

SPECTROSCOPICALLY EVALUATED RATES AND ENERGIES  
FOR PROTON TRANSFER AND BJERRUM DEFECT  
MIGRATION IN CUBIC ICE

By

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## CHAPTER I

### LITERATURE REVIEW

#### Introduction

Ice is a peculiar solid. Its unique properties have fascinated scientists for decades but surprisingly its behavior and structure is still often ambiguous. The importance of ice to chemistry, physics, geophysics, atmospheric science, hydrology, cryobiology, molecular biology, and other fields has grown and with it an immense volume of published research. Any review of ice must be necessarily selective and the present work is no exception. The main focus will be the areas of ice research related to the spectroscopy of ice with special emphasis on proton transfer.

Ice can exist in several solid phases. The low pressure forms are amorphous or vitreous ice (ice Ia or Iv), cubic ice (ice Ic), and the naturally occurring hexagonal ice (ice Ih). Sometimes cubic and hexagonal ice mixtures or hexagonal ice alone is referred to as polycrystalline ice I. The presently known high pressure polymorphs are ice II-VII and are accessible at pressures in the kilobar range.

As early as 1917 (1) x-ray diffraction studies on ice were pursued. In 1922 Bragg (2) deduced that each oxygen atom is located at the center of gravity of its four neighboring oxygen atoms for hexagonal ice. Since then modern x-ray and electron diffraction data (3) have confirmed that hexagonal ice is composed of tetrahedrally

coordinated oxygen atoms stacked such that it gives a hexagonal crystal structure. Figure 1 gives the oxygen lattice structure of hexagonal ice with its dimensions. Notice that the oxygen atoms can be grouped into a series of rough planes called basal planes that are normal to the crystal c-axis. By looking parallel to the c-axis the hexagonal holes of ice Ih are readily seen. This open structure gives hexagonal ice its lower than liquid water density.

Since each oxygen atom has two bonded hydrogen atoms, the oxygen lattice must somehow accommodate these hydrogens. One early suggestion by Barnes (4) was to put the hydrogens midway between each oxygen-oxygen link. This would preserve the symmetry and stoichiometry of the crystal. However this arrangement was ruled out since the water molecule would be ionized into  $H^+$  and  $O^{2-}$  ions, and water is a poor electrical conductor. Bernal and Fowler (5) suggests that the hydrogens were distributed between the oxygen-oxygen bonds but displaced such that a given hydrogen atom was about 1 angstrom from one oxygen atom and 1.76 angstroms from the other with the arrangement of hydrogen atoms in the complete lattice such that each oxygen atom has two hydrogens chemically bonded at 1 angstrom and two hydrogen bonds to two adjacent water molecules with the hydrogen at 1.76 angstroms. Figure 2 illustrates Bernal and Fowler's suggestion. A closer examination of this model reveals a water subunit with a HOH angle of approximately  $109^\circ$  compared with the gas phase value of approximately  $105^\circ$ . The OH bond length is 1 angstrom compared with .97 angstrom for the gas phase. Thus the distortion of a water molecule transcending the gas to ice transition is slight. Some researchers have even suggested that the  $H_2O$  subunit retains its gas phase angle in the hexagonal phase ice such that the

