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Deep exclusive π^+ electroproduction off the proton at CLAS

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Abstract. The exclusive electroproduction of π^+ above the resonance region was studied using the CEBAF Large Acceptance Spectrometer (CLAS) at Jefferson Laboratory by scattering a 6 GeV continuous electron beam off a hydrogen target. The large acceptance and good resolution of CLAS, together with the high luminosity, allowed us to measure the cross section for the $\gamma^*p \rightarrow n\pi^+$ process in 140 (Q^2, x_B, t) bins: $0.16 < x_B < 0.58$, $1.6 \text{ GeV}^2 < Q^2 < 4.5 \text{ GeV}^2$ and $0.1 \text{ GeV}^2 < -t < 5.3 \text{ GeV}^2$. For most bins, the statistical accuracy is on the order of a few percent. Differential cross sections are compared to four theoretical models, based either on hadronic or on partonic degrees of freedom. The four models can describe the gross features of the data reasonably well, but differ strongly in their ingredients. In particular, the model based on Generalized Parton Distributions (GPDs) contain the interesting potential to experimentally access transversity GPDs.

1 Introduction

One of the major challenges in contemporary nuclear physics is the study of the transition between hadronic and partonic pictures of the strong interaction. At asymptotically short distances, the strong force is actually weak and the appropriate degrees of freedom are the quarks and gluons (partons) whose interaction can be quantified very precisely by perturbative Quantum Chromodynamics (pQCD). However, at larger distances on the order of one Fermi, effective theories that take hadrons as elementary particles whose interactions are described by the exchange of mesons appear more applicable. The connection between these two domains is not well understood. In order to make progress, a systematic study of a series of hadronic reactions probing these intermediate distance scales is necessary. The exclusive electroproduction of a meson (or of a photon) from a nucleon, $\gamma^*N \rightarrow N'M$, is particularly interesting. Indeed, it offers two ways to vary the scale of the interaction and therefore to study this transition regime. One can vary the virtuality of the

incoming photon $Q^2 = -(p_e - p'_e)^2$, which effectively represents the transverse size of the probe, or the momentum transfer to the nucleon $t = (p_N - p'_N)^2$, which effectively represents the transverse size of the target. Here, p_e and p'_e are the initial and scattered electron four-momenta and p_N and p'_N are the initial and final nucleon four-momenta, respectively. Figure 1 sketches the transition regions that have been experimentally explored until now (lightly shaded areas) as a function of these two variables, Q^2 and $|t|$. In this figure, we keep, quite arbitrarily, only the experiments for which $|t| > 3 \text{ GeV}^2$ in photoproduction (SLAC [1,2] and JLab [3]) and $Q^2 > 1.5 \text{ GeV}^2$ in electroproduction (Cornell [4,5], JLab [6–8] and HERMES [9]). These are the domains for which, we believe, there are chances to observe first signs that partonic degrees of freedom play a role in the reactions. The darkly shaded area in fig. 1 represents the phase space covered by the present work. It is divided into 140 $(Q^2, x_B \text{ or } W, t)$ bins, to be compared to only a few $(Q^2, x_B \text{ or } W, t)$ bins in the lightly shaded areas for the previous electroproduction experiments.

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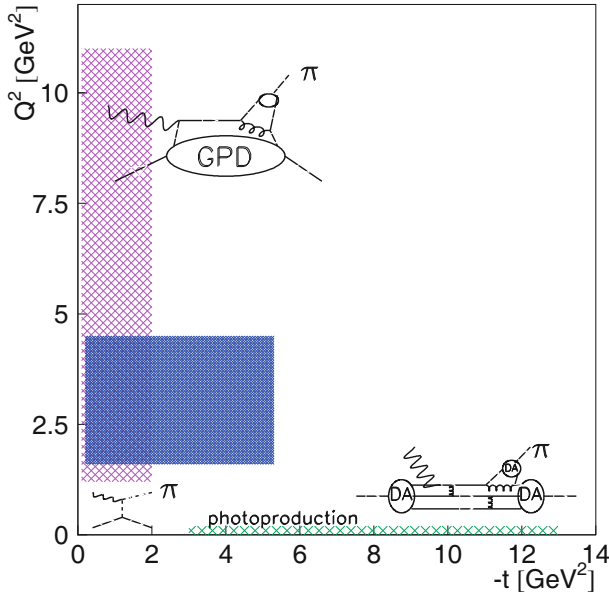


Fig. 1. (Color online) Schematic representation of the $\gamma^* N \rightarrow N' \pi$ process (above the resonance region) in different regions of the (Q^2, t) plane. The Feynman diagrams describe the reaction in terms of meson exchanges at low Q^2 and $|t|$, in terms of GPDs at large Q^2 and small $|t|$, and in terms of hadron distribution amplitudes (DA) at large $|t|$. The lightly shaded areas (magenta and green online) represent approximately the experimentally explored regions up to now. The darkly shaded area (blue online) represents the phase space covered by this work.

We also display in fig. 1 three Feynman-type diagrams illustrating the mechanisms believed to be at stake for the $\gamma^* N \rightarrow N' M$ process: at asymptotically large Q^2 , asymptotically large $|t|$ (both in terms of partonic degrees of freedom) and at low Q^2 and low $|t|$ (in terms of hadronic degrees of freedom).

At asymptotically large Q^2 and small $|t|$ (along the vertical axis in fig. 1), the exclusive electroproduction of a meson should be dominated by the so-called “handbag diagram” [10–15]. The initial virtual photon hits a quark in the nucleon and this same quark, after a single gluon exchange, ends up in the final meson. A QCD factorization theorem [15] states that the complex quark and gluon non-perturbative structure of the nucleon is described by the Generalized Parton Distributions (GPDs). For the π^+ channel at leading twist in QCD, *i.e.* at asymptotically large Q^2 , the longitudinal part of the cross section σ_L is predicted to be dominant over the transverse part σ_T . Precisely, $d\sigma_L/dt$ should scale as $1/Q^6$ at fixed x_B and $|t|$, while $d\sigma_T/dt$ should scale as $1/Q^8$. It is predicted that σ_L is sensitive to the helicity-dependent GPDs \tilde{E} and \tilde{H} [15] while, if higher-twist effects are taken into account and factorization is assumed, σ_T is sensitive to the transversity GPDs, H_T and $\bar{E}_T = 2\tilde{H}_T + E_T$ [16].

At large values of $|t|$, in photoproduction (*i.e.* along the horizontal axis in fig. 1) but also presumably in electroproduction, the $\gamma^{(*)} N \rightarrow N' M$ process should be dominated by the coupling of the (virtual) photon to one of the va-

lence quarks of the nucleon (or of the produced meson), with minimal interactions among the valence quarks. In this regime, a QCD factorization theorem states that the complex structure of the hadrons is described by distribution amplitudes (DA) which at small distances (large $|t|$) can be reduced to the lowest Fock states, *i.e.* 3 quarks for the nucleon and $q\bar{q}$ for the meson [17]. At sufficiently high energy, constituent counting rules (CCR) [18–20] predict an s^{-7} scaling of the differential cross section $d\sigma/dt$ at fixed center-of-mass pion angles, provided $|s|$, $|t|$, and $|u|$ are all large. Here $s = W^2$ is the squared invariant mass of the $\gamma^* p$ system and $u = (p_\gamma^* - p'_N)^2$ is given in terms of the four-vectors $p_\gamma^* = p_e - p'_e$ and p'_N for the final-state nucleon. The large $|t|$ and $|u|$ region corresponds typically to a center-of-mass pion angle $\theta_{\text{cm}} \approx 90^\circ$. In this domain, the CCR predict $d\sigma/dt = f(\theta_{\text{cm}})s^{2-n}$ for the energy dependence of the cross section, where $f(\theta_{\text{cm}})$ depends on details of the dynamics of the process and n is the total number of point-like particles and gauge fields in the initial and final states. For example, our reaction $\gamma^* p \rightarrow n\pi^+$ should have $n = 9$, since there is one initial photon, three quarks in the initial and the final nucleons, and two in the final pion.

Many questions are open, in particular at which Q^2 and s do such scaling laws start to appear. Even if these respective scaling regimes are not reached at the present experimentally accessible Q^2 and s values, can one nevertheless extract GPDs or DAs, provided that some corrections to the QCD leading-twist mechanisms are applied? Only experimental data can help answer such questions.

2 Insights from previous experiments with respect to partonic approaches

The two most recent series of experiments that have measured exclusive π^+ electroproduction off the proton, in the large- Q^2 , low- $|t|$ regime where the GPD formalism is potentially applicable, have been conducted in Hall C at Jefferson Lab (JLab) [6–8] and at HERMES [9].

The Hall C experiments, with 2 to 6 GeV electron beam energies, separated the σ_L and σ_T cross sections of the $\gamma^* p \rightarrow n\pi^+$ process using the Rosenbluth technique for $0.17 < x_B < 0.48$ and Q^2 up to 3.91 GeV^2 . The term σ_L dominated the cross section for $|t| < 0.2 \text{ GeV}^2$, while σ_T was dominant for larger $|t|$ values. These data were compared to two GPD-based calculations, hereafter referred to as VGG [21] and GK [16, 22] from the initials of the models’ authors. The comparison of the data with the VGG model can be found in the Hall C publications [6, 7] while the comparison with the GK model can be found in the GK publications [16, 22]. For σ_L , which should be the QCD leading-twist contribution, these GPD calculations were found to be in general agreement with the magnitude and the Q^2 and t dependencies of the experimental data. In these two calculations the main contribution to σ_L stems from the \tilde{E} GPD, which is modeled either entirely as pion-exchange in the t -channel [21] or is at least dominated by it [16, 22] (see refs. [23, 24] for the connection

between the t -channel pion-exchange and the \tilde{E} GPD). This term is also called the “pion pole”, and the difference between the two calculations lies in the particular choice made for the t -channel pion propagator (Reggeized or not) and the introduction of a hadronic form factor or not at the πNN vertex. In both calculations, σ_L contains higher-twist effects because the pure leading-twist component of the pion pole largely underestimates the data. Only the GK model, which explicitly takes into account higher-twist quark transverse momentum, is able to calculate σ_T . Agreement between data and calculation is found only if the H_T transversity GPD is introduced, which makes up most of σ_T .

The HERMES experiment used 27.6 GeV electron and positron beams to measure the $\gamma^*p \rightarrow n\pi^+$ cross section at four (x_B, Q^2) values, with x_B ranging from 0.08 to 0.35 and Q^2 from 1.5 to 5 GeV². Since all data were taken at a single beam energy, no longitudinal/transverse separation could be carried out. The differential cross section $d\sigma/dt$ was compared to the same two GPD models mentioned above. The GK model, which calculates both the longitudinal and transverse parts of the cross section, displays the same feature as for the lower energy JLab data, *i.e.* a dominance of σ_L up to $-t \approx 0.2$ GeV², after which σ_T takes over. The sum of the transverse and longitudinal parts of the cross section calculated by the GK model is in very good agreement with the data over most of the t range measured at HERMES [16,22]. The VGG model, which calculates only the longitudinal part of the cross section, is in agreement with the data only for low t values [9]. Again, in both calculations, σ_L is dominated by the \tilde{E} GPD, modeled essentially by the pion pole term, and σ_T , in the GK model, is due to the transversity GPDs. The HERMES experiment also measured the transverse target spin asymmetry A_{UT} for the $\gamma^*p \rightarrow n\pi^+$ process, which indicate [16,22] that the transversity GPDs H_T or \tilde{E}_T indeed play an important role in the process, confirming the approach of the GK group.

