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PHOTON EXCITATION FOR SATELLITE FREE X-RAY SPECTROSCOPY: **INSTRUMENTATION CHALLENGES***

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1. I**nt**ro**duct**io**n**,

In an attempt to produce radiation which would travel through matter opaque to ordinary visible light, Roentgen¹ in 1895 discovered x-rays. The potential of applications of x-rays in Science, Technology and Medicine were rea**l**ized immediately.² Within a year, x-rays were widely applied in medical and industrial radiography. In 1910-1920*,* M. Sieghbahn 3 and co-workers developed the techniques for measuring wavelengt**h**s of x-ray absorption and emission spectra of the chemica**l** elements. The effect of chemical combination on x-ray emission and absorption spectra was noted and well realized that the effect was small except for the lighter elements since valence electrons are more involved in chemical bonding than core electrons. Hence K and L x-ray emission and absorption spectra of light elements $(Z = 3$ to 18), that occur in the low energy x-ray region were investigated and soft xray spectroscopy was developed.

Lindh⁴, in summarizing earlier studies of S-K β emission spectra, showed that there is a great inconsistency between these measu*r*ements*,* mainly because measurements were obtained using direct electron bombardment, suggested using fluo*r*escent x-ray excitation for x-ray emission spectral studies. The advantages of photon excitation over electron excitation were realized and the occurrence of satellites (or extra lines) in x-ray emission spectra was widely known. The extra features that appear in x-ray emission spectra in addition to the single vacancy (parent lines) transitions were known as satellites. Wentzel⁵ proposed the theory of satellites in x-ray spectra based upon a multivacancy processes, in which two or more electrons are excited or ionized, result in extra features. Later Langer⁶ and Wolfe⁷ studied the satellites in K α emission spectra to demonstrate the importance of the exchange interactions in determining the energy levels of the doubly ionized systems.

First systematic study of satellites in x-ray emission spectra was perfo*r*med by Deslattes⁸ using quasi-monochromatic photon excitation from a group of L α x-ray sources lying close to the K edge of C**I**. He observed significant alterations in the CI $K\beta$ spectrum of KCl depending on the character of the excitation radiation and identified the initial state of these satellites as a double vacancy state. Recently, the valence electronic structure of the chlorofluoromethanes were analyzed by chlorine K x-ray emission under satellite-free conditions.⁹ These studies were based on the use of synchrotron radiation to eliminate the multivacancy effects that are inherent in conventional x-ray spectroscopy.¹⁰ In this report, satellite free x-ray emission spectra from chlorofluoromethanes will be presented to demonstrate that the simplified spectra can be obtained using selective photon excitation. Results from va*r*ious research groups world wide, utilizing the tunable photon excitation form synchrotron sources to eliminate the obscuring features in x-ray emission spectra of rare-gas solids (RGS) and metals will be discussed. Also, the technical challenges in utilizing the small phase-space attributes of high brightness from third generation SR sources producing soft x-ray and vacuum ultra-violet wavelengths to study weak features like satellites in x-ray emission spectra will be presented*.*

ייצט א∰טרי בארץ אין איש און און האוט מועד מועד און השמה המערי הודע אש המכות האוט מערכה מה מאמה משכה מכונה המכות המכונה בין היום היום בין המכונה בין המכונה היום בין המכונה היום בין המכו

