

# FMOS near-IR spectroscopy of *Herschel*-selected galaxies: star formation rates, metallicity and dust attenuation at $z \sim 1$

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## ABSTRACT

We investigate the properties (e.g. star formation rate, dust attenuation, stellar mass and metallicity) of a sample of infrared (IR) luminous galaxies at  $z \sim 1$  via near-IR spectroscopy with Subaru-FMOS. Our sample consists of *Herschel* SPIRE and *Spitzer* MIPS selected sources in the COSMOS field with photometric redshifts in the range of  $0.7 < z_{\text{phot}} < 1.8$ , which have been targeted in two pointings ( $0.5 \text{ deg}^2$ ) with FMOS. We find a modest success rate for emission-line detections, with candidate  $H\alpha$  emission lines detected for 57 of 168 SPIRE sources (34 per cent). By stacking the near-IR spectra we directly measure the mean Balmer decrement for the  $H\alpha$  and  $H\beta$  lines, finding a value of  $\langle E(B - V) \rangle = 0.51 \pm 0.27$  for  $\langle L_{\text{IR}} \rangle = 10^{12} L_{\odot}$  sources at  $\langle z \rangle = 1.36$ . By comparing star formation rates estimated from the IR and from the dust-uncorrected  $H\alpha$  line we find a strong relationship between dust attenuation and star formation rate. This relation is broadly consistent with that previously seen in star-forming galaxies at  $z \sim 0.1$ . Finally, we investigate the metallicity via the N2 ratio,

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finding that  $z \sim 1$  IR-selected sources are indistinguishable from the local mass–metallicity relation. We also find a strong correlation between dust attenuation and metallicity, with the most metal-rich IR sources experiencing the largest levels of dust attenuation.

**Key words:** galaxies: evolution – submillimetre: galaxies.

## 1 INTRODUCTION

Accurate measurements of the characteristic properties of galaxies, star formation rate (SFR), metallicity and stellar mass are central to our understanding of their evolution. Significant progress in our ability to measure these properties in distant ( $z > 1$ ) galaxies has been made in the last two decades. Deep surveys with the *Hubble Space Telescope* (*HST*) have enabled the SFRs of large numbers of galaxies up to  $z \sim 7$  to be measured (e.g. Madau et al. 1996; Bunker et al. 2004; Bouwens et al. 2006, 2009; McLure et al. 2010). Large-scale optical and near-infrared (IR) spectroscopic surveys have targeted key emission lines allowing the study of gas metallicity to  $z < 4$  (e.g. Tremonti et al. 2004; Erb et al. 2006; Mannucci et al. 2009; Zahid, Kewley & Bresolin 2011; Cresci et al. 2012). Finally, deep optical and near-IR photometric surveys have allowed an accurate assessment of the stellar mass contained in galaxies, and its build-up with redshift out to  $z \sim 5$  (e.g. Fontana et al. 2006; Ilbert et al. 2010; Caputi et al. 2011).

While these advances have reshaped our understanding of galaxy formation and evolution, they typically rely observations in a single wavelength window, i.e. ultraviolet (UV), optical/near-IR, far-IR, etc. Meanwhile, at low and intermediate redshifts it is becoming clear that large-scale, multiwavelength studies of galaxies are needed to determine unbiased estimates of their properties. In the case of SFR estimates, where the corrections for dust attenuation tend to be large, comparisons of UV and IR SFR estimates at  $z \sim 0$  (Hao et al. 2011),  $z \sim 1$  (Buat et al. 2010) and  $z \sim 2$  (Reddy et al. 2010) show that widely used *in band* (i.e. in the same waveband as the SFR estimate) dust attenuation estimators (e.g. the UV continuum slope; Meurer, Heckman & Calzetti 1999) have large errors ( $\delta \log_{10} \text{SFR} \sim 0.3$  dex) and can have significant systemic biases for certain populations of galaxies [low SFR spirals, ultraluminous infrared galaxies (ULIRGs)  $10^{12} L_{\odot}$  and young starbursts]. By comparison SFR estimators based on combinations of the far-IR,  $\text{H}\alpha$  line or radio have much smaller errors ( $\delta \log_{10} \text{SFR} \sim 0.1$  dex; Kennicutt et al. 2009; Hao et al. 2011) and are universally valid.

Measurements of other galaxy properties also benefit from a multiwavelength approach. Metallicity estimates from single tracers e.g. the  $N2$  or  $R_{32}$  methods (Pettini & Pagel 2004) can disagree by up to  $\Delta[\log_{10}(\text{O}/\text{H})] = 0.7$  dex (Kewley & Ellison 2008). Stellar mass estimates obtained via the fitting the stellar population models require multiband observations in the optical and near-IR to be reliable; omitting near-IR observations introduces an error of  $\sim 0.1$  dex to stellar mass estimates at  $z \sim 1$  (Pozzetti et al. 2007; Ilbert et al. 2010).

In order to put galaxy evolution at high  $z$  on a firm footing multiwavelength observations at the same rest-frame wavelengths as our low- $z$  benchmarks [e.g. Sloan Digital Sky Survey (SDSS), *Spitzer* and *IRAS*] for a large number of high- $z$  galaxies is needed. This necessitates wide-field imaging and spectroscopy in the IR.

Here we investigate the rest-frame optical-to-far-IR properties of a sample of *Herschel*<sup>1</sup> (Pilbratt et al. 2010) sources which were targeted for near-IR spectroscopy with FMOS (Kimura et al. 2010). The key goal of this work is to determine the key galaxy properties (SFR, dust attenuation, stellar mass and metallicity) between a sample of high- $z$  ( $0.8 < z < 1.7$ ), IR luminous ( $> 10^{11} L_{\odot}$ ) sources using the same tracers commonly used for low- $z$  samples. In this way we can be sure that our results are fully consistent (in terms of both calibration and selection effects) with those at low  $z$ . The data sets used in this work are described in Section 2, Section 3.1 presents the detection rate of  $\text{H}\alpha$ , Section 3.2 presents the aggregate near-IR spectral properties and Section 3.3 presents a comparison of the SFRs from the IR and  $\text{H}\alpha$  line. In Section 3.4 we investigate the stellar mass and metallicity of our sample and, finally, Section 4 summarizes our conclusions. Throughout we assume a  $\Lambda$  cold dark matter ( $\Lambda$ CDM) cosmology with  $\Omega_{\Lambda} = 0.7$ ,  $\Omega_{\text{m}} = 0.3$  and  $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$ .

## 2 DATA

### 2.1 Pre-existing COSMOS data

The starting point for this work is the SPIRE observations of the COSMOS field (Scoville et al. 2007) taken as part of the *Herschel* Multi-tiered Extragalactic Survey (HerMES; Oliver et al. 2012). The SPIRE instrument, its in-orbit performance and its scientific capabilities are described by Griffin et al. (2010); its calibration methods and accuracy are outlined in Swinyard et al. (2010). Here we make use of SPIRE maps as described in Levenson et al. (2010). At the time of writing HerMES observations of COSMOS cover  $\sim 4.8 \text{ deg}^2$  to a  $1\sigma$  instrumental noise of  $\sim 2 \text{ mJy beam}^{-1}$  at the three SPIRE wavelengths of 250, 350 and 500  $\mu\text{m}$ .

As the SPIRE data offer an instrumental noise significantly lower than the confusion noise ( $\sim 6 \text{ mJy}$ ; Nyugen et al. 2010), we make use of prior source positions from higher angular resolution data to extract SPIRE photometry. The MIPS 24  $\mu\text{m}$  channel is the most obvious prior for SPIRE data as it offers a significant improvement in angular resolution [6 arcsec full width at half-maximum (FWHM) for MIPS 24  $\mu\text{m}$  versus 18.6 arcsec FWHM for SPIRE 250  $\mu\text{m}$ ] while also being able to account for  $> 80$  per cent of the flux at SPIRE wavelengths at the 24  $\mu\text{m}$  depths now available in a large fraction of HerMES fields (Bethersin et al. 2012; Oliver et al. 2012).

To construct our prior catalogue for SPIRE photometry we begin with the MIPS 24  $\mu\text{m}$  imaging from the *Spitzer* COSMOS survey (Le Flocc’h et al. 2009). Here we make use of the publicly available imaging, performing source extraction via the `STARFINDER IDL`

<sup>1</sup> *Herschel* is an ESA space observatory with science instruments provided by Principal Investigator consortia. It is open for proposals for observing time from the worldwide astronomical community.

package (Diolaiti et al. 2000). The resulting catalogue covers  $\sim 2.1$  deg<sup>2</sup> and has a typical  $1\sigma$  sensitivity of  $\sigma = 15$   $\mu$ Jy.

In order to provide the most accurate positional information, for both our SPIRE photometry and FMOS fibre positioning, we cross-match our 24  $\mu$ m catalogue to the publicly available *HST* Advanced Camera for Surveys (ACS)  $I_{F814W}$ -band catalogue of Leauthaud et al. (2007). This catalogue covers 1.64 deg<sup>2</sup> to a limiting magnitude of  $I_{F814W} < 26.5$ . In addition to improving the positional accuracy this matching helps eliminate spurious 24  $\mu$ m sources produced by artefacts in the image, in particular those located close to bright ( $> 1$  mJy) sources. Of the 35 914 24  $\mu$ m sources located well within the ACS  $i$ -band coverage, 33 071 (92 per cent) have  $I_{F814W}$ -band counterparts within 2 arcsec.

SPIRE photometry is performed using the  $I_{F814W}$ -band positions of the 24  $\mu$ m sources as a prior, following the algorithms described in Roseboom et al. (2010, 2012). All  $> 3\sigma$  ( $\sim 60$   $\mu$ Jy) 24  $\mu$ m sources are considered as potential SPIRE counterparts. Using the residual map statistics we estimate that our prior-driven SPIRE catalogue reaches a typical point-source sensitivity of  $\sigma_{\text{tot}} = 2.7, 3.5$  and 3.2 mJy at 250, 350 and 500  $\mu$ m, including the contribution from source confusion.