The comparison between the JLab Hall C and HERMES experiments and the two GPD-based calculations yields very encouraging signs that, although higher-twist contributions definitely play a major role, these data can be interpreted in terms of GPDs, in particular transversity GPDs. More precise and extensive data would be highly useful to confirm these findings. Firstly, the present CLAS experiment extends somewhat the (x_B, Q^2) phase space previously covered by the JLab Hall C experiments and secondly, it covers 20 (x_B, Q^2) bins (with statistical errors of a few percent on average) which doubles the number of bins of the JLab Hall C experiments (and triples the HERMES number of bins). These new data are important to test the present GPD-based model calculations and, if successful, bring more stringent constraints on the current GPD parametrizations.

The large- $|t|$ (large- $|u|$) domain, where the DA formalism is asymptotically applicable for $\gamma^{(*)}p \rightarrow n\pi^+$, has so far been explored only in high-energy photoproduction at SLAC [1,2] and intermediate-energy photoproduction at JLab [25,26]. While the SLAC data tend to follow the s^{-7}

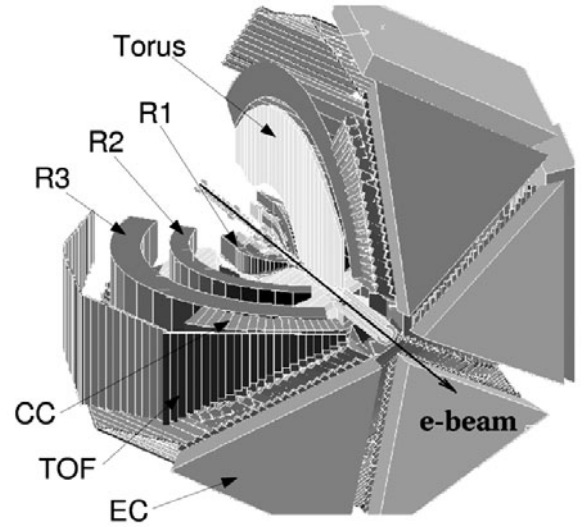


Fig. 2. Three-dimensional view of the CLAS detector.

scaling asymptotic prediction, for a 90° center-of-mass angle, the more recent JLab data, which are compatible with the SLAC data but are more precise, actually reveal some large oscillations around this s^{-7} behavior.

In recent years a similar trend, *i.e.* “global” scaling behavior, has been observed in deuteron photodisintegration experiments [27–33], and also in hyperon photoproduction [34]. It would be interesting to see this in exclusive pion electroproduction and if so, whether the oscillations disappear as Q^2 increases. The measurement presented in this article is the first one to explore this large- $|t|$, large- $|u|$ domain ($\theta_{\text{cm}} \approx 90^\circ$) for $\sqrt{s} > 2$ GeV in π^+ exclusive electroproduction off the proton. The present CLAS electroproduction experiment covers a t -range up to ≈ 5 GeV² while the largest $|t|$ values measured by Hall C are ≈ 0.9 GeV² and by HERMES ≈ 2 GeV².

3 The experiment

The measurement was carried out with the CEBAF Large Acceptance Spectrometer (CLAS) [35]. A schematic view of CLAS is shown in fig. 2. CLAS has a toroidal magnetic field generated by six flat superconducting coils (main torus), arranged symmetrically around the azimuth. Six identical sectors are independently instrumented with 34 layers of drift cells for particle tracking (R1, R2, R3), plastic scintillation counters for time-of-flight (TOF) measurements, gas threshold Cherenkov counters (CC) for electron and pion separation, and electromagnetic calorimeters (EC) for photon and neutron detection. To aid in electron/pion separation, the EC is segmented into an inner part closer to the target and an outer part further away from the target. CLAS covers on average 80% of the full 4π solid angle for the detection of charged particles. The azimuthal acceptance is maximum at a polar angle of 90° and decreases at forward angles. The polar angle coverage ranges from about 8° to 140° for the detection of π^+ . The scattered electrons are detected in the CC and EC, which extend from 8° to 45°.

The target is surrounded by a small toroidal magnet (mini-torus). This magnet is used to shield the drift chambers closest to the target from the intense low-energy electron background resulting from Møller scattering.

A Faraday cup, composed of 4000 kg of lead and 75 radiation lengths thick, is located in the beam dump, ≈ 29 meters downstream the CLAS target. It completely stops the electrons and thus allows to measure the accumulated charge of the incident beam and therefore the total flux of the beam [35].

The specific experimental data set “e1-6” used for this analysis was collected in 2001. The incident beam had an average intensity of 7 nA and an energy of 5.754 GeV. The 5 cm long liquid-hydrogen target was located 4 cm upstream of the CLAS center. This offset of the target position was found to optimize the acceptance of forward-going positively charged particles. The main torus magnet was set to 90% of its maximum field, corresponding to an integral magnetic field of ≈ 2.2 Tm in the forward direction. The torus current during the run was very stable ($< 0.03\%$). Empty-target runs were performed to measure contributions from the target cell windows.

In this analysis, the scattered electron and the produced π^+ were detected and the final state neutron determined from missing mass. The continuous electron beam provided by CEBAF is well suited for measurements involving two or more final-state particles in coincidence, leading to very small accidental coincidence contributions, smaller than 10^{-3} , for the instantaneous luminosity of $10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ of the present measurement.

Raw data were subjected to the calibration and reconstruction procedures that are part of the standard CLAS data analysis sequence. Stringent kinematic cuts were applied to select events with one electron candidate and only one positively charged track. These events were then subjected to further selection criteria described in the following Section. Throughout the analysis, the experimental data distributions were compared to the output of our Monte Carlo program GSIM (see sect. 4).

A schematic illustration of electron scattering off a nucleon target producing an outgoing nucleon and one pion is shown in fig. 3. The scattered electron angle θ_e is given in the laboratory frame. The angle between the virtual photon three-momentum and the direction of the pion is denoted as θ_π^* and the angle between the electron scattering plane and hadronic production plane is denoted as ϕ_π^* . These two angles are defined in the center-of-mass frame of the hadronic system. The angle ϕ_π^* is defined so that the scattered electron lies in the $\phi_\pi^* = 0^\circ$ half plane with the z -axis pointing along the virtual photon momentum. For exclusive single π^+ production from the proton, we request the simultaneous detection of one single electron and of one single π^+ in CLAS and the final state neutron will be identified by the missing mass squared $((p_e + p_N) - (p'_e + p_\pi))^2$, where p_π is the four-momentum of the detected π^+ . The kinematic range and bin sizes are chosen to provide reasonable statistics in each bin. These are summarized in table 1.

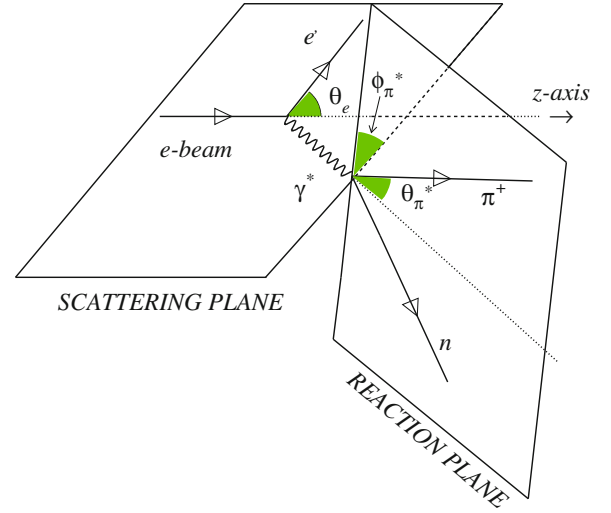


Fig. 3. Kinematics of exclusive single π^+ electroproduction from a proton target.

Table 1. Kinematic bins used in this analysis.

Variable	Number of bins	Range	Bin size
x_B	7	0.16–0.58	0.06
Q^2	5	1.6–3.1 GeV ²	0.3 GeV ²
	3	3.1–4.5 GeV ²	0.5 GeV ²
$-t$	6	0.1–1.9 GeV ²	0.3 GeV ²
	3	1.9–4.3 GeV ²	0.8 GeV ²
	1	4.3–5.3 GeV ²	1.0 GeV ²

Our aim is to extract the three-fold differential cross section $\frac{1}{T} \frac{d^3\sigma}{dQ^2 dx_B dt}$, where

$$\frac{1}{T} \frac{d^3\sigma}{dQ^2 dx_B dt} = \frac{n_w(Q^2, x_B, t)}{\mathcal{L}_{int} \Delta Q^2 \Delta x_B \Delta t}, \quad (1)$$

with

- $n_w(Q^2, x_B, t)$ is the weighted number of $ep \rightarrow e'n\pi^+$ events in a given bin (Q^2, x_B, t) ; in particular, $n_w(x_B, Q^2, -t)$ contains the detector's acceptance correction factor $\text{Acc}(x_B, Q^2, -t, \phi_\pi^*)$ (see sect. 5.1) and the correction factor due to radiative effects $F_{rad}(x_B, Q^2, -t)$ (see sect. 5.2);
- \mathcal{L}_{int} is the effective integrated luminosity;
- ΔQ^2 , Δx_B and Δt are the corresponding bin widths (see table 1); for bins not completely filled, because of W or E' cuts on the electron, for instance (see fig. 8), the phase space $\Delta Q^2 \Delta x_B \Delta t$ includes a volume correction and the Q^2 and x_B central values are modified accordingly.

In the following three sections, we detail the various cuts and correcting factors entering the definition of $n_w(Q^2, x_B, t)$.

4 Data analysis

4.1 Particle identification and event selection

4.1.1 Electron identification

The electrons are identified at the trigger level by requiring at least 640 MeV energy deposited in the EC in coincidence with a signal in the CC (which triggers on one photoelectron).

Additional requirements for particle identification (PID) were used in the off-line analysis to refine the electron identification. First, we required that the EC and CC hits matched with a reconstructed track in the drift chambers (DC). Second, we correlated the energy deposited in the EC and the momentum obtained by the track reconstruction in the DC. This is aimed at removing the pion contamination. Electrons deposit energy in proportion to their incident energy in the calorimeter whereas pions are minimum ionizing and deposit a constant amount of energy in the calorimeter. The ratio of the total deposited energy in the EC to the momentum of the particle is called the sampling fraction. For electrons, approximately 30% of the total energy deposited in the EC is directly measured in the active scintillator material. The remainder of the energy is deposited in the lead sheets interleaved between the scintillators. Figure 4 shows the sampling fraction E/p_e versus particle momentum p_e . The average sampling fraction for electrons was found to be 0.291 for this experiment. The solid lines in fig. 4 show the $\pm 3\sigma$ sampling fraction cuts used in this analysis.

To further reject pions, we required the energy deposited in the inner EC to be larger than 50 MeV. Minimum ionizing particles lose less than this amount in the 15 cm thickness of the inner EC.

Fiducial cuts were applied to exclude the EC detector edges. When an electron hit is close to an edge, part of the shower leaks outside the device; in this case, the energy cannot be fully reconstructed from the calorimeter information alone. This problem can be avoided by selecting only those electrons lying inside a fiducial volume within the EC that excludes the edges. A GEANT-based simulation (GSIM) was used to determine the EC-response with full electron energy reconstruction. The calorimeter fiducial volume was defined by cuts that excluded the inefficient detector regions.

