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# THE REPRESENTATION OF ORIENTATION DISTRIBUTIONS $\dagger$ 

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#### Abstract

1. Mirmoduction

It is widely acknowledged that texture is the prime cause of anisotropy in polycrystalline metals: the nonrandom distribution of the crystallographic orientations of the grains ('preferred orientation', 'texture', or 'fabric') transfers some of the anisotropic properties of single crystals to the aggregate. Whereas some properties show little rr no anisotropy even in single crystals, many properties are strongly anisotropic even in materials with a cubic lattice.

A nonrandomnese of the orientation diatribution is virtially everpresent in metals, because all processes involved in producing such materials (casting, deformation, recrystallization) are locally orientation dependent. Texture studies are, in fact, frequently used by metallurgista to help identify the crystallographic details of ach processes.

Despite this general recognition of the importance of texture for a description of macroscopic properties, it seems that quantitative evaluations of texture are rarely used in enginearing practice or even in academic physical metallurgy, outside amall community of texture experts. In general applications, one or two pole figures are given at bect, or some idealized orientations ('texture components'), with m qualitative interpretation of the expected effects. This is so even though sophiaticated quantitative descriptions of three-dimensional orientation distributiong have been avallable for twenty yeari [1-3].

Why has there been such inertia in using quantitative texture descriptions as a general tool of deformation studies? We submit that this is primarily due to some unfortunate choices that were made by the pionesrs of this development, with respect to the graphic repxesentation of orientation dietributions. of course, any represantation may appear


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easy once it has become familier; but some provide a significant
activation barrier upon firat contact.
    A number of alternative representations have been suggested in the
literature but have, for some reason, not becone common. We find a
particular combination of these aspecially easy to visualize for the
uninitiated, easy to assess qualitatively, and easy to evaluate
quantitatively. These plots use polar rather than Cartesian coordinates
(as proposed by Williams [3] and by Pospiech and Lucke [4,5]), and
\bulletqual-area projections (as they are commonly used in geology [6]). Wr
will deacribe these reprasentations in the present paper, and also review
In tutorial form some of the basic concepts of crientation distributlons
in uniform terminology.
    In addition to graphic representations of orientation distributions,
algebraic onea were introduced right at the beginning of this development;
particularly, the definition of a continuovs orientstion distributaon
function (OUF), arid ita expanaion in generalized apherical harmonics.
Th1s provides an elegant and concise deacription for some applications
(such as tensor properties), but becomes cumbersome when very many terms
are needed. In ouch ceses, a useful graphic representation of the data is
casential for a quantatative analyais.
    A complete deacription of orientation diatributions zequirea a three=
-dimensional orientation epare. On che other hand, two-dimenilonal pro-
jectiong of this epace ie what la measured. There are varioue ways of
Inferring a 3-dimenional digtributicn from a number of 2-dimenaional pro-
jections, and it has been a topic of inteneive recenc debate to what
extent, under what circimetances, thia can be done unambiguoualy {7-11).
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This current interest provides a further reason for us to clarify the descriptions. For the purpose of this paper, we will assume that the three-dimensional orientation distribution is known, by theoretical prediction or by some deconvolution of experimental measurements; we only concern ourselves with the representation of known distribution.

We will illustrate the principles presented with a particular experimental texture: from the surface layer of a copper polycrystal cold-rolled to 60\% reduction in thickness. Four incomplete pole figures (200, 220, 222 , and 113) were determined by x-ray diffraction in reflection geometry. The measured pole figures nearly exhibited orthorhombic symmetry (as expected), which was then atrictly enforced by averaging the four quadrants of the pole figure. The orientation distribution function was obtained using the expansion in spherical harmonics (with only even-order coafficients up to $\ell=18$ ). Inasmuch as this is oaly meant to serve as an example, neither the detailed sample history nor the pole figure deconvolution procedure are of esselice.

## 2. DIRECTION SPACE AND ORIENTATION SPACR

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    While the description of (two-dimensional) directions is relatively
trivial, we will review it in detail, because we find it important to pre-
serve some continuity betwean the two- und three-dimenaional cases -- more
contlnuity than has generally bean made use of.
In some cases, only the distribution of aingle diraction is of
interest: e.g, the distribution of c-axes in n haxagonal materiai with
respect to the normal of sheet; or the distribution of the tenaile axes
In the atereographic triangle characterizing each grain of a cubic roly-
cryatal. Such dietributions are easy to treat quantitatively: the pafnt
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density on the surface of a unit sphere characterizes the distribution of directions in a uniform way. It may be represented by various projections, and it may be simplified by the application of symmetry principles-which we shall summarize in Section 3.

A particular symmetry deserves mention: namely when the aign of the direction is of no concern. Unsigned directions are often called 'poles' or 'axes'; we shall use pole for the unsigned normal to a crystallographic plane, as is common usage; and we shall use axis specifisally for an unsigned sample cooxdinate. The distribution of either use with respect to some reference system can be dsscribed by point densities on the surface of hemisphere. (Signed directions may be called 'unit vectors'; they are represented on the surface of whole sphere.)

