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CONF-8510209--1

DE86 003084

CONF-8510209-1

PHYSICS DESIGN OPTIONS FOR COMPACT IGNITION EXPERIMENTS*

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Invited paper
Presented at

U.S.-Japan Workshop on
Next Step Machine Design
JAERI, Japan
October 21-25, 1985

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*Research sponsored by the Office of Fusion Energy, U.S. Department of Energy, under Contract No. DE-AC05-84OR21400 with Martin Marietta Energy Systems, Inc.

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PHYSICS ASSESSMENTS - DESIGN IMPACT

<u>Physics Area</u>	<u>Design/Engr Impact</u>	<u>Evaluation</u>
Confinement	$a, R, B_t, I_p \dots$ Ignition Margin Auxiliary Power	0-D, 1-D & 1-1/2-D Analyses
Magnetics	Divertor/Limiter Assessments PF Control System	MHD Equilibria Positional Stability
Startup & Operating Scenarios	Pulse Lengths PF & TF Power Supplies	1-D & 1-1/2-D Analyses
β -limit	$a, R, B_t, I_p \dots$ Wall/Limiter/ Divertor Loadings	Troyon Limit MHD Stability (Near Walls)
ICRF Heating	Power Port Size	Ant. Coupling Minority Species 2nd Harmonic
α Physics	Flat-top Time Diagnostics	1-D & 1-1/2-D Analyses

O-D Confinement Analysis

Features included in the model:

- Plasma profiles and geometry.
- Neoclassical resistivity enhancement.
- Various forms of χ_e and χ_i .
- Physics constraints
$$\beta_{crit} \sim I/aB_0,$$
$$n_{murakami} \sim B_0/R_0$$
- Equilibrium plasma current profile.
- Fuel, alphas and impurities (Z_{eff}).

Steady-state global analysis is a useful complement to full 1-1/2-D transport code calculations.

- A simple analytic global model is developed to establish ignition conditions and plasma parameters operating regimes over large regions of parameter space (R_0/a , b/a , aB_0^2/q_* , etc.) under various physics assumptions (x_e , x_i , q_ψ , β_{crit} , n_{crit} , etc.)
- This simple model includes detailed enough physics information to be useful in exploring the options and reproduces many global features and trends of the 1-1/2-D WHIST transport code calculations (especially those of POPCON's).
- For many of the confinement scalings considered, it was possible to generate nearly universal contour plots of ignition, auxiliary power requirements, optimal path to ignition, plasma heating and operating windows, etc. in terms of a small number of parameters (such as aB_0^2/q_* , $\langle n \rangle/n_{mu}$, $\langle T \rangle$, etc.)
 - this is found to be useful for rapid assessment of a particular device and/or classes of devices with equivalent performance
- These contour plots are used to analyze potential physics design space, operating regimes, and plasma performance characteristics of small ($R_0 \sim 1-2$ m), high field ($B_0 \sim 8-13$ T) tokamak ignition experiments

Calculations are based on global analysis

- Local energy balance equation

$$\frac{\partial W}{\partial t} = \frac{\partial}{\partial t} \left[\frac{3}{2} n_e k T_e + \frac{3}{2} n_i k T_i \right] = P_{\text{aux}} + P_{\text{OH}} + P_{\alpha} - P_{\text{con}} - P_{\text{rad}}$$

averaged over a flux surface for a given plasma profiles.

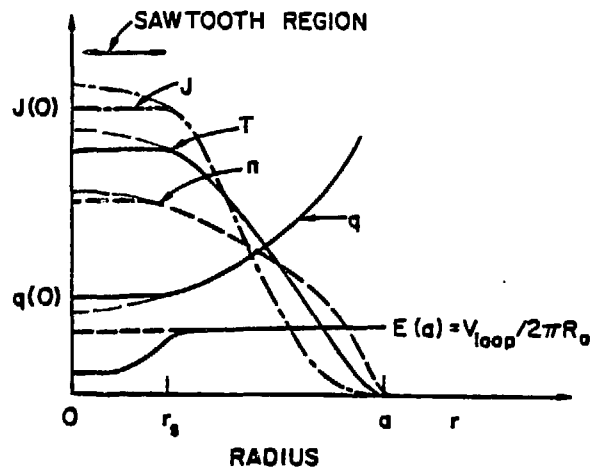
- Geometry:

Concentric flux surfaces with elongation $\kappa = b/a$
and triangularity δ .

- Profiles: Typical profiles assumed are as shown

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- $r \leq r_s$:
 $T = T_0$; $J = J(0)$; $q = q(0)$
- $r > r_s$:
 $n, T \sim (1 - r^2/a^2)^{\alpha_{n,T}}$
 $\alpha_n \sim 0.5$; $\alpha_T \sim 1.0$
 $J \sim T_e^{3/2}$



Definitions:

Ignition Margin

$$M = \frac{P_{\alpha} + P_{OH}}{P(\text{all losses})}$$

MHD safety factor

$$q_{\psi} = q_{*} f(\epsilon, \beta_p, \dots)$$

Equivalent cylindrical safety factor

$$q_{*} = \frac{5aB_0}{I} \left(\frac{a}{R_0} \right) \left[\frac{1 + \kappa^2(1 + 2\delta^2)}{2} \right]$$

"Figure of merit"

$$aB_0^2/q_{*}$$

Constraints/Limits:

Density limit

$$\bar{n}_{20} \leq n_{\text{murakami}} = \nu_{\text{mu}} B_0 / (q_{*} R_0)$$

$\nu_{\text{mu}} \sim 1.5$ (OH)
 $\nu_{\text{mu}} > 1.5$ (aux)

Beta limit

$$\beta_{\text{total}} < \beta_{\text{crit}} = 3I / (aB_0) (\%)$$

Units:

mks units with T in keV,

n_{20} in 10^{20} m^{-3} , I in MA

P in MW.

The flux-surface-averaged energy balance equation (ions plus electrons) is

$$\frac{\partial}{\partial t} \left[\frac{3}{2} n_e k T_e + \frac{3}{2} n_i k T_i \right] = -P_{\text{con}} - P_{\text{rad}} + P_{\alpha} + P_{\text{OH}} + P_{\text{aux}}$$

where

$$P_{\text{con}} = -\frac{1}{V'(\rho)} \frac{\partial}{\partial \rho} \left[A(\rho) \left(n_e \chi_e \frac{\partial k T_e}{\partial \rho} + n_i \chi_i \frac{\partial k T_i}{\partial \rho} \right) \right]$$

$$P_{\text{rad}} \text{ (MW/m}^3\text{)} = 5.3 \times 10^{43} n_e^2 z_{\text{eff}}^2 T_e^{1/2} = 1.68 \times 10^2 n_{e20}^2 z_{\text{eff}}^2 T_{e10}^{1/2}$$

$$P_{\alpha} \text{ (MW/m}^3\text{)} = 0.155 [4 f_D (1-f_D) f_{DT}^2] n_{e20}^2 T_{i0}^3$$

$$P_{\text{OH}} \text{ (MW/m}^3\text{)} \approx 10^6 \eta_{\parallel} \text{ (ohm.m)} J^2 \text{ (MA/m}^2\text{)} \\ \approx 5.22 \times 10^5 \gamma_{\text{NC}} z_{\text{eff}} \ln \Lambda J^2 / T_{e10}^{3/2}$$

$$V(\rho) = 2\pi^2 R_0 \rho^2 \kappa, \quad V' = \partial V / \partial \rho$$

$$A(\rho) = V'(\rho) \langle (\nabla \rho)^2 \rangle, \quad \langle (\nabla \rho)^2 \rangle = (1+\kappa^2) / 2\kappa^2$$

$$z_{\text{eff}} = \sum n_i z_i^2 / n_e, \quad n_e = \sum n_i z_i$$

$$f_{DT} = n_{DT} / n_e = (n_D + n_T) / n_e, \quad f_D = n_D / n_{DT}$$

$$\eta_{\parallel} = \gamma_{\text{NC}} \eta_{\text{spitzer}}, \quad \gamma_{\text{NC}} \approx [1 - 1.95 (r/R_0)^{1/2} + 0.95 (r/R_0)]^{-1}$$

$$\langle \sigma v \rangle_{DT} \approx 1.1 \times 10^{-22} T_{i0}^3 \text{ (m}^3/\text{s)} \begin{cases} s \sim 3 & 4 < T_i < 10 \text{ keV} \\ s \sim 2 & 10 < T_i < 20 \text{ keV} \end{cases}$$

Global power balance

- multiply energy balance eq. with $V'(r) dr$ integrate over r with assumed profiles

$$n = n_0 (1 - r^2/a^2)^{\alpha_n}, \quad T = T_0 (1 - r^2/a^2)^{\alpha_T}, \quad J = J_0 (1 - r^2/a^2)^{\alpha_j}$$

$$\frac{\partial W}{\partial t} = - P_{\text{con}} - P_{\text{rad}} + P_{\alpha} + P_{\text{OH}} + P_{\text{aux}} \quad (\text{MW})$$

Here:

$$W = W_e + W_i = 0.24 n_{e20} T_{e10} \left(1 + \frac{n_i T_i}{n_e T_e}\right) V \quad (\text{MJ})$$

$$P_{\text{con}} (\text{MW}) = 0.16 n_{e20} T_{e10} (\chi_e + \frac{n_i T_i}{n_e T_e} \chi_i) \left[\frac{4}{a^2} \left(\frac{1+k^2}{2k^2} \right) g_c \left(\frac{r_*}{a}, \alpha_n, \alpha_T \right) \right]$$

$$-- g_c = \alpha_T (1 + \alpha_n + \alpha_T) \left(\frac{r_*^2}{a^2} \right) \left(1 - \frac{r_*^2}{a^2} \right)^{\alpha_n + \alpha_T - 1}$$

[Temperature gradients are evaluated at $r_*/a < 1$

$$-- \text{typically } r_*/a \sim 0.6 - 0.8 \Rightarrow g_c \sim O(1)]$$

$$P_{\text{rad}} (\text{MW}) = C_B n_{e20}^2 T_{e10}^{1/2} z_{\text{eff}} V; \quad C_B(\alpha_n, \alpha_T) \sim (1.6 - 1.9) \times 10^{-2}$$

$$P_{\alpha} (\text{MW}) = C_{\alpha} n_{e20}^2 T_{e10}^5 V; \quad C_{\alpha}(\alpha_n, \alpha_T) \sim 0.22 \quad (s \sim 3)$$

for $z_{\text{eff}} \sim 1.5$
 $\alpha_n \sim 0.5, \alpha_T \sim 1.0$

$$P_{\text{OH}} (\text{MW}) = C_{\text{OH}} z_{\text{eff}} \bar{\gamma}_{\text{Ne}} T_{e10}^{-3/2} \left(\frac{B_0^2}{q_0^2 R_0^2} \right) \left(\frac{1+k^2}{2k} \right)^2 V$$

$$-- C_{\text{OH}} \sim 5.4 \times 10^{-4} \quad (\ln \Lambda \sim 16, \alpha_n \sim 0.5, \alpha_T \sim 1.0)$$

$$-- \bar{\gamma}_{\text{Ne}} \sim 2.8 - 2.2 \quad \text{for } A = R_0/a \sim 2.5 - 3.5$$

Also $P_{\text{con}} \sim \int \frac{nT}{\tau} dV \Rightarrow \chi_j = \frac{3}{8} \frac{a^2}{\tau_{ej}} \left(\frac{2k^2}{1+k^2} \right) g_{\chi} \left(\frac{r_*}{a} \right) \quad j=e, i$

typically: $g_{\chi} \sim 0.8 - 1.1$

Confinement Models

(Representative thermal diffusivities)

- Ions: Neoclassical - Chang-Hinton

$$\chi_{ich} \text{ (m}^2/\text{s)} \approx 2 \times 10^{-2} K_2^* \frac{n_{e20} z_{eff} q_*^2}{\epsilon^{3/2} T_{i0}^{1/2} B_0^2} \left(\frac{2}{1+k^2} \right)$$

$$K_2^* = (0.66 + 1.88 \epsilon^{1/2} - 1.54 \epsilon) (1 + 1.5 \epsilon^2) \quad \epsilon = r/R_0$$

- Electrons: Empirical

- Ohmic -- NeoAlcator

$$\tau_{NA} \approx 0.07 n_{e20} a R_0^2 q_*$$

$$\chi_{NA} \text{ (m}^2/\text{s)} \approx \frac{4.3 a}{n_{e20} R_0^2 q_*} \left(\frac{2k^2}{1+k^2} \right)$$

- Auxiliary -- various L- and H-mode scalings

Typically $\tau_E \sim I^{(1-1.5)} [P^{-(1/3-2/3)} \text{ or } b + a/P]$

$$\chi_e \text{ (L-mode)} \sim (2-4) \chi_e \text{ (OH)}$$

$$\chi_e \text{ (H-mode)} \sim 0.5 \chi_e \text{ (L-mode)}$$

- For example: GMS / Mirnov scaling

$$\tau_{E(mir)}^L \sim 0.12 a I k^{1/2} \Rightarrow \chi_{E(mir)}^L \approx \frac{2.5 a}{I k^{1/2}} \left(\frac{2k^2}{1+k^2} \right)$$

$$\tau_{E(mir)}^H \approx (2-3) \tau_{E(mir)}^L$$

- For example: Kaye-Goldston scaling

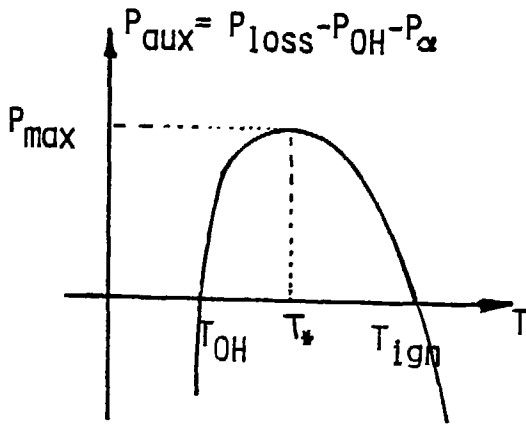
$$\tau_{EKG}^L \approx 0.056 I^{1.24} P^{-0.58} R_0^{1.65} a^{-0.49} k^{0.28} n_{e20}^{0.26} B_0^{-0.09} A_i^{0.5}$$

$$\tau_{EKG}^H \sim 2 \tau_{EKG}^L$$

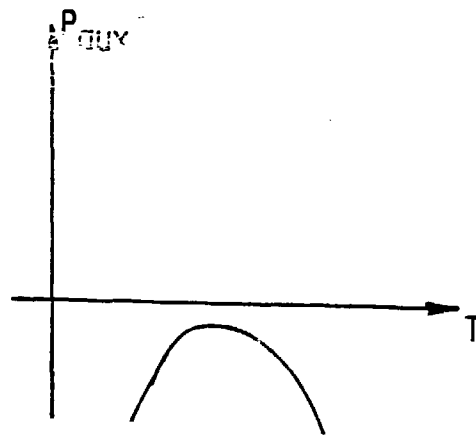
Steady-state (equilibrium) : $\partial W / \partial t = 0$

$$P_{\alpha} + P_{OH} + P_{aux} = P_{con} + P_{rad} = P_{losses}$$

There are two possibilities



$P_{aux} \geq P_{max}$ is required
for ignition



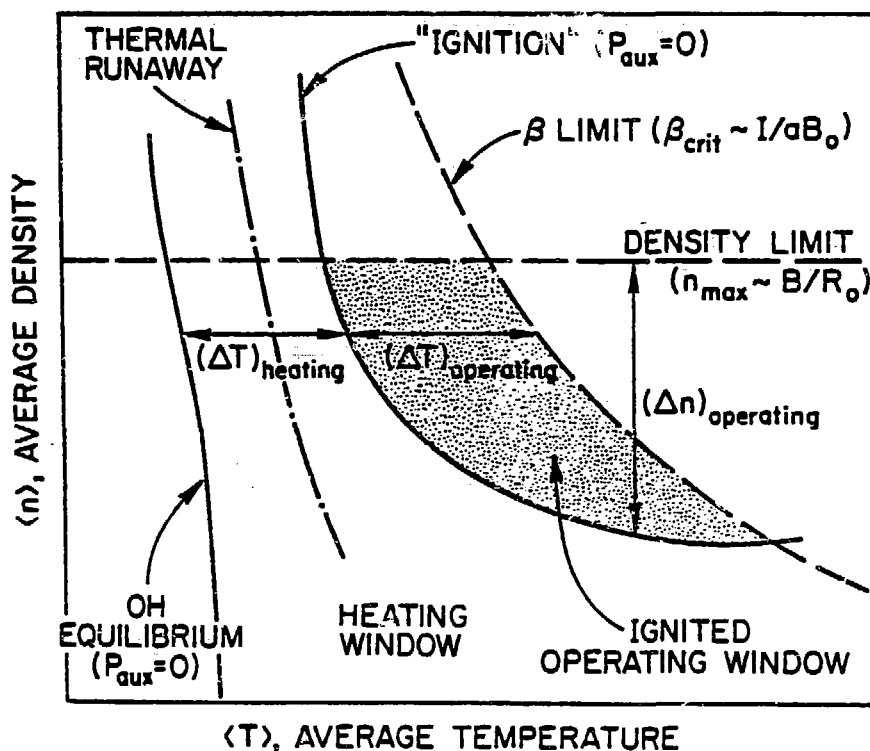
Ohmic ignition

$$P_{max} = P_{aux}(T_*) \rightarrow \partial P_{aux} / \partial T = 0 \text{ at } T = T_*$$