Particle tracks were reconstructed using the drift chamber information, and each event was extrapolated to the target center to obtain a vertex location. We demanded that the reconstructed z -vertex position (distance along the beam axis from the center of CLAS, with negative values indicating upstream of the CLAS center) lies in the range $-80 \text{ mm} < z_{\text{vtx}} < -8 \text{ mm}$. This is slightly larger than the target cell size in order to take into account the resolution effects on the vertex reconstruction.

Finally, a lower limit on the number of photoelectrons detected in the photomultiplier tubes of the CC provided

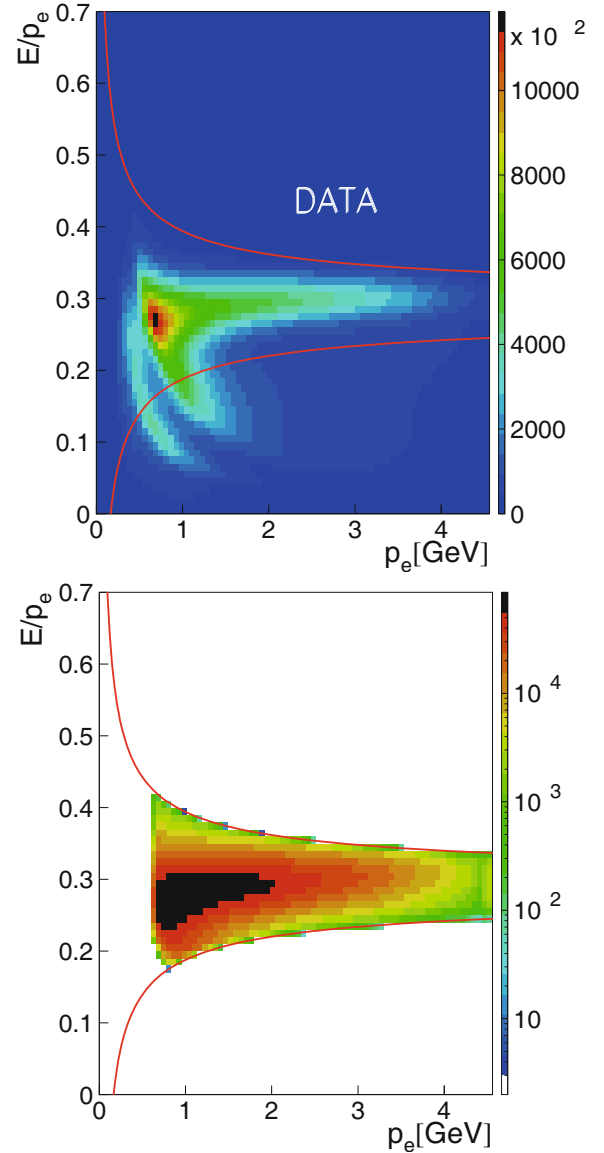


Fig. 4. (Color online) EC sampling fraction versus particle momentum for the experimental data before (top) and after (bottom) EC energy cuts. The solid curves show the $\pm 3\sigma$ sampling fraction cuts which are applied to select electrons.

an additional cut to improve electron identification. The number of photoelectrons detected in the CC follows a Poisson distribution modified for irregularities in light collection efficiency for the individual elements of the array. For this experiment, a good electron event was required to have 3 or more photoelectrons detected in the CC. The efficiency of the CC cut was determined from the experimental data. We fit the number of photoelectrons using the modified Poisson distribution. The efficiency range after the CC cut is 78% to 99% depending on the kinematic region. The correction is then the integral below the cut divided by the total integral of the resulting fit function.

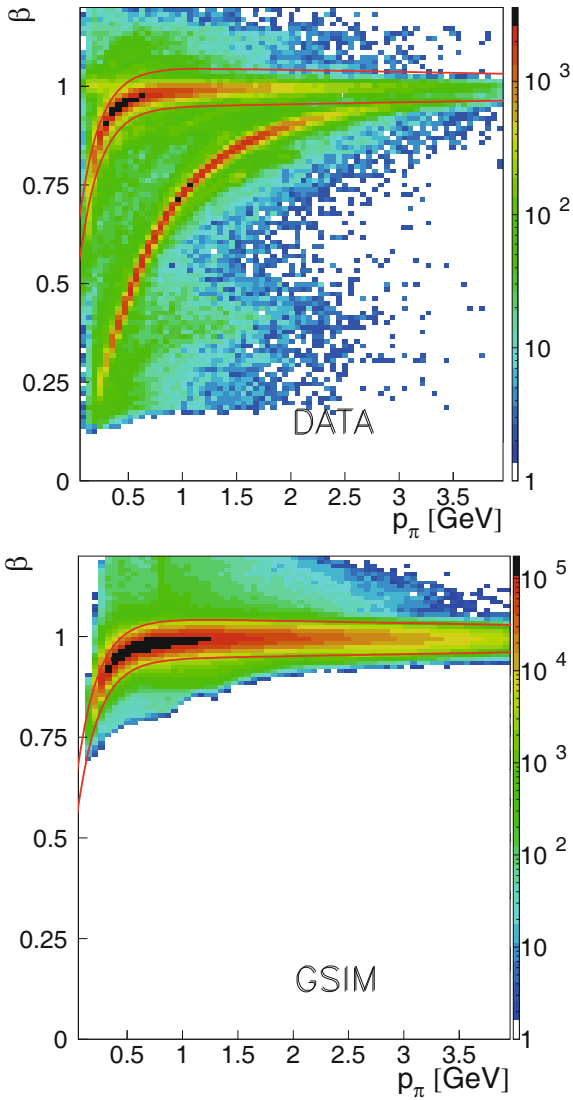


Fig. 5. (Color online) Velocity β versus momentum for π^+ candidates using experimental data (top) and the GSIM Monte Carlo simulation (bottom). The solid curves are $\pm 1.5\sigma$ β cut lines used to select π^+ from positron ($\beta = 1$ band) and proton ($\beta < 0.8$ band) backgrounds. (The data are issued from a skimmed file which selected one electron and at least one positively charged particle and one neutral particle).

4.1.2 Positively charged pion identification

The main cuts to select the π^+ are based on charge, z -vertex, fiducial cuts and velocity versus momentum correlations. The velocity β is calculated from the ratio of the path length of the reconstructed track, to the time of flight.

Figure 5 shows the β versus p distribution for positively charged particles from experimental data (top) and from the GSIM Monte Carlo simulation (bottom). A Gaussian is fit to β for bins in momentum p_π . A $\pm 1.5\sigma$ cut on β is chosen for pion candidates as shown in fig. 5 (solid curves in the plot). Pions and positrons ($\beta = 1$) are well separated below $p_\pi = 250$ MeV/c of momentum in the experimental data, but this is no longer the case at momenta

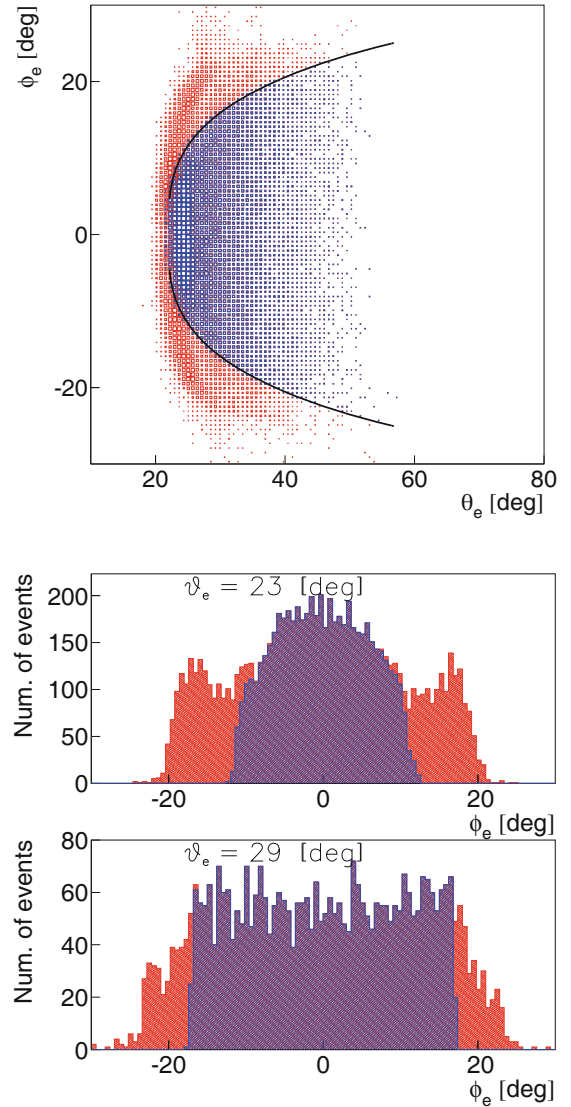


Fig. 6. (Color online) Example of electron fiducial cuts for the electron momentum bin ($p_e = 1.437 \pm .025$ GeV) in Sector 2. See the detailed explanation in the text.

larger than 400 MeV/c. For this reason, positrons can be mis-identified as pions, which increases the background. At higher momenta, there can also be some particle mis-identification from protons and kaons. We estimated that the missing mass and vertex cuts reduce this mis-identification to the 5–10% level. This residual background contamination was subtracted as described in sect. 6.

4.2 Fiducial cuts

4.2.1 Electron fiducial cuts

The fiducial cuts for electrons were developed to exclude regions with non-uniform detector efficiency such as the edges of a sector in the CC and EC. The fiducial cut is a function of the angles θ_e , ϕ_e , and momentum p_e of the electron. An example of such fiducial cut can be seen in fig. 6 for a given electron momentum bin. The solid line in

the top plot shows the boundary of the fiducial region for the central momentum in that bin. Only electron events inside the curve (blue area) were used in the analysis. This curve was determined by selecting the flat high-efficiency areas in the θ_e -sliced ϕ_e distributions. The histograms on the bottom of fig. 6 show examples of such ϕ_e distributions at two values of $\theta_e = 23^\circ \pm 0.5^\circ$ and $29^\circ \pm 0.5^\circ$. One sees a central, uniform area, flanked by two fringes. The highlighted area in the center indicates the selected fiducial range. In addition, a set of θ_e versus p_e cuts was used to eliminate the areas with low detection efficiency due to problematic time-of-flight counters, photomultiplier tubes in Cherenkov counters, or drift chamber areas.

4.2.2 Pion fiducial cuts

The fiducial cuts for pions depend on the angles θ_π , ϕ_π and the momentum p_π . The pion momentum is scanned in 100 MeV steps from 0.3 to 1.7 GeV. The uniform detector efficiency region was determined by selecting a flat high-efficiency ϕ_π region in each θ_π -sliced momentum bin, and the bad TOF counters and the inefficient DC areas were excluded by additional software cuts (the same procedure as was applied to electrons). Figure 7 shows an example for the fiducial cuts for pions. The low-efficiency DC regions (between the black solid lines) and the bad TOF paddles (between red solid lines) are removed in both experimental (top) and simulated (bottom) data as part of the fiducial cuts.