Whereas the term 'orientation' is sometimes loosely used in the same sense as 'direction' was used above (such as in 'the orientation of the tensile axis'), we wish to reserve the term orientarion for the relation between an entire coordinate system (triad) and some reference triad. This requires the specificetion of three (perhaps perpendicular) axes or, more commonly, of two unit vectors and a handedness.

It tusk long time for an appropriate description of orientation. distributions to be developed [1,2]: 1.e., an orientation apace in which a uniform point density corresponds to aniform (e.g. random) distribution of orientations, analogous to the surface of sphere in the case of directions. We will call such epeces 'homochoric'. One easy solution is as followo (Fig. 1): deacribe the firet direction (usualiy called $Z$ or $\times 3$ or [001]) Just like any direction, and the second (perpendicular) direction (usually called $X$ ) by en azimuth around the firat. A useful image Ls that of boat at location 2 on the surface of the aerth, which
is heading in direction $X$. Specification of the first direction requires two numbers (the ship's longitude and latitude), that of the second direction requires one more number (the ship's hesding, or azimuth). There is some arbitrariness in the precise specification of these three angles and their zeroes, but Euler angles are most commonly used and quite adequate. We shall apply the term Euler space to any three-dimensional space in which the coordinates are determined in some way by these three angles [5]; it is the precise structure of this space that we shall discuss, with the aim of making it homochoric.

There is an aesthetic flaw in this Euler space description of orientations: namely, that one direction (the 'first' direction used above) must be preferred by the observer, even if no such preference is inherent in the physics. For example, the roll: plane normal is commonly preferred ovar the rolling direction, although both have equivalent statur. This artificiailty can only be alleviated by prefering various poles or axes in succession, perhaps until one is found that is most easily visualizod --or by showing two or three of these rerresentations in parallel.

Two other descriptions of orientation relations are someimes used, for specific reasons [5]. One is in cerms of a paricular axis in space (or a number of symetrically equivalent ones) around which a single rotation brings the two triads inte coincidence. This has heen useful, e.g., In mechanistic discussions of racrystallization [12]; it is quite different from the description we use (in terms of a direction in one of the coordinate syems ard animuth around it), which is completely equivalent to the three-rotations cheme of introducing Euler angles [13].

Fiually, the orientation of one triad with respect to another can be describad by a rication matrix $[5,14]$. This deacription is casiest in

# computer codes. It does not prefer any one axis; but the three numbers necessary for a complete specification do not form a homochoric space. 

## 3. REPRESEITATION OF DIRECTIONS

### 3.1 Equal-area and stereographic projections

The location of a point on a sphere is easily described in the familiar geographic terminology of longitude and latitude. Following Bunge [13], we will use the longitude* $\beta$ and the co-latitude $\alpha$ (also called pole distance).

A point on the sphere is then projected onto a plane by means of some standard projection. Stereographic and equal-area projections are most widely used. Projections always distorta true representation. In the stereographic projection (almost exclusively used in metallurgy) equal great-circle segments have the same length wherever they appear on the sphere. In equal-area projestion (commonly used in geophysics) equal areas on the aphere have the same size in projection. This is best illustrated by projections of the coordinate system, which provides a Wulff net (Fig. 2a) and a Schmidt net (Fig. 2b) respectively. Either net is casily used for graphic constructions and the determination of angular relationships. While stereographic projection is most appropilate when the angles between crystal directions are of prime concern, equal-area projection would seem more appropriate when population densities are to be descrited. (An auvatage of the stereographic projection is that circles remain circles--but their centers do not remain their centers.)

[^1]Figure 3 shows a \{100\} polc figure of a typical copper texture, both in stereographic (Fig. 3a) and equal-area projection (Fig. 3b). Pole densi:ies are experimentally averaged into a continuous distribution and contours are expressed in multiples of a random distribution ('m.r.d."). Note that the ( 100 ) pole figure of a cubic matarial displays all three \{100\} orientations for each grain--but it does not give direct information on which $\{100\}$ is in fact the specific (001) pole for a specific (100), for one grain. It is for this reason that pole figures are usually insufficient for a quantitative description of textures; hut they are the most quantitative experimental information available.

### 3.2 Pole figures and 'inverse pole figures'

In the last section, the sampla axes were chosen as a reference system, and the crystallographic poles were described in this frame. A completely equivalent description is its inverse (or dual) that of sample axes in terms of crystal axes. Sometimes one is more appropriate, sometimes the other. For example, when the sample is a wire, there is no natural sample triad, and all pole figures degenerate into circles; but a description of the wire axis distribution in terms of the crystal axes -an 'inverse pole figure'-- is illustrative. Conversely, when the most appropriate sample axes (e.g., the most symmet ic ones) are not known sufficiently accurately beforehand, a pole figure provides information on the actual sample symmetry, and thus guidance in selecting 'sample' axes. For cubic materials, which we emphasize in the present paper, the three <100> directions provide a ready reference frame for inverse pole figurea. Cubic symetry reduces the area on the sphere that is neressary for a conplete description by factor 24 : only 2 of the 48 unit triangles are needed, even if the plotted acis has no symmetry; if it lies in a mirror plane of the sample, single triangle suffices [15].

