# Planetary system and star formation science with non-redundant masking on JWST

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# ABSTRACT

Non-redundant masking (NRM) is a high contrast high resolution technique that is relevant for future space missions dedicated to either general astrophysics or extrasolar planetary astronomy. NRM mitigates not only atmospheric but instrument-induced speckle noise as well. The recently added mask in the Fine Guidance Sensor Tunable Filter Imager (FGS-TFI) on the James Webb Space Telescope (JWST) will open up a search space between 50 and 400 mas at wavelengths longer than  $3.8\mu$ m. Contrast of  $10^4$  will be achievable in a 10 ks exposure of an M = 7 star, with routine observing, target acquisition, and data calibration methods. NRM places protoplanets in Taurus as well as Jovians younger than 300Myr and more massive than  $2M_J$  orbiting solar type stars within JWST's reach. Stars as bright as M = 3 will also be observable, thus meshing well with next-generation ground-based extreme adaptive optics coronagraphs. This parameter space is inaccessible to both JWST coronagraphs and future 30-m class ground-based telescopes, especially in the mid-IR. We show that NRM used on future space telescopes can deliver unsurpassed image contrast in key niches, while reducing mission risk associated with active primary mirrors.

Keywords: Extrasolar planets, JWST, PSF subtraction, high-contrast imaging

## 1. INTRODUCTION

Landmark discoveries such as dusty disks imaged around young stellar objects, mass-loss shells of evolved stars and the fascinating time-varying spiral plumes surrounding dusty Wolf-Rayet systems have been reported amongst the 50-odd peer-reviewed papers describing results produced by this technique (Ref. 1–15). However, the atmosphere places limitations on non-redundant masking (NRM) performance.

NRM on a space telescope has not been described in the literature to date. We present detailed simulations of NRM performance on the James Webb Space Telescope (JWST) to demonstrate the exciting planetary science enabled by the recent addition of NRM to JWST's suite of established instruments, and to emphasize the importance of this technique to other future space-based observatories. On JWST NRM will widen the telescope's science reach to include a unique combination of wavelength and angular resolution regimes inaccessible from the ground, even with the advent of extreme adaptive optics (ExAO) and coronagraphs behind 30 m extremely

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Figure 1. The parameter space available to the newly-added non-redundant sparse aperture mask on JWST's FGS-TFI instrument. It is designed to operate at 3.8 microns and longer.

large telescopes of the future. For example, NRM used at  $4\mu$ m will place warm extrasolar jovians within 4 to 30 AU of F, G, and K dwarfs 30pc from the Sun within JWST's purview (Figure 1).

Indeed, NRM can add high resolution capability to large filled-aperture space telescopes without sacrificing their wide utility for general astrophysical observations. Furthermore, the NRM search space is complementary to future coronagraphic space missions dedicated to extrasolar planet imaging and characterization, as well as to ground-based ExAO instruments,<sup>16,17</sup> in that it yields moderate contrasts between 0.5 and  $4\lambda/D$  ( $\lambda$  being the observing wavelength, D the telescope diameter). Diffraction-limited stellar coronagraphy, on the other hand, typically covers a search space at higher contrast, at separations larger than ~  $4\lambda/D$  (Ref. 18). Techniques such as Angular or Spectral Differential Imaging (*e.g.*, Ref. 19,20 and references therein) do not work well at these small angular separations. In addition, NRM data has a high degree of *self-consistency* that is absent in full-aperture direct or coronagraphic imaging.

We examine the scientific merit and feasibility of NRM on JWST's Fine Guidance Sensor's Tunable Filter Imager (FGS-TFI),<sup>21</sup> between  $3.8\mu$ m and  $5\mu$ m, operating at a spectral resolution of 100. These contrast predictions are relevant to JWST science planning for its newly-added high angular resolution capability. Our NRM designs work well in 10% bandpass filters, making NRM a possible technique for JWST's near-IR imager NIRCam and mid-IR imager MIRI.<sup>22</sup> Our results can be scaled with both wavelength and optical path difference stability when planning NRM on future missions.