To complete our multiwavelength COSMOS data set we add multiband optical/near-IR data and photometric redshifts from the catalogue of Ilbert et al. (2009). This catalogue is limited to  $i_{\text{AB}}^+ < 25$ , and hence our HerMES-COSMOS sample is similarly restricted.

## 2.2 FMOS observations and emission-line measurements

IR-selected sources were targeted in two pointings (0.5 deg<sup>2</sup>) located within the COSMOS field with FMOS as part of the Guaranteed time observations program. The FMOS instrument (Kimura et al. 2010) consists of 400 1.2 arcsec diameter fibres which can be placed within a 30 arcmin diameter field of view. We used the low-resolution mode ( $R \sim 600$ ), allowing instantaneous coverage of both the  $J$  and  $H$  band ( $0.9 < \lambda < 1.8$   $\mu$ m), with cross-beam switching, i.e. two fibres for each target: one placed on the sky and one on the target, with the target/sky ‘switched’ between them at regular intervals.

Potential targets for FMOS fibre allocation were selected from our HerMES-COSMOS parent catalogue by requiring a photometric redshift in the range of  $0.65 < z_{\text{phot}} < 1.75$  from the catalogue of Ilbert et al. (2009). This restriction was introduced to ensure that the  $H\alpha$  line was likely within the FMOS wavelength coverage.

Fibre allocation preference was given to sources detected at both 24 and 250  $\mu$ m ( $> 3\sigma_{\text{tot}}$ ), followed by 24  $\mu$ m only sources. As well as science targets, a number (typically 2–4) of Two Micron All Sky Survey (2MASS; Skrutskie et al. 2006) selected stars were included in the observations for flux calibration purposes.

The first of our FMOS pointings was dedicated to solely HerMES-COSMOS targets, while for the second pointing (2010 November 24 and 25) we shared fibres with the evolSMURF project (Bunker et al., in preparation). While the split between the samples was roughly 50–50, this was aided by the overlap between the samples (42 sources).

Table 1 details the exposure times, together with the number of 24 and 250  $\mu$ m detected (henceforth referred to as 24  $\cap$  250  $\mu$ m), and 24  $\mu$ m only sources in each pointing. In total 241 fibres were allocated to IR-selected sources, with four sources appearing in both configurations, resulting in 237 unique targets (168 unique 24  $\cap$  250  $\mu$ m targets). All data were reduced using the standard FMOS pipeline (Iwamuro et al. 2011).

**Table 1.** Summary of FMOS observations.

	Date	$T_{\text{exp}}$	$N_{24 \mu\text{m}}$	$N_{24 \cap 250 \mu\text{m}}$
P1	2010 November 22	$8 \times 900$ s	136	102
P2	2010 November 24 and 25	$14 \times 900$ s	105	67
	Total (unique)		241 (237)	169 (168)

Emission lines were identified in the 2D-reduced frames, after flux calibration, via a semi-automated procedure. At each pixel the line profile was fitted to the surrounding  $9 \times 9$  pixel. The pixel scale was 5  $\text{\AA}$  in the spectral direction and 0.13 arcsec in the fibre direction. We only considered pixels in the wavelength ranges 1.1–1.36 and 1.42–1.7  $\mu$ m. Pixels within 5  $\text{\AA}$  of an OH line were excluded from consideration. The line profile was assumed to be Gaussian with  $\text{FWHM} = \lambda/600$   $\text{\AA}$  in the spectral direction, and 6.9 pixel in the fibre direction. The noise was estimated by taking the variance of all illuminated pixels at that wavelength on the detector. Regions where the noise is exceptionally high ( $> 10$   $\mu$ Jy pixel<sup>-1</sup>) were excluded. The local continuum was estimated by taking the median pixel value in a window of 20 pixel (200  $\text{\AA}$ ), excluding the closest 7 pixel. For each fit the line signal-to-noise ratio (S/N), peak S/N and correlation coefficient between the line profile and 2D spectrum were measured. The line S/N ( $\sigma_{\text{line}}$ ) was calculated as

$$\sigma_{\text{line}} = \sum_i (d_i - c_i) P_i / \sigma_i^2 \left[ \sum_i (P_i^2 / \sigma_i^2) \right]^{-1/2}, \quad (1)$$

where  $d_i$  is the pixel intensity at position  $i$ ,  $c_i$  is the continuum at pixel  $i$ ,  $P_i$  is the line profile at position  $i$  and  $\sigma_i$  is the noise estimate at position  $i$ .

The peak S/N is defined as the ratio of the peak flux density, taken to be the mean flux density in a  $3 \times 3$  pixel window less the local continuum, to the standard deviation of the surrounding pixels in the spectrum.

Finally, the correlation coefficient ( $\rho_{\text{line}}$ ) is calculated as

$$\rho_{\text{line}} = \frac{\sum_i (d_i - \bar{d})(P_i - \bar{P})}{\sigma_P \sigma_d}, \quad (2)$$

where  $\sigma_P$  and  $\sigma_d$  are the standard deviation of the line profile and data values, respectively.

Line fits which have  $\sigma_{\text{line}} > 4$ ,  $\sigma_{\text{peak}} > 2.5$  and  $\rho_{\text{line}} > 20$  were considered as candidate emission lines.

All candidate emission lines that have a wavelength within the range  $(1 + z_{\text{phot}} - 0.16) \times 6563.4 < \lambda < (1 + z_{\text{phot}} + 0.16) \times 6563.4$   $\text{\AA}$  were considered to be  $H\alpha$ . The window of  $\delta z = 0.16$  equates to  $4\sigma_{\text{phot-z}}$  for the typical photo- $z$  error quoted by Ilbert et al. (2009) at  $z \sim 1$ .

A total of 85 candidate  $H\alpha$  emission lines were found from the sample of 237 unique 24  $\mu$ m targets. We assessed the reliability of our line identification technique in two ways. First the line identification was repeated, but with the proposed redshifts (and hence wavelength search window) shifted. To ensure that the mock search windows are sufficiently far away from real lines, but still within the wavelength coverage of FMOS, sources at  $z_{\text{phot}} < 1.1$  were given  $z_{\text{mock}} = z_{\text{phot}} + 0.16 + \delta$ , while those at  $z_{\text{phot}} > 1.1$  were set to  $z_{\text{mock}} = z_{\text{phot}} - 0.15 - \delta$ , where  $\delta$  is a random number between 0 and 0.15. As a result of this process seven lines were identified, giving an estimate of the false positive line detection rate of 7/85 or  $8 \pm 3$  per cent. No false lines were returned with  $\sigma_{\text{line}} > 8$ .

We compare the redshifts, as determined by the wavelength of our candidate  $H\alpha$  lines, and the known spectroscopic redshifts. From our sample of 85  $H\alpha$  line emitters, 28 are found to have reliable

**Table 2.** Completeness of line identification process as a function of flux. Completeness is estimated via injection of mock emission lines into the 2D spectra at random wavelengths/redshifts. Note that this is for flux contained within the 1.2 arcsec diameter fibre of FMOS; no aperture effects are considered.

Line flux ( $\times 10^{-16}$ erg cm $^{-2}$ s $^{-1}$ )	Completeness	
	P1 (per cent)	P2 (per cent)
0.5	4.5	1.4
1.	23.2	25.3
1.5	52.9	57.8
2.	65.8	67.8
2.5	71.0	72.1
3.	74.2	74.7
3.5	77.4	77.6
4.	77.6	78.7

( $z_{\text{qual}} > 3$ ) spectroscopic redshifts in the zCOSMOS bright ( $i < 22.5$ ) sample of Lilly et al. (2007). Of these, 27 (96 per cent) are found to be within  $\delta z = 0.01$  of our assumed H $\alpha$  redshift. No incorrect redshifts are found amongst the 11 sources with candidate H $\alpha$  lines at  $\sigma_{\text{line}} > 8$  and a spectroscopic redshift from Lilly et al. (2007). While this result is encouraging, it is likely that the reliability of our line identification is a function of brightness, and the zCOSMOS bright sample is limited to  $i_{\text{AB}}^+ < 22.5$ ; roughly 54 per cent (46/85) of our candidate H $\alpha$  line emitters have  $22.5 < i_{\text{AB}}^+ < 25$ .

Via a similar process we can estimate the completeness of our line identification process. Assuming the random redshifts,  $z_{\text{mock}}$ , described above, we inject mock emission lines into our data at a wavelength corresponding to  $1 + z_{\text{mock}} 6564.3 \text{ \AA}$ . We estimate the completeness for the two pointings independently. Table 2 details the completeness (i.e. the ratio of sources detected to those injected) as a function of flux. No aperture correction is assumed, injected line fluxes are considered to be those contained within the 1.2 arcsec diameter fibre of FMOS. The completeness never reaches 100 per cent as lines at certain wavelengths will always be undetectable due to the gap in wavelength coverage from 1.36 to 1.42  $\mu\text{m}$ , as well as the masking of OH sky lines. In total 125 OH lines are suppressed which, combined with the gap due to atmospheric absorption, remove 24 per cent of the potential wavelength coverage. It can be seen that our identification process reaches this maximum level above a line flux of  $f_{\text{H}\alpha} \gtrsim 4 \times 10^{-16} \text{ erg cm}^{-2} \text{ s}^{-1}$ , while we are  $\gtrsim 50$  per cent complete above a line flux of  $f_{\text{H}\alpha} \gtrsim 1.5 \times 10^{-16} \text{ erg cm}^{-2} \text{ s}^{-1}$ . The completeness estimates for the second pointing are marginally higher at  $f_{\text{H}\alpha} \lesssim 3 \times 10^{-16} \text{ erg cm}^{-2} \text{ s}^{-1}$  due to the increased exposure time (12 600 s for the second pointing versus 7200 s for the first). Note that these line fluxes consider only the flux contained within the 1.2 arcsec diameter fibres used on FMOS; no aperture correction has been applied yet.