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Figure \(4 a\) shows one quadrant of the unit hemisphere, as it is usually projected: with the preferred pole ( 001 ) at the center. This makes the center a special poirt. One effect is that the mesh units have very small area there; if the constant- \(\Phi\) lines were drawn such as tc make each mesh unit encompass equal area, the units near the center would attain a very anisotropic shape. The unit triangles under cubic symmetry are also shown: any one of the three triangle-pairs (labelled \(I, ~ I I a / b\), and II) suffices for the most general csse; but the special nature of the origin would seem to make pair I an especially poor choice.
Figure 4b shows a different scheme: the preferred pole (001) has been moved to the periphery [3]. In this case, it is possible to have similar and equi-axed unit areas everywhere. (Fig. 4 b is an equal-area projection, see Fig. 2b.) The preferred triangles are shown in what appears to us to be the most convenient way (and rather conventional). For a numerical description, it would seem easier to use the latitude \(\lambda\) (rather than the co-latitude \(\alpha^{\prime}\) ), and a longitude \(\mu\) (for meridian) counted from (100).
Figure 5 demonstrates the case of the particular copper sample whose measured pole figure was shown in Fig. 3. Use is now made of the expected orthorhombic symmetry of the rolled sample (which was essentially verified in Fig. 3), to plot only one quadrant of the \(\{100\}\) pole figure. Also, this is complemented by a \(\{111\}\) pole figure (also measured and folded into one quadrant). Finally, two inverse pole figures, for the \(X^{s}\) and \(Z^{s}\) axes (rolling direction \(R D\) and rolling plane normal \(N D, ~ r e s p e c t y e l y)\), were derived from these pole figures and are plottedin Fig. 5.
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## 4. REPRESETTATION OF ORIETTATIOAS

### 4.1 Frevious representations

In the last figure (Fig. 5c, d), we showed in parallel two inverse pole figures, one for $X^{8}=R D$ and one for $Z^{S}=N D$. This gives almost all information about the three-dimensionsl relation between the axis aystem $X^{s}, Y^{s}, Z^{s}$ and the crystallographic reference system $X^{c}=$ $\langle 100\rangle, Y^{c}=\langle 010\rangle, Z^{c}=\langle 001\rangle$, and such two figures together are ir. fact of ten sufficient to characterize an orientation distribution. What is missing is information on the correlation between che $X^{3}$-axis cf each particular grain and its $2^{s}-a x i s$.

It is for this reason that a truly three-ámensional representation is necessary in goneral. One easy way to introduce it is to plot, for example, an inverse pole figure for $\mathbb{Z}^{8}$ and perpendicular to it the amount of rotation around this preferred direction [3,4]. Another good methed is to attach ticks to the points in an inverse pole figure to indicate the rotation around this point [16] much like in the "boats on the earth" picture introduced above; the disadvantage of the latter is that it works only for discrete points, not for continuous distributions. Neither of these procedures is easily applicable to pole figures, if there is more than one equivalent pole. Nevertheless, the basic idea, we feel, is of compelifig aimplicity: to represent an orientation by one direction (that of the arbitrarily preferred axis) and a rotation around it; this (last) rotation we shall call an azimuth (and count it from the equator). This description makes open use of the need to prefer one direction; and it retains continuity with the description of directions when only directions are important. We shall elaborate on two specific representations besed on this principle below.

In texture research, another description, introduced by Bunge [1] and Roe [2], has become common, which is entirely equivalent, only differently phrased and differently represented. Here, three successive rotations are performed around the coordinate axes of the reference system (in a certain sequence), by the three Euler angles $\phi_{1}, \Phi, \phi_{2}$ (Fig. 6). * In the sample reference frame, the first two rotations are exactly equivalent to the angles $\beta$ and $\alpha$ used before: they do describe a cirection, and $\phi_{2}$ neasures an azimuth around this direction --just as in the description used above. However, $\left\{\Phi_{1}, \Phi, \Phi_{2}\right\}$ are usually treated as three parameters on equal footing, and plotted in a three-dimensional space. The problem is that this space was chosen to be Cartesian [17]. This is equivaient to projecting the hemisphere on which the directions ( $\phi_{1}, \Phi$ ) are uniformly distributed onto a square: the entire ajvantage of using a homochoric space (the Euler space) has been lost in its repesentation. An adrantage is that the periodicity in all three rotations can be seen: the Cartesian Euler space can in fact be represented by a space lattice with orthorhombic symmetry (but no inversion center) $[18,19]$.

