# CLOUD ATLAS: HUBBLE SPACE TELESCOPE NEAR-INFRARED SPECTRAL LIBRARY OF BROWN DWARFS, PLANETARY-MASS COMPANIONS, AND HOT JUPITERS

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(Accepted December 7, 2018)

# ABSTRACT

Bayesian atmospheric retrieval tools can place constraints on the properties of brown dwarfs and hot Jupiters atmospheres. To fully exploit these methods, high signal-to-noise spectral libraries with well-understood uncertainties are essential.

We present a high signal-to-noise spectral library (1.10-1.69  $\mu$ m) of the thermal emission of 76 brown dwarfs and hot Jupiters. All our spectra have been acquired with the Hubble Space Telescope's Wide Field Camera 3 instrument and its G141 grism.

The near-infrared spectral types of these objects range from L4 to Y1. Eight of our targets have estimated masses below the deuterium-burning limit. We analyze the database to identify peculiar objects and/or multiple systems, concluding that this sample includes two very-low-surface-gravity objects and five intermediate-surface-gravity objects. In addition, spectral indices designed to search for composite atmosphere brown dwarfs, indicate that eight objects in our sample are strong candidates to have such atmospheres. None of these objects are overluminous, thus their composite atmospheres are unlikely a companion-induced artifact. Five of the eight confirmed candidates have been reported as photometrically variable, suggesting that composite atmospheric indices are useful in identifying brown dwarfs with strongly heterogeneous cloud covers.

We compare hot Jupiters and brown dwarfs in a near-infrared color-magnitude diagram. We confirm that the coldest hot Jupiters in our sample have spectra similar to mid-L dwarfs, and the hottest hot Jupiters have spectra similar to those of M-dwarfs. Our sample provides a uniform dataset of a broad range of ultracool atmospheres, allowing large-scale, comparative studies, and providing a HST legacy spectral library.

Keywords: Brown dwarfs - stars: atmospheres

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# 1. INTRODUCTION

Over the past decade, increasingly detailed observations are available on a wide range of objects: spectroscopic information is now available on hot Jupiters (e.g., Stevenson et al. 2014a; Ranjan et al. 2014; Line et al. 2016; Evans et al. 2017; Sheppard et al. 2017), directly imaged exoplanets (e.g., Samland et al. 2017; Rajan et al. 2017), and over a thousand brown dwarfs (e.g., Cushing et al. 2005; Kirkpatrick 2005; Burgasser et al. 2006a; Apai et al. 2013; Buenzli et al. 2014; Schneider et al. 2015). These datasets have enabled major steps in the complexity and quantitative evaluation of atmospheric models. A particularly significant advancement has been the adaptation of Bayesian modeling framework first for hot Jupiters (e.g., Madhusudhan & Seager 2009; Line et al. 2013; Lee et al. 2014; Gandhi & Madhusudhan 2017; Pinhas et al. 2018; Fisher & Heng 2018) and smaller transiting planets (Benneke & Seager 2012), then for directly-imaged exoplanets (Todorov et al. 2016; Lavie et al. 2017), and, most recently, for brown dwarfs (Line et al. 2015, 2017; Madhusudhan et al. 2016).

The Bayesian modeling framework – although often less detailed than forward models - has two key advantages: first, it provides a probabilistic assessment of the fitted parameters and degeneracies, even if the parameter space is highly complex. Second, it enables systematic, comprehensive, unbiased modeling of large number of atmospheres, allowing for detailed comparative studies of the posterior distributions of the model parameters (e.g., C/O ratios, molecular abundances, surface gravities). Although the information provided by posterior distributions is very powerful, it must be remembered that the probabilities derived for individual model components are only correct under the assumption that the data and uncertainties are correctly represented by the priors, and that the modeling framework itself is complete and correct. For example, data with hidden biases (resulting in incorrect priors) will yield systematically incorrect posteriors. In this sense, due to typical observational biases, it is particularly challenging to compare objects over a broad range of parameters (e.g., very different temperatures or surface gravities). In fact, no spectral library with well understood systematics exists for ultracool atmospheres (hot Jupiters, directly imaged exoplanets, brown dwarfs). In short, to exploit the potential of atmospheric retrievals and enable rigorous comparative studies of atmospheres, homogeneous spectral datasets with well-understood systematics are required.

Comprehensive spectral libraries exist for brown dwarfs (Kirkpatrick et al. 1999, 2000; Leggett et al. 2000; Burgasser et al. 2002; McLean et al. 2003; Cush-

ing et al. 2005; Kirkpatrick 2005; Burgasser et al. 2006b, the SpeX spectral library<sup>1</sup>, the Montreal spectral library<sup>2</sup>, and references therein) built from ground-based spectroscopy of hundreds of brown dwarfs in dozens of studies. These libraries have played and continue to play essential roles in a broad range of brown dwarf studies. However, existing ground-based spectral libraries were built from data that are non-uniform in terms of instruments, observing conditions, setups, and usually reduced slightly differently by different groups.

While these spectral libraries remain powerful, these datasets have several limitations for atmospheric retrieval studies: first, it is not possible to reliably capture the variety of differences in data acquisition, quality, and reduction with priors due to the unknowns involved. Second, ground-based observations unavoidably are influenced by telluric absorption, most notably by water bands. Although it is possible to correct for these to some extent, their time-varying nature and the optical depth in the bands lead to limited reliability in these bands. In fact, quantitative comparisons (Apai, priv. comm.) of some brown dwarfs with SpeX spectra and Hubble Space Telescope Wide Field Camera 3 near-infrared grism spectra, revealed mismatches in water band shape, and overall color (wavelength-dependent slope). We show these differences in Fig. 1. In the left column, we show the direct comparison of the near infrared SpeX (black) and HST/WFC3 spectra (blue) for randomly selected brown dwarfs with spectra in the SpeX and in our HST/WFC3 near infrared library, and with spectral types between L5.5 and T6. In the right column, we show the ratio of the SpeX and the HST/WFC3 near infrared spectra for the object in the left column. We show a green line indicating where a perfect match between the SpeX and the HST/WFC3 spectra should be (ratio Spex vs HST/WFC3 equal 1). In addition, we fit a line to the slope of the ratio between the two spectra, avoiding the water band at 1.4  $\mu$ m, (see black line) showing that in most of the cases, the slope is non-zero, indicating color trends on the SpeX spectra. In these plots a common mismatch between the SpeX and the HST/WFC3 near infrared spectra at the 1.4  $\mu$ m

 $<sup>^1</sup>$  The SpeX Prism Library is composed by low-resolution, near-infrared spectra, primarily of low-mass stars and brown dwarfs, obtained with the SpeX spectrograph mounted on the 3 m NASA Infrared Telescope Facility on Mauna Kea, Hawaii. The data provided here have been obtained using the prism-dispersed mode of SpeX with an average resolution of  $\sim\!120$  and spectra spanning 0.90-2.50  $\mu\rm m$ : http://pono.ucsd.edu/~adam/browndwarfs/spexprism/library.html

 $<sup>^2 \</sup>quad \text{https://jgagneastro.wordpress.com/the-montreal-spectral-library/}$ 

water band is also evident, due to imperfect telluric correction. Given that the photometric precision and instrumental systematics of the HST/WFC3 instrument are very well understood, and that the HST/WFC3 near infrared spectra are not affected by tellurics, the comparison reveals that low-level biases exist in the ground-based spectral libraries. While these corrections are well-suited for forward-modeling and object-to-object comparisons, they are often limiting for retrieval studies.

For example, Line et al. (2015) applied Bayesian atmospheric retrieval tools to SpeX spectra to derive thermal structures and molecular abundances of some brown dwarfs. Nevertheless, they could only reach convergence in their Markov Chain Monte Carlo retrievals if they assumed that the SpeX spectral uncertainties were underestimated. Thus, they artificially increased the SpeX spectral uncertainties that could reach a maximum of a factor of 100. In this case, it is impossible to disentangle if Line et al. (2015) models were incomplete, or the uncertainties of the datasets were not accurately estimated and understood. Therefore, to properly test retrieval models, there is an obvious need for a uniform, space-based spectral library with well-understood spectral uncertainties.

In this paper, we are presenting a high signal-to-noise spectral library with 76 near-infrared WFC3/HST spectra of brown dwarfs, low-mass companions to stars, and hot Jupiters. Our study supplements the HST Cloud Atlas Treasury program data (GO 14241, PI: D. Apai) with other published datasets (see Section A), carefully analyzing and correcting for (typically very small) data reduction differences. The advantage of the HST/WFC3 instrument is that it provides near-infrared spectroscopy (1.10–1.69  $\mu$ m, S/N up to 3,000 in the J-band) where the SED (Spectral Energy Distribution) of these objects peak, and we observe the dominant absorbing species for brown dwarfs.

Besides the presentation of our HST/WFC3 near infrared spectral library, the goal of this study is to provide a comprehensive characterization of the objects of our sample. This spectral characterization is important to validate if the results provided by the retrieval models match with the expectations for a given object. In fact, Line et al. (2015) used two well-characterized T dwarfs, Gl 570D and HD 3651B, to test their retrievals, confirming that the effective temperatures, surface gravities, masses, radii, etc., were consistent with the expected values for those T-dwarfs. Therefore, we aim to provide a basic characterization of our sample, and also to identify peculiar objects: extremely red or blue brown dwarfs, revealing low surface gravity or low metallicity objects, and overluminous brown dwarfs, potentially re-

vealing multiple systems. In addition, we use spectral indices (Burgasser et al. 2006b, 2010; Bardalez Gagliuffi et al. 2014) to search for spectral binaries. As a byproduct of this analysis, we found that these spectral indices can also be useful to search for variable brown dwarfs. In addition, we show the potential of our spectral library by performing a novel direct photometric and spectroscopic comparison between hot Jupiters and brown dwarfs, that can be only be accomplished at this level of accuracy with data acquired from the space. This comparison confirms that some hot Jupiters share effective temperatures and spectra with some M and L dwarfs.

In Section 2, we describe the targets that we include in the near infrared HST/WFC3 spectral library. In Section 3 we derive spectral types for our sample using the SpeX Prism Spectral Library. In Section 4 we compare the L and T dwarfs with other brown dwarfs from Dupuy & Liu (2012) using a color-magnitude diagram. In Section 5 we search for low surface gravity objects in our sample using low surface gravity spectral indices. In Section 6, we search for composite atmosphere objects within our sample. In Section 7 we measure the water and methane bands on objects of our sample, to trace the change on the depth of those bands with the near infrared spectral type. In Section 8 we directly compare colors and spectra of brown dwarfs and hot Jupiter in a color-magnitude diagram and finding the best matching brown dwarf to every of the hot Jupiters in our sample. Finally, in Section 9 we present our conclusions.

# 2. TARGETS

We compiled all spectra of brown dwarfs (including planetary-mass brown dwarfs), brown dwarf companions to stars, and hot Jupiter with emission spectra with published data from HST/WFC3 and its G141 grism (MacKenty et al. 2010). In addition, we present seven unpublished spectra observed as part of the Cloud Atlas treasury program (GO 14241, PI Apai), and other two brown dwarf spectra that belong to the HST program GO 13299 and 14051 (PI Radigan). These spectra cover the wavelength range between  $\sim 1.10$  and 1.69  $\mu$ m, with a spectral resolving power  $R=\frac{\lambda}{\Delta\lambda}=130$  at 1.4  $\mu$ m. The image scale of WFC3/IR is 0.13 arcsec/pixel. In Tables 1 and 2 we provide the list of objects with names, celestial coordinates, and HST program identifiers in which the objects were observed, as well as the most relevant references in which this spectra were first published (Buenzli et al. 2012, Buenzli et al. 2014, Buenzli et al. 2015, Yang et al. 2015, Yang et al. 2016, Lew et al. 2016, Manjavacas et al. 2018, Peña Ramírez et al. 2015, Zhou et al. 2018, Biller et al. 2018, Apai et al.

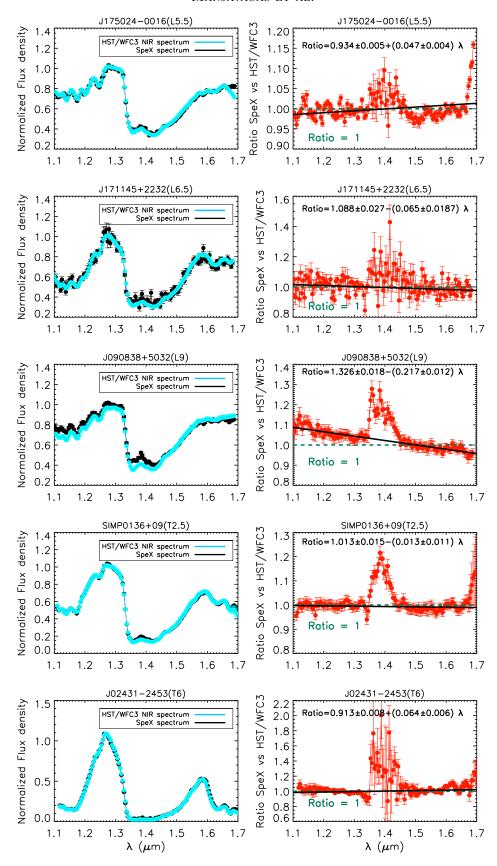


Figure 1. In the left column, we show a direct comparison between the SpeX and the HST/WFC3 near infrared spectra. In the right column we show the ratio between the SpeX and the HST/WFC3 near infrared spectra. These plots reveal clear differences in the water band at 1.4  $\mu$ m between both spectra, and color trends shown by non-zero slopes of the linear fits.

2013, Schneider et al. 2015, Line et al. 2016, Haynes et al. 2015, Cartier et al. 2017, Stevenson et al. 2014a, Ranjan et al. 2014, Beatty et al. 2017, Stevenson et al. 2014b, Evans et al. 2017, Sheppard et al. 2017, and this work). In Table 1, for brown dwarfs and substellar companions, we specify their spectral types, Two Micron All Sky Survey (2MASS, Cutri et al. 2003) photometry, and trigonometric parallaxes, the signal-to-noise as measured at 1.25  $\mu$ m from the corresponding HST programs, the HST program for which each object was observed, and the references in which these spectra were published. For hot Jupiters (Table 2) we also list the spectral types of the host stars, the star–planet separations, and the radii of the planets. In Figures 2, 3, and 4 we show the spectra of the objects used in our study,

76 in total (22 L dwarfs, 28 T dwarfs, 16 Y dwarfs and 10 hot Jupiters), from which nine are presented here for the first time. Some of the objects in our sample in Table 1 have been classified as planetary-mass objects (M < 13  $\rm M_{Jup}$ ). We list these objects in Table 3 along with their estimated masses, ages, young moving group membership (if known), and key references.

The observing log, including the observing dates, number of orbits per visit, exposure time of each single exposure, and the number of single exposures per orbit is compiled in Table A1 in the Appendix.

The datasets and the data reduction by the authors that published the spectra presented in Tables 1 and 2 are described in Appendix A.