Plasma parameter operating space is constrained by:

- Confinement (ion & electron transport, impurities).
- MHD effects (β limit, q_ψ , etc).
- Density limit (Murakami limit).

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Special Cases:

$\left\{ \begin{array}{l} (\Delta T)_{\text{heating}} \\ (\Delta T)_{\text{operating}} \\ (\Delta n)_{\text{operating}} \end{array} \right\} \rightarrow 0$	Ohmic Ignition
	No Ignition
	No Ignition

Desirable Characteristics:

$\left\{ \begin{array}{l} (\Delta T)_{\text{heating}} \\ (\Delta T)_{\text{operating}} \\ (\Delta n)_{\text{operating}} \end{array} \right\} \rightarrow 0\text{-few keV}$	Low P_{aux}
$\left\{ \begin{array}{l} (\Delta T)_{\text{operating}} \\ (\Delta n)_{\text{operating}} \end{array} \right\} \rightarrow \text{several keV}$	} Ignition margin
$\left\{ \begin{array}{l} (\Delta n)_{\text{operating}} \end{array} \right\} \rightarrow (0.2\text{-}0.5)n_{\text{max}}$	

To cover the range of possibilities for the confinement scalings and to bracket the options for ignition experiments, we choose

- "Optimistic", ohmic-like scaling
 - NeoAlcator ($\tau \sim n$)
- "Pessimistic", L-mode-like scaling
 - Kaye-Goldston ($\tau \sim P^{-\alpha}$)
 with $P = P_{\alpha} + P_{OH} + P_{aux}$

A generalized ohmic / auxiliary scaling

$$\frac{1}{\tau_E^2} = \frac{1}{\tau_{EOH}^2} + \frac{1}{\tau_{EAUX}^2}$$

$\tau_{EOH} \Rightarrow$ neo-Alcator

$\tau_{EAUX} \Rightarrow f \tau_{EKG}^L$; $f = 1$ L-mode
 $= 2$ H-mode

Steady-state power balance equation

$$F = -P_{\text{con}} - P_{\text{rad}} + P_{\alpha} + P_{\text{OH}} + P_{\text{aux}} = 0$$

can be reformulated in terms of a small number of parameters, thus identifying classes of devices with equivalent performance.

One such formulation

$$F = F(m, T, \bar{X})$$

$$m = \frac{\langle n \rangle}{n_{\text{murakami}}} = \text{normalized density} \\ \text{("Murakami parameter")}$$

$$T = \langle T \rangle = \langle nT \rangle / \langle n \rangle \\ = \text{density-averaged temperature}$$

$$\bar{X} = \text{"Figure of merit" parameter} \sim (n\tau)$$

For neo-Alcator, Kaye-Goldston, Mirnov, etc.-like scalings

$$\bar{X} \equiv aB_0^2 / q_*$$

Similar parameterization (in terms of these or any other similar quantities) is also possible for other confinement scalings

Examples for ohmic-like scaling

a) $Z_E = Z_{\text{neoclassical}}$

b) $\chi_e = f_{\text{ex}} \chi_{\text{neoclassical}}$

$$\chi_i = f_{\text{ix}} \chi_{\text{chang-Hinton}}$$

• Functional form :

$$(Z \sim n a R_0^2 q_*)$$

• Typical parameter range ("standard" tokamak)

$$A \sim 3 \pm 0.5$$

$$k \sim 1.6 - 1.7$$

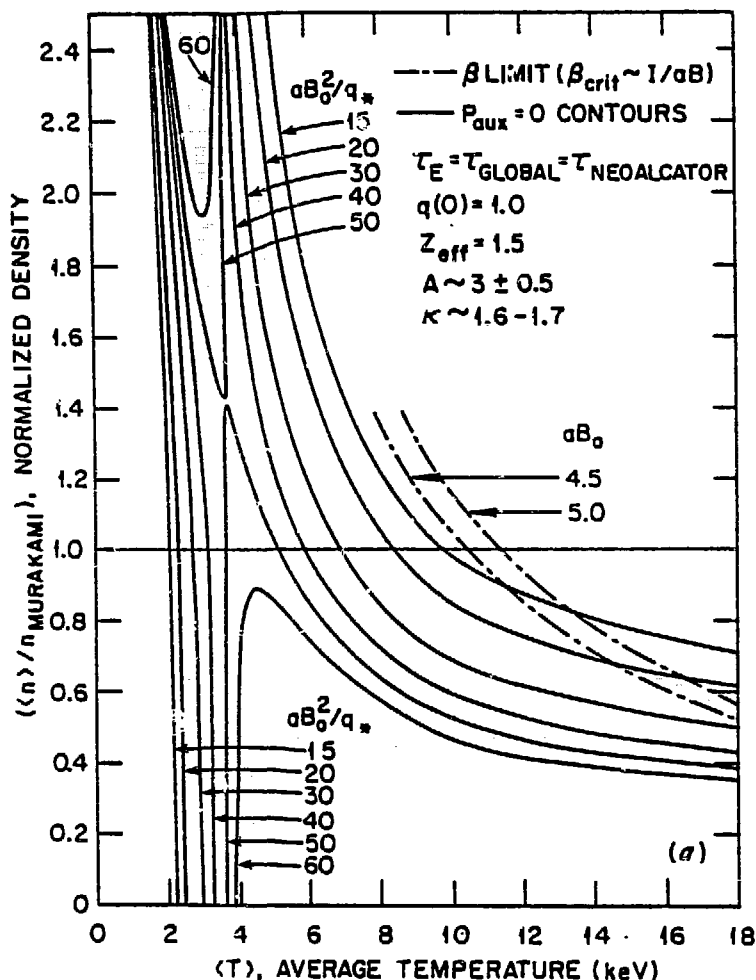
$$q_w \sim 2.6$$

$$z_{\text{eff}} \sim 1.5$$

$$(a B_0 \sim 4.5 - 5.0)$$

Steady-state $P_{aux} = 0$ contours for various values of aB_0^2/q_* showing:

- Requirements for ohmic ignition.
- Relative size of heating and operating windows.
- Optimal route (min P_{aux}) to ignition.

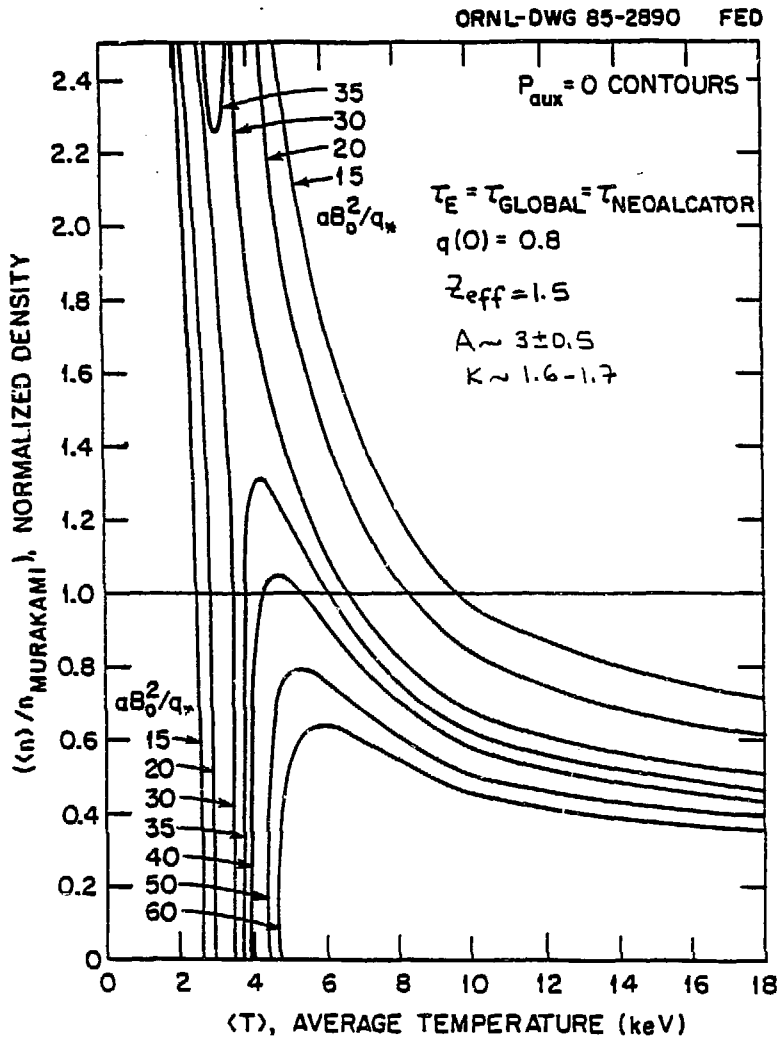


OH ignition
if $aB_0^2/q_* > 57$

Confinement model: $T_E = T_{Global}$
 $= T_{Neoalculator}$
 (ohmic-like scaling) $q(0) = 1.0$

Steady-state $P_{aux} = 0$ contours for various values of aB_0^2/q_* showing:

- Requirements for ohmic ignition.
- Relative size of heating and operating windows.
- Optimal route (min P_{aux}) to ignition.



OH ignition
 f_i
 $\frac{aB_0^2}{q_*} > 41$

Confinement model: $T_E = T_{Global}$
 (ohmic-like scaling) $= T_{Neoalculator}$
 $q(0) = 0.8$

Summary -- Steady-state $P_{aux} = 0$ Contours

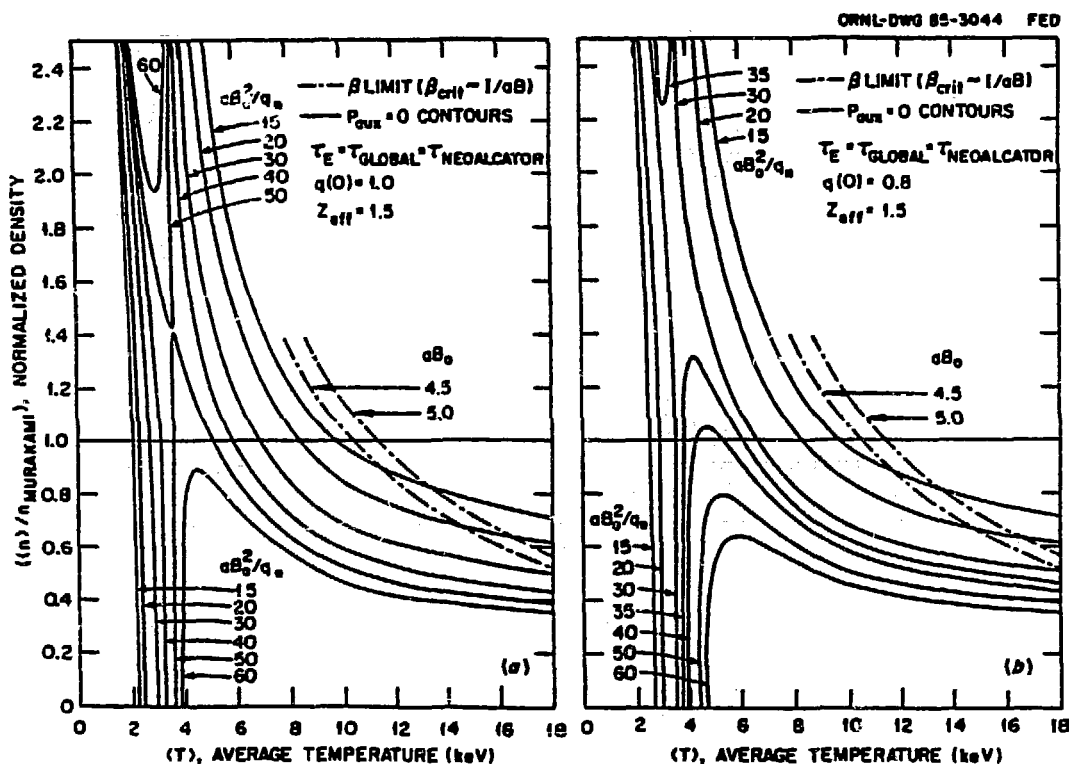
Conf. model : $\tau_E = \tau_{global} = \tau_{neoclassical}$

(a) $q(0) = 1.0$ (b) $q(0) = 0.8$

Superimposed are $\beta = \beta_{crit} = 3 I / a B_0$ (%) contours

(Geometry: $A \sim 3 \pm 0.5$, $\kappa \sim 1.6-1.7$)

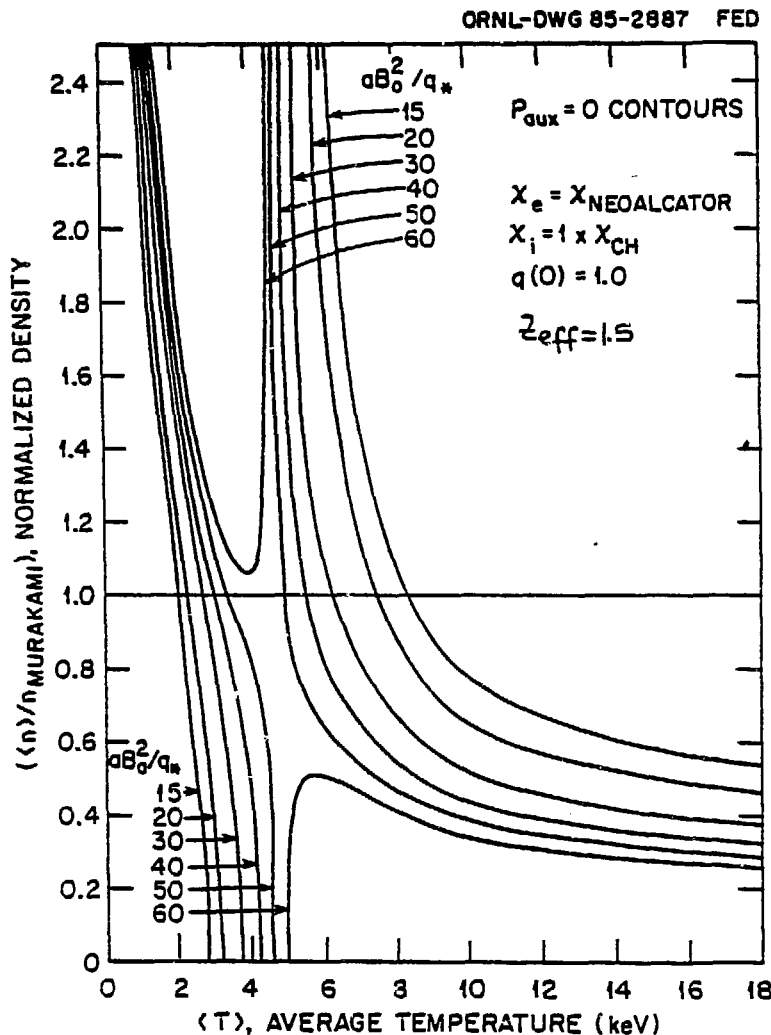
(ohmic-like scaling)



$\tau_{global} = \tau_E = \frac{W}{P_{con}}$ (global energy conf. time -- due to conduction)

Steady-state $P_{aux} = 0$ contours for various values of aB_0^2/q_* showing:

- Requirements for ohmic ignition.
- Relative size of heating and operating windows.
- Optimal route (min P_{aux}) to ignition.

