



316
 317 Fig. 20. How the spatial distribution of wind speed varies with time. Left: partially coherent wind field; right: completely
 318 coherent wind field.

319 In the unsteady uniform wind field, any two points are completely coherent whereas a realistic
 320 turbulent wind field is partially coherent. To demonstrate the coherence structure quantitatively, the
 321 proper orthogonal decomposition (POD) method is used to decompose the above two wind fields. In
 322 the present research, 21×21 points are uniformly distributed across the space so that we have total 441
 323 measurements of wind speed at each time instant. We are to decompose 1000 s of wind inflow time
 324 series with the time step being 0.1 s. Consequently, the overall wind flow data $\mathbf{U} = [\mathbf{u}_1, \mathbf{u}_2, \dots, \mathbf{u}_N]$ has
 325 a dimension of 441×10000 (the 21×21 points have been re-organized). The auto-covariance matrix of
 326 \mathbf{U} is

$$327 \quad \mathbf{R} = \mathbf{U}^T \cdot \mathbf{U} \quad (11)$$

328 Then we have the orthonormal eigenvectors $\mathbf{G} = [\mathbf{g}_1, \mathbf{g}_2, \dots, \mathbf{g}_N]$ and the eigenvalue matrix \mathbf{A}

$$329 \quad \mathbf{R}\mathbf{G} = \mathbf{R}\mathbf{A} \quad (12)$$

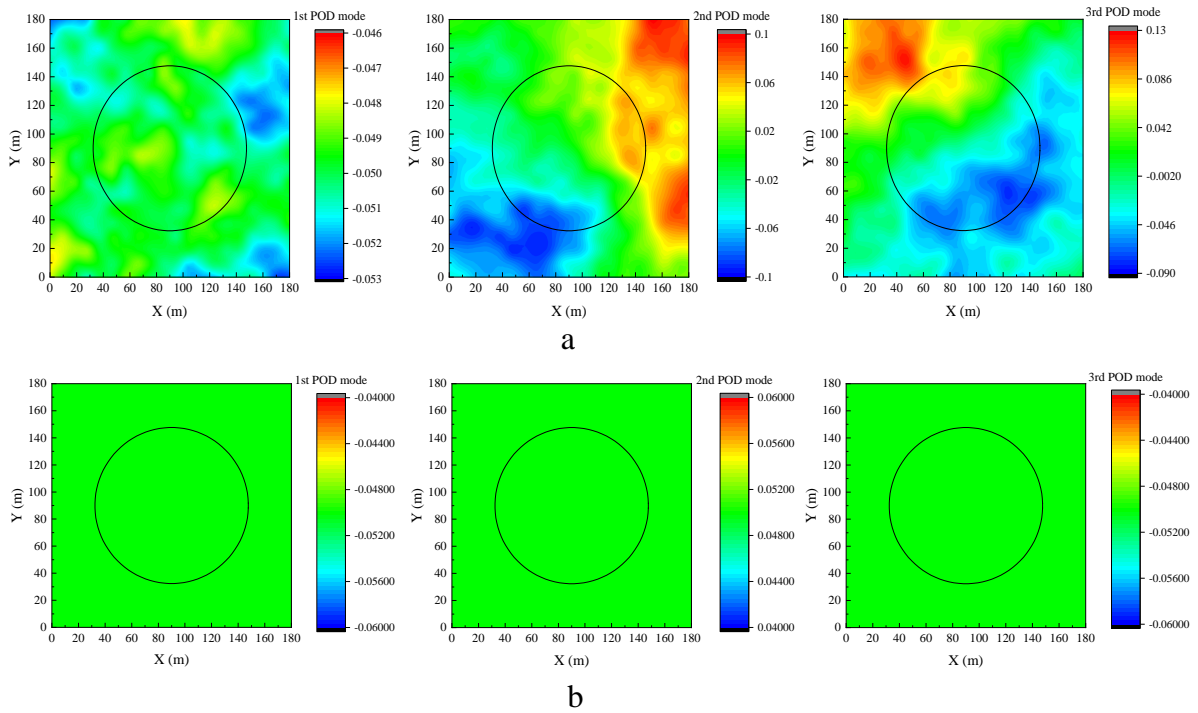
330 Please note that eigenvalue matrix \mathbf{A} should be re-organized if necessary, to satisfy