For each H $\alpha$  line emitter we attempt to measure the flux of the neighbouring [N II] 6584 line. Line fluxes are estimated in the same manner as H $\alpha$ , but with the central wavelength fixed at  $(1 + z) 6584 \text{ \AA}$ . Of the 85 H $\alpha$  line emitters, 33 also have [N II] at S/N  $> 3$ .

Finally, both H $\alpha$  and [N II] line fluxes are corrected for the limited aperture and unknown stellar continuum. The aperture correction is determined via the ratio of  $I_{F814W}$  flux within the FMOS fibre (1.2 arcsec diameter) to that within the Kron radius. The typical aperture correction is  $\sim 2$ –3. For the continuum correction, the near-IR spectra are not deep enough to detect the continuum emission

near the H $\alpha$  line. Thus, continuum emission at the wavelength of the line is estimated, and removed, using the available broad-band optical and near-IR imaging from the COSMOS survey (Capak et al. 2007; McCracken et al. 2010). For sources where the line lies at  $\lambda < 1.4 \mu\text{m}$  we use the  $J$ -band magnitude, and the  $[z^+ - J]$  colour to estimate the continuum flux. For those lines where  $\lambda > 1.4 \mu\text{m}$  the  $K$ -band magnitude and  $[J - K]$  colour are used. In all cases we assume that the H $\alpha$  line is coincident with stellar absorption of equivalent width  $\text{EW} = 4.4$  (Moustakas & Kennicutt 2006).

Appendix A contains the line fluxes and photometric properties for our 85 H $\alpha$  line emitters. While we list the full catalogue of H $\alpha$  line detections resulting from our observations, in the following analysis we consider only the 57 H $\alpha$  line emitters which have robust detections at 250  $\mu\text{m}$ . This restriction is implemented as 24  $\mu\text{m}$  alone is not a good tracer of total IR luminosity at these redshifts (Elbaz et al. 2010). Enforcing 250  $\mu\text{m}$  detections also helps to minimize contamination by active galactic nucleus (AGN) activity (Hatziminaoglou et al. 2010).

### 2.3 SDSS comparison sample

To establish a low-redshift baseline for our  $z \sim 1$  measurements we identify a sample of *Spitzer* 160- $\mu\text{m}$ -selected sources from the SWIRE survey (Lonsdale et al. 2003) in the Lockman Hole field with SDSS Data Release 6 (DR6) (Adelman-McCarthy et al. 2008) spectroscopy. Sources are required to be brighter than 50 mJy at 160  $\mu\text{m}$  and to lie at  $z < 0.4$ . The 160  $\mu\text{m}$  photometry is complemented with SWIRE MIPS 24, 70  $\mu\text{m}$  data, and SPIRE photometry at 250, 350 and 500  $\mu\text{m}$  from HerMES. H $\alpha$  line fluxes are calculated using the GANDALF package (Sarzi et al. 2006), with aperture corrections based on the ratio of the SDSS Petrosian-to-fibre  $r$ -magnitudes. We retain sources with a peak line flux density to noise ratio of greater than 3. The typical limiting H $\alpha$  line flux ( $3\sigma$ ) is  $5.8 \times 10^{16} \text{ erg cm}^{-2} \text{ s}^{-1}$ . Our final low- $z$  sample consists of 156 160- $\mu\text{m}$ -detected sources with reliable H $\alpha$  line flux measurements, at a mean redshift of  $z = 0.1$ .

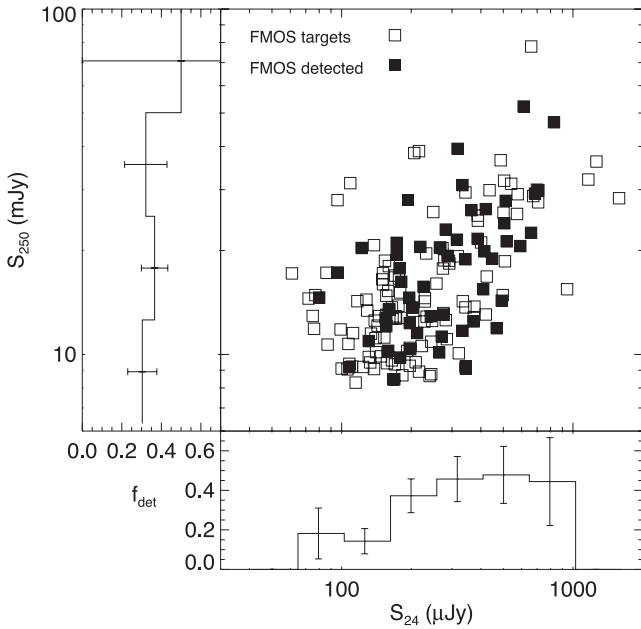
## 3 RESULTS

### 3.1 H $\alpha$ detection rate for IR-selected sources

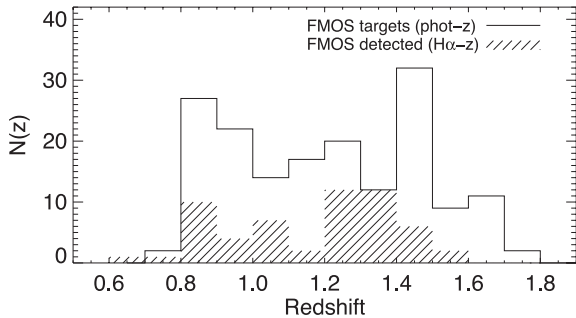
The raw H $\alpha$  detection rate for IR-sources ( $24 \cap 250 \mu\text{m}$ ) is 57/168 ( $34 \pm 4$  per cent). In Fig. 1 we present the H $\alpha$  detection rate as a function of 250 and 24  $\mu\text{m}$  flux densities. While the H $\alpha$  detection rate appears to be insensitive to 250  $\mu\text{m}$  flux density, a modest increase in the detection rate is seen for sources with  $S_{24 \mu\text{m}} > 200 \mu\text{Jy}$ .

Fig. 2 shows the redshift distribution of IR sources targeted with FMOS, and those with H $\alpha$  line detections. The contrast between the two distributions highlights the visibility of H $\alpha$  with FMOS as a function of redshift. Prior to FMOS observations, only photometric redshifts were available for the vast majority of targets. Thus while targets were photo- $z$  selected to be in the redshift range where H $\alpha$  is visible, some fraction of sources were expected to have H $\alpha$  fall at wavelengths outside the FMOS wavelength range.

Despite this the observed detection rate compares well with what would be predicted given the known completeness (Section 2.2). For each detected source we calculate the completeness at that line flux,  $c_i$ , by interpolating the values in Table 2. If the H $\alpha$  line fluxes of the undetected sources are distributed similar to the detected ones (i.e. the bulk of the incompleteness is due to OH sky lines), then the sum  $\sum_i 1/c_i$  should be equal to the number of targeted objects. For



**Figure 1.**  $H\alpha$  line detection rate of FMOS targeted sources as a function of 24 and 250  $\mu\text{m}$  flux density. The detection rate is relatively insensitive to 250  $\mu\text{m}$  flux density, while a modest gain is seen for sources with  $S_{24\ \mu\text{m}} > 200\ \mu\text{Jy}$ .



**Figure 2.** Redshift distribution of FMOS targeted sources and sources with  $H\alpha$  line detections. For FMOS targets we use photometric redshifts from Ilbert et al. (2009). For detected sources the redshift implied by the location of the  $H\alpha$  line is used.

our  $24 \cap 250\ \mu\text{m}$  sample  $\sum_i 1/c_i = 177$ , close to the actual number of sources targeted (168).

We can compare our detection rate for IR-selected sources with FMOS to optically selected sources with Very Large Telescope (VLT) VIMOS. For the VVDS-DEEP survey (Le Fèvre et al. 2005), the detection rate for  $i$ -selected sources with similar magnitudes (median for our sample is  $i_{AB}^+ \sim 23.5$ ) and redshifts ( $z \sim 1.2$ ) to our sample is  $\sim 70$  per cent (for 4.5 h exposures; Ilbert et al. 2005). While our FMOS detection rate is almost one-half of this, IR sources have quite a low areal density and cannot make the most of the large multiplex of VIMOS. For our parent sample of 24 and 250  $\mu\text{m}$  detected,  $i_{AB}^+ < 25$  and  $0.65 < z_{\text{phot}} < 1.75$  sources, the areal density is  $\sim 1500\ \text{deg}^{-2}$ . Using these numbers as a guide, the expected number of redshifts recovered in a single 2 h FMOS pointing ( $0.36 \times 200$  fibres, i.e. 72 sources) is comparable to the number expected from a single 4.5 h VIMOS pointing ( $0.7 \times 91$ , i.e. 64 targets). Thus for *Herschel*-selected targets FMOS is a competitive facility for redshift recovery in the range of  $0.7 < z < 1.8$ .