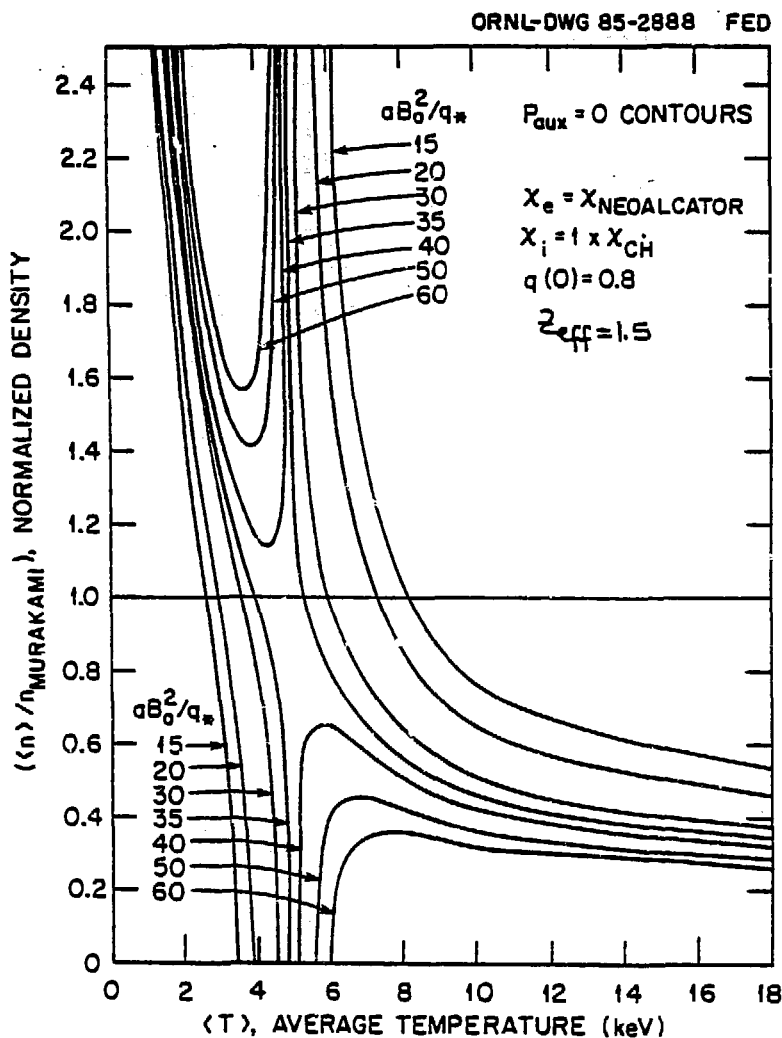


• OH ignition
if
 $\frac{aB_0^2}{q_*} > 52$

Confinement model: $X_e = X_{Neoalculator}$
 (ohmic-like) $X_i = 1 \times X_{Chang-Hinton}$
 $q(0) = 1.0$

Steady-state $P_{aux} = 0$ contours for various values of aB_0^2/q_* showing:

- Requirements for ohmic ignition.
- Relative size of heating and operating windows.
- Optimal route (min P_{aux}) to ignition.

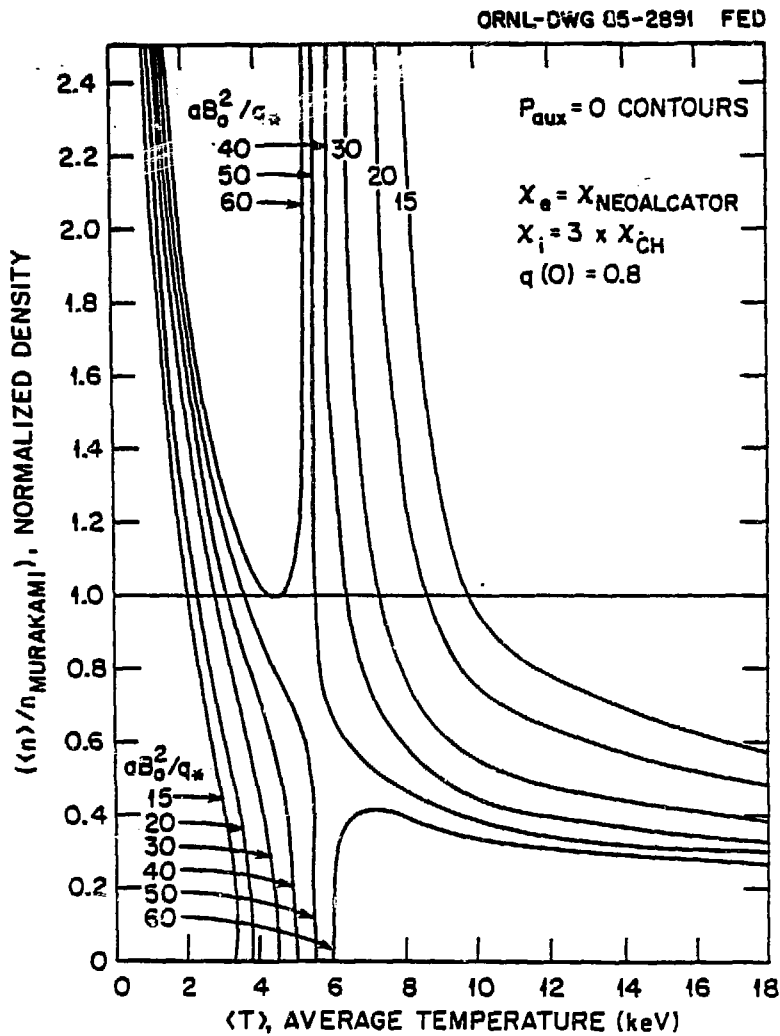


• OH ignition
if
 $\frac{aB_0^2}{q_*} > 37$

Confinement model: $X_e = X_{Neoalculator}$
 (ohmic-like) $X_i = 1 \times X_{Chang-Hinton}$
 $q(0) = 0.8$

Steady-state $P_{aux} = 0$ contours for various values of aB_0^2/q_* showing:

- Requirements for ohmic ignition.
- Relative size of heating and operating windows.
- Optimal route (min P_{aux}) to ignition.



OH ignition
if
 $\frac{aB_0^2}{q_*} > 55$

Confinement model: $X_e = X_{Neoalcator}$
 (ohmic-like) $X_i = 3 \times X_{Chang-Hinton}$
 $q(0) = 0.8$

SUMMARY -- Steady-state $P_{aux} = 0$ contours

Conf. models (a) $\chi_e = \chi_{neo-alcator}$ $\chi_i = 1 \times \chi_{chang-hinton}$

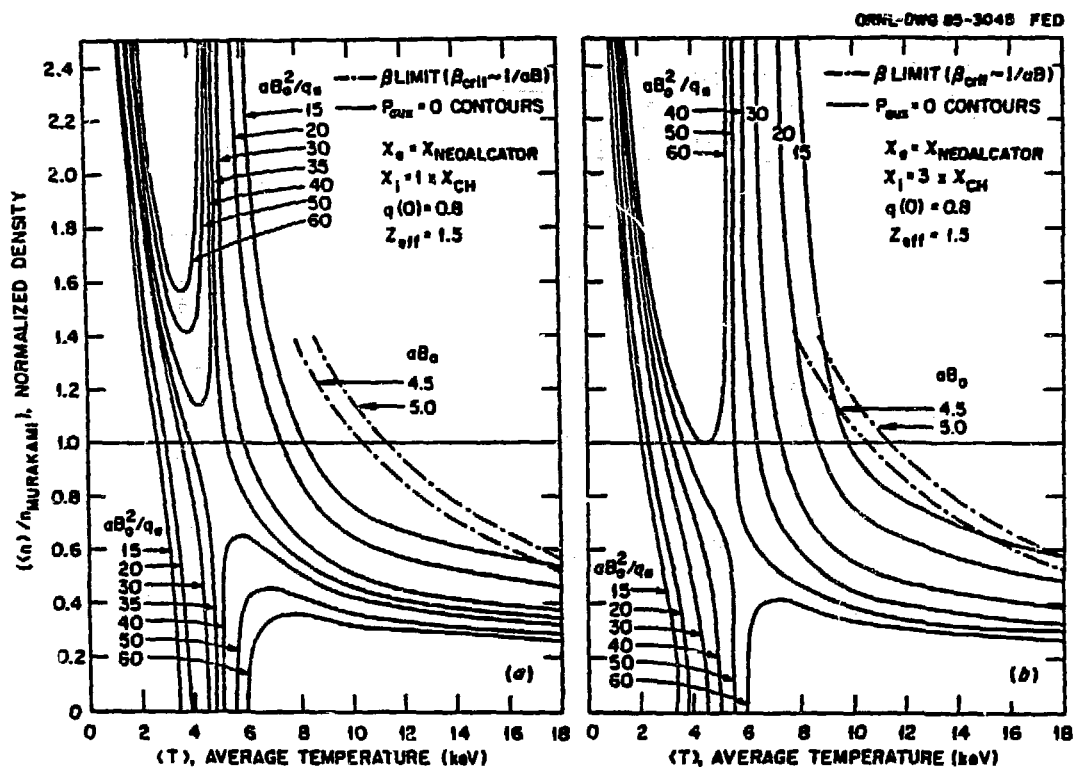
(b) $\chi_e = \chi_{neo-alcator}$ $\chi_i = 3 \times \chi_{chang-hinton}$

Superimposed are $\beta = \beta_{crit} = 3I/aB_0$ (%) contours.

Optimal density path to ignition (narrowest ΔT heating window) is clearly evident and it occurs at

(a) $\langle n \rangle \approx 0.85 n_{murakami}$

(b) $\langle n \rangle \approx 0.65 n_{murakami}$

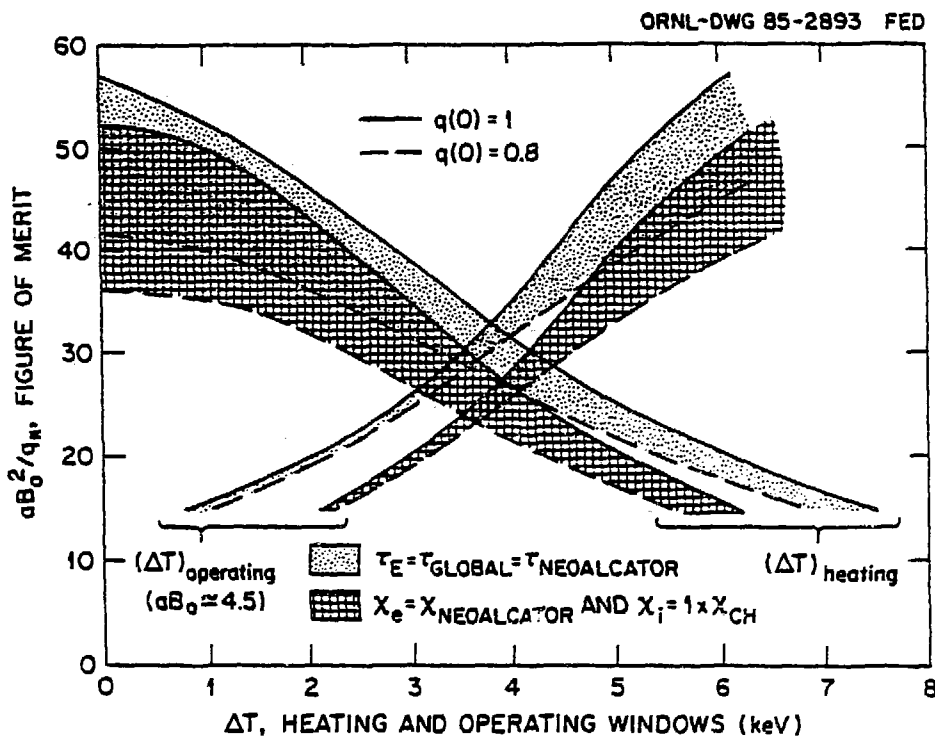


(ohmic-like-scaling)

Devices with large aB_0^2/q_* have favorable heating and operating windows.

As aB_0^2/q_* increases

- (ΔT) heating window decreases
 - ▶ small P_{aux} requirements
 - ▶ margin against uncertainties
- (ΔT) operating window increases
 - ▶ margin against uncertainties
 - ▶ separation of ignition physics from β and density limits



(ohmic-like scaling)

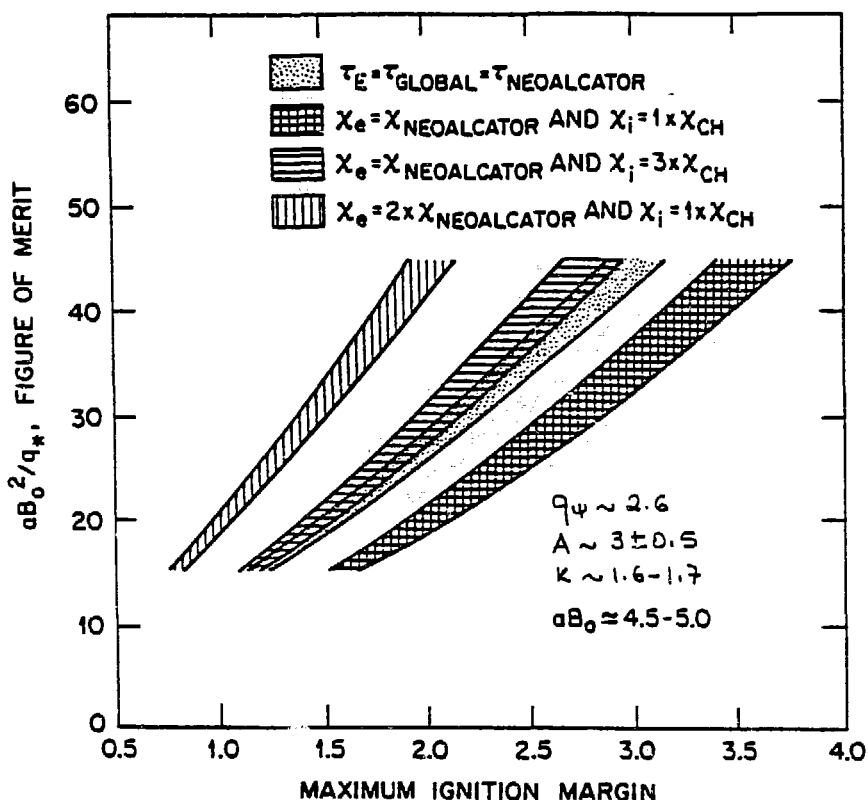
$(\Delta T)_{op}$ in the figure is measured at the Murakami limit.

