

Discrete Routh reduction

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Abstract

This paper develops the theory of Abelian Routh reduction for discrete mechanical systems and applies it to the variational integration of mechanical systems with Abelian symmetry. The reduction of variational Runge–Kutta discretizations is considered, as well as the extent to which symmetry reduction and discretization commute. These reduced methods allow the direct simulation of dynamical features such as relative equilibria and relative periodic orbits that can be obscured or difficult to identify in the unreduced dynamics. The methods are demonstrated for the dynamics of an Earth orbiting satellite with a non-spherical J_2 correction, as well as the double spherical pendulum. The J_2 problem is interesting because in the unreduced picture, geometric phases inherent in the model and those due to numerical discretization can be hard to distinguish, but this issue does not appear in the reduced algorithm, where one can directly observe interesting dynamical structures in the reduced phase space (the cotangent bundle of shape space), in which the geometric phases have been removed. The main feature of the double spherical pendulum example is that it has a non-trivial magnetic term in its reduced symplectic form. Our method is still efficient as it can directly handle the essential non-canonical nature of the symplectic structure. In contrast, a traditional symplectic method for canonical systems could require repeated coordinate changes if one is evoking Darboux’ theorem to transform the symplectic structure into canonical form, thereby incurring additional computational cost. Our method allows one to design reduced symplectic integrators in a natural way, despite the non-canonical nature of the symplectic structure.

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1. Introduction

This paper addresses reduction theory for discrete mechanical systems with Abelian symmetry groups and its relation to variational integration. To establish the setting of the problem, a few aspects of the continuous theory are first recalled (see [29] for general background).

1.1. Continuous reduction theory

Consider a mechanical system with configuration manifold Q and a symmetry group G (with Lie algebra \mathfrak{g}) acting freely and properly on Q and hence, by cotangent lift on T^*Q , with the corresponding (standard, equivariant) momentum map $\mathbf{J} : T^*Q \rightarrow \mathfrak{g}^*$. Recall from reduction theory (see [25, 36], and references therein) that, under appropriate regularity and nonsingularity conditions, the flow of a G -invariant Hamiltonian $H : T^*Q \rightarrow \mathbb{R}$ naturally induces a Hamiltonian flow on the reduced space $P_\mu = \mathbf{J}^{-1}(\mu)/G_\mu$, where G_μ is the isotropy subgroup of a chosen point $\mu \in \mathfrak{g}^*$. In the Abelian case, if one chooses a connection \mathcal{A} on the principal bundle $Q \rightarrow Q/G$, then P_μ is symplectically isomorphic to $T^*(Q/G)$ carrying the canonical symplectic structure modified by magnetic terms, that is terms induced from the μ -component of the curvature of \mathcal{A} .

The Lagrangian version of this theory is also well developed. In the Abelian case, it goes by the name of *Routh reduction* (see, for instance, [29], section 8.9). The reduced equations are again equations on $T(Q/G)$ and are obtained by dropping the variational principle, expressed in terms of the Routhian, from Q to Q/G . The non-Abelian version of this theory was originally developed in [32, 33], with important contributions and improvements given in [15, 30].

Of course, reduction has been enormously important for many topics in mechanics, such as stability and bifurcation of relative equilibria, integrable systems, etc. We need not review the importance of this process here as it is extensively documented in the literature.

1.2. Purpose, main results and examples

This paper presents the theory and illustrative numerical implementation for the reduction of discrete mechanical systems with Abelian symmetry groups. The discrete reduced space has a similar structure as in the continuous theory, but the curvature will be taken in a discrete sense. The paper studies two examples in detail, namely, satellite dynamics in the presence of the bulge of the Earth (the J_2 effect) and the double spherical pendulum (which has a non-trivial magnetic term). In each case the benefit of studying the numerics of the reduced problem is shown. Roughly, the reduced computations reveal dynamical structures that are hard to pick out in the unreduced dynamics in a way that is reminiscent of the phenomena of pattern evocation, as in [34, 35]. Another interesting application of the theory is that of orbiting multibody systems, studied in [41, 42].

We refer to [37] for a review of discrete mechanics, its numerical implementation, some history, as well as references to the literature. The value of geometric integrators has been documented in a number of references, such as [9]. In the present paper, we shall focus, to be specific, on discrete Euler–Lagrange and variational symplectic Runge–Kutta schemes and their reductions. One could, of course, use other schemes as well, such as Newmark, Störmer–Verlet or Shake schemes. However, we wish to emphasize that *without theoretical guidelines, coding algorithms for the reduced dynamics need not be a routine procedure since the reduced equations are not in canonical form because of non-trivial magnetic terms*. For example, using Darboux’ theorem to put the structure into canonical form so that standard algorithms

can be used is not practical. We also remind the reader that there are real advantages to taking the variational approach to the construction of symplectic integrators. For example, as in [23], the variational approach provides the design flexibility to take different time steps at spatially different points in an asynchronous way and still retain all the advantages of symplecticity even though the algorithms are not strictly symplectic in the naive sense; such an approach is well known to be useful in molecular systems, for instance.

1.3. Motivation for discrete reduction

Besides its considerable theoretical interest, there are several practical reasons for carrying out discrete Routh reduction. These are as follows.

- (i) Features that are clear in the reduced dynamics, such as relative equilibria and relative periodic orbits, can be obscured in the unreduced dynamics, and appear more complicated through the process of reconstruction and associated geometric phases. This is related to the phenomenon of *pattern evocation* that is an important practical feature of many examples, such as the double spherical pendulum [34, 35] and the stepping pendulum [12]. Going to a suitable (but non-obvious) rotating frame can ‘evolve’ such phenomena (see the movie at <http://www.cds.caltech.edu/~marsden/research/demos/movies/Wendlandt/pattern.mpg>). This is essentially a window to the reduced dynamics, which the theory in the present paper allows one to compute directly.
- (ii) While directly studying the reduced dynamics can yield some benefits, it can be difficult to code using traditional methods. In particular, the presence of magnetic terms in the reduced symplectic form, as is the case with the double spherical pendulum, means that traditional symplectic methods for canonical systems do not directly apply; if one attempts to do so, it may result (and has in the literature) in many inefficient coordinate changes when evoking Darboux’ theorem to put things into canonical form.
- (iii) Although simulating the reduced dynamics involves an initial investment of time in computing geometric quantities symbolically, these additional terms do not appreciably affect the sparsity of the system of equations to be solved. As such, direct coding of the reduced algorithms can be quite efficient, due to its reduced dimensionality.

1.4. Two obvious generalizations

The free and proper assumption that we make on the group action means that we are dealing with *nonsingular*, that is, *regular* reduction (see [38] for the general theory of singular reduction and references to the literature). It would be interesting to extend the work here to the case of singular reduction but already the regular case is non-trivial and interesting. While our examples have singular points and the dynamics near these points is interesting, there is no attempt to study this aspect in the present paper.

Secondly, it would be interesting to generalize the present work to the case of non-Abelian groups and to develop a discrete version of non-Abelian Routh reduction (as in [15, 30]). We believe that such a generalization will require the further development of the theory of discrete connections, which is currently part of the research effort on discrete differential geometry (see [20], and references therein). Other future directions are discussed in the conclusions.

1.5. Other discrete reduction results

We briefly summarize some related results that have been obtained in the area of reduction for discrete mechanics. First of all, there is the important case of *discrete Euler–Poincaré* and

Lie–Poisson reduction that were obtained in [2, 27, 28]. This theory is appropriate for rigid body mechanics, for instance.

Another important case is that of *discrete semidirect product reduction* that was obtained in [3, 4] and applied to the case of the heavy top, with interesting links to discrete elastica. This case is of interest in the present study since with the heavy top, as with the general theory of semidirect product reduction (see [13, 31]), one can view the S^1 reduction of this problem as Routh reduction. Linking these two approaches is an interesting topic for future research.

1.6. Outline

After recalling the notation from continuous reduction theory, section 2 develops discrete reduction theory, derives a reduced variational principle and proves the symplecticity of the reduced flow. The relationship between continuous- and discrete-time reduction is also discussed. How the variational (and hence symplectic) Runge–Kutta algorithm induces a reduced algorithm in a natural way is shown in section 3. In section 4, we put together in a coherent way the main theoretical results of the paper up to that point. In section 5, the numerical example of satellite dynamics about an oblate Earth is given, and in section 6, the example of the double spherical pendulum, which has a non-trivial magnetic term, is given. Lastly, in section 7, we address some computational and efficiency issues.

2. Discrete reduction

In this section, it is assumed that the reader is familiar with continuous reduction theory as well as the theory of discrete mechanics; reference is made to the relevant parts of the literature as needed. It will be useful to first recall some facts about discrete mechanical systems with symmetry (see, for instance, [37] for proofs).

2.1. Discrete mechanical systems with symmetry

Let G be a Lie group (which shortly will be assumed to be Abelian) that acts freely and properly (on the left) on a configuration manifold Q . Given a discrete Lagrangian $L_d : Q \times Q \rightarrow \mathbb{R}$ that is invariant under the diagonal action of G on $Q \times Q$, the corresponding *discrete momentum map* $\mathbf{J}_d : Q \times Q \rightarrow \mathfrak{g}^*$ is defined by

$$\mathbf{J}_d(q_0, q_1) \cdot \xi = D_2 L_d(q_0, q_1) \cdot \xi_Q(q_1), \quad (1)$$

where D_2 denotes the derivative in the second slot and where ξ_Q is the infinitesimal generator associated with $\xi \in \mathfrak{g}$. The map \mathbf{J}_d is equivariant with respect to the diagonal action of G on $Q \times Q$ and the coadjoint action on \mathfrak{g}^* . The *discrete Noether theorem* states that the discrete momentum is conserved along solutions of the DEL (discrete Euler–Lagrange) equations,

$$D_2 L_d(q_{k-1}, q_k) + D_1 L_d(q_k, q_{k+1}) = 0. \quad (2)$$

Note that

$$\mathbf{J}_d(q_0, q_1) \cdot \xi = \mathbf{J}(D_2 L_d(q_0, q_1)) \cdot \xi,$$

where $\mathbf{J} : T^*Q \rightarrow \mathfrak{g}^*$ is the momentum map on T^*Q ; i.e., $\mathbf{J}_d = \mathbf{J} \circ \mathbb{F}L_d$, where $\mathbb{F}L_d = D_2 L_d : Q \times Q \rightarrow T^*Q$ is the discrete Legendre transform. Thus, for $\mu \in \mathfrak{g}^*$, we have $\mathbb{F}L_d(\mathbf{J}_d^{-1}(\mu)) \subset \mathbf{J}^{-1}(\mu)$. The symplectic algorithm (usually called the *position-momentum form of the algorithm*) obtained on T^*Q from that on $Q \times Q$ via the discrete Legendre transform thus preserves the standard momentum map \mathbf{J} .