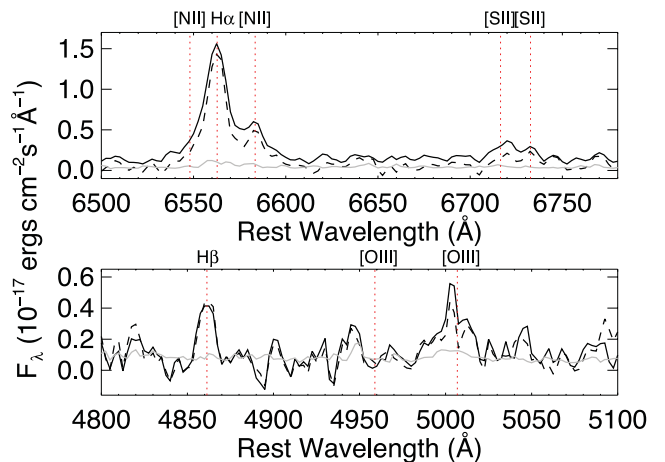
### 3.2 Composite spectrum for $H\alpha$ -detected sources

While our near-IR spectra are of sufficient quality to robustly measure fluxes for bright emission lines, very little additional information can be extracted from the individual spectra. This is partly due to the low S/N, but also because of the large number of OH sky lines in the near-IR, which significantly limit the wavelength coverage. However, using the spectral coverage unaffected by OH sky lines from each of our  $H\alpha$ -detected sources we can build a composite spectrum across a reasonably wide and continuous range of rest-frame wavelengths.

Fig. 3 shows the composite rest-frame spectrum in the region of  $H\alpha$  and  $H\beta$  produced from our sample of 57 sources. The composite spectrum is produced by first adopting a grid in rest-frame wavelength and calculating the contribution of each observed near-IR spectrum via a Gaussian kernel. The FWHM of the Gaussian kernel is taken to be the spectral resolution,  $\lambda/600\ \text{\AA}$ . Pixels within  $10\ \text{\AA}$  of an OH-suppressed line are excluded, as are those which lie outside the wavelength range  $1.1\text{--}1.36$  or  $1.42\text{--}1.7\ \mu\text{m}$ . The variance in the composite spectrum is estimated via jackknife resampling of the sources contributing to each wavelength bin. The absolute flux calibration of the composite is corrected by comparing the continuum level in the composite to the mean broad-band photometry in the  $J$  and  $K$  bands.

Several well-known spectral lines can be seen in addition to  $H\alpha$ :  $[\text{N II}] 6584$ , although this is blended with  $H\alpha$ ; a blend of the  $[\text{S II}]$  doublet at  $6716$  and  $6731\ \text{\AA}$ ; and weak signatures of  $[\text{O III}] 5008$  and  $H\beta$ .

The line ratio  $H\alpha/[\text{N II}] 6584$  compared to the line ratio  $[\text{O III}] 5008/H\beta$  is often used as diagnostic of AGN activity (Baldwin, Phillips & Terlevich 1981, hereafter BPT). Measuring the strengths of these lines from Fig. 3 we determine  $f_{[\text{N II}] 6584}/f_{H\alpha} = 0.24 \pm 0.07$ , while  $f_{[\text{O III}] 5008}/f_{H\beta} = 1.3 \pm 0.6$ . Comparing these ratios with the BPT diagnostic plots suggests that, on average, the primary origin of emission lines, and by proxy the IR luminosity, in our sample is star formation.



**Figure 3.** Composite rest-frame spectrum for the 57  $24 \cap 250\ \mu\text{m}$  sources with robust line detections in the region of  $H\alpha$  (top) and  $H\beta$  (bottom). The dashed line is the composite spectrum for only those 24 sources which are visible at the rest-frame wavelength of  $H\beta$ . The grey solid line represents the  $1\sigma$  variance in the composite. Well-known spectral features are marked. Encouragingly, our composite spectrum recovers several other spectral lines, such as  $[\text{N II}] 6584$ , the blended doublet of  $[\text{S II}] 6716$  and  $[\text{S II}] 6731$ , and weak signatures of  $[\text{O III}] 5008$  and  $H\beta$ .

The standard approach to estimating the dust attenuation in star-forming galaxies is to use the flux ratios of Balmer lines. Under typical conditions (i.e.  $T_e \sim 10^4$  K), and in the absence of dust, the line ratio  $f_{H\alpha}/f_{H\beta} = 2.86$  (Osterbrock 1989). Observed differences in this ratio, also known as the Balmer decrement, can be attributed to differential dust attenuation at the rest-frame wavelengths of the Balmer lines.

Using our composite spectrum we estimate the aggregate value of the Balmer decrement for SPIRE sources at  $z \gtrsim 1$ . The  $H\beta$  line is visible in the FMOS wavelength coverage at  $1.26 < z < 2.5$ , excluding sky lines. To account for this we build a second composite spectrum, using only those sources which have ‘clean’ (i.e. no overlapping sky lines) FMOS wavelength coverage at the wavelength of both  $H\beta$  and  $H\alpha$ ; only 24 sources satisfy this criterion. The composite spectrum from these sources in the region of  $H\alpha$  and  $H\beta$  is shown in Fig. 3. The  $H\beta$  observable sources have a mean redshift of  $\langle z \rangle = 1.36$  and a mean IR luminosity of  $\langle L_{IR} \rangle = 10^{45.5}$  erg s $^{-1}$  ( $10^{12} L_{\odot}$ ). We measure line fluxes of  $\langle f_{H\alpha} \rangle = 6.8 \pm 0.6 \times 10^{-16}$  erg cm $^{-2}$  s $^{-1}$  and  $\langle f_{H\beta} \rangle = 1.3 \pm 0.4 \times 10^{-16}$  erg cm $^{-2}$  s $^{-1}$ , after correcting both lines for stellar absorption ( $EW_{H\alpha} = 4.4$  and  $EW_{H\beta} = 2.8$  Å; Moustakas & Kennicutt 2006) and applying an aperture correction of 2.8 (the mean value for these 24 sources; see Section 3.3). This gives a Balmer decrement of  $\langle R \rangle = f_{H\alpha}/f_{H\beta} = 5.2 \pm 1.6$ , resulting in a dust attenuation of  $\langle E(B - V) \rangle = \log_{10}(R/2.86)/0.4[k(\lambda_{H\alpha}) - k(\lambda_{H\beta})] = 0.51 \pm 0.26$ , where  $k(\lambda_{H\alpha}) = 4.596$  and  $k(\lambda_{H\beta}) = 3.325$  (Calzetti et al. 2000). This is equivalent to  $A_v = 2.1$  mag, similar to that found for local IR luminous galaxies (e.g. Hopkins et al. 2001; Wijesinghe et al. 2011).

### 3.3 Relationship between $H\alpha$ and IR star formation rate estimates

In the absence of AGN activity or strongly non-solar metallicity, differences between the  $H\alpha$  estimated SFR ( $SFR_{H\alpha}$ ) and the best estimate of the total SFR ( $SFR_{tot}$ ) can be attributed to the effect of dust attenuation. Thus, the ratio  $SFR_{tot}/SFR_{H\alpha}$  can be used as an estimator of the level of dust attenuation (e.g. Hopkins et al. 2001; Kewley et al. 2002). In order to calculate  $SFR_{tot}$  and  $SFR_{H\alpha}$  the following steps were taken.

IR luminosities are calculated by fitting template spectral energy distributions (SEDs) to the 24  $\mu$ m and SPIRE data, then integrating the best-fitting template in the range of 8–1000  $\mu$ m. IR template SEDs are taken from Rieke et al. (2009). Both IR and  $H\alpha$  luminosities are converted into SFR via the relations presented in Kennicutt (1998), assuming a Chabrier (2003) initial mass function (IMF), i.e.

$$SFR_{IR} = 2.61 \times 10^{-44} L_{IR} \text{ (erg s}^{-1}\text{)}, \quad (3)$$

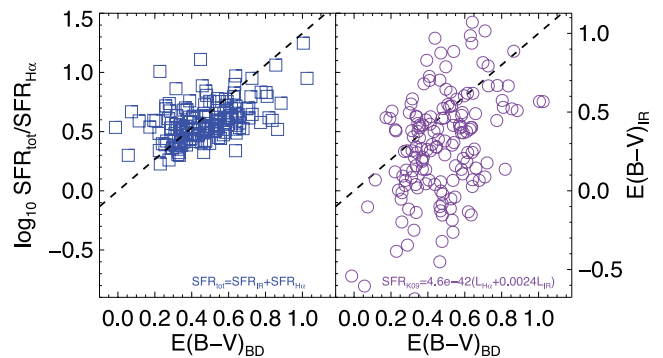
and

$$SFR_{H\alpha} = 4.61 \times 10^{-42} L_{H\alpha} \text{ (erg s}^{-1}\text{)}. \quad (4)$$

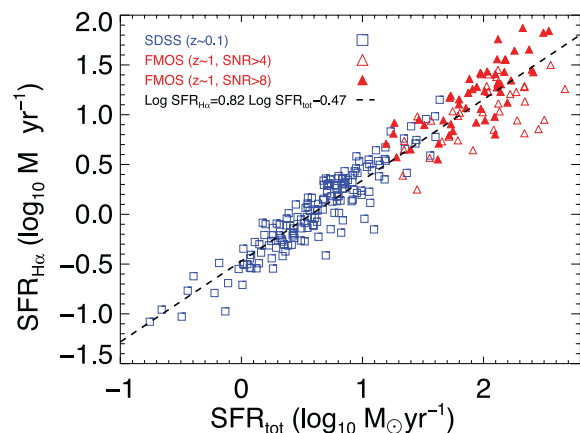
Finally,  $SFR_{tot}$  is calculated by adding together the IR and  $H\alpha$  SFRs, i.e.

$$SFR_{tot} = SFR_{IR} + SFR_{H\alpha}.$$

While other combined  $H\alpha$ +IR estimators of SFR exist (e.g. Kennicutt et al. 2009), these may not give good results for the class of IR luminous galaxies (nor *Herschel*-derived IR luminosities). We can confirm our assumption that  $SFR_{tot}/SFR_{H\alpha}$  is a good tracer of dust attenuation by comparing to estimates of the attenuation from the Balmer decrement for our SDSS sample. All 156 of our SDSS comparison sample have reliable ( $>3\sigma$ )  $H\beta$  line flux estimates. We produce estimates of the dust attenuation independent of the



**Figure 4.** Left-hand panel: comparison of dust attenuation estimates from the Balmer decrement  $[E(B - V)_{BD}]$  and from the ratio of  $SFR_{tot}/SFR_{H\alpha}$   $[E(B - V)_{IR}]$  for SDSS galaxies. The dashed line represents  $E(B - V)_{IR} = E(B - V)_{BD}$ . The two estimates are in excellent agreement, with an rms difference of 0.17. Right-hand panel: comparison of dust attenuation estimates from the Balmer decrement  $[E(B - V)_{BD}]$  and from the ratio of  $SFR_{K09}/SFR_{H\alpha}$ , where  $SFR_{K09}$  is the combined  $H\alpha$ +IR SFR estimator presented in Kennicutt et al. (2009, K09). Again, the dashed line represents  $E(B - V)_{IR} = E(B - V)_{BD}$ . It can be seen that the K09 SFR calibration gives much poorer agreement with the Balmer decrement estimated attenuations (rms = 0.27).