Table 1. Sample of L, T and Y dwarfs with HST/WFC3 spectroscopy

Ш	MANJAVACAS ET AL.																																		
Ref.	LS	1,13	TS, 14	2, 15	2	1	1, 13	1, 16	3, 17	4, 18	TS, 19	1, 20	1, 19	1	10, 21	2, 22	5, 23	1, 24	1	1, 24	1, 16	1	LS	П	П	11, 6, 25	1, 6	LS	6	6, 11, 26	$^{\mathrm{LS}}$	1, 16	1	1, 22	1
HST GO°	14241	12550	14241	13176	13176	12550	12550	12550	14241	14241	14241	12550	12550	12550	14188	14241	13280	12550	14241	12550	12550	12550	13299	12550	12550	12314	12550	14051	14241	12314	14241	12500	12550	12500	12500
$\mathrm{SNR}^\mathrm{b}$	45	2000	322	370	526	1250	2000	1111	370	333	208	1666	1428	1428	285	47	200	2000	296	1428	1428	1428	3333	1250	1250	172	11111	3333	2222	370	100	606	299	833	606
$\pi_T \; [{ m mas}]$	44.63±0.03	$45.16\pm0.16$	$108.70\pm2.36$	$122.00\pm13.00$	$106.50\pm0.19$		$81.97 \pm 2.69$	$33.11\pm4.71$	$82.00\pm 3.00$	$29.60\pm 2.80$	$64.13\pm4.51$		$86.21 \pm 4.46$		$45.10\pm1.70$	$86.20\pm1.10$	$501.40\pm0.09$	$93.46\pm0.87$		$65.79\pm2.16$	$36.90 \pm 3.40$					$101.50\pm2.00$			$54.37\pm0.85$	$162.87 \pm 1.06$	$21.00\pm0.07$	$36.23 \pm 4.46$	$14.84\pm0.12$	$96.15\pm0.96$	
$K_s \; [{ m mag}]$		$11.84\pm0.02$	$11.53\pm0.02$	$11.65\pm0.02$	$11.31\pm0.02$	$13.44\pm0.04$	$12.52\pm0.02$	$14.72 \pm 0.09$	$13.05\pm0.03$	$15.14\pm0.13$	$13.71\pm0.04$	$12.50\pm0.02$	$12.59\pm0.03$	$14.53\pm0.10$	$14.74\pm0.12$	$12.02\pm0.02$	$9.44\pm0.07$	$13.03\pm0.03$	$12.95\pm0.03$	$14.00\pm0.05$	$14.31\pm0.07$	$14.31 \pm 0.07$	$13.88\pm0.06$	$15.15\pm0.16$	$15.17\pm0.15$	$13.58\pm0.04$	$14.05\pm0.06$	$14.04\pm0.03$	$15.12\pm0.03$	$12.56\pm0.02$		$15.48\pm0.189$		$13.57\pm0.06$	$16.14\pm0.31$
H [mag]		$12.41\pm0.02$	$12.53\pm0.03$	$12.39\pm0.02$	$11.89\pm0.02$	$14.28\pm0.04$	$13.10\pm0.03$	$15.79\pm0.11$	$13.97\pm0.04$	$15.89\pm0.14$	$14.51\pm0.04$	$13.09\pm0.03$	$13.33\pm0.03$	$15.45\pm0.14$	$15.72\pm0.17$	$12.81 \!\pm\! 0.03$	$10.37\pm0.04$	$13.79\pm0.03$	$13.48\pm0.03$	$14.61\pm0.04$	$14.93\pm0.07$	$14.91\pm0.07$	$14.11\pm0.04$	$15.34\pm0.11$	$15.21 \pm 0.09$	$14.16\pm0.05$	$14.57\pm0.06$	$14.48\pm0.03$	$15.40\pm0.03$	$12.77 \pm 0.03$		$15.95\pm0.13$	$14.73\pm0.07$	$13.67 \pm 0.04$	$15.82\pm0.15$
$J [\mathrm{mag}]$		$13.29\pm0.02$	$14.05\pm0.02$	$13.43\pm0.02$	$12.83\pm0.03$	$15.56\pm0.05$	$14.03\pm0.03$	$17.08\!\pm\!0.18$	$15.60\pm0.07$	$17.22\pm0.21$	$15.82\pm0.06$	$14.11\pm0.03$	$14.48\pm0.03$	$16.27\pm0.13$	$16.71\pm0.19$	$14.07\pm0.03$	$11.5\pm0.04$	$15.10\pm0.03$	$14.55\pm0.02$	$15.86\pm0.07$	$16.02\pm0.08$	$15.91\pm0.08$	$14.95\pm0.04$	$16.41\pm0.15$	$16.03\pm0.09$	$14.71\pm0.01$	$15.59\pm0.07$	$15.29\pm0.04$	$15.86\pm0.03$	$13.46\pm0.03$		$16.34\pm0.10$	$15.06\pm0.04$	$13.80\pm0.02$	16.24±0.11
Dec. (J2000)	-35 49 31.20	-00 16 15.11	$+11\ 33\ 43.70$	$+14\ 14\ 01.04$	-16 27 38.62	-63 06 02.25	-00 59 01.90	$+22\ 32\ 04.41$	$+68\ 03\ 52.10$	$+35\ 58\ 09.80$	$+00\ 41\ 56.40$	+48 47 41.69	-45 21 54.89	+46 28 49.84	-22 51 35.84	-01 58 52.14	-53 19 10.08	$+21\ 15\ 52.12$	$+50\ 32\ 08.82$	$+19\ 04\ 40.71$	$+16\ 48\ 15.60$	$+31\ 28\ 49.71$	$+32\ 47\ 18.39$	$+32\ 56\ 26.40$	$+65\ 25\ 27.57$	$+02\ 20\ 22.70$	$+63\ 58\ 28.15$	$+03\ 35\ 35.01$	$+14\ 46\ 07.80$	$+09\ 33\ 47.30$	+17 04 31.80	+175904.30	$+25\ 54\ 18.10$	-14 04 48.88	+13 52 28.50
R. A. (J2000)	06 09 19.21	17 50 24.84	03 55 23.37	18 21 28.15	15 07 47.69	04 21 07.19	$05\ 39\ 52.00$	17 11 45.73	00 47 01.06	09 51 04.60	$01\ 07\ 52.42$	15 15 00.83	06 24 45.95	08 01 40.56	21 14 08.03	22 24 43.82	10 49 18.92	$08\ 25\ 19.69$	09 08 38.04	$16\ 32\ 29.11$	$03\ 10\ 59.87$	$12\ 19\ 51.56$	07 58 40.03	$10\ 39\ 31.38$	98.00 60 60	21 39 26.77	$13\ 24\ 35.54$	$16\ 29\ 18.62$	21 44 28.47	$01\ 36\ 56.62$	$01\ 12\ 36.48$	$17\ 50\ 32.94$	$00\ 00\ 13.54$	$05\ 59\ 19.14$	23 39 10.25
$\mathrm{SpT}^{\mathrm{a}}_{\mathrm{SpeX}}$	L7.5	L5	L5.5	L5.5	L5.5	T2	T2		L5.5	L4.5	L5.5	P7	L5	P7	L7.5	L4.5		P7	L7.5	F8	T2	L8		T2	T1	T1.5	T0.5	T2	T2.5	T2.5	T1.5	T3.5	T4.5	T4.5	T5
$_{ m PpT}$	L4	L4.5	L5	L5	L5	L5	L5	L5.0+T5.5	L6	Te	Fe	Te	L6.5	L6.5	L7	L7.5	L7.5+T0.5	L7.5	L8	L8	L8	F3	T0+T3.5	T1	T1.5	T2	T2	T2	T2.5	T2.5	T3.5	T3.5	T4.5	T4.5	T5.4
Name	CD-352722b	$2 \rm MASS\ J17502484-0016151$	$2MASS\ J03552337+1133437$	$2 MASS\ J18212815 + 1414010$	$2MASSW\ J1507476-162738$	$2 \rm MASSI~J0421072-630602$	$2 \rm MASS\ J05395200-0059019$	2MASSI J1711457 + 223204	2MASS J00470038+6803543	$\rm LP261-75B$	$2MASS\ J01075242+0041563$	2MASSW J1515008+484742	$2 MASS\ J06244595-4521548$	2MASSW J0801405 + 462850	PSO J318.5-22	$2 \rm MASSW~J 2224438-015852$	Luh 16AB	2MASSI J0825196+211552	2MUCD 10802	$2MASS\ J16322911+1904407$	2MASSW J0310599+164816	2MASS J12195156+3128497	SDSS J075840.33+324723.4	2MASS J10393137+3256263	$2 MASS\ J09090085 + 6525275$	$2MASS\ J21392676+0220226$	$2MASS\ J13243553+6358281$	$2 MASS\ J16291840 + 0335371$	HN PEG B	SIMP J013656.5+093347.3	GU PSC B	$2MASS\ J17503293+1759042$	$2 \rm MASS\ J00001354 + 2554180$	$2 \rm MASS~J05591914{-}1404488$	2MASS J2339101+135230
Num.	1	2	3	4	2	9	7	œ	6	10	11	12	13	14	15	16	17	18	19	20	21	22	23	24	25	26	27	28	29	30	31	32	33	34	35

Table 1 continued

Table 1 (continued)

	l					Η	ST	\\	VF	C3	B I	NFI	RAI	REI	s S	PE	СТ	RA	LΙ	ĹΙΒ	RA	RY	Ol	F S	UB	ST	ELI	LΑF	? О	BJ	ECT 
Ref.	TS, 27	12, 6, 25	1, 19	7	1, 22	1, 28	TS, 29	7	$^{\mathrm{LS}}$	∞	∞	<b>∞</b>	∞	œ	8, 27	8, 34	8, 34	8, 34	8, 34	8, 34	8, 34	8, 34	8, 34	8, 34	8, 34	8, 34	8, 34	8, 34	8, 34	8, 34	∞
$\mathrm{HST}\ \mathrm{GO}^{\mathrm{c}}$	14241	12314	12550	12217	12550	12550	14241	12217	14241	13178	12970	13178	12970	12970	12230	12970	12970	12970	13178	12970	12230	12230	12970	12970	12970	12230	12970	12970	12230	12970	13178
$\mathrm{SNR}^\mathrm{b}$	147	243	1000	22	299	606	6.7	τċ	65	∞	22	23	4	2	2	2	2	2	2	2	2	ಬ	က	1	1	1	2	1	1	1	П
$\pi_T \; [{ m mas}]$	$52.10\pm1.20$	$92.10\pm 2.60$	$204.08 \pm 12.49$		$93.46 \pm 3.49$	$90.91 \pm 1.65$	$18.30 \pm 1.80$		$85.54 \pm 1.53$		$14.52\pm0.06$					$75.40\pm6.62$	$153.40\pm4.00$	$67.60\pm 8.70$	$85.10\pm9.30$	$167.10\pm4.20$	$128.50\pm6.30$	$138.30\pm3.90$	$84.10\pm5.90$		$139.00\pm4.30$	$144.30\pm 8.60$	$228.10\pm 8.90$	$79.50\pm 8.80$	$168.80\pm 8.50$	$83.70\pm5.70$	
$K_s \; [{ m mag}]$	$52.10\pm1.20$	$15.29\pm0.21$	$13.52\pm0.04$		$15.22\pm0.17$	$11.00\pm0.20$		$20.91 \pm 0.15$	$16.89\pm0.06$	>12.918		>12.297	>12.278		$96.00\pm41.00$	$15.38\pm0.05$	$14.11\pm0.05$		$15.06\pm0.05$			$11.73\pm0.25$	$14.71\pm0.06$	$154.40\pm5.70$	$14.58\pm0.06$			>12.349	$14.74\pm0.04$		$15.01\pm0.09$
$[{ m mag}]$	$15.92 \pm 0.14$	$15.36\pm0.12$	$13.53\pm0.03$		$15.14\pm0.11$	$15.52\pm0.10$		$20.83\pm0.12$	$17.01\pm0.04$	$16.21 \pm 0.15$		$15.73\pm0.06$	$14.87\pm0.06$	$20.37\pm0.20$	$21.80\pm0.80$					$20.99\pm0.52$	$20.39\pm0.33$	$13.84\pm0.04$		$22.98\pm0.31$		$21.45\pm0.41$		$14.90\pm0.05$			
$J [{ m mag}]$	$16.34\pm0.12$	$15.66\pm0.07$	$13.61\pm0.02$		$15.38\pm0.05$	$15.49\pm0.05$		$20.91\pm0.07$	$16.68\pm0.01$	$18.43\pm0.26$	>18.652	$18.44\pm0.18$	$17.26\!\pm\!0.12$	$19.74\pm0.05$	$20.25 \!\pm\! 0.13$	>19.031	>18.170	$20.13\pm0.08$	>18.734	$20.99\pm0.03$	$19.47\pm0.08$	$16.48\pm0.07$	>18.772	$22.58\pm0.14$	>18.444	$21.06\pm0.06$	$20.57\pm0.05$	$17.94\pm0.14$	>18.699	$22.45\pm0.07$	>18.263
Dec. (J2000)	+01 16 13.09	-43 10 26.27	-61 55 15.82	-02 36 26.00	-24 53 29.80	$+00\ 29\ 15.82$	-04 03 08.90	-02 45 11.80	$+12\ 21\ 14.72$	-50 44 03.00	+43 10 44.70	-64 20 30.00	-69 31 21.60	+36 07 23.57	$+22\ 30\ 05.20$	-54 01 54.80	$+15\ 02\ 47.90$	-71 57 43.80	+84 01 10.60	-22 50 24.99	$+27\ 32\ 59.10$	$+14\ 59\ 53.22$	-36 28 17.50	$+27\ 11\ 44.00$	$+28\ 05\ 48.56$	$+55\ 34\ 21.40$	-68 47 38.60	-75 00 24.60	-56 58 30.50	-62 32 35.40	+02 40 14.10
R. A. (J2000)	11 10 09.99	22 28 28.89	08 17 30.01	05 38 10.10	02 43 13.72	16 24 14.37	21 49 47.20	05 38 14.19	13 00 41.94	$03\ 25\ 04.52$	03 35 15.07	04 04 43.50	22 12 16.27	09 43 05.99	15 42 14.00	03 59 34.07	$04\ 10\ 22.75$	07 34 44.03	12 06 04.38	$15\ 41\ 51.66$	17 38 35.53	20 56 28.92	22 20 55.32	22 09 05.73	08 25 07.36	14 05 18.39	16 39 40.83	$05\ 35\ 16.87$	035000.31	06 47 23.24	23 54 02.79
$\mathrm{SpT}^{\mathrm{a}}_{\mathrm{SpeX}}$	T5.5	$^{9}L$	$^{9}L$	T7.5	$^{9}L$	$^{16}$		T7.5	T7.5																						
$^{\mathrm{SpT}}$	T5.5	9L	9L	T4.5	$^{16}$	9L	$^{17}$	$^{17}$	8L	AT	6L	6L	6L	T9.5	T9.5	$_{ m A0}$	$_{ m A0}$	$_{\rm Y0}$	$_{\rm Y0}$	$_{\rm Y0}$	$^{\mathrm{Y0}}$	$V_0$	$V_0$	Y0:	Y0.5	Y0.5	Y0pec	$Y_1$	$Y_1$	Y1	Y1
Name	2MASS J1110100+0116130	2MASS J2228288-4310262	2MASS J0817300–6155158	S Ori 73	2MASS J0243137-245329	2MASS J1624143+0029158	${\rm CFBDSIR}{2149-0403}$	S Ori 70	ROSS458C	WISEA $J032504.5-504403.0$	WISEA J033515.1+431044.7	WISEA J040443.5-642030.0	WISEA J221216.3–693121.6	WISEA J094306.0+360723.3	WISEA J154214.0+223005.2	WISEA J035934.1–540154.8	WISEA J041022.8+150247.9	WISEA J073444.0-715743.8	WISEA J120604.3+840110.5	WISE J154151.7 $-225024.9$	WISEA J173835.5+273258.8	WISEA J205628.9+145953.6	WISEA J222055.3–362817.5	WISEA J220905.8+271143.6	WISEA J082507.4+280548.2	WISEA J140518.3+553421.3	WISEA J163940.8-684739.4	WISEA $J053516.9-750024.6$	WISEA J035000.3-565830.5	WISEA $J064723.2-623235.4$	WISEA J235402.8+024014.1
Num.	36	37	38	39	40	41	42	43	44	45	46	47	48	49	50	51	52	53	54	55	26	57	28	59	09	61	62	63	64	65	99

Table 1 continued

# Table 1 (continued)

Num.