There will be a standing assumption in this paper, namely that the given discrete Lagrangian L_d is *regular*; that is, for a point $(q, q) \in Q \times Q$ on the diagonal, the iterated derivative $D_2 D_1 L_d(q, q) : T_q Q \times T_q Q \rightarrow \mathbb{R}$ is a non-degenerate bilinear form. By the implicit function theorem, this implies that a point (q_{k-1}, q_k) near the diagonal and the DEL equations (2) uniquely determine the subsequent point q_{k+1} in a neighbourhood of the diagonal in $Q \times Q$ (or, if one prefers, for small time steps); in other words, the DEL algorithm is well defined by the DEL equations. Regularity also implies that the discrete Legendre transformation $\mathbb{F}L_d = D_2 L_d : Q \times Q \rightarrow T^*Q$ is a local diffeomorphism from a neighbourhood of the diagonal in $Q \times Q$ to a neighbourhood in T^*Q . For a detailed discussion, see [37].

2.2. Reconstruction

The following lemma gives a basic result on the reconstruction of discrete curves in the configuration manifold Q from those in *shape space*, defined to be $S = Q/G$. The lemma is similar to its continuous counterpart, as in, for example, [15], lemma 2.2. The natural projection to the quotient will be denoted by $\pi_{Q,G} : Q \rightarrow Q/G$; $q \mapsto x = [q]_G$ (the equivalence class of $q \in Q$). Let $\text{Ver}(q)$ denote the *vertical space* at q , namely the space of all vectors at the point q that are infinitesimal generators $\xi_Q(q) \in T_q Q$ or, in other words, the tangent space to the group orbit through q . We say that the discrete Lagrangian L_d is *group-regular* if the bilinear map $D_2 D_1 L_d(q, q) : T_q Q \times T_q Q \rightarrow \mathbb{R}$ restricted to the subspace $\text{Ver}(q) \times \text{Ver}(q)$ is nondegenerate. In addition to regularity, we shall make group-regularity a standing assumption in the paper. The following result is fundamental for what follows.

Lemma 2.1 (Reconstruction lemma). *Fix $\mu \in \mathfrak{g}^*$ and let x_0, x_1, \dots, x_n be a sufficiently closely spaced discrete curve in S . Let $q_0, q_1 \in Q$ be such that $[q_0]_G = x_0$, $[q_1]_G = x_1$ and $\mathbf{J}_d(q_0, q_1) = \mu$. Then there is a unique closely spaced discrete curve q_1, q_2, \dots, q_n such that $[q_k]_G = x_k$ and $\mathbf{J}_d(q_{k-1}, q_k) = \mu$, for $k = 1, 2, \dots, n$.*

Proof. We must construct a point q_2 close to q_1 such that $[q_2]_G = x_2$ and $\mathbf{J}_d(q_1, q_2) = \mu$; the construction of the subsequent points q_3, \dots, q_n then proceeds in a similar fashion.

To do this, pick a local trivialization of the bundle $\pi_{Q,G} : Q \rightarrow Q/G$, where locally $Q = S \times G$, and write points in this trivialization as $q_0 = (x_0, g_0)$ and $q_1 = (x_1, g_1)$, etc. Given the points $q_0 = (x_0, g_0)$, $q_1 = (x_1, g_1)$ with $\mathbf{J}_d(q_0, q_1) = \mu$, we seek a near identity group element $k \in G$ such that $q_2 := (x_2, kg_1)$ satisfies $\mathbf{J}_d(q_1, q_2) = \mu$. According to equation (1), this means that we must satisfy the condition $D_2 L_d(q_1, q_2) \cdot \xi_Q(q_2) = \langle \mu, \xi \rangle$ for all $\xi \in \mathfrak{g}$. In the local trivialization, this reads

$$D_2 L_d((x_1, g_1), (x_2, kg_1)) \cdot (0, TR_{kg_1} \xi) = \langle \mu, \xi \rangle, \quad (3)$$

where R_g is right translation on G by g . Consider solving the equation

$$D_2 L_d((\bar{x}_1, \bar{g}_1), (\bar{x}_2, k\bar{g}_1)) \cdot (0, TR_{k\bar{g}_1} \xi) = \langle \mu, \xi \rangle, \quad (4)$$

for k as a function of the variables $\bar{g}_1, \bar{x}_1, \bar{x}_2$ with μ fixed. By assumption, there is a solution for the case $\bar{x}_1 = x_0, \bar{x}_2 = x_1$ and $\bar{g}_1 = g_0$, namely $k = k_0 = g_1 g_0^{-1}$ (a near identity group element). The implicit function theorem shows that when the point g_0, x_0, x_1 is replaced by the nearby point g_1, x_1, x_2 , there will be a unique solution for k near k_0 provided that the derivative of the defining relation (3) with respect to k at the identity is invertible, which is true by group regularity. Since group regularity is G -invariant, the above argument remains valid as k_i drifts from the identity. \square

Note that the above lemma makes no hypotheses about the sequences x_n or q_n satisfying any discrete evolution equations.

To carry out reconstruction in the continuous case, in addition to the requirements that the lifted curve in TQ lie on the momentum surface, and that it projects to the reduced curve $x(t) \in S$ under $\pi_{Q,G}$, one also requires that it be second order, which is to say that it is of the form $(q(t), \dot{q}(t))$. If a connection is given, then the lifted curve is obtained by integrating the *reconstruction equation*—again, see [15] for details. The discrete analogue of the second-order curve condition is explained as follows. Consider a given discrete curve as a sequence of points, $(x_0, x_1), (x_1, x_2), \dots, (x_{n-1}, x_n)$, in $S \times S$. Lift each of the points in $S \times S$ to the momentum surface $\mathbf{J}_d^{-1}(\mu) \subset Q \times Q$. This yields the sequence, $(q_0^0, q_1^0), (q_0^1, q_1^1), \dots, (q_0^{n-1}, q_1^{n-1})$, which is unique up to an overall diagonal group action. The discrete analogue of the second-order curve condition is that this sequence in $Q \times Q$ defines a discrete curve in Q , which corresponds to requiring that $q_1^k = q_0^{k+1}$, for $k = 0, \dots, n-1$, which is clearly possible in the context of the reconstruction lemma.

Discrete reconstruction naturally leads to issues of discrete geometric phases, and it would be interesting to express the discrete geometric phase in terms of the discrete curvature on shape space; this will surely involve some ideas from the currently evolving subject of discrete differential geometry and so we do not attempt to push this idea further at this point.

While many of the computations we present in this paper are in the setting of local trivializations, the results are valid globally through the construction given below.

2.3. Identification of the quotient space

Now assume that G is Abelian so that $G = G_\mu$ acts on $\mathbf{J}_d^{-1}(\mu)$ and that the quotient space $\mathbf{J}_d^{-1}(\mu)/G$ makes sense. We assume the above regularity hypotheses and freeness and properness of the action of G so that this quotient is a smooth manifold. It is clear that the map $\varphi_\mu : \mathbf{J}_d^{-1}(\mu)/G \rightarrow S \times S$ given by $[(q, q')]_G \rightarrow (x, x')$ is well defined, where the square brackets denote the equivalence class with respect to the given G action, and where $x = [q]_G$. The argument given in the reconstruction lemma shows that for a point $(q_0, q_1) \in \mathbf{J}_d^{-1}(\mu)$, φ_μ is a local diffeomorphism in a neighbourhood of the point $[(q_0, q_1)]_G$. In fact, the uniqueness part of that argument shows that for two nearby points (q_1, q_2) and (q'_1, q'_2) in $\mathbf{J}_d^{-1}(\mu)$, if $q_1 = g_1 q'_1$ and $q_2 = g_2 q'_2$, then $g_1 = g_2$. Thus, there is a neighbourhood U of a given chain of closely spaced points lying in $\mathbf{J}_d^{-1}(\mu)$ with this property. Saturating this neighbourhood with the group action, we can assume that U is G -invariant. Restricted to U , φ_μ becomes a diffeomorphism to a neighbourhood of the diagonal of $S \times S$.

Assume, as above, that $L_d : Q \times Q \rightarrow \mathbb{R}$ is a discrete Lagrangian that is invariant under the action of an Abelian Lie group G on $Q \times Q$. In view of the preceding discussion, L_d restricted to $\mathbf{J}_d^{-1}(\mu)$ (and in the neighbourhood of a given chain of points in this set) induces a well-defined function $\hat{L}_d(x_0, x_1)$ of pairs of points (x_0, x_1) in $S \times S$. This *discrete reduced Lagrangian* will play an important role in what follows.

2.4. Discrete reduction

Let $\mathbf{q} := \{q_0, \dots, q_n\}$ be a solution of the discrete Euler–Lagrange (DEL) equations. Let the value of the discrete momentum along this trajectory be μ . Let $x_i = [q_i]_G$, so that $\mathbf{x} := \{x_0, \dots, x_n\}$ is a discrete shape space trajectory. Since \mathbf{q} satisfies the discrete variational principle, it is appropriate to ask if there is a reduced variational principle satisfied by \mathbf{x} .

An important issue in dropping the discrete variational principle to the shape space is whether we require that the varied curves are constrained to lie on the level set of the momentum

map. The constrained approach is adopted in [15], and the unconstrained approach is used in [30]. In the rest of this section, we will adopt the unconstrained approach of [30] and show that the variations in the discrete action sum, when evaluated on a solution of the discrete Euler–Lagrange equations without assuming that the variations at the endpoints vanish, depends only on the quotient variations, and therefore drop to the shape space without constraints on the variations.

If \mathbf{q} is a solution of the DEL equations, the interior terms in the variation of the discrete action sum vanish, leaving only the boundary terms; that is,

$$\delta \sum_{k=0}^{n-1} L_d(q_k, q_{k+1}) = D_1 L_d(q_0, q_1) \cdot \delta q_0 + D_2 L_d(q_{n-1}, q_n) \cdot \delta q_n. \quad (5)$$

Given a principal connection \mathfrak{A} on Q , there is a horizontal–vertical split of each tangent space to Q denoted by $v_q = \text{hor } v_q + \text{ver } v_q$ for $v_q \in T_q Q$. Thus,

$$D_2 L_d(q_{n-1}, q_n) \cdot \delta q_n = D_2 L_d(q_{n-1}, q_n) \cdot \text{hor } \delta q_n + D_2 L_d(q_{n-1}, q_n) \cdot \text{ver } \delta q_n.$$

As in continuous Routh reduction, we will rewrite the terms involving vertical variations using the fact that we are on a level set of \mathbf{J}_d . Namely, write the vertical variation as $\text{ver } \delta q_n = \xi_Q(q_n)$, where $\xi = \mathfrak{A}(\delta q_n)$ and use definition (1) of \mathbf{J}_d to give

$$\begin{aligned} D_2 L_d(q_{n-1}, q_n) \cdot \text{ver } \delta q_n &= D_2 L_d(q_{n-1}, q_n) \cdot \xi_Q(q_n) = \mathbf{J}_d(q_{n-1}, q_n) \cdot \xi \\ &= \langle \mu, \xi \rangle = \langle \mu, \mathfrak{A}(\delta q_n) \rangle = \mathfrak{A}_\mu(q_n) \cdot \delta q_n. \end{aligned} \quad (6)$$

Thus, the boundary terms can be expressed as

$$D_2 L_d(q_{n-1}, q_n) \cdot \delta q_n = D_2 L_d(q_{n-1}, q_n) \cdot \text{hor } \delta q_n + \mathfrak{A}_\mu(q_n) \cdot \delta q_n, \quad (7)$$

$$D_1 L_d(q_0, q_1) \cdot \delta q_0 = D_1 L_d(q_0, q_1) \cdot \text{hor } \delta q_0 - \mathfrak{A}_\mu(q_0) \cdot \delta q_0, \quad (8)$$

and so (5), the variation of the discrete action sum, becomes

$$\begin{aligned} \delta \sum_{k=0}^{n-1} L_d(q_k, q_{k+1}) &= D_1 L_d(q_0, q_1) \cdot \text{hor } \delta q_0 + D_2 L_d(q_{n-1}, q_n) \cdot \text{hor } \delta q_n \\ &\quad + \mathfrak{A}_\mu(q_n) \cdot \delta q_n - \mathfrak{A}_\mu(q_0) \cdot \delta q_0, \end{aligned} \quad (9)$$

when restricted to solutions of the discrete Euler–Lagrange equations.