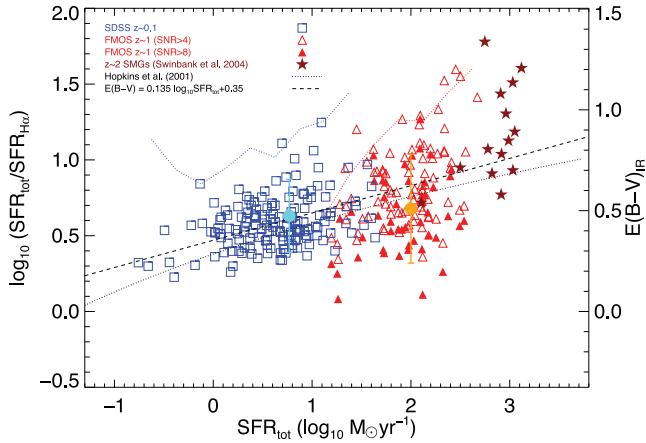
**Figure 5.** Comparison of  $H\alpha$  to total ( $H\alpha$  + far-IR) SFRs, assuming the conversions of Kennicutt (1998) with a Chabrier (2003) IMF. FMOS-detected sources (red triangles) and the low- $z$  sample of 160- $\mu$ m-selected SDSS galaxies (blue squares) are shown. The dashed line shows the best-fitting log-linear relation to the data.

far-IR measurements using the Balmer decrement. Fig. 4 compares the dust attenuation estimated from the Balmer decrement  $[E(B - V)_{BD}]$  to the inferred value assuming that  $SFR_{tot}/SFR_{H\alpha}$  is a good estimate of  $A_{H\alpha}$ .

Fig. 5 shows the relationship between  $SFR_{H\alpha}$  and  $SFR_{tot}$  for our sample of 57 sources detected at 24, 250  $\mu$ m and  $H\alpha$ . Fitting this observed correlation with a log-linear function results in the best fit:

$$\log_{10} SFR_{H\alpha} = (0.82 \pm 0.08) \log_{10} SFR_{IR} + 0.47 \pm 0.11 \text{ (} M_{\odot} \text{ yr}^{-1}\text{)}.$$

Using the Calzetti et al. (2000) model for the variation of dust attenuation with wavelength, we can use this result to convert the ratio  $SFR_{tot}/SFR_{H\alpha}$  to  $E(B - V)$ . Taking our best-fitting correlation between these values, we derive the following relationship between



**Figure 6.** Ratio of IR to dust uncorrected  $H\alpha$ -based SFR versus IR-based star formation rate (SFR). The equivalent  $E(B - V)$  is also given, assuming a Calzetti et al. (2000) attenuation curve. FMOS-detected sources (red triangles) and the low- $z$  sample of 160- $\mu\text{m}$ -selected SDSS galaxies (blue squares) are shown. The dashed line shows the local relation between SFR and  $E(B - V)$  from Hopkins et al. (2001). The  $E(B - V)$  determined from the Balmer decrement of our FMOS composite spectrum (Fig. 3), and mean value for SDSS spectra are shown as cyan and orange circles, respectively. Also shown are the positions of  $z \sim 2$  SMGs with  $H\alpha$  line measurements from Swinbank et al. (2004). Dotted lines show the effect of the  $H\alpha$  detection limit ( $5\sigma$ ) at the mean redshift for sources with that  $\text{SFR}_{\text{tot}}$  for the SDSS (blue) and FMOS (red) samples.

SFR and dust attenuation:

$$E(B - V)_{\text{IR}} = (0.135 \pm 0.06) \log_{10} \text{SFR}_{\text{tot}} + 0.35 \pm 0.08 (M_{\odot} \text{yr}^{-1}). \quad (5)$$

Fig. 6 presents dust-uncorrected  $\text{SFR}_{\text{tot}}/\text{SFR}_{H\alpha}$  versus  $\text{SFR}_{\text{tot}}$ . We also show in Fig. 6 our best-fitting relationship between  $E(B - V)$  and  $\text{SFR}_{\text{tot}}$ , and the empirical relationship for low- $z$  IRAS galaxies as derived by Hopkins et al. (2001). These two relations show reasonable agreement with the Hopkins et al. (2001) relation slightly below our simple log-linear fit.

Alternative estimates of the mean (and standard deviation) of dust attenuation measured from the Balmer decrement of individual SDSS spectra and the aggregate FMOS value from our composite spectrum (Section 3.2) are also shown in Fig. 6. Encouragingly the direct estimates of the mean dust attenuation are in good agreement with that inferred from  $\text{SFR}_{\text{IR}}/\text{SFR}_{H\alpha}$ , and both the dust attenuation and SFR relations.

Fig. 6 also shows *unlensed* 850- $\mu\text{m}$ -selected sources (SMGs) at  $z \gtrsim 2$ , with  $H\alpha$  line measurements from Swinbank et al. (2004). Here we make use of the  $\text{SFR}(H\alpha)$  and  $L(\text{FIR})$  quantities given in table 2 of Swinbank et al. (2004), converting  $\text{SFR}(H\alpha)$  to the Chabrier (2003) IMF used here and calculating  $\text{SFR}_{\text{IR}}$  from  $L(\text{FIR})$  using the equation given above.  $\text{SFR}(H\alpha)$  as quoted by Swinbank et al. (2004) includes corrections for slit loss and so should be compatible with the values we derive from FMOS and SDSS data. Interestingly the  $z \gtrsim 2$  SMGs appear slightly above both the Hopkins et al. (2001) and our best-fitting dust attenuation–SFR relation, suggesting that SMGs experience enhanced attenuation. However, it is worth noting that the  $L(\text{FIR})$  estimates for the SMGs come from pre-*Herschel* submillimetre radio estimates and hence may be overestimated (see Magnelli et al. 2012). Future studies with FMOS in this SFR range, as well as a re-assessment of the SMG population with *Herschel* photometry, will allow this trend to be confirmed.

The modest  $H\alpha$  detection limits achievable with FMOS mean we will not recover IR sources which have very large  $\text{SFR}_{\text{tot}}/\text{SFR}_{H\alpha}$ . To quantify this we calculate the typical maximum observable limit of  $\text{SFR}_{\text{tot}}/\text{SFR}_{H\alpha}$  as a function of  $\text{SFR}_{\text{tot}}$  for both the SDSS and FMOS samples. For the SDSS sample we assume a detection limit of  $f_{H\alpha} = 5.8 \times 10^{16} \text{ erg cm}^{-2} \text{ s}^{-1}$ , while for the FMOS sample we assume a detection limit of  $f_{H\alpha} = 1 \times 10^{16} \text{ erg cm}^{-2} \text{ s}^{-1}$ . In both cases these limits include a correction for the mean loss due to the limited aperture of the fibres (1.2 arcsec diameter for FMOS and 3 arcsec diameter for SDSS). While the maximum limit for the SDSS sample is significantly higher ( $\sim 0.2$  dex) than the observed values of  $\text{SFR}_{\text{tot}}/\text{SFR}_{H\alpha}$ , the limits for the FMOS data set appear quite close to the observed data points.

Given that  $\sim 65$  per cent of our parent sample is undetected in  $H\alpha$  a potential explanation for this large incompleteness is a significant population of sources with  $\text{SFR}_{\text{tot}}/\text{SFR}_{H\alpha}$  above these selection limits. Hence, the observed consistency with the Hopkins et al. (2001) and our equation (5) may be a result of a bias towards low  $\text{SFR}_{\text{tot}}/\text{SFR}_{H\alpha}$ .

While we cannot rule out the existence of large  $\text{SFR}_{\text{tot}}/\text{SFR}_{H\alpha}$  sources (as we cannot detect them), we can estimate the *observed* completeness for our parent sample assuming that our best fit to the  $E(B - V)$ – $\text{SFR}_{\text{tot}}$  relation is a good description for the whole population. For each source in our parent sample of  $168 \ 24 \cap 250 \ \mu\text{m}$  sources, we first estimate  $L_{\text{IR}}$ , assuming the photo- $z$  from Ilbert et al. (2009) and the SED-fitting process described above. The  $H\alpha$  line flux is then predicted from  $L_{\text{IR}}$  (assuming  $\text{SFR}_{\text{IR}} = \text{SFR}_{\text{tot}}$ ) using the best-fitting log-linear relation from Fig. 6 (with 0.5 dex of intrinsic scatter) and the mean loss due to the fibre aperture (2.8). Applying the completeness curves from Table 2 we would expect a completeness of 35 per cent, in good agreement with the observed completeness of  $34 \pm 4$  per cent, and consistent with a similar assessment of the completeness presented in Section 3.1.

