CCNF. 740(1.2.-26

# ANALYSIS OF THE MULTIPHASE INDUCTOR-CONVERTER BRIDGE

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Prepared for

Second International Pulsed Power Conference

Lubbock, Texas

12-14 June 1979

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#### ANALYSIS OF THE MULTIPHASE INDUCTOR-CONVERTER BRIDGE\*

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#### Abstract

Analytical derivations are presented for inductorconverter bridge (ICB) circuits in which energy is transferred from a storage inductor to a load inductor with solid state bridges. These derivations provide complete analytical circuit solution in contrast to previously available numerical (non-analytical) procedures. The analysis is based on two parallel methods: (1) Fourier expansion of the inverter waveforms and (2) a novel method based on the inherent waveforms of the ICB, labeled square functions. Our analytical values of power flow, inductor currents, and voltages compare favorably with the results of a three-phase ICB experiment at Argonne National Laboratory.

## Introduction

The inductor-converter bridge (ICB) is a solid state dc-ac-dc converter system for reversible energy transfer between two inductors. This system is especially suitable for pulsed power supply applications greater than several hundred megawatts and durations from a fraction of  $\epsilon$ second to many seconds. Two such applications are the superconductive equilibrium field coils of the projected tokamak fusion power reactors and superconductive magnets to be used in future particle accelerators.

The ICB system is inherently efficient, controllable in real time and allows isolation of large pulsed reactive loads from the power grid. Thus, only the average system losses are drawn from the grid.

#### Operation of the ICB Circuit

Detailed operation of the ICB may be found in references 1 and 2. Figure 1 shows a three-phase ICB where the storage and the load coil is represented by L<sub>S</sub> and L<sub>1</sub>, respectively. At a typical instant during the energy transfer, dc currents is and i, will be flowing in the storage and load coils, respectively. The SCR's of the left hand side (storage side) are fired in the normal Graetz bridge sequence:  $S_{L1}^{}$   $S_{L5}^{}$ ,  $S_{L1}^{}$   $S_{L6}^{}$ , S<sub>L2</sub> S<sub>L6</sub>, S<sub>L2</sub> S<sub>L4</sub>, S<sub>L3</sub> S<sub>L4</sub>, S<sub>L3</sub> S<sub>L5</sub>, S<sub>L1</sub> S<sub>L5</sub>, .... The SCR's of the right-hand side (load side) follow the same sequence but may be out of step with respect to the storage side. The direction and the level of power flow is determined by the relative timing between the source and the load bridge switching sequences such that a load bridge lead will cause power flow into the load and vice versa. The Y-connected capacitors in the middle serve as the intermediate energy store between the storage and load coils and they provide the reverse voltages to commutate the inductor circuits from one SCR to the next. Thus, no external counterpulse circuit is needed.

Only a very small fraction of each coil energy is extracted in each bridge cycle. Therefore, by varying the relative timing between the source and load switching sequences and/or the frequency of operation, very fine control over the rate of energy transfer can be achieved. The functional dependence of the power<sup>9</sup> on the relative timing

<sup>\*</sup>Work supported by the U.S. Department of Energy. †Also at the University of Wisconsin.

(relative phase difference, and the frequency will be illustrated in the following sections.

Circuit Analysis Based on Fourier Components

Figure 2 shows an m-phase ICB. For this analysis, the SCR's have been replaced by ideal switches. This idealization implies a lossless system and operation within the successful commutation bounds. Since each bridge cycle changes the coll energies by a very small amount, the inductor currents are nearly constant in one cycle. Viewing such inductor-converter half from the mphase ac lines, one will see an m-phase square wave current source system, each phase being  $\frac{2\pi}{m}$ radians displaced from the next and the wave amplitude being equal to the instantaneous coll current, Fig. 3. Since the average power is divided equally between the phases, it may be calculated from a one-phase diagram, Fig. 4.