Motivated by the preceding equation, introduce the 1-form \mathcal{A} on $Q \times Q$ defined by

$$\mathcal{A} = \pi_2^* \mathfrak{A}_\mu - \pi_1^* \mathfrak{A}_\mu, \quad (10)$$

where $\pi_1, \pi_2 : Q \times Q \rightarrow Q$ are projections onto the first and the second components respectively. This allows us to expand the boundary terms involving \mathfrak{A}_μ into a telescoping sum, and rewrite (9) in terms of the 1-form \mathcal{A} as

$$\sum_{k=0}^{n-1} (DL_d - \mathcal{A})(q_k, q_{k+1}) \cdot (\delta q_k, \delta q_{k+1}) = D_1 L_d(q_0, q_1) \cdot \text{hor } \delta q_0 + D_2 L_d(q_{n-1}, q_n) \cdot \text{hor } \delta q_n. \quad (11)$$

We now drop (11) to the reduced space $S \times S$. Consider the projection maps $\pi : Q \times Q \rightarrow (Q \times Q)/G$ and $\pi_{\mu,d} : \mathbf{J}_d^{-1}(\mu) \rightarrow \mathbf{J}_d^{-1}(\mu)/G$, and the inclusion maps $\iota_{\mu,d} : \mathbf{J}_d^{-1}(\mu) \hookrightarrow Q \times Q$ and $\tilde{\iota}_{\mu,d} : \mathbf{J}_d^{-1}(\mu)/G \hookrightarrow (Q \times Q)/G$. Then clearly, the following diagram commutes:

$$\begin{array}{ccc} \mathbf{J}_d^{-1}(\mu) & \xrightarrow{\iota_{\mu,d}} & Q \times Q \\ \pi_{\mu,d} \downarrow & & \downarrow \pi \\ \mathbf{J}_d^{-1}(\mu)/G & \xrightarrow{\tilde{\iota}_{\mu,d}} & (Q \times Q)/G \end{array}$$

By the G -invariance, L_d drops to a function \tilde{L}_d on the quotient $(Q \times Q)/G$ so that $L_d = \tilde{L}_d \circ \pi$. The pullback (in this case, the restriction) of \tilde{L}_d to $\mathbf{J}_d^{-1}(\mu)/G$ is called the *discrete reduced Lagrangian* and is denoted by as \hat{L}_d . Thus, $\hat{L}_d = \tilde{L}_d \circ \iota_{\mu,d}$; identifying $\mathbf{J}_d^{-1}(\mu)/G$ with $S \times S$ this definition agrees with \hat{L}_d as defined earlier.

Lemma 2.2. *The 1-form \mathcal{A} on $Q \times Q$ restricted to $\mathbf{J}_d^{-1}(\mu)$ drops to a 1-form $\hat{\mathcal{A}}$ on $\mathbf{J}_d^{-1}(\mu)/G$ and induces (for closely spaced points), via the map φ_μ , a 1-form that we denote by the same letter, on $S \times S$. Similarly, the 1-form $D\tilde{L}_d - \mathcal{A}$ on $Q \times Q$ restricted to the momentum level set $\mathbf{J}_d^{-1}(\mu)$ then drops to the 1-form $D\hat{L}_d - \hat{\mathcal{A}}$ on $\mathbf{J}_d^{-1}(\mu)/G$ and induces, via the map φ_μ , a 1-form that we denote by the same letter, on $S \times S$.*

Proof. The equivariance of the projections π_i with respect to the diagonal action on $Q \times Q$ and the given action on Q , together with the invariance of the 1-form \mathfrak{A}_μ on Q , implies that \mathcal{A} is invariant. Since \mathfrak{A}_μ vanishes on vertical vectors for the bundle $Q \rightarrow Q/G$, it follows that \mathcal{A} vanishes on vertical vectors for the bundle $Q \times Q \rightarrow (Q \times Q)/G$. Therefore, there is a 1-form $\tilde{\mathcal{A}}$ on $(Q \times Q)/G$ such that $\mathcal{A} = \pi^* \tilde{\mathcal{A}}$.

Since $L_d = \tilde{L}_d \circ \pi$, and the exterior derivative commutes with pullback, it follows that $dL_d = \pi^* d\tilde{L}_d$. From $\pi \circ \iota_{\mu,d} = \tilde{\iota}_{\mu,d} \circ \pi_{\mu,d}$, we get $\iota_{\mu,d}^* \mathcal{A} = \iota_{\mu,d}^* \pi^* \tilde{\mathcal{A}} = \pi_{\mu,d}^* \tilde{\iota}_{\mu,d}^* \tilde{\mathcal{A}}$. Thus, the 1-form \mathcal{A} restricted to $\mathbf{J}_d^{-1}(\mu)$ drops to the 1-form $\hat{\mathcal{A}} = \tilde{\iota}_{\mu,d}^* \tilde{\mathcal{A}}$ on $\mathbf{J}_d^{-1}(\mu)/G$. Similarly, $\iota_{\mu,d}^* dL_d = \pi_{\mu,d}^* \tilde{\iota}_{\mu,d}^* d\tilde{L}_d$ and so dL_d restricted to $\mathbf{J}_d^{-1}(\mu)$ drops to the 1-form $\tilde{\iota}_{\mu,d}^* d\tilde{L}_d = d\hat{L}_d$ on $\mathbf{J}_d^{-1}(\mu)/G$. These 1-forms push forward under the map $\varphi_\mu : \mathbf{J}_d^{-1}(\mu)/G \rightarrow S \times S$ in the manner that was explained earlier. \square

With the preceding lemma, and equation (11), we conclude that

$$\begin{aligned} \sum_{k=0}^{n-1} (D\hat{L}_d - \hat{\mathcal{A}})(x_k, x_{k+1}) \cdot (\delta x_k, \delta x_{k+1}) &= D_1 L_d(q_0, q_1) \cdot \text{hor } \delta q_0 \\ &\quad + D_2 L_d(q_{n-1}, q_n) \cdot \text{hor } \delta q_n. \end{aligned} \quad (12)$$

Assuming that $\delta \mathbf{x}$ vanishes at the endpoints, $\text{hor } \delta q_0 = 0$, and $\text{hor } \delta q_1 = 0$ and consequently, the boundary terms vanish and we obtain the *reduced discrete variational principle*

$$\delta \sum_{k=0}^{n-1} \hat{L}_d(x_k, x_{k+1}) = \sum_{k=0}^{n-1} \hat{\mathcal{A}}(x_k, x_{k+1}) \cdot (\delta x_k, \delta x_{k+1}). \quad (13)$$

In an analogous fashion to rewriting $D\hat{L}_d(x_k, x_{k+1}) \cdot (\delta x_k, \delta x_{k+1})$ as $D_1 \hat{L}_d(x_k, x_{k+1}) \cdot \delta x_k + D_2 \hat{L}_d(x_k, x_{k+1}) \cdot \delta x_{k+1}$, we do the same for the $\hat{\mathcal{A}}$ term by defining

$$\hat{\mathcal{A}}(x_0, x_1) \cdot (\delta x_0, \delta x_1) = \hat{\mathcal{A}}_1(x_0, x_1) \cdot \delta x_0 + \hat{\mathcal{A}}_2(x_0, x_1) \cdot \delta x_1.$$

Then, equating terms involving δx_k on the left-hand side of (13) to the corresponding terms on the right-hand side, we get the *discrete Routh (DR) equations* giving dynamics on $S \times S$:

$$D_2 \hat{L}_d(x_{k-1}, x_k) + D_1 \hat{L}_d(x_k, x_{k+1}) = \hat{\mathcal{A}}_2(x_{k-1}, x_k) + \hat{\mathcal{A}}_1(x_k, x_{k+1}). \quad (14)$$

Note that these equations depend on the value of momentum μ . Thus, if \mathbf{q} is a discrete curve satisfying the discrete Euler–Lagrange equations, the curve \mathbf{x} obtained by projecting \mathbf{q} down to S satisfies the DR equations (14).

Now consider the converse, the discrete reconstruction procedure: given a discrete curve \mathbf{x} on S that satisfies the DR equations, is \mathbf{x} the projection of a discrete curve \mathbf{q} on Q that satisfies the DEL equations?

Let the pair (q_0, q_1) be a lift of (x_0, x_1) such that $\mathbf{J}_d(q_0, q_1) = \mu$. Let $\mathbf{q} = \{q_0, \dots, q_n\}$ be the solution of the DEL equations with initial condition (q_0, q_1) . Note that \mathbf{q} has momentum

μ . Let $\mathbf{x}' = \{x'_0, \dots, x'_n\}$ be the curve on S obtained by projecting \mathbf{q} . By our arguments above, \mathbf{x}' solves the DR equations. However, \mathbf{x}' has the initial condition (x_0, x_1) , which is the same as the initial condition of \mathbf{x} . By uniqueness of the solutions of the DR equations, $\mathbf{x}' = \mathbf{x}$. Thus, \mathbf{x} is the projection of a solution \mathbf{q} of the DEL equations with momentum μ . Also, for a given initial condition q_0 , there is a unique lift of \mathbf{x} to a curve with momentum μ . Such a lift can be constructed using the method described in lemma 2.1. Thus, lifting \mathbf{x} to a curve with momentum μ yields a solution of the discrete Euler–Lagrange equations, which projects down to \mathbf{x} .

We summarize the results of this section in the following theorem.

Theorem 2.3. *Let \mathbf{x} be a discrete curve on S , and let \mathbf{q} be a discrete curve on Q with momentum μ that is obtained by lifting \mathbf{x} . Then the following are equivalent.*

1. \mathbf{q} solves the DEL equations.
2. \mathbf{q} is a solution of the discrete Hamilton's variational principle

$$\delta \sum_{k=0}^{n-1} L_d(q_k, q_{k+1}) = 0$$

for all variations $\delta \mathbf{q}$ of \mathbf{q} that vanish at the endpoints.

3. \mathbf{x} solves the DR equations

$$D_2 \hat{L}_d(x_{k-1}, x_k) + D_1 \hat{L}_d(x_k, x_{k+1}) = \hat{A}_2(x_{k-1}, x_k) + \hat{A}_1(x_k, x_{k+1}).$$

4. \mathbf{x} is a solution of the reduced variational principle

$$\delta \sum_{k=0}^{n-1} \hat{L}_d(x_k, x_{k+1}) = \sum_{k=0}^{n-1} \hat{A}(x_k, x_{k+1}) \cdot (\delta x_k, \delta x_{k+1})$$

for all variations $\delta \mathbf{x}$ of \mathbf{x} that vanish at the endpoints.