The graphic representation of orientation distributions in this Cartesian space of Euler angles can be shown by contour lines in a series of two-dimenional sections, and this diagram has become so common as to be virtually identified with the abbreviation ODF. Figure 7 shows such a plot, for the same copper sample illustrated before. Note that the representation is severely distorted: the single orientation, in each section,

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    We are staying as close as possible to Bunge's conventions [13]. Most
type fonts contain only the straight or only the curly version of lower-
case phi: they should be treated as completely equivalent. All three rot-
ations increase counterclockwise (looking down on the positive axis) in
the sample frame, clockwioe in the crystal frame. In the sample frame, $l
measures the angle from - }\mp@subsup{Y}{}{8}\mathrm{ to + + 'c, 利 from the equator to + }\mp@subsup{X}{}{c}\mathrm{ . In
the crysial frame, }\mp@subsup{\phi}{2}{}\mathrm{ counts from +Yc to +Zs}\mathrm{ and }\mp@subsup{\phi}{1}{}\mathrm{ from the equator
to +Xs. (The asymmetry in thesc relations stems from keeping }\Phi\geq0.
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for which $\Phi=0$ is represented by a line.
In the inverse description, an orientation may be specified as that of the sample triad in the reference frame of the crystal. The preferred direction is usually taken to be $Z^{s} ;\left\{\phi_{2}, \Phi\right\}$ become its longitude and latitude, respectively, and $\phi_{1}$ the azimuthal rotation around $Z^{3}$. (For the convential signs, see the last footnote.) Figure 8 sketches one constant $\phi_{1}$-section. This figure is the exact analog to Fig. 4a: it is evident how the singular point in Fig. 4 a has been stretched out into the line $\Phi=0$ in Fig. 8 . It is also clear that a symmetry reduction to region $I$ is an especially unfortunate choice. Region III, on the other rand, is close to the fariliar description in terms of triangles.

### 4.2 Polar coordingtes: the COD

Mosi ef the difficulties of visualization disappear when the angle pair $\left\{\phi_{1}, \Phi\right\}$ is represented in polar coordinntes; fust as in a pole figure; the third angle, $\phi_{2}$, can then be represented in a third dimension, perpendicular to the pular plot $[5,20]$. We call this a COD (irystal orientation distribution). For a two-dimensional image, a set of sections through this space at constant values of $\phi_{2}$, with contour lines, is adequate.

A $\phi_{2}$-section of the $C O D$ may be viewed as a "partial" \{001\} pole figure which showe the distribution of only those (001) poles that have corresponding (100) poles rotated $\phi_{2}$ degrees away from the equatur. Figure 9 displays a set of such partial pole figures for the copper specimen. These $\phi_{2}-\sec t i \mathrm{O}_{\mathrm{u}}$, of the COD contain the same information as those in Fig. 7 , but they are represented in polar rather than Cartesian coordinaten, and in equal-area projection. It appears to us that such a polar representation of the $C O D$ is considarably easier to read thar the traditional Cartesian way, which represents the ODF as a density function with three equivalent rotations. Let is list some of these advantages.
a) Orientations can be readily identified

Firstly, this is trivially true for directions, i.e. when only one axis is of importance: this is the advantage of having chosen a deacription in which first axis is explicitly given in the classical way. More generally, consider the maximum labelled $C$ at $\phi_{1}=40^{\circ}, \Phi=66^{\circ}$ in the $\Phi_{2}=25^{\circ}$ section. It contains (001) axes of those grains whose (100) axes lie $25^{\circ}$ away from the equator. By means of an equal-area net we can construct the full crystal orientation (Fig. loa). We see that (l2l) colncides with ND and [lī1] with RD; this is the well-known "copper" texture component [5]. For another case, consider the maximum labelled Bat $\phi_{1}=35^{\circ}, \infty=45^{\circ}$ in the $\phi_{2}=0$ section. The analysis in Fig. lob shows that this orientation has (011) pirallel to $N D$ and [2II] parallal to RD: the "brass" texture component.
b) Angles can be directly measured in the diagram

Agsume we would like to know the relation between the two orientations that are associated with the $C$ and $B$ maxima discussed above, expressed as a rotation about a single axis. This is demonstrated in Fig. loci construct the two orientations as in (a), superpose the two diagrams, and find the two bisectors (dotted): their interiection marks the axis of rotation $(\langle 210\rangle)$, and the angle $u)\left(-35^{\circ}\right.$; of the rotation around this axis that brings the two oriantations into coincidence is easily rad off. (This is not anique solution because of the nigh crystal symmetry [12]).
c) The symmetry is clearly displayed

Crystal and sample ammetry cause certain orientations in the $C O D$ tu be equivalent. For example, four-fold symmetry axis in [001] causes tz to repeat every $90^{\circ}$, and there is no need to extend $\phi_{2}-\sec t i o n s$ through a
full $360^{\circ}$ span. In order to be complete, the sector shown in the COD must contain at least one orientation of each aymmetrically equivalent set. A summary of equivalent orientations for important crystal and sample symmetries is shown in Table 1 . In the case of cubic crystal aymmetry and orthorhombic specimen symmetry, a range of $\phi_{1}$ from $0^{\circ}$ to $180^{\circ}$, of $\Phi$ from $0^{\circ}$ to $90^{\circ}$, and of $\phi_{2}$ from $0^{\circ}$ to $45^{\circ}$ is sufficient. A unit with all three angles going from $0^{\circ}$ to $90^{\circ}$ would also suffice, resulting in smaller sectors but more sections, which is less convenient for printing and also makes it more awkward to visualize angular relations. Both of these schemes atill contain three equivalent orientations due to the three-fold axis along <lll>; these are not easily recognizad in the COD. Sometimes it is useful to choose to plot more than one irreducible unit to Lllustrate symmetry relationships.