4.3 Kinematic corrections

Due to effects that are not included in the reconstruction software (deviations of the magnetic field from perfect toroidal symmetry, misalignment of the tracking system, ...), we have to apply some empirical corrections to the measured angles and momenta of both electrons and pions. For electrons, the kinematic corrections are applied using the elastic $ep \rightarrow e'p'$ process for which the kinematics is overconstrained. The goal is to correct the three-momentum of the electron so as to minimize the constraints due to the equations of conservation of energy and momentum. The same procedure is applied to the π^+ 's three-momentum using our reaction $ep \rightarrow e'\pi^+n$ under study, minimizing the deviation of the missing mass peak position from the neutron mass. The same correction factors are used for all events having the same kinematics. In this way we keep the spatial resolution of the drift chamber systems and multiple scattering effects and the missing mass resolution approaches its intrinsic limitations. The corrections were most sizable ($\approx 5\%$) for the pion momentum. They resulted in an improved missing mass resolution, from 23 to 35 MeV depending on kinematics. The corrections were most sizable for the high-momentum and forward-angle pions at high W which are of interest in this experiment. We then applied additional *ad hoc* smearing factors for the tracking and timing resolutions to the Monte Carlo so that they match the experimental data.

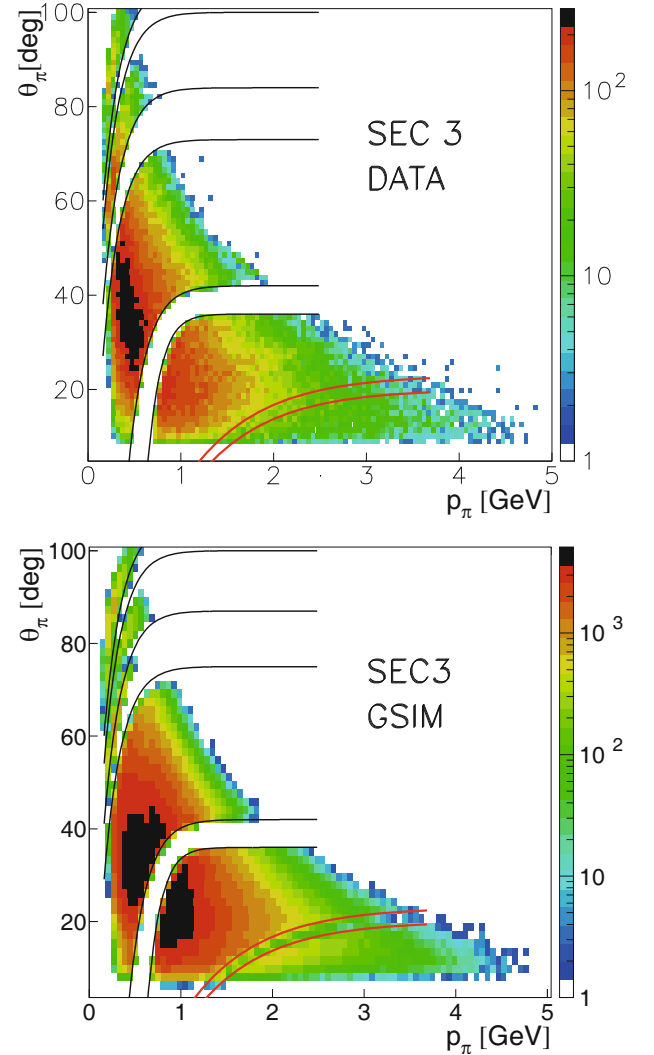


Fig. 7. (Color online) Pion polar angular distribution as a function of momentum in Sector 3. The low detector response areas are removed by empirical cuts for experimental (top) and simulated data (bottom). The black thin solid curves are fiducial cuts based on DC inefficiencies and the red thick solid curves are cuts for bad TOF counters.

5 Monte Carlo simulation

In order to calculate the CLAS acceptance for $ep \rightarrow e'\pi^+n$, we simulated electron and pion tracks using the CLAS GEANT3-based Monte Carlo Package GSIM. For systematic checks, we used two Monte Carlo event generators. Our approach is that by comparing the results of simulations carried out with two very different event generators, a conservative and reliable estimation of systematic effects, such as finite bin size effects, is obtained. The first event generator, GENEV (see ref. [36] for the original publication dedicated to photoproduction processes), generates events for various exclusive meson electroproduction reactions for proton and neutron targets (π , ω , ρ^0 , and ϕ), including their decay, radiative effects, and

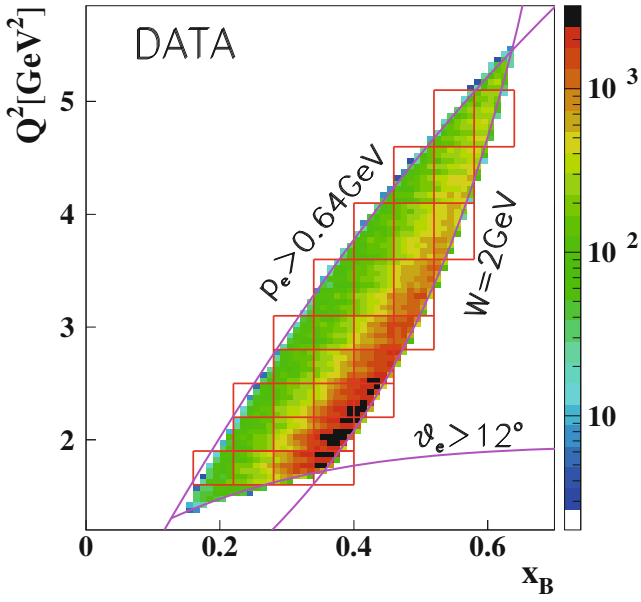


Fig. 8. (Color online) Kinematic coverage and binning (red boxes) as a function of x_B and Q^2 (integrated over all other variables) for the experimental data. Only events with $W > 2$ GeV are shown.

resonant and non-resonant multi-pion production, with realistic kinematic distributions. GENEV uses cross section tables based on existing photoproduction data and extrapolates to electroproduction by introducing a virtual photon flux factor (Γ) and the electromagnetic form factors. Radiative effects, based on the Mo and Tsai formula [37], are part of this event generator as an option. Although the formula is exact only for elastic e - p scattering, it can be used as a first approximation to simulate the radiative tail and to estimate bin migration effects in our pion production process, as will be discussed in sect. 5.2. The second event generator, FSGEN [38], distributes events according to the $ep \rightarrow e'\pi^+n$ phase space.

Electrons and positive pions were generated under the “e1-6” experimental conditions. Events were processed through GSIM. As already mentioned, additional *ad hoc* smearing factors for the tracking and timing resolutions are applied after GSIM so that they match the experimental data. The low-efficiency regions in the drift chambers and problematic TOF channels were removed during this procedure. Acceptance and radiative corrections were calculated for the same kinematic bins as were used for the yield extraction as shown in table 1. Figure 8 shows the binning in Q^2 and x_B applied in this analysis. However, some bins will be dropped at some later stage in the analysis, in particular due to very low acceptances (see following subsection). Our cross sections will be defined at the $(x_B, Q^2, -t)$ values given by the geometrical center of the three-dimensional bins. To account for non-linear variations of the cross section within a bin, a correction to our cross sections is determined by fitting with a simple *ad hoc* three-variable function the simultaneous $(x_B, Q^2, -t)$ dependence of our cross sections. This correction comes out at the level of a couple of percent in average.

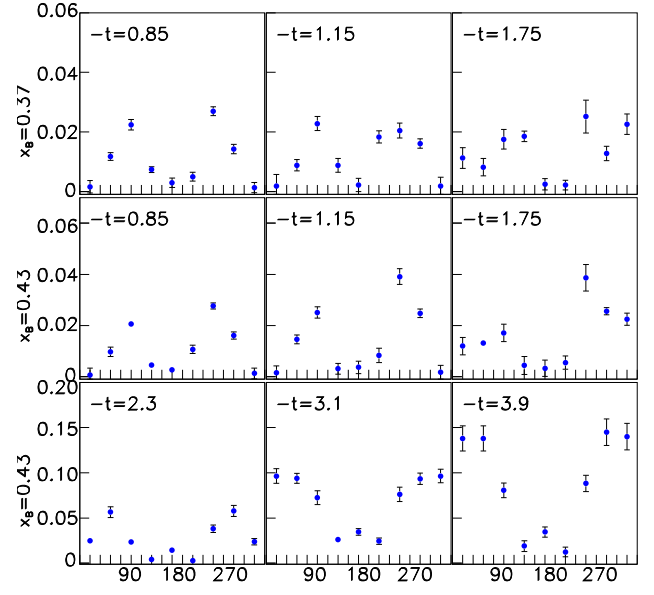


Fig. 9. (Color online) Examples of acceptance as a function of ϕ_π^* for various t and x_B bins at $Q^2 = 2.35$ GeV². The dips at $\phi_\pi^* = 0^\circ$ and 180° are due to sector boundaries in CLAS.

5.1 Acceptance corrections

We related the experimental yields to the cross sections using the acceptance, including the efficiency of the detector. The acceptance factor (Acc) compensates for various effects, such as the geometric coverage of the detector, hardware and software inefficiencies, and resolution from track reconstruction. We generated approximately 850 million events, taking radiative effects into account. This results in a statistical uncertainty for the acceptance determination of less than 5% for most bins, which is much lower than the systematic uncertainty that we have estimated (see sect. 7).

We define the acceptance as a function of the kinematic variables,

$$\text{Acc}(x_B, Q^2, -t, \phi_\pi^*) = \frac{N^{REC}(x_B, Q^2, -t, \phi_\pi^*)}{N_{rad.}^{GEN}(x_B, Q^2, -t, \phi_\pi^*)}, \quad (2)$$

where N^{REC} is the number of reconstructed particles and $N_{rad.}^{GEN}$ is the number of generated particles in each kinematic bin (the meaning of the subscript *rad.* will become clear in the next section). The kinematic variables in N^{GEN} refer to the generated values (after “radiation” but before propagation through CLAS) so that bin migration effects are taken into account in the definition of our acceptance. The acceptances are in general between 1 and 9%. Figure 9 shows examples of acceptances, determined with the GENEV+GSIM packages, as a function of the angle ϕ_π^* at a given Q^2 for various x_B and t bins. Bins with an acceptance below 0.2% were dropped. For the integration over the ϕ_π^* angle, in order to obtain our three-fold cross sections, we fitted the acceptance-corrected ϕ_π^* distributions, so that any hole in the ϕ_π^* distribution would be replaced by its fit value.

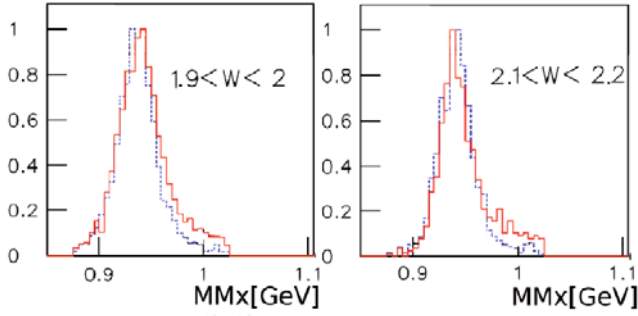


Fig. 10. (Color online) The simulated neutron missing mass distributions for two W bins with $\Delta W = 100$ MeV at $W = 1.95$ (left) and 2.15 (right) GeV integrated over ϕ_π^* , $\cos\theta_\pi^*$, and Q^2 . Normalized yields are shown with (solid red) and without (dashed blue) radiative effects.