Note that for smooth group actions, the order of accuracy will be equal for the reduced and unreduced algorithms.

2.5. Preservation of the reduced discrete symplectic form

The DR equations define a discrete flow map $\hat{F}_k : S \times S \rightarrow S \times S$. We already know that the flow of the DEL equations preserves the symplectic form Ω_{L_d} on $Q \times Q$. In this section, we show that the reduced flow \hat{F}_k preserves a *reduced* symplectic form $\Omega_{\mu,d}$ on $S \times S$, and that this reduced symplectic form is obtained by restricting Ω_{L_d} to $\mathbf{J}_d^{-1}(\mu)$ and then dropping to $S \times S$. In other words, $\pi_{\mu,d}^* \Omega_{\mu,d} = \iota_{\mu,d}^* \Omega_{L_d}$. The continuous analogue of this equation is $\pi_\mu^* \Omega_\mu = i_\mu^* \Omega_Q$.

Recall from continuous reduction theory on the Hamiltonian side that in the case of cotangent bundles, the projection $\pi_\mu : J^{-1}(\mu) \rightarrow T^*S$ can be defined as follows: if $\alpha_q \in J^{-1}(\mu)$, then the *momentum shift* $\alpha_q - \mathfrak{A}_\mu(q)$ annihilates all vertical tangent vectors at $q \in Q$, as shown by the following calculation:

$$\langle \alpha_q - \mathfrak{A}_\mu(q), \xi_Q(q) \rangle = J(\alpha_q) \cdot \xi - \langle \mu, \xi \rangle = \langle \mu, \xi \rangle - \langle \mu, \xi \rangle = 0.$$

Thus, $\alpha_q - \mathfrak{A}_\mu(q)$ induces an element of T_x^*S and $\pi_\mu(\alpha_q)$ is defined to be this element.

Let $\mathbb{F}' : \mathbf{J}_d^{-1}(\mu) \rightarrow J^{-1}(\mu)$ be the restriction of $\mathbb{F}L_d$ to $\mathbf{J}_d^{-1}(\mu)$. Thus $\mathbb{F}' \circ \iota_\mu = \iota_{\mu,d} \circ \mathbb{F}L_d$, where $\iota_\mu : J^{-1}(\mu) \rightarrow T^*Q$ and $\iota_{\mu,d} : \mathbf{J}_d^{-1}(\mu) \rightarrow Q \times Q$ are inclusions. Define the map $\hat{\mathbb{F}} : S \times S \rightarrow T^*S$ by $\hat{\mathbb{F}}(x_0, x_1) = D_2 \hat{L}_d(x_0, x_1) - \hat{A}_2(x_0, x_1)$. By equation (6) and lemma 2.2, this map is well defined. The map $\hat{\mathbb{F}}$ will play the role of a reduced discrete Legendre transform,

and in contrast to the continuous theory, the momentum shift appears in the map $\hat{\mathbb{F}}$, as opposed to the projection $\pi_{\mu,d}$. As in the continuous theory, $\Omega_{\mu,d}$ does involve magnetic terms.

Lemma 2.4. *The following diagram commutes:*

$$\begin{array}{ccc} \mathbf{J}_d^{-1}(\mu) & \xrightarrow{\mathbb{F}'} & J^{-1}(\mu) \\ \pi_{\mu,d} \downarrow & & \downarrow \pi_\mu \\ S \times S & \xrightarrow{\hat{\mathbb{F}}} & T^*S \end{array}$$

Proof. Let $(q_0, q_1) \in \mathbf{J}_d^{-1}(\mu)$. Thus $D_2L_d(q_0, q_1) \in J^{-1}(\mu)$, and

$$\pi_\mu(\mathbb{F}'(q_0, q_1)) = \pi_\mu(D_2L_d(q_0, q_1)).$$

As noted above, $\pi_\mu(D_2L_d(q_0, q_1))$ is the element of $T_{x_1}^*S$ determined by $(D_2L_d(q_0, q_1) - \mathfrak{A}_\mu(q_1))$. For $\delta q_1 \in T_{q_1}Q$, we have

$$\begin{aligned} \langle D_2L_d(q_0, q_1) - \mathfrak{A}_\mu(q_1), \delta q_1 \rangle &= (DL_d - \mathcal{A})(q_0, q_1) \cdot (0, \delta q_1) \\ &= (D\hat{L}_d - \hat{\mathcal{A}})(x_0, x_1) \cdot (0, \delta x_1) \\ &= D_2\hat{L}_d(x_0, x_1) \cdot \delta x_1 - \hat{\mathcal{A}}_2(x_0, x_1) \cdot \delta x_1, \end{aligned}$$

where in the second equality, we used lemma 2.2. Thus,

$$\pi_\mu(D_2L_d(q_0, q_1)) = D_2\hat{L}_d(x_0, x_1) - \hat{\mathcal{A}}_2(x_0, x_1),$$

which means $\hat{\mathbb{F}} \circ \pi_{\mu,d} = \pi_\mu \circ \mathbb{F}'$. □

Theorem 2.5. *The flow of the DR equations preserves the symplectic form*

$$\Omega_{\mu,d} = \hat{\mathbb{F}}^*(\Omega_S - \pi_{T^*S,S}^*\beta_\mu),$$

where β_μ is the 2-form on S obtained by dropping $\mathfrak{d}\mathfrak{A}_\mu$.

Furthermore, $\Omega_{\mu,d}$ can be obtained by dropping to $S \times S$ the restriction of Ω_{L_d} to $\mathbf{J}_d^{-1}(\mu)$. In other words, $\pi_{\mu,d}^*\Omega_{\mu,d} = \iota_{\mu,d}^*\Omega_{L_d}$.

Proof. We give an outline of the steps involved; the details are routine to fill in. The strategy is to first show that the restriction to $\mathbf{J}_d^{-1}(\mu)$ of the symplectic form Ω_{L_d} drops to a 2-form $\Omega_{\mu,d}$ on $S \times S$. The fact that the discrete flow on $Q \times Q$ preserves the symplectic form Ω_{L_d} is then used to show that the reduced flow preserves $\Omega_{\mu,d}$. The steps involved are as follows.

- (i) Consider the 1-form Θ_{L_d} on $Q \times Q$ defined by $\Theta_{L_d}(q_0, q_1) \cdot (\delta q_0, \delta q_1) = D_2L_d(q_0, q_1) \cdot \delta q_1$. The 1-form Θ_{L_d} is G -invariant, and thus the Lie derivative $\mathcal{L}_{\xi_{Q \times Q}} \Theta_{L_d}$ is zero.
- (ii) Since $\Omega_{L_d} = -\mathfrak{d}\Theta_{L_d}$, Ω_{L_d} is G -invariant. If $\iota_{\mu,d} : \mathbf{J}_d^{-1}(\mu) \rightarrow Q \times Q$ is the inclusion, $\Theta'_{L_d} = \iota_{\mu,d}^*\Theta_{L_d}$ and $\Omega'_{L_d} = \iota_{\mu,d}^*\Omega_{L_d}$ are the restrictions of Θ_{L_d} and Ω_{L_d} respectively to $\mathbf{J}_d^{-1}(\mu)$. One checks that Θ'_{L_d} and Ω'_{L_d} are invariant under the action of G on $\mathbf{J}_d^{-1}(\mu)$.
- (iii) If $\xi_{\mathbf{J}_d^{-1}(\mu)}$ is an infinitesimal generator on $\mathbf{J}_d^{-1}(\mu)$, then

$$\xi_{\mathbf{J}_d^{-1}(\mu)} \lrcorner \Omega'_{L_d} = -\xi_{\mathbf{J}_d^{-1}(\mu)} \lrcorner \mathfrak{d}\Theta'_{L_d} = -\mathcal{L}_{\xi_{\mathbf{J}_d^{-1}(\mu)}} \Theta'_{L_d} + \mathfrak{d}\xi_{\mathbf{J}_d^{-1}(\mu)} \lrcorner \Theta'_{L_d} = 0.$$

This follows from the G -invariance of Θ'_{L_d} , and the fact that $\Theta'_{L_d} \cdot \xi_{\mathbf{J}_d^{-1}(\mu)} = \langle \mu, \xi \rangle$.

- (iv) By steps 2 and 3, the form Ω'_{L_d} drops to a reduced form $\Omega_{\mu,d}$ on $\mathbf{J}_d^{-1}(\mu)/G \approx S \times S$. Thus, if $\pi_{\mu,d} : \mathbf{J}_d^{-1}(\mu) \rightarrow S \times S$ is the projection, then $\pi_{\mu,d}^*\Omega_{\mu,d} = \Omega'_{L_d}$. Note that the closure of $\Omega_{\mu,d}$ follows from the fact that Ω'_{L_d} is closed, which in turn follows from the closure of Ω_{L_d} and the relation $\Omega'_{L_d} = \iota_{\mu,d}^*\Omega_{L_d}$.

- (v) If $F_k : Q \times Q \rightarrow Q \times Q$ is the flow of the DEL equations, let F'_k be the restriction of this flow to $\mathbf{J}_d^{-1}(\mu)$. We know that F'_k drops to the flow \hat{F}_k of the DR equations on $S \times S$. Since F_k preserves Ω_{L_d} , F'_k preserves Ω'_{L_d} . Using this, it can be shown that \hat{F}_k preserves $\Omega_{\mu,d}$. Note that it is sufficient to show that $\pi_{\mu,d}^*(\hat{F}_k^* \Omega_{\mu,d}) = \pi_{\mu,d}^* \Omega_{\mu,d}$.
- (vi) It now remains to compute a formula for the reduced form $\Omega_{\mu,d}$. Using lemma 2.4, it follows that

$$\begin{aligned} \pi_{\mu,d}^* \Omega_{\mu,d} &= \iota_{\mu,d}^* \Omega_{L_d} = \iota_{\mu,d}^* \mathbb{F} L_d^* \Omega_Q = (\mathbb{F}')^* \iota_{\mu}^* \Omega_Q \\ &= (\mathbb{F}')^* \pi_{\mu}^* (\Omega_S - \pi_{T^*S,S}^* \beta_{\mu}) = \pi_{\mu,d}^* \hat{\mathbb{F}}^* (\Omega_S - \pi_{T^*S,S}^* \beta_{\mu}). \end{aligned}$$

Thus $\pi_{\mu,d}^* \Omega_{\mu,d} = \pi_{\mu,d}^* \hat{\mathbb{F}}^* (\Omega_S - \pi_{T^*S,S}^* \beta_{\mu})$, from which it follows that $\Omega_{\mu,d} = \hat{\mathbb{F}}^* (\Omega_S - \pi_{T^*S,S}^* \beta_{\mu})$. Incidentally, this expression shows that $\Omega_{\mu,d}$ is nondegenerate provided the map $\hat{\mathbb{F}} = D_2 \hat{L}_d - \hat{A}_2$ is a local diffeomorphism. \square

One can alternatively prove symplecticity of the reduced flow directly from the reduced variational principle—see section 2.3.4 of [21].