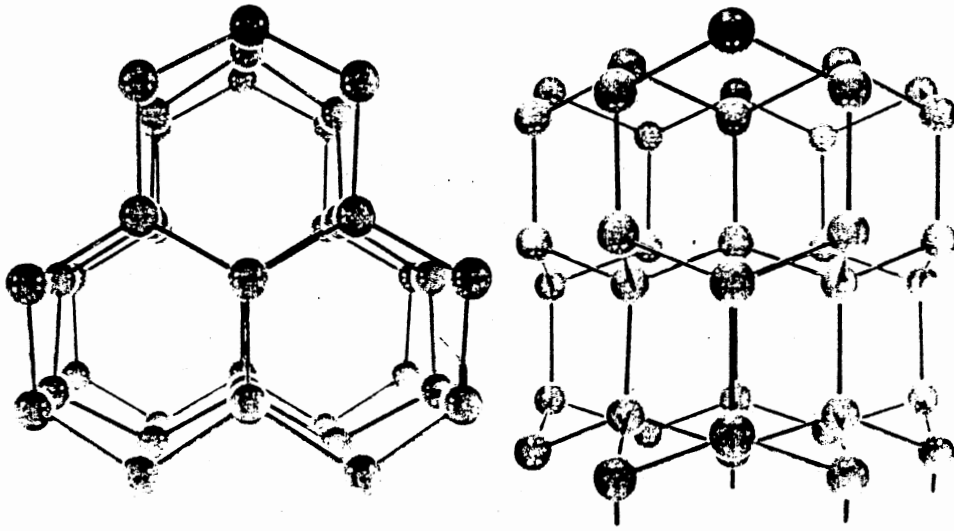


Figure 1. Oxygen Lattice Structure of Hexagonal Ice

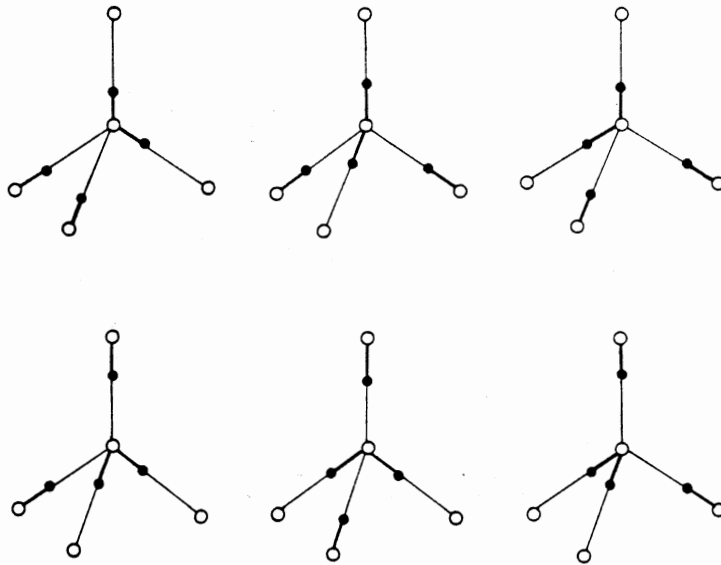


Figure 2. Possible Hydrogen Configuration in Ice Ih



hydrogens lie slightly off center of the O-O lines.

One interesting implication of the Bernal and Fowler model is that around any given oxygen atom there are six possible configurations of hydrogens as Figure 2 illustrates. Thus even at 0 K ice Ih may freeze into a number of different possible hydrogen configurations leaving ice with a nonzero thermodynamic probability and hence a nonzero residual entropy at 0 K. This led Pauling (6) to propose a statistical model for ice Ih with the following assumptions.

- (1) Each oxygen atom has two hydrogen atoms attached at distances of about 0.95 angstroms, thus forming a water molecule.
- (2) Each water molecule is oriented so that its two hydrogen atoms are directed approximately towards two of the four oxygen atoms which surround it tetrahedrally.
- (3) The orientations of neighboring water molecules are such that one hydrogen atom normally lies between each pair of oxygen atoms.
- (4) Under normal conditions ice Ih can exist in any one of a large number of configurations, each of which corresponds to a particular hydrogen atom distribution among the oxygen lattice (p. 2680).

These four assumptions are often called the Bernal and Fowler rules and ice crystals which follow them are called ideal crystals.