5.2 Radiative correction

Our goal is to extract the so-called Born cross section (tree-level) of the $p(e, e'\pi^+)n$ which can thus be compared to models. However, we measure a process which is accompanied by higher-order radiative effects. Our experimental cross section must therefore be corrected. Radiative corrections are of two types: “virtual” corrections where there is no change in the final state of the $p(e, e'\pi^+)n$ process and “real” ones where there is in addition one (or several) Bremsstrahlung photon(s) in the final state. Such real Bremsstrahlung photons can originate either from the primary hard scattering at the level of the target proton (internal radiation) or from the interaction of the scattered or the initial electron with the various material layers of the CLAS detector that it crosses (external radiation).

We have dealt with these corrections in two steps. The effects of the radiation of hard photons (for instance, the loss of events due to the application of a cut on the neutron missing mass) are taken into account by the Monte Carlo acceptance calculation described in the previous section. Indeed, as mentioned earlier, the GENEV code has the option to generate radiative photons according to the Mo and Tsai formula and the N_{rad}^{GEN} events in eq. (2) were generated with this option turned on. Figure 10 shows examples of the simulated neutron missing mass with and without radiative effects in two W bins, obtained with the GENEV event generator and GSIM. Again, the Monte Carlo simulations were carried out with the same cut procedures and conditions as used in the data analysis.

Then, the correction due to soft photons and virtual corrections is determined by extracting the ratio between the number of events without radiative and with radiative effects at the level of GENEV for each three-dimensional kinematic bin. We therefore apply the following additional correction factor to our data:

$$F_{rad}(x_B, Q^2, -t) = \frac{N_{norad}^{GEN}(x_B, Q^2, -t)}{N_{rad}^{GEN}(x_B, Q^2, -t)}. \quad (3)$$

As a check, these radiative-correction factors were also calculated with the EXCLURAD code [39], which contains

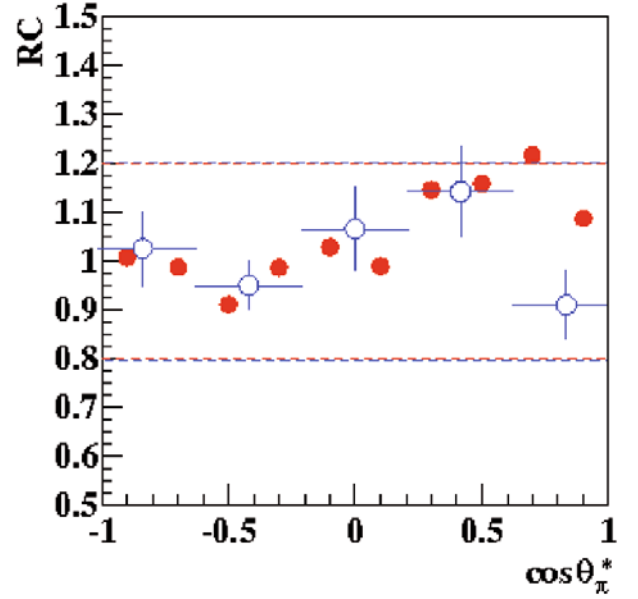


Fig. 11. (Color online) Radiative-correction factors (RC) as a function of $\cos\theta_\pi^*$ from EXCLURAD (red solid points) at $W = 1.74$ GeV, $Q^2 = 3$ GeV², and $\phi_\pi^* = 112.5^\circ$ and GENEV plus GSIM (blue open circle) at $W \approx 1.75$ GeV, $Q^2 \approx 3$ GeV², and $80^\circ < \phi_\pi^* < 120^\circ$.

a complete description of all internal radiative effects in exclusive processes, but is currently valid only up to $W = 2$ GeV. We compare the two different radiative-correction methods in a kinematic region where both methods are valid. Figure 11 shows the results for radiative-correction factors in the region $W \approx 1.75$ GeV and $Q^2 \approx 3$ GeV² as a function of $\cos\theta_\pi^*$.

The radiative correction factors from EXCLURAD are within $\pm 20\%$ of unity over the full $\cos\theta_\pi^*$ range (red solid points). The radiative corrections from GENEV+GSIM also fluctuate around 1.0 with a similar structure (blue open circles). The GENEV+GSIM error bars are due to Monte Carlo statistics (EXCLURAD is a theoretical code which has therefore no statistical uncertainty). The agreement between the two approaches is important because EXCLURAD is believed to be the most reliable of the two methods because it does not have the limitations of Mo and Tsai. Building on this reasonable agreement in this part of the phase space, we rely on the GENEV+GSIM radiative-correction factors for our data. In sect. 7, we discuss the systematic uncertainty associated with these radiative corrections.

6 Background subtraction

There are two main sources of background in our reaction. One consists of the mis-identification of pions with other positively charged particles (protons, kaons, positrons). This is particularly important for the pion-proton separation at high momenta ($p > 2$ GeV), see sect. 4.1. The other consists of multi-pion production. To subtract both

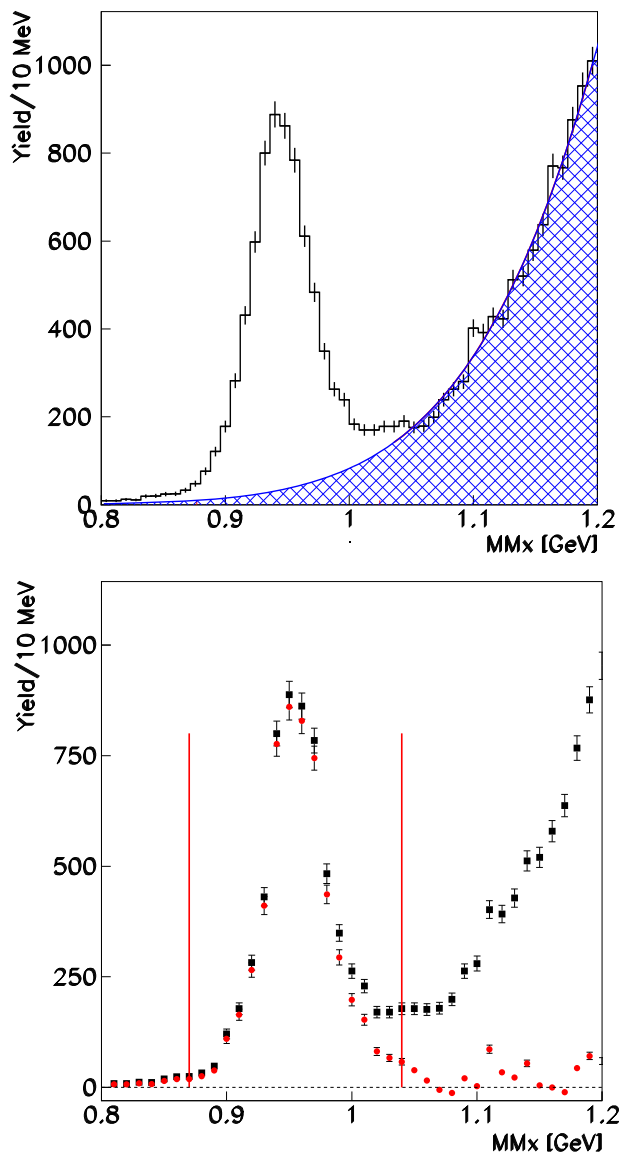


Fig. 12. (Color online) Example of the missing mass peak plus background at $Q^2 = 2.65 \text{ GeV}^2$, $-t = 1.15 \text{ GeV}^2$, and $x_B = 0.43$. The top plot shows the fitted background distribution (hatched region). The bottom plot shows the neutron missing mass distribution before (black squares) and after (red solid points) background subtraction.

backgrounds, we fit the neutron missing mass distribution bin by bin. We used many methods to fit these spectra: fit of only the background, fit of the signal plus background, with different functional forms both for the signal and the background, variation of the fitted range, etc. . . , from which we extracted a systematic uncertainty (see sect. 7).

Figure 12 (top) shows an example of a fit based on only the background, with an exponential plus a Gaussian. The former function was determined from simulations of the multi-pion spectra in the neutron missing mass region $> 1.02 \text{ GeV}$. A comparison of the missing mass (MMx) spectrum is shown in the bottom plot of fig. 12 before (black squares) and after (red solid circles)

background subtraction. In the range of the neutron missing mass cut, shown by the two vertical lines at 0.877 GeV and 1.0245 GeV , the background is small, and the remaining radiative tail becomes visible after the background is subtracted.

7 Systematic uncertainties

Several sources of systematic uncertainty that can affect our measurements have been studied by changing various cuts and using different event generators.

We varied the criteria used for the particle identification to provide more and less stringent particle selection simultaneously for experimental and GSIM data and then reran the complete analysis. The cuts on EC energy deposition and CC amplitude for the electron, as well as cuts on the TOF timing for the pion, have been varied. The EC sampling fraction cut was varied from $\pm 3\sigma_{\text{EC}}$ to $\pm 2\sigma_{\text{EC}}$ which led to a 5% uncertainty for electron identification. Changing the TOF β cut from $\pm 2\sigma_{\text{TOF}}$ to $\pm 2.5\sigma_{\text{TOF}}$ for pion identification gives a 1.7% uncertainty. The various cuts for channel identification such as fiducial, missing mass, and vertex cuts produced 3%, 1%, and 1.6% systematic uncertainties, respectively.

Acceptance and radiative corrections are the biggest sources of systematic errors. The systematic uncertainty from the acceptance is evaluated by comparing our results using the GENEV and FSGEN event generators. In the limit of infinitely large statistics and infinitely small bin size, our acceptances should be model-independent (up to the bin-migration effects). But these conditions are not reached here and we find differences between 2 and 8%. The systematic uncertainty for radiative corrections is estimated similarly by comparing the radiative-correction factors (GENEV and EXCLURAD). We calculated the difference between the cross sections corrected for radiative effects using either GENEV-GEANT simulation or the W -expanded EXCLURAD (where EXCLURAD was linearly extrapolated to $W > 2 \text{ GeV}$). An average 8% systematic uncertainty was found. Acceptance and radiative corrections are actually correlated, but after a combined analysis we estimated an averaged range 4–12% total uncertainty for both of these effects together.

Concerning the background subtraction procedure under the neutron missing mass (see sect. 6), we used various fitting functions (Gaussian plus exponential, Gaussian plus polynomial, exponential plus polynomial, etc.) and various fitting ranges. These various fitting functions and ranges eventually produced small differences and we estimated a 3% systematic uncertainty associated with this procedure.

To take into account the model dependency of our bin-centering correction (see sect. 4.3), we also introduce an error equal to the correction factor itself which is, we recall, at the level of a couple of percent in average.

These latter systematic uncertainties were determined for each bin. Concerning overall scale uncertainties, the target length and density have a 1% systematic uncertainty and the integrated charge uncertainty is estimated

Table 2. Average systematic uncertainties for the differential cross sections.

Source	Criterion	Estimated contribution
Type	Point to point	
e^- PID	Sampling fraction cut in EC ($3\sigma_{SF} \rightarrow 2\sigma_{SF}$)	5%
e^- fiducial cut	Fiducial volume change (10% reduced)	2.5%
π^+ PID	β resolution change ($2\sigma_{TOF} \rightarrow 2.5\sigma_{TOF}$)	1.7%
π^+ fiducial cut	Width (10% reduced)	3.5%
Missing mass cut	Neutron missing mass resolution ($3\sigma_{MMx} \rightarrow 3.5\sigma_{MMx}$)	1%
Vertex cut	z -vertex width (5% reduced)	1.6%
Acceptance	GENEV <i>vs.</i> FSGEN	
Radiative corrections	GENEV <i>vs.</i> EXCLURAD	4–12%
Background subtraction	Various fit functions exponential, Gaussian and high-order polynomials	3%
Bin-centering effect	Toy model	2–4%
Type	Overall scale/normalization	
LH2 target	Density/length	1%
Luminosity	Integrated charge	2%
Total		9–14%

at 2%. The background from the target cell was subtracted based on the empty-target runs and amounted to $0.6 \pm 0.2\%$ of our $e'\pi^+X$ events. The total systematic uncertainties, averaged over all bins, is then approximately 12%. Table 2 summarizes the main systematic uncertainties in this analysis averaged over all the accessible kinematic bins seen in fig. 8.