2.6. Relating discrete and continuous reduction

As shown in [37], if the discrete Lagrangian L_d approximates the Jacobi solution of the Hamilton–Jacobi equation, then the DEL equations give us an integration scheme for the EL equations. In our commutative diagrams we will denote the relationship between the EL and DEL equations by a dashed arrow as follows:

$$(TQ, EL) \dashrightarrow (Q \times Q, DEL).$$

This arrow can thus be read as ‘the corresponding discretization’. By the continuous and discrete Noether theorems, we can restrict the flow of the EL and DEL equations to $J_L^{-1}(\mu)$ and $\mathbf{J}_d^{-1}(\mu)$, respectively. The flow on $J_L^{-1}(\mu)$ induces a reduced flow on $J_L^{-1}(\mu)/G \approx TS$, which is the flow of the Routh equations. Similarly, the discrete flow on $\mathbf{J}_d^{-1}(\mu)$ induces a reduced discrete flow on $\mathbf{J}_d^{-1}(\mu)/G \approx S \times S$, which is the flow of the discrete Routh equations. Since the DEL equations give us an integration algorithm for the EL equations, it follows that the DR equations give us an integration algorithm for the Routh equations.

To numerically integrate the Routh equations, we have two options.

- (i) First solve the DEL equations to yield a discrete trajectory on Q , which can then be projected to a discrete trajectory on S .
- (ii) Solve the DR equations to directly obtain a discrete trajectory on Q .

Either approach yields the same result, which we express by the commutative diagram:

$$\begin{array}{ccc} (J_L^{-1}(\mu), EL) & \dashrightarrow & (\mathbf{J}_d^{-1}(\mu), DEL) \\ \pi_{\mu,L} \downarrow & & \downarrow \pi_{\mu,d} \\ (TS, R) & \dashrightarrow & (S \times S, DR) \end{array} \quad (15)$$

3. Reduction of the symplectic Runge–Kutta algorithm

A well-studied class of numerical schemes for Hamiltonian and Lagrangian systems is the symplectic partitioned Runge–Kutta (SPRK) algorithms (see [10, 11] for history and details).

We will adopt a local trivialization to express the SPRK method in which the group action is addition. Given a connection \mathfrak{A} on Q , it can be represented in local coordinates as $\mathfrak{A}(\theta, x)(\dot{\theta}, \dot{x}) = A(x)\dot{x} + \dot{\theta}$. By rewriting the symplectic partitioned Runge–Kutta algorithm in terms of this local trivialization, and using the local representation of the connection, we obtain the following algorithm on T^*S :

$$x_1 = x_0 + h \sum b_j \dot{X}_j \quad (16a)$$

$$s_1 = s_0 + h \sum_j \tilde{b}_j \dot{S}_j + \left[h \sum_j \left(\tilde{b}_j \mu \frac{\partial A}{\partial x}(X_j) \dot{X}_j \right) - (\mu A(x_1) - \mu A(x_0)) \right] \quad (16b)$$

$$X_i = x_0 + h \sum a_{ij} \dot{X}_j \quad (16c)$$

$$S_i = s_0 + h \sum_j \tilde{a}_{ij} \dot{S}_j + \left[h \sum_j \left(\tilde{a}_{ij} \mu \frac{\partial A}{\partial x}(X_j) \dot{X}_j \right) - (\mu A(X_i) - \mu A(x_0)) \right] \quad (16d)$$

$$S_j = \frac{\partial \hat{R}^\mu}{\partial \dot{x}}(X_j, \dot{X}_j) \quad (16e)$$

$$\dot{S}_j = \frac{\partial \hat{R}^\mu}{\partial x}(X_j, \dot{X}_j) - i_{\dot{X}_j} \beta_\mu(X_j), \quad (16f)$$

This system of equations is called the *reduced symplectic partitioned Runge–Kutta (RSPRK)* algorithm. A detailed derivation can be found in section 2.5 of [21]. As we obtained this system by dropping the symplectic partitioned Runge–Kutta algorithm from $\mathbf{J}^{-1}(\mu)$ to T^*S , it follows that this algorithm preserves the reduced symplectic form $\Omega_\mu = \Omega_S - \pi_{T^*S, S}^* \beta_\mu$ on T^*S .

Since the SPRK algorithm is an integration algorithm for the Hamiltonian vector field X_H on T^*Q , the RSPRK algorithm is an integration algorithm for the reduced Hamiltonian vector field X_{H_μ} on T^*S . The relationship between the cotangent bundle reduction and the reduction of the SPRK algorithm can be represented by the following commutative diagram:

$$\begin{array}{ccc} (J^{-1}(\mu), X_H) & \dashrightarrow & (J^{-1}(\mu), \text{SPRK}) \\ \pi_\mu \downarrow & & \downarrow \pi_\mu \\ (T^*S, X_{H_\mu}) & \dashrightarrow & (T^*S, \text{RSPRK}) \end{array}$$

The dashed arrows here denote the corresponding discretization, as in (15). The SPRK algorithm can be obtained by pushing forward the DEL equations by the discrete Legendre transform. See, for example, [37]. By lemma 2.4, this implies that the RSPRK algorithm can be obtained by pushing forward the DR equations by the reduced discrete Legendre transform $\hat{\mathbb{F}} = D_2 \hat{L}_d - \hat{\mathcal{A}}_2$. These relationships are shown in the following commutative diagram:

$$\begin{array}{ccc} (\mathbf{J}_d^{-1}(\mu), \text{DEL}) & \xrightarrow{\mathbb{F}L_d} & (J^{-1}(\mu), \text{SPRK}) \\ \pi_{\mu, d} \downarrow & & \downarrow \pi_\mu \\ (S \times S, \text{DR}) & \xrightarrow{\hat{\mathbb{F}}} & (T^*S, \text{RSPRK}) \end{array}$$

4. Putting everything together

The relationship between Routh reduction and cotangent bundle reduction can be represented by the following commutative diagram:

$$\begin{array}{ccc} (J_L^{-1}(\mu), EL) & \xrightarrow{\mathbb{F}L} & (J^{-1}(\mu), X_H) \\ \pi_{\mu, L} \downarrow & & \downarrow \pi_{\mu} \\ (TS, R) & \xrightarrow{\mathbb{F}\hat{R}^{\mu}} & (T^*S, X_{H_{\mu}}) \end{array}$$

We saw in section 2.6 that if L_d approximates the Jacobi solution of the Hamilton–Jacobi equation, discrete and continuous Routh reduction is related by the following diagram:

$$\begin{array}{ccc} (J_L^{-1}(\mu), EL) & \dashrightarrow & (J_d^{-1}(\mu), DEL) \\ \pi_{\mu, L} \downarrow & & \downarrow \pi_{\mu, d} \\ (TS, R) & \dashrightarrow & (S \times S, DR) \end{array}$$

The dashed arrows mean that the DEL equations are an integration algorithm for the EL equations, and that the DR equations are an integration algorithm for the Routh equations.

Pushing forward the DEL equation using the discrete Legendre transform $\mathbb{F}L_d$ yields the SPRK algorithm on T^*Q , which is an integration algorithm for X_H . This is depicted by

$$\begin{array}{ccc} (J_L^{-1}(\mu), EL) & \dashrightarrow & (J_d^{-1}(\mu), DEL) \\ \mathbb{F}L \downarrow & & \downarrow \mathbb{F}L_d \\ (J^{-1}(\mu), X_H) & \dashrightarrow & (J^{-1}(\mu), SPRK) \end{array}$$

The SPRK algorithm on $J^{-1}(\mu) \subset T^*Q$ induces the RSPRK algorithm on $J^{-1}(\mu)/G \approx T^*S$. As we saw in section 3, this reduction process is related to cotangent bundle reduction and to discrete Routh reduction as shown in the following diagram:

$$\begin{array}{ccccc} (J^{-1}(\mu), X_H) & \dashrightarrow & (J^{-1}(\mu), SPRK) & \xleftarrow{\mathbb{F}L_d} & (J_d^{-1}(\mu), DEL) \\ \pi_{\mu} \downarrow & & \downarrow \pi_{\mu} & & \downarrow \pi_{\mu, d} \\ (T^*S, X_{H_{\mu}}) & \dashrightarrow & (T^*S, RSPRK) & \xleftarrow{\mathbb{F}} & (S \times S, DR) \end{array}$$

Putting all the above commutative diagrams together into one diagram, we obtain figure 1.

5. Example: J_2 satellite dynamics

5.1. Configuration space and Lagrangian

An illustrative and important example of a system with an Abelian symmetry group is that of a single satellite in orbit about an oblate Earth. The general aspects and background for this problem are discussed in [40], and some interesting aspects of the geometry underlying it are discussed in [7].

The configuration manifold Q is \mathbb{R}^3 , and the Lagrangian is

$$L(q, \dot{q}) = \frac{1}{2} M_s \|\dot{q}\|^2 - M_s V(q),$$

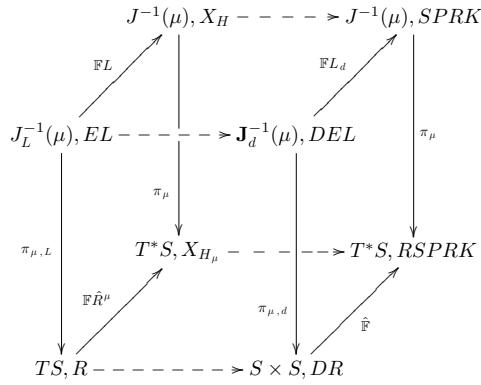


Figure 1. Complete commutative cube. The dashed arrows represent discretization from the continuous systems on the left face to the discrete systems on the right face. The vertical arrows represent reduction from the full systems on the top face to the reduced systems on the bottom face. The front and back faces represent Lagrangian and Hamiltonian viewpoints, respectively.

where M_s is the mass of the satellite and $V : \mathbb{R}^3 \rightarrow \mathbb{R}$ is the gravitational potential due to the Earth truncated at the first term in the expansion in the ellipticity

$$V(q) = \frac{GM_e}{\|q\|} + \frac{GM_e R_e^2 J_2}{\|q\|^3} \left(\frac{3}{2} \frac{(q^3)^2}{\|q\|^2} - \frac{1}{2} \right).$$

Here, G is the gravitational constant, M_e is the mass of the Earth, R_e is the radius of the Earth, J_2 is a small non-dimensional parameter describing the degree of ellipticity and q^3 is the third component of q . In non-dimensional coordinates,

$$L(q, \dot{q}) = \frac{1}{2} \|\dot{q}\|^2 - \left[\frac{1}{\|q\|} + \frac{J_2}{\|q\|^3} \left(\frac{3}{2} \frac{(q^3)^2}{\|q\|^2} - \frac{1}{2} \right) \right]. \quad (17)$$

This corresponds to choosing space and time coordinates in which the radius of the Earth is 1 and the period of orbit at zero altitude is 2π when $J_2 = 0$ (spherical Earth).