Polar COD's are particularly useful to check how closely an assumed specimen symmetry is approached. For example, in our rolled copper npecimens, the firct-measured pole figures, based on the "given" coordinate axes $N D$ and $R D$ were clearly not close enough to orthorhombic symmetry. We redefined the sample axes until the pole figure exhibited satisfactory orthorhombic symmetry, and only then averaged the four quadrants as described above. Often $1 t$ would be preferable to manipulate actual measurements as little as possible and to compare these "raw data" With theoretical predictions. In this case, variations that are amaller than the observed deviations from the appropriate sample symmetry should be treated as meaningless.

## d) There is visible connection to the pole fipure

ilewing the $C O D$ as at of partial \{001\} pole flgures, with each $\phi_{2}$-section contafiling the subset of (001) poles for a certain (100)


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elevation, implies that the aver. ge of all $\phi_{2}$-sections constitutes the complete $\{001\}$ pole figure. This is the same as a projection along $\phi_{2}$; it is shown in the last diagram of Fig. 9 and compares favorably with the measured (002) pole figure (Fig. 3b).


e) The representation is uniform

The space selected for representation and for projection (Euler space in pnlar equal-area projection) is homochoric with orientation space: equal densities seen are equal densities present. For thiu reason, we recommend equal-area projection also for pole figures [21]: it provides a better representation of the total fraction of crystals that contribute to a maximum (see Fig. 3a va Fig. 3b).

The points (a) through (d) made above illustrate that a COD in polar coordinates is indeed easy to visualize and to analyze quantitatively. It displays all the inform tion needed to derive full orientation relations using simple geometric constructions. If, in addition, one chooses equal-area plots, the visual impression is representative of the actual distribution (point e). The concept of viewing the $C O D$ as a pole figure deconvoluted into partial pole figures in orientation space is close in philosophy to the vector method [22].

### 4.3 The sample orientation distribution (SOD)

A deacription of sample axes in terms of cryatallographic axes can make use of the same Euler angles. Now the pair $\left\{\phi_{2}, \Phi\right\}$ describes the direction (typically of the $z^{s}-a x i s$ ) in the cryatal system, just like In any inverse pole figure, and $\phi_{1}$ describes the azimuthal rotation around this direction. An appropriate space for the $S O D$ is thus an equal-area projection (or the part of it that is necessary according to the aymetry of the crystal, as in Fig. 4b) and a perpendicular dimension along which
$\phi_{1}$ is plotted. It is seen that $\phi_{1}$ and $\phi_{2}$ have switched their roles as azimuth and longitude. Again, we can represent the three-dimensional SOD by a series of sections, this time at constant $\phi_{1}$ (Fig. 11). A projection along $\phi_{1}$ gives a $2^{8}$-axis figure: an 'inverse pole figure'. In analogy to the term 'partial pole figure' which we tentatively introduced above


A similar concept was introduced by Williams in 1968 [3]: he represented three-dimensional orientation distributions by a series of partial inverse pole figures (which he called "axial pole figures"), and also ahowed how these can be derived from pole figures by the "matrix method" (quite similar to the "vector method" [22]). Williams used a definition for the azimuthal angle (his $\beta$ ) that differs from the Euler angle $\phi_{1}$; while this has some advantages [3], s.t has the disadvantage that the duality between $\phi_{1}$ and $\phi_{2}, 1 . e$, between the $S O D$ and the $C O D$, gets lost. More significantly, using $\beta$ as the third dimension in the SOD would not make homochoric space. We therefore do not follow Williams in this respect.

In Fig. 11 , each $\phi_{1}$-saction displays the distribution of those apecimen normals $Z^{s}=N D$ which have their $X^{s}=R D$ axis $\phi_{1}$ degrees off the equator (clockwise rotation, cf. Fig. 6b). As explained in detall before, a range of $0^{\circ}$ to $90^{\circ}$ in each of the three angles would be afficient; in fact, there are still three equivalent orientations due to the three-fold <lll>-axis. They are not easily recognized in the $S O D$ sections (although they are evident in the projection). For example, the maxima at $\left\{\phi_{1}, \Phi, \phi_{2}\right\}=\left\{0,45^{\circ}, 0\right\},\left\{0,45^{\circ}, 90^{\circ}\right\}$, and $\left\{90^{\circ}, 90^{\circ}, 45^{\circ}\right\}$ are equivalent. (For an analytical expresion of this ammetry relation, see e.g. [19].) Because of this, only one unit triangle in Fig. 11 needs to be rapresented (an example is shown emphasized); but again, a larger
sector is of enen cas to visualize. In Fig. 11 , we have chosen to show twice the necessary range in $\phi_{2}$ (but it is not possible to compensate for this by reducing the range of $\phi_{1}$ to be from $0^{\circ}$ to $45^{\circ}$, as it could be done for $\phi_{2}$ in Fig. 9).

