

H<sub>2</sub>O AND OH GAS IN THE TERRESTRIAL PLANET–FORMING ZONES OF PROTOPLANETARY DISKSCOLETTE SALYK,<sup>1</sup> KLAUS M. PONTOPPIDAN,<sup>1</sup> GEOFFREY A. BLAKE,<sup>1</sup> FRED LAHUIS,<sup>2,3</sup>EWINE F. VAN DISHOCK,<sup>3</sup> AND NEAL J. EVANS II<sup>4</sup>*Received 2007 November 15; accepted 2008 January 31; published 2008 March 6*

## ABSTRACT

We present detections of numerous 10–20  $\mu\text{m}$  H<sub>2</sub>O emission lines from two protoplanetary disks around the T Tauri stars AS 205A and DR Tau, obtained using the InfraRed Spectrograph on the *Spitzer Space Telescope*. Follow-up 3–5  $\mu\text{m}$  Keck NIRSPEC data confirm the presence of abundant water and spectrally resolve the lines. We also detect the P4.5 (2.934  $\mu\text{m}$ ) and P9.5 (3.179  $\mu\text{m}$ ) doublets of OH and <sup>12</sup>CO/<sup>13</sup>CO  $v = 1 \rightarrow 0$  emission in both sources. Line shapes and LTE models suggest that the emission from all three molecules originates between  $\sim 0.5$  and 5 AU, and so will provide a new window for understanding the chemical environment during terrestrial planet formation. LTE models also imply significant columns of H<sub>2</sub>O and OH in the inner disk atmospheres, suggesting physical transport of volatile ices either vertically or radially, while the significant radial extent of the emission stresses the importance of a more complete understanding of nonthermal excitation processes.

*Subject headings:* astrochemistry — circumstellar matter — ISM: molecules — planetary systems: protoplanetary disks

## 1. INTRODUCTION

One of the most intriguing questions in the study of the formation of planets, and of terrestrial planets in particular, is how water is transported to their surfaces, and whether or not water is a common ingredient during their formation and early evolution. Spatially and spectrally resolved observations of water in extrasolar planetary systems and protoplanetary disks would be instrumental in resolving these questions. Unfortunately, detecting water in extrasolar planetary systems remains difficult (Ehrenreich et al. 2007) and water in protoplanetary disks has been elusive, with only a few definitive ice (Malfait et al. 1998; Terada et al. 2007) and vapor (Carr et al. 2004) detections. Recently, however, the *Spitzer* InfraRed Spectrograph (IRS) has begun to reveal warm water vapor emission lines from disks around classical T Tauri stars (CTTSs; Carr & Najita 2008).

In this Letter, we present high-resolution spectroscopy of numerous water vapor emission lines from protoplanetary disks around two CTTSs: DR Tau and AS 205A. Spectra were obtained using both the *Spitzer* IRS and NIRSPEC (McLean et al. 1998) at the Keck II telescope. Accretion rates onto these stars are high, which causes the inner disks to be significantly hotter than in more quiescent systems. DR Tau is known to have variable mass accretion, with measured rates between 0.3 and  $79 \times 10^{-7} M_{\odot} \text{ yr}^{-1}$  (Gullbring et al. 2000; Johns-Krull & Gafford 2002). AS 205A has an accretion rate of  $7.2 \times 10^{-7} M_{\odot} \text{ yr}^{-1}$  and is the primary component of a 1.3'' triple system, with the secondary being a spectroscopic binary (Eisner et al. 2005). However, it is likely that the physical separation of AS 205A and B is significantly larger than the projected separation and that, therefore, AS 205B does not affect the (inner) disk of AS 205A (Andrews & Williams 2007).