2. Results an**d D**i**scus**s**ion**

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2.1 Molecular Gasses

j _ C1 K absorpti*o*n and C1 K-V emission spectra of several chlorofluoromethanes were measured using monochromatic radiation f*r*om X-24A beamline at National Synch*r*otron Light Source (NSLS) as the s**o**urce of excitation. Detailed description of the soft x-ray beamline¹¹, monochromator¹² and experimental procedures⁹ are given els**e**where. As an example, results from CF3CI (gas phase) will be presented. In order to estimate contributions from satellite features in the x-ray e**m**ission spectra, C1 K-V fluorescence spectra were recorded at several incident x-ray energies in the vicinity of the single **-** and multiple- vacancy C1 K ionization threshold energies. Using the present C1 K-V emission data and UPS final state binding **e**nergies 13, the C1 ls ionization thr**e**shold of CF3C1 is est**i**mated to be 2830.2 eV. Hence, the 2833 eV photon excitation is an appropriate photon, energy for producing satellite free **s**p**ec**trum**, s**in**c**e the ex**c**it**a**tion energy is below all double vacancy t**h**reshold energies. The CI K-V emission spectrum obtained using 2833 eV photon excit**a**tion is presented in Fig. 1 along with the C1 K-V emission spectrum obtained using 2880 eV photon excitation. The 2880 eV photon energy is above most [KV] double vacancy threshold energies, which correspond to production of initial states required for xray emiss**i**on satellites in this energy range. As seen in Fig. 1, the 2880 eV spectrum (dots) **c**ontains extra contributions resulting from multi-vacancy transition*s*. The main satellite features thus obtained are located on the high energy si**d**e of the single-vacancy spectrum.

To facilitate interpretation of the experimental fe**a**tures in the pres**e**nt x**-**ray emission spectra, the spectra were modeled empirically¹⁰ using a SIMPLEX minimization technique. For the modeling*,* the instrumental broa**d**ening was assumed to be Gaussian and the natural line width to be L**o**rentzian, resulting in Voigt spectral profiles. Th**e** Gaussian instrumental broadening was estimat**e**d to be 1.0 eV by analyzing a series of x-ray emission spectra in t**h**is *r*egion. This value is consistent with direct measurements of the secondary**-**spectrometer resolution. The 2833 eV C1 K-V spectrum is presented again in Fig. 2 along with the mode**l**ed Voigt spectral components and their sum. Modeling of the x-ray emission spectra permitted extraction of energies and relative intensities of the emission peaks.
Assignments were then determined by comparison with known valence molecular orbital (MO) orderings from UPS measurements. 13

I ClK-V spectrum in Fig. 2 have a common initial state; namely, a single Cl 1s Because of sele**c**tive excitation of CF3CI at 2833 eV, all emission features in the $\frac{1}{2}$, and the ze and tual valence MO's, respectively, filling the initial CI is vacancy. The electrons fr**o**m valence MO's with C1 3p character. **T**he stronger e**m**ission **c**omponents, peaks **C** and B, are assigned to dipole-allowed transitions of elect*r*ons weaker emission peaks B' and A are assigned to allowed dipole transitions form the 5e and the $9a_1 + 4e$ valence MO's, respectively.

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2.2 Rare**-Gas Solids**

RGS **o**ften used in th**e** de**v**e**lo**pment of solid**-s**tate theory are bonded only by the weak Van der Waals forces in a FCC latti**c**e. Therefore, uncomplicated satellite free x-ray emission spe**c**tra of RGS, which provide information about the partial density of states of the occupied valence bands will enable to test the rigorous band structure calculations¹⁴ performed on various solids. Perera and Henke¹⁵ measured the photon excited L2*,*3 (3s-2p) emission spectra f**r**om gas and solid argon. Recently, high resolution soft x-ray emission measur*e*ments 16 were performed by Jia *et al*. 17 on solid Kr $M_{4,5}$ (4p-3d) and Xe $N_{4,5}$ (5p-4d) using monochromatic radiation from U-10A beamline at NSLS as the source of excitation. As **a**n example of their work, the solid Kr M_{4.5} (4p-3d) x-ray emission spectra excited with various photon energies are presented in Fig. 3. They observed that the previously reported 18 satellite feature at about 82 eV is prominent when white light excitation is used and the satellite intensity decreases as the excit**a**tion energy is decr**e**ased. When a photon energy corresponding to a excitation energy below the multiple ionization threshold is used (103 eV photon excitation), the measured intensity in that energy region was extremely small. Furthermore, the previously observed 18 satellite feature at 79.99 eV resulting from *m*ultiple ionization in krypton *via* Coster-Kronig decay of the L1 hole state is energetically not possible. Therefore, the emission intensities observed in there 103 eV photon excited spectrum is a true measure of the p-like partial density of states of the valence band in solid Kr, a result that will stimulate the theoretical efforts on RGS. This is first time satellite free x-ray emission spectra from RGS were obtained using monochromatic synchrotron radiation to suppress the complicating satellites.