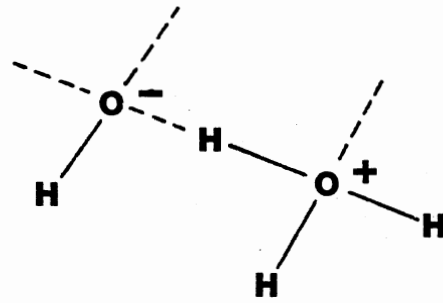
Bjerrum (7) in 1951 proposed ice crystal defects that could facilitate proton transfer. These defects named Bjerrum defects are created when the third Bernal and Fowler rule is violated. A D defect is formed when two hydrogen atoms occupy an O-O bond and an L defect is formed when an O-O bond is left vacant. In conjunction with ionization defects, these could allow positive and negative ionic charges to migrate through the ice crystal. If a hydrogen ion in an ideal hexagonal ice crystal shifts from one OH bonded position to the OH

bonded position of the adjacent oxygen, i.e. a movement of approximately .24 angstroms of a hydrogen to the adjacent hydrogen bonded oxygen, then two ion defects will be formed. One will be the  $\text{H}_3\text{O}^+$  ion about the oxygen the defect moved toward, the other a hydroxyl  $-\text{OH}$  ion left behind. Figure 3 illustrates these defects.

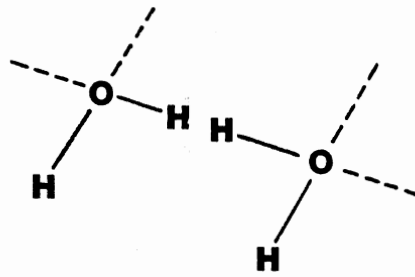
As mentioned earlier, there are presently two other known low pressure phases of ice, amorphous ice (ice Ia or Iv) and cubic ice (ice Ic). If water vapor is condensed on a substrate below 130 K ice Ia is formed. If the temperature is above 130 K and below approximately 170 K - 190 K then cubic ice is formed. Above 190 K hexagonal ice tends to form, though it should be noted that the exact temperatures are disputed.

Amorphous ice has a disordered structure associated with it since the energy of the substrate is not sufficient to allow a deposited water molecule to reorient itself to a more energetically favorable position. Because of the lack of structural order amorphous ice has been studied as a possible model for liquid water structure.

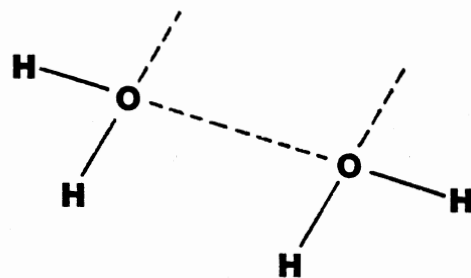
Cubic ice can be made in situ by direct vapor deposition, or by warming amorphous ice to above the cubic ice transition temperatures as noted by Honjo et al. (8). The oxygen lattice of cubic ice is almost identical to the hexagonal lattice except that every alternate basal plane is rotated  $60^\circ$  about the crystal c-axis relative to the adjacent nonrotated basal planes. This preserves the tetrahedral oxygen lattice but removes the hexagonal holes along the c-axis. Honjo and Shimaoko (9) investigated the possible hydrogen atom positions of cubic ice using electron diffraction and found that the cubic equivalent of Pauling's statistical model of ice Ih still holds.



ION PAIR DEFECT



D ROTATIONAL DEFECT



L ROTATIONAL DEFECT

Figure 3. Ice Crystal Lattice Defects

Further evidence of the cubic hydrogen atom positions comes from infrared spectroscopy where, to date, no definite difference has been detected between the cubic and hexagonal ice IR spectra. Likewise it appears that the Bjerrum and ionization defects of hexagonal ice also apply to cubic ice. The oxygen lattice structure of cubic ice is shown in Figure 4.

The  $\text{H}_2\text{O}$  molecule possesses a  $C_{2V}$  symmetry with three normal modes of vibration as shown in Figure 5. The  $\nu_3$  vibration is the asymmetric stretch belonging to the  $B_1$  irreducible representation and occurs in the  $3000 - 3500 \text{ cm}^{-1}$  region for  $\text{H}_2\text{O}$  and  $2200 - 2600 \text{ cm}^{-1}$  for  $\text{D}_2\text{O}$ . The  $\nu_1$  vibration is the symmetric stretching mode belonging to the  $A_1$  irreducible representation, and has a wavenumber range similar to  $\nu_3$ . The  $\nu_2$  bending vibration also belongs to the  $A_1$  irreducible representation and usually is found around  $1500 - 1800 \text{ cm}^{-1}$  for  $\text{H}_2\text{O}$  and  $1100 - 1300 \text{ cm}^{-1}$  for  $\text{D}_2\text{O}$ .