The calculated instantaneous average power from the left-hand source (storage) to the right-hand source (load) is

$$e_{\text{PSa}}(\mathbf{t}) = \sum_{n} \frac{1}{2n/c} (a_{n\text{Sa}} b_{n\text{La}} - b_{n\text{Sa}} a_{n\text{La}})$$

where the a's and b's are the Fourier coefficients of the storage and load source waveforms as indicated by the subscripts. For illustration, the Fourier coefficients of the symmetrical waveform of Fig. 5 which is for when m is odd will be used.

$$\begin{cases} I_{Sa}(\cdot) = \sum_{n} \frac{2^{1} S}{m} \left[ 1 - (-1)^{n} \right] \sin \frac{n}{m} \cos n \, \text{ut} \\ I_{La}(t) = \sum_{n} \frac{2^{1} L}{m\pi} \left[ 1 - (-1)^{n} \right] \sin \frac{n\pi}{m} \cos n \, (\text{ut+i}) \end{cases}$$

where  $\alpha$  is the angle by which the load bridge leads the storage bridge. After substitution, the average power per phase will simplify to

$$< p_{Sa}(t) > = \sum_{n} \frac{2^{I} S^{I} L}{n^{3} \pi^{2} \omega c} \left[ 1 - (-1)^{n} \right]^{2} (\sin \frac{n \pi}{m})^{2} \sin n \omega$$

The total instantaneous average power delivered from the storage coil is m times the above, or

The effect of the relative phase, 1, and the bridge switching frequency, 1, on the power flow is evident in the above relationship. Note also that the contribution of higher harmonics is attenuated by the  $\frac{1}{n^3}$  term. Figure 6 represents the plot of the instantaneous average power as a function of  $\alpha$  with the number of phases, m, as the parameter. For m = 3, over 29.5% of the met power is delivered by the fundamental frequency, making the m = 3 curve in Fig. 6 almost purely sinusoidal.

The time functions of average coll currents, power, and voltages may now be calculated. For simplicity, let

$$-k = \sum_{n=n} \frac{2r}{\frac{2r}{r^2}} \left[ 1 - (r)^n \right]^2 \left( \sin \left(\frac{n}{r^2}\right)^2 \sin n \right),$$

Then from

$$\langle \mathbf{i}_{S}(\mathbf{t}) \rangle = -\frac{d^{-1}S(t)}{dt} = \frac{d^{-1}L(t)}{dt}$$

$$\begin{cases} \mathbf{k} \langle \mathbf{i}_{S} \rangle \langle \mathbf{i}_{L} \rangle = -\frac{d}{dt} \left(\frac{1}{2} L_{S} \langle \mathbf{i}_{S} \rangle^{2}\right) \\ \mathbf{k} \langle \mathbf{i}_{S} \rangle \langle \mathbf{i}_{L} \rangle = \frac{d}{dt} \left(\frac{1}{2} L_{L} \langle \mathbf{i}_{L} \rangle^{2}\right) \end{cases}$$

Solving these equations with the initial conditions

$$\begin{cases} {}^{ci}S^{>} \mid_{t=o} = I_{o} = \text{ initial storage current} \\ {}^{o}_{} \mid_{t=o} = o \end{cases},$$

we obtain

$$\begin{cases} I_{S} > (t) = I_{o} \cos \frac{K}{\sqrt{L_{S}L_{L}}} t \\ I_{L} > (t) = I_{o} \sqrt{\frac{L_{S}}{L_{L}}} \sin \frac{K}{\sqrt{L_{S}L_{L}}} t. \end{cases}$$