Ref.	
HST GO°	
$SNR^b$	al. (2015) 2015), [8] 016), [16] al. (2013) 2013), [27 016), [32]
$\pi_T \; [{ m mas}] \qquad { m SNR}^{ m b} \;\; { m HST \; GO}^{ m c}$	i, [2] - Yang et amirez et al. (? (2012). [Mann et al. (2) s), [21] - Liu et s), [21] - Liu et am & Esplin (2)
$K_s [{ m mag}]$	zii et al. (2014) 6), [7] - Peña R 6), [7] - Peña R 013), [15] - Sah 7 (2018) 1013), [26] - Wei 1013), [26] - Wei 3), [31] - Luhm
H [mag]	udy, [1] - Buer ang et al. (2013), [12], aberty et al. (2013), [12], aberty et al. (2012), [20] - Warocco et al. (2013 rash et al. (2013 r
$J [{ m mag}]$	: TS - This st [2015), [6] - Y [11] - Apai et a 2011), [14] - F 2011), [14] - F 2010, [25] - M. 3017), [30] - Mi 3018)
Dec. (J2000)	nfrared spectra: Buenzii et al., et al., (2018), [ Colló,, [19] - [ Dahn et al. (2 lorme et al. (2 Martin et al. (2)
SpT Sp $\mathrm{T_{SpeX}^a}$ R. A. (J2000) Dec. (J2000) J [mag] H [mag] $K_s$ [mag]	OTE—References for first publication of HST/WFC3 near infrared spectra: TS - This study, [1] - Buenzli et al. (2014), [2] - Yang et al. (2015), [4] - Manjavacas et al. (2018), [5] - Buenzli et al. (2015), [6] - Yang et al. (2016), [7] - Peña Ramírez et al. (2015), [8] - Schneider et al. (2015), [9] - Zhou et al. (2018), [10] - Biller et al. (2018), [11] - Appi et al. (2013), [12] - Buenzli et al. (2012). [7] - Kripa et al. (2014), [17] - Gizis et al. (2015), [18] - Lin et al. (2011), [14] - Faherty et al. (2013), [15] - Sahlmann et al. (2016), [16] - Kripa et al. (2012), [23] - (Bedin et al. 2017), [24] - Dahn et al. (2012), [25] - Marocco et al. (2013), [26] - Weinberger et al. (2013), [27] - Linney et al. (2014), [28] - Tinney et al. (2013), [29] - Delorme et al. (2017), [30] - Marsh et al. (2013), [31] - Luhman & Esplin (2016), [32] - Kirkpatrick et al. (2011), [33] - Leggett et al. (2017), [34] - Martin et al. (2018)
$\mathrm{SpT}^{\mathrm{a}}_{\mathrm{SpeX}}$	hication of HS Manjavacas et Zhou et al. (20 parallax measi s et al. (2015) 23] - (Bedin et Tinney et al.
$_{\mathrm{P}}$	s for first put (2016), [4] - (2015), [9] - rigonometric 4), [17] - Gizi Liu (2012), [28] - 1. (2014), [28] - 1. (2011), [33]
Name	Nore—References for first publication of HST/WFC3 near infrared spectra: TS - This study, [1] - Buenzli et al. (2014), [2] - Yang et al. (2015), [4] - Manjavacas et al. (2018), [5] - Buenzli et al. (2015), [6] - Yang et al. (2016), [7] - Peña Ramírez et al. (2015), [8] - Schneider et al. (2015), [9] - Zhou et al. (2018), [10] - Biller et al. (2018), [11] - Appi et al. (2013), [12] - Buenzli et al. (2012). References for trigonometric parallax measurements: [13] - Andrei et al. (2011), [14] - Faherty et al. (2013), [15] - Sahlmann et al. (2016), [16] - Vrba et al. (2014), [17] - Gizis et al. (2015), [18] - Lin et al. (2011), [14] - Faherty et al. (2012), [20] - Wang et al. (2018), [21] - Lin et al. (2013), [22] - Dahne et al. (2002), [25] - Marocco et al. (2013), [26] - Weinberger et al. (2013), [27] - Delorme et al. (2017), [30] - Marsh et al. (2013), [31] - Luhman & Esplin (2016), [32] - Kirkpatrick et al. (2011), [33] - Leggett et al. (2017), [34] - Martin et al. (2018)

 $<sup>^</sup>a\mathrm{SpT}^a_\mathrm{SpeX}$  is the spectral type of each the spectra compiled in this work derived by comparison to the SpeX Prism Spectral Library

 $<sup>^</sup>b\,\mathrm{SNR}$  measured at 1.25  $\mu\mathrm{m}$ 

<sup>&</sup>lt;sup>c</sup> Further details of the program-specific observation plan can be found at: http://www.stsci.edu/cgi-bin/get-proposal-info?id=######submit=Go&observatory=HST, where ##### should be replaced by the given GO program number.

# 3. SPECTRAL TYPES

The spectral types provided in column 3 of Table 1 for brown dwarfs and substellar companions are those given in the literature from the different sources. To provide a homogeneous spectral type classification, we compared our HST/WFC3 spectra to the spectra in the SpeX Prism Spectral Library. We compared our HST/WFC3 spectra using a modified  $\chi^2$  metric as presented in Cushing et al. (2008):

$$G = \sum_{\lambda} \left[ \frac{C(\lambda) - \alpha T(\lambda)}{\sigma_c(\lambda)} \right]^2, \tag{1}$$

where  $C(\lambda)$  is the spectrum of our object,  $T(\lambda)$  is the comparison spectrum,  $\alpha$  is a scaling factor that minimizes G, and  $\sigma_c(\lambda)$  are the uncertainties of the spectrum. We additionally checked the best spectral matches by visual inspection.

In column 4 of Table 1 we show the resulting spectral types for each object. We found that the spectral types derived using the SpeX library are consistent with those published for each object in the literature, matching typically to within +/- 1.5 spectral sub-types. The only exceptions are some of the known intermediate- or low-surface gravity objects in our sample. For these objects (CD-352722, LP 261-75b, 2MASS J2224438-015852, 2MASS J0310559+164816, S Ori 73) the difference in spectral types with respect to the literature values can be up to  $\pm 3$  spectral sub-types. These differences are expected, as the SpeX spectral library is mostly composed of field, i.e., high-surface gravity brown dwarfs.

# 4. COLOR-MAGNITUDE DIAGRAMS

We use near-infrared color-magnitude diagrams (CMDs) for a simple, yet quantitative comparisons of the L and T brown dwarfs in this study to those published in Dupuy & Liu (2012), with the aim of identifying peculiar objects (extremely red or extremely blue dwarfs), or multiple systems. In Figure 5 we show the CMD plot, 2MASS J-H color versus J-band absolute magnitude, for L and T brown dwarfs of Table 1 and objects from Dupuy & Liu (2012) as a comparison. We calculated the absolute J-band magnitude using trigonometric parallaxes, when available. Black stars represent objects from our sample listed in Table 1 with trigonometric parallaxes available in the literature. We do not include Y dwarfs in this CMD, as there are few other Y dwarfs for comparison. Red dots represent L dwarfs, green dots represent L-T brown dwarfs, and blue dots represent T dwarfs with trigonometric parallaxes published in Dupuy & Liu (2012). The solid grey line represents the color-absolute magnitude relationship for

brown dwarfs, and the dotted grey line represents the rms (root-mean-square) of that relation.

The object 2MASS J00470038+6803543 (hereafter W0047, L7, object 9) stands out outside the rms of the color-absolute magnitude relation with red J-H color index. Objects 2MASS J17503293+1759042 (hereafter 2M1750+1759, object 32) and 2MASS J05591914-1404488 (hereafter 2M0559-1404, object 34) are overluminous with respect to the other L-T transition objects, as they are above the rms of the color-magnitude relation for brown dwarfs (Dupuy & Liu 2012). The cause of their overluminosity is unknown, as no multiplicity has been reported previously for these objects.

### 5. LOW SURFACE GRAVITY OBJECTS

# 5.1. Gravity Index Determination

We use the low surface gravity indicators presented in Allers & Liu (2013) that are applicable to our sample up to spectral type L7. We aim to search for as-of-yetunidentified low surface gravity objects and to confirm those that have been classified as low surface gravity objects previously. Due to the spectral wavelength coverage of our HST/WFC3 spectra and their resolution, only the H- continuum and the  $KI_{\rm J}$  indices are applicable to our spectra. The H-continuum measures the shape of the H-band, that has been found to be triangular for most of the very low-gravity brown dwarfs (<100 Myrs) with spectral types between M6 and L7. Intermediate and field gravity brown dwarfs show a "shoulder" at 1.57  $\mu$ m indicative of the appearance of the FeH molecular absorption and the collisionally-induced absorption by H<sub>2</sub> molecules (Borysow et al. 1997; Allard et al. 2012). Allers & Liu (2013) warned, though, that this index needs to be used in combination with others, as the H-band is also triangular for some objects that do not have low surface gravity. The KI<sub>J</sub> index measures the alkali line doublet depth at 1.244 and 1.253  $\mu m$ . These have been found to be weaker for low surface gravity objects of spectral types between M5 to L7. The continuum, center, and bandwidth of these indices are described in Table 4. Their values are calculated using equation 1 from Allers & Liu (2013). Different values for the KI<sub>I</sub> and H-continuum indices correspond to different gravity scores: 0, 1 and 2, corresponding to field gravity (FLD-G), intermediate (INT-G), or low surface gravity objects (VL-G), respectively. The ranges of values of the KI<sub>J</sub> and the H-continuum indices that correspond to different gravity scores are given in Table 9 of Allers & Liu (2013). The  $KI_J$  and the H-continuum indices and their gravity scores obtained for our objects are shown in Table 5. In Figure 6 we show the spectral type versus the  $KI_J$  and the H-continuum indices for

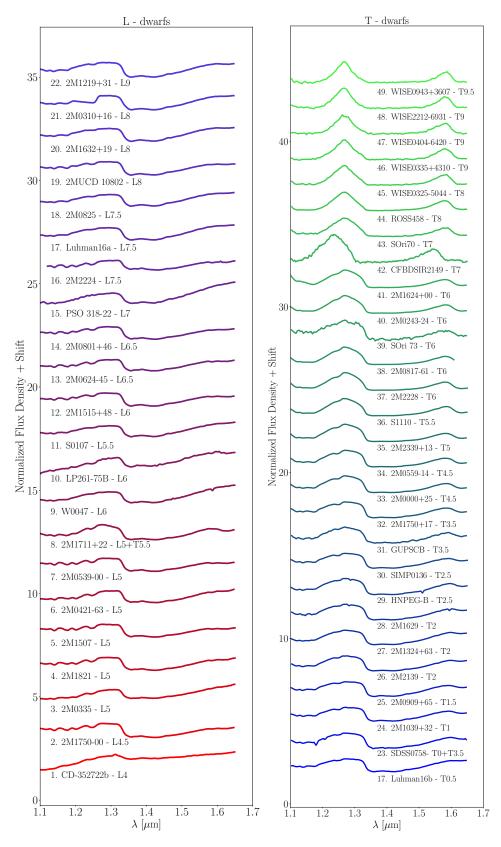


Figure 2. L and T dwarfs with HST/WFC3 spectra used as a part of this study. Flux is normalized at 1.25  $\mu$ m.

Num.	Name	RA	DEC	SpT host	$\pi_T^a \text{ [mas]}$	Separation [au]	Planet radius [R <sub>Jup</sub> ]	$SNR^b$	HST Prog.	Ref.
67	WASP-18b	01 37 25.03	-45 40 40.39	F6V	$7.91 \pm 0.30$	$0.02014 \pm 0.00034$	$1.165 {\pm} 0.077$	37.8	13467	9
68	WASP-33b	$02\ 26\ 51.06$	$+37\ 33\ 01.73$	A5	$8.51 {\pm} 0.24$	$0.02555 {\pm} 0.00043$	$1.497 {\pm} 0.045$	38.9	12495	2
69	$\mathrm{HD}\ 209458\mathrm{B}$	$22\ 03\ 10.77$	$+18\ 53\ 03.54$	G0V	$20.47 {\pm} 0.02$	$0.04723 \pm 0.00079$	$1.359 \pm 0.015$	6.4	13467	1
70	WASP-12b	$06\ 30\ 32.79$	$+29\ 40\ 20.29$	G0V	$2.57{\pm}0.27$	$0.02253 {\pm} 0.00038$	$1.790 \pm 0.090$	12.3	13467	4
71	WASP-121b	$07\ 10\ 24.06$	-39 05 50.55	F6V	$3.82 {\pm} 0.25$	$0.02544 {\pm} 0.00050$	$1.865 {\pm} 0.044$	13.3	14767	8
72	WASP- $43b$	10 19 38.01	-09 48 22.59	K7V	$11.49 \pm 0.04$	$0.01424 {\pm} 0.00041$	$0.930 \pm 0.070$	12.5	13467	7
73	WASP-103b	$16\ 37\ 15.57$	$+07\ 11\ 00.07$	F8V	$2.13 \pm 0.16$	$0.01987 \pm 0.00033$	$1.528 {\pm} 0.073$	8.3	14050	3
74	TrES-3b	$17\ 52\ 07.02$	$+37\ 32\ 46.18$	K0V	$4.29 \pm 0.02$	$0.02272 \pm 0.00038$	$1.336 \pm 0.031$	2.5	12181	5
75	Kepler-13Ab	19 07 53.15	$+46\ 52\ 05.91$	A0	$1.91 \pm 0.01$	$0.04171 \pm 0.00078$	$1.406 {\pm} 0.038$	10.8	13308	6
76	WASP-4b	23 34 15 08	-42 03 41 14	G7V	$3.63 \pm 0.70$	0.02304+0.00042	1.341+0.023	4.4	12181	5

Table 2. Sample of hot Jupiters with HST/WFC3 spectroscopic data

Note—[1] - Line et al. (2016), [2] - Haynes et al. (2015), [3] - Cartier et al. (2017), [4] - Stevenson et al. (2014a), [5] - Ranjan et al. (2014), [6] - Beatty et al. (2017), [7] - Stevenson et al. (2014b), [8] - Evans et al. (2017), [9] - Sheppard et al. (2017)

**Table 3.** Age and mass estimates for the planetary-mass objects of our sample.

Num.	Name	$\operatorname{SpT}$	$Mass (M_{Jup})$	Age (Myr)	YMG	Ref.
3	2MASS J03552337+1133437	L5.0	13-30	~120	AB-Doradus	1
9	2MASS J00470038+6803543	L6.0	$20^{+3}_{-7}$	$\sim 120$	AB-Doradus	2
15	PSO J318.5-22	L7.0	$8.3 {\pm} 0.5$	$23\pm3$	$\beta ext{-Pic}$	3, 4
30	SIMP J013656.5+093347.3	T2.5	$12.7{\pm}1.0$	$\sim 120$	AB-Doradus	5
31	GU PSC B	T3.5	$11.9^{+2}_{-1.5}$	$\sim 120$	AB-Doradus	2
36	2MASS J11101001+0116130	T5.5	10 – 12	$\sim 120$	AB-Doradus	6
42	CFBDSIR2149-0403	T7.0	2-13	< 500	None	7
44	ROSS458C	T8.0	5-20	<1000	None	8

NOTE—References: [1] - Faherty et al. (2013), [2] - Aller et al. (2016), [3] - Liu et al. (2013), [4] - Allers et al. (2016) [5] - Gagné et al. (2017), [6] - Gagné et al. (2015a), [7] - Delorme et al. (2017), [8] - Burningham et al. (2011).

 $<sup>^</sup>a\operatorname{Trigonometric}$  parallaxes delivered by Gaia Collaboration et al. (2018)

 $<sup>^</sup>b$  The SNR is measured between 1.05 and 1.65  $\mu\mathrm{m}$ 

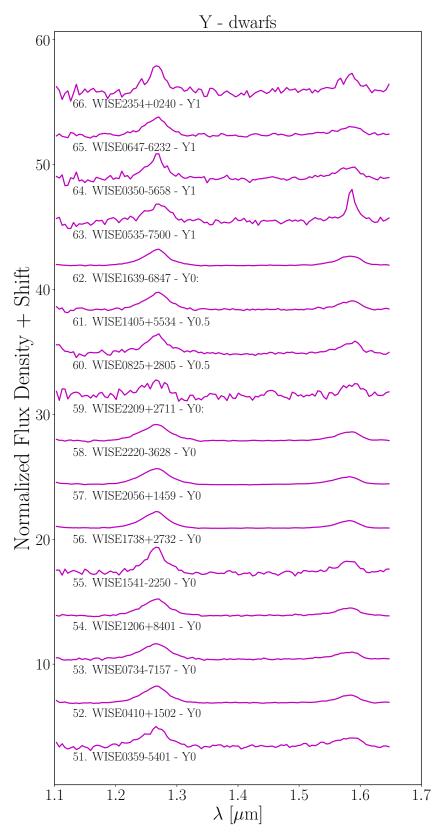


Figure 3. Y dwarfs with HST/WFC3 spectra used as a part of this study. Flux is normalized at 1.25  $\mu$ m.

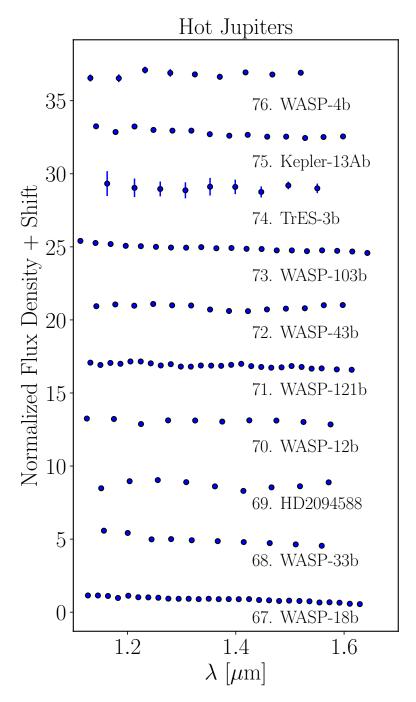


Figure 4. Hot Jupiters with HST/WFC3 emission spectra used as a part of this study. Some error bars are smaller that the symbol for the measurement.

