

Kettering University Digital Commons @ Kettering University

Physics Publications

Physics

3-2009

Introducing Students to Raman Spectroscopy as a Research Tool

Bahram Roughani *Kettering University*

David Warner *Kettering University*

Uma Ramabadran *Kettering University*

Follow this and additional works at: https://digitalcommons.kettering.edu/physics_facultypubs Part of the <u>Physics Commons</u>

Recommended Citation

Roughani, Bahram; Warner, David; and Ramabadran, Uma, "Introducing Students to Raman Spectroscopy as a Research Tool" (2009). *Physics Publications*. 39. https://digitalcommons.kettering.edu/physics_facultypubs/39

This Conference Proceeding is brought to you for free and open access by the Physics at Digital Commons @ Kettering University. It has been accepted for inclusion in Physics Publications by an authorized administrator of Digital Commons @ Kettering University. For more information, please contact digitalcommons@kettering.edu.

Introducing Students to Raman Spectroscopy as a Research Tool

Bahram Roughani (Physics Professor & Dept Head)

David Warner (Undergraduate Student) & Uma Ramabadran (Physics faculty)



MRS Spring Meeting- San Francisco, CA – April 2, 2013

Outline

- Introduction
- Raman Spectroscopy Overview
- Raman Spectroscopy of Si (100)
 - Model: Polarization Analysis \$ Raman Intensity
 - Measurements: Experimental Si Peak Intensity
- Discussion and Summary
- Questions

In Light Scattering "Sky is the limit"

Raman Scattering





http://www.doitpoms.ac.uk/tlplib/raman/printall.php

Raman Scattering Process





Anti-Stokes event

Kinematic

Stokes event

 $\hbar \omega_{s} = \hbar \omega_{I} - \hbar \omega \quad \text{(Energy)} \quad \hbar \omega_{As} = \hbar \omega_{I} + \hbar \omega$ $\hbar k_{s} = \hbar k_{I} - \hbar q \quad \text{(Momentum)} \quad \hbar k_{As} = \hbar k_{I} + \hbar q$



IR Active – Raman Inactive Change in Dipole Moment Sign

Raman Active - IR Inactive Change in Polarizability Sign

http://www.doitpoms.ac.uk/tlplib/raman/printall.php



Si (100) Intensity: Theoretical Model: 0° Rotation

$$R_{x} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & b \\ 0 & b & 0 \end{bmatrix}, \quad R_{y} = \begin{bmatrix} 0 & 0 & b \\ 0 & 0 & 0 \\ b & 0 & 0 \end{bmatrix}, \quad R_{z} = \begin{bmatrix} 0 & b & 0 \\ b & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$
$$I \propto \begin{bmatrix} e_{i} \ R_{j} \ e_{s} \end{bmatrix}^{2} \quad \begin{array}{c} \text{e.g., for x, y} \ I \ \text{Backscattering}} I \propto \begin{bmatrix} 1 & 0 & 0 \end{bmatrix} \begin{bmatrix} 0 & b & 0 \\ b & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \\ 0 \end{bmatrix}^{2} = b^{2}$$

Polarization	Phonon modes		
Configuration	(direction of atomic displacement)		
	LO (z-direction)	TO (x-direction)	TO (y-direction)
Z (x, x) Z	0	0	0
Z (x, y) Z	b²	0	0
Z (y, y) Z	0	0	0

Where; x = < 1 0 0 >, y = < 0 1 0 >, and z = < 0 0 1 >

Si (100) Intensity: Theoretical Model: 45° Rotation

$$R_{x} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & b \\ 0 & b & 0 \end{bmatrix}, \qquad R_{y} = \begin{bmatrix} 0 & 0 & b \\ 0 & 0 & 0 \\ b & 0 & 0 \end{bmatrix}, \qquad R_{z} = \begin{bmatrix} 0 & b & 0 \\ b & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

Polarization Configuration	Phonon modes (direction of atomic displacement)			
	LO (z-direction)	TO (x-direction)	TO (y-direction)	
Z (x', x') Z	b²	0	0	
Z (x', y') Z	0	0	0	
Z (y', y') Z	b²	0	0	