The average power will become

$$e_{P_{S}} > (t) = 1/2 \text{ K} l_{0}^{2} \sqrt{\frac{L_{S}}{L_{L}}} \sin \frac{2\kappa}{L_{S}L_{L}} t$$

and the average coil voltages from

$$v_{V} = L \frac{d \cdot i}{dt},$$

$$\begin{cases}
v_{S} \cdot (t) = K I_{o} \sqrt{\frac{L_{S}}{L_{L}}} \sin \frac{k}{\sqrt{L_{S}L_{L}}} t \\
v_{L'}(t) = K I_{o} \cos \frac{k}{\sqrt{L_{S}L_{L}}} t
\end{cases}$$

# Circuit Analysis Based on Square Function Calculations

The preceding Fourier method shows the influence of harmonics in the system. However, for real time control, solving the equations of the form

$$=\sum_{n} f_{n}(\alpha)$$

for a is time consuming and possibly inaccurate. The following technique will provide closed form equations that are efficiently solved by a microcomputer in real time control.

The calculations are based on specially tailored functions symbolized by Sq (X) and Tr (X), Fig. 7. These waveforms are inherent in the operation of the ICB circuits. The Sq function is a good mathematical representation of the ac phase currents of the system. The Tr function is the integral form of Sq and is a good representation of the capacitor voltages due to the phase currents. A brief mathematical development of the Sq and Tr functions is presented in the Appendix. Figure 8 shows how a phase current of the equivalent three-phase circuit may be decomposed into two Sq functions. Thus, the currents  $i_{Sa}$  and  $i_{La}$  of Fig. 4 may be written as

$$\begin{cases} \mathbf{i}_{Sa}(t) = \frac{\mathbf{I}_{S}}{2} \left[ Sq \left( \tau + \frac{T}{6} - \tau_{o} \right) + Sq \left( \tau - \tau_{o} \right) \right] \\ \mathbf{i}_{La}(t) = \frac{\mathbf{I}_{L}}{2} \left[ Sq \left( \tau + \frac{T}{6} \right) + Sq \left( \tau \right) \right] , \\ \mathbf{o} \le t \ge T, \mathbf{o} \le \tau \le T, \mathbf{o} \le \tau \ge 1/2 \end{cases}$$

where  $\tau_{\rm O}$  is the relative lag time of the storage bridge.

These current representations for one period are adequate for instantaneous average power calculations. The net power out of storage per phase, results from the interaction of  $i_{Sa}$  and  $v_{La}$ , the capacitor voltage due to the load phase current:

$$v_{La}(t) = \frac{1}{C} \int_{0}^{t} i_{La} d\tau = \frac{1}{2C} \left[ Tr (\tau + \frac{T}{6}) + Tr (t) \right].$$

This power is

$$p_{Sa}(t) = i_{Sa} v_{La} = \frac{I_{S}I_{L}}{4C} \left[ Sq (\tau + \frac{T}{6} - \tau_{o}) \cdot Tr (\tau + \frac{T}{6}) + Sq (\tau + \frac{T}{6} - \tau_{o}) \cdot Tr (\tau) + Sq (\tau - \tau_{o}) + Tr (\tau + \frac{T}{6}) + Sq (\tau - \tau_{o}) \cdot Tr (\tau) \right].$$

The instantaneous average power per phase is

$$\langle \mathbf{p}_{Sa} \rangle = \frac{1}{T} \int_{0}^{T} \mathbf{i}_{Sa} \mathbf{v}_{La} d\mathbf{t} = \frac{\mathbf{I}_{L} \mathbf{I}_{S}}{4TC} \int_{0}^{T} [Sq (\tau + \frac{T}{6} - \tau_{o}) \\ \cdot \mathbf{Tr} (\tau + \frac{T}{6}) + Sq (\tau + \frac{T}{6} - \tau_{o}) \cdot \mathbf{Tr} (\tau) + \\ Sq (\tau - \tau_{o}) \cdot \mathbf{Tr} (\tau + \frac{1}{6}) + Sq (\tau - \tau_{o}) \\ \cdot \mathbf{Tr} (\tau) ] d\tau.$$