Let us finally mention that this same data set “e1-6” was analyzed earlier for the measurement of the five-fold differential cross section in semi-inclusive π^+ electropro-

duction on protons (ref. [40]). In this work, as checks of the normalization procedure, the inclusive proton structure function F_2 and the elastic cross section in electron-proton coincidence were extracted (figs. 7 and 8 of ref. [40]). The agreement with the world’s data of these two quantities was found to be very good with a maximum systematic uncertainty estimated to 10%. Although the groups involved in that earlier work and the present ones are different, as the data set analyzed and many software tools used are identical, we believe that this gives further confidence in our normalization procedure.

8 Results and discussion

In this section, we present our results for the cross sections of the $p(e, e'\pi^+)n$ reaction in the invariant mass region $W > 2\text{ GeV}$. We have extracted the differential cross sections as a function of several variables (t , Q^2 , and W or x_B). The angle ϕ_π^* is always integrated over in the following. The extraction of the interference cross sections σ_{TT} and σ_{TL} is the subject of an ongoing analysis and will be presented in a future article. The error bars on all cross sections include both statistical and systematic uncertainties added in quadrature. All values of our cross sections and uncertainties can be found on the CLAS database web page: <http://clasweb.jlab.org/cgi-bin/clasdb/db.cgi>.

8.1 $d\sigma/dt$ as a function of t

Figure 13 shows the differential cross section $d\sigma/dt$ as a function of t for different (x_B, Q^2) bins. We define the reduced differential cross section,

$$\frac{d\sigma}{dt} = \frac{1}{I} \frac{d^3\sigma}{dQ^2 dx_B dt}, \quad (4)$$

where the virtual photon flux factor [41] has been factored out. We have included in fig. 13 the JLab Hall C data, which cover only the very small t domain. We note that there is generally reasonable agreement between the results of the two experiments. However, care must be taken in comparing the Hall C and Hall B measurements as the central (t , Q^2 , and W or x_B , ϵ) values do not exactly match each other. For instance, the most important discrepancy seems to appear in the bin $(x_B, Q^2) = (0.49, 3.35)$ where the Hall C measurement was carried out at $\epsilon = 0.45$ [8] while ours corresponds to $\epsilon = 0.58$ (the x_B and Q^2 values being almost similar). According to the value of σ_L relative to σ_T , the Hall C cross section should then be renormalized: between a factor of $1.58/1.45 \approx 9\%$ (if $\sigma_L \approx \sigma_T$ which the Hall C separated data [6, 7] indicate, although at a slightly different kinematics) and a factor $0.58/0.45 \approx 30\%$ (if σ_L dominates over σ_T which the Laget model predicts). For better visualization, which is also relevant for the comparison with the models, we also show fig. 14 which concentrates on the low $|t|$ range of fig. 13.

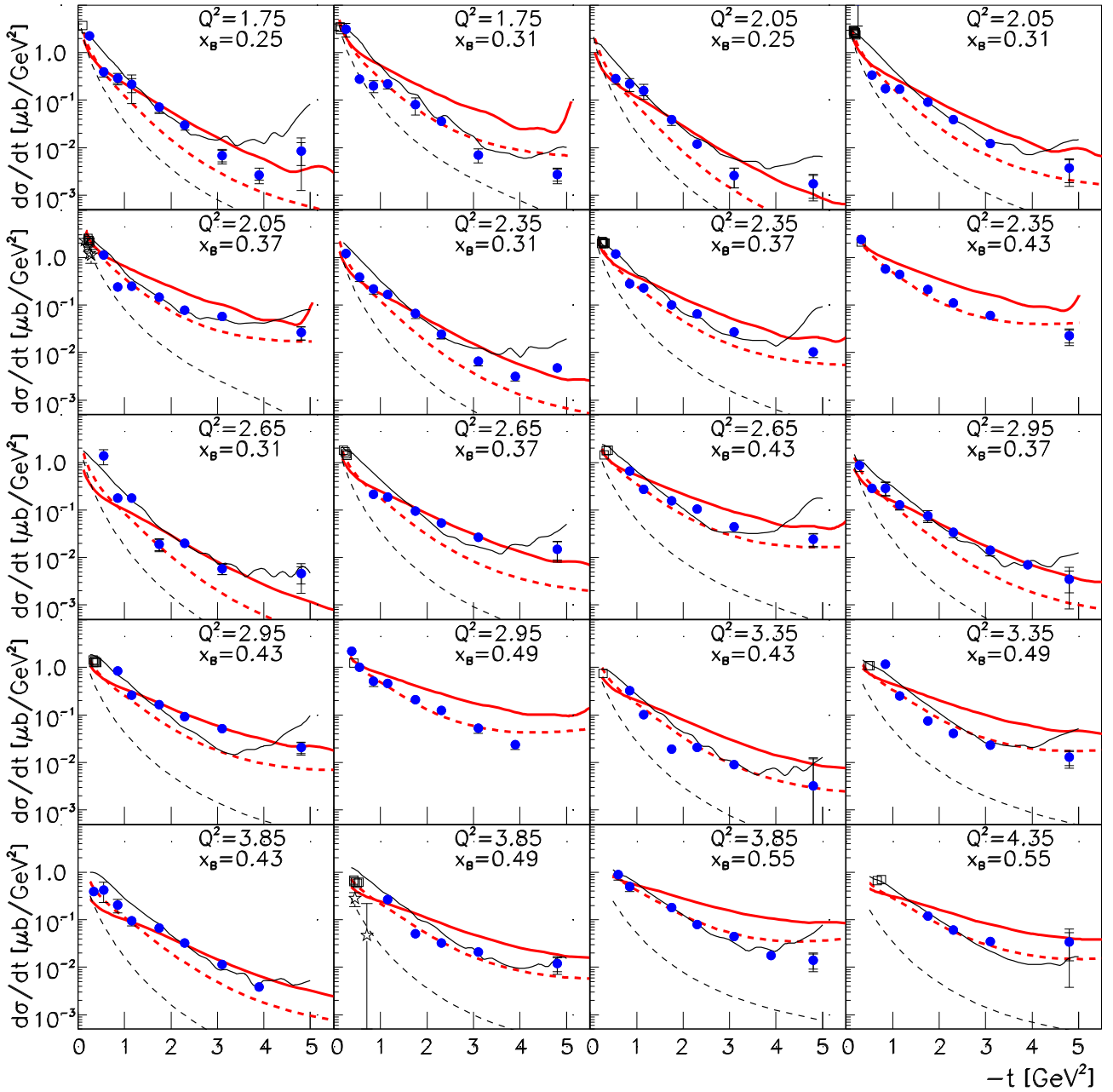


Fig. 13. (Color online) Differential cross sections $d\sigma/dt$ [$\mu\text{b}/\text{GeV}^2$] integrated over ϕ_π^* for various (Q^2, x_B) bins. The blue solid points are the present work. The error bars (outer error) on all cross sections include both statistical (inner error) and systematic uncertainties added in quadrature. The black open squares ($d\sigma/dt$) [8] and open stars ($d\sigma_L/dt$) [6] are JLab Hall C data. The red thick solid ($d\sigma/dt$), and dashed ($d\sigma_L/dt$) curves are the calculations from the Laget model [42] with (Q^2, t) -dependent form factors at the photon-meson vertex. The black thin solid ($d\sigma/dt$) and dashed ($d\sigma_L/dt$) curves are the calculations from the Kaskulov *et al.* model [45]. This model does not provide the calculation at $(x_B, Q^2) = (0.43, 2.35)$, $(0.49, 2.95)$ due to the kinematic limit in the model.

The $d\sigma/dt$ cross sections fall as $|t|$ increases, with some flattening at large $|t|$, which is a feature that is also observed in photoproduction [1, 2, 25, 26]. For several bins, for instance $(x_B, Q^2) = (0.31, 1.75)$ or $(0.37, 2.05)$, we notice a structure in $d\sigma/dt$ for $|t| \approx 0.5 \text{ GeV}^2$. The origin of this dip is not known. We note that the JLab Hall C experiment [8] also measured such a structure in $d\sigma/dt$ (see fig. 13 in ref. [8] for bin $(W, Q^2) = (1.8, 2.16)$).

We first compare our data to calculations using hadronic degrees of freedom. The first one with which we will compare our data is the Laget model [42] based on Reggeized π^+ and ρ^+ meson exchanges in the t -channel [43, 44]. The hadronic coupling constants entering the calculation are all well known or well constrained, and the main free parameters are the mass scales of the electromagnetic form factors at the photon-meson vertices.