JWST's target acquisition methods, pointing stability, and data calibration pipeline processing are aimed at general purpose imaging with an undemanding operating protocol. They will suffice for the high dynamic range NRM observations we describe here. In contrast, coronagraphic observations require specialized target acquisition and peak-up,<sup>23</sup> very small pointing errors,<sup>24</sup> exquisite wavefront flatness and highly uniform pupil illumination. In comparison, NRM relies on long-term *stability* of the telescope's wavefront during a several minute or hour long exposure, rather than on any stringent requirements on wavefront *quality*. It therefore mitigates some of the risk associated with the complex mechanisms and sequences involved in co-phasing JWST's 18-segment, deployable 6.5 m primary mirror (PM) as the telescope orbits the second Earth-Sun Lagrangian equilibrium point (*e.g.*, Ref. 25).

#### 2. NRM INTERFEROMETRY AND ITS LIMITS

Interferometry with non-redundant baselines was developed for radio astronomy,<sup>26</sup> and subsequently adapted to optical wavelengths.<sup>27</sup> Today it is used in IR and optical bandpasses, often behind instruments not originally developed for sparse aperture interferometry.<sup>1</sup>

Refs. 4,28–30 and references therein present a fuller description of NRM. In brief, a non-redundant array of subapertures is achieved with an N-holes ( $\{h_1, h_2, ..., h_N\}$ ) placed so no vector (baseline) between the centers of two holes is repeated (see Figure 2, left). The resulting PSF is created by several coherent fringe patterns (each with different angular periods on the detector, thanks to non-redundancy) overlaid across one another. A fringe formed by holes  $h_i$  and  $h_j$  is quantified by a complex visibility, which has an amplitude (degree of modulation) and a phase (the offset of the fringe's center from the centroid of the PSF). The fringe visibility is the complex



Figure 2. Left: The FGS-TFI 7-hole non-redundant pupil mask (holes) overlaid on the JWST primary mirror. The secondary mirror supports are shown as three thick black lines. The mask has undersized holes to mitigate against pupil shear. Center: A 4.05  $\mu$ m FGS-TFI image of a point source, with a 1% bandwidth filter, on a square root logarithmic stretch. The pixel scale (65 mas) is that of FGS-TFI. The first Airy ring diameter is 2.4". Generally data out to the first Airy ring is analyzed. Right: the power spectrum of the image (also on a square root scale), showing fringe power on the 21 baselines passed by the mask.

product of a Fourier component of the object and a system visibility. When observing a point-source, this phase  $\phi_{ij}$  measures the piston wavefront error between the two holes. A *closure phase* is formed by the addition of three fringe phases created by a triangle of holes,  $\{h_i, h_j, h_k\}$ , notably  $\phi_{ij} + \phi_{jk} + \phi_{ki}$ . Although piston wavefront errors change the fringe phases, the closure phases are insensitive to these wavefront errors and only measure source structure. There are N(N-1)/2 independent closure phases measurable in the image.

NRM is a form of *imaging* with a PSF that looks unusual. The NRM PSF displays multiple sharp peaks, whereas traditional diffraction-limited PSFs are dominated by one peak (Figure 2, center). However, the NRM PSF possesses some useful properties. First, through the use of closure phase, it is insensitive to large scale wavefront errors over the pupil. Second, its core is more than twice as narrow as a traditional PSF's (the Michelson criterion,  $0.5\lambda/D$ , cf. Rayleigh's criterion,  $1.22\lambda/D$ ). Third, NRM data sets produce results that can be averaged to reduce noise, even in the presence of slowly-varying speckles. By comparison, co-adding PSFs in the presence of speckle noise is always limited by a speckle noise floor.<sup>31–33</sup>

On the ground NRM filters out aberrations that pass through an AO system. It is insensitive to the noncommon path (NCP) aberrations between the science and sensing arms of the instrument. The NCP issue has only recently been addressed by specialized calibration systems in future ExAO coronagraphic systems utilizing full-aperture pupils.<sup>34</sup> Today's ground-based NRM routinely achieves stability of 0.5 degrees on the closure phase, hence passively stabilizing the phase at the level of  $\lambda/500$  to  $\lambda/1000$ .