2.3 Metals

A seri**e**s of experiments in x-ray emission spectra of metals were carried out by Nordgren¹⁹ and his co-workers²⁰ at the Hamburg synchrotron radiation laboratory (HASYLAB) at DESY, using radiation from the FLIPPER multipole wiggler beamline. The measured photon flux at the sample position of this beamline is about 10**1**2 ph*/*s with energy resolutions between 0.5 eV and 10 eV in the 20 eV to 1 keV photon region. They measured²⁰ the Cu L_{2.3} (3s-2p) spectra of copper metal excited using different photon energies as shown in Fig. 4. The 970 eV photon energy is sufficient to excite L_2 (2p_{1/2}) and L_3 (2p_{3/2}) hole states and the corresponding spectrum shows two bands *r*esulting from the transitions from the 3s band to 1,11the 2pl*/*2 and 2p3*/*2 vacancies respectively. The 935 eV photon excitation is below the L₂ (2p_{1/2}) threshold energy but nearly above the L₃ (2p_{3/2}) threshold energy, therefore, should energetically r**e**move the emission peak at about 950 eV corresponding to the transitions to the 2p_{1/2} hole state. The reduction in intensity of the emission peak at about 930 eV when 935 eV photon excitation is used, suggested the occurrence of satellite features²⁰ resulting from Coster-Kronig decay of the L₂ (2pl*/*2) hole state, giving rise to L3 (2p3*/*2) multi-vacancy states. '**]**'he difference spectrum shown in Fig. 4 clearly reflects the contributions from such satellite features yielding information about the shake-off probabilities.

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Th**e** a**d**vantages of **se**le**c**tiv**e e**x**c**it**a**ti**o**n for x**-**ray emission spectroscopy was also de*m*onstrated 1**9** when separate N K (392 eV) and Ti L3 (394 eV) x**-**ray emission spectra of TiN was obtained²¹ even though they overlap in this energy range. They measured the N K emission spect*r*um of TiN using 410 eV monochromatized **,** synchrotron radiation as the source of excitation (see Fig. 5)*.* Since, the corresponding excitation energy is below the Ti L_3 threshold energy (453 eV), the spectrum presented in Fig. 5 is free of contributions from overlapping Ti L₃ x-ray emission. Furthe**r**more, use of photon excitation closer to the N K threshold ene*r*gy (409 eV) will eliminate the satellite features, resulting a true N K x**-**ray emission spectr**u**m of TiN. This will enable the band spectrum mapping of the N 2p contributions to the valence **b**and of TiN in an unambiguous way. Also shown in Fig. 5 is the Ti L_3 emission spectrum obtained using 458 eV photon excitation, an excitation energy lie between L₂ and L₃ threshold energies. The spectra presented in Fig. 5 are aligned to the Fermi energy by means of XPS binding energy **d**ata. They also measured²¹ Ti L emission spectra using photon energies well above and well below the L2 threshold en**e**rgy of 453 eV, and the differ**e**nce spectrum presented in Fig. 5 thus represent the Ti L_2 (3s - 2p_{1/2}) emission spectrum of TiN. This is the first time undistorted N K, Ti L2 and Ti L3 em**i**ssion spectra of TiN was measured, a measurement that will facilitate the ban**d** structure investigations of this compound. They also obse*r*ved a new type of resonance in the Ti L emission spectra, resulting from the decay of a quasiatomic intermediate state²¹.