## 2. OBSERVATIONS

The *Spitzer* IRS Short-High (SH) module spectrum of AS 205A was taken as part of the “cores to disks” (c2d) Legacy program (Evans et al. 2003; Kessler-Silacci et al. 2006), while the spectrum of DR Tau, observed as part of the guaranteed time observations (Houck et al. 2004), was obtained from the *Spitzer* archive. The 2D basic calibrated data images were reduced using the c2d pipeline as described in Lahuis et al. (2006).

To confirm the detections as well as measure the gas emission-line profiles, we obtained high-resolution ( $\lambda/\Delta\lambda = 25,000$ ) L-band spectra using NIRSPEC in a single grating setting, of which 2 orders, centered at 2.92 and 3.16  $\mu\text{m}$ , respectively, have identifiable H<sub>2</sub>O and OH emission features. We also collected CO  $v = 1 \rightarrow 0$  M-band spectra ( $\sim 5 \mu\text{m}$ ) as part of an ongoing survey (see Blake & Boogert 2004; Salyk et al. 2007). The data were reduced in a standard way. Division by the spectra of standard stars (HR 1620 [A7 V] for DR Tau; HR 5993 [B1 V] for AS 205A) was used to correct for telluric absorption, which can be poor when the absorption is high. Thus, regions of the spectrum with  $<70\%$  atmospheric transmission were removed. Fluxes were calibrated using the photometry and spectral types of the standards, and are generally accurate to  $\leq 20\%$ .

The NIRSPEC lines are narrow ( $\text{FWHM} \lesssim 35 \text{ km s}^{-1}$ ), spatially unresolved, and centered at the stellar  $v_{\text{LSR}}$  (Herbig & Bell 1988; Thi et al. 2001) to within our uncertainty (a few  $\text{km s}^{-1}$ )—all consistent with disk emission. Because the disks are optically thick and heated from above, the line emission likely arises from the upper disk atmosphere (Najita et al. 2003).

## 3. RESULTS

3.1. *Spitzer* IRS

The IRS spectra of AS 205A and DR Tau show a large number of water emission lines (see Fig. 1), the vast majority pure rotational transitions with quantum numbers in the range 20–50. The  $R = 600$  resolution of the SH module results in unresolved, and significantly blended, lines. Consequently, the identification with water is achieved via a simple, isothermal LTE ring model computed using the HITRAN 2004 database (Rothman et al. 2005). The ring model consists of a single

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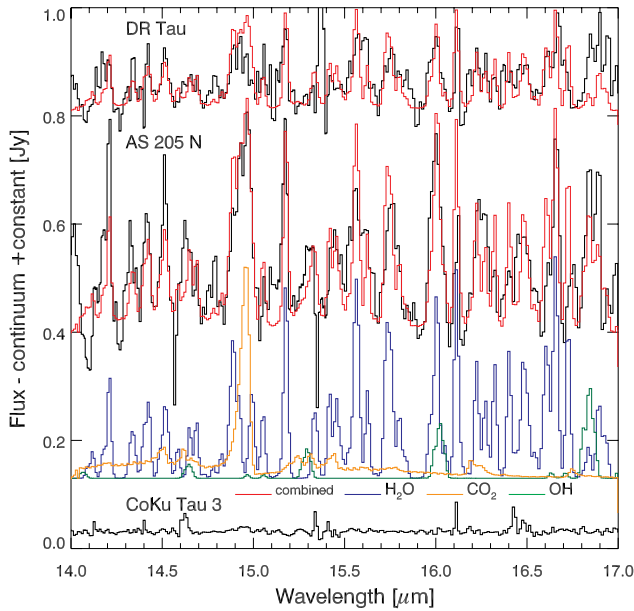


FIG. 1.—Portion of the *Spitzer* IRS data (black) and molecular emission models (red). The three components of the AS 205A model are shown in color below. The largely featureless spectrum of CoKu Tau/3 is included to demonstrate the dynamic range achievable by the IRS and the level of systematics, such as residuals from the defringing process. For both spectra, splines have been used to remove broad continuum features from silicates. Line-to-continuum excesses are  $\sim 5\%$ – $10\%$ .