For a point source any closure phase must be zero. A measured non-zero closure phase in the data provides spatial information on the object. A consistent, systematic, instrument-induced departure from a zero closure phase can be measured using a calibration star, and subtracted from a target's closure phase. Conceptually, closure phase data analysis is a model fit of the intensity distribution on sky, the observations being closure phases from all the closed triangles that can be constructed with the given set of baselines. Fringe visibility data makes solving for the actual image (or its measured Fourier components) possible. The inner working angle (IWA) of such an image is  $0.5\lambda/D$ , the outer working angle (OWA) is set by the shortest non-redundant baseline available. Recovering the measured Fourier components of source structure with NRM imaging using both closure phases and fringe visibilities is a well-posed problem. This is not the case with full aperture imaging. Ground-based fringe visibilities suffer from temporal instabilities due to atmospheric scintillation and transparency variations, though closure phases still extremely useful. On JWST fringe visibilities should be stable since segment reflectivity is unlikely to vary perceptibly during a single exposures, or between science and calibrator exposures.

In practice dynamic range may be limited by effects such as a non-isotropic guiding error, detector flat field errors, and pupil wander. Target placement repeatable to an arcsecond on JWST's HAWAII-2RG detectors will

#### NRM in FGS-TFI projected to PM (V2,V3)



Figure 3. The pupil plane in FGS-TFI is decentered, demagnified, stretched slightly in one dimension, and slightly rotated relative to JWST's primary mirror (PM). The NRM in the JWST FGS-TFI pupil wheel is shown in PM coordinates, after being projected back to the PM using JWST's optical design. Each of the 18 segment centers are indicated with light green circles labelled with their segment names (A1,A2,..., C6). The conceptual design NRM with holes placed in a regular hexagonal grid in (V2,V3) space is shown by small black solid circles at the nominal segment centers' locations (*i.e.*, assuming segments form a regular hexagonal grid in (V2,V3) space). The projected mask's holes in PM space are the 7 green hexagons. Slight misalignments of the holes with their respective PM segment centers will not affect the scientific capabilities of the NRM. However it is preferable that the NRM's holes do not span a segment-to-segment discontinuity. Undersizing of the mask's holes relative to the segments themselves allows for pupil shear up to 3.8% of the pupil diameter before the mask's holes cross segment boundaries.

improve data calibration, given the spatial frequency of the detector' flat field structure.

#### 3. NRM IN SPACE

Without AO, ground-based NRM, like speckle interferometry, must freeze temporally-varying fringe patterns. AO stabilizes relative piston differences between subapertures in the aperture mask. This enables longer exposure times, limited by pixel well-depth, thermal background rates or instrument stability. Once deployed, JWST's primary mirror segments' positions relative to each other are expected to drift very slowly. Exposure times will likely be limited by cosmic rays, with a maximum single exposure time measured in hours. Under such conditions, NRM opens up a search space completely beyond the reach of ground-based telescopes. Thermal background and atmospheric opacity also limit ground-based imaging longward of about  $2\mu$ m. Furthermore, JWST's unusual primary mirror geometry and articulation are unlikely to cause problems for NRM interferometry.

For a  $0.5m^2$  subaperture at  $4\mu m$ , a 1% filter bandwidth and 50% net quantum efficiency,  $M \simeq 7$  is the limit for  $10^4$  contrast in 10ks, and M = 12 the limit for  $10^3$  contrast in 10ks. A 10% bandwidth filter will reduce exposure times by a factor of ten for the same contrast. At bandwidths in excess of 10% the 7-hole FGS-TFI designs' contrast starts to degrade. In FGS-TFI, sub-array readouts will enable observations of objects as bright as M = 3.