3. Im_t**rume**n**tatio**n **Challe**n**ges**

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The technical requirements to study weak features like satellites in x-ray emission spectra are: high flux delivered to a small spot size, stable beamline optical **sys**t**em,** tun**abili**t**y o**f t**he x-**r**ay sou**r**ce** (m**o**n**och**r**omat**i**zed**) **ove**r a lar**ge e**n**e**r**gy** ra**nge** and a high resolution x-ray spectrometer to analyze the emitted photons. Third **gene**r**a**t**ion sy**n**ch**r**o**tr**on** r**ad**i**a**ti**o**n **sou**r**ces** l**ike Ad**v**a**n**ced Ligh**t **So**ur**ce** (**ALS**) in B**e**r**keley designed 22 fo**r **dedica**t**ed expe**r**imen**t**s by** ph**o**t**o**n u**se**rs **23 is op**t**i**m**ized fo**r high spectral brightness in the soft x**-**ray and vacuu**m** ultra-violet wavelengths. The undulator radiation frorn ALS can deliver high photon flux to a small spot size provided the d**e**sign of the beamline do not diminis**h**, the spectral brightness through aberrations associated with the imaging properties or through the figure err**o**rs associated with the optical components. Because of the high photon fluxes and pow**e**r loading, it is quite difficult to design a stable beamline optical system consisting of collecting mir*r*ors, gratings, focussing elements, etc. that does not diminish the spectral brightness.

An insertion device beamline²⁴ at ALS is under construction utilizing **.** radiation**2**5**,** 22 from a 5.**0 c**m peri**o**d un**d**ulator (U5.0) to cover t**h**e 1**00-**]5**0**0 eV ph**o**t**o**n range. Spectral flux in the central curve (0.1% band pass) as a function of photon energy for the U5.*0* undulator is shown in Fig. 6. Each undulator c**u**rve is the locus , of narrow peaks of' radiation, tuned by altering the undulato*r* gaps. Separate curves are shown in Fig. 6 f**o**r the first, third, an**d** fifth h**a**rmonics of the undulator. A spherical grating monochromator (SGM), is similar to the one constructed 27 at ALS for another U5.0 undulator beamline with fixed grating rotationa**l** axis and fixed incoming and outgoing ray directions. Simple rotation of the grating p*r*ovides the

energy tuning. The Rowland-Circle condition can be *s*atisfied throughout the scanning range by moving the entrance*/*exit slits along the beam di*r*ection providing a good focus. The energy resolution of SGM is shown in Fig. 7 when the exit slit motion is about 0.7m. The monochromator utilizing three gratings, 200 1*/*mm, 500 1*/*mrn and 1200 1*/*mm to cover the first*,* third*,* and fifth harmonics of the undulator respectively, results in a energy resolution $E/\Delta E$ better than 10⁴ for the 100-200 eV region. The beamline output flux 26 at the exit slit for the U5.0 undulator beamline consist of a SGM and 10 μ m entrance and exit slits are presented in Fig. 8*.* The flux transmission through the slits and grating efficiencies are included in these calculations. As seen from Fig. 8, using refocussing optics, over 1012 ph*/*s in a 0.1% bandwidth can be delivered to a 10μ m spot size.

One of the five end stations under construction to utilize this radiation is a high resolution soft x-ray emission (SXE) spectrometer²⁸, a grazing incidence, Rowland Circle instrument where the diameter of the Rowland Circle can vary from 5 to 10m. This spectrometer is designed to match the small phase-space attributes of high spectral brightness of the beamline optical system. Special features of the SXE spect*r*ometer are four toroidal gratings on a carousel mount to cover the energy range of 40-1400 eV and a photon counting area detector in place of the output slit. The energy resolution $(E/\Delta E)$ of the SXE spectrometer in this energy range is about 1000, sufficient to resolve features in x-ray emission spectra.

4. Conclusions

Tunable radiation from third generation synchrotron radiation sources like ALS when combined with a SXE spectrometer designed to match the emittance properties of these sources will enable to perform more difficult experiments discussed in this paper in a routine manner. Also, will able to extend these measurements to more difficult systems like gas phase studies, buried interfaces etc*.* and to obtain spectra with a higher resolution. It was noted 15 several years ago that the satellite features in molecular gases and rare gases were less prominent in the sold phase. A comparison of satellite features in solid and gas phases will give a insight to the production of multi-vacancy states as well as the decay of the multivacancy hole states resulting from solid state effects. Also tunable radiation has the ability to "turn off" the multi-electron channels producing satellite features in x-ray emission spectra. Simplified spectra thus obtain can provide information about radiation yields, line widths and peak shapes in an unambiguous way and no other technique can attain this results.