$(\Delta T)_{op}$ is larger along the optimal density path.

Note: Assumption of $q(0)=1.0$ or $q(0)=0.8$ has very little or no impact on $(\Delta T)_{op}$ for $aB_0^2/q_* < 25$.

Maximum attainable margin for ignition within the plasma operating window increases with increasing aB_0^2/q_* .

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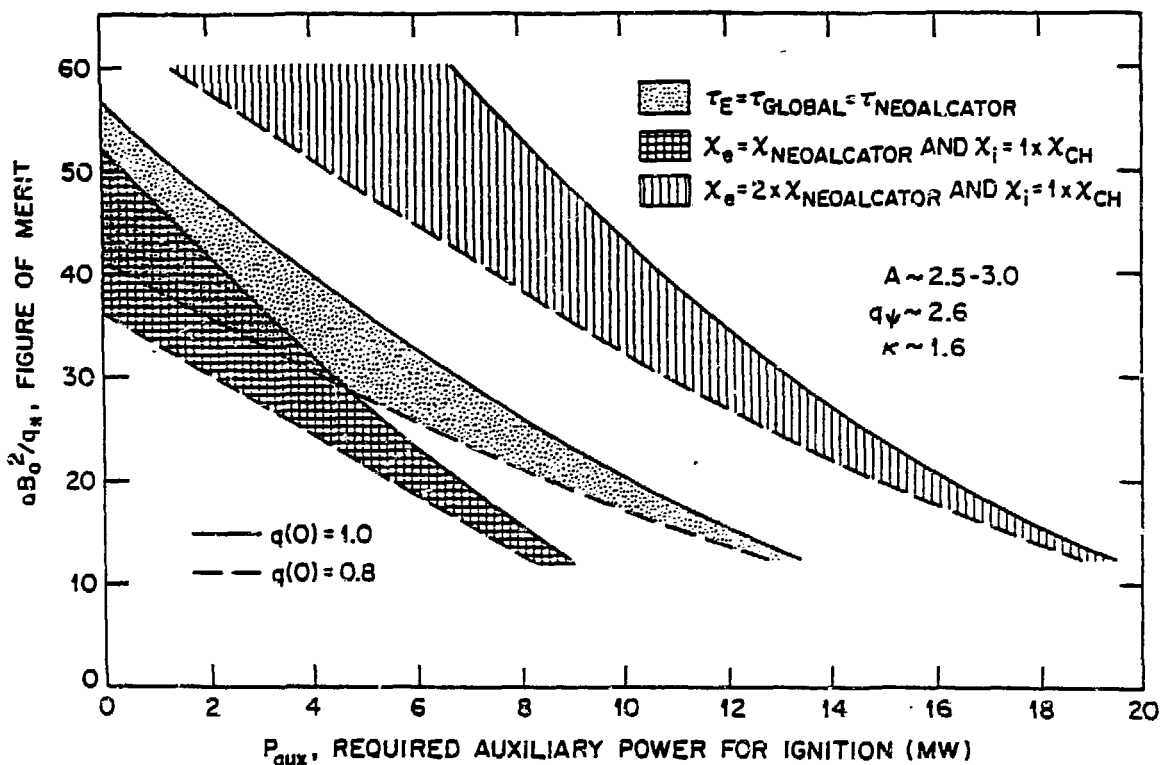


Based on the ohmic-like confinement scaling devices with $aB_0^2/q_* \sim 20 \pm 5$ appear ignitable with a margin $M \sim 1.5 \pm 0.5$. Required auxiliary power $P_{aux}(\text{equilibrium}) \sim 10 \pm 5$ MW.

Auxiliary power required for ignition decreases as aB_0^2/q_* increases.

- Figure shows the minimum auxiliary power required for ignition as determined by the maximum equilibrium power along the optimal density path for several confinement models.
- To raise the temperature to ignition in a finite time more auxiliary power is required

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Based on the ^{ohmic-like} confinement scalings considered, possibility of ohmic ignition exists for devices with

$$aB_0^2/q_* \sim 40 \pm 10$$

(much of the uncertainty can be removed by some amount of added auxiliary power)

Ohmic + Auxiliary Scaling*

$$\frac{1}{\tau_E^2} = \frac{1}{\tau_{EOH}^2} + \frac{1}{\tau_{EAUX}^2}$$

$$\tau_{OH} = \tau_{neoclassical} \sim n a R_0^2 q_*$$

$$\tau_{aux} = f \tau_{Kaye - Goldston} \sim I^{1.24} P^{-0.58} R^{1.65} a^{-0.49}$$

$$\Rightarrow \tau_{KG} \sim I^{2.95}$$

- Strong dependence on plasma shape (κ, δ)

Consider 3 Geometries with Increasing κ

- I. "Standard" tokamak : $\kappa \sim 1.6-1.7$; $\delta \sim 0.3-0.1$
- II. "Elongated" " : $\kappa \sim 1.9-2.1$; $\delta \sim 0.2-0.0$
- III. "More-elongated" " : $\kappa \sim 2.3-2.5$; $\delta \sim 0.2-0.0$

Typical Parameters

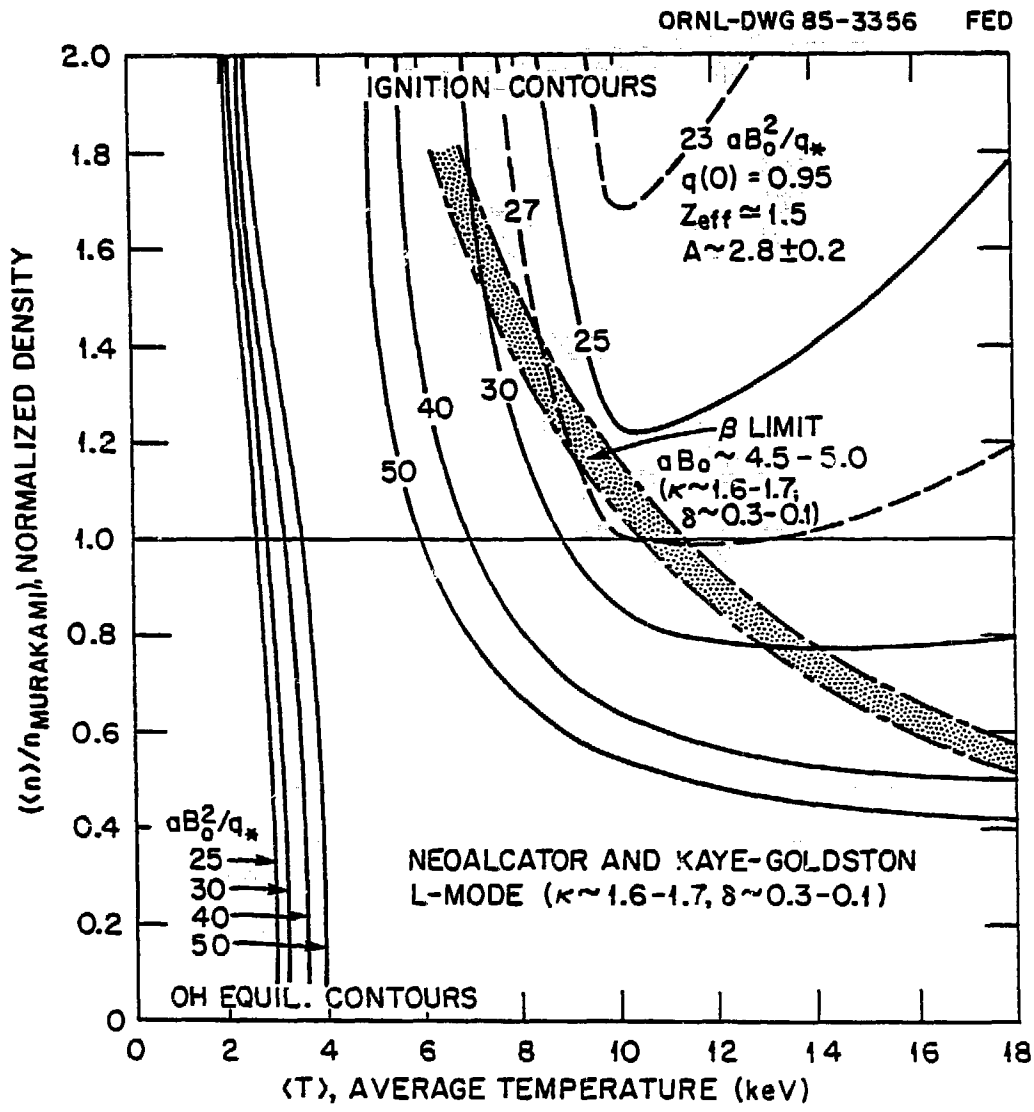
$$A \sim 2.8 \pm 0.2$$

$$q_\psi \sim 2.6$$

$$z_{eff} \sim 1.5$$

* Note : P_α and P_{OH} is included in P , in addition to P_{aux} , thus the model can be considered as pessimistic.

Steady-state ignition / OH equilibrium contours
for various values of aB_0^2/q_* .

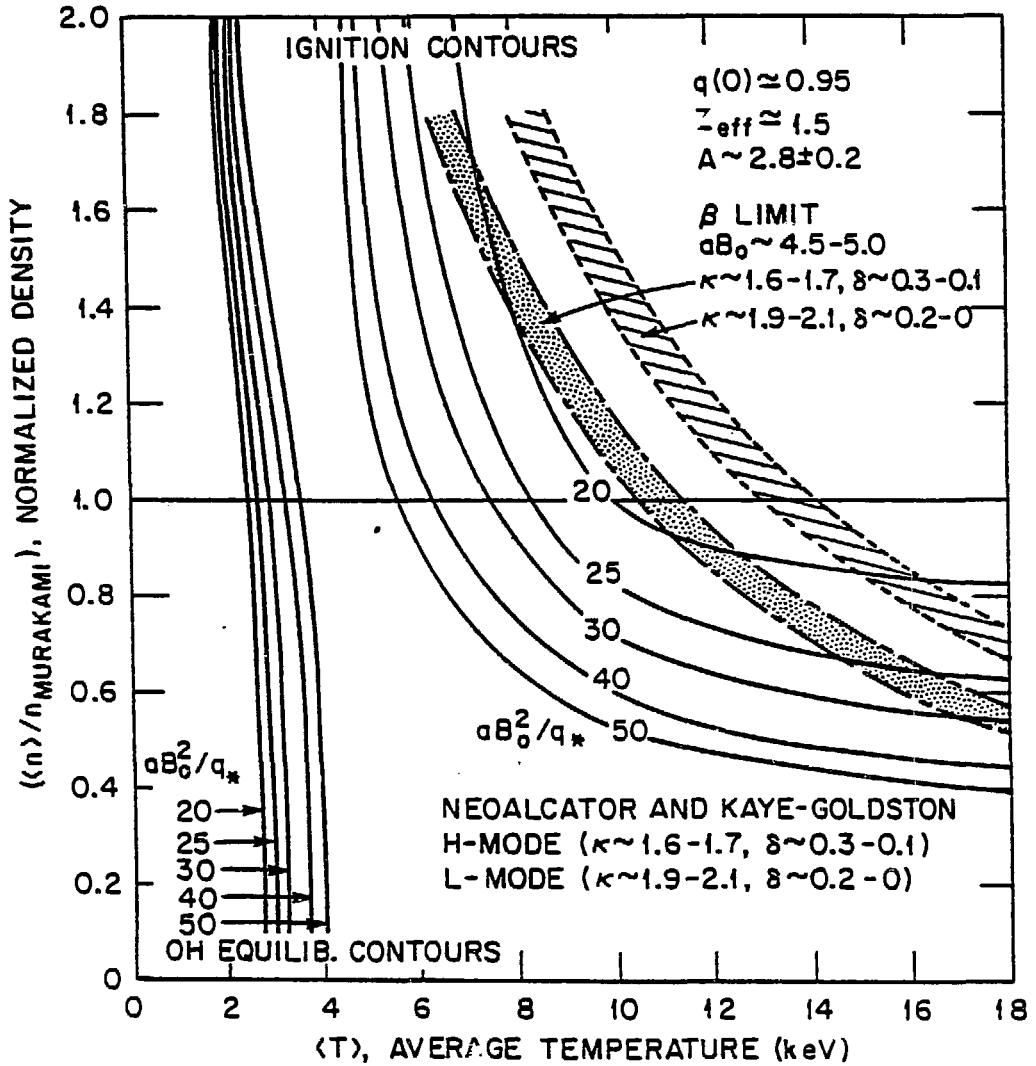


Confinement model : NeoAlcator + Kaye-Goldston
L - MODE

Geometry : "Standard" tokamak

Steady-state $P_{aux}=0$ (ignition/OH equilibrium) contours for various values of aB_0^2/q_* in two different geometries. Corresponding $\beta = \beta_{crit}$ contours are shown.