temperature ( $T$ ) and column density ( $N$ ) gas in Keplerian rotation. Model parameters include  $T$ ,  $N$ , solid angle ( $\Omega$ ), stellar mass ( $M_*$ ), inclination ( $i$ ), and characteristic radius ( $r$ ). In addition, the local line broadening is assumed to be Gaussian, with specified width ( $\sigma$ ).  $M_*$  and  $i$  were fixed to values taken from the literature (Clarke & Bouvier 2000; Kitamura et al. 2002; Eisner et al. 2005; Andrews & Williams 2007). The spectrally unresolved IRS data cannot completely constrain the

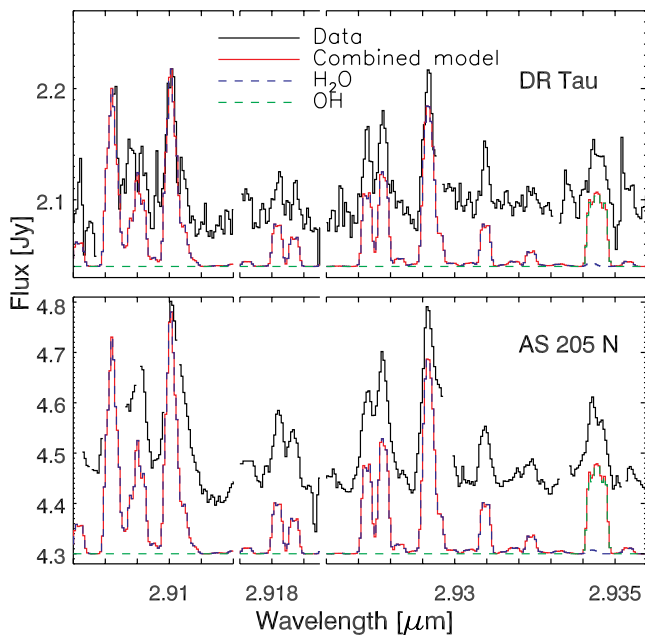


FIG. 2.—Comparison between portions of the NIRSPEC L-band data and best-fit  $\text{H}_2\text{O}/\text{OH}$  disk models (offset).