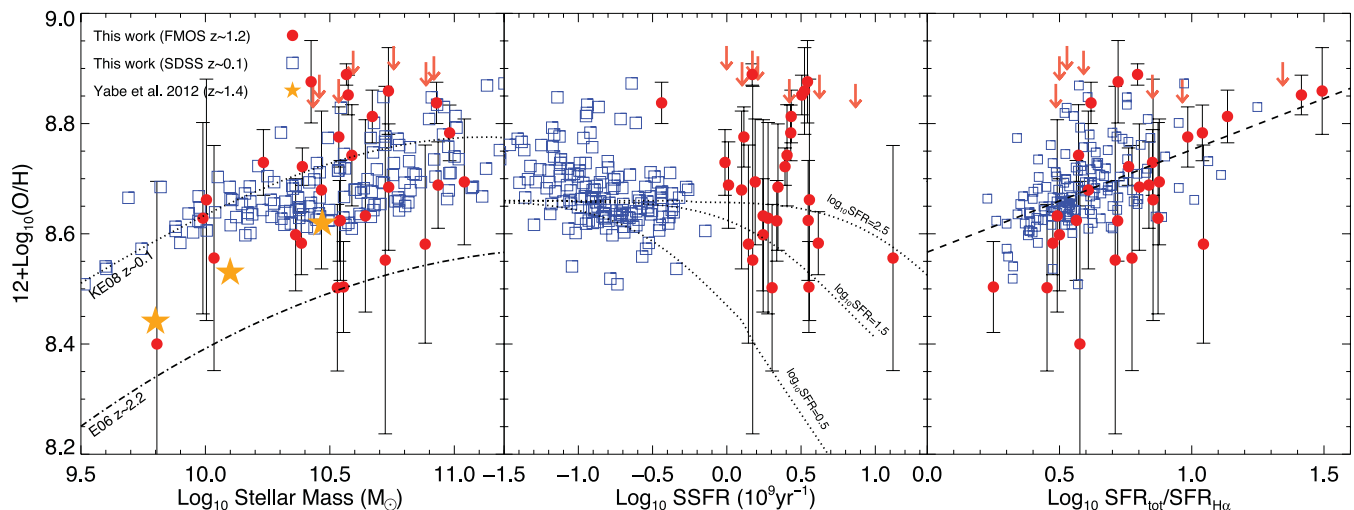
If a large fraction of our parent IR-selected sample had dust attenuation levels (i.e.  $\text{SFR}_{\text{tot}}/\text{SFR}_{H\alpha}$ ) significantly above that predicted by equation (5), we would expect much lower completeness in our FMOS observations than achieved. Given the good agreement between the observed and expected completeness, we conclude that equation (5) must hold for the bulk of IR luminous sources at  $z \sim 1$ .

### 3.4 The mass–metallicity relation for IR galaxies

A significant fraction (28/57) of our sample have robust ( $S/N > 3$ ) measurements of the  $[\text{N II}]$  6584 line. This allows us to investigate the gas-phase metallicity ( $12 + \log_{10} \text{O/H}$ ) of our IR sources via the ratio of the  $[\text{N II}]$  6584 to the  $H\alpha$  line (N2 method; Kewley & Dopita 2002; Pettini & Pagel 2004). For each source with a robust  $[\text{N II}]$  measurement we estimate the metallicity via equation (1) of Pettini & Pagel (2004, hereafter PP04):

$$12 + \log_{10} \text{O/H} = 8.9 + 0.57 \times \log_{10} f_{[\text{N II}]} / f_{H\alpha}.$$

The existence of a relationship between stellar mass and metallicity has been confirmed across a wide range in redshift ( $0 < z < 3$ ; Lequeux et al. 1979; Tremonti et al. 2004; Erb et al. 2006; Maiolino et al. 2008; Zahid et al. 2011), with a steady trend towards lower metallicity for a given stellar mass with increasing redshift. To investigate the mass–metallicity relation for our  $z \sim 1$  IR sources we combine our metallicity estimates with stellar masses as derived by Wang et al. (2012). Stellar mass estimates from Wang et al. (2012) are calculated by finding the best-fitting stellar population model to the observed multiwavelength photometry using the *LE PHARE*



**Figure 7.** Relationship between stellar mass (left-hand panel), specific SSFR (middle panel) and dust attenuation (right-hand panel) versus gas-phase metallicity as seen in our IR-selected sample. All sources with  $[N\ II]$  line flux measurements with  $S/N > 3$  are shown, while limits are shown for those sources without robust  $[N\ II]$  line measures. Sources from our *Herschel*-FMOS sample at  $z \sim 1.2$  are shown as red dots, while the  $\sim 0.1$  *Herschel*-SDSS samples are blue squares. Previous estimates of the stellar mass–metallicity relation at  $z \sim 0.1$  (Kewley & Ellison (2008, KE08) and  $z \sim 2.2$  (Erb et al. 2006, E06) are shown as dotted and dot–dashed lines in the left-hand panel, respectively. Results from a previous FMOS study for star-forming galaxies at  $z \sim 1.4$  (Yabe et al. 2012) are shown as orange stars. Where necessary, comparison samples have been converted to use the PP04 calibration of the  $N2$  metallicity tracer (KE08) and a Chabrier (2003) IMF. In the middle panel the position of the ‘fundamental metallicity relation’ of Mannucci et al. (2010) is shown for a range of SFRs representative of our sample as the dotted lines. In the right-hand panel the best-fitting log-linear relation between metallicity and dust attenuation (as traced by the ratio  $SFR_{tot}/SFR_{H\alpha}$ ) is shown as a dashed line.

software (Arnouts et al. 2002; Ilbert et al. 2006), combined with stellar population synthesis models from Bruzual & Charlot (2003) and assuming a Chabrier (2003) IMF.

The left-hand panel of Fig. 7 compares the stellar mass to the metallicity for both our sample of  $z \sim 1$  *Herschel*-FMOS and comparison sample of  $z \sim 0.1$  SDSS sources. A tentative trend of metallicity with stellar mass can be seen in both samples, although with significant scatter. No discernable evolution is seen in the metallicity between the  $z \sim 1$  and 0.1 IR-selected samples.

Also shown in the left-hand panel of Fig. 7 are studies using stacking of near-IR spectra by Erb et al. (2006) with Keck NIRSPEC at  $z \sim 2.2$  and Yabe et al. (2012) with FMOS at  $z \sim 1.4$ . (We convert the Yabe et al. 2012 stellar masses to a Chabrier 2003 IMF for consistency.) Given the similar redshift range between our *Herschel*-FMOS sample and the Yabe et al. (2012) study, the difference in the metallicities is surprising, although likely due to a combination of the different mass ranges probed, the stacking nature of the Yabe et al. (2012) result and our bias towards dusty galaxies.

In the middle panel of Fig. 7, we show the relationship between specific SFR (SSFR;  $SFR/M_*$ ) and metallicity. Over the range in SSFR well sampled by our *Herschel*-FMOS and SDSS samples we see no discernable trends, although there are hints of decreasing metallicity with increasing SSFR above  $SSFR = 0.5\text{ Gyr}^{-1}$ . It has been proposed that the metallicity is a function of both SFR and stellar mass, with a ‘fundamental mass relation’ (FMR) linking the three parameters (Mannucci et al. 2010). In Fig. 7 we show the fit to the FMR from Mannucci et al. (2010) for values of SFR representative of our study. The metallicities of both our *Herschel*-FMOS and SDSS are consistent with the prediction from the FMR.

Finally, we consider the relationship between dust attenuation and metallicity. In the right-hand panel of Fig. 7 is shown the metallicity as a function of dust attenuation (as traced by the ratio  $SFR_{tot}/SFR_{H\alpha}$ ). A clear trend between dust attenuation and metallicity is seen in both the *Herschel*-FMOS and SDSS samples, with

the most metal-rich galaxies experiencing the largest levels of attenuation. The best-fitting log-linear relation between attenuation and metallicity is found to be

$$12 + \log_{10} O/H = 0.19 \log_{10} SFR_{tot}/SFR_{H\alpha} + 8.57.$$

Converting to  $E(B - V)$  via the Calzetti et al. (2000) attenuation curve we find

$$12 + \log_{10} O/H = 0.24 E(B - V) + 8.57. \quad (6)$$

This correlation is somewhat expected; the dust grains responsible for attenuating starlight (and nebular line emission) are synthesized from metals in the interstellar medium (ISM). Similar studies, using the IR to UV luminosity ratios as a proxy for dust attenuation, have also found a correlation between metallicity and attenuation for nearby galaxies (Heckman et al. 1998) and  $z \sim 2$  (Reddy et al. 2010). Given the correlation between metallicity with attenuation the high metallicity of our *Herschel*-FMOS sample (when compared to Yabe et al. 2012, at  $z \sim 1.4$ ) is to be expected as our IR-selected sample must be biased towards the most obscured sources.

#### 4 DISCUSSION AND CONCLUSIONS

We have investigated the properties of  $z \sim 1$  IR luminous galaxies by performing near-IR spectroscopy of a sample of *Spitzer* and *Herschel*-selected sources with FMOS. Candidate emission lines were identified in the 2D-reduced FMOS frames via a semi-automated procedure. Via comparison with known spectroscopic redshifts, and direct testing of the line detection algorithm, we estimate that our  $H\alpha$  line sample is  $\gtrsim 90$  per cent reliable. Our scientific conclusions are as follows.

- (i) Robust detections of  $H\alpha$  were found for 57 of 168  $24 \times 250\ \mu\text{m}$  sources, resulting in a detection rate of  $34 \pm 4$  per cent. This detection rate is consistent with the expected incompleteness measured by simulating the line detection process. For sparse targets such as



*Herschel*-selected sources, FMOS is a competitive redshift recovery instrument with equivalent optical multi-object spectrograph instruments on other 8-m class telescope (e.g. VIMOS on VLT).

(ii) The mean dust attenuation, estimated via the Balmer decrement for the  $H\alpha$  and  $H\beta$  emission lines for a composite spectrum of  $H\alpha$  detected sources, is  $E(B - V) = 0.51 \pm 0.27$  for  $L_{\text{IR}} = 10^{12} L_{\odot}$  sources at  $z = 1.36$ .