Where

 $I \propto \left[e_i \ R_j \ e_s \right]^2$

here
$$x' = \frac{1}{\sqrt{2}} < 1 \ 1 \ 0 >$$
, $y' = \frac{1}{\sqrt{2}} < 1 \ \overline{1} \ 0 >$, and $z' = < 0 \ 0 \ 1 >$

General Configuration – Back Scattering Geometry

- θ Angle between <001> & the scattered wave vector \mathbf{k}_{s} ,
- ϕ Angle between <100> & the projection of \mathbf{k}_{s} onto the (001)
- Ψ_{o} Angle between the x-axis of laboratory coordinate and the projection of the <001> axis onto the x-y plane,
- Ψ Angle between the x-axis and unit polarization vectors \mathbf{e}_i
- Ψ ' Angle between the x-axis and unit polarization vectors \mathbf{e}_{s} ,



Incident polarization vector (e_i) & Scattered light polarization vector (e_s)

$$e_i^x = \cos\theta \ \cos \phi \ \cos[\psi + \psi_0] - \sin\phi \ \sin[\psi + \psi_0]$$
(3a)

$$e_i^y = \cos\theta \ \sin \phi \ \cos[\psi + \psi_0] + \cos\phi \ \sin[\psi + \psi_0]$$
(3b)

$$e_i^z = -\sin\theta \ \cos[\psi + \psi_0]$$
(3c)

And

$$e_s^x = \cos\theta \ \cos \phi \ \cos[\psi' + \psi_0] - \sin\phi \ \sin[\psi' + \psi_0]$$
(4a)

$$e_s^y = \cos\theta \ \sin \phi \ \cos[\psi' + \psi_0] + \cos\phi \ \sin[\psi' + \psi_0]$$
(4b)

$$e_s^z = -\sin\theta \ \cos[\psi' + \psi_0]$$
(4c)

Raman Intensity Calculations follows:

 $I \propto \begin{bmatrix} e_i & R_j & e_s \end{bmatrix}^2$ $R_x = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & b \\ 0 & b & 0 \end{bmatrix}, \qquad R_y = \begin{bmatrix} 0 & 0 & b \\ 0 & 0 & 0 \\ b & 0 & 0 \end{bmatrix}, \qquad R_z = \begin{bmatrix} 0 & b & 0 \\ b & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$

Si (100) Raman Tensors R_x, R_y & R_z

Simplified Version – Back Scattering Geometry

- θ Angle between <001> & the scattered wave vector \mathbf{k}_{s} ,
- ϕ Angle between <100> & the projection of \mathbf{k}_{s} onto the (001)
- Ψ_{o} Angle between the x-axis of laboratory coordinate and the projection of the <001> axis onto the x-y plane,
- Ψ Angle between the x-axis and unit polarization vectors \mathbf{e}_i
- Ψ' Angle between the x-axis and unit polarization vectors \mathbf{e}_{s} ,

```
Simplified version assuming:

Lab z-axis || < 001 >

Thus,

\theta = 0

\phi = 0,

And, \Psi_0 = 0

With \Psi and \Psi' as variable angles
```



General Case:

$$e_i^x = \cos\theta \ \cos\phi \ \cos[\psi + \psi_0] - \sin\phi \ \sin[\psi + \psi_0]$$
(3a)

$$e_i^y = \cos\theta \ \sin\phi \ \cos[\psi + \psi_0] + \cos\phi \ \sin[\psi + \psi_0]$$
(3b)

$$e_i^z = -\sin\theta \ \cos[\psi + \psi_0]$$
(3c)

And

$$e_s^x = \cos\theta \ \cos\phi \ \cos[\psi' + \psi_0] - \sin\phi \ \sin[\psi' + \psi_0]$$
(4a)

$$e_s^y = \cos\theta \ \sin\phi \ \cos[\psi' + \psi_0] + \cos\phi \ \sin[\psi' + \psi_0]$$
(4b)

$$e_s^z = -\sin\theta \ \cos[\psi' + \psi_0]$$
(4c)