In ice and liquid water  $\nu_3$  and  $\nu_1$  are found in the OH stretching region as a very broad almost featureless band whose width has traditionally been attributed to hydrogen bonding. Certain lattice modes are found below  $\nu_2$  and their combinations with other modes are often interspersed among the fundamentals.

#### Pure Ice Spectra

The spectra of pure ice has been studied for many years. Bertie and Whalley have figured significantly in the early characterizations of the ice IR spectra.

In 1968 Bertie, Labbe', and Whalley (10) examined the absorptivity of ice I in the  $4000$  to  $30 \text{ cm}^{-1}$  range. True absorptivity profiles of

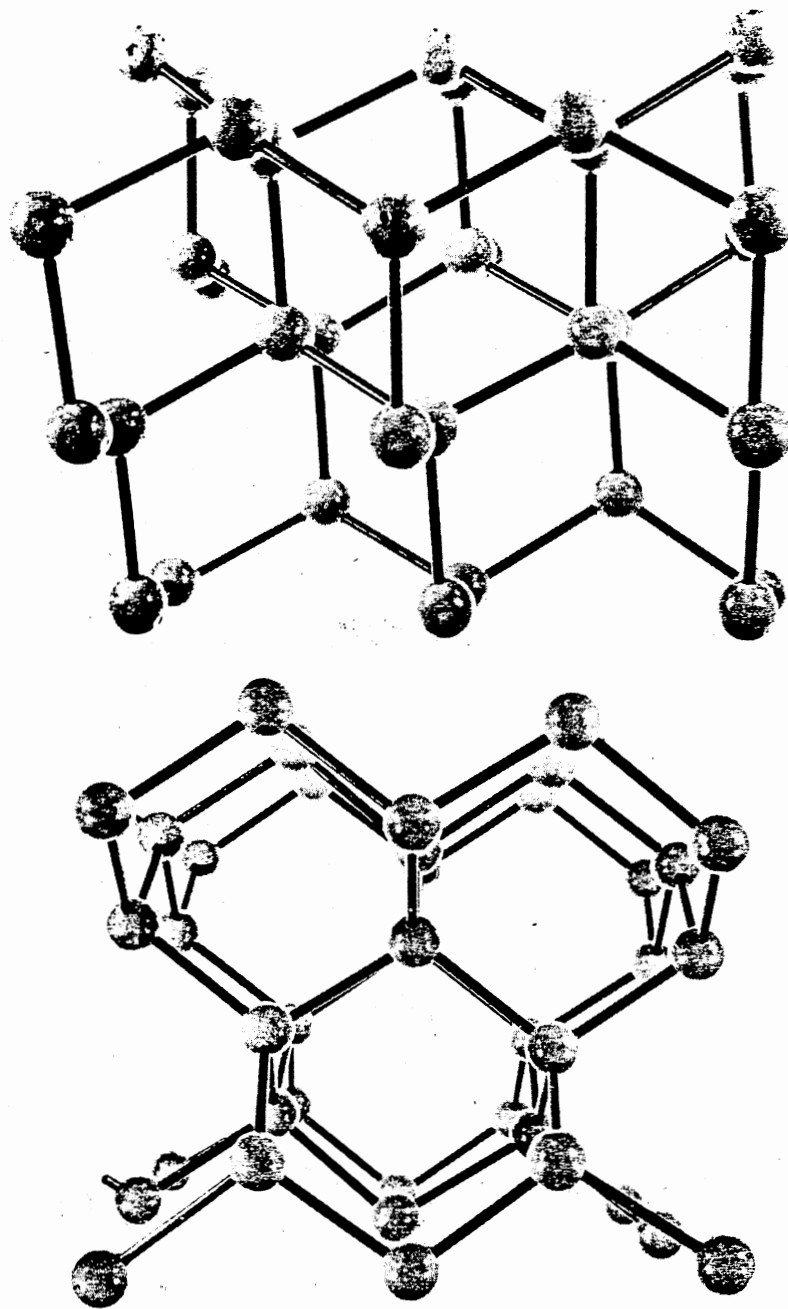
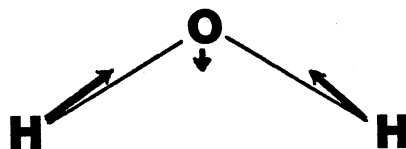
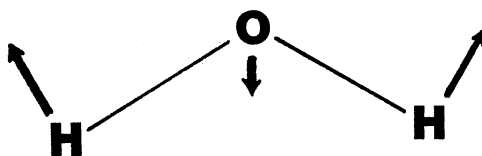