331
$$A_{11} \geq A_{22} \geq \dots \geq A_{NN} \quad (13)$$

332 The normalized POD modes (the orthonormal basis) is then given by

333
$$\phi_k = \frac{\mathbf{g}_k \cdot \mathbf{u}_i}{\|\mathbf{g}_k \cdot \mathbf{u}_i\|} \quad (14)$$

334 The first three POD modes on based 1000 seconds of numerical realizations are displayed in Fig. 21,
 335 where the two wind fields exhibit distinctive coherence structures. For the partially coherent wind field,
 336 the first POD mode is characteristic of a single major coherent structure. In higher order modes, more
 337 patterns are observed and the inhomogeneity becomes more significant. However, all the POD modes
 338 of the completely coherent wind field are exactly flat. Apparently, the coherence structure is to have an
 339 influence on the performance of the platform wind turbine.



340 Fig. 21. Normalized POD modes of the wind field. (a) partial coherence; (b) complete coherence.
 341

342 Fig. 22 illustrates the coherence structure effects on the platform motions. As shown, the pitch
 343 motion increases substantially when the points in the wind field become completely coherent with each
 344 other. According to the FFT analysis result, the resonant platform motion is amplified implying that the
 345 floating wind turbine is subject to more aerodynamic loads.

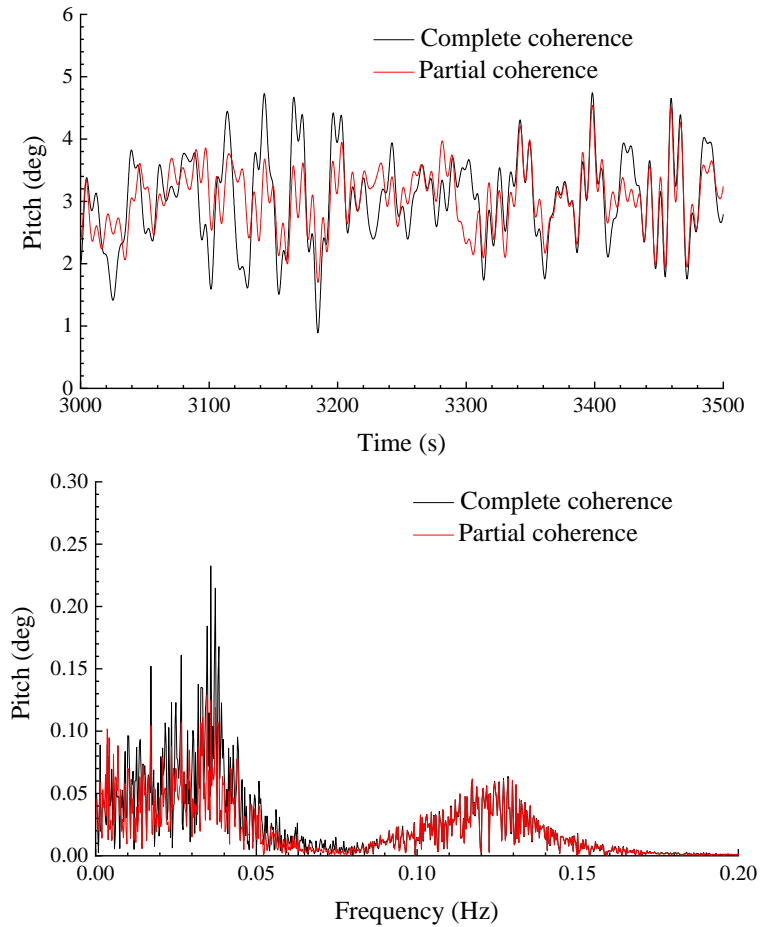


Fig. 22. Platform pitch motion, rated operation condition.