(iii) Good agreement was found between the dust attenuation estimated from the Balmer decrement and that inferred from the ratio of  $H\alpha$  estimated to best estimate of the total SFR. Using  $\text{SFR}_{\text{tot}}/\text{SFR}_{H\alpha}$  as an indicator of attenuation in the  $H\alpha$  line we derive a relationship between dust attenuation and SFR of  $E(B - V) = (0.135 \pm 0.06)\log_{10}\text{SFR}_{\text{tot}} + 0.35 \pm 0.08$ . These results are broadly consistent with the relationship between SFR and dust attenuation seen both in low- $z$  star-forming galaxies (e.g. Hopkins et al. 2001).

(iv) The gas-phase metallicity relation was investigated for the subset of *Herschel*-FMOS, and SDSS, sources with robust measurements of the  $[\text{N II}] 6584$  line. For IR-selected sources with  $M_{\star} \sim 10^{10.5} M_{\odot}$  the typical metallicity is found not to evolve between  $z \sim 0.1$  and 1.2. No discernable trend with SSFR is seen, but a strong correlation between metallicity and dust attenuation is seen, described by a best-fitting log-linear relation:  $12 + \log_{10} \text{O/H} = 0.24 E(B - V) + 8.57$ .

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The data presented in this paper will be released through the *Herschel* Database in Marseille HeDaM ([hedam.oamp.fr/HerMES](http://hedam.oamp.fr/HerMES)).

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## APPENDIX A: DETAILS OF THE HERMES-FMOS SAMPLE

**Table A1.** Details of *Herschel* 250  $\mu\text{m}$  and *Spitzer* 24  $\mu\text{m}$  selected sources with reliable  $\text{H}\alpha$  line detections.

RA <sup>a</sup> (°)	Dec. (°)	Redshift <sup>b</sup>	$I_{F814W}$ AB mag	$S_{24\mu\text{m}}$ (mJy)	$S_{250\mu\text{m}}$ <sup>c</sup> (mJy)	$S_{350\mu\text{m}}$ <sup>c</sup> (mJy)	$S_{500\mu\text{m}}$ <sup>c</sup> (mJy)	$f_{\text{H}\alpha}$ <sup>d</sup> ( $\times 10^{16} \text{ erg cm}^{-2} \text{ s}^{-1}$ )	$f_{[\text{N III}]}$ <sup>e</sup> ( $\times 10^{16} \text{ erg cm}^{-2} \text{ s}^{-1}$ )
149.94205	2.32077	1.03	21.94 ± 0.01	0.20 ± 0.01	9.80 ± 2.73	<12.22	<15.98	4.32 ± 0.30	<3.95
149.96689	2.44185	0.90	22.44 ± 0.02	0.08 ± 0.01	<8.20	<13.15	<9.28	2.84 ± 0.32	<5.78
149.97292	2.49010	1.36	23.64 ± 0.03	0.43 ± 0.01	15.44 ± 2.73	19.04 ± 3.22	<15.41	1.52 ± 0.41	<6.17
149.95179	2.48627	1.03	22.00 ± 0.01	0.28 ± 0.01	10.13 ± 2.74	11.98 ± 3.38	<14.81	2.33 ± 0.17	<2.62
149.99931	2.45197	1.51	23.13 ± 0.03	0.33 ± 0.01	21.61 ± 2.73	15.34 ± 3.61	18.14 ± 4.87	2.96 ± 0.40	<3.05
150.03484	2.26353	0.90	21.46 ± 0.01	0.25 ± 0.01	11.26 ± 2.73	<26.43	<33.24	5.44 ± 0.37	<6.56
150.04760	2.62123	1.38	22.73 ± 0.02	0.23 ± 0.01	20.49 ± 2.73	13.12 ± 3.25	<17.00	3.20 ± 0.27	<2.19
150.02476	2.35211	0.93	21.76 ± 0.01	0.15 ± 0.01	<8.20	<13.09	<35.23	4.55 ± 0.38	<5.37
150.03650	2.31730	1.46	22.36 ± 0.02	0.44 ± 0.01	26.35 ± 2.73	25.45 ± 3.48	<17.85	11.43 ± 0.33	<4.66
150.02440	2.44388	1.03	21.72 ± 0.02	0.39 ± 0.01	<41.18	<32.57	<25.69	10.69 ± 0.56	<7.19
150.04820	2.46775	1.17	23.33 ± 0.03	0.34 ± 0.01	11.71 ± 2.73	<9.31	<14.30	7.69 ± 0.34	2.57 ± 0.78
150.03637	2.49426	1.03	21.86 ± 0.01	0.41 ± 0.01	9.23 ± 2.74	<9.59	<15.54	6.86 ± 0.28	<3.75
150.08336	2.53619	1.42	23.70 ± 0.04	0.30 ± 0.01	22.99 ± 2.73	25.93 ± 4.79	33.71 ± 9.86	7.97 ± 0.28	<3.39
150.05928	2.51858	1.03	22.10 ± 0.01	0.38 ± 0.01	12.46 ± 2.74	<9.61	<16.43	1.92 ± 0.13	<2.17
150.05386	2.58972	0.70	18.89 ± 0.00	3.45 ± 0.01	<8.19	<10.02	<14.47	68.67 ± 0.31	47.99 ± 1.39
150.10975	2.60274	0.98	21.80 ± 0.01	0.32 ± 0.01	<8.21	<33.17	<9.27	10.12 ± 0.19	–
150.10019	2.48157	0.89	22.11 ± 0.02	0.08 ± 0.01	<8.23	<9.27	<14.32	4.88 ± 0.48	<5.28
150.13340	2.26201	0.75	21.63 ± 0.01	0.08 ± 0.01	<8.19	<9.94	<14.36	3.31 ± 0.32	<4.71
150.14563	2.29341	0.88	22.27 ± 0.02	0.10 ± 0.01	<8.20	<9.29	<4.77	3.77 ± 0.37	<4.26
150.13445	2.61448	0.89	22.44 ± 0.01	0.15 ± 0.01	<8.22	<16.66	<16.33	2.93 ± 0.16	<1.78
150.16162	2.69151	1.48	23.51 ± 0.03	0.52 ± 0.01	21.25 ± 2.73	24.13 ± 3.95	17.66 ± 5.36	5.05 ± 0.61	–
150.15563	2.67708	1.04	22.28 ± 0.02	0.21 ± 0.01	10.36 ± 2.74	<29.40	<32.03	3.14 ± 0.15	<2.70
150.12472	2.66871	1.28	22.56 ± 0.01	0.19 ± 0.01	21.05 ± 2.73	<12.30	<14.92	5.82 ± 0.12	1.43 ± 0.45
150.15229	2.21933	0.92	21.86 ± 0.01	0.36 ± 0.01	<8.21	<12.19	<7.61	3.60 ± 0.19	<2.96
150.21805	2.52182	1.18	21.72 ± 0.01	0.30 ± 0.01	13.05 ± 2.74	22.89 ± 3.56	<16.66	5.98 ± 0.34	<2.59
150.21025	2.56547	1.40	22.29 ± 0.02	0.34 ± 0.01	30.92 ± 2.73	22.22 ± 3.69	<36.15	5.34 ± 0.15	2.24 ± 0.66
150.25980	2.29235	0.99	22.42 ± 0.01	0.14 ± 0.01	<8.20	<16.07	<18.66	3.48 ± 0.27	–
150.22858	2.31620	0.90	21.25 ± 0.01	0.30 ± 0.01	<8.22	<9.56	<17.02	8.68 ± 0.17	<2.55
150.22203	2.62003	0.69	21.50 ± 0.01	0.65 ± 0.01	27.91 ± 2.73	22.13 ± 7.22	<21.83	9.14 ± 0.48	<9.70
150.24813	2.39912	0.68	21.03 ± 0.01	0.15 ± 0.01	<2.73	<3.35	<17.64	9.01 ± 0.50	<6.14
150.29832	2.46967	0.85	21.75 ± 0.01	0.18 ± 0.01	<8.19	<10.17	<25.53	2.91 ± 0.20	<3.79
150.29408	2.51691	0.84	21.89 ± 0.01	0.16 ± 0.01	<8.21	<18.92	<12.42	3.02 ± 0.30	<0.84
150.25672	2.48480	1.24	23.13 ± 0.03	0.19 ± 0.01	17.77 ± 2.74	12.31 ± 3.66	<30.69	4.84 ± 0.64	4.32 ± 0.66
150.29414	2.57633	0.78	22.40 ± 0.02	0.21 ± 0.01	<8.21	<11.11	<19.23	1.98 ± 0.20	<5.92
150.31964	2.61813	1.00	21.59 ± 0.01	0.24 ± 0.01	<30.62	<31.33	<24.63	5.03 ± 0.36	–
150.27952	2.59666	1.49	23.03 ± 0.02	0.17 ± 0.01	13.49 ± 2.73	13.22 ± 3.15	<15.18	3.81 ± 0.47	<2.73
150.34435	2.73311	0.85	21.56 ± 0.01	0.26 ± 0.02	<8.19	15.09 ± 3.95	<48.57	2.33 ± 0.12	1.10 ± 0.14
150.32742	2.68368	0.96	21.62 ± 0.01	0.45 ± 0.01	<8.21	<18.05	<15.67	9.08 ± 0.22	3.84 ± 0.28
150.34866	2.45537	1.02	22.12 ± 0.02	0.22 ± 0.01	11.54 ± 2.74	<9.68	<21.07	5.77 ± 0.25	1.67 ± 0.27
150.31731	2.50830	0.98	22.31 ± 0.01	0.09 ± 0.01	14.61 ± 2.74	12.64 ± 3.38	<33.01	2.73 ± 0.29	1.01 ± 0.30