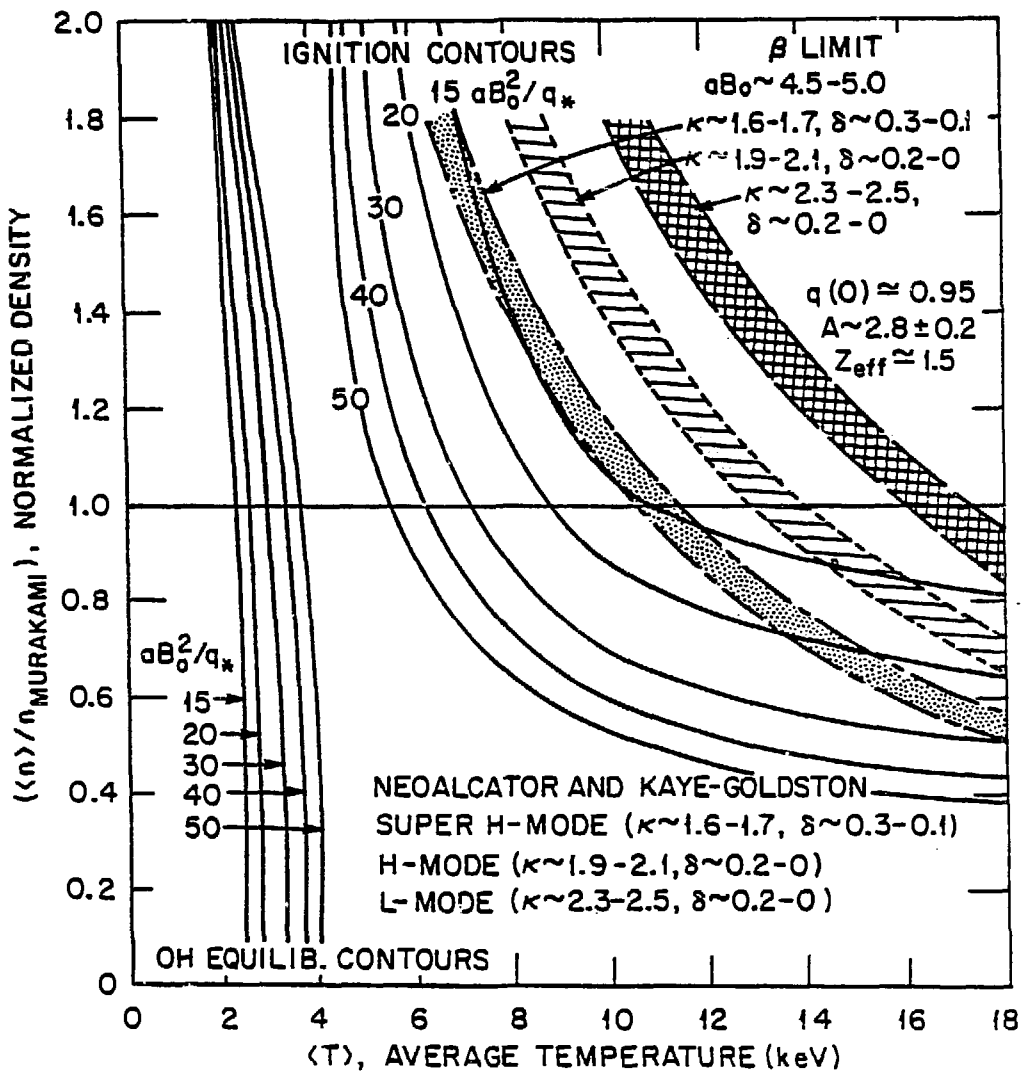
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Confinement model: Neo-Alcator + Kaye-Goldston
 Geometry: "Standard" tokamak - H-mode
 "Elongated" tokamak - L-mode

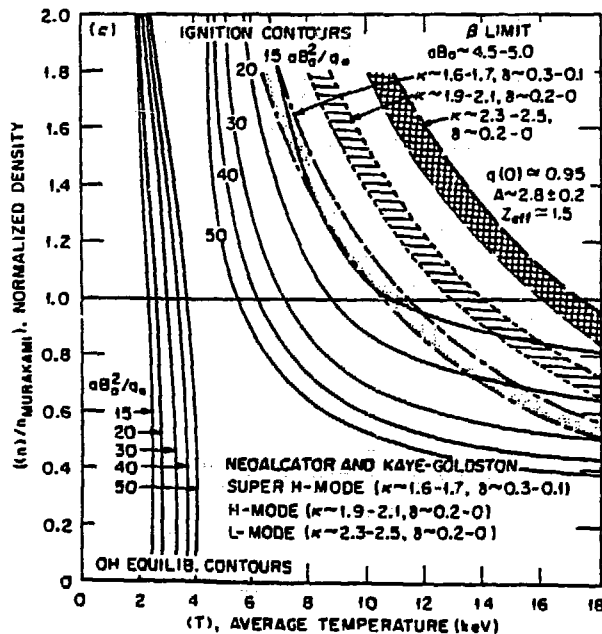
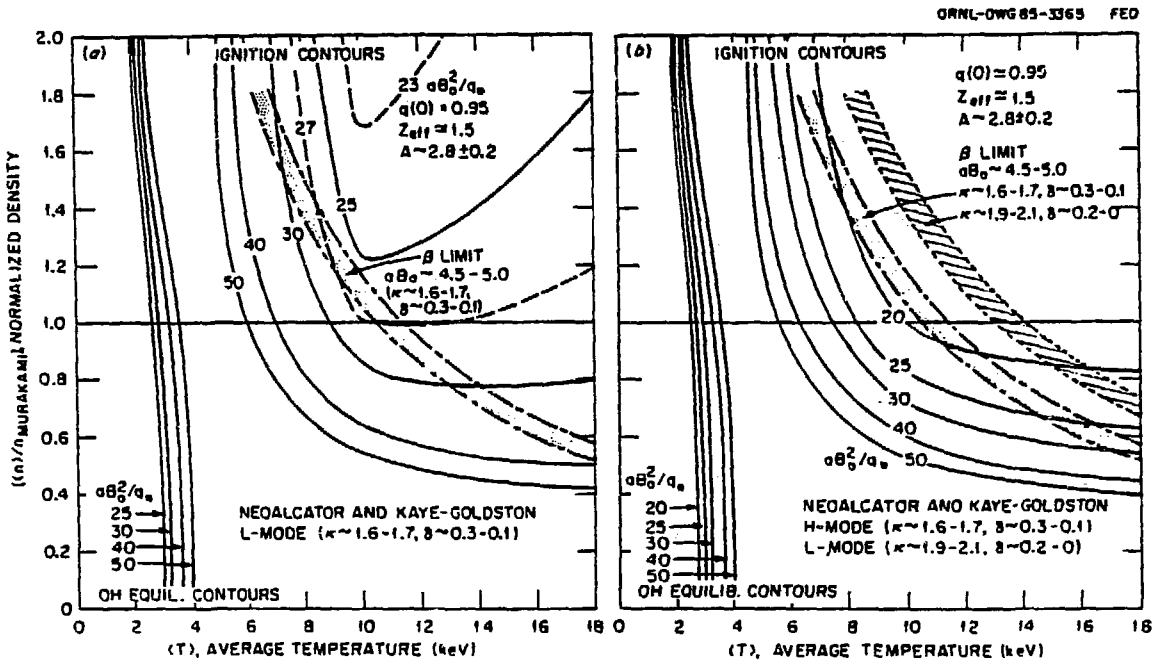
Steady-state $P_{aux}=0$ (ignition/DH equilib.) contours for various values of aB_0^2/q_* in 3 different geometries. Corresponding $\beta = \beta_{crit}$ contours are shown.

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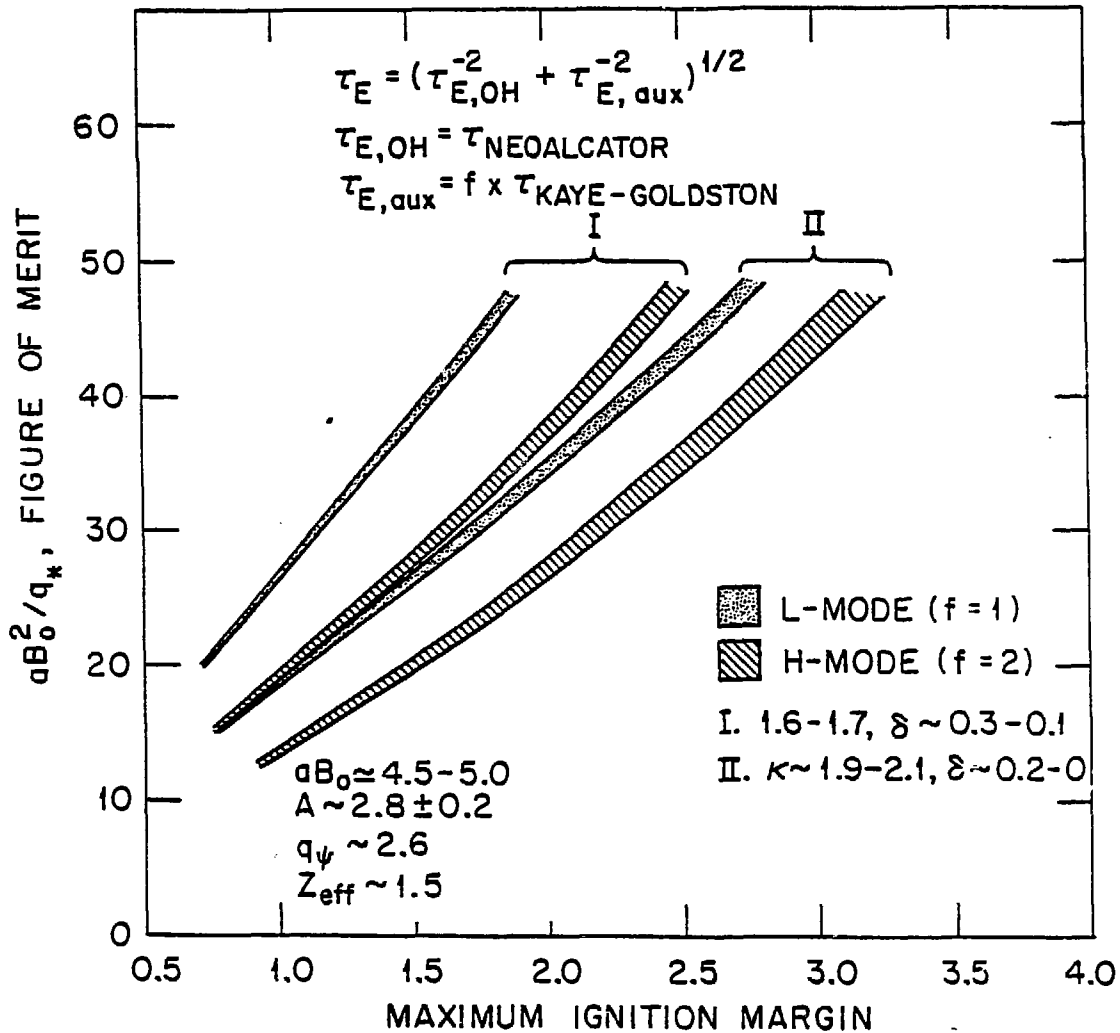
Confinement model: Neo-Alcator + Kaye-Goldston
 Geometry : Standard : Super-H-mode
 Elongated : H-mode
 More-Elongated: L-mode

Summary: Ohmic + Auxiliary scaling $P_{aux}=0$ contours and $\beta = \beta_{crit}$ contours for various values of $a\beta_0^2/q_*$ (classes of devices with equivalent performance) in three different tokamak geometry with increasing k .



Maximum attainable margin for ignition within the plasma operating window increases with increasing aB_0^2/q_* .

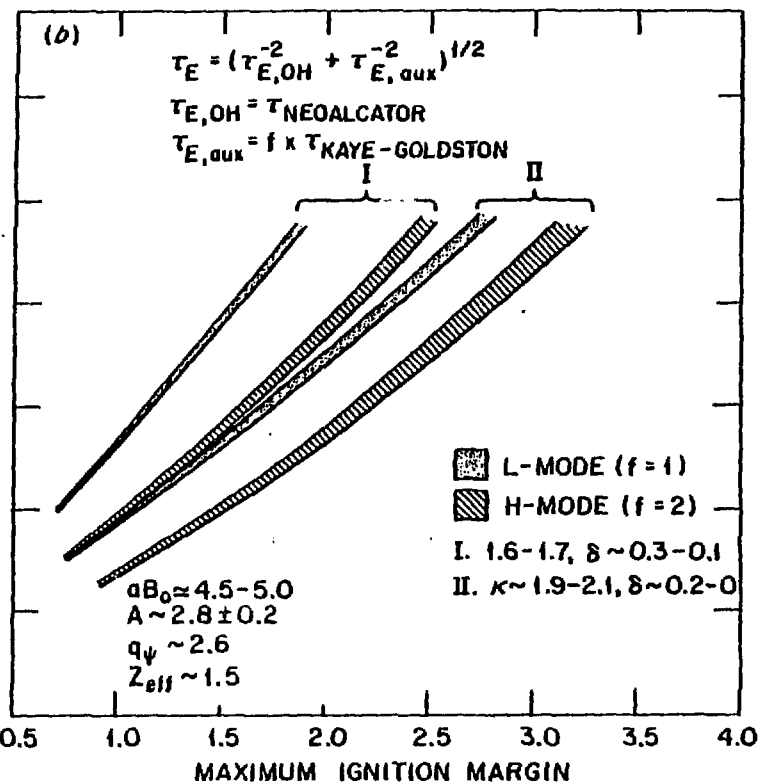
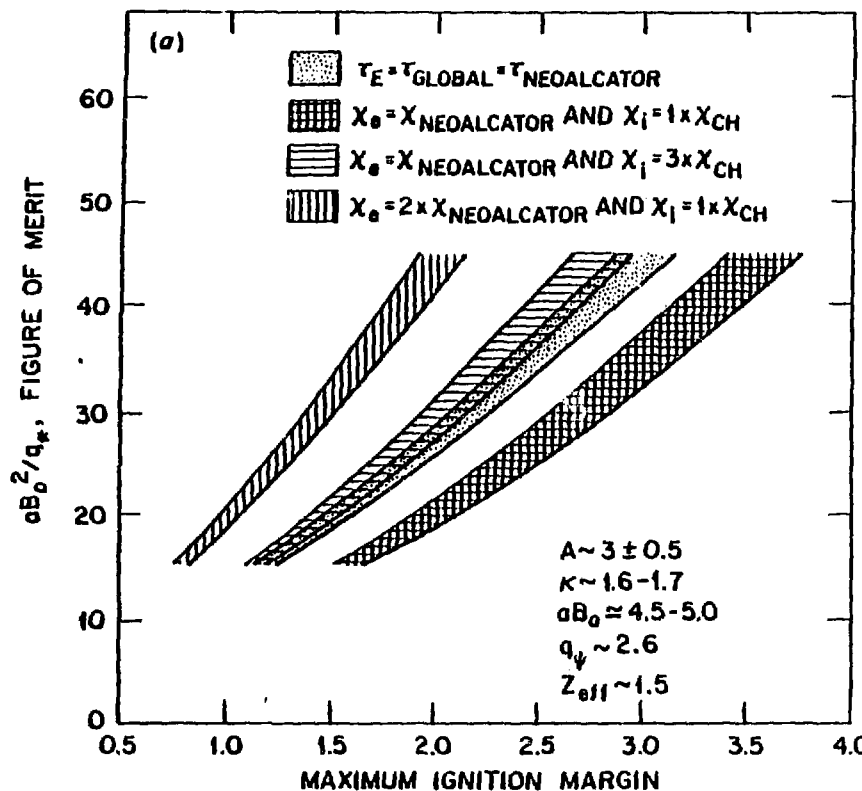
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Based on ohmic + auxiliary scalings considered devices with $aB_0^2/q_* \sim 25 \pm 5$, $\kappa \sim 1.8 \pm 0.2$ appear ignitable.

Summary
 Comparison of ignition margins for "ohmic-like"
 and "ohmic+ auxiliary" scalings.