Figure 4. Oxygen Lattice Structure of Cubic Ice

$\nu_1$  Symmetric Stretching Mode



$\nu_2$  Bending Mode



$\nu_3$  Asymmetric Stretching Mode

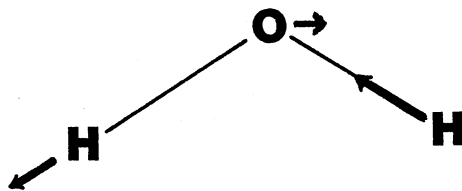


Figure 5. Free  $\text{H}_2\text{O}$  Normal Modes of Vibration

solid state ice can be difficult to obtain, especially near large absorptivities where the infrared index of refraction changes dramatically. Therefore a true absorptivity profile that accounted for reflection and scattering losses would facilitate interpretation of ice spectra.

The vitreous to cubic ice I phase transition was studied by Hardin and Harvey (11) by monitoring the bandshape changes of several bands over a 92 to 210 K temperature range. They reported the well known sharpening of the vitreous ice bands upon transformation to the cubic phase. One exception noted was the vitreous ice  $\nu_2$  bending mode whose half-height bandwidth was reported as  $350 \pm 10 \text{ cm}^{-1}$  in the amorphous phase but  $365 \pm 10 \text{ cm}^{-1}$  in the cubic phase. The transition for this band was between  $123 \pm 5 \text{ K}$  and  $137 \pm 5 \text{ K}$ . However the  $800 \text{ cm}^{-1}$   $\nu_r$  showed a  $-25 \text{ cm}^{-1}$  decrease in bandwidth at 92 K as expected. Ice Iv transformed over a  $120 \pm 5$  to  $135 \pm 5 \text{ K}$  temperature range. They concluded that vitreous (or ice Iv or Ia) ice is stable up to  $120 \pm 5 \text{ K}$  and undergoes an irreversible transformation to cubic ice between 120 and 140 K. Hardin and Harvey felt confident that their spectra represented cubic ice and did contain some small discernable differences from the spectra of hexagonal ice.

In 1973 Li and Devlin reported the glassy water Raman spectra from a trapped laser beam. Both polarized and depolarized spectra in the  $3570$  to  $2950 \text{ cm}^{-1}$  region were reported where  $\nu_1$  was located at  $3110 \text{ cm}^{-1}$  and  $2\nu_2$  at approximately  $3250 \text{ cm}^{-1}$ . In 1975 Narten, Venkatesh, and Rice (13) obtained x-ray diffraction data on amorphous  $\text{H}_2\text{O}$  and compared it with neutron diffraction data. They suggested a high density 10 K ice and a low density 77 K ice amorphous ice phase

to explain their data. Their model predicted a randomized ice I structure, with interstitial molecules, having a density of 1.1 grams/cm<sup>3</sup>. Their model was also consistent with the predictions of supercooled water; with no structural changes other than a contraction of the O-O separations. Thus Narten et al. suggested that the low temperature phase of amorphous ice may serve as a model for liquid water.