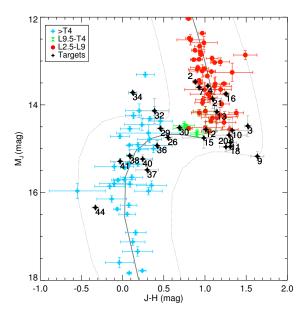


Figure 5. 2MASS J-H color versus  $M_J$  2MASS absolute magnitude diagram. The black stars correspond to L and T dwarfs of our sample with measured trigonometric parallaxes published in the literature. Red dots are L dwarfs, green dots are L-T transitions dwarfs and blue dots are T dwarfs with parallaxes published in Dupuy & Liu (2012). The solid grey lines represent the color-absolute magnitude relationships for brown dwarfs. The dotted grey lines represent the rms of that relation. Objects lying outside the rms of the relation, are catalogued extremely red or blue, or overluminous.

our sample (up to L7 spectral type) and for objects that belong to young moving groups,  $\gamma$  and  $\beta$  dwarfs<sup>3</sup>, young companions (Allers & Liu 2013; Bonnefoy et al. 2014), and field brown dwarfs (McLean et al. 2003; Cushing et al. 2005) for comparison.

### 5.2. Results: gravity class determination

Among the 16 objects in our sample with spectral types up to L7 for which the  $\mathrm{KI_J}$  and the H-continuum indices could be measured, two have gravity scores consistent with very low surface gravities, namely: CD-352722B (object 1), and 2MASS J03552337+1133437 (object 3). Five objects had gravity scores consistent with intermediate gravities, namely: 2MASSI J0421072-630602 (object 6), 2MASS J00470038+6803543 (object 9), 2MASS J01075242+0041563 (object 11), PSO J318.5-22 (object 15), and 2MASSW J2224438-015852 (object 16).

CD-352722B and 2MASS J03552337+1133437 had been reported before as low surface gravity objects

**Table 4.** Spectral indices to segregate young brown dwarfs from Allers & Liu (2013).

Index	$\lambda_{line} \; (\mu \mathrm{m})$	$\lambda_{cont1} \; (\mu \mathrm{m})$	$\lambda_{cont2} \; (\mu \mathrm{m})$	Width in $\lambda$
$\mathrm{KI_{J}}$	1.244	1.220	1.270	0.0166
$H\text{-}\mathrm{cont}$	1.560	1.470	1.670	0.0208

(Wahhaj et al. 2011 and Faherty et al. 2013, respectively), and they are also high-probability members of the AB-Doradus moving group with an estimated age of  $\sim$ 120 Myr (Zuckerman et al. 2004).

Object 2MASSI J0421072-630602 (object 6) was classified as intermediate-surface gravity (L5 $\beta$ ) object by Cruz et al. (2009) in optical wavelengths.

2MASS J00470038+6803543 (object 9) is an extremely red L-dwarf, and a bona fide member of the AB Doradus moving group, for which intermediate surface gravity characteristics have been previously reported (Gizis et al. 2012; Allers & Liu 2013).

PSO J318.5-22 is also an extremely red L-dwarf for which low surface gravity signatures have been found Liu et al. (2013) and is a bona fide member of the  $\beta$ -Pictoris moving group (age = 23±3 Myr, Zuckerman et al. 2001; Mamajek & Bell 2014).

Object 2MASS J01075242+0041563 (object 11) does not clearly show low surface-gravity spectral characteristics (Gagné et al. 2015b), and it was found to be a possible member of the Hyades association (age  $\sim$ 625 Myr, Bannister & Jameson 2007). If confirmed, its age would not be consistent with its intermediate gravity classification (expected for objects with ages between 50 and 200 Myr, see Allers & Liu 2013).

Finally, 2MASSW J2224438-015852 (object 16) is an extremely red L4.5 dwarf, that was classified as a field dwarf by Liu et al. (2016) using the BANYAN II tool. Martin et al. (1996) classified it as a field gravity object using NIRSPEC/Keck II spectra to obtain Allers & Liu (2013) indices in the J-band.

# 6. CANDIDATES FOR COMPOSITE ATMOSPHERES

Our high-quality spectra are well-suited for exploring the diversity of ultracool atmospheres, including the identification of potentially composite (multicomponent) spectra. An obvious source of such composite spectra are unresolved binaries with different spectral types. However, given the very high occurrence rate of heterogeneous cloud cover in brown dwarfs (Buenzli et al. 2014; Metchev et al. 2015), it is expected that brown dwarfs with strong spectral heterogeneity

<sup>&</sup>lt;sup>3</sup> Optical classification for very low and intermediate surface gravity, respectively (Cruz et al. 2009).

Object	Name	H-continuum	$\mathrm{KI_{J}}$	Gravity score per index	Gravity class
1	CD-352722B	$1.028 \pm 0.001$	$1.013 \pm 0.001$	2 2	VL-G
2	2MASS J17502484-0016151	$0.829 {\pm} 0.001$	$1.106 \pm 0.001$	0 0	FLD-G
3	$2 {\rm MASS}\ J03552337{+}1133437$	$0.982 {\pm} 0.001$	$1.046 {\pm} 0.001$	2 2	VL-G
4	2MASS J18212815+1414010	$0.867{\pm}0.001$	$1.101 {\pm} 0.001$	0 0	FLD-G
5	2MASSW J1507476-162738	$0.829 {\pm} 0.001$	$1.109 \pm 0.001$	0 0	FLD-G
6	2MASSI J0421072-630602	$0.892 {\pm} 0.001$	$1.070\!\pm\!0.001$	1 1	INT-G
7	$2 {\rm MASS}~ J05395200\text{-}0059019$	$0.806 \pm 0.001$	$1.109 \pm 0.001$	0 0	FLD-G
8	2MASSI J1711457+223204	$0.734 {\pm} 0.001$	$1.052 {\pm} 0.001$	0 –	FLD-G
9	$2 {\rm MASS\ J} 00470038{+}6803543$	$0.934 {\pm} 0.001$	$1.006 \pm 0.001$	1 –	INT-G
10	LP261-75B	$0.713 \pm 0.004$	$1.076 \!\pm\! 0.001$	0 1	FLD-G
11	$2 {\rm MASS}\ J01075242{+}0041563$	$0.883 {\pm} 0.001$	$1.015{\pm}0.001$	1 2	INT-G
12	2MASSW J1515008+484742	$0.809 \pm 0.001$	$1.072 {\pm} 0.001$	0 –	FLD-G
13	2MASS J06244595-4521548	$0.839 {\pm} 0.001$	$1.086{\pm}0.001$	0 1	FLD-G
14	2MASSW J0801405+462850	$0.869 {\pm} 0.001$	$1.059 \pm 0.001$	0 –	FLD-G
15	PSO J318.5-22	$0.918 {\pm} 0.002$	$0.982 {\pm} 0.001$	1 -	INT-G

 $1.084 \pm 0.001$ 

 $0.898 \pm 0.001$ 

Table 5. Values obtained for the H-continuum and KI<sub>J</sub> indeces for our objects with their corresponding gravity scores.

will also contribute to the population of composite atmospheres.

2MASSW J2224438-015852

# 6.1. Search for composite atmosphere brown dwarfs

To identify potential composite spectra candidates in our sample, we used the spectral indices from Burgasser et al. (2006b, 2010) and Bardalez Gagliuffi et al. (2014). These indices examine peculiar spectral characteristics of spectroscopic L-plus-T and M-plus-T composite atmosphere brown dwarfs. L plus T composite spectra have bluer SEDs (spectral energy distributions) in the near-infrared than field objects of similar spectral type. In cases of L-plus-T spectroscopic binaries atomic and molecular features are blended, affecting the H<sub>2</sub>O  $(1.15 \mu m)$  and CH<sub>4</sub>  $(1.32 \mu m)$  molecular features. At 1.55  $\mu$ m spectroscopic L+T binaries show larger flux from the T dwarf. Burgasser et al. (2006b, 2010) and Bardalez Gagliuffi et al. (2014) combine different pairs of indices, defining those that segregate brown dwarfs with composite spectra more efficiently.

The criteria to select potential brown dwarf composite atmospheres, and the delimiters of the areas within the plots comparing different indices that segregate spectroscopic brown dwarfs with composite spectra are defined in Tables A3, A4 and A5 of the Appendix. We do not measure indices that are outside the wavelength range of the HST/WFC3 G141 near-infrared spectra. Thus, only 13 out of 18 available plots comparing indices from Bur-

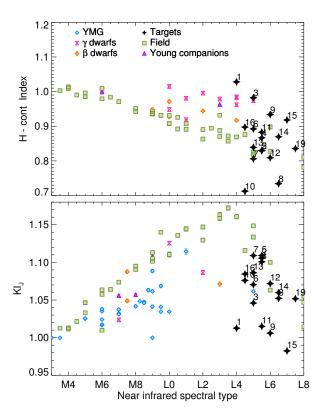
gasser et al. (2006b, 2010) and Bardalez Gagliuffi et al. (2014) are applicable to in our data (see Tables A4 and A5). In Figures A2, A3 and A4 of the Appendix we show the comparison of the spectral indices listed in Tables A4 and A5 of the Appendix. To be considered a weak candidate to have a composite brown dwarf atmosphere, Burgasser et al. (2006b, 2010), and Bardalez Gagliuffi et al. (2014) established that an object needs to appear within the selection area of minimum four plots, and to be considered a strong candidate, it needs to appear at least in 8 plots. Although due to the wavelength coverage of the HST/WFC3 data we were not able to perform 5 of the index comparisons in Burgasser et al. (2006b, 2010), and Bardalez Gagliuffi et al. (2014), we conservatively use the same criteria to select weak and strong spectroscopic binaries.

1 1

INT-G

### 6.2. Results: composite atmosphere candidates

We found that ten objects in our sample were selected by the indices as weak composite spectra candidates: 2M0310+16 (object 21), 2M0624-45 (object 13), 2MUCD 10802 (object 19), 2M0909+65 (object 25), 2M1039+32 (object 24), 2M1324+63 (object 27), 2M1515+48 (object 12), 2M1632+19 (object 20), 2M1711+22 (object 8), 2M1750-00 (object 2). In addition, we found 3 strong candidates to have composite spectra: 2M1507-16 (object 5), 2M1219+31 (object 22), and SIMP0136+09 (object 30).



**Figure 6.** Spectral types versus H-continuum and KI<sub>J</sub> indices from Allers & Liu (2013). For comparison, we show the value of these indices objects of the field (green squares), young companions (purple triangles),  $\beta$  dwarfs (intermediate surface gravity, orange crosses),  $\gamma$  dwarfs (very low gravity, pink hourglass symbol) and young moving group members (blue diamonds). Black stars belong to objects of the sample presented in this work with spectral types between L3 and L7. Numbers identify the objects from Table 1.

To confirm or reject the candidates, we compared our HST/WFC3 spectra to single template spectra from the SpeX Prism Spectral Library, and independently, to synthetic L plus T composite spectra created using single spectra from those libraries (see Table 6). To create the synthetic composite spectra, we scaled the fluxes of the components to 10 pc using the color-magnitude relation from Dupuy & Liu (2012), and coadded the two component fluxes.

Following Burgasser et al. (2006b, 2010) and Bardalez Gagliuffi et al. (2014), we compared the goodness of the fit of the HST/WFC3 spectra to the single, and independently to synthetic L plus T composite spectra using a modified  $\chi^2$  (G) using equation 1 from Cushing et al. (2008).

Finally, we tested if the fit to a composite template spectra was significantly better than the fit to a single template using one-sided F-test statistic. The distribution statistic ratio we used:

$$\eta_{SB} = \frac{min(G_{single}) \times df_{composite}}{min(G_{composite}) \times df_{single}}$$
(2)

where  $\min(G_{single})$  and  $\min(G_{composite})$  are the minimum G for the best match to a single or a composite template, and  $df_{composite}$  and  $df_{single}$  are the degrees of freedom for the composite template spectra fit and the single template fit. The degrees of freedom are the number of data points used in the fit minus one, to account the scaling between our spectra and the template spectra. To rule out the null hypothesis, meaning that the candidate spectrum is not best described by a single template at the 99% confidence level, we require  $\eta_{SB} > 1.41$ . The F-statistic analysis rejected five of our candidates: 2M0310+16 (object 21), 2M0909+65 (object 25), 2M1750-00 (object 2), 2M1507-16 (object 5) and SIMP0136+09 (object 30).

Unresolved binaries should appear overluminous compared to single brown dwarfs. To test for evidence of overluminosity in our targets, we compared the absolute magnitudes, derived using the trigonometric parallaxes for our targets, with the spectro-photometric absolute magnitudes derived using the relation published by Dupuy & Liu (2012) (see Table 7). To obtain spectro-photometric absolute magnitudes we used the spectral type of the principal component. We found that none of the sources with trigonometric parallaxes are overluminous. Actually, we find that some of them are slightly underluminous. This fact does not support the multiplicity hypothesis.

# 6.3. Rotational modulation and composite atmosphere candidates

We also searched for published rotational modulation detections for our candidates as a potential marker for composite atmospheres. We found that nine of the thirteen composite atmosphere candidates have reported photometric variability due to cloud patterns (see Table 6). In fact, from the eight objects that satisfied the criteria for composite spectra candidates, six are known to have rotational modulations: 2M0624-45 (object 13), 2MUCD 10802 (object 19), 2M1039+32 (object 24), 2M1324+63 (object 27), 2M1219+31 (object 22), and 2M1632+19 (object 20). One is a confirmed binary: 2M1711+22 (object 8), and one has been reported as non-variable (2M1515+48, object 12).

These results suggest that Burgasser et al. (2006b, 2010), and Bardalez Gagliuffi et al. (2014) spectral indices are biased towards L and T brown dwarfs that show photometric variability due to rotational modu-

0.95

Yes [3]

Num.	Name	SpT	Single Component	Binary Components	$\eta_{SB}$	Variable?
21	2M0310+16	L8	L8 (SD J121951.45+312849.4)	(L8) SD J085758.45 $+$ 570851.4 + T0 (SD J042348.57-041403.5)	1.27	Yes [1]
13	2M0624-45	L6.5	L5 (2M J23512200+3010540)	${\rm L3.5~(2M~J2224438\text{-}015852)+T1~(SD~J163239.34\text{+}415004.3)}$	2.18	Yes [1]
19	$2 \mathrm{MUCD}\ 10802$	L7.5	L7.5 (SD J115553.86+055957.5)	L7.5  (SD J115553.86+055957.5) + T0  (Gl337CD)	1.43	No [1]
25	2M0909+65	T1.5	T2 (2M J11220826-3512363)	L9.5 (SD J082030.12+103737.0) + T3.5 (SD J175032.96+175903.9)	1.03	No [1]
24	2M1039+32	T1	T1.5 (SD J090900.73+652527.2)	$L8 \; (SD \; J121951.45 + 312849.4) \; + \; T4 \; (2M \; J2254188 + 312349)$	3.25	Yes [1]
22	2M1219+31	L8	L8 (2M J0328426+230205)	${\rm L9~(2M~J0310599+164816)+T0~(Gl~337CD)}$	2.07	Yes [1]
27	2M1324+63	T2	T2 (SD J125453.90-012247.4)	L9 (SD J083008.12+482847.4) + T7.5 (2M J11145133-2618235)	3.94	Yes [1]
12	2M1515+48	L6	L9 (2M J0908380+503208)	$ L9 \ (2M \ J0908380 + 503208) \ + \ T0 \ (SD \ J152039.82 + 354619.8) $	1.69	No [1]
20	2M1632+19	L8	L8 (Gl584C)	L6(2M J0825196+211552)+T0(SD J204749.61-071818.3)	2.71	Yes [1]
8	2M1711+22	L5.0+T5.5	T1 (SD J085834.42+325627.7)	L6(2M J0825196+211552)+T3(SD J102109-030420)	2.30	Binary
2	2M1750-00	L5.5	L5 (2M J18131803+5101246)	${\rm L4.5~(2M~J15200224\text{-}4422419B)} + {\rm T0.5~(SD~J151643.01+305344.4)}$	0.79	Yes[1]
5	2M1507-16	L5.5	L5 (2M J17461199 + 5034036)	${\rm L5~(2M~J10461875+4441149)}  +  {\rm T0~(2M~J0920122+351742)}$	1.18	Yes [2]

 $T2 \ (2MASS \ J11220826-3512363) \quad SD \ J213154.43-011939.3 \ (L9) + SD \ J092615.38+584720.9 \ (T4.5)$ 

Table 6. Candidates for composite spectra selected by Burgasser et al. (2006b, 2010) and Bardalez Gagliuffi et al. (2014) spectral indices.