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Applications to Compact Tokamak Ignition Experiments

At present several candidate design options for a compact, high-field ignition experiment are being considered by the U.S. Tokamak Ignition Studies Design Teams. These options include

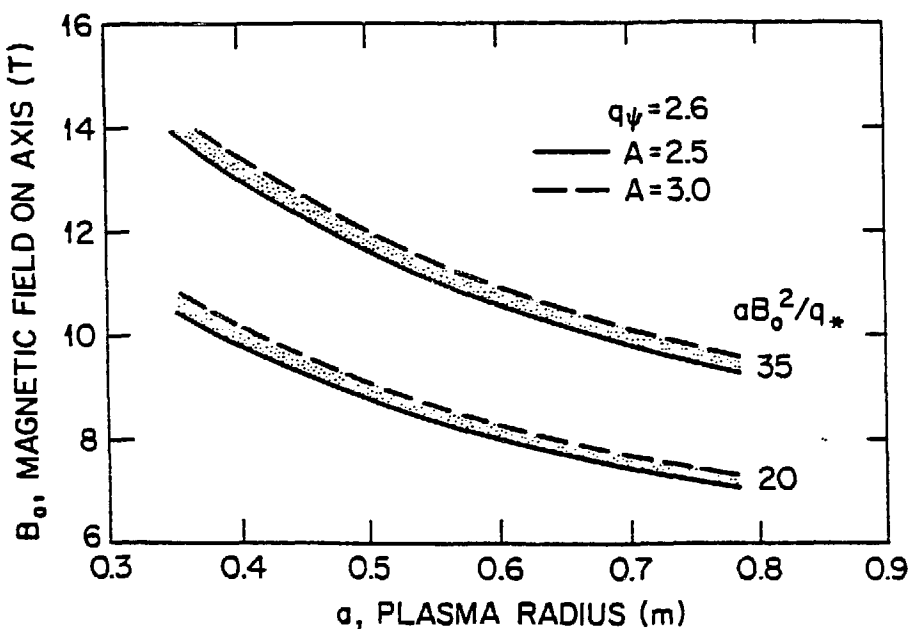
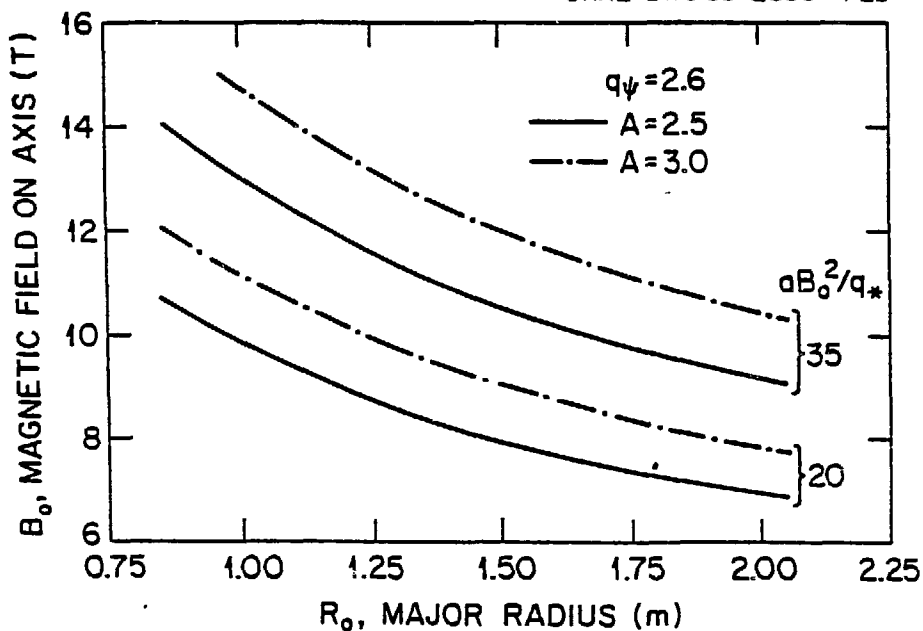
- Ignitor-A \rightarrow $a \sim 0.39 \text{ m}$, $R_0 \sim 1 \text{ m}$, $B_0 \sim 12.6 \text{ T}$
 $I \sim 10 \text{ MA}$, $q_\psi \sim 2.6$, $k \sim 1.67$, $\delta \sim 0.25$
 $\Rightarrow aB_0^2/q_* \sim 32$
- PPPL - ISP \rightarrow $a \sim 0.53 \text{ m}$, $R_0 \sim 1.6 \text{ m}$, $B_0 \sim 9 \text{ T}$
 (0424) $I \sim 8 \text{ MA}$, $q_\psi \sim 2.6$, $k \sim 1.6$, $\delta \sim 0.4$
 $\Rightarrow aB_0^2/q_* \sim 20$
- MIT - LITE \rightarrow $a \sim 0.55 \text{ m}$, $R_0 \sim 1.75 \text{ m}$, $B_0 \sim 8.5 \text{ T}$
 $I \sim 7 \text{ MA}$, $q_\psi \sim 2.6$, $k \sim 1.6$, $\delta \sim 0.3$
 $\Rightarrow aB_0^2/q_* \sim 19$

Can be categorized in two classes

- I. Devices with $aB_0^2/q_* \sim 20$
- II. Devices with $aB_0^2/q_* \sim 32$

For typical range of aspect ratios ($A \sim 2.5 - 3$) and MHD safety factor ($q_\psi \sim 2.6$), the range of device parameters and plasma current corresponding to $aB_o^2/q_* \sim 20 - 35$ are: $R_o \sim 1 - 2$ m ($a \sim 0.35 - 0.65$ m)
 $B_o \sim 8 - 14$ T

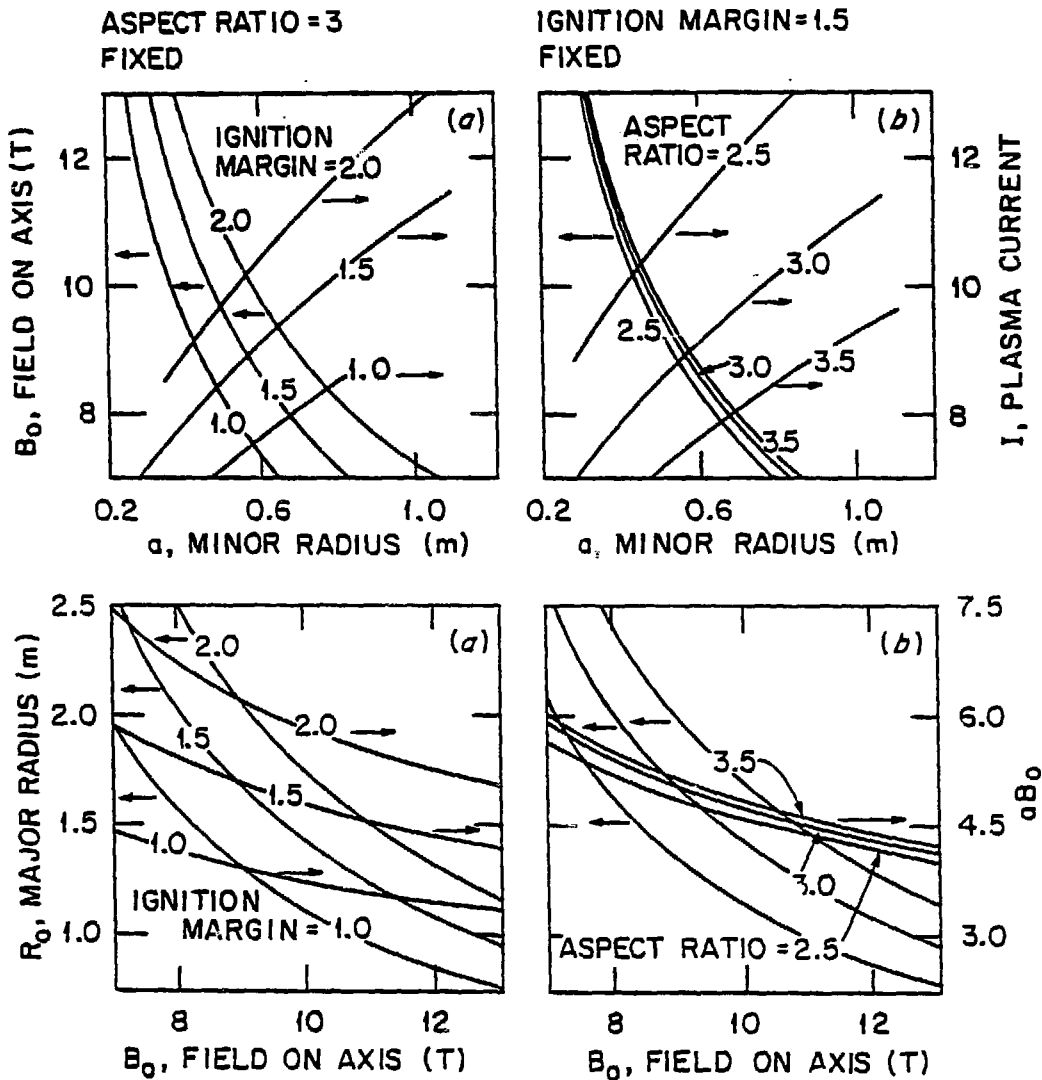
$I \sim 7-10$ MA ($K \sim 1.6-1.7$); $\sim 9-13$ MA ($K \sim 1.9-2.1$); $\sim 12-17$ MA ($K \sim 2.3-2.5$)
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For most parameter range of interest $aB_o \sim 4.5 - 5$ (mT)

Variations of parameters with ignition margin and aspect ratio

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- a) $A = \text{fixed}$: increasing ignition margin increases a, B, I
 \Rightarrow require larger aB_0^2/q_*
- b) $M_I = \text{fixed}$: increasing A increases R , decreases I
 but little effect on a, B
 $\Rightarrow aB_0 \sim \text{const.}$ $aB_0^2/q_* \sim \text{const.}$

'STEADY-STATE' AND TIME-DEPENDENT ANALYSES

'Steady-state' contours generated by slow time evolution of 1-1/2-D WHIST transport code and 0-D global model.

- IGNITOR-A, PPPL 0424, LITE, and "MX".
- Sensitivity to χ_e (ohmic & auxiliary).

Time-dependent simulation of 'flat-top':

- Time-to-ignition.
- Auxiliary power requirements.

Full startup analysis:

- TF compression with current and density ramps.
- Volt-second consumption.

GENERAL PARAMETERS

R_o	1 - 1.75 m
a	0.4 - 0.6 m
B_o	8 - 13 T
I	8 - 13 MA
k	1.6 - 2.0

SPECIFIC PARAMETERS

	IGNITOR-A	PPPL-ISP (O424)	MIT-LITE	"MX" (*)
R_o (m)	1.01	1.62	1.76	1.4
a (m)	0.39	0.53	0.55	0.5
$A = R_o/a$	2.6	3.0	3.2	2.8
k	1.67	1.6	1.6	1.8
δ	0.25	0.4	0.3	0.2
B_o (T)	12.6	9.0	8.6	10.0
I	10.0	8.0	7.0	10.0
q_ψ	2.6	2.6	2.6	2.6
z_{eff}	1.5	1.5	1.5	1.5

* Not an engineering design. A reference physics device (Middle of the Box) for calculational purpose.

Example 1 (O-D & 1-1/2-D)

Plasma performance contours for devices with

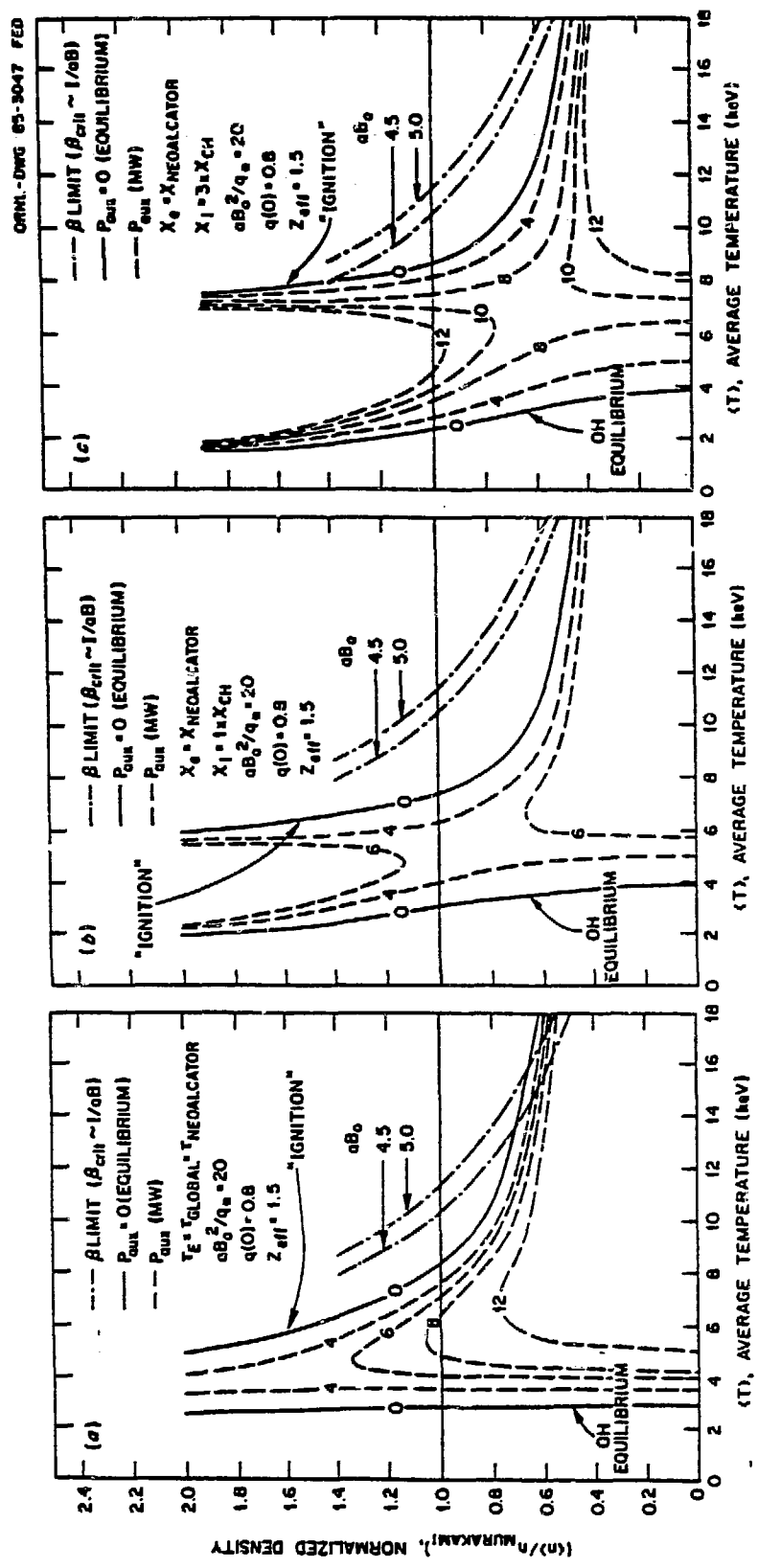
$$aB_0^2/q_* = 20$$

under different scaling assumptions.

Optimal path, P_{aux} , and Max. ignition margins are given.