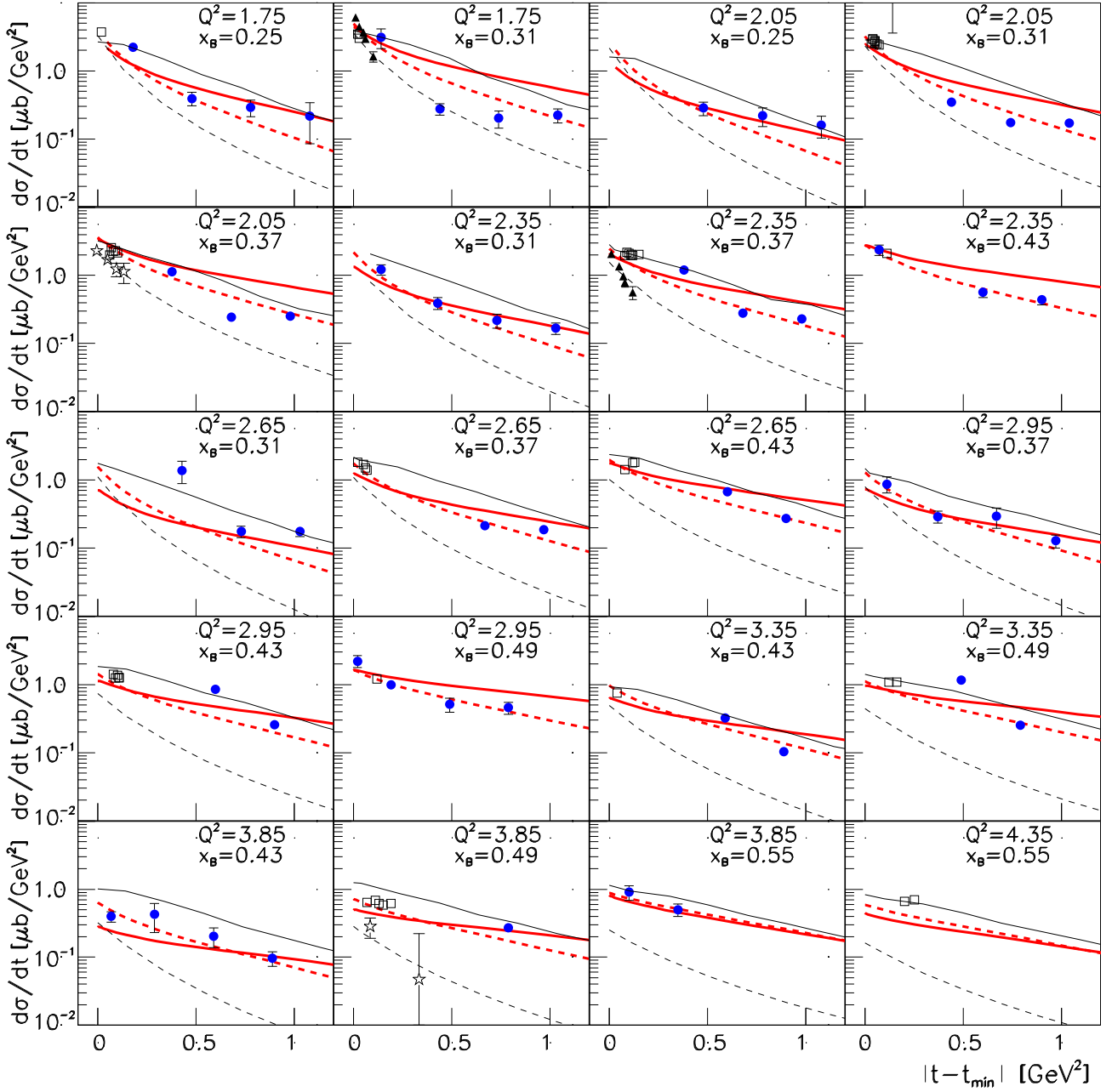


Fig. 14. (Color online) Same as fig. 13 except with an expanded low $|t - t_{min}|$ scale, where $-t_{min}$ is the minimum kinematically possible four-momentum transfer. In addition, the black solid triangles [7] show the JLab Hall C extracted $d\sigma_L/dt$ data.

If one considers only standard, monopole, Q^2 -dependent form factors, one obtains much steeper t -slopes than the data. An agreement with the data can be recovered by introducing a form factor mass scale that also depends on t according to the prescription of ref. [42]. This form factor accounts phenomenologically for the shrinking in size of the nucleon system as t increases. The size of the effect is quantitatively the same as in the $p(e, e'\omega)p$ channel (see fig. 1 of ref. [42]), which is dominated by pion exchange in the same energy domain as in our study. The results of this calculation with (Q^2, t) -dependent meson electromagnetic form factors are shown, for $d\sigma_L/dt$ and $d\sigma/dt = d\sigma_T/dt + \epsilon d\sigma_L/dt$, in figs. 13 and 14 by the red curves. The Laget model gives a qualitative description of

the data, with respect to the overall normalization at low t and the x_B, Q^2 and t dependencies. We recall that this model already gives a good description of the photoproduction data (SLAC, JLab) and of the HERMES electroproduction data, and that the form factor mass scale [42] has not been adjusted to fit our data.

In the framework of this model, $d\sigma_L/dt$ dominates at low $|t|$, while $d\sigma_T/dt$ takes over around $|t| \approx 2 \text{ GeV}^2$, this value slightly varying from one (Q^2, x_B) bin to another. This dominance of σ_L at low $|t|$ is a consequence of the t -channel π^+ -exchange (pion pole). At larger $|t|$, the ρ^+ meson exchange, which contributes mostly to the transverse part of the cross section, begins to dominate. The Laget Regge model, in addition to t -channel meson exchanges,

also contains u -channel baryon exchanges. It thus exhibits an increase of the cross section in some (Q^2, x_B) bins at the largest $|t|$ values, corresponding to low- $|u|$ values. We have additional data at larger $|t|$ (lower $|u|$) that are currently under analysis.

The second model with which we compare our data is the “hybrid” two-component hadron-parton model proposed in refs. [45, 46]. Like in the Laget model, it is based on the exchange of the π^+ and ρ^+ Regge trajectories in the t -channel. However, the model complements these hadron-like interaction types, which dominate in photoproduction and low Q^2 electroproduction, by a direct interaction of virtual photons with partons at high values of Q^2 followed by string (quark) fragmentation into π^+n . The partonic part of the production mechanism is described by a “deep inelastic”-like electroproduction mechanism where the quark knockout reaction $\gamma^*q \rightarrow q$ is followed by the fragmentation process of the Lund type. The transverse response is then treated as the exclusive limit of the semi-inclusive reaction $p(e, e'\pi^+)X$. Figures 13 and 14 show the results of this model compared to our data where very good agreement is found. This calculation was also found to give a good description of the L/T-separated Hall C and unseparated HERMES data [45, 46].

The third model that we wish to discuss, the GK model, is based purely on partonic degrees of freedom and is based on the handbag GPD formalism. In this model $d\sigma_L/dt$ is also mostly generated by the pion pole, similar to the two previous models. There are, however, a couple of differences in the treatment of this pion pole in the GK calculation. For instance, the Laget model has an intrinsic energy dependence. It is “Reggeized”, so the t -channel propagator is proportional to $s^{\alpha_\pi(t)}$, where $\alpha_\pi(t)$ is the pion Regge trajectory. In addition, it uses a (Q^2, t) -dependent electromagnetic form factor. These two features change the s , x_B , and t dependencies of the pion pole with respect to the GK treatment. Indeed, in the latter case, the t -channel pion propagator is proportional to $1/(t - m_\pi^2)$, so it has no s dependence, and the hadronic form factor at the πNN vertex is only t -dependent.

Figure 15 shows the results of the GK calculation (in blue) for $d\sigma_L/dt$ and $d\sigma/dt$. We recall that the GK model is applicable only for small values of $-t/Q^2$. Outside this regime, higher-twist contributions that are not taken into account in the GK handbag formalism are expected. The GK model describes qualitatively our low- t unseparated cross sections over our whole (x_B, Q^2) domain. This is remarkable since the GK model was optimized for higher-energy kinematics (HERMES) and no further adjustments were made for the present CLAS kinematics. We see that $d\sigma_L/dt$ has a non-negligible contribution only in the very low $|t|$ domain and only for a few (x_B, Q^2) bins, in particular at the lowest x_B and the largest Q^2 values (for instance, the (x_B, Q^2) bins (0.25, 1.75) and (0.31, 2.35)). This is in line with the observation that at HERMES kinematics, *i.e.* at lower x_B and larger Q^2 values, the longitudinal part of the cross section dominates in the GK model at low $|t|$. For the larger x_B values, one sees that the dominance of $d\sigma_L/dt$ at low $|t|$ is not at all systematic in the GK calculation. The ratio of $d\sigma_L/dt$ to $d\sigma/dt$ strongly de-

pends on x_B . Specifically, it decreases as x_B increases and at $x_B = 0.49$, $d\sigma_L/dt$ is only a few percent of $d\sigma/dt$, even at the lowest t values. This is a notable difference from the Laget Regge model, for instance.

In particular, one can remark in fig. 15, where we display in two (x_B, Q^2) bins ((0.31, 1.75) and (0.37, 2.35)) the longitudinal part of the cross section as extracted from Hall C [7], that the longitudinal part of the GK calculation is not in good agreement with the experimental data. This can be attributed to the way the pion pole and/or the pion-nucleon form factor, which are the main contributors to the longitudinal part of the cross section, are modeled in the GK approach. A Reggeization (like in the Laget model) or a change in the pion-nucleon form factor parametrization could possibly enhance the pion pole contribution at JLab kinematics and provide better agreement with our data (without damaging the agreement with the HERMES data) [47]. We recall that the GK model for which the GPD parameters were fitted to the low x_B HERMES data, was simply extrapolated to the kinematics of the present article without any optimization and thus the present disagreement observed in $d\sigma_L/dt$ should not be considered as definitive.

In the GK model, the transverse part of the cross section is due to transversity GPDs. In fig. 15, the GK calculation predicts that the transverse part of the cross section dominates essentially everywhere in our kinematic domain. Although the GK L/T ratio probably needs to be adjusted as we just discussed, the GK calculation opens the original and exciting perspective to access transversity GPDs through exclusive π^+ electroproduction.

Finally, at the kinematics of our experiment, in spite of our $W > 2$ GeV cut, it cannot be excluded that nucleon resonances contribute. In ref. [48], Kaskulov and Mosel identify these high-lying resonances with partonic excitations in the spirit of the resonance-parton duality hypothesis and invoke the continuity in going from an inclusive final deep inelastic state to exclusive pion production. During this transition one expects that the inclusion of resonance excitations enhances the transverse response while leaving the longitudinal strength originating in the t -channel meson exchanges intact. Thus, in this work, the t -channel exchange part of the production amplitude is again described by the exchange of the Regge trajectories (π^+ , ρ^+ and a_1^+) to which it is added a nucleon resonance component that is described via a dual connection between the resonance and partonic deep inelastic processes. The parameters of this model have been tuned using the forward JLab Hall C data. Figure 15 shows the results of this calculation with our data and a reasonable agreement is found.

The four models that we just discussed, although they give a reasonable description of the unseparated cross sections, display rather different L/T ratios. The precise measurement of this ratio as a function of x_B , Q^2 and t appears thus as essential to clarify the situation. For instance, in order to validate and/or tune the GK approach, it would be interesting to study the Q^2 dependence at fixed x_B and

