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The dependence of optical polarisation of blazars on the synchrotron peak frequency

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The *RoboPol* instrument and the relevant program was developed in order to conduct a systematic study of the optical polarisation variability of blazars. Driven by the discovery that long smooth rotations of the optical polarisation plane can be associated with the activity in other bands and especially in gamma rays, the program was meant to investigate the physical mechanisms causing them and quantify the optical polarisation behaviour in blazars. Over the first three nominal observing seasons (2013, 2014 and 2015) *RoboPol* detected 40 rotations in 24 blazars by observing a gamma–ray-loud and gamma–ray-quite unbiassed sample of blazars, providing a reliable set of events for exploring the phenomenon. The obtain datasets provided the ground for a systematic quantification of the variability of the optical polarisation in such systems. In the following after a brief review of the discoveries that relate to the gamma-ray loudness of the sources we move on to discuss a simple jet model that explains the observed dichotomy in terms of polarisation between gamma–ray-loud and quite sources and the dependence of polarisation and the stability of the polarisation angle on the synchrotron peak frequency.

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1. Introduction

The first detection of a smooth and long rotation of the optical polarisation plane i.e. the electric vector position angle (EVPA), was reportedly found in OJ287 [1]. The attention of the community was however intrigued when [2] reported a similar event in BL Lacertae that was associated with a radio-to-gamma-ray outburst. On the basis of their observation that "a bright feature in the jet" induced "a double flare" observed from optical bands to TeV energies, as well as "a delayed outburst at radio wavelengths", they concluded that the event started in a region with a helical magnetic field configuration. They identified that region "with the acceleration and collimation zone predicted by the theories". The latest brightening of the feature was attributed to its crossing "a standing shock wave" which appeared as the bright radio "core". The importance of the discovery lies mostly on the fact that it offers a way to probe the inner jet of Active Galactic Nuclei in the region where the the jet gets accelerated and collimated in an an environment of coiled magnetic field. Thus, address the very question of the jet formation. Later studies (e.g. [3], [4] etc.) proposed several alternative explanations; all however inevitably based on hand-picked events given the transit character of the phenomenon. The RoboPol program [5], [6] was developed to conduct a systematic study of the phenomenon alongside with quantifying the optical polarisation of blazars and its variability.

Starting from the Second *Fermi* Source Catalog (2FGL) [7] and after excluding 10° on either side of the galactic plane, setting a minimum integrated photon flux at around 10^{-8} cm⁻² s⁻¹ and complying with observational constrains, we composed a sample of 62 gamma–ray-loud (GL) blazars to be monitored (for details c.f. [6]). A comparison gamma–ray-quite (GQ) sample was drawn from the 15 GHz OVRO monitoring sample [8]. The selection was based on a the amplitude of radio variability (> 5 %), the radio mean flux density (at least 60 mJy), the absence from the 2FGL, and again the observability from Skinakas telescope.

In the following we start with a review of angle domain findings that relate with the gammaray activity of the sources. Later we focus our attention on the amplitude domain and examine the dichotomy between GL and GQ sources. Finally we propose a jet model that can explain the observations and especially the dependence of both the degree of polarisation and the uniformity of the angle on the synchrotron peak frequency.

The convention we adopt for the definition of a rotations includes: a continuous EVPA change larger than 90°, comprised by at least four consecutive measurements with significant swings between them.

2. EVPA rotations and gamma-ray activity

The first three seasons of nominal *RoboPol* operation (2013, 2014 and 2015) revealed a total of 40 EVPA rotations that happened in 24 blazars summarised in [9] and [10]. Before robopol a mere total of 16 events had been found in 10 sources. As it is discussed there the EVPA rotations do not seem to be favouring any of the source classes that were examined. High and Low-Synchrotron Peaked (LSP, HSP) sources can show rotations alike as much as BL Lacs and quasars or TeV and non-TeV emitting ones. Interestingly, the rotation rate (deg day⁻¹) can vary noticeably for the same source from one event to the other and so is the sense of rotation. Given the adopted conditions

needed for a sequence of measurements to constitute a rotation, all detected events are in GL sources. This already indicates a relevance with the gamma-ray loudness. Within the uncertainties, there is no indication that for an "orphan" rotation which is unassociated with some activity in gamma rays. There has not been any rotation that occurred with a non-zero lag from the associated gamma-ray event (Blinov et al. submitted). Finally, Figure 8 of [11] shows that the "rotators" are both more luminous and more variable in gamma-rays.