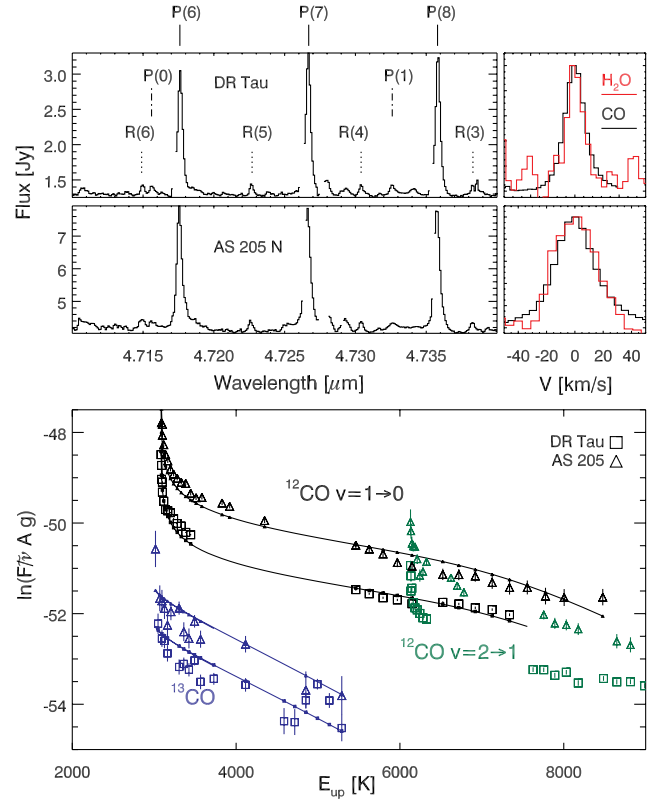


FIG. 3.—*Top left*: Portion of the NIRSPEC M-band spectra.  $^{12}\text{CO}/^{13}\text{CO}$   $v = 1 \rightarrow 0$  transitions are marked with solid and dashed lines, respectively.  $^{12}\text{CO}$   $v = 2 \rightarrow 1$  transitions are marked with dot-dashed lines. *Top right*: CO and  $\text{H}_2\text{O}$  emission lines are overlaid. *Bottom*: Excitation diagram with  $1\sigma$  error bars and best-fit models to the  $v = 1 \rightarrow 0$  transitions, with  $F$  in  $\text{W m}^{-2}$  and  $\bar{v}$  in  $\text{cm}^{-1}$ .

remaining model parameters, but do require high  $T$  ( $T \geq 500$  K) and  $N$  ( $N \geq 10^{17} \text{ cm}^{-2}$  for  $\sigma \sim 2 \text{ km s}^{-1}$ ). Figure 1 compares the IRS spectra with a model in which  $T$  and  $N$  are constrained using the NIRSPEC data described below, but where  $\Omega$  has been increased by a factor of 2.

The addition of  $\text{CO}_2$  (and, for AS 205A, OH) improves our fits, as shown by the AS 205A model in Figure 1. With  $T$  and  $\Omega$  fixed, we find  $\text{H}_2\text{O} : \text{CO}_2 = 40$  for DR Tau and 20 for AS 205. Detections of these and other molecules will be discussed in detail elsewhere.

### 3.2. NIRSPEC

While tropospheric water vapor blocks significant portions of incoming radiation near  $3 \mu\text{m}$ , there are a handful of rovibrational lines that are optically thin in the atmosphere but readily excited at temperatures typical of the inner  $\sim 1$  AU of circumstellar disks. This is illustrated in Figure 2, in which spectral regions with  $<70\%$  atmospheric transmission have been removed, yet many prominent (line-to-continuum  $\sim 5\%$ – $10\%$ )  $\text{H}_2\text{O}$  emission features remain. These are primarily  $v_3 = 1 \rightarrow 0$  lines, with upper level energies of  $\sim 5000$ – $10,000$  K. A few weaker (line-to-continuum  $\sim 1\%$ ) features are nominally detected from  $3.13$  to  $3.18 \mu\text{m}$  and are used as model constraints. We also detect the P4.5 ( $2.9344 \mu\text{m}$ ;  $T_{\text{up}} \sim 5400$  K) and P9.5 ( $3.1788 \mu\text{m}$ ;  $T_{\text{up}} \sim 7500$  K) OH doublets. The M-band spectra include portions of the  $^{12}\text{CO}/^{13}\text{CO}$   $v = 1 \rightarrow 0$  and  $^{12}\text{CO}$   $v = 2 \rightarrow 1$   $P$ - and  $R$ -branches (Fig. 3).

We fit all spectra with the ring models described in § 3.1.

Because the CO lines span a large range of excitation energies, across which they transition from optically thick to optically thin,  $N$ ,  $T$ , and  $\Omega$  can be determined uniquely. Also, a constraint on  $N_{\text{CO}}$  is provided by the <sup>13</sup>CO lines (assuming a <sup>12</sup>CO : <sup>13</sup>CO ratio, here fixed at 77; see Blake & Boogert 2004). Thus, we began by fitting the CO emission by comparison with an excitation diagram of the  $v = 1 \rightarrow 0$  transitions (Fig. 3), as described in Salyk et al. (2007). For our nominal H<sub>2</sub>O and OH fits, we then fixed  $\Omega$  and  $T$  to the CO-derived values to determine  $N$ —a reasonable, although not ideal, assumption, given the similarity in line shapes (see Fig. 3).