Whalley (14) studied the spectra of ice Ic and Ih and presented a detailed assignment based on the spectrum of ice VIII. Whalley noted that assignments can be difficult because the vibrations can not be treated as arising from a single molecule, but due to the strong intermolecular coupling that exists in ice, the coupled crystalline lattice. He split the ice I  $\nu_1$  mode into in-phase and out-of-phase components at 3083 cm<sup>-1</sup> and 3420 cm<sup>-1</sup>, respectively and the  $\nu_3$  band into the TO and LO components of 3209 cm<sup>-1</sup> and 3323 cm<sup>-1</sup>, respectively. Whalley contended that there should be strong similarities between the spectra of ordered and disordered ice Ic, and that the spectra of ice Ic and Ih are identical.

Scherer and Snyder (15) measured the Raman intensities of a single crystal of ice Ih at various crystallographic orientations. Their analysis suggested that the hydrogen bonds parallel to the crystallographic c-axis are more linear than the equatorial ones. They disagreed with Whalley's (14) 3209 cm<sup>-1</sup> TO  $\nu_3$ , and 3420 out-of-phase  $\nu_1$  assignments, but instead preferred assignments using various combinations of  $\nu_1$ ,  $\nu_3$ , and  $\nu_t$ .

The reduced vibration (RVM) model for describing the amorphous and polycrystalline ice Ih spectra was introduced by Bergren, Schuh,



Sceats, and Rice (16). They claimed that the inter and intramolecular coupling of OH oscillators is almost equal, thus any vibration of the condensed phase does not necessarily have a one-to-one parentage with vibrations of the free molecule. From their data Bergren et al. determined the optical constants of  $H_2O$  and  $D_2O$  polycrystalline ice Ih and amorphous ice.

Sivakumar, Rice, and Sceats (17) examined the Raman spectra of the OH stretching region of low density amorphous water and polycrystalline ice Ih. They found that the low temperature amorphous phase is stable at 10 K for at least two months. The spectra support a structure derived from ice Ih which is created by introducing a distribution of O-O-O angles. This angle and separation dispersion is partly responsible for the observed spectral width of the bands.

Scherer (18) published an extensive review on the vibrational spectroscopy of water. Early in the review, Scherer discussed the importance of Fermi resonance and the formation of Evans holes. Much controversy has resulted from whether or not there is Fermi resonance between  $\nu_1$  and  $2\nu_2$  in the liquid and condensed phases of ice. Scherer pointed out that Fermi resonance between two vibrational states is not as simple as a frequency displacement between two interacting states. Often the magnitude of the interaction constant and the bandshapes of the interacting states can combine to give unusual bandshapes. When a broad band Fermi interacts with a narrow band, the resultant is often a broad band with a spike or hole (Evans hole) that is not a simple sum of the two component bands. This can complicate the interpretation of the possible  $\nu_1$  and  $2\nu_2$  Fermi interaction.

He pointed out that many liquid water bandshape analyses and

bandfitting attempts are based on false assumptions because the ice data and widths of infrared bands do not support a spectral breadth arising from a small number of molecular species, but from a broad distribution of OH bonds strengths, a disorder in the hydrogen bonding environment, and intermolecular coupling.

Scherer stated that the ice Ic to Ih transition seems to be dependent on the thermal history of the particular sample, and lies between  $-120^{\circ}\text{C}$  and  $-63^{\circ}\text{C}$ . He also raised the question over how much does the disorder from true oxygen tetrahedral coordination widens the OH stretch band versus intermolecular coupling. Ice II which has an ordered proton structure yields four sharp crystallographically distinct  $\nu_{\text{OD}}$  bands. Scherer explained the wide OH stretch region of ice as follows: an uncoupled OH or OD ice oscillator is split under the intramolecular coupling into two vibrations that would be equivalent to  $\nu_3$  or  $\nu_1$  in the gas phase in their in-phase and out-of-phase character. Strong intermolecular coupling present in the ice lattice would further split each of these into two more components