3. The polarisation of GL and GQ sources

Already on the basis of single-shot measurement obtained during the *RoboPol* instrument commissioning, [6] noted that GL sources as a population are significantly more polarised that GQ ones. Using the datasets collected during the 2013 and 2014 observing seasons and after implementing a maximum likelihood analysis that accounts for uneven sampling and measurement uncertainties, [12] confirmed beyond a 4σ level that GL sources are on average more polarised (~ 9 %) than GQ ones (~ 3 %).

4. The jet model

A jet populated by a helical magnetic field and possibly an underlying turbulent one, was thought to be adequate for explaining the observed dichotomy. Occasional mild shocks would locally energise the particles and increase the local field uniformity. The particles in the small volume immediately downstream the shock will emit the high energy photons. Since this emission originates at an environment of increased field uniformity it will be associated with higher polarisation. As the particles propagate downstream losing energy, the associated low energy photons are integrated over larger volumes with different orientations of the field. This naturally being associated with lower polarisation. Hence, high energy emission should be more polarised that low energy one.

In this context, GL which are know to be jet dominated and intensely variable, show often impulsive events of particle acceleration. They hence often reach the necessary high energy limit to produce high energy and high polarisation emission as well as reach gamma-ray detection limits and be classified as GL ones. On the contrary, GQ sources are deprived of such activity. The rarely managed to produce high energy photons and hence are characterised by lower polarisation. In parallel they barely ever pass the detection threshold that would classify them as GL sources.

Though simple, this model includes all the basic components of relativistic jets as they are thought today. Beyond the explanation it offers for the dichotomy of the GL and GQ sources, it also provides a natural explanation of the observed dependence of the polarisation on the synchrotron component peak frequency as we discuss in the next.

5. The polarisation dependence on the synchrotron peak frequency

As we show in Figure 5 of [12], the degree of optical polarisation drops with the frequency of the peak in the synchrotron SED component. So that HSP sources deliver less polarised optical emission compared to sources with their peak of the synchrotron component at much lower energies

(LSP). For the LSP sources the synchrotron peak appears in the vicinity of IR so that the optical emission observed by *RoboPol* constitutes the high energy part of their spectrum. So as we said earlier high energy photons must be associated with higher polarisation. For the HSP sources on the other hand the optical emission corresponds to low energy part of their SED as their synchrotron peak lives in UV or even higher bands. Hence the low energy emission will be characterised by lower polarisation.

This picture further implies a predictable behaviour for the EVPA whose "randomness" should depend also on the synchrotron peak frequency. The topic is discussed immediately.

6. The dependence of the EVPA randomness on the synchrotron peak frequency

As we discussed earlier, the optical emission from HSP sources should be coming from a larger region downstream the shock than in the LSP sources. In a larger region, the helical component of the field would dominate over the turbulent one. Further more, any variability would be happening slower that for a smaller region which would include a smaller number of random field cells. Statistically then, the EVPA of HSPs would tend to align with a preferred direction. That is, a plot of EVPA "stability" would increase with increasing synchrotron peak frequency. Figure 6 in [12] shows that this prediction is indeed seen. Although the data points there do not show a correlation of high significance, the fitted slope is so. For some example cases we show the gradual increase if the EVPA stability with the synchrotron peak position in Figure 1. In the plot only cases with a minimum of 20 measurements are shown. As it becomes apparent, towards higher frequencies the EVPA tends to concentrate around a preferred direction thus its PDF appears less uniform (or equivalently the reduced χ^2 increases).

7. The orientation of the non-uniform EVPAs

According to our model which predicts the dominance of the helical over the turbulent field component in HSP sources and given the small angles to the line of site for blazars, one would expect that the projected magnetic field would orient itself perpendicularly to the jet axis. As a consequence, the EVPA – a proxy of the projected magnetic field – would appear in close alignment to the jet axis in HSPs as the optical emission is in the optically thin part of the synchrotron component.