346
347

348 At first sight, this conclusion appears contradictory since the spatial inhomogeneity is totally
349 removed when all the points in the wind field are completely coherent. To interpret the underlying

350 philosophy, we plot the average wind speed $\bar{u} = \sum_{i=1}^N u_i / N$ of the wind grid in Fig. 23 (There are $N =$

351 421 points in the wind field). As shown, when all the points in the wind field are completely coherent,

352 the average wind inflow seems to become more ‘turbulent’. In partially coherent wind field (see Fig.

353 20, left side), the wind speed at each point oscillates around the mean level of 11.4 m/s. Due to the

354 phase difference, the instantaneous speeds at some points are higher than 11.4 m/s whereas lower than

355 the mean level at others. In this circumstance, the variation of the average wind speed seen by the rotor

356 is relieved due to the phase lag at different points. Consequently, the average wind speed across the

357 rotor plane is close to 11.4 m/s, although it also varies with time. On the contrary, the wind speeds at

358 all points are in phase with each other if the wind field is completely coherent. If the speed at the

359 reference point (hub centre) exceeds 11.4 m/s, then the speeds at all other point are higher than 11.4

360 m/s as well. Apparently, the inflow seen by the rotor is more unstable in the presence of complete

361 coherence. The result is that the rotor will be subject to more violent aerodynamic loads.

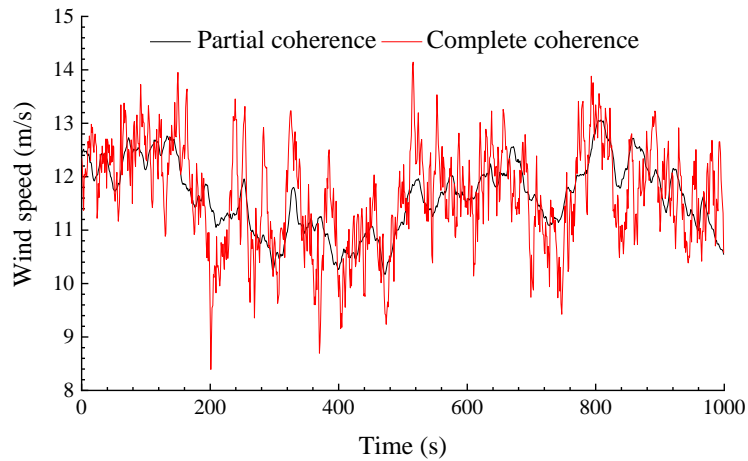


Fig. 23. Average wind speed seen by the rotor, rated operation condition.

362
363

364 6. Conclusions

365 This work aims to investigate how offshore floating wind turbines react to the wind shear, the
366 turbulence intensity and the coherence structure of a turbulent wind field.

367 The platform is not sensitive to the wind shear since the resultant aerodynamic loads applied on the
368 three blades will not differ much. Nevertheless, the structural load at each individual blade becomes
369 unstable in the wind shear. During the rotation process, a blade experiences low-speed and high-speed
370 regions alternately due to the wind shear. Consequently, load at each blade is excited at the 1P frequency.

371 The turbulence intensity effect on the floating wind turbine is quite considerable. According to the
372 wind spectrum, the air inflow mainly carries low-frequency components and thereby the low-frequency
373 platform motion responses are excited. Moreover, the response is not sensitive to the turbulence
374 intensity at high frequency range (higher than the resonant frequency). Regarding the structural loads,
375 the floating wind turbine is more likely to exceed the limit state with high turbulence intensity.

376 The coherence structure of the turbulent wind field is interpreted quantitatively using the proper
377 orthogonal decomposition method. Although a completely coherent wind field removes the spatial
378 inhomogeneity, the time-scale variation increases. In a partially coherent wind field, the average wind
379 inflow seen by the rotor is more stable. Therefore, the floating wind turbine is safer in a partially
380 coherent wind field.

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