Our nominal ring model is shown in Figure 2, with best-fit parameters in Table 1, and uncertainties of order  $\Delta N \sim 30\%$ ,  $\Delta T \sim 15\%$ , and  $\Delta \Omega \sim 10\%$ . We have fixed  $r = 3$  AU to roughly match the emission profiles since the line shape does not affect the integrated flux; we discuss line shapes below.  $\Omega$  corresponds to areas of 0.2 and 0.4 AU<sup>2</sup> for DR Tau and AS 205, respectively, at distances of 140 and 120 pc. Uncertainties in the nominal H<sub>2</sub>O and OH fits are linked to CO uncertainties, but if  $T$  is instead allowed to vary, the H<sub>2</sub>O spectra are consistent with  $T \sim 900$ – $1200/900$ – $1100$  K for DR Tau/AS 205A and ranges in  $N$  of about 1 order of magnitude. For OH, we can only constrain  $T \geq 900$  K.

The model results depend crucially on the assumed local line width ( $\sigma$ ), which is unknown. For our nominal model, we have adopted  $\sigma = 2$  km s<sup>-1</sup>—the sound speed for H<sub>2</sub> at 1000 K. However, we have also tested values between  $\sim 0.3$  and 10 km s<sup>-1</sup>. We find that the opacity needs to remain similar from fit to fit, so that a change in  $\sigma^2$  requires a proportional change in  $N$ . To maintain the flux level, the total number of molecules must remain similar, so an increase in  $\sigma$  implies a decrease in  $\Omega$ . Accordingly, the CO : H<sub>2</sub>O ratio changes by no more than a factor of  $\sim 2$  between models. (If  $T_{\text{H}_2\text{O}}$  is not fixed, however, this ratio may change). In addition,  $T$  is only weakly affected by  $\sigma$ , since it is set by the overall shape of the spectrum, rather than by the absolute flux levels.

NIRSPEC spectrally resolves the molecular emission, allowing an investigation of the disk gas kinematics. In Figure 3 we compare the unblended H<sub>2</sub>O transition at 2.931  $\mu\text{m}$  with a CO line composite constructed from an average of all  $v = 1 \rightarrow 0$  lines. In general, Keplerian disks should produce double-peaked emission profiles, with peak separation set by the inner and outer emission radii ( $r_{\text{in}}$ ,  $r_{\text{out}}$ ), as well as the temperature gradient (Horne & Marsh 1986). The CO and H<sub>2</sub>O line profiles, however, are more Gaussian, or even Lorentzian in shape, with relatively wide wings and a narrow peak. The narrow core could be produced by a ring with a very large  $r_{\text{in}}$  ( $>10$  AU) if the peak-to-peak separation is equal to the NIRSPEC resolution (given the  $M_*$  and  $i$  in Table 1). However, this is difficult to reconcile with the high  $T$  derived from the model fits and the observed line wings.

Another way to produce a single peak would be to have the line emission extend to significant radii ( $r_{\text{out}}$  of order 7 AU, for a temperature profile that is constant with radius), such that the contribution from low-velocity portions of the disk infills the center of the line profile. However, this model produces too much flux to match observations. In addition, passive disk models predict a steep temperature decline with radius. A more promising solution may be that some flux is produced via fluorescence, in which molecules in the flared outer disk (typically out to  $\sim 5$  AU) intercept and reradiate inner disk IR continuum, thus filling in the line profiles at low velocity (Blake & Boogert 2004). With canonical disk flaring and temperature structure, we find that a resonance fluorescence model with 100% scattering efficiency