Table A1 – *continued*

RA <sup>a</sup> (°)	Dec. (°)	Redshift <sup>b</sup>	$I_{F814W}$ AB mag	$S_{24\ \mu\text{m}}$ (mJy)	$S_{250\ \mu\text{m}}^c$ (mJy)	$S_{350\ \mu\text{m}}^c$ (mJy)	$S_{500\ \mu\text{m}}^c$ (mJy)	$f_{\text{H}\alpha}^d$ ( $\times 10^{16}\ \text{erg cm}^{-2}\ \text{s}^{-1}$ )	$f_{[\text{N II}]}^e$ ( $\times 10^{16}\ \text{erg cm}^{-2}\ \text{s}^{-1}$ )
150.36745	2.55928	0.91	22.20 ± 0.02	0.13 ± 0.01	<8.19	<8.07	<16.49	2.94 ± 0.26	<4.69
150.35415	2.51916	0.96	21.32 ± 0.01	0.75 ± 0.01	29.87 ± 2.74	<23.46	<31.47	6.18 ± 0.50	1.63 ± 0.30
150.35350	2.55854	1.04	21.63 ± 0.01	0.17 ± 0.01	10.25 ± 2.73	<13.46	<14.98	4.17 ± 0.24	3.08 ± 0.30
150.34722	2.57026	0.92	21.48 ± 0.01	0.30 ± 0.01	<8.19	<20.10	<22.18	3.01 ± 0.19	<3.12
150.36110	2.60817	1.30	22.89 ± 0.02	0.11 ± 0.01	<8.19	<20.03	<24.99	3.96 ± 0.21	<0.58
150.33580	2.64975	1.33	22.58 ± 0.02	0.12 ± 0.01	20.33 ± 2.73	23.03 ± 3.78	<14.14	3.66 ± 0.21	1.49 ± 0.20
150.34937	2.36957	1.57	23.49 ± 0.04	0.13 ± 0.02	<8.23	<18.35	<15.06	7.85 ± 0.70	2.70 ± 0.64
150.39351	2.44529	0.92	21.67 ± 0.01	0.28 ± 0.01	<8.19	<11.84	<15.06	7.62 ± 0.25	1.18 ± 0.26
150.37070	2.49822	0.82	20.52 ± 0.01	0.43 ± 0.01	19.88 ± 2.73	11.93 ± 3.54	<17.47	8.12 ± 0.28	7.45 ± 0.66
150.41783	2.55859	1.21	22.13 ± 0.01	0.76 ± 0.01	47.01 ± 2.74	34.07 ± 3.15	19.10 ± 4.88	2.47 ± 0.32	1.97 ± 0.30
150.42288	2.58332	0.82	20.43 ± 0.01	0.69 ± 0.01	29.20 ± 2.74	<28.48	<15.74	30.07 ± 0.65	3.95 ± 1.03
150.42188	2.32331	0.83	23.16 ± 0.02	0.61 ± 0.01	20.59 ± 2.74	<9.86	<18.94	16.98 ± 1.98	–
150.40063	2.33453	1.21	22.24 ± 0.01	0.36 ± 0.01	26.17 ± 2.74	21.99 ± 3.20	<15.97	5.93 ± 0.47	2.43 ± 0.43
150.42068	2.62304	1.29	23.56 ± 0.04	0.22 ± 0.01	14.60 ± 2.73	<15.63	<16.45	4.58 ± 0.21	–
150.44572	2.76095	1.35	23.11 ± 0.03	0.18 ± 0.01	12.08 ± 2.74	<10.95	<21.60	2.02 ± 0.18	<0.74
150.49513	2.37789	0.89	21.62 ± 0.01	0.62 ± 0.01	52.22 ± 2.73	40.12 ± 3.56	22.55 ± 6.23	3.88 ± 0.25	<1.12
150.48572	2.71969	0.89	21.95 ± 0.01	0.58 ± 0.01	27.84 ± 2.72	17.44 ± 3.12	<15.71	7.59 ± 0.33	<1.65
150.51527	2.56268	1.27	23.40 ± 0.03	0.20 ± 0.01	12.75 ± 2.74	<51.16	<46.81	6.56 ± 0.64	3.43 ± 0.42
150.51142	2.57740	0.79	21.35 ± 0.01	0.14 ± 0.01	10.93 ± 2.74	13.80 ± 3.40	<16.00	2.88 ± 0.24	0.91 ± 0.21
150.50345	2.65037	1.29	23.02 ± 0.03	0.22 ± 0.02	13.67 ± 2.73	13.00 ± 3.09	15.19 ± 4.77	1.52 ± 0.11	–
150.55036	2.73248	0.85	20.93 ± 0.01	0.52 ± 0.01	14.30 ± 2.74	<10.32	<18.21	10.68 ± 0.49	3.41 ± 0.32
150.56193	2.47613	0.82	20.86 ± 0.01	0.48 ± 0.01	11.91 ± 2.73	<19.96	<16.95	9.00 ± 0.28	4.25 ± 0.34
150.54374	2.49755	0.88	21.73 ± 0.01	0.15 ± 0.02	<8.21	<28.80	<33.36	5.60 ± 0.48	3.26 ± 0.61
150.59733	2.61799	1.41	21.75 ± 0.00	0.28 ± 0.01	20.36 ± 2.73	<45.00	18.81 ± 5.39	1.72 ± 0.15	–
150.56845	2.64125	1.26	23.55 ± 0.03	0.19 ± 0.01	16.23 ± 2.75	<10.33	<11.99	8.51 ± 0.46	–
150.59931	2.38340	0.89	21.78 ± 0.01	0.37 ± 0.01	9.10 ± 2.73	<12.41	<15.17	3.58 ± 0.22	2.07 ± 0.25
150.59031	2.53405	1.50	22.60 ± 0.01	0.18 ± 0.01	8.46 ± 2.74	10.75 ± 3.54	<18.56	3.20 ± 0.25	–
150.62185	2.55274	0.88	21.45 ± 0.01	0.46 ± 0.01	18.95 ± 2.74	15.05 ± 3.73	16.41 ± 5.12	6.22 ± 0.66	2.49 ± 0.43
150.63049	2.55622	1.29	22.29 ± 0.02	0.12 ± 0.01	9.19 ± 2.74	<45.35	<52.71	5.96 ± 0.25	–
150.63901	2.61703	1.33	22.80 ± 0.02	0.22 ± 0.01	12.32 ± 2.72	<9.65	<15.44	15.27 ± 0.41	3.06 ± 0.32
150.62021	2.62575	1.32	21.14 ± 0.01	0.34 ± 0.01	21.49 ± 2.74	<11.06	<18.35	3.00 ± 0.20	1.99 ± 0.19
150.62429	2.72533	1.21	22.73 ± 0.02	0.17 ± 0.01	8.48 ± 2.73	<10.81	<11.59	6.74 ± 0.26	1.34 ± 0.24
150.66639	2.44691	0.87	21.51 ± 0.01	0.32 ± 0.01	<8.23	<16.02	<16.12	5.26 ± 0.33	–
150.63611	2.74441	1.38	23.63 ± 0.04	0.35 ± 0.01	39.39 ± 2.73	<34.02	<29.14	5.12 ± 0.31	3.12 ± 0.33
150.66376	2.49340	1.34	22.26 ± 0.02	0.27 ± 0.01	12.90 ± 2.73	<11.30	<15.90	3.48 ± 0.24	–
150.67631	2.52723	1.38	24.10 ± 0.05	0.69 ± 0.01	22.50 ± 2.73	20.77 ± 3.66	<21.77	3.43 ± 0.41	<0.76
150.68868	2.61353	0.96	22.15 ± 0.02	0.36 ± 0.01	18.86 ± 2.73	<11.03	<12.07	3.97 ± 0.21	–
150.66590	2.64705	1.40	22.45 ± 0.02	0.18 ± 0.01	13.48 ± 2.74	<9.29	<4.78	6.61 ± 0.16	1.81 ± 0.20
150.70877	2.52972	0.98	21.54 ± 0.01	0.49 ± 0.01	23.98 ± 2.73	11.34 ± 3.28	<16.55	3.34 ± 0.19	2.58 ± 0.26
150.70577	2.56658	0.90	22.51 ± 0.02	0.07 ± 0.01	<20.36	<11.73	<17.71	1.02 ± 0.11	<0.52
150.72242	2.61250	1.20	22.76 ± 0.03	0.23 ± 0.01	15.70 ± 2.74	<12.59	<17.81	4.58 ± 0.46	1.10 ± 0.31
150.31065	2.31922	1.43	25.53 ± 0.13	0.17 ± 0.01	<8.21	<18.13	<19.54	3.28 ± 0.26	<2.85
150.37685	2.45953	1.30	25.53 ± 0.14	0.36 ± 0.01	19.29 ± 2.74	22.76 ± 4.10	<19.70	4.51 ± 0.38	<2.60
150.46438	2.63802	1.46	24.59 ± 0.10	0.09 ± 0.01	17.27 ± 2.74	25.30 ± 3.15	26.53 ± 4.91	1.80 ± 0.32	<1.50
150.68825	2.45364	1.38	25.95 ± 0.15	0.18 ± 0.01	19.43 ± 2.73	16.53 ± 3.66	<17.36	6.78 ± 0.45	2.20 ± 0.27

<sup>a</sup>Position from  $I_{F814W}$  catalogue.<sup>b</sup>Redshift from detected H $\alpha$  line.<sup>c</sup>Limits are  $3\sigma$ , including confusion noise.<sup>d</sup>All values corrected for the limited aperture (1.2 arcsec diameter) of the FMOS fibre.<sup>e</sup>All values corrected for the limited aperture (1.2 arcsec diameter) of the FMOS fibre. Limits are  $3\sigma$ , no value given in cases, where the [N II] 6584 line falls on an OH sky line.This paper has been typeset from a  $\text{\TeX}/\text{\LaTeX}$  file prepared by the author.