Simplified Case:

$$e_{i}^{x} = \cos[\psi]$$

$$e_{i}^{y} = \sin[\psi] \quad \text{Or}, \quad e_{i} = \begin{bmatrix} \cos\psi \\ \sin\psi \\ 0 \end{bmatrix} \quad (5)$$

$$e_{i}^{z} = 0$$
And,
$$e_{s}^{x} = \cos[\psi']$$

$$e_{s}^{y} = \sin[\psi'] \quad \text{Or}, \quad e_{s} = \begin{bmatrix} \cos\psi' \\ \sin\psi' \\ 0 \end{bmatrix} \quad (6)$$

$$e_{s}^{z} = 0$$

$$e_{i} = \begin{bmatrix} \cos \psi \\ \sin \psi \\ 0 \end{bmatrix}$$
When Fixing
The Analyzer
Along The
$$e_{i} = \begin{bmatrix} \cos \psi \\ \sin \psi \\ 0 \end{bmatrix}$$
, and, $e_{s} = \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix}$

$$e_{s} = \begin{bmatrix} \cos \psi' \\ \sin \psi' \\ 0 \end{bmatrix}$$
X-direction

TO Mode:
$$I \propto [e_i \ R_x \ e_s]^2 = \begin{bmatrix} [\cos\psi \ \sin\psi \ 0] \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & b \\ 0 & b & 0 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix}^2 = 0$$
 (7)
TO Mode: $I \propto [e_i \ R_y \ e_s]^2 = \begin{bmatrix} [\cos\psi \ \sin\psi \ 0] \begin{bmatrix} 0 & 0 & b \\ 0 & 0 & 0 \\ b & 0 & 0 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix}^2 = [b \ \cos\psi]^2$ (8)
LO Mode: $I \propto [e_i \ R_z \ e_s]^2 = \begin{bmatrix} [\cos\psi \ \sin\psi \ 0] \begin{bmatrix} 0 & b & 0 \\ b & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix}^2 = 0$ (9)



Three dimensional theoretical Raman intensity plot for backscattered signal from single crystalline Si from (100) plane is shown based on Raman intensity calculations.



Three dimensional experimental Raman intensity plot for backscattered signal from single crystalline Si from (100) plane is shown based on Raman measurements.

Similarly For backscattering along <110> we have:



Raman tensors for x_2 , y_2 , and z_2 axes are:

$$R_{x_2} = \begin{bmatrix} 0 & b & 0 \\ b & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}, \qquad R_{y_2} = \begin{bmatrix} 0 & 0 & -b \\ 0 & 0 & b \\ -b & b & 0 \end{bmatrix}, \qquad R_{z_2} = \begin{bmatrix} 0 & 0 & b \\ 0 & 0 & b \\ b & b & 0 \end{bmatrix}$$

Raman Scattering rules for Backscattering from <110> Si:

Polarization	Phonon Modes			
Configuration				
	TO (x ₂)	TO(y ₂)	LO(z ₂)	
$z_2(x_2,x_2)\overline{z_2}$	0	0	0	
$z_2(x_2, y_2)\overline{z_2}$	0	b ²	0	
$z_2(y_2,y_2)\overline{z_2}$	\mathbf{b}^2	0	0	
	TO (X'2)	TO(y'2)	LO(Z'2)	
$\mathbf{z}_2(\mathbf{x}'_2,\mathbf{x}'_2)\overline{\mathbf{z}_2}$	$\frac{2}{3}b^2$	$\frac{1}{3}b^2$	0	
$\mathbf{z}_2(\mathbf{x'}_2,\mathbf{y'}_2)\overline{\mathbf{z}_2}$	$\frac{1}{3}b^2$	0	0	
$z_2(y'_2,y'_2)\overline{z_2}$	0	$\frac{4}{3}b^2$	0	
1				

Discussion and Summary:



Close agreement between theory and experiment Same analysis could be done for other Si samples; (110), (111), etc. Same analysis could be done for any single crystalline samples

An effective approach for introducing students to Raman Spectroscopy applications in scientific research!





Angle dependence of first order Raman intensity measurements for the (100) single crystals Si



Model

Experiment

Thank you for your attention!

"For the things we have to learn before we can do them, we learn by doing them." — Aristotle