Note—[1] - Buenzli et al. (2014), [2] - Yang et al. (2015), [3] - Artigau et al. (2009).

30

SIMP0136+09

T2.5

lations. Thus, these indices should also be useful to search for brown dwarfs candidates with heterogeneous cloud patterns in their atmospheres. In Table 8, we list all objects with L4 to T4 spectral types in our sample for which the method presented in this Section are applicable. We specify which of them have reported rotational modulations in the literature, and we compare with those that have been found by the indices as composite spectra candidates.

We conclude that 21 out of the 32 objects listed in Table 7 have rotational modulations reported in the literature (see Table 8 for details); nonetheless, only nine are detected by the indices as candidates for composite spectra, with spectral types from L4 to T2 (spectral types from the literature). None of the low-gravity brown dwarfs found in Section 5 with reported rotational modulations have been detected by indices. In addition, the indices have detected three other candidates to have composite spectra, but are not known to show rotational modulations.

# 7. METHANE (1.2 $\mu$ m) AND WATER (1.4 $\mu$ m) ABSORPTION BANDS FOR SPECTRAL CLASSIFICATION

The most prominent molecular absorption bands in the near-infrared spectra of brown dwarfs and substellar companions are  $\rm H_2O$ ,  $\rm CO$ , and  $\rm CH_4$ . The absorption bands of the different molecules are controlled by the availability of C and O (Marley et al. 2010). Within the temperature range (approx. 1800 K to 600 K and below) of substellar objects in our sample – corresponding to spectral types from L4 to Y1 – carbon appears mainly in the form of CO in L dwarfs and as  $\rm CH_4$  in T dwarfs

(Marley & Robinson 2015). The equilibrium reaction that takes place is:

$$CO + 3H_2 \leftrightarrow CH_4 + H_2O$$
 (3)

At higher temperature (L dwarfs), the left side of the reaction is favored; thus there is an overabundance of CO, that implies an underabundance of H<sub>2</sub>O. At lower temperatures, below the L/T transition, the right hand-side of the reaction is favored, leading to higher abundances of CH<sub>4</sub> and H<sub>2</sub>O (Marley & Robinson 2015). The depth of some of these molecular bands in our HST/WFC3 spectra can serve to provide a robust spectral classification of substellar objects. In the HST/WFC3 1.10-1.69  $\mu$ m spectra of brown dwarfs and substellar companions, we are able to measure the depth of the CH<sub>4</sub> band at approximately 1.2  $\mu m$  and the H<sub>2</sub>O band at approximately 1.4  $\mu$ m, and trace their change in depth with near-infrared spectral types (see Figures 7 and 8). We measured the depths of the H<sub>2</sub>O and CH<sub>4</sub> bands using equation 1 from Allers & Liu (2013):

$$index = \left(\frac{\lambda_{line} - \lambda_{cont1}}{\lambda_{cont2} - \lambda_{cont1}} F_{cont2} + \frac{\lambda_{cont2} - \lambda_{line}}{\lambda_{cont2} - \lambda_{cont1}} F_{cont1}\right) / F_{line}$$
(4)

In Table 9 we show the wavelengths in which the continuum, the center and the width of the bands are defined. The minimum value of this index is 1, implying that there is no absorption feature. In addition, we derive an exponential function to relate the near-infrared spectral types and the depths of the  $CH_4$ , and the  $H_2O$  bands. To calculate the best-fit exponential function, we used the IDL function COMFIT.PRO, that

Num.	Name	SpT principal component	Variable?	$\pi_{\mathrm{Trig}}$ (mas)	$M_J \text{ (mag)}$	$M_{J,SP}^a$ (mag)
13	2M0624-45	L6.5	Yes [1]	86.21±4.46	$14.16 \pm 0.12$	14.11±0.40
19	2MUCD 10802	L7.5	No [1]			
24	2M1039+32	T1	Yes [1]			
22	2M1219+31	L8	Yes [1]			
27	2M1324+63	T2	Yes [1]			
12	2M1515+48	L6	No [4]	$123.8{\pm}5.0$	$14.56 {\pm} 0.09$	$13.94 \pm 0.40$
20	2M1632+19	L8	Yes [2]	$65.19{\pm}2.16$	$14.96 \pm 0.10$	$14.51 \pm 0.40$
8	2M1711+22	L6.5	Binary [3]	$33.11 \pm 4.71$	$14.69 \pm 0.36$	$14.11 \pm 0.40$

Table 7. Photometric variability reported for final weak and strong candidates for composite spectra.

NOTE—[1] - Buenzli et al. (2014), [2] - Metchev et al. (2015), [3] - Burgasser et al. (2010), [4] - Bardalez Gagliuffi et al. (2014), [5] - Artigau et al. (2009), [6] - Yang et al. (2015).

NOTE—[a] - Absolute magnitude given by the empirical spectro-photometric relation by Dupuy & Liu (2012).

fits an exponential equation of the form:  $y = c_0 * c_1^x + c_2$ using a gradient-expansion least-squares method (Marquardt 1963). The exponential function was preferred over polynomial functions, as it provides a smaller reduced  $\chi^2$ . We did not include Y dwarfs in the fit due to their lower quality data (see Table 1), nor hot Jupiters, as their atmospheres are, in general, physically differ-To obtain the CH<sub>4</sub> index, we discarded as well some companions (CD-352722 B, HN-Peg B, and GU PSC B), as they were outliers in the SpT vs. CH<sub>4</sub> index plot, probably due to contamination of the star at the wavelength range in which the CH<sub>4</sub> index is measured. The best fits to exponential functions are displayed in Figures 7 and 8 with a dashed thick black line. The grey dashed lines in those figures represent the standard deviation of the data points with respect the fitted function. The values of the coefficients for each best fit exponential functions are displayed in Table 10. In addition, in the former Table, we also show the function that provides spectral types of brown dwarfs given the value of the H<sub>2</sub>O and CH<sub>4</sub> bands. In Table 11, we present the typical dimensionless values for the H<sub>2</sub>O and CH<sub>4</sub> bands calculated using the exponential functions indicated in Table 10. The depths of the  $H_2O$  and  $CH_4$  bands can provide robust spectral classification for brown dwarfs with high-quality near-infrared spectra that includes the 1.4  $\mu$ m water band. This is especially true for brown dwarfs with spectral types later than T2, for which the change of the band width is more abrupt with spectral type.

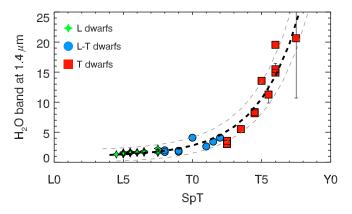


Figure 7. Evolution of the depth of the  $H_2O$  band at 1.4  $\mu m$  with near-infrared spectral types calculated using equation 4. The value of the depth of the  $H_2O$  band is dimensionless.

# 8. COMPARISON OF BROWN DWARFS AND HOT JUPITERS PHOTOMETRY AND SPECTRA

Color-magnitude diagrams have been traditionally used to directly compare the colors and absolute magnitudes of low-mass stars and brown dwarfs (Burgasser et al. 2008; Faherty et al. 2012; Dupuy & Liu 2012, and references therein), revealing that different parameters influence the atmospheres of these objects. Beside effective temperatures, other secondary parameters that influence brown dwarfs colors include surface gravity, metallicity, dust sedimentation, and non-equilibrium chemistry. Brown dwarfs and hot Jupiters share similar effective temperatures and size ranges. Nevertheless, direct comparisons are usually challenging, as hot Jupiters

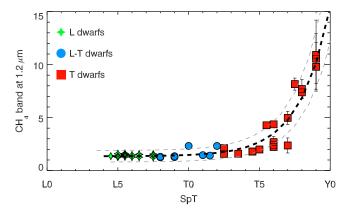
Table 8. L4–T4 dwarfs with reported rotational modulations and composite atmosphere candidates.

Num.	Name	Variable?	Detected by indexes	Confirmed as candidate
1	CD-352722b	_	No	No
2	2MASS J17502484-0016151	Yes [1]	Yes	No
3	2MASS J03552337+1133437	_	No	No
4	2MASS J18212815+1414010	Yes [2]	No	No
5	2MASSW J1507476-162738	Yes [2]	Yes	No
6	$2 {\rm MASSI} \ J0421072\text{-}630602$	No [1]	No	No
7	$2 {\rm MASS} \ J05395200\text{-}0059019$	No [1]	No	No
8	$2 {\rm MASSI~J1711457} + 223204$	No, binary [1]	Yes	Yes
9	$2 {\rm MASS\ J00470038}{+}6803543$	Yes [3]	No	No
10	LP261-75B	Yes [4]	No	No
11	$2 {\rm MASS\ J} 01075242 + 0041563$	Yes [5]	No	No
12	2MASSW J1515008+484742	No [1]	Yes	No
13	$2 {\rm MASS} \ J06244595\text{-}4521548$	Yes [1]	Yes	Yes
14	$2 {\rm MASSW}~ {\rm J}0801405{+}462850$	No [1]	No	No
15	PSO J318.5-22	Yes [6]	No	No
16	$2 {\rm MASSW} \ J2224438\text{-}015852$	No [7]	No	No
17	Luh 16AB	Yes [8]	No	No
18	$2 {\rm MASSI~J0825196} {+} 211552$	No [1]	No	No
19	2MUCD 10802	No [1]	Yes	No
20	$2 {\rm MASS\ J} 16322911{+}1904407$	Yes [1]	Yes	Yes
21	$2 {\rm MASSW}~ {\rm J}0310599{+}164816$	Yes [1]	Yes	No
22	$2 {\rm MASS\ J} 12195156 + 3128497$	Yes [1]	Yes	Yes
23	SDSS J075840.33+324723.4	Yes [13]	No	No
24	2MASS J10393137+3256263	Yes [1]	Yes	Yes
25	$2 {\rm MASS\ J09090085} {+} 6525275$	No [1]	Yes	No
26	2MASS J21392676+0220226	Yes [9,14]	No	No
27	$2 {\rm MASS\ J} 13243553 + 6358281$	Yes [1,14]	Yes	Yes
28	$2 {\rm MASS\ J} 16291840 + 0335371$	Yes [13]	No	No
29	HN PEG B	Yes [10]	No	No
30	SIMP J013656.5+093347.3	Yes [11,14]	No	No
31	GU PSC B	Yes [12]	No	No
32	2MASS J17503293+1759042	Yes [1]	No	No

Note—References: [1] - Buenzli et al. (2014), [2] - Yang et al. (2015), [3] - Lew et al. (2016), [4] - Manjavacas et al. (2018), [5] - Apai et al. in prep, [6] - Biller et al. (2018), [7] - Metchev et al. (2015), [8] - Buenzli et al. (2015), [9] - Radigan et al. (2012), [10] - Zhou et al. (2018), [11] - Artigau et al. (2009), [12] - Naud et al. (2017), [13] - Radigan et al. (2014), [14] - Apai et al. (2017).

**Table 9.** Wavelength for the continuum and central wavelengths in which the depth of the  $CH_4$  and  $H_2O$  bands are measured. Wavelength units are  $\mu m$ .

Index	$\lambda_{line}$	$\lambda_{cont1}$	$\lambda_{cont2}$	Band width
${\rm H_2O}$	1.40	1.31	1.47	0.08
$\mathrm{CH}_4$	1.18	1.10	1.30	0.12



**Figure 8.** Evolution of the depth of the CH<sub>4</sub> band at  $1.2 \mu m$  with near-infrared spectral types calculated using equation 4. The value of the CH<sub>4</sub> band is dimensionless.

orbit at close distance to their host stars, and they are, therefore, highly irradiated and difficult to observe directly.

In this Section, we compare the hot Jupiters HST/WFC3 near-infrared day-side emission photometry and spectra (see Table 2) to similar photometry and spectra of field and young brown dwarfs, to explore differences and similarities between these two classes of substellar objects.

Before proceeding to the spectral comparison, we need to transform the relative flux density typically given for eclipse depth (ED) (a ratio between the flux densities of the planet and the host star,  $F_{\lambda, planet}/F_{\lambda, star}$ ), to absolute (physical) flux density. If the flux of the planet is given in relative eclipse depth, we transform first those units to relative flux given in  $F_{planet}/F_{star}$  by:

$$\frac{F_{\lambda, planet}}{F_{\lambda, planet} + F_{\lambda, star}} = ED \to \frac{F_{\lambda, planet}}{F_{\lambda, star}} = \frac{1}{1/ED - 1}$$
(5)

Once the flux of the planet is given in  $F_{\lambda, planet}/F_{\lambda, star}$ , to transform to actual physical flux density units, we use a model spectrum for the temperature and surface gravity of the spectral type of each the parent stars given in Table 2. We used the BT-Settl atmospheric models (Allard et al. 2012), scaling the model absolute flux by

 $(R/d)^2$ , where R is the radius of the star in m, and d is the distance of the system to Earth. The star radii were obtained from Table 5 from Pecaut & Mamajek (2013). We used trigonometric distances available either from the Gaia DR2 (Gaia Collaboration et al. 2018) or in the literature. Finally, we binned the stellar model spectra to match the corresponding HST/WFC3 spectra bins for each hot Jupiter, and obtained the physical flux density for each planet solving equation 5 as:

$$F_{\lambda, planet} = \frac{F_{\lambda, star}}{1/ED - 1} \tag{6}$$

Once we transformed the units of hot Jupiter spectra to physical units  $(F_{\lambda,\,planet}$  in erg s<sup>-1</sup> cm<sup>-2</sup> Å<sup>-1</sup>), we obtained HST/WFC3 photometric magnitudes in the J (1.10-1.35  $\mu m$ ) and  $H_s$ -band (1.50-1.69  $\mu m$ ) for each hot Jupiter and brown dwarf in our sample. The total flux in each band was calculated by integrating the flux densities in the relevant wavelength ranges. To obtain J and  $H_s$  HST/WFC3 magnitudes as:

$$STmag_{J/H_s} = -2.5 \log F_{\lambda, planet} - ZP_{J/H_s}$$
 (7)

Where  $F_{\lambda, planet}$  is given in erg s<sup>-1</sup> cm<sup>-2</sup> Å, and ZP is the zeropoint in the J or  $H_s$ -band. We use the zero points for the F125W ( $ZP_{F125W} = 25.3293$ , in Vega magnitude) and the F160W filters ( $ZP_{F160W} = 24.6949$ , in Vega magnitude) that are centered at those bands<sup>4</sup>.

To compare the colors of highly-irradiated hot Jupiters and isolated brown dwarfs, we plot a J- $H_s$  HST color vs.  $M_J$  HST magnitude in a CMD diagram with all brown dwarfs with available trigonometric parallax and all hot Jupiters in Figure 8. For comparison, we include M-Y dwarfs presented in Dupuy & Liu (2012) after transforming their 2MASS photometry to HST photometry using the polynomials presented in Section B of the Appendix. Grev squares represent M dwarfs, grey dots are L dwarfs, and grey hourglass symbols correspond to T dwarfs. In addition, we add the targets from our sample with trigonometric parallax: red dots are L dwarfs, green hourglass symbols are T dwarfs, blue crosses are Y dwarfs, black stars are hot Jupiters, and pink stars are hot Jupiters after removing the contribution from the reflected light of the host star. The observed flux from the hot Jupiter is:

$$F_{\lambda, planet} = F_{\lambda, thermal} + F_{\lambda, reflected}$$
 (8)

Where  $F_{\lambda, \text{ planet}}$  is the observed hot Jupiter spectra,  $F_{\lambda, thermal}$  is the thermal flux from the planet, and

<sup>&</sup>lt;sup>4</sup> http://www.stsci.edu/hst/wfc3/ir\_phot\_zpt

Table 10. Exponential functions relating the depth of the CH<sub>4</sub> band at 1.2  $\mu$ m and H<sub>2</sub>O band at 1.4  $\mu$ m with near-infrared spectral types.