TABLE 1  
FITS TO NIRSPEC DATA

Parameter	DR Tau	AS 205A
CO $v_{\text{in}}, v_{3\sigma}$ (km s <sup>-1</sup> )	27, 32	35, 40
CO $r_{\text{in}}, r_{3\sigma}$ (AU)	0.8, 0.6	0.5, 0.4
H <sub>2</sub> O $v_{\text{in}}, v_{3\sigma}$ (km s <sup>-1</sup> )	24, 28	36, 42
H <sub>2</sub> O $r_{\text{in}}, r_{3\sigma}$ (AU)	1.0, 0.7	0.4, 0.3
OH $v_{\text{in}}, v_{3\sigma}$ (km s <sup>-1</sup> )	28, 32	34, 39
OH $r_{\text{in}}, r_{3\sigma}$ (AU)	0.7, 0.5	0.5, 0.4
$r^a$ (AU)	3	3 <sup>b</sup>
$\sigma^a$ (km s <sup>-1</sup> )	2	2
$M_*^a$ ( $M_{\odot}$ )	0.76	1.2
$i^a$ (deg)	67	47
$\Omega_{\text{CO,H}_2\text{O,OH}}$ (sr)	$2 \times 10^{-16}$	$6 \times 10^{-16}$
$T_{\text{CO,H}_2\text{O,OH}}$ (K)	1000	1000
$N_{\text{CO}}$ (cm <sup>-2</sup> )	$7 \times 10^{18}$	$6 \times 10^{18}$
$N_{\text{H}_2\text{O}}$ (cm <sup>-2</sup> )	$8 \times 10^{17}$	$6 \times 10^{17}$
$N_{\text{OH}}$ (cm <sup>-2</sup> )	$2 \times 10^{17}$	$2 \times 10^{17}$

<sup>a</sup> Fixed.

<sup>b</sup> Chosen to approximately match line shapes.

can indeed produce fluxes comparable to those of thermal emission; a more thorough investigation is left as future work.

In any case, the amount of flux in the line profile wings is consistent with emission arising in a Keplerian disk with a small inner radius. With this assumption, the line wing velocity can be used to estimate  $r_{\text{in}}$ . Because turbulent motion and disk structure can both have nonnegligible effects on the line shapes, we adopt here the velocity,  $v_{\text{in}}$ , at 2.2 times the half-width at half-maximum—a value found to be appropriate for the moderate-inclination disk around GM Aur (Salyk et al. 2007). To compute  $v_{\text{in}}$  we use Gaussian fits to the CO and H<sub>2</sub>O lines discussed above, as well as the average of two-Gaussian fits to the OH emission doublets, deconvolved with the NIRSPEC instrument response function. If turbulent velocities are low, radii associated with  $v_{\text{in}}$  will be overestimates, and so we also compute  $v$  at  $3\sigma$ , which should represent a lower limit to  $r_{\text{in}}$  (see Table 1).

#### 4. DISCUSSION

We have convincingly detected H<sub>2</sub>O emission in the 10–20  $\mu\text{m}$  region with *Spitzer* IRS as well as H<sub>2</sub>O and OH emission near 3  $\mu\text{m}$  with NIRSPEC, arising from the disk atmospheres surrounding two CTTSs. By combining these data with constraints from spectrally resolved CO  $v = 1 \rightarrow 0$  emission, we find that the excitation temperatures are typical of terrestrial planet-forming regions ( $\sim 1000$  K). In addition, line wing velocities imply inner emission radii no larger than 1 AU for all molecules. Therefore, further observations and analyses of these species may provide a new window into terrestrial planet-forming regions.

Observations of H<sub>2</sub>O in disks may also provide constraints on disk evolution and water transport. In a disk in which water vapor transport is controlled by diffusion, the presence of a condensation front (“snow line”) can potentially dry out an inner disk in as little as  $10^5$  yr, for typical disk viscosities (Stevenson & Lunine 1988). The snow line for the early solar system is estimated to be near 3 AU, but would be farther out for these high accretion rate stars (Lecar et al. 2006). Therefore, the presence of significant water vapor inside 1 AU may be evidence for inward radial migration, or upward mixing, of icy solids, followed by evaporation. With both solid migration and diffusion controlling water transport, water vapor concentrations become sensitive to an array of parameters, including disk viscosity and planetesimal growth rates, and water concentrations may trace disk evolution (Ciesla & Cuzzi 2006).