A special point must be made regarding the center of the polar plots: the singular point where $Z^{8}$ and $Z^{c}$ are parallel, the 'North pole'. Here, the meridian is multivalued ( $\phi_{1}$ in the $C O D, \phi_{2}$ in the $\operatorname{SOD}$ ). This correspunds to the physical situation (and was one of the reasons for us to abandon the Cartesian plots of Euler space). However, the heading of the boat is defined; thus, $\phi_{2}$ has a meaning in the COD, $\phi_{1}$ in the SOD. The trouble is that each is measured along the azimuthal great circle (the horizon) 'from the equator'; but in the special case $\Phi=0$, the horizon and the equator are the same thing, and thus the value of the azimuth is undefined. It is easy to overcome this apparent difficulty, when it is realized that $\phi_{1}$ and $\phi_{2}$ are sompletely equivalent rotations in this case;
 singuiar point. Thus, when we plot sections at constant azimuth, the value of the azimuth at the center point is meant to be that corresponding to the zero meridian.

Finally, we must emphasize that we have actually stayed with the conventional and expediant type of polar plot: with the singular point at the center-not, as suggested in Fig. $4(b)$, at the periphery. For cubic materials, the partial inverse pole figures (Fig. 11) could be converted by merely aperimposing on them a net of the kind used in Fig. 4(b); one would then be able to choose a reduced region with minimal distortion. For the partial pole figures, one would have to write a new ploting routine, which would then be incompatible with conventional pole plots.

## 5. CONCLUSIOAS

In the following, we summarize the points made in the present paper.

1) We have emphasized the distinction between directions (2-D quantities) and (3-D) orientations, but kept some continuity in their raspective quantative descriptions.
2) We have re-emphasized the need for representations of direction space and orientation space that do not distort densities (and labelled such spaces 'homochoric'). The surface of a unit sphere is a homochoric direction apace, Euler space is a homochoric orientation space. Unfortunately, it requires the preference, in the description, of one of the three coordinate axes defining an orientation-or, equivalently, of one direc:ion (around which a rotation defines the third parameter). The preferred direction should be described like any direction: as a point on the surface of unit sphere (e.g., by its longitude and co-latitude).
3) Two-dimensional representations of either the surface of a sphere or of a thraedimensional orientation space should preserve the homochoricity so strenuously achieved. Thus, the surface of aphere should be projected in equal-area projection, which is just as ensy to une as the stereographic projection, but much casier to visualize in terma of relative densities of directions. Similarly, three-dimensional orientatiou space should first be projected onto the surface of aphere (to describe the preferred direction, e.g. In terme of its longitude and co-latitude); then, this urface should be projected onto a circle in equal-area projection; and finally, the third dimension may be represented by means of constant-azimuth sections. We find the analog of the positions and headings of boats on the surface of the earth halpful.


#### Abstract

4) In application of these principles, we propose to use two dual orientation distribution representations: the $C O D$, which represents the orientation of a crystal coordinate system with respect to ample frame in terms of a set of partial pole figures (each containing those of the preferred poles that have a particular azimuth $\phi_{2}$ ); and the $S O D$, which conversely diaplays the orientations of sample coordinate axes with respect to a crystal frame in terms of a set of partial inverse pole figures (each containing those preferred axes that have a particular azimuth $\phi_{1}$ ). The final figure represented in each $C O D$ is the average partial pole figure (or $\phi_{2}$-projection), which is a pole figure; similarly, the final figure shown in a $S O D$ is the average partial inverse pole figure (or $\phi_{1}$-projection): it is an 'inverse fole figure'. A summary of these conventions is given in Table 2 . We find these polar representations of the orientation distribution much easier to visualize and evaluate than the Cartesian ODF representations that have become commen. 5) The net to be used in any of the representations should, while consisting of constant-area units, be also equiaxed (rather than having very anisotropic shapes). This is possible when the least symetric unit needed is chisen so as not to include the (North) 'pole' of the representation (see Fig. 4b).

Our principal concern has been to find two-dimensional graphic representations of three-dimensional orientation distributions that allow one to "see" something (with minimal distortion). We "like" two sets of circular sections; each set may be atacked up as a cylinder. Unfortunately, it is difficult to visualize three-dimensional apace through which the $\operatorname{CODR}$ and the $\operatorname{SODR}$ are difierent sections. This was easy in the conventional Cartesian Euler space.


#### Abstract

In all of our discussion, we have assumed that the threc-dimensional orientation distributions are knewn. To illustrate the proposed representations of orientation distributions, we have used a sample of a cubic metal deformed in rolling (and an oversimplified pole figure deconvolution). The same concepts have been applied to cases of lower crystal and sample symmetries, where the advantages are even more striking [23].

We end our discussion by outining the specific, very minor cns: es that would have to be made in existing computer codes to implement unr suggeftions. The calculation of ODF'a from pole figures is unaffected; but instead of using the contouring routine for a rectangular grid, one should use, for each ODF section, the pole-figure contouring routine (such es Vedon's URFPD). To convert from stereographic to equal-area projection, one simply must cirange the formula for the projection of the pole distance $\alpha$ from $\tan (\alpha / 2)$ to $\gamma 2 \cdot \sin (\alpha / 2)$.