		Exponential fit			
x	У	$c_0$	$c_1$	$c_2$	$\chi^2_{red}$
NIR SpT	$\mathrm{CH}_4$	$(1.13\pm1.09)\times10^{-3}$	$1.60 \pm 0.08$	$1.35 \pm 0.20$	1.32
NIR SpT	${\rm H_2O}$	$(6.51\pm1.90)\times10^{-2}$	$1.40 \pm 0.03$	$(9.71\pm2.57)10^{-1}$	1.31

Note—The exponential function is defined as:  $y = c_0 * c_1^x + c_2$ 

Note—To obtain NIR spectral types from the value of the index :  $x = \log_{c_1} \left( \frac{y - c_2}{c_0} \right)$ 

**Table 11.** Typical dimensionless values for the  $H_2O$  and  $CH_4$  bands per spectral type calculated using the corresponding exponential function in Table 10.

SpT	Value CH <sub>4</sub> index	Value H <sub>2</sub> O index
L4	$1.35 \pm 0.54$	$1.21 {\pm} 1.01$
L5	$1.36 {\pm} 0.54$	$1.31 \pm 1.01$
L6	$1.37 {\pm} 0.54$	$1.45 \pm 1.01$
L7	$1.38 {\pm} 0.54$	$1.64 \pm 1.01$
L8	$1.40 {\pm} 0.54$	$1.90 \pm 1.01$
L9	$1.43 \pm 0.54$	$2.27{\pm}1.01$
T0	$1.47 {\pm} 0.54$	$2.79 \pm 1.01$
T1	$1.55 {\pm} 0.54$	$3.50 {\pm} 1.01$
T2	$1.67 {\pm} 0.54$	$4.50 {\pm} 1.01$
T3	$1.87 {\pm} 0.54$	$5.89 \pm 1.01$
T4	$2.18 {\pm} 0.54$	$7.84 {\pm} 1.01$
T5	$2.68 {\pm} 0.54$	$10.54{\pm}1.01$
T6	$3.48 {\pm} 0.54$	$14.32 {\pm} 1.01$
T7	$4.77 {\pm} 0.54$	$19.60 \pm 1.01$
Т8	$6.83 {\pm} 0.54$	$26.95{\pm}1.01$
Т9	$10.13 \pm 054$	$37.20 \pm 1.01$

 $F_{\lambda, \, {\rm reflected}}$  is the flux reflected from the star by the planet in the near-infrared, that depends on its albedo at those wavelengths.

$$F_{\lambda, thermal} = F_{\lambda, planet} - F_{\lambda, reflected} \tag{9}$$

Where,

$$F_{\lambda, reflected} = A \times \frac{F_{\lambda, star}}{4\pi a_{star-planet}^2} \times 4\pi R_{Planet}^2 \quad (10)$$

A is the geometrical albedo in the near-infrared. The geometrical albedo is wavelength dependent, and varies depending on multiple factors, including the composi-

tion of the planetary atmosphere, particle sizes in its atmosphere, surface gravity, etc. (Marley et al. 1999). As it is non-trivial to determine the wavelength dependency of hot Jupiter geometrical albedo, we assume a maximum near-infrared constant albedo of 0.1 for the estimation of the reflected flux from the star, as predicted by Marley et al. (1999).  $F_{\lambda, star}$  is the flux density of the star given by the corresponding model spectra (Allard et al. 2012), scaled to the star's distance from Earth. The scaling was done by multiplying the model spectrum by  $(R_{star}/d_{star-\oplus})^2$ , where  $R_{star}$  is the radius of the star (obtained from Pecaut & Mamajek 2013) and  $d_{star-\oplus}$  is the distance between the Earth and the star (based on the trigonometric parallaxes of the host stars, see references in Table 2).  $a_{star-planet}$  is the star-planet distance, available in the literature for all hot Jupiters (see Table 2). Finally,  $R_{Planet}$  is the radius of each planet (Table 2).

As seen in Figure 8, the contribution of the albedo-assumed reflected light to the observed HST/WFC3 emission spectra is almost negligible, and it does not significantly change the colors and/or absolute magnitudes of the hot Jupiters considered in this study.

A sub-group within hot Jupiters has been identified recently as ultra-hot Jupiters. The six brightest hot Jupiters in our sample belong to this category (WASP-33B, Kepler-13Ab, WASP-18b, WASP-121b, WASP-103b, and WASP-12b). Lothringer et al. (2018) and Parmentier et al. (2018) among others proposed that under extreme irradiations strong molecular dissociations and H<sup>-</sup> opacity will significantly reduce or even eliminate the molecular absorption bands in the dayside emission spectra of hot Jupiters. In our spectral library, these ultra-hot Jupiters all appear to lack the  $1.4\mu m$  water absorption band (Figure 10), while they have consistent J-H colors with the color sequence defined by M dwarfs (Figure 8). This agrees with the  $T_{\text{eff}}$  of 2,500–3,000 K estimated by Haynes et al. (2015), Beatty et al. (2017), Sheppard et al. (2017), Evans et al. (2017), Cartier et al. (2017) and Stevenson et al.

(2014a), respectively for the ultra-hot Jupiters in our sample. WASP-4b and TrES-3b have similar  $\rm M_{J}$  to those of early-type L dwarfs (Dupuy & Liu 2012), with similar  $\rm T_{eff}$  of  $\sim\!2,\!000$  K (Ranjan et al. 2014). Finally, HD 209458B and WASP-43b have similar estimated  $\rm T_{eff}$  of  $\sim\!1,\!500\text{-}1,\!700$  K (Line et al. 2016; Stevenson et al. 2014b, respectively), than mid-L dwarfs, and actually lie among other mid-L dwarfs in Figure 8.

Finally, we compared the HST/WFC3 hot Jupiter emission spectra compiled in this work, to spectra collected in the SpeX Spectral Library. chose the best-fits based on the value of their modified  $\chi^2$  (G), as obtained using equation 1, and visual inspection. We found that only three of the ten hot Jupiters in our study had best matches to mid-L dwarfs: HD 209458B was matched to SDSS J104335.08+12131 (L7), WASP-43b was matched to SDSS J140023.12+43382 (L7), and TrES-3b was matched to 2MASS J21513979+34024 (L7 peculiar). The other seven hot Jupiters are best matched to M-dwarfs: WASP-33b, WASP-103b and WASP-18b are best matched by M3-type stars (NLTT 6012a, 2MASSJ13032137+23511, and LSPMJ0734+5810, respectively). Kepler-13Ab is best matched by 2MASS J11070582+28272 (M7), that is consistent with the result of Beatty et al. (2017), who found a best match to an M8 brown dwarf. WASP-4b best matched to 2MASS J02481204+24451 (M8). Finally, we did not find a best match for WASP-12b and WASP-121b (see Figure 10).

These results are in general consistent with the positions of hot Jupiters within the CMD in Figure 8, and with the temperatures predicted for those objects by their respective authors. In addition, these results also agree with predictions made by the atmospheric models presented in Fortney et al. (2008), in which they suggested that there are two classes of hot Jupiter day-side atmospheres analogous to the M- and L dwarfs spectral types, that they called pM and pL, respectively. The pM class planets have hot stratospheres due to the high irradiation of their parent star ( $T_{\rm eff} > 2000 \text{ K}$ ) with temperature inversion in their atmospheres, and molecular bands in emission. The models predict that the temperature differences between their day and night side are high due to radiative time constants at photospheric pressures are shorter than advective timescales. In contrast, the pL class planets are less irradiated by their parent star. The incident flux from the parent star is absorbed in the atmosphere, and redistributed easily, as there is no thermal inversion in their photospheres. Thus, they have cooler day sides and warmer night sides. Their spectra are dominated by H<sub>2</sub>O, in the near-infrared, and by Na and K absorptions in the optical.

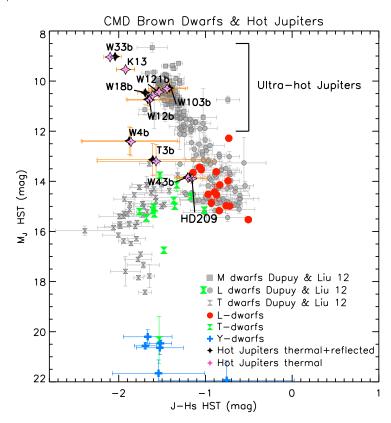
Our results are consistent with, and extend on, a previous study by Triaud (2014), who compared in a CMD plot transiting planet dayside emission measured in two Spitzer/IRAC bands ([3.6] and [4.5] filters) to those of brown dwarfs. This comparison suggested overall similarity, with a few possible outliers. The Spitzer colormagnitude comparison is particularly sensitive to the presence/absence of methane absorption. In a follow-up work Triaud et al. (2014) extended their study to nearinfrared continuum bands, and to a larger sample using photometric distances to compare their objects with brown dwarfs in a CMD plot. They found that, for a given luminosity, hot Jupiters' daysides show larger ranges of colors than brown dwarfs, specially with decreasing intrinsic luminosity. In contrast, our study suggests that brown dwarfs and hot jupiters tend to be similar in the shorter-wavelength continuum emission (J and H bands) bands and that, for many hot jupiters, brown dwarfs can provide surprisingly good spectral matches.

Fortney et al. (2008) used these atmospheric models to predict the classes of several hot Jupiters, in which TrES-3b and HD 209458B were included. He found that TrES-3b should belong to the pM class, as we consistently obtained. He predicted HD 209458B to be in the transition zone between both classes, but, as Line et al. (2016) also found, we conclude that this hot Jupiter matches better to an L-dwarf spectrum, thus, to the pL class.

### 9. SUMMARY AND CONCLUSIONS

We present a very high-quality HST/WFC3 near-infrared spectral library of brown dwarfs (field and companion to stars), planetary-mass objects, and hot Jupiters, to enable quantitative comparative studies. In this paper, we provide an initial characterization and analysis of these HST/WFC3 near-infrared spectra:

1. In Section 3 we derive near-infrared spectral types for the brown dwarfs and the substellar companions to stars uniformly, using as comparison the SpeX Spectral Library templates. We conclude that their spectral types are mostly consistent with the spectral types provided in the literature within ±1.5 subspectral type. The only exception is for low surface-gravity objects, for which the differences found in spectral types are ±3 spectral types. This is expected, as the SpeX Spectral Library templates are mostly composed of field gravity low-mass stars and brown dwarfs.



**Figure 9.** CMD diagram showing brown dwarfs and hot Jupiters together. Red dots represent L dwarfs, green hour-glasses represent T dwarfs, and blue crosses represent Y dwarfs. hot Jupiters are shown as black stars, and hot Jupiter after removing the contribution of their albedo (thermal flux) are shown as pink stars.

- 2. In Section 4 we plot a M<sub>J</sub> vs. J − H color-magnitude diagram to compare our sample to other substellar objects, with the objective to identify brown dwarfs with peculiar colors/brightness, including red or blue objects, low-surface gravity objects, binaries, etc. We find that objects 32 (2M1750+1759, known binary), and 34 (2M0559-34, overluminous), are overluminous in the color-magnitude diagram suggesting that they are potential multiple systems.
- 3. In Section 5 we obtain the *H*-continuum and KI<sub>J</sub> near-infrared spectral indices from Allers & Liu (2013) to search for potential L4 to L8 low-surface gravity substellar objects in our sample. We found two very low-gravity dwarfs: CD-352722B (object 1) and 2M0355+1133 (object 3). In addition, we found five intermediate surface-gravity objects: 2M0421-6306 (object 6), W0047 (object 9), 2M0107+0041 (object 11), PSO J318.5-22 (object 15) and 2M2224-0158 (object 16).
- 4. In Section 6 we apply the method from Burgasser et al. (2006b, 2010) and Bardalez Gagliuffi et al. (2014) to search for candidates for composite spec-
- tra in our sample. Their spectral indices selected 13 composite spectra candidates, from which eight were selected from the F-statistic analysis described in Section 6. None of these eight objects are overluminous, as would be expected for binary or multiple brown dwarfs systems. In addition, we found that five of the eight selected objects have been reported in the literature as photometrically variable. Thus, this method might be useful to find potential variable late L and early T dwarfs. Nevertheless, we also found that not all objects in our sample with reported photometric variability have been detected by Burgasser et al. (2006b, 2010), and Bardalez Gagliuffi et al. (2014) method. The indices itself detected 9 out of 19 variables and 3 no variable objects in our sample, with spectral types between L4 and T2.
- 5. In Section 7 we measure the depths of the water band at  $\sim 1.4~\mu \mathrm{m}$  and the methane band at  $\sim 1.2~\mu \mathrm{m}$  for brown dwarfs and substellar companions to stars. We derive a relation between their near-infrared spectral types and the depths

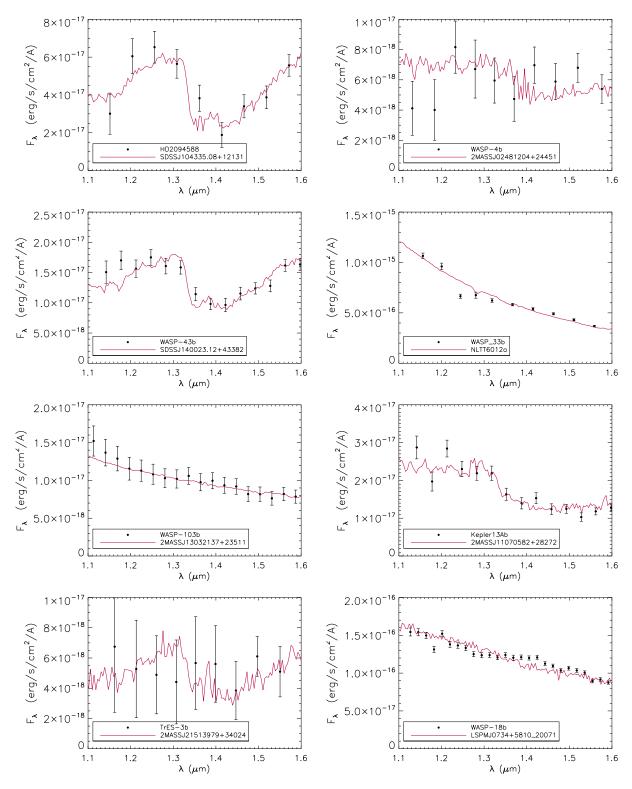


Figure 10. Best matches of hot Jupiters to brown dwarf spectra and M-dwarfs.

of those bands, providing a tool for spectral classification of other substellar objects.

6. In Section 8 we compare the emission spectra of the dayside of hot Jupiters to the spectra of brown dwarfs and substellar companions to stars in our sample and low-mass stars and brown dwarfs from the SpeX spectral library. We found best matches to either L or M-dwarfs for eight out of the ten hot Jupiters of our sample. In addition, we plot a color-magnitude diagram using J and  $H_{short}$ HST bands for all our sample. The hottest hot Jupiters, WASP-33B, Kepler-13Ab, WASP-18b, WASP-121b, WASP-103b, and WASP-12b have similar M<sub>J</sub> magnitudes to mid-M dwarfs (Dupuy & Liu 2012), which agrees with the  $T_{\rm eff}$  of 2,500– 3,000 K estimated by their respective authors. WASP-4b and TrES-3b have similar  $M_J$  that early L dwarfs (Dupuy & Liu 2012), with similar  $T_{\rm eff}$ of  $\sim 2000$  K. Finally, HD 209458B and WASP-43b have similar estimated  $T_{\rm eff}$  of  $\sim 1500\text{-}1700~{\rm K}$  (Line et al. 2016; Stevenson et al. 2014b, respectively), than mid-L dwarfs, and actually lie among other mid-L dwarfs in the CMD diagram.

The HST/WFC3 near-infrared spectra presented in this work will be available as csv files that will include wavelength, flux and uncertainty in flux. The spectral library will be available in a machine-readable format at the High-level Science Products website at the MAST archive under the Cloud Atlas program's page doi:doi:10.17909/t9-asft-6k38.

We thank our anonymous referee for his/her useful comments that helped to improve our paper. Based on observations made with the NASA/ESA Hubble Space Telescope, obtained at the Space Telescope Institute, which is operated by AURA, Inc., under NASA contract NAS 5-26555, under GO-13241, GO-14241, GO-12550, GO-13176, GO-12550, GO-13299, GO-13280, GO-13281, GO-12314, GO-14051, GO-12217, GO-13178, GO-12970, GO-12230, GO-13467, GO-12495, GO-14050, GO-13467, GO-12181, GO-13308, and GO-14767. This publication makes use of data products from the Two Micron All Sky Survey, which is a joint project of the University of Massachusetts and the Infrared Processing and Analysis Center/California Institute of Technology, funded by the National Aeronautics and Space Administration and the National Science Foundation. This work makes use of the SpeX Prism Spectral Library. This work has made use of data from the European Space Agency (ESA) mission Gaia (https://www.cosmos.esa.int/gaia), processed by the Gaia Data Processing and Analysis Consortium (DPAC, https://www.cosmos.esa.int/web/ gaia/dpac/consortium). Funding for the DPAC has been provided by national institutions, in particular, the institutions participating in the Gaia Multilateral Agreement. We have done use of the Matplotlib python library (Hunter 2007).