The detection of strong OH emission is also interesting, for OH is known to be an important ingredient to the chemistry in the inner regions of protoplanetary disks. It also acts as a coolant near the disk surface (e.g., Dullemond et al. 2007), and photolysis of OH to O(<sup>1</sup>D) has been invoked to explain the observed strengths of the 6300 Å O I line (Acke et al. 2005). Detections of OH in these and other disks will be discussed in greater detail in A. M. Mandell et al. (in preparation).

An understanding of the local line broadening in disks will be crucial for obtaining accurate absolute column densities of molecular species, as smaller local line widths require lower column densities and vice versa. Nevertheless, we find that molecular ratios are more robust, and remain similar when utilizing different disk models with the nominal fits having CO : H<sub>2</sub>O ~ 9–10 and H<sub>2</sub>O : OH ~ 3–4, provided  $T_{\text{H}_2\text{O}} = T_{\text{OH}} = T_{\text{CO}}$ . Stronger constraints on both  $T$  and  $N$  will require high line-to-continuum observations across a wider wavelength range, but the available data reveal that H<sub>2</sub>O : CO is similar to, or slightly higher than, that in the dense clouds out of which disks form (Boogert et al. 2004). Vertically integrated H<sub>2</sub>O : OH ratios have been predicted to be much higher (>10<sup>5</sup>) near 1 AU (Markwick et al. 2002), but significant OH abundances are expected near the C<sup>+</sup>/C/CO transition zone at large scale heights (Kamp & Dullemond 2004). Also, it is important to keep in mind that disk transport is not included in current chemical models.

Our analysis of the L- and M-band spectra thus far has assumed thermal excitation of the vibrationally excited states, but this need not be the case.  $\Delta v = 1$  emission can also be driven by the stellar UV output or by the absorption of IR photons from the star or disk, followed by fluorescence (Dello Russo et al. 2004; Blake & Boogert 2004). For OH, vibrationally excited emission can also result from the (photo)dissociation of H<sub>2</sub>O followed by prompt emission (Bonnev et al. 2006). In fact, fluorescence may be necessary to explain observed line profile shapes (see § 3.2), and could explain the surprisingly strong OH emission. A better understanding of these processes may therefore prove essential for relating observed line fluxes to disk column densities.

If emission in the line wings truly represents emission from the disk inner radius, relatively large inner emission radii are implied for CO, OH, and H<sub>2</sub>O. Both DR Tau and AS 205A have had their dust inner rim sizes measured with the Keck interfer-

ometer. At  $0.11 \pm 0.03$  and  $0.07 \pm 0.01$  AU, respectively, they are consistent with the location of dust sublimation (Akeson et al. 2005; Eisner et al. 2005), while the inner emission radii we derive are at least 3 times larger. Perhaps these high accretion rate stars can stir up the inner disk enough to erase the vertical temperature gradients necessary for formation of emission lines. Another possibility may be that replenishment rates cannot keep up with photodissociation at the disk surface.

The ease with which H<sub>2</sub>O and OH are now detected in the terrestrial planet-forming zones of certain disks holds the promise of more extensive and detailed studies. Although the ground-based echelle observations presented here probe only the inner disk atmospheres, future high spectral resolution follow-up studies of *Spitzer*-detected water emission from protoplanetary disks will provide constraints at a variety of disk radii and vertical depths. In addition, spectro-astrometric observations are being used to directly constrain emitting location and structure of the molecular gas (Pontoppidan et al. 2008). Eventually, a suite of constraints combined with non-LTE radiative transfer models will allow us to rigorously address such far-reaching questions as, How does the water vapor abundance vary within and between disks? What constraints can be put on models for the chemistry and transport of volatiles in disks, and what implications will these have for the early evolution of the solar system?

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