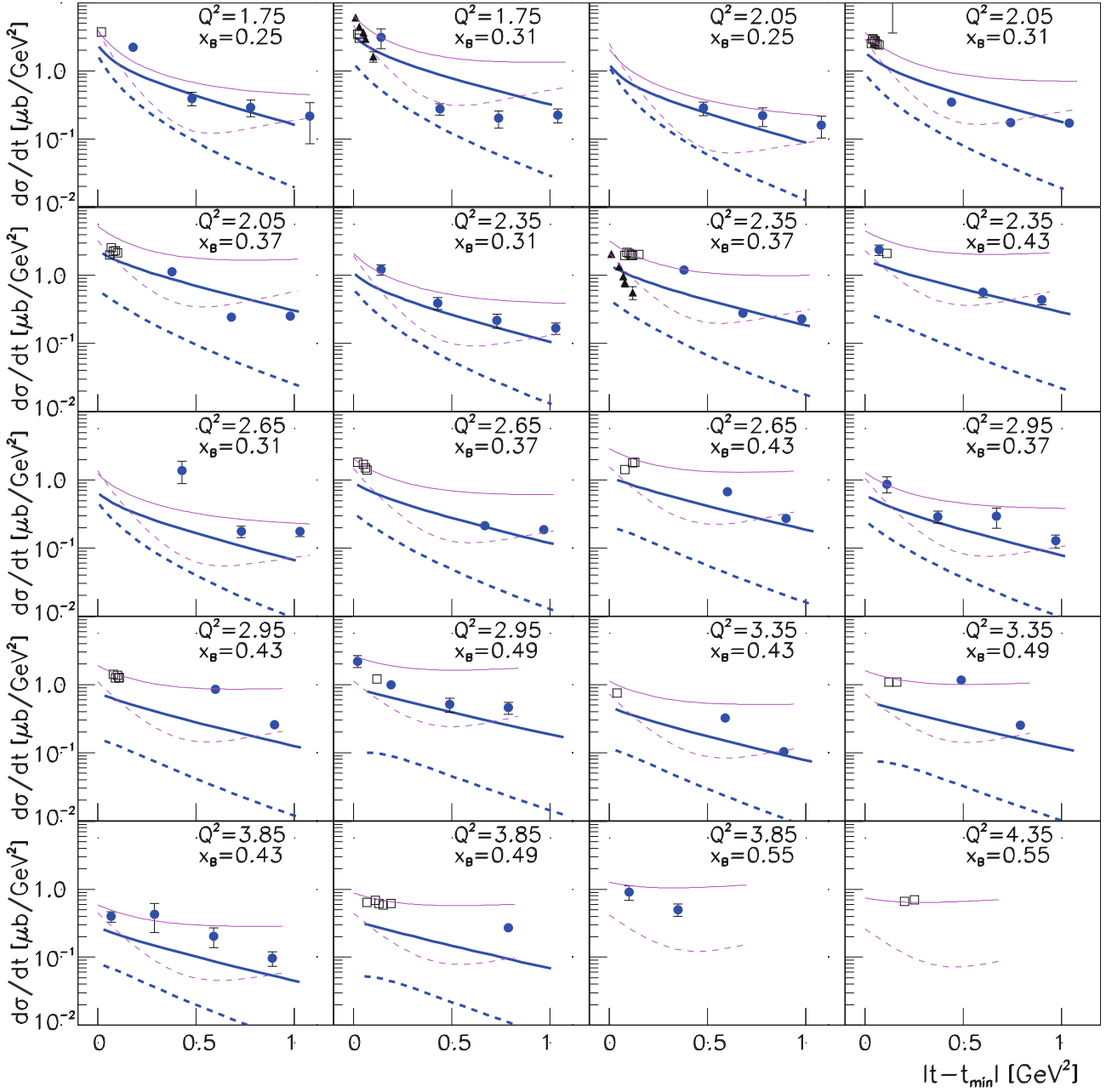


Fig. 15. (Color online) Same as fig. 14 except with $|t - t_{min}|$ scale. The blue thick solid and dashed curves are $d\sigma/dt$ and $d\sigma_L/dt$, respectively, from the GK model [22]. The magenta thin solid ($d\sigma/dt$) and dashed ($d\sigma_L/dt$) curves are the calculations from the 2nd Kaskulov *et al.* model [48].

t of the longitudinal and transverse cross sections. They should approach, as Q^2 increases, a $1/Q^6$ and $1/Q^8$ scaling behavior, respectively, as mentioned in the introduction of this article. In contrast, the Laget Regge model, for which x_B is not a “natural” variable (it is rather W) should not predict such a Q^2 scaling at fixed x_B . Although we are probably very far from such an asymptotic regime, the measurement of the Q^2 dependence in the transition region accessible with the upcoming JLab 12 GeV upgrade should provide some strong constraints and in particular some checks on the way the higher-twist corrections are treated in the GK model. Such a program is already planned at JLab [49].

8.2 $d\sigma/dt$ as a function of Q^2 at fixed t

Figures 16 and 17 show the differential cross section $d\sigma/dt$ as a function of Q^2 at fixed x_B for various t values. In fig. 16, our data are fitted with a $1/Q^n$ function and are compared to the GK model. We recall that, at asymptotically large Q^2 , the handbag mechanism predicts a dominance of σ_L which should scale as $1/Q^6$ at fixed t and x_B . The resulting exponents n of our fit indicates a flatter Q^2 dependence than $1/Q^6$. At the relatively low Q^2 range accessed in this experiment, higher-twist effects are expected to contribute and hence the leading-twist $1/Q^6$ dependence of σ_L is no longer expected. We note that such

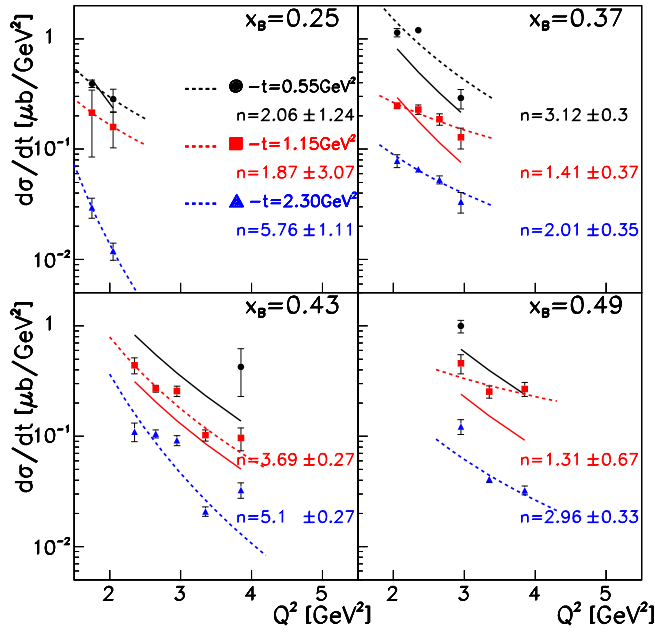


Fig. 16. (Color online) Differential cross sections $d\sigma/dt$ [$\mu\text{b}/\text{GeV}^2$] versus Q^2 at fixed x_B for various t values. The dashed curves are the results of a fit to the function A/Q^n . The solid curves are the results of the GK calculations [22]. The GK calculations are only valid for $-t < \approx 1 \text{ GeV}^2$ so we do not display those results for $-t = 2.3 \text{ GeV}^2$.

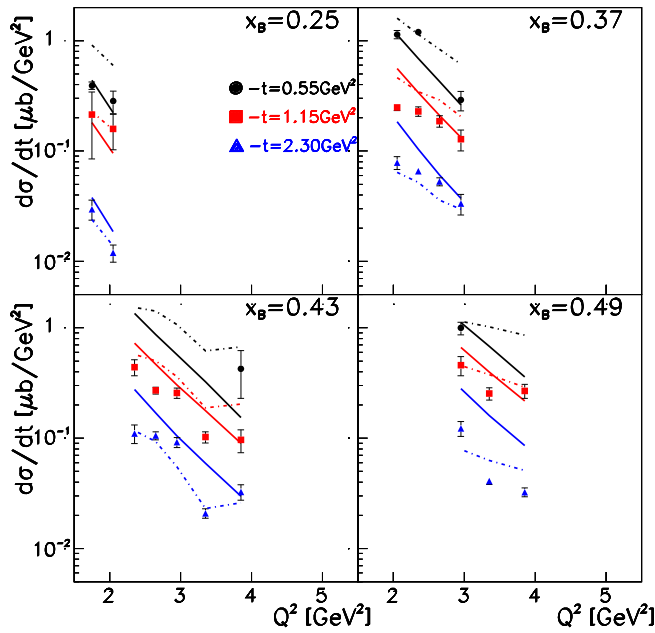


Fig. 17. (Color online) Differential cross sections $d\sigma/dt$ [$\mu\text{b}/\text{GeV}^2$] versus Q^2 at fixed x_B for various t values. The solid curves are the results of the Laget calculations [42] and the dash-dotted curves of the “hybrid” two-component hadron-parton model of refs. [45,46].

higher-twist contributions are part of the GK calculation and the GK model also does not show this scaling behavior at the present Q^2 values. Although the GK model tends to underestimate the normalization of our data, its Q^2 dependence agrees reasonably well with our data.

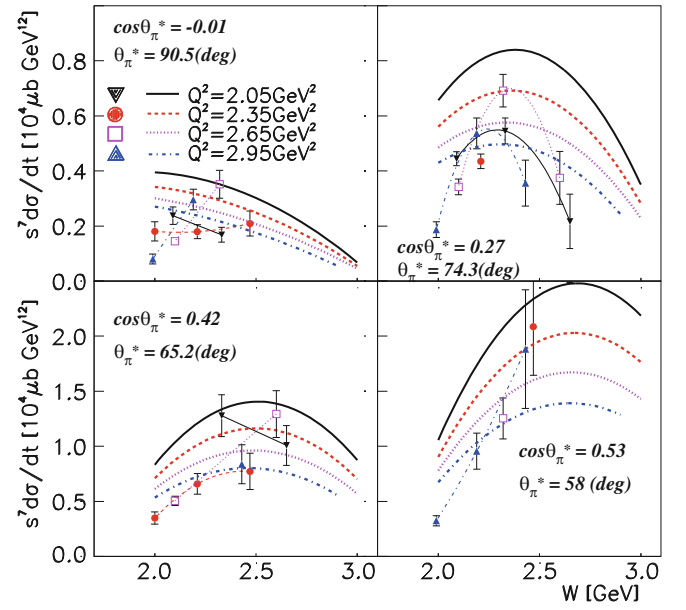


Fig. 18. (Color online) Scaled cross sections $s^7 d\sigma/dt$ [$10^4 \mu\text{b GeV}^{12}$] versus W for $\theta_\pi^* \approx 60^\circ \dots 90^\circ$ and various Q^2 bins. Thick curves are from the Laget model [42]. The thin curves are to guide the eye, connecting points with the same Q^2 values.

In fig. 17, we compare our data to the Laget [42] and the Kaskulov *et al.* [45,46] models. The Laget calculation gives a reasonable description of the data although it seems to have a slightly steeper Q^2 dependence than our data (particularly in the $x_B=0.37$ bin). We note that in the $x_B=0.43$ bin, our data seem to display a structure (dip) for Q^2 values between 3 and 4 GeV^2 , which is certainly intriguing. We have at this stage no particular explanation for this. We just observe that the “hybrid” two-component hadron-parton model of refs. [45,46] displays apparently also such structure which should therefore be further investigated.

8.3 $d\sigma/dt$ as a function of W at fixed θ_π^*

Figure 18 shows our scaled cross sections, $s^7 d\sigma/dt$, as a function of W for four Q^2 values and four bins in $\cos\theta_\pi^*$: -0.01 ± 0.16 , 0.27 ± 0.1 , 0.42 ± 0.05 and 0.53 ± 0.06 . The lever arm in W is limited. At $\theta_\pi^* = 90^\circ$, where the scaling behavior is expected to set in most quickly, we have only 2 or 3 data points in W , depending on the Q^2 bin. It is therefore difficult to draw precise conclusions at this stage for the W dependence at fixed Q^2 . Nevertheless, with these limited (but unique) data, one can say that, at $\theta_\pi^* = 90^\circ$, except for the 3 data points at $Q^2 = 2.35 \text{ GeV}^2$, the W dependence of $s^7 d\sigma/dt$ does not appear to be constant. We also display in fig. 18 the result of the Laget model. It gives, within a factor two, a general description of these large-angle data. The W dependence of our data is very similar to the energy dependence that was observed in photoproduction [3]. In the same energy range as covered by the present study, real-photon data exhibit strong deviations from scaling. Within the Laget model, these

deviations are accounted for by the coupling between the $n\pi^+$ and the ρN channels [50]. The JLab 12 GeV upgrade will allow us to increase the coverage in W and check whether the hints of oscillations that we observe remain in the virtual-photon sector.

9 Summary

We have measured the cross sections of exclusive electroproduction of π^+ mesons from protons as a function of $-t = 0.1\text{--}5.3\text{ GeV}^2$, $x_B = 0.16\text{--}0.58$, and $Q^2 = 1.6\text{--}4.5\text{ GeV}^2$. We have compared our differential cross sections to four recent calculations based on hadronic and partonic degrees of freedom. The four models give a qualitative description of the overall strength and of the t , Q^2 and x_B dependencies of our unseparated cross sections. There is an obvious need for L-T separated cross sections in order to distinguish between the several approaches. These separations will be possible with the upcoming JLab 12 GeV upgrade. In particular, if the handbag approach can accommodate the data, the $p(e, e'\pi^+)n$ process offers the outstanding potential to access transversity GPDs.

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