The unique aspects of photon-in photon-out measurements are: threshold excitation provides spectral purity and sensitivity, photons are less damaging compared to other particles, provide bulk sensitivity and ultra high vacuum is not required. Photon-out phenomena can also be used in the presence of charged particles and strong electric and magnetic fields associated with the various deposition methods enabling to monitor on-line processing of thin films and layered structures.

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5. Acknowledgements

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Figure Captions:

- **Figure 1. Co**m**parison of 2880 eV photon excited Cl K-V emission spect**r**um of CF3***C***1** (**c**ir**cles**) **with a 2833** e**v excited spec**tru**m i**n**d**i**cated by the** s**olid curve. The inc**r**eased** i**n**te**nsity con**tri**b**u**tio**n**s** in **the 2880 eV excited s**pectru**m** are attributed to multi**-**va**c**ancy transitions.
- Fi**g**ur**e 2**. Exp**e**rim**e**ntal 2833 **e**V ph**o**ton **e**x**c**it**ed C**1 K-V **e**mi**s**sion **s**pe**c**tru**m o**f CF**3**C1 (dots) along with the fitted spe**c**tral components (dashed curves) and the sum of the derived components (solid curves).
- Figure 3. The solid krypto**n** M**4***,*5 (4p-3d) x-ray emission spectra obtained using monchromatized synchrotron radiation. The corresponding excitation energies are indicated above each spectrum. (J.J. Jia *et al*. Ref.17).
- Figure 4. The Cu L_{2.3} (3s-2p) emission spectra of copper metal obtained using 935 and 970 eV photon excitations and the corresponding difference spectrum. (J Nor**d**gren *et al*. Refs. 19 and 20).
- Figure 5. Nitrogen K and titanium L x-ray emission spectra of TiN. (J Nordgren *et al*. Refs**.** 19 and 20).
- Figure 6. Spectral flux in the central curve (0.1% band pass) as a function of the photon energy for the U5.0 undulator at ALS for the first*,* third, and fifth h**a**rmonics of the undulator. Each curve is the locus of narrow peaks of radiation, tuned by altering the undulator gaps. (P. Heimann Ref. 26)
- Figure 7. The energy resolution of the monochromator utilizing three gratings, 200 1*/*mm, 500 1*/*mm and 1200 1*/*mm to cover the first, thi*r*d*,* and fifth harmonics of the undulator respectively. The exit slit motion is about 0.7m. (A.Warwick *et al*. Ref.27).
- Figure 8. The beamline output flux at the exit slit for the U5.0 undulator beamline consist of a SGM and 10 μ m entrance and exit slits. The flux tr**a**nsmission through the slits and grating efficiencies are included in these calcu**l**ations.