The best study of this prediction came from [13] (Figures 7 and 8) who conducted careful study of the optical polarisation of high energy BL Lacertae objects in order to identify possible systematic differences with sources that appear in the TeV sky. For a limited sample of sources for which the EVPA as well as the jet orientation was known, they found that indeed six out of nine sources orient their EVPA within 20 degrees from the jet axis implying that the projected magnetic field is indeed perpendicular to the jet. This is the best test of our hypothesis so far as these sources would populate the right-hand part of Figure 1. Unfortunately, the lack of high angular resolution radio observations will make the expansion of this sample challenging. In Figure 2, we show the stacked radio VLBA images from the MOJAVE 15-GHz program [14], for all the sources in Figure 1 with synchrotron peak frequency above 10¹⁵ Hz. The red double-headed arrows there mark the direction of the optical EVPA. For the cases with a clear jet presence a preferentially

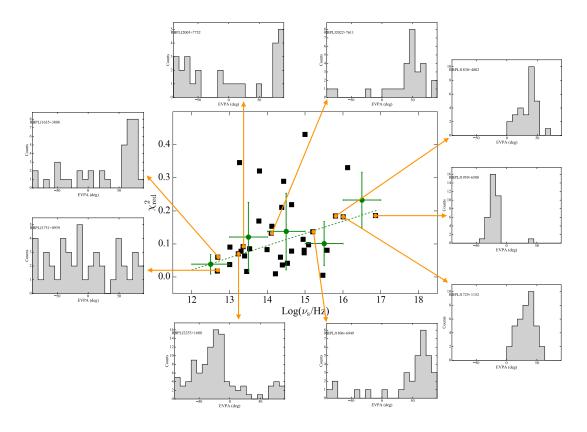


Figure 1: The increase of the EVPA stability as a function of the synchrotron peak frequency. The EVPA stability is measured with the reduced χ^2 of the comparison with a uniform distribution. The satellite panels show the PDFs for some example cases of progressively (from left to right) increasing EVPA stability.

close alignement of the EVPA and the jet axis is observed. Hence, the magnetic field is oriented perpendicular to the jet, in support of our model.

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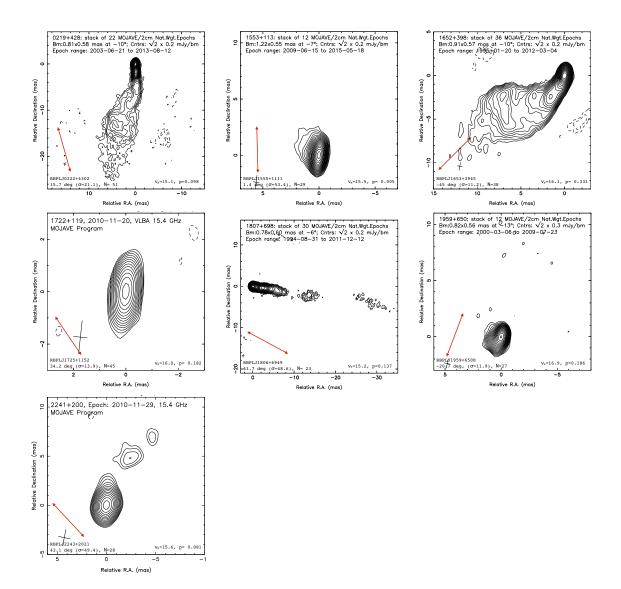


Figure 2: The 15 GHz MOJAVE VLBA images [14] of the sources with $\log(v_s/\text{Hz}) \ge 15$. For those our model forecasts a close alignment of the EVPA and the radio jet. Except for the case of 2241+200 the maps have been subjected to natural weighting. The red double headed arrow denotes the median EVPA. In the lower part of each plot we also give the median EVPA, its sigma and the number of measurements (lower left), along with the logarithm of the synchrotron peak frequency and the median polarisation fraction (lower right).

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