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Table 1: Equivalent orientations for some crystal and sample symmetries

| \$1 | $\Phi$ | ¢2 |  |
| :---: | :---: | :---: | :---: |
| $\phi_{1}+\pi$ | -Ф | $\phi_{2}+\pi$ | identity |
| $\phi_{1}$ | $\Phi$ | $\phi_{2}+\pi / 3$ | 6-fold axis in $\mathrm{Z}^{\text {c }}$ (hexagonal) |
| $\phi_{1}$ | $\Phi$ | $\$_{2}+\pi / 2$ | 4-fold axis in $\mathrm{Z}^{\text {c }}$ (cubic, tetrag.) |
| $\Phi_{1}$ | $\Phi$ | $\phi_{2}+2 \pi / 3$ | 3-fold axis in $\mathrm{Z}^{\text {c }}$ (trigonal) |
| $\phi_{1}+\pi$ | $\pi-\Phi$ | - $\phi_{2}$ | 2-fold axis $\mathrm{Y}^{\text {c }}$ |
| $-\phi_{1}$ | $\pi-\Phi$ | $\phi_{2}+\pi$ | 2-fold axis in $\mathrm{Y}^{\text {S }}$ |
| $\Phi_{1}+\pi$ | $\Phi$ | $\phi_{2}$ | 2-fold axis in $\mathrm{Z}^{8}$ |
| $\pi-\oplus_{1}$ | $\pi-\Phi$ | $\phi_{2}+\pi$ | 2-fold axis in $\mathrm{X}^{s}$ |

Table 2: Representations of the three-dimensional orientadion distribution $\operatorname{ODF}\left(\phi_{1}, \Phi, \phi_{2}\right)$

| 2-D sections* <br> COD $\left(\phi_{1}, \Phi\right) ; ~ \phi_{2}=$ const <br> "partial pole figures" <br> SOD $\left(\phi_{2}, \Phi\right) ; \phi_{1}=$ const <br> "partial inv. pole figures" <br> pole figure $(\beta, \alpha)$ <br> For $\Phi=0, \phi_{1}+\phi_{2}=$ const |  |
| :--- | :--- |

## CAPTIOHS

Fig. 1 A direction 2 is represented as a location on the surface of a sphere: by its longitude $\phi_{1}$ and its pole distance $\Phi$ (perspective drawing). An orientation $g\left(\phi_{1}, \Phi_{,} \phi_{2}\right)$ is represented by $Z\left(\phi_{1}, \Phi\right)$ and an azimuth $\phi_{2}$ around it. An orientation distribution corresponds to a distribution of boats (with specified headings) on the surface of the earth.

Fig. 2 Projeriton of the coordinate grid on the sphere: (a) stereographic projectiou (Wulff net), (b) equal-area projection (Schmidt net).

Fig. 3 Experimentally determined $\{200\}$ pole figure of copper rolle to $50 \%$ reduction at room temperature. Transverse (TD) and rolling (RD) direction are indicated. (a) stereographic projection, (b) equal area projection.

Fig. 4 (a) The quairant of an inverse pole figure for cubic crystals ps it is conventionaily drawn. (b) An equivalent quadrant in equalarea projection, not in: uding the sperial pole (001).

Fig. $5\{100\}$ and $\{111\}$ pole figures ( $\mathrm{a}, \mathrm{b}$ ) and inverse pole figure for the $N D$ and $R D$ axis ( $c, d$ ) for copper rcalculated from the ODF with the harmonic analysis. Equal-ares prijection. Contour intervals are $0.5 \mathrm{~m} . \mathrm{r}_{\mathrm{d}} \mathrm{d} \cdot ;$ stippled below 0.5 m .1 .d.

Fig. 6 Lefinition of Euler angles $\phi_{1}, \Phi, \phi_{2}$ using the convention of Bunge !1], based on the sampie coordinate system (a) and the crystal foordinate system (b).

Fig. 7 ODF of rolled copper, represented in conventional Cartesian coordinates. Contour i:tervals $0.5 \mathrm{~m} . \mathrm{r}_{\mathrm{d}} \mathrm{d} .$, stippled below 0.5 m.r.d. Constant- $\phi_{2}$ sections, extending Grom $0^{\circ}$ to $45^{\circ}$. The range of $\Phi$ is from 0 to $90^{\circ}$ (down), that of $\phi_{1}$ is from 0 to $180^{\circ}$ (right).

Fig. 8 A section of conventional Cartesian Euler space (constant $\phi_{1}$ ) for cubic materials with irreducible regions $I$, IIa/b, and III. Compare Fig. 4a. Region III is least distorted.

Fig. 9 COD of rolled copper represented as partial pole figures in equalarea projection, corresponding directly to Fig. 7. The last diagram is an average over all partial pole figures and corresponds to a ( 001 ) pole figure (Fig. 3 and Fig. 5a). The most common components of the f.c.c. rolling, texture are indicated, by one representative: $C$ - the 'copper' component $\{121\}\langle 1 \bar{I} 1\rangle ; s-$ - the ' $S$ ' component $\{132\}\langle 6 \overline{4} 3\rangle ; B=$ the 'brass' component $\{011\}\langle 2 \bar{I} 1\rangle ; C$ the 'Goss' component $\{011\}\langle 100\rangle$; and also the ' $\alpha$-fibre'.