APPENDIX

Table A1. Log of the sample of L, L-T, T and Y dwarfs with HST/WFC3 spectroscopy.

Num.	Name	Obs. dates	Num. orbits×visits	$T_{exp}$ single expo. (s)	Num. single expo./orbit
1	CD-352722b	2015 Sept 07	2x1	29.6	86
2	$2MASS\ J17502484-0016151$	2012 Jun 15	1x1	22.6	64
3	2MASS J03552337+1133437	2015 Oct 06	2x1	9.68	46
4	2MASS J18212815+1414010	2013  Jun  09  & Jun  27	3x2	112.0	19
2	$2MASSW\ J1507476-162738$	2013 Apr 30 & May 12	4x2	67.3	30
9	$2 \rm MASSI~J0421072\text{-}630602$	$2012~\mathrm{Mar}~20$	1x1	112.0	19
7	$2 \rm MASS\ J05395200\text{-}0059019$	$2012~\mathrm{Mar}~01$	1x1	45.0	37
œ	2MASSI J1711457+223204	2012  Aug  01	1x1	223.7	6
6	2MASS J00470038+6803543	2016  Jun  06/07	6x1	201.3	65
10	LP261-75B	$2016 \mathrm{Dec}\ 20$	6x1	201.4	99
11	2MASS J01075242 + 0041563	2017 Jan 02	6x1	201.4	22
12	2MASSW J1515008+484742	2012 Feb 25	1x1	45.0	36
13	2MASS J06244595-4521548	2012 May 08	1x1	45.0	38
14	2MASSW J0801405+462850	2011 Nov 10	1x1	223.7	11
15	PSO J318.5-22	2016 Sept 08	5x1	278.0	6
16	$2 \rm MASSW~J 2224438-015852$	2015 Sept 09	2x1	9.68	26
17	Luh 16AB	2013 Nov 8	5x1	76.2	100
18	2MASSI J0825196+211552	2012 May 09	1x1	112.0	21
19	2MUCD 10802	2011 Dec 09	1x1	45.0	40
20	$2MASS\ J16322911+1904407$	2012 Aug 11	1x1	223.7	6
21	$2MASSW\ J0310599+164816$	2012 Aug 25	1x1	223.0	10
22	$2MASS\ J12195156+3128497$	2012 Jun 18	1x1	223.7	6
23	SDSS $J075840.33 + 324723.4$	2014  Apr  12	5x1	112.0	22
24	$2MASS\ J10393137 + 3256263$	2012 May 08	1x1	223.7	11
25	$2MASS\ J09090085+6525275$	2012  Aug  21	1x1	223.7	10
26	$2MASS\ J21392676+0220226$	2010 Oct 21	6x1	22.3	11
27	$2MASS\ J13243553+6358281$	2012 Feb 25	1x1	112.0	21
28	$2MASS\ J16291840 + 0335371$	2015 June 06	4x1	112.0	21
29	HNPEGB	2017 May 16	6x1	201.4	65
30	SIMP J013656.5+093347.3	2013 Sep 28 & Oct 7	4x2	112.0	19
31	GUPSCB	2018 Jan 08	6x1	201.4	43
32	$2MASS\ J17503293+1759042$	2012 Oct 05	1x1	223.7	10
33	$2 {\rm MASS\ J00001354}{+}2554180$	2012 Sept 13	1x1	45.0	40
34	$2 {\rm MASS\ J05591914\text{-}}1404488$	2011  Oct  16	1x1	22.6	62
35	2MASSI J2339101+135230	$2012~\mathrm{Aug}~21$	1x1	223.7	11

Table A1 continued

Table A1 (continued)

Num.	Name	Obs. dates	Num. orbits×visits	T <sub>exp</sub> single expo. (s)	Num. single expo./orbit
36	2MASS J11101001+0116130	2016 Feb 10	2x1	201.4	21
37	2MASS J22282889-4310262	2013  Jul  20 & 27	4x2	201.3	39
38	2MASS J08173001-6155158	2011 Oct 09	1x1	22.6	29
39	S Ori J053814.5-024512	2010 Sept 05	1x1	602.7	4
40	$2MASSI\ J0243137-245329$	2011 Dec 31	1x1	112.0	19
41	$2MASS\ J16241436+0029158$	2012 Jul 13	1x1	112.0	17
42	CFBDSIR2149-0403	2015 Sept 09 & Nov 18	4x2	290.7	16
43	S Ori J053814.5-024512	2010 Oct 06	1x1	602.7	4
44	ROSS458C	2018  Jan  05 & 06	7x1	201.4	21
45	WISEA J032504.52504403.0	2013 Aug 04	1x1	403.0	4
46	WISEA J033515.07+431044.7	2013 Aug 30	1x1	453.0	4
47	WISEA J040443.50642030.0	2013 Apr 09	1x1	453.0	4
48	WISEA J221216.27693121.6	2013 Sept 11	1x1	453.0	4
49	WISEA J094306.00+360723.3	2013 Feb 20	1x1	503.0	4
20	WISEA J154214.00+223005.2	2012 Mar 04	1x1	503.0	4
51	WISEA J035934.07540154.8	2011 Aug 10	1x1	553.0	4
52	WISEA J041022.75+150247.9	2012 Sept 01	1x1	503.0	4
53	WISEA J073444.03715743.8	2013 May 20	1x1	453.0	4
54	WISEA J120604.25+840110.5	2013 July 15	1x1	453.0	4
55	WISE J154151.65225024.9	2013 May 09	1x1	453.0	4
26	WISEA J173835.52+273258.8	2011 May 12	1x1	503.0	4
22	WISEA J205628.88+145953.6	2011 Sept 04	1x1	503.0	4
28	WISEA J222055.34362817.5	2013 Jun 20	1x1	1103.0	4
59	WISEA J220905.75+271143.6	2012 Sept 15	1x1	503.0	4
09	WISEA J082507.37+280548.2	2014 Jan 17	1x1	2406.0	3
61	WISEA J140518.32+553421.3	2011 Mar 14	1x1	553.0	4
62	WISEA J163940.84684739.4	2013 Oct 29	1x1	602.9	4
63	WISEA J053516.87750024.6	2011 Sept 27	1x2	553.0	4
63	WISEA J053516.87750024.6	2012 Sept 17	1x2	553.0	4
63	WISEA J053516.87750024.6	2013 Sept 27	1x1	1269.0	9
64	WISEA $J035000.31565830.5$	2011 Aug 13	1x1	553.0	4
65	WISEA J064723.24623235.4	$2013~\mathrm{May}$ 13 & $2013~\mathrm{Dec}$ 29	1x2	1203.0	9
99	WISEA $J235402.79+024014.1$	2013 Sept 22	1x1	806.0	4
29	WASP-18b	2014 Apr-Jun & Aug	6x4	73.74	8?
89	WASP-33b	2012 Nov $25~&~2013$ Jan $14$	5x2	51.7	119
69	WASP-12b	$2011~\mathrm{Apr}~12$	5x1	7.35	188
70	WASP-121b	$2016 \; \mathrm{Nov} \; 10$	5x2	103	16

Table A1 continued

Table A1 (continued)

Num.	Name	Obs. dates	Num. orbits×visits	$T_{exp}$ single expo. (s)	Num. orbits×visits $T_{exp}$ single expo. (s) Num. single expo./orbit
7.1	WASP-43b	2013 Nov 09 & Dec 5	14x2	103.129	19
72	WASP-103b	2015  Jun  17 & 17	5x2	81.089	12
73	TrES-3b	2011 March 02	4x1	36.02	219
74	Kepler-13Ab	2014 Apr 28 & Oct 13	5x2	9.2	101
75	HD 209458B	Sept-Dec 2014	5x5	14.971	43
92	WASP-4h	2010 Nov 25	5x1	36.02	268

# A. DATASET DESCRIPTION AND DATA REDUCTION

In the following we summarize the key steps and references for the different data reductions performed on the spectra compiled in this work.

# A.1. Brown Dwarf and Low-mass Companion Spectra A.1.1. Time-resolved spectra

We present the datasets with time-resolved spectroscopy taken for several of the brown dwarfs compiled in this work.

The Apai et al. (2013) dataset consists of the first two brown dwarfs (SIMP J013656+093347 and 2MASS J21392676+0220226) observed in time-resolved observations, and obtained in the GO-12314 program (PI Apai). Each objects was observed in six consecutive HST orbits. Apai et al. (2013) provides a detailed summary of the reduction procedure. In this spectral library, we present the median of the time-resolved spectra for each object.

The Buenzli et al. (2012) study presented near-infrared, time-resolved, six orbit-long spectroscopy of a single target (2M2228, object 37) and was reduced with a method identical to that in Apai et al. (2013). These data were also taken in program GO-12314 (PI: Apai), and the target was also a known variable brown dwarf. The observations showed spectral variability with pressure-dependent phase shifts Buenzli et al. (2012). We took the median of the time-resolved spectra for our library.

The Buenzli et al. (2014) sample consists of 22 brown dwarfs with spectral types between L5 and T6. These data were collected in an HST SNAP program (PO 12550, PI Apai). Basic reduction followed the same steps as for the previous programs. The ramp effect was corrected using an analytical function fitted to the flux of a non-variable star as in Apai et al. (2013), in addition to removing the first 180 s of each time series where the scatter in the ramp effect is substantial.

The Buenzli et al. (2015) study presented spatially and temporally resolved spectroscopy for the Luhman 16 A and B binary brown dwarf components. Reduction followed the steps described in Apai et al. (2013); Buenzli et al. (2014). We present the combined Luhman 16 A and B spectra in this paper.

The time-domain programs described above focused on relative variations and did not correct for wavelength-dependent aperture losses, which is not relevant in the related studies. However, for our purposes these corrections are necessary. We performed a uniform aperture correction on all sources from the above studies to correct for flux loss due to the finite width of the spectral extraction windows. We corrected for the missing flux

per wavelength on the basis of measured wavelengthdependent flux losses, performing a bi-linear interpolation in wavelength and aperture width of the values of the aperture corrections tabulated in Table 6 from Hartig (2009).

Within the Cloud Atlas HST treasury program (HST) GO 14241), time-resolved spectroscopy observations for eight L4 to T7 high- and low-surface gravity brown dwarfs were obtained. The data were collected between September 2015 and September 2018. The Cloud Atlas program uses time-resolved spectroscopy to probe the spatial distribution and properties of condensate clouds. A publication in preparation (Apai et al., in prep.) will provide an overview of the program and its key results from the time-resolved spectroscopy. Results for three objects have already been published (WISE0047: Lew et al. 2016; LP261-75B: Manjavacas et al. 2018; HN Peg B: Zhou et al. 2018), while other papers are in preparation (S0107: Apai et al. 2018, in prep.). Here we present the spectral results based on time-averaged spectra for all objects. We performed the data reduction using very similar methods as described above for previous studies (Apai et al. 2013; Buenzli et al. 2012, 2014, 2015, and references therein). An important difference, however, is the use of a significantly improved WFC3 ramp correction method. Zhou et al. (2017) identified charge trapping and delayed release as the cause of the "ramp effect" and developed a solid-state physics-based model capable of reliably correcting this effect in a wide variety of WFC3 data. Most of the Cloud Atlas datasets published use the ramp effect correction by Zhou et al. (2017). The uncertainty level for our spectra after the data reduction is 0.1–0.3\% per spectral bin, measured using the reduced individual spectra. These uncertainties are due to photon noise, errors in the sky subtraction, and the read-out noise. Finally, we performed aperture corrections following the same procedure as for the other spectra mentioned previously in this Section.

Within the HST GO 13299 and 14051 (P.I. Radigan), time-resolved near-infrared spectra observations were obtained with HST/WFC3 to study the rotational modulations of two unusually blue L dwarfs. The objective of this project was making spectrally and spatially resolved maps of these objects. These objects are SDSS J075840.33+324723.4 (object 23), and 2MASS J16291840+0335371 (object 28). SDSS J075840.33+324723.4 was observed during five consecutive orbits, and 2MASS J16291840+0335371 was observed during four consecutive orbits. The data reduction was performed using a similar procedure as for the *Cloud Atlas* treasury program data. In this paper, we present the median combined spectra of all the time-

resolved near-infrared spectra taken during the consecutive orbits in which these objects were observed. The uncertainty level for these spectra after median combine all time-resolved spectra is  $\sim 0.03\%$  at 1.25  $\mu$ m. These uncertainties are due to photon noise, errors in the sky subtraction, and the read-out noise.

Finally, Biller et al. (2018) presents time-resolved spectroscopy of the red L7 dwarf, PSO 318-22. A difference with the previous studies is that the ramp correction was corrected using four background stars in the field of view (2-3 times brighter than the target). They median combined and normalized the while light curves of the background stars to produce a calibration curve. Then they divided the target's light curve by the calibration curve to eliminate the ramp effect and other systematics, following a similar approach as done by previous ground studies (Biller et al. 2015; Radigan 2014).

## A.1.2. Single spectra

S Ori 70 and S Ori 73 (Peña Ramírez et al. 2015) are T7±0.5 and T4.5±0.5 dwarfs, respectively. S Ori 70 and S Ori 73 were observed with HST/WFC3 (PI Lucas, HST-GO-12217). Details on the data reduction can be found in Peña Ramírez et al. (2015).

In addition, we include 22 T8 to Y2 brown dwarf spectra presented and analyzed in Schneider et al. (2015). The observations were carried out within the P.I. Kirkpatrick, programs 12330, 13178 and P.I. Cushing, programs HST-GO-12544 and HST-GO-12970. As G141 with which the observations were performed is slitless grism, the source spectra are sometimes contaminated by photons from nearby sources. To address this problem, Schneider et al. (2015) developed a source extraction routine to define source apertures and background regions on the individual images. After the best aperture is defined, aperture corrections and flux calibrations are performed following Kuntschner et al. (2011). For objects with multiple visits, the images have been median combined to produce a final spectroscopic image. Finally, spectra are extracted as indicated above. The published spectra are time-averaged spectra.

# A.2. Hot Jupiter emission spectra

In this Section, we summarize the different reduction methods performed by the respective authors in which hot Jupiter's emission spectra were published.

### A.2.1. WASP-18b

The emission spectrum of WASP-18b was presented in Sheppard et al. (2017). Observations of three individual eclipse events were obtained during three epochs as part of the program GO 13467. At a forth epoch observations were obtained with two eclipses within an orbital phase curve. Grism observations were taken in spatial scan mode with forward-reverse cadence (Dressel 2018). Further details in the data reduction are found in Sheppard et al. (2017). Finally, a forward scan and a reverse scan light curve were obtained and analyzed separately. To correct non-astrophysical effects, the systematic trends were removed using parametric marginalization (Wakeford et al. 2016), and then further detrending was performed by the subtraction of scaled bandintegrated residuals from wavelength bins (Haynes et al. 2015). The wavelength bins of the spectrum are given in Table 1 of Sheppard et al. (2017).

# A.2.2. WASP-33b

The emission spectrum of WASP-33b was first published by Haynes et al. (2015). WASP-33b is orbiting a  $\delta$ -Scuti star (Herrero et al. 2011), and its modulations were model with sine functions. To produce the 2D spectral frames from the files provided by the standard pipeline, a top hat mask was applied in the spatial direction of each read of a width of 20 pixels tall (Herrero et al. 2011); then, subsequent reads were subtracted and added to differenced frames to create one scanned image (Deming et al. 2013). To correct bad pixels, the method by Mandell et al. (2013) was used within the combined spectral frames, and combine the images into 1D spectra. To perform the wavelength and flat-field wavelength dependent calibrations, the coefficients from Wilkins et al. (2014) were used.

# A.2.3. WASP-12b

The emission spectrum of WASP-12b was published in Stevenson et al. (2014a). The observations of WASP-12b were taken in five consecutive in staring mode. Further details on the observations can be found in Swain et al. (2013). Data were reduced using the standard HST pipeline as explained in detailed in Stevenson et al. (2014a).