#### 4. SIMULATIONS

We used ten Monte Carlo realizations of JWST's exit pupil wavefront that conform to JWST's wavefront error budget.<sup>39</sup> These OPDs,<sup>40</sup> provided with the JWPSF<sup>41</sup> software distribution, represent the wavefront in JWST NIRCam's short-wavelength channel. They assume particular low-, mid-, and high-spatial frequency wavefront



Figure 4. Predicted JWST NRM and coronagraph performance. Dynamic range using only closure phase (CLP) as well as closure phase in combination with fringe visibility (CLP+VIS) data for 1% bandpass NRM imaging at 4.81 $\mu$ m (using the 7-hole mask shown in Figure 1) in JWST's FGS Tunable Filter Imager (TFI) is plotted (heavy curves). A 1 day delay between the NRM target and calibrator stars, with worst-case attitude-dependent thermally induced primary mirror configuration changes (breathing), guiding errors, and measured detector intra-pixel sensitivity are simulated. In a 10 ks exposure on an M = 7 target, photon noise sets a 10 magnitude dynamic range limit. 10% bandwidths will produce similar results in 1ks. The range of estimated contrasts between a solar type star and 1 and 10 Jupiter mass planets, at ages of 1 and 10 Myr (Ref. 35, 36) are shown (vertical blue lines placed at arbitrary separations). This mask will enable JWST to detect protoplanets in Taurus. The same mask in JWST NIRCam will produce 1 to 2 magnitudes less contrast at 2  $\mu$ m, at half the corresponding angular separation. In MIRI NRM can produce even higher contrast than in FGS-TFI. NRM performance at 4.8  $\mu$ m drops outside about 600mas. NIRCam's linear band-limited coronagraph performance using planned JWST roll angles<sup>37</sup> complements the FGS-TFI NRM search space. The JWST MIRI four quadrant phase mask coronagraph 11  $\mu$ m contrast curve,<sup>38</sup> assuming a static 5mas pointing error between target and calibrator, no guiding error, unrestricted telescope roll capability, and no error in calibration star placement is also plotted. Jovian planets are brighter at this wavelength than at 3-5  $\mu$ m.



Figure 5. NRM detection probabilities for FGS-TFI at 4.44  $\mu$ m, drawing from a large sample of young nearby stars. Companion mass is cut off above 10 Jupiter masses. Contours show the estimated probability of finding a planet with a given mass (in Jupiter masses, log scale), on the horizontal axis, and a given semi-major axis (AU, vertical axis) within nominal performance (a 10 magnitude contrast floor imposed on the curves in Fig. 4). The exposure time assumed is about 1 hour per source. Better detection probabilities would result from longer exposures on an individual object.

error allocations to JWST's Optical Telescope Element (OTE) and Integrated Science Instrument Module (ISIM). These realizations yield PSFs at  $2\mu$ m with Strehl ratios above 80%. They are estimates, given the expected residual wavefront error 14 days after active tuning of the JWST primary mirror (PM) figure has been performed, and the telescope slewed from a spacecraft attitude with the most thermal load from the Sun to one with the least such thermal load. To preserve the individual segment wavefront errors associated with one realization of the 18 segments in JWST's PM, we extracted the individual segment piston errors from each of the ten wavefront maps, and used them to create ten instances of a single thermally perturbed set of PM segments.

Using methods described in Refs. 42–45, we simulated images on a 9-fold finer grid than the 65 mas FGS-TFI detector pixel pitch. We dithered the oversampled PSFs by one sample spacing to model pointing errors of 7mas per axis. Intra-pixel quantum efficiency was modeled as a quadratic, ranging from unity at a pixel center to 0.8 at a pixel corner. We then binned the  $9 \times 9$  dithered subsampled PSFs to create an image on the FGS-TFI pixel scale. We did not explicitly simulate photon noise, read out noise (~  $15e^-$  rms for a double-correlated sample), or detector flat field errors, although our detection threshold calculations have taken some experience-based account of these processes.