6. References

- 1**.** W.C. R6nt**g**en, Sitzungsber **d**er W_irzburger Physik**-**Medic. Gesellsch. Jahrg 1895; W.C. R6ntgen (Tran. by A. Stanton), Nature, **5**3, 274 (1896).
- 2. M.I. Pupin, Science, **3**, 544 (1896).
- 3. M. **S**ie**g**bahn**, Spectroskopie der Ro**ent_enstrahlen (Sprin**g**er Verla**g**, Berlin**,** 1931); Ann. der Phys. 4*,* 59 (1919).
- 4. A.E. Lindh. Roentgenspektroskopie (Akad Verlag, Leipzig, 1930).
- 5. G. Wentzel, Z. Physik 3**1**, 445 (1925).
- 6. R.M. Langer, Phys. Rev. 3**7**, 457 (1931).
- 7. H.C. Wolfe*,* Phys. Rev. 43, 221 (1933).
- 8. R.D. Deslattes, Phys**.** Rev. **1**33, A399 (1964).
- 9. R.C.C. Perera, P.L. Cowan, D.W. Windl**e**, R.E. LaVilla, and R.D. Deslattes, Phys. Rev. A 4**3**, 3609 (1991).
- 10. R.D. Deslattes, R.E. LaVilla, P.L. Cowan, and A. Henins, Phys. Rev. A **27** 923 (1983).
- 11. P.L. Cowan, S. B**r**ennan, R.D. Deslattes, A. Henins, T. Jach, and E.G**.** Kessler, Nucl. Instr. Meth. A **2**46, 154 (1986); P.L. Cowan, S. Brennan, T. Jach, D.W. Lindle, and B.A. Karlin*,* Rev. Sci. Instrum. 6**0**, 1603 (1989).
- 12. P.L. Cowan, J.B. Hastings, T. Jach, and J.P. Kirland*,* Nucl. Instr_ Meth. **20**8, 349 (1983).
- 13. J. Doucet*,* P. Sauvageau, and C**.** Sandorfy, J. Chem. Phys. **58***,* 3708 (1973); R. Jadrny, L. Karlsson, L. Matts**o**n, and K. Siegbahn, Physica Scripta **1**6, 235 (1977); **T. C**vitas*,* H. Gusten**,** and L. Klasnic**,** J. C**he**m. Phys. **67,** 2687 (**1**977).
- 14. U. Rossler, Phys. Status Soilds 4**2**, 345 (1970).
- 15. R.C.C. Perera and B.L. Henke, X-ray Spectrometry 9, 81 (1980).
- 16. T.A. Callcott, C.H. Zhang, K.L. Tsang, E.T. Arakawa, D.L. Ederer, and J. Kern**e**r, Nucl**.** Inst*r*. Meth. B 40*/*41, 398 (1989).
- 17. J.J. Jia*,* W.L. O'Brien, T.A. Callcott, Q.Y. Dong, J.-E. Rubensson, D.R. Mueller, and D.L. Ederer, Phys. Rev. Lett. 6**7**, 731 (1991).
- 18. R.C.C. Perera, M.C. Hettrick and D.W. Lindle, J. de Physique 48, C9-645 (1987).
- **.** 19. **J**N**o**rd**g**ren*,* J. **d**e Ph**y**sique **48, C9-**693 (1**9**87).

and of the

- 20. N. Wasskahl, J.-E. Rubensson, G. Bray, P. Glans, P. Bleckert, R. Nyholm, S. **C**r**amm, N. Ma**r**te**n**s**so**n, and J. No**r**dg**ren**,** P**hys. Re**v**. Le**tt**.** 6**4, 280**7 (**1990**)**.** ⁱ
- 21. J.-E. Rubensson, N. Wassdahl, G. Bray, J. Rindstedt, R. Nyholm, S. Cramm, N. **M**a**rtensso**n*,* **and J. No**r**dg**r**en,** P**hys. Rev.** L**e**tt**. 60***,* **1**75**9** (**1988).**

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- 22 ALS Handbook, Lawrence Berkeley Laboratory, Berkeley, CA. PUB-643 Rev. (1989)
- 23. D.W. Lindle and R.C.C. Perera, Eds., Proceedings of the Workshop on Applications of Photon-in Photon-out Spectroscopy with Third Generation Synchrotron Radiation Sources, Washington, D.C. April 25, 1991, Lawrence Berkeley Laboratory, Berkeley, CA. LBL-31323, Conf-9104248 (1991)
- 24. Approved as an insertion device participating research team, a c*o*llaboration between IBM, LBL, NIST and Univ. of Tennessee.
- 25. E. Hoyer, J. Chin, K. Halbach, W.V. Hassenzahl, D. Humphries, B. Kincaid, H. Lancaster, and D. Plate, Nucl. Instrum. Meth. A291, 383 (1990).
- 26. P. Heimann (unpublished work) 1991.
- 27. A.Warwick, R. DiGennaro and M. Howells, Lawrence Berkeley Laboratory*,* Berkeley, CA. LSBL-040 (1989) unpublished.
- 28. T.A. Calcott, D.L. Ederer, E.T. Arakawa and R.C.C. Perera, NSF Proposal (1990) unpublished.

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