SUMMARY -- Steady-State Auxiliary power Contours and critical beta Contours for devices with $aB_0^2/q_* = 20$ and $q(0) = 0.8$

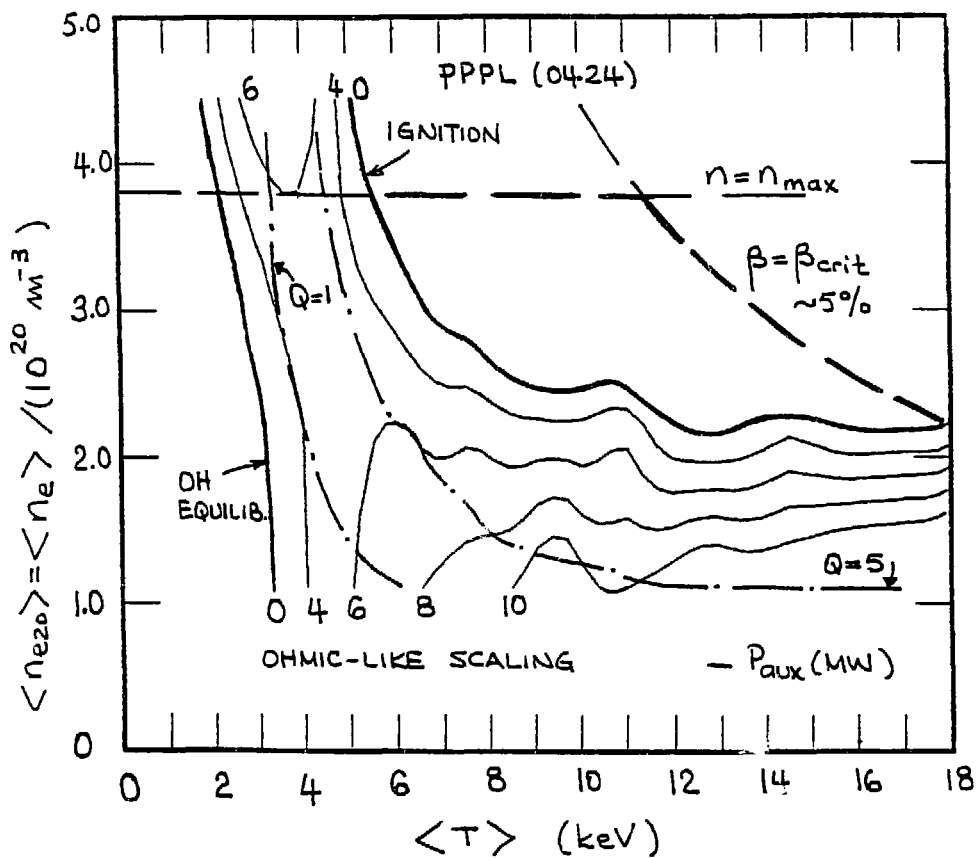
Conf. models: (a) $Z_E = Z_{NEOALCATOR}$, $\chi_i = 1 \times \chi_{CH}$
 (b) $\chi_e = \chi_{NEOALCATOR}$ plus $\chi_i = 3 \times \chi_{CH}$



Steady-state POPCON plots (1-1/2-D code) showing P_{aux} and Q contours for a typical $aB_0^2/q_* = 20$ (PPPL-0424) device.

$n = n_{max}$ (murakami limit) } are shown
 $\beta = \beta_{crit} (\sim 5\%)$

Confinement model: "ohmic-like" scaling
 $\chi_e = 0.8 \chi_{neo-Alcator} g(r)$
 $\chi_i = 0.2 \chi_e + 1 \times \chi_{CH}$

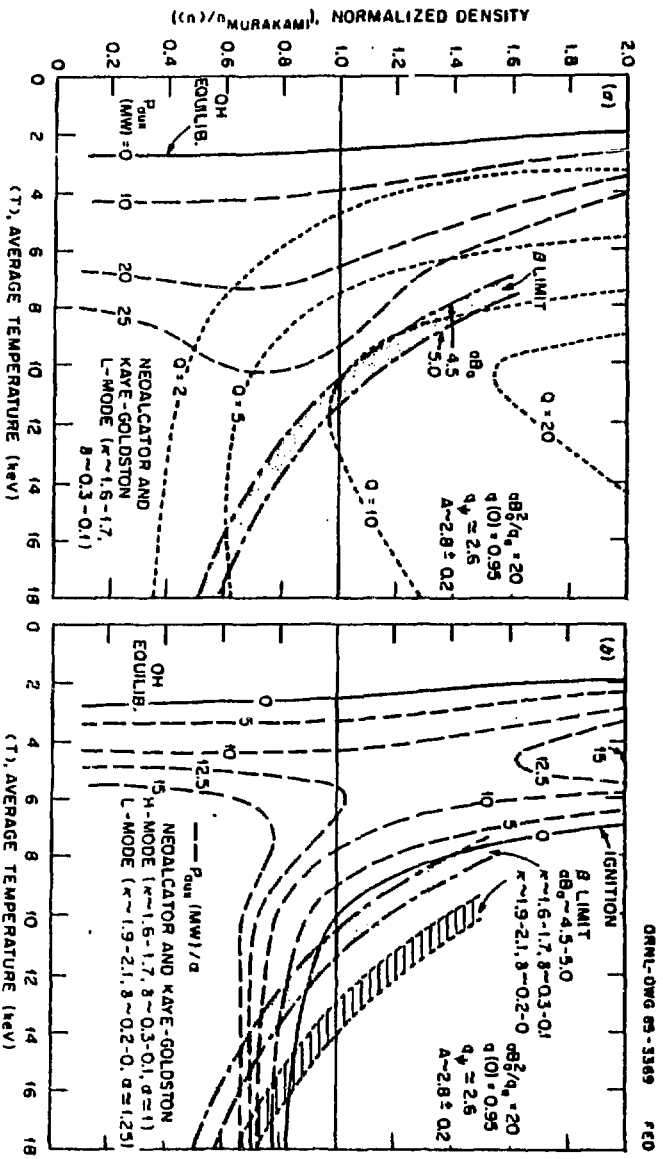


Optimal path $\sim 0.8 n_{max}$
 $P_{aux} \sim 5$ MW

(fluctuations in P_{aux} contours reflect sawtooth activity)

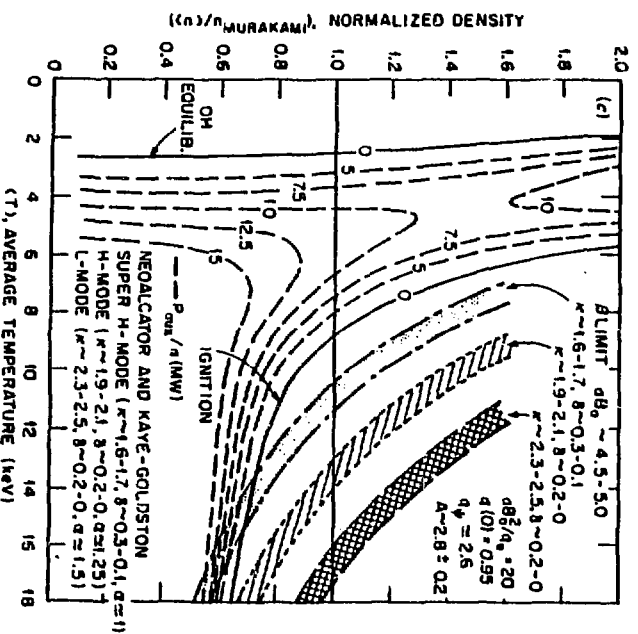
SUMMARY (D-D, analytic model)

Neo-Alcator & Kaye-Goldston scaling in 3 different geometry with increasing elongation for a class of devices with $aB_0^2/q \approx 20$



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120

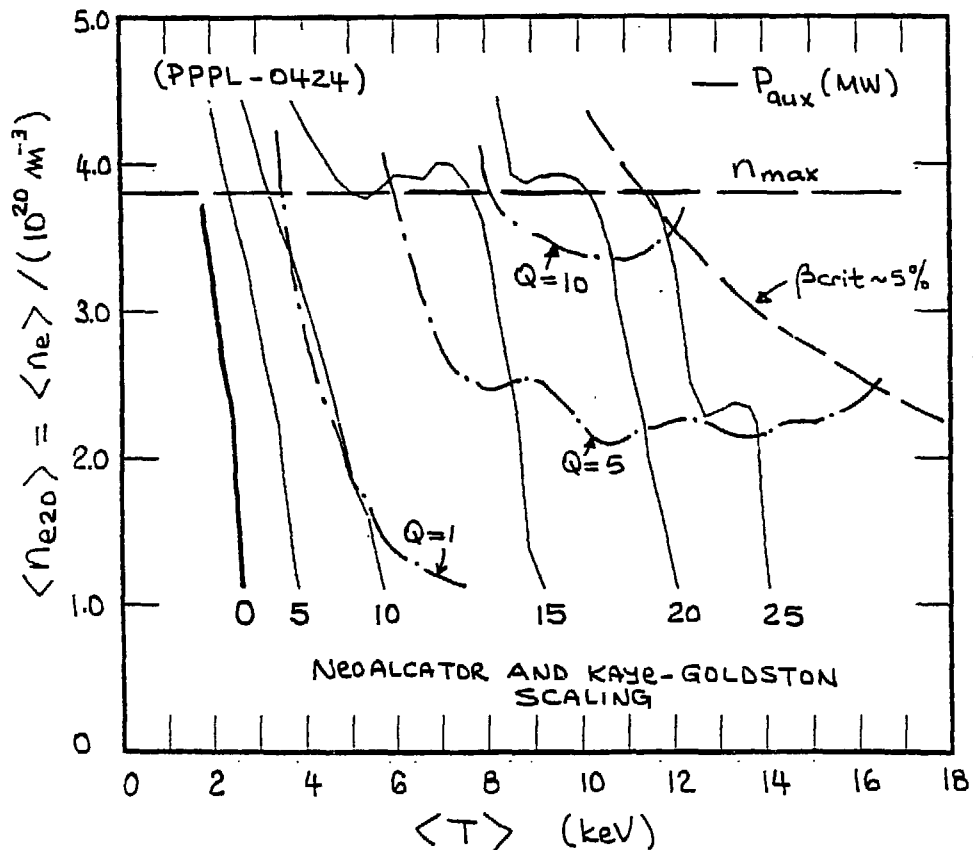


Steady-state POPCON (1-1/2-D) plots showing P_{aux} and Q contours for a typical: $aB_0^2/q_* \approx 20$ (PPPL-0424) device.

Confinement model: "ohmic + auxiliary scaling"

OH - neoAlcator

aux - Kaye-Goldston (L-mode)



Ignition is not accessible in this device with L-mode KG scaling. There is a reasonable $Q > 5$ window ($P_{aux}^{eq} \sim 20 \pm 5$ MW required to access this window).

Example 2 (0-D & 1-1/2-D)

Plasma performance contours for devices with

$$aB_0^2/q_* = 32$$

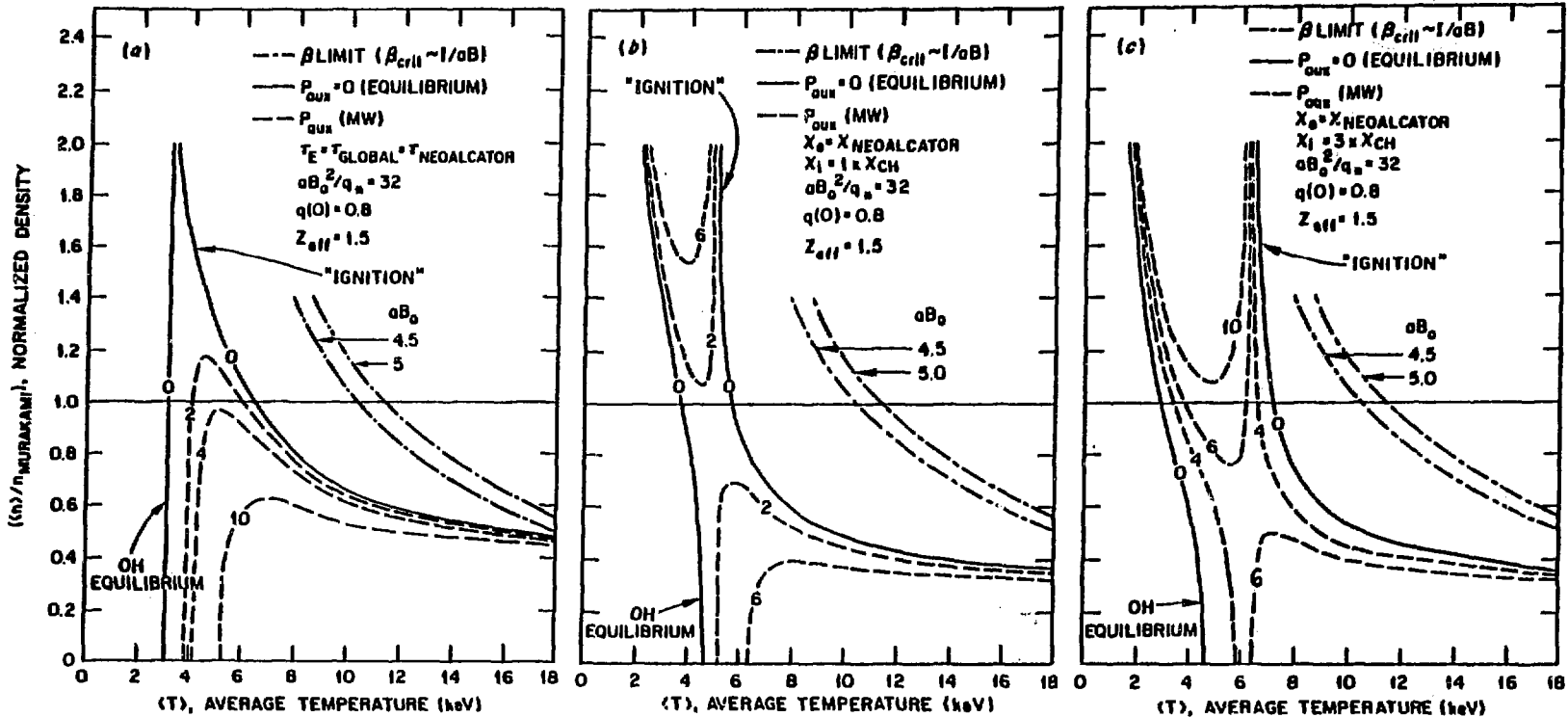
under different scaling assumptions.

Optimal path, P_{aux} , and max. ignition margin are given.

SUMMARY -- Steady-state auxiliary power contours and critical beta contours for devices with $aB_0^2/q_0 = 32$ and $q(0) = 0.8$

Conf. models: (a) $\tau_E = \tau_{\text{neoculator}}$; (b) $\chi_e = \chi_{\text{neoculator}}$ plus $\chi_i = 1 \times \chi_{\text{CH}}$; (c) $\chi_e = \chi_{\text{neoculator}}$ plus $\chi_i = 3 \times \chi_{\text{CH}}$.