Simulated data were analyzed with a pipeline used for Keck, Palomar, and VLT data. (e.g., Ref. 46). Nine of the ten PSF realizations in a set were taken to be calibrator observations, and one of the ten was taken to be the target observation. This process was repeated for several choices of target PSF. Rather than replicating the complex Monte-Carlo process of Ref. 46, we chose to use a simple detection threshold of  $5\sigma$ , with closure-phase and visibility standard deviations calculated from scatter within the 9 calibrator observations. Note that typical 99.9% detection thresholds in Ref. 46 were ~4 to  $4.5\sigma$  with a more complete Monte-Carlo, which was verified by the lack of numerous false detections in that paper. We therefore believe our  $5\sigma$  limit is conservative.

Figure 4 shows the contrast achieved with dynamic range calculated two ways: first, using only the closure phases (labelled CLP) for a comparison with current ground-based results, and second, using both closure phases and fringe visibilities (labelled CLP+VIS), to estimate the full range of contrast available to JWST. Contrast between  $M_J$  and  $10M_J$  planets 1 and 10 Myr old and a parent G2V dwarf are shown in the figure. We find our mask designs produce these estimated contrasts even when the detector's pixel pitch is slightly worse than



Figure 6. Distribution of stellar ages with distance in the sample used to generate the detection probability shown in Fig. 5. The sample comprises stars younger than 5Myr and closer than 140pc from the Sun, and stars older than 10Myr that are within 50pc of the Sun.

Nyquist. The very highest contrast shown is obtained using both closure phases as well as visibilities.

Dithering and drizzling images<sup>47</sup> may further enable undersampled data recovery. Information from routine wavefront sensing observations<sup>45</sup> may help improve the dynamic range of NRM data. We have not examined these two refinements.

### 5. BENEFITS OF NRM IN SPACE

NRM brings both scientific and operational advantages to JWST, given the telescope's segmented architecture and planned co-phasing methods during commissioning.<sup>48–50</sup> Every NRM image measures the relative piston between each hole in the aperture mask, with a capture range set by the filter coherence length. Images taken through two masks with  $\sim 5\%$  throughput can measure segment pistons and tilts to interferometric precision, without the need for any alteration of the optical train, whether they be defocus, segment tilt, or indeed any other active mirror motion. Coarse co-phasing JWST can be accomplished using 1% bandpass filters and NRMs. Masks in separate cameras mitigate instrument failure-induced risk, and enables campaign mode observing with the best image quality possible in any camera equipped with two NRM's. Two such masks in a camera can measure its field-dependent and chromatic aberrations. Camera-specific wavefront knowledge can feed into data reduction, benefiting science that requires thorough understanding of temporal, spatial, and chromatic variations in the PSF. In addition, NRM observations determine stellar multiplicity at the highest resolution and contrast possible. Such observations can eliminate inappropriate wavefront sensing or guide star choices. Images taken as the mask is stepped across the pupil (or vice versa) will measure pupil location without specialized pupil imaging optics. Such information is relevant for IR instrument and telescope maintenance.

#### 6. CONCLUSION

Detailed simulations with time-varying mirror figure errors and existing data reduction methods suggest that placing a non-redundant aperture mask on any of the JWST instruments brings exciting high resolution high contrast imaging within reach. The 7-hole NRM in the Fine Guidance Sensor's Tunable Filter Imager will image protoplanets in Taurus (see Figure 5 and 6). Such aperture masking provides JWST with alternative, risk-reducing wavefront measuring techniques, while increasing the science output of JWST substantially, and paving the way for future missions with NRM instrumentation. Future missions can increase the science payoff with more optimized NRM modes, making the technique interesting for galactic, extragalactic, and cosmological observations.

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