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Fig. 10 Derivation of Orientation Relationships,
    (a) Reconstruction of the crystal orientation for the maximum at
        \phi}=4\mp@subsup{0}{}{\circ},\Phi=6\mp@subsup{6}{}{\circ},\mp@subsup{\phi}{2}{}=2\mp@subsup{5}{}{\circ}\mathrm{ using the COD of Fig. }9\mathrm{ and an
        equal-area net (Cu component).
    (b) Same as (a) for the orientation }\mp@subsup{\phi}{1}{}=3\mp@subsup{5}{}{\circ},\phi=4\mp@subsup{5}{}{\circ},\mp@subsup{\phi}{2}{}=\mp@subsup{0}{}{\circ
        (B component).
    (c) Superposition of (a) and (b) to determine the rotation axis
        and angle w which brings the two orientations into
        coincidence.
Fig. 11 SOD of rolled copper represented as partial inverse pole figures.
    The test diagram is an average over partial inverse pole figures
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    figure for the ND axis (Fig. 5c). The letters indicate f,c.c.
    rolling components as in Fig. 9, and the ' }\beta\mathrm{ -fibre'.
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Pig. 1 A direction 2 is represented ad location on the ourface of a aphere: by ite loagitude $\phi_{1}$ und ita pole distance $\Phi$ (perapective drawiog). An orientation $g\left(\phi_{1}, \Phi_{1} \phi_{2}\right)$ is represented by $Z\left(\phi_{1}, \phi\right)$ and an aefmuth $\phi_{2}$ around it. An orientation distribution correspouds to a distribution of boate (with epecified headinge) on the surface of the carth.
(a)

(b)


Fig. 2 Projaction of the coordinate grid on the aphera: (a) stereographic projection (Wulff net). (b) equal-area prujection (Schmidt net).


Fig. 3 Experimentally deternined $\{200\}$ pole figure of copper rolled to $50 \%$ reduction at room temperature. Transverse (TD) and roling (RD) Airection are indicated. (a) stereographic projection, (b) equal aree projection.



Pig. 5 \{ 100 \} and $\{111\}$ pole figure e ( $a, b$ ) and favarae pole figure for the ND and ND axis ( $c, d$ ) for copper recalculated from the ODF with the harmonic mayas. Equal-aran projection. Contour latervale are 0.5 m.r.d.l stippled below 0.5 E.r.d.




Yis. 7 ODF of rolled copper, represented in conventional Cartenian coordinates. Contour intarvale 0.5 n. r.d., etippled below 0.5 arrod. Constant-it acectiona, extendiag frem $0^{\circ}$ to $45^{\circ}$. The range of If $^{10}$ from 0 to $90^{\circ}$ (down), thet of id th fron 0 to $180^{\circ}$ 1-.-…


Fip. A A section of conventional Certesian Euler apace (constant $\phi_{1}$ ) for cubic materiale with irreducible resione $I$, IIa/b, and III. Compare Fis. 4a. Region III la least dietorted.


T18. 9 COD of rolled copper rupresented en partial pole figures in equalarea projection, corresponding directly to Fig. 7. The last dian gram 19 an average over all partial pole elgures and corresponds to (001) pole figure (Fig. 3 and Fig. 5a). The moit common componente of the f.c.c. rolling texeura are indiceted, by one repreaentativei $C$ - the 'copper' component $\{121\}<1 I \mid>;$. the ' $S$ ' component $|132|<643>1, L$ - the 'brane' comporent $\{011 \mid<2[1\rangle ; 6$ the 'Gose' component $\{011\}\langle 100\rangle$ ' and also the ' $\alpha-110$ ere'.


Fig. 10 Derivation of Orientation Relationshif.
(a) Reconstruction of the crystal orientation for the maximum at $\varphi_{1}=40^{\circ}, \Phi=86^{\circ}, \varphi_{2}-25^{\circ}$ uating the COD of Fig .', and an equal-area eet ( Cu compoent).
(b) Same (a) for the orientation $\phi_{1}=35^{\circ}, 1=45^{\circ}, \varphi_{2}-0^{\circ}$ (B component).
(c) Saperpoiition of (a) and (b) to deternine the rotation axie and angle which bilage the two orientations into colncidence.


Fis. 11 SOD of rolled copper represented ae partial inverse pole figures. The test diagram is an avarage over partial inverse pole figures from $\Phi_{1}$ - $\theta^{\circ}$ to $\phi_{1}$ - $180^{\circ}$ and corresponds to an inverse pole figure for the ND axis (Fig. Sc). The letters indicate f.c.c. rolling componente at in Fig. 9 , and the ' $\beta$-fibre'.


[^0]:    t Work supported by the U.S. Department of Energy.

[^1]:    * $\beta$ is also sometimes called an azimuth, but we reserve this term for the "heading" introduced above; see also sec. 4.1 .