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NAU 6/85

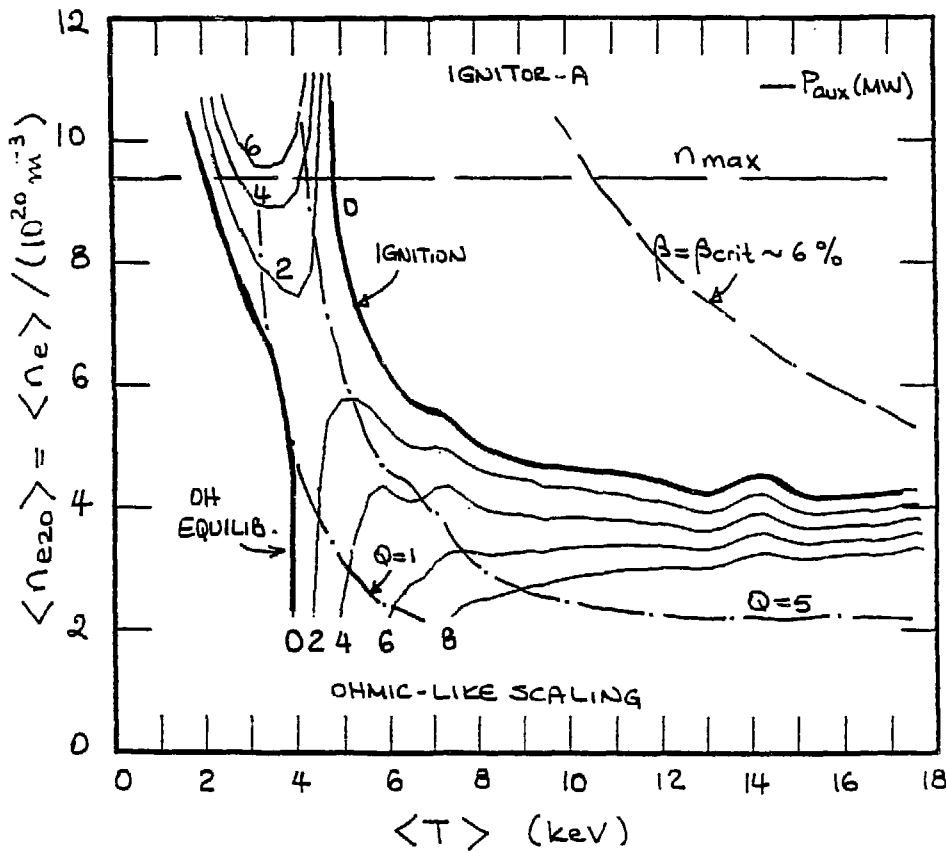
Steady-state POPCON (1-1/2-D) plots

-- P_{aux} and Q contours for Ignitor-A ($a\beta_0^2/q_* \sim 32$)

Confinement model: Ohmic-scaling

$$\chi_e \propto \chi_N$$

$$\chi_i = 0.2\chi_e + 1 \times \chi_{CH}$$

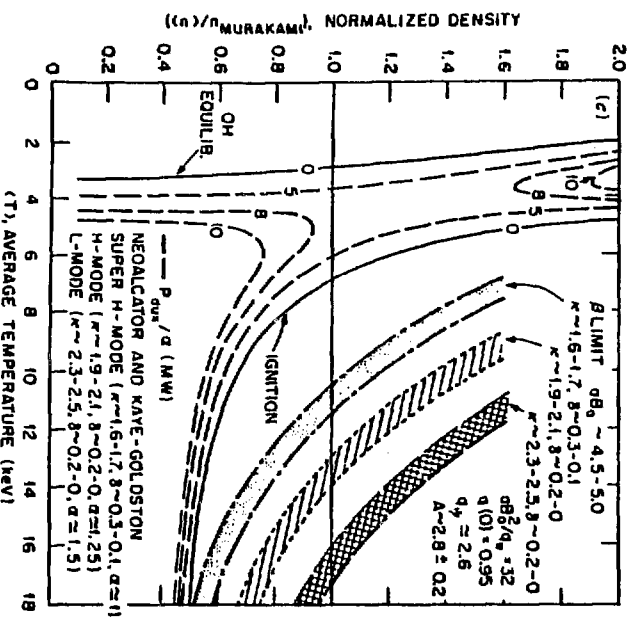
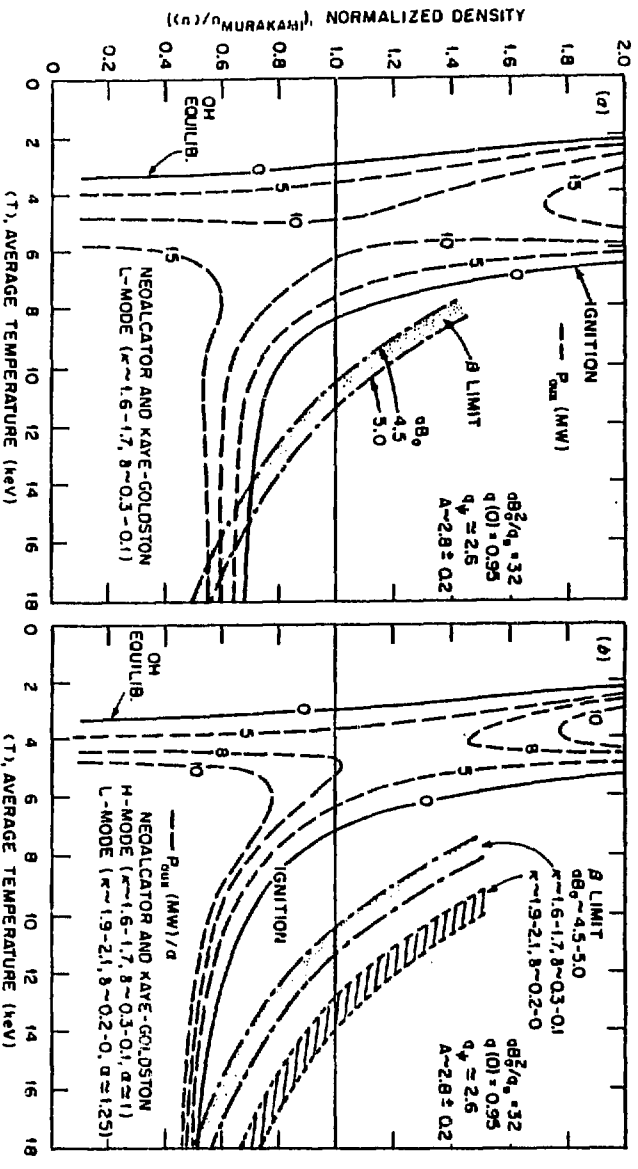


Optimal path $\sim 0.7 n_{max}$

$P_{aux} \sim 1.5 MW$

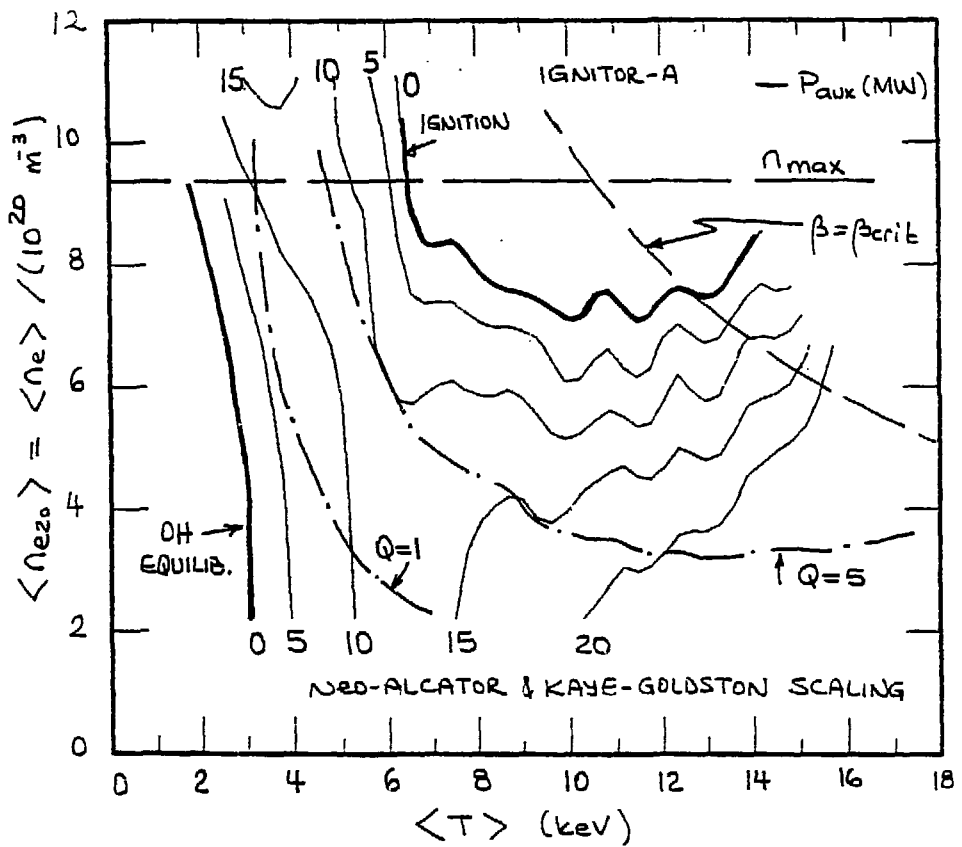
Summary : Neo-Alcator & Kaye - Goldston scaling in
 3 different geometry with increasing κ
 for a class of devices with $\alpha B_0^2/q_0 \sim 32$

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Steady-state POPCON (1-1/2-D) plots
- P_{aux} and Q contours for IGNITOR-A ($aB_0^2/q_* \sim 32$)

Confinement model: ohmic + auxiliary
OH - neo-Alcator
aux - Kaye-Goldston
L-mode



(Fluctuations in P_{aux} contours reflect the effects of discrete sawtooth activity.)

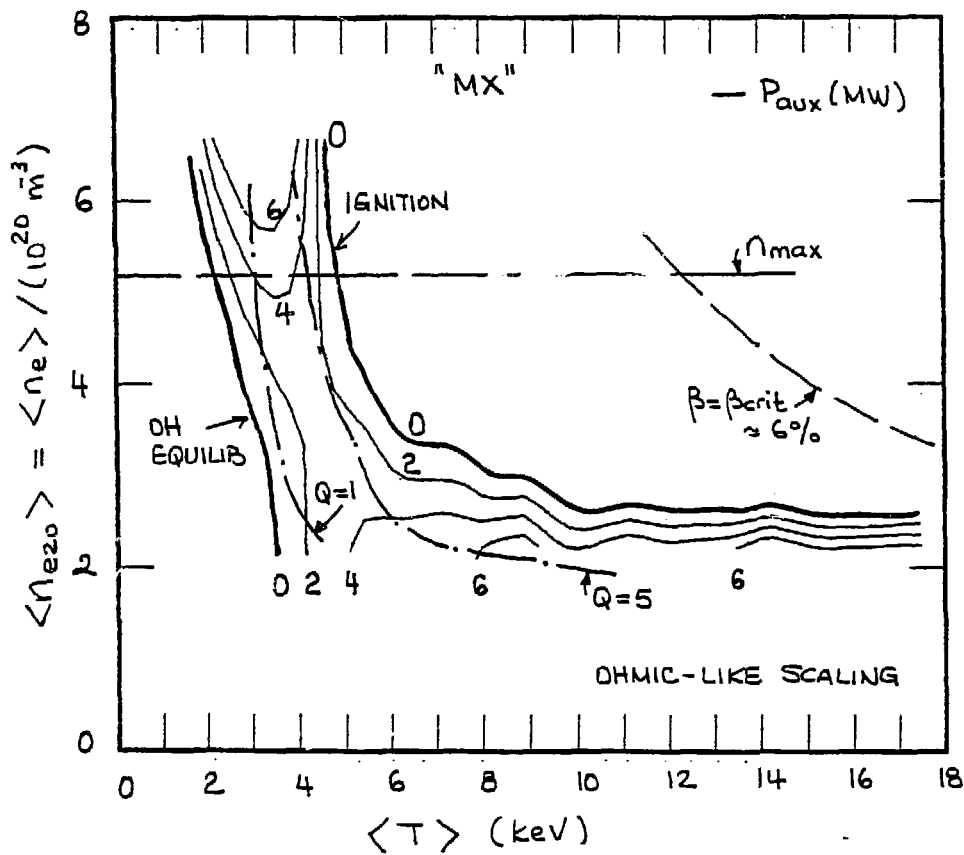
Steady-state POPCON (1-1/2-D) plots

-- P_{aux} and Q contours for "MX"

"MX" -- $aB_0^2/q_* \approx 25$, $\kappa \sim 1.8$, $A = 2.8$

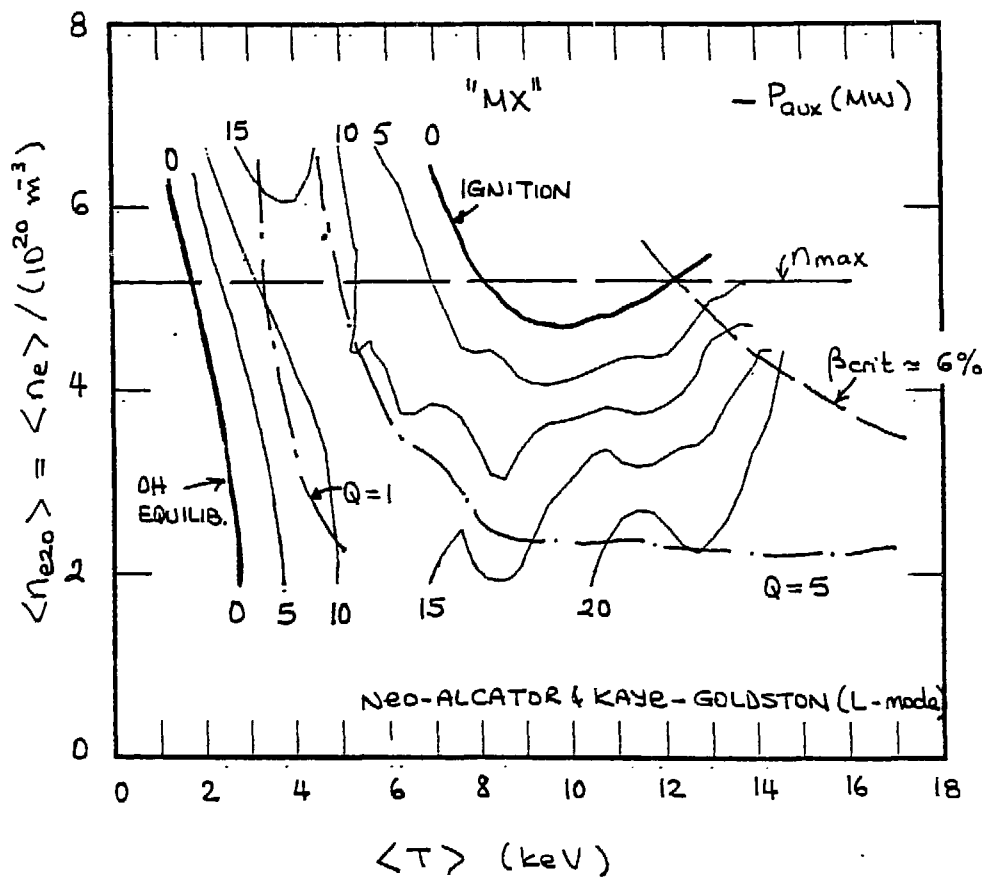
Confinement model: Ohmic scaling

$$\chi_e \sim \chi_{NA} \quad \chi_i \sim 0.2\chi_e + 4\chi_{CH}$$



Steady-state POPCON plots (1-1/2-D)
-- P_{aux} and Q contours for "MX"

Confinement model: ohmic + auxiliary
OH - Neo Alcator
Aux - Kaye-Goldston
L-Mode



OUTLINE

- Physics assessments - design and engineering impact
- 0-D confinement studies
 - physics requirements and options for ignited plasmas
 - classes of devices with equivalent performance
 - sensitivity to variations in confinement models
- 1-1/2-D confinement studies
 - dynamic simulations
 - critical physics issues (sawteeth, α - slowing down, local transport, etc.)
 - startup analysis
 - volt-second consumption