5.2. Symmetry action

The symmetry of interest to us is that of rotation about the vertical (q^3) axis, so the symmetry group is the unit circle S^1 . Using cylindrical coordinates $q = (r, \theta, z)$ for the configuration, the symmetry action is $\phi : (r, \theta, z) \mapsto (r, \theta + \phi, z)$. Since $\|q\|$, $\|\dot{q}\|$, and $q^3 = z$ are all invariant under this transformation, so too is the Lagrangian.

This action is clearly not free on all of $Q = \mathbb{R}^3$, as the z -axis is invariant for all group elements. This is not a serious obstacle as the lifted action is free on $T(Q \setminus (0, 0, 0))$ and this is enough to permit the application of the intrinsic Routh reduction theory. Alternatively, one can simply take $Q = \mathbb{R}^3 \setminus \{(0, 0, z) \mid z \in \mathbb{R}\}$ and then the theory literally applies.

The shape space $S = Q/G$ is thus the half-plane $S = \mathbb{R}^+ \times \mathbb{R}$ and we will take coordinates (r, z) on S . In doing so, we are implicitly defining a global diffeomorphism $S \times G \rightarrow Q$ given by $((r, z), \theta) \mapsto (r, \theta, z)$.

The Lie algebra \mathfrak{g} for $G = S^1$ is the real line $\mathfrak{g} = \mathbb{R}$, and we will identify the dual with the real line itself $\mathfrak{g}^* \cong \mathbb{R}$. For a Lie algebra element $\xi \in \mathfrak{g}$, the corresponding infinitesimal generator is given by $\xi_Q : (r, \theta, z) \mapsto ((r, \theta, z), (0, \xi, 0))$. The Lagrange momentum map $\mathbf{J}_L : TQ \rightarrow \mathfrak{g}^*$ is given by $\mathbf{J}_L(v_q) \cdot \xi = \langle \mathbb{F}L(v_q), \xi_Q(q) \rangle$, which in our case is a scalar quantity, the vertical component of the standard angular momentum $J_L((r, \theta, z), (\dot{r}, \dot{\theta}, \dot{z})) = r^2 \dot{\theta}$.

Consider the Euclidean metric on \mathbb{R}^3 , which corresponds to the kinetic energy norm in the Lagrangian. From this metric we define the mechanical connection $\mathfrak{A} : TQ \rightarrow \mathfrak{g}$ given by $\mathfrak{A}((r, \theta, z), (\dot{r}, \dot{\theta}, \dot{z})) = \dot{\theta}$. The 1-form \mathfrak{A}_μ on Q is thus given by $\mathfrak{A}_\mu = \mu \mathbf{d}\theta$. The exterior derivative of this expression gives $\mathbf{d}\mathfrak{A}_\mu = \mu \mathbf{d}^2\theta = 0$, and so the reduced 2-form is $\beta_\mu = 0$.

5.3. Equations of motion

The Euler–Lagrange equations for the Lagrangian (17) give the equations of motion,

$$\ddot{q} = -\nabla_q \left[\frac{1}{\|q\|} + \frac{J_2}{\|q\|^3} \left(\frac{3}{2} \frac{(q^3)^2}{\|q\|^2} - \frac{1}{2} \right) \right].$$

To calculate the reduced equations, we begin by calculating the Routhian

$$\begin{aligned} R^\mu(r, \theta, z, \dot{r}, \dot{\theta}, \dot{z}) &= L(r, \theta, z, \dot{r}, \dot{\theta}, \dot{z}) - \mathfrak{A}_\mu(r, \theta, z) \cdot (\dot{r}, \dot{\theta}, \dot{z}) \\ &= \frac{1}{2} \|(\dot{r}, \dot{\theta}, \dot{z})\|^2 - \left[\frac{1}{r} + \frac{J_2}{r^3} \left(\frac{3}{2} \frac{z^2}{r^2} - \frac{1}{2} \right) \right] - \mu \dot{\theta}. \end{aligned}$$

We choose a fixed value μ of the momentum and restrict ourselves to the space $J_L^{-1}(\mu)$, on which $\dot{\theta} = \mu$. The reduced Routhian $\hat{R}^\mu : TS \rightarrow \mathbb{R}$ is the restricted Routhian dropped to the tangent bundle of the shape space. In coordinates this is

$$\hat{R}^\mu(r, z, \dot{r}, \dot{z}) = \frac{1}{2} \|(\dot{r}, \dot{z})\|^2 - \left[\frac{1}{r} + \frac{J_2}{r^3} \left(\frac{3}{2} \frac{z^2}{r^2} - \frac{1}{2} \right) \right] - \frac{1}{2} \mu^2.$$

Recalling that $\beta_\mu = 0$, the Routh equations can be evaluated to give

$$(\ddot{r}, \ddot{z}) = -\nabla_{(r,z)} \left[\frac{1}{r} + \frac{J_2}{r^3} \left(\frac{3}{2} \frac{z^2}{r^2} - \frac{1}{2} \right) \right],$$

which describes the motion on the shape space.

To recover the unreduced Euler–Lagrange equations from the Routh equations one uses the procedure of reconstruction. This is covered in detail in [25, 26, 30].

5.4. Discrete Lagrangian system

We discretize this system with a high-order discrete Lagrangian. Recall that the pushforward discrete Lagrange map associated with this discrete Lagrangian is a symplectic partitioned Runge–Kutta method.

Given a point $(q_0, q_1) \in Q \times Q$, we will take (q_0, p_0) and (q_1, p_1) to be the associated discrete Legendre transforms. As the discrete momentum map is the pullback of the canonical momentum map, we have that $J_{L_d}(q_0, q_1) = (p_\theta)_0 = (p_\theta)_1$. Take a fixed momentum map value μ and restrict L_d to the set $J_{L_d}^{-1}(\mu)$. Dropping this to $S \times S$ now gives the reduced discrete Lagrangian $\hat{L}_d : S \times S \rightarrow \mathbb{R}$. More explicitly, L_d depends on (r_k, θ_k, z_k) and $(r_{k+1}, \theta_{k+1}, z_{k+1})$, but group invariance implies that the group variables only enter in the combination $(\theta_{k+1} - \theta_k)$. The condition $\mathbf{J}_d(q_k, q_{k+1}) = \mu$ can be inverted to eliminate the dependence of L_d on $(\theta_{k+1} - \theta_k)$, and L_d can be expressed in terms of $\mu, (r_k, z_k)$ and (r_{k+1}, z_{k+1}) , which yields \hat{L}_d .

As discussed in section 3, the fact that we have taken coordinates in which the group action is addition in θ means that the pushforward discrete Lagrange map associated with the reduced discrete Lagrangian is the reduced method given by (16a)–(16f). In fact, as the mechanical connection has $A(r, z) = 0$ and $\beta_\mu = 0$, the pushforward discrete Lagrange map is exactly a partitioned Runge–Kutta method with Hamiltonian equal to the reduced Routhian. These are generically related by a momentum shift, rather than being equal.

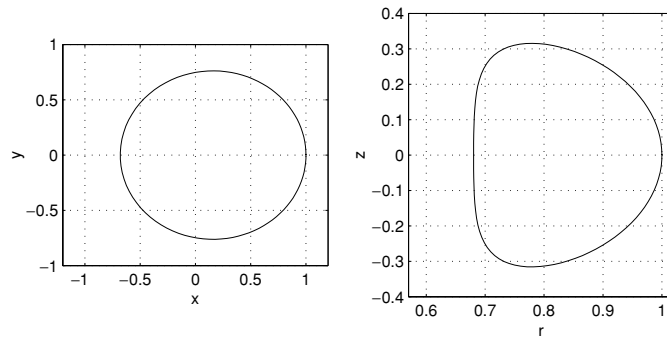


Figure 2. Unreduced (left) and reduced (right) views of an inclined elliptic trajectory for the continuous time system with $J_2 = 0$ (spherical Earth).

Given a trajectory of the reduced discrete system, we can reconstruct the unreduced discrete trajectory by solving for θ . Correspondingly, a trajectory of the unreduced discrete system can be projected onto shape space to give a trajectory of the reduced discrete system.

5.5. Example trajectories

We compute the unreduced trajectories using the fourth-order SPRK algorithm, and the reduced trajectories using the corresponding fourth-order RSPRK algorithm.

5.6. Solutions of the spherical Earth system

Consider initially the system with $J_2 = 0$. This corresponds to the case of a spherical Earth, and so the equations reduce to the standard Kepler problem. A slightly inclined circular trajectory is shown in figure 2, in both the unreduced and reduced pictures. Note that the graph of the reduced trajectory is a quadratic, as $\|q\| = \sqrt{r^2 + z^2}$ is a constant.

We will now investigate the effect of two different perturbations to the system, one due to taking non-zero J_2 and the other due to the numerical discretization.

5.7. The J_2 effect

Taking $J_2 = 0.05$ (which is close the actual value for the Earth), the system becomes near-integrable and experiences breakup of the KAM tori. This can be seen in figure 3, where the same initial condition is used as in figure 2.

Due to the fact that the reduced trajectory is no longer a simple curve, there is a geometric-phase-like effect which causes precession of the orbit. This precession can be seen in the thickening of the unreduced trajectory.

5.8. Solutions of the discrete system for a spherical Earth

We now consider the discrete system with $J_2 = 0$, for the second-order Gauss–Legendre discrete Lagrangian with the step size of $h = 0.3$. The trajectory with the same initial condition as above is given in figure 4.

As can be seen from the reduced trajectory, the discretization has caused a similar breakup of the periodic orbit as was produced by the non-zero J_2 . This induces precession of the orbit in the unreduced trajectory, in a way which is difficult to distinguish from the perturbation

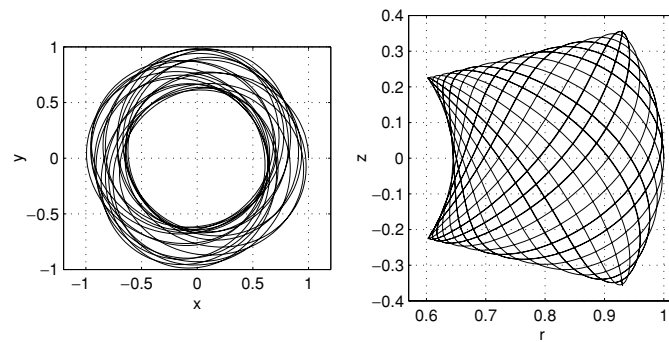


Figure 3. Unreduced (left) and reduced (right) views of an inclined elliptic trajectory for the continuous time system with $J_2 = 0.05$. Observe that the non-spherical terms introduce precession of the near-elliptic orbit in the symmetry direction.