To trace the first order spectra, the direct image was located using a two-dimensional Gaussian and then use Table 1 in Kuntschner et al. (2009) to provide a direct-to-dispersed image offset. The wavelength calibration is performed using the coefficients provided in Table 5 from Kuntschner et al. (2009). The flat field was modeled using the standard calibration flat files. The spectral extraction was performed within a box of  $150\times150$  pixels centered on the spectrum. The spectral extraction was performed along 40 pixels in the spatial direction, and the remaining pixels in the box were used for the background subtraction, generating eleven light curves.

# A.2.4. WASP-121b

The emission spectrum of WASP-121b was first presented in Evans et al. (2017). Data reduction performed using the HST/WFC3 standard pipeline. Details explained in detailed in Evans et al. (2017). The target flux was extracted taking the difference between successive non-destructive reads. The background was measured as a median count of a box of 110 columns along the dispersion axis and 20 rows along the cross-dispersion axis. To remove flux contributions from nearby stars and cosmic ray hits, all pixels above and below 35 pixels from the center of the spectrum along the cross-dispersion axis were set to zero. Finally, all frames were added together. The spectrum was then extracted by summing the flux within a rectangular aperture across the dispersion axis with apertures from 100 to 200 pixels. The data taken during the first HST orbit were discarded due to a strong ramp effect, as well as the first exposure of the remaining HST orbits.

### A.2.5. WASP-43b

The emission spectrum of WASP-43b was first presented in Stevenson et al. (2014b). The observations were performed during 13-14 HST orbits on each primary transit or secondary eclipse visit, each of them consisting of four orbits. Due to the ramp effect, the first orbit of each visit was removed from the analysis. For the rest, the ramp was fitted with an exponential ramp model. Further details on the data reduction can be found in Stevenson et al. (2014b).

# A.2.6. WASP-103b

The emission spectrum of WASP-103b was first presented in Cartier et al. (2017). WASP-103b was observed with ten HST orbits in two visits. The first orbit of both visits was discarded.

To remove the background, images were created using sequential pairs of up-the-ramp readouts within each exposure. For those subframes, a conservative mask was used to determine the background region and measure the sky background level, assuming that is spatially flat and uniform due to the short exposure times. This background was subtracted from all subframes. In addition, a smaller mask was defined (Deming et al. 2013; Knutson et al. 2014) and all pixels outside of the mask were zeroed. This helps to reduce noise and exclude cosmic rays (CRs) in the background area when later combining all subframes to determine the flux for each exposure. Special flat fields were created for the data reduction using the determined centroids in the spectral direction (X) and scan direction (Y) direct image frame, assuming that every column has the same wavelength (Cartier

et al. 2017). Finally, to remove additional cosmic rays and bad pixels, a moving median filter was applied. The final extracted spectrum was binned to 22 wavelength channels.

Instrumental effects and systematics due to the ramp correction were removed using Gaussian Processes (GP) regression (Rasmussen & Williams 2006) that does not need to pre-specify a parametric model. To find the best fit light curve eclipse model GP regression was used (Cartier et al. 2017).

# A.2.7. TrES-3b and WASP-4b

The emission spectra of TrES-3b and WASP-4b were first presented in Ranjan et al. (2014). The observations were carried out during four consecutive orbits during the eclipse of TrES-3b, and five consecutive orbits during the eclipse of WASP-4b. The first orbit of each observation was discarded to avoid the most prominent ramp effect systematics.

The details on the data reduction can be found in Ranjan et al. (2014). The background subtraction was performed by choosing a fixed area on the detector, matching the wavelength range of the spectrum free of object flux in the individual 2D images. These background columns are scaled to match the spectral extraction aperture. Finally, the extracted spectra were binned in wavelength to enhance the signal-to-noise per resolution element (see the bin's wavelength ranges in Ranjan et al. 2014 for each object).

# A.2.8. Kepler-13Ab

The emission spectra of Kepler-13Ab was first presented in Beatty et al. (2017). The Kepler-13Ab system is composed by three stars: the planet host, Kepler-13A, and the unresolved binary Kepler-13BC, with the two components separated by 1."15 (Shporer et al. 2014).

The observations were carried out during two visits were composed of a total of five HST orbits. The planet host star is in a close binary system. The data reduction includes primary subtraction. All details of the data reduction can be found in Beatty et al. (2017). The cosmic-rays hits were removed separately in an area around the stellar spectra, and the area dominated by the sky background. Finally, the background was subtracted from each exposure by defining two background regions across the bottom and top of each of these images.

To perform a spectral extraction of Kepler-13Ab, the contribution of Kepler-13BC needs to be subtracted first. Using the WAYNE simulator (Varley et al. 2017), the artificial 2D spectra of Kepler-13BC was created and subtracted to create an undiluted 2D spectrum of

Kepler-13A. Finally, to perform the light-curve extraction, the spectral trace of Kepler-13A was fitted with a Gaussian profile along the detector columns. Then the columns along the detector were summed using an extraction aperture with a half width of 4.5 pixels, to generate a 1D spectrum of Kepler-13A.

The wavelength calibrations was done using the direct image taken at the beginning of each of the visits. The X- and Y- location of both Kepler-13A and Kepeler-13BC where determined on the detector subarray, and then Kuntschner et al. (2009) wavelength calibration method was implemented to calculate a wavelength solution for each star. The Paschen- $\beta$  line visible at 1.282  $\mu$ m was used to verified the accuracy of the wavelength calibration.

## A.2.9. HD 209458B

The emission spectra of HD 209458B was first presented in Line et al. (2016). HD 209458b was observed as part of the GO 13467 HST treasury program. It was observed during secondary eclipse over five visits, each with five HST orbits. The first orbit of each visit were excluded from the analysis to minimize the impact of the ramp effect on the dataset. A direct image was taken at

the beginning of each orbit to aid the wavelength calibration.

All details about the data reduction can be found in Line et al. (2016). To extract the 1D spectra, and optimal extraction was used (Horne 1986) with a extraction window of 110 pixel rows centered on the spectra and flanked by additional 110 pixels rows for background extraction. The spectral of all frames were combined. Finally, the combined spectra was divided into 10 spectroscopic bins.

# B. TRANSFORMATION BETWEEN 2MASS TO HST MAGNITUDES

We derive empirical relations to transform L and T brown dwarf magnitudes from the 2MASS to the HST photometric system. To obtain the HST/WFC3 near-infrared magnitudes, we follow the same procedure than in Section 8, equation 7. In Figure A1 we show the relation between the J and H-band 2MASS magnitudes, and the J and  $H_s$ -bands in the HST photometric system. We do not include the T9-T9.5 dwarfs with high phometric uncertainties. Finally, we calculate a linear relationship between both photometric systems for the J and the H-band independently. The coefficients for both relations are presented in Table A2.

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**Table A2.** Linear functions relating the J and H-band 2MASS magnitudes and the J and Hs respectively for L and T brown dwarfs.

		Linear fit	
x	у	$c_0$	$c_1$
$J~{ m 2MASS}$	$J~\mathrm{HST}$	$-0.15535 \pm 0.00254$	$1.00938 \pm 0.00017$
H 2MASS	Hs HST	$2.17362 \pm 0.00206$	$0.98229 \pm 0.00014$

Note— The linear function is defined as:  $y = c_0 + c_1 x$ 

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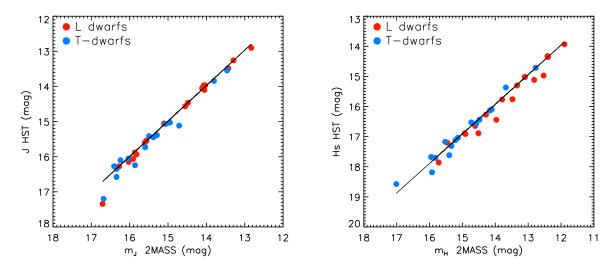


Figure A1. Relation between J and H 2MASS magnitudes, and J and Hs HST magnitudes for L and T dwarfs. L dwarfs are shown as red points, and T dwarfs are shown as blue points. Photometric uncertainties are smaller than the symbols.

Table A3. Spectral indices to select L plus T brown dwarf binary candidates.

Index	Numerator Range $^a$	Denominator Range <sup>a</sup>	Feature	Reference
$\mathrm{H}_2\mathrm{O}\text{-}\mathrm{J}$	1.140-1.165	1.260-1.285	$1.150~\mu\mathrm{m}~\mathrm{H}_2\mathrm{O}$	1
$\mathrm{CH_{4}\text{-}J}$	1.315-1340	1.260-1285	$1.320~\mu\mathrm{m}~\mathrm{CH_4}$	1
${ m H}_2{ m O-H}$	1.480-1.520	1.560-1.600	$1.400~\mu\mathrm{m~H_2O}$	1
$\mathrm{CH_{4} ext{-}H}$	1.635-1.675	1.560-1.600	$1.650~\mu\mathrm{m~CH_4}$	1
H-dip	1.610-1.640	$1.560 \text{-} 1.590 + 1.660 \text{-} 1.690^b$	$1.650~\mu\mathrm{m}~\mathrm{CH_4}$	2
J-slope	1.27-1.30	1.30-1.33	$1.28~\mu\mathrm{m}$ flux peak shape	4
J-curve	$1.04\text{-}1.07\text{+}1.26\text{-}1.29^{c}$	1.14-1.17	Curvature across J-band	4
H-bump	1.54-1.57	1.66-1.69	Slope across H-band peak	4
Derived NIR SpT			near-infrared spectral ${\rm type}^d$	1

Note—a: Wavelength range in nm over which flux density is integrated; b: denominator is the sum of the flux in the two wavelength ranges; c: numerator is the sum of the two ranges; d: near-infrared spectral type derived using comparison to SpeX spectra. References: 1 – Burgasser et al. (2006b); 2 – Burgasser et al. (2010); 3 – Burgasser et al. (2002); 4 –Bardalez Gagliuffi et al. (2014).

Table A4. Index criteria for the selection of potential brown dwarf binary systems

Abscissa	Ordinate	Inflection Points
${ m H_2O ext{-}H}$	H-dip	(0.5, 0.49), (0.875, 0.49)
$\mathrm{Spex}\;\mathrm{SpT}$	$\rm H_2O ext{-}J/H_2O ext{-}H$	(L8.5,0.925),(T1.5,0.925),(T3,0.85)

<b>Table A5.</b> Delimiters for selection regions of potential brown dwarf binary system	Table A5.	Delimiters	for selection	regions of	potential brown	dwarf bin	arv systems
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Abscissa	Ordinate	Limits
SpT	СН4-Н	Best fit curve: $y = -4.3x10^{-4}x^2 + 0.0253x + 0.7178$
${ m H_2O} ext{-J}$	$\mathrm{CH}_{4} ext{-H}$	Intersection of: $-0.08x+1.09$ and $x = 0.90$
$\mathrm{H}_2\mathrm{O}\text{-}\mathrm{J}$	H-bump	Intersection of: $y = 0.16x+0.806$ and $x = 0.90$
$\mathrm{CH_{4} ext{-}J}$	$\mathrm{CH}_{4} ext{-H}$	Intersection of: $y = -0.56x + 1.41$ and $y = 1.04$
$\mathrm{CH}_{4}\text{-}\mathrm{J}$	H-bump	Intersection of: $y = 1.00x + 0.24$ , $x = 0.74$ and $y = 0.91$
$\mathrm{CH_{4} ext{-}H}$	J-slope	Intersection of: $y = 1.250x$ -0.207, $x = 1.03$ and $y = 1.03$
$\mathrm{CH_{4} ext{-}H}$	J-curve	Best fit curve: $y = 1.245x^2 - 1.565x + 2.224$
$\mathrm{CH_{4} ext{-}H}$	H-bump	Best fit curve: $y = 1.36x^2 - 4.26x + 3.89$
J-slope	H-dip	Intersection of $y = 0.20x + 0.27$ and $x = 1.03$
J-slope	H-bump	Intersection of: $y = -2.75x + 3.84$ and $y = 0.91$
J-curve	H-bump	Best fit curve: $y = 0.269x^2 - 1.326 + 2.479$

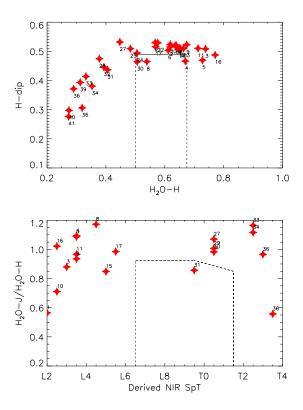


Figure A2. Spectral index selection. Numbers correspond to our objects. The boxes shown with dashed lines mark the areas where the selection criteria of Table  $\frac{A4}{A5}$  are valid.

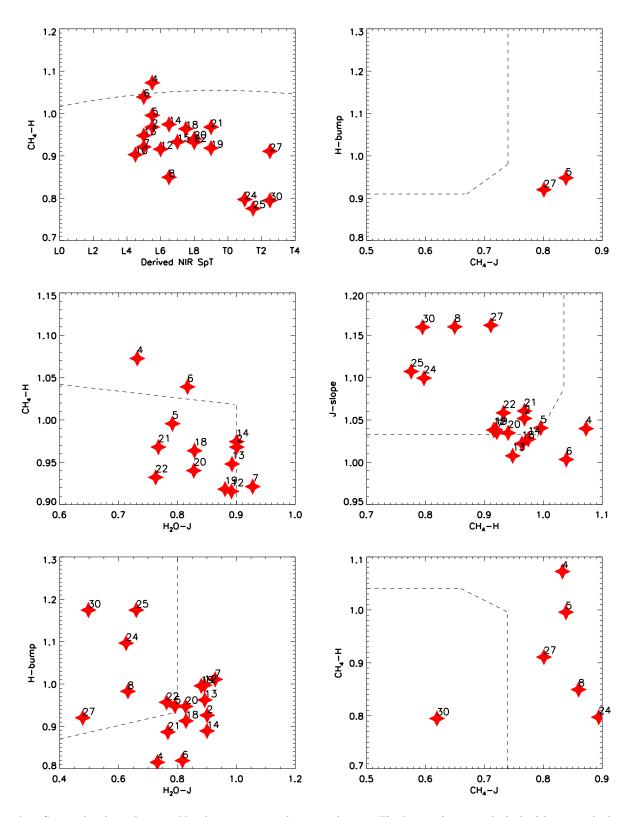


Figure A3. Spectral index selection. Numbers correspond to our objects. The boxes shown with dashed lines mark the areas where the selection criteria of Table  $\frac{A4}{A5}$  and  $\frac{A5}{A5}$  are valid.

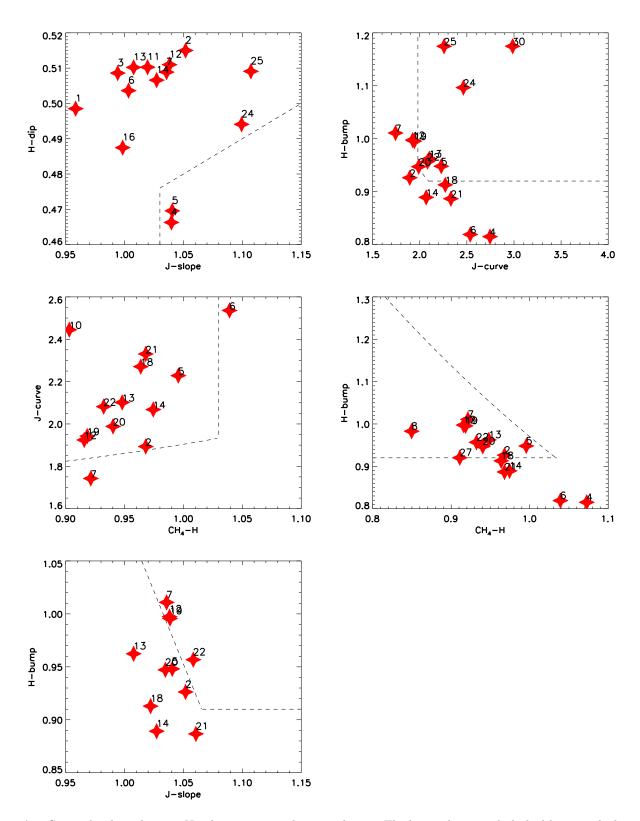


Figure A4. Spectral index selection. Numbers correspond to our objects. The boxes shown with dashed lines mark the areas where the selection criteria of Table  $\frac{A4}{A5}$  and  $\frac{A5}{A5}$  are valid.

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