The four terms in the integrand are first transformed to the normalized variable functions shown in the Appendix, then each term is evaluated by using the proper integral identity in the Appendix. The result is then multiplied by three for the total three-phase power

$$\langle \mathbf{p}_{Sa} \rangle = \begin{cases} \frac{T \cdot \mathbf{I}_{S} \mathbf{I}_{L}}{C} (2 \cdot \mathbf{y}_{0} - 3 \cdot \mathbf{y}_{0}^{2}), & 0 \le \mathbf{y} \le \frac{1}{6} \\ \frac{T \cdot \mathbf{I}_{S} \mathbf{I}_{L}}{C} (3 \cdot \mathbf{y}_{0} - 6 \cdot \mathbf{y}_{0}^{2} - \frac{1}{12}), & \frac{1}{6} \le \mathbf{y} \le \frac{1}{3} \\ \frac{T \cdot \mathbf{I}_{S} \mathbf{I}_{L}}{C} (\mathbf{y}_{0} - 3 \cdot \mathbf{y}_{0}^{2} + \frac{1}{4}), & \frac{1}{3} \le \mathbf{y} \le \frac{1}{2} \end{cases}$$

where  $\gamma_0 = \frac{\tau_0}{T} = \frac{2}{360^\circ}$  is the normalized lag time of the storage bridge relative to the load bridge and the frequency is represented by the period T. This is the closed form of the 'p' vs a curve for m = 3 in Fig. 6. The time functions of the average coil currents, power and voltages may be calculated as before. The only difference being that K ( $\gamma_0$ ) is in closed form.

Expressing the actual circuit waveforms analytically, allows other useful calculations such as the actual capacitor voltages throughout the transfer cycle, the study of commutation throughout the cycle, and the actual coil voltages and instantaneous currents, without resorting to humerical procedures.

### Comparison with Test Results

A model three-phase ICB has been built and tested at Argonne National Laboratory. This system uses two identical superconducting coils capable of storing 125 kJ at 250 f, as the storage and load inductors. Other system parameters are:

$$L_{S} = L_{L} = 4 \text{ H}, I_{o} = 100 \text{ A}$$
  
 $C = 10^{-4} \text{ F}, \quad \alpha = 90^{\circ}$   
 $\omega = 4084 \text{ rad/s}$ 

The system equations are derived from substitution of these parameters and m = 3 into the time functions shown in the Fourier analysis section:

$(1_{S})$ (t) = 100 cos 0.5609 t	A
<i<sub>L&gt; (t) = 100 sin 0.5609 t</i<sub>	A
<v<sub>S&gt; (t) = 224.4 sin 0.5609 t</v<sub>	v

<v_l></v_l>	(t)	-	224.4	C 0 5	0.5609	t	v
Ps'	(t)	-	11218	sín	1.1218	t	¥

A plot of these equations appears in Fig. 9.

Figure 10 shows the average coil voltages and currents obtained experimentally. Good agreement exists between the analytical and the experimental results. Note, however, that the experimental initial storage current is somewhat higher that the final load current. This is due to the losses in the system which is neglected in this analysis, but may be incorporated in the differential equations leading to the average time functions.

#### Conclusions

The behavior of the multiphase inductor-converter bridges have been studied by two analytical techniques. The conventional Fourier technique produces the average circuit power, currents, and voltages as a function of time. It also shows the effect of the existing harmonics in the circuit behavior. The square function technique is particularly devised for the ICB and other S + circuits in which rectangular waveforms appear. The identities defined on the special functions Sq (X) and Tr (X) operate directly on the circult waveforms. Thus, much more information about instantaneous behavior of the circuit is available for analysis, system design, and the development of real time control algorithms. Preliminary tests with open loop microcomputer control have been conducted with satisfactory results. The development of an optimal closed loop control algorithm is currently in progress.