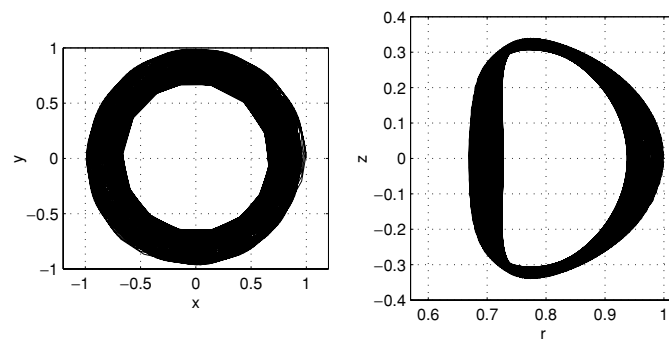


Figure 4. Unreduced (left) and reduced (right) views of an inclined trajectory of the discrete system with step size $h = 0.3$ and $J_2 = 0$. The initial condition is the same as that used in figure 2. The numerically introduced precession means that the unreduced picture looks similar to that of figure 3 with non-zero J_2 , whereas only by considering the reduced picture we can see the correct resemblance to the $J_2 = 0$ case of figure 2.

above due to non-zero J_2 if only the unreduced picture is considered. If the reduced pictures are consulted, however, then it is immediately clear that the system is much closer to the continuous time system with $J_2 = 0$ than to the system with non-zero J_2 .

5.9. Solutions of the discrete system with J_2 effect

Finally, we consider the discrete system with non-zero $J_2 = 0.05$. The resulting trajectory is shown in figure 5, and it is clearly not easy to determine from the unreduced picture whether the precession is due to the J_2 perturbation, the discretization, or some combination of the two.

Taking the reduced trajectories, however, immediately shows that this discrete time system is structurally much closer to the non-zero J_2 system than to the original $J_2 = 0$ system. This confusion arises because both the J_2 term and the discretization introduce perturbations which act in the symmetry direction.

While this system is sufficiently simple that one can run simulations with such small time steps that the discretization artefacts become negligible, this is certainly not generally possible. This example demonstrates how knowledge of the geometry of the system can be important in

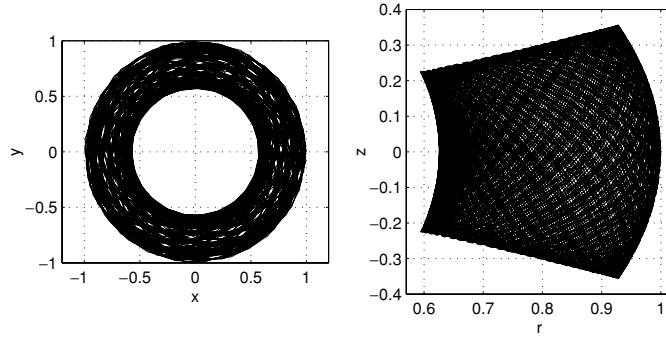


Figure 5. Unreduced (left) and reduced (right) views of an inclined trajectory of the discrete system with step size $h = 0.3$ and $J_2 = 0.05$. The initial condition is the same as that used in figure 3. The unreduced picture is similar to those of figures 3 and 4. By considering the reduced picture the correct resemblance to 3.

understanding the discretization process, and how this can give insight into the behaviour of numerical simulations. In particular, understanding how the discretization interacts with the symmetry action is extremely important.

5.10. Coordinate systems

In this example we have chosen cylindrical coordinates, thus making the group action addition in θ . One can always do this, as an Abelian Lie group is isomorphic to a product of copies of \mathbb{R} and S^1 , but it may sometimes be preferable to work in coordinates in which the group action is not addition. For example, Cartesian coordinates in the present example. Reasons for choosing a different coordinate system might include ease of computation, or simplicity of the expressions.

If we adopt a coordinate system wherein the group action is not expressed in terms of addition, the RSPRK method is not applicable, but we can still apply the Discrete Routh equations to obtain an integration scheme on $S \times S$. The push forward of this under $\hat{\mathbb{F}}$ yields an integration scheme on T^*S . The trajectories on the shape space that we obtain in this manner could be different from those we would get with the RSPRK method. However, in both cases we would have conservation of symplectic structure, momentum, and the order of accuracy would be the same. One could choose whichever approach is cheaper and easier.

6. Example: double spherical pendulum

6.1. Configuration space and Lagrangian

We consider the example of the double spherical pendulum which has a non-trivial magnetic term and constraints. The configuration manifold Q is $S^2 \times S^2$, and the embedding linear space V is $\mathbb{R}^3 \times \mathbb{R}^3$. The position vectors of each pendulum with respect to the pivot point are denoted by \mathbf{q}_1 and \mathbf{q}_2 . These vectors are constrained to have lengths l_1 and l_2 respectively, and the pendula masses are denoted by m_1 and m_2 . The Lagrangian is

$$L(\mathbf{q}_1, \mathbf{q}_2, \dot{\mathbf{q}}_1, \dot{\mathbf{q}}_2) = \frac{1}{2}m_1\|\dot{\mathbf{q}}_1\|^2 + \frac{1}{2}m_2\|\dot{\mathbf{q}}_1 + \dot{\mathbf{q}}_2\|^2 - m_1g\mathbf{q}_1 \cdot \mathbf{k} - m_2g(\mathbf{q}_1 + \mathbf{q}_2) \cdot \mathbf{k},$$

where g is the gravitational constant, and \mathbf{k} is the unit vector in the z direction. The constraint function $c : V \rightarrow \mathbb{R}^2$ is given by $c(\mathbf{q}_1, \mathbf{q}_2) = (\|\mathbf{q}_1\| - l_1, \|\mathbf{q}_2\| - l_2)$. Using cylindrical

coordinates $\mathbf{q}_i = (r_i, \theta_i, z_i)$, L becomes

$$\begin{aligned} L(q, \dot{q}) = & \frac{1}{2}m_1(\dot{r}_1^2 + r_1^2\dot{\theta}_1^2 + \dot{z}_1^2) + \frac{1}{2}m_2\{\dot{r}_1^2 + r_1^2\dot{\theta}_1^2 + \dot{r}_2^2 + r_2^2\dot{\theta}_2^2 \\ & + 2(\dot{r}_1\dot{r}_2 + r_1r_2\dot{\theta}_1\dot{\theta}_2)\cos\varphi + 2(r_1\dot{r}_2\dot{\theta}_1 - r_2\dot{r}_1\dot{\theta}_2)\sin\varphi + (\dot{z}_1 + \dot{z}_2)^2\} \\ & - m_1gz_1 - m_2g(z_1 + z_2), \end{aligned}$$

where $\varphi = \theta_2 - \theta_1$. Furthermore, we can automatically satisfy the constraints by performing the substitutions, $z_i = (l_i^2 - r_i^2)^{1/2}$, and $\dot{z}_i = -r_i\dot{r}_i(\sqrt{l_i^2 - r_i^2})^{-1/2}$.

6.2. Symmetry action

The symmetry of interest to us is the simultaneous rotation of the two pendula about the vertical (z) axis, so the symmetry group is the unit circle S^1 . Using cylindrical coordinates $\mathbf{q}_i = (r_i, \theta_i, z_i)$ for the configuration, the symmetry action is $\phi : (r_i, \theta_i, z_i) \mapsto (r_i, \theta_i + \phi, z_i)$. Since $\|\mathbf{q}_i\|$, $\|\dot{\mathbf{q}}_i\|$, $\|\dot{\mathbf{q}}_1 + \dot{\mathbf{q}}_2\|$ and $\mathbf{q}_i \cdot \mathbf{k}$ are all invariant under this transformation, so too is the Lagrangian.

This action is clearly not free on all of $V = \mathbb{R}^3 \times \mathbb{R}^3$, as the z -axis is invariant for all group elements. However, this does not pose a problem computationally, as long as the trajectories do not pass through the downward hanging configuration, corresponding to $r_1 = r_2 = 0$. To treat the downward hanging configuration properly, we would need to develop a discrete Lagrangian analogue of the continuous theory of singular reduction described in [38].

We will now show the geometric structures involved in implementing the RSPRK algorithm. See section 2.10.2 of [21] for a detailed derivation. The Lie algebra \mathfrak{g} for $G = S^1$ is the real line $\mathfrak{g} = \mathbb{R}$, and we will identify the dual with the real line itself $\mathfrak{g}^* \cong \mathbb{R}$. For a Lie algebra element $\xi \in \mathfrak{g}$, the corresponding infinitesimal generator is given by $\xi_Q : (r_1, \theta_1, z_1, r_2, \theta_2, z_2) \mapsto ((r_1, \theta_1, z_1, r_2, \theta_2, z_2), (0, \xi, 0, 0, \xi, 0))$. As such, the momentum map is given by

$$\begin{aligned} J_L((r_1, \theta_1, z_1, r_2, \theta_2, z_2), (\dot{r}_1, \dot{\theta}_1, \dot{z}_1, \dot{r}_2, \dot{\theta}_2, \dot{z}_2)) \\ = (m_1 + m_2)r_1^2\dot{\theta}_1 + m_2r_2^2\dot{\theta}_2 + m_2r_1r_2(\dot{\theta}_1 + \dot{\theta}_2)\cos\varphi + (r_1\dot{r}_2 - r_2\dot{r}_1)\sin\varphi, \end{aligned}$$

which is simply the vertical component of the standard angular momentum.

The locked inertia tensor is given by [25]

$$\mathbb{I}(\mathbf{q}_1, \mathbf{q}_2) = m_1\|\mathbf{q}_1^\perp\|^2 + m_2\|(\mathbf{q}_1 + \mathbf{q}_2)^\perp\|^2 = m_1r_1^2 + m_2(r_1^2 + r_2^2 + 2r_1r_2\cos\varphi).$$

The mechanical connection as a 1-form is given by

$$\begin{aligned} \alpha(\mathbf{q}_1, \mathbf{q}_2) = [m_1r_1^2 + m_2(r_1^2 + r_2^2 + 2r_1r_2\cos\varphi)]^{-1}[(m_1 + m_2)r_1^2\mathbf{d}\theta_1 + m_2r_2^2\mathbf{d}\theta_2 \\ + m_2r_1r_2(\mathbf{d}\theta_1 + \mathbf{d}\theta_2)\cos\varphi + (r_1\mathbf{d}r_2 - r_2\mathbf{d}r_1)\sin\varphi]. \end{aligned}$$

The μ -component of the mechanical connection is given by

$$\begin{aligned} \alpha_\mu(\mathbf{q}_1, \mathbf{q}_2) = \mu[m_1r_1^2 + m_2(r_1^2 + r_2^2 + 2r_1r_2\cos\varphi)]^{-1} \\ \times \{[(m_1 + m_2)r_1^2 + m_2r_1r_2\cos\varphi]\mathbf{d}\theta_1 + [m_2r_2^2 + m_2r_1r_2\cos\varphi]\mathbf{d}\theta_2\}. \end{aligned}$$

Taking the exterior derivative of this 1-form yields a non-trivial magnetic term on the reduced space, which drops to the quotient space to yield

$$\begin{aligned} \beta_\mu = \mu m_2[2(m_1 + m_2)r_1r_2 + (m_1r_1^2 + m_2(r_1^2 + r_2^2))\cos\varphi] \\ \times [m_1r_1^2 + m_2(r_1^2 + r_2^2 + 2r_1r_2\cos\varphi)]^{-2}\mathbf{d}\varphi \wedge (r_2\mathbf{d}r_1 - r_1\mathbf{d}r_2). \end{aligned}$$

The local representation of the connection is given by

$$A(r_1, r_2, \varphi) = m_2(m_1 r_1^2 + m_2(r_1^2 + r_2^2 + 2r_1 r_2 \cos \varphi))^{-1} \begin{bmatrix} -r_2 \sin \varphi \\ r_1 \sin \varphi \\ r_2^2 + r_1 r_2 \cos \varphi \end{bmatrix}^T,$$

and the amended potential V_μ has the form

$$V_\mu(q) = -m_1 g(l_1^2 - r_1^2)^{1/2} - m_2 g[(l_1^2 - r_1^2)^{1/2} + (l_2^2 - r_2^2)^{1/2}] \\ + 2^{-1} \mu^2 [m_1 r_1^2 + m_2(r_1^2 + r_2^2 + 2r_1 r_2 \cos \varphi)]^{-1}.$$

The Routhian on the momentum level set is given by $R^\mu = \frac{1}{2} \|\text{hor}(q, v)\|^2 - V_\mu$. Recall that $\text{hor}(v_q) = v_q - \xi_Q(v_q)$, where $\xi = \alpha(v_q)$ and $\xi_Q(v_q) = (0, \xi, 0, 0, \xi, 0)$. Then we obtain

$$\text{hor}(v_q) = v_q - (0, \alpha(v_q), 0, 0, \alpha(v_q), 0) = (\dot{r}_1, \dot{\theta}_1 - \alpha(v_q), \dot{z}_1, \dot{r}_2, \dot{\theta}_2 - \alpha(v_q), \dot{z}_2).$$

The kinetic energy metric has the form

$$\begin{bmatrix} m_1 + m_2 & 0 & 0 & m_2 \cos \varphi & -m_2 r_2 \sin \varphi & 0 \\ 0 & (m_1 + m_2) r_1^2 & 0 & m_2 r_1 \sin \varphi & m_2 r_1 r_2 \cos \varphi & 0 \\ 0 & 0 & m_1 + m_2 & 0 & 0 & 0 \\ m_2 \cos \varphi & m_2 r_1 \sin \varphi & 0 & m_2 & 0 & 0 \\ -m_2 r_2 \sin \varphi & m_2 r_1 r_2 \cos \varphi & 0 & 0 & m_2 r_2^2 & 0 \\ 0 & 0 & 0 & 0 & 0 & m_2 \end{bmatrix}.$$

This together with the expression for $\text{hor}(v_q)$ allows us to compute $\frac{1}{2} \|\text{hor}(q, v)\|^2$, and when combined with the formula for the amended potential V_μ gives the Routhian R^μ . Note that all our expressions are in terms of the reduced variables on TS , so they trivially drop to yield \hat{R}^μ . These expressions can then be directly substituted into the RSPRK algorithm in equations (16a)–(16f) to obtain the example trajectories presented in the following subsection.

6.3. Example trajectories

We have computed the reduced trajectory of the double spherical pendulum using the fourth-order RSPRK algorithm on the Routh equations. We first consider the evolution of r_1, r_2 and φ using the RSPRK algorithm on the Routh equations, as well as the projection of the relative position of m_2 with respect to m_1 onto the xy plane as seen in figure 6.

Figure 7 illustrates that the energy behaviour of the trajectory is very good, as is typical of variational integrators, and does not exhibit a spurious drift. In contrast, the non-symplectic fourth-order Runge–Kutta applied to the unreduced dynamics has a systematic drift in the energy, even when using time steps that are smaller by a factor of 4.

7. Computational considerations

7.1. Reduced versus unreduced simulations

The reduced dynamics can either be computed directly, by using the discrete Routh or RSPRK equations, or by computing in the unreduced space, and projecting onto the shape space. We discuss the relative merits of these approaches.

Given a configuration space and symmetry group of dimensions n and m , respectively, we can either use a simpler algorithm in $2n$ dimensions, or a more geometrically involved algorithm in $2(n - m)$ dimensions. The reduced algorithm involves curvature terms that need to be symbolically precomputed, but these do not affect the sparsity of the system of equations,

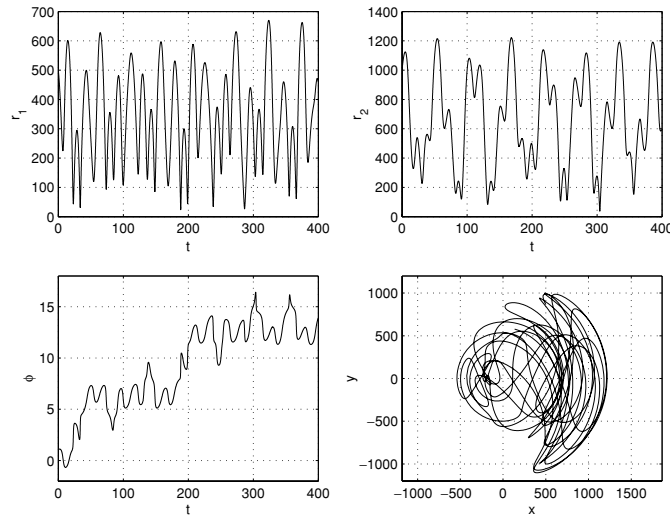


Figure 6. Time evolution of r_1 , r_2 , ϕ , and the trajectory of m_2 relative to m_1 using RSPRK.

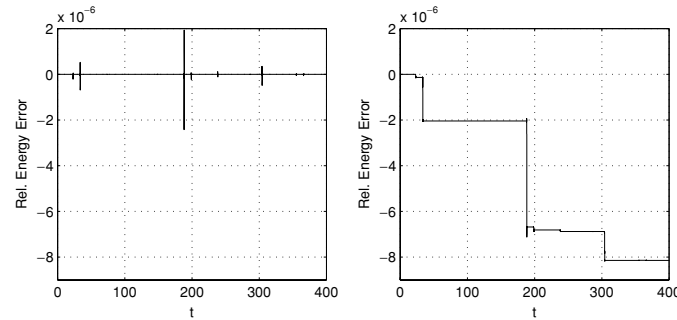


Figure 7. Relative energy drift $(E - E_0)/E_0$ using RSPRK (left) compared to the relative energy drift in a non-symplectic RK (right).

and the lower dimension results in computational saving that are particularly evident in the repeated, or long-time simulation of problems with a shape space of high codimension.

An example which is of current engineering interest is the dynamics of connected networks of systems with their own internal symmetries, such as coordinated clusters of satellites modelled as rigid bodies with internal rotors. If the systems to be connected are all identical, the geometric quantities that need to be computed, such as the mechanical connection, have a particularly simple repeated form, and the small additional upfront effort in implementing the reduced algorithm can result in substantial computational savings due to the lower dimension of the reduced system.

7.2. Intrinsic versus non-intrinsic methods

Non-intrinsic numerical schemes, such as SPRK applied to the classical Routh equations, can have undesirable numerical properties due to the need for coordinate-dependent local trivializations and the presence of coordinate singularities in these local trivializations, as is the case when using Euler angles for rigid body dynamics.

In the case of non-canonical symplectic forms, frequent computationally expensive coordinate changes are necessary when using standard non-intrinsic schemes, as documented in [39, 44], due to the need to repeatedly apply Darboux' theorem to put the symplectic structure into canonical form. In contrast, intrinsic methods do not depend on a particular choice of coordinate system, and allow for the use of global charts through the use of containing vector spaces with constraints enforced using Lagrange multipliers.

Coordinate singularities can affect the quality of the simulation in subtle ways that may depend on the numerical scheme. In the simulation of the double spherical pendulum, we notice spikes in the energy corresponding to times when r_1 or r_2 are close to 0. These errors accumulate in the non-symplectic method, but remain well behaved in the symplectic method. Alternatively, sharp spikes can be avoided altogether by evolving the equations on a constraint surface in $\mathbb{R}^3 \times \mathbb{R}^3$, as opposed to choosing local coordinates that automatically satisfy the constraints. Here, the increased cost of working in the containing linear space with constraints is offset by not having to transform between charts of $S^2_{l_1} \times S^2_{l_2}$, which can be significant if the trajectories are chaotic. An extensive discussion of the issue of representations and parametrizations of rotation groups and its implications for computation can be found in [5].

Methods which exhibit local conservation properties on each chart may still exhibit a drift in the conserved quantity across coordinate changes. As discussed in [1], only methods in which the local representatives of the algorithm commute with coordinate changes exhibit geometric conservation properties that are robust. In particular, intrinsic methods retain their conservation properties through coordinate changes.

Another promising approach is to use the exponential map to update the numerical solution, which is the basis of Lie group integrators, see [14]. The Lie group approach can yield rigid body integrators that are embedded in the space of 3×3 matrices but automatically evolve on the rotation group, without the use of constraints or reprojection. In [16, 17], analogues of the explicit Newmark and midpoint Lie algorithm are presented that automatically stay on the rotation group, and exhibit good energy and momentum stability. A method based on generating functions and the exponential on Riemannian manifolds is introduced in [19]. In the variational setting, a Lie group variational integrator for rigid body dynamics is described in [18], and higher-order generalizations can be found in [21].

8. Conclusions and future work

This paper derives the discrete Routh equations on the discrete shape space $S \times S$, which are symplectic with respect to a non-canonical symplectic form, and retains the good energy behaviour typically associated with variational integrators. Furthermore, when the group action can be expressed as addition, we obtain the reduced symplectic partitioned Runge–Kutta algorithm on T^*S , which can be considered as a discrete analogue of cotangent bundle reduction. By providing an understanding of how the reduced and unreduced formulations are related at a discrete level, we enable the user to freely choose whichever formulation is most appropriate, and provides the most insight to the problem at hand.

Certainly one of the obvious things to do in the future is to extend discrete reduction to the case of non-Abelian symmetry groups following the non-Abelian version of Routh reduction given in [15, 30]. There are also several problems, including the averaged J_2 problem, in which one can carry out discrete reduction by stages and relate it to the semidirect product work of [4]. This is motivated by the fact that the semidirect product reduction theory of [13] is a special case of reduction by stages (at least without the momentum map constraint), as was shown in [6].

In further developing discrete reduction theory, the discrete theory of connections on principal bundles developed in [22] is particularly relevant, as it provides an intrinsic method of representing the reduced space $(Q \times Q)/G$ as $(S \times S) \oplus \tilde{G}$, where $\tilde{G} = Q \times_G G$ with G acting by conjugation on G . Indeed, such a construction can be viewed as a connection on a Lie groupoid, and it is natural to express discrete mechanics on $Q \times Q$ in the language of pair groupoids, as originally proposed in [43]. Generalizations of this approach to arbitrary Lie groupoids, as well as a discussion of the role of discrete connections in yielding a discrete analogue of the Lagrange–Poincaré equations, can be found in [24].

Another component that is needed in this work is a good discrete version of the calculus of differential forms. Note that in our work we found, being directed by mechanics, that the right discrete version of the magnetic 2-form is the difference of two connection 1-forms. It is expected that we could recover such a magnetic 2-form by considering the discrete exterior derivative of a discrete connection form in a finite discretization of spacetime, and taking the continuum limit in the spatial discretization. Developing a discrete analogue of Stokes' theorem would also provide insight into the issue of discrete geometric phases. Some preliminary work on a discrete theory of exterior calculus can be found in [8].

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