#### Appendix

The Sq function is defined as the sum of unit ste; functions as follows,

Sq 
$$(\gamma + \gamma_0) \stackrel{\ell}{=} u(\gamma) - 2 u(\gamma - \frac{1}{2} + \gamma_0) + 2$$
  
 $u(\gamma - 1 + \gamma_0) - u(\gamma_0 - 1),$   
 $0 \le \gamma_0 \le \frac{1}{2}, \ 0 < \gamma \le 1.$ 

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The Tr function is defined as the integral value of the Sq functions,

$$Tr (\gamma + \gamma_0) \stackrel{i}{=} \int_0^\infty Sq (\xi + \gamma_0) d\xi$$

$$= -i u (\gamma) - 2 (i - \frac{1}{2} + \gamma_0)$$

$$u (i - \frac{1}{2} + \gamma_0) + 2 (i - 1 + \gamma_0)$$

$$u (\gamma - 1 + \gamma_0) - (\gamma - 1) u (\gamma - 1)$$

The following integral is extensively used in calculating the average power flow in the circuits of interest. Therefore, it will be stated as an identity which may be directly verified,

$$\int_{-0}^{-1} \{ Sq_{-}(x_{1} + x_{1}) \} + \begin{cases} -2x_{1}^{2} - x_{1} + 4x_{1}x_{2} + x_{2} \\ -2x_{2}^{2}, x_{1} + x_{2} \\ 2x_{2}^{2} - x_{1} + 4x_{1}x_{2} + x_{2} \\ 2x_{2}^{2} - x_{1} - 4x_{1}x_{2} + x_{2} \\ +2x_{2}^{2}, x_{1} + 2x_{2} \end{cases}$$

$$0 \ge x \le \frac{1}{2}, \quad 0 \le x_{2} \le -\frac{1}{2}.$$

The following identities will also be of considerable value

$$Sq (\gamma - \gamma_{0}) = -Sq [\gamma + (1/2 - \gamma_{0})],$$
  
Tr (\gamma - \gamma\_{0}) = -Tr [\gamma + (1/2 - \gamma\_{0})], 0 - \gamma\_{0} - \frac{1}{2}

which may be verified by direct substitution. Note that by using these identities, we can easily evaluate the above integral for any combination of leading and lagging functions.

The above Sq and Tr functions have been defined for a period equal to 1. For periods other than I, the argument should be multiplied by the appropriate constant:

Sq (t + t<sub>0</sub>) = Sq (Y + Y<sub>0</sub>), 
$$\begin{cases} 0 \le t \le T, \ 0 \le t_0 \le \frac{T}{2} \\ 0 \le Y \le 1, \ 0 \le Y_0 \le \frac{1}{2} \end{cases}$$

$$\begin{aligned} & \mathrm{Tr} \ (t + t_{o}) = \int_{0}^{t} \ \mathrm{Sq} \ (c + t_{o}) \ \tilde{u}^{*}, \begin{cases} 0 \leq t \leq T \\ 0 \leq t_{o} \leq -\frac{T}{2} \end{cases}, \\ & \text{where } t = \frac{t}{T} \text{ and } t_{o} = -\frac{to}{T} \end{aligned}$$

Substitution of  $\cdot$  and  $\cdot_{\rm O}$  in the above integral will give

$$Tr (t+t_{0}) = T \int_{0}^{Y} Sq (i + v_{0}) dt$$

$$= T [Tr (v + v_{0})] = t u (t) - 2$$

$$(t - \frac{T}{2} + t_{0}) U (t - \frac{T}{2} + t_{0})$$

$$+ 2 (t - T + t_{0}) u (t - T + t_{0})$$

$$- (t - T) u (t - T).$$

For time base calculations, these functions may be used,

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Fig. 3. Equivalent Diagram of an m-phase ICB.



Fig. 4. One Phase Equivalent Circuit of an m-phase ICB.



Fig. 5. Single Phase Current Source Waveform for an m-phase ICB Circuit.



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Fig. 6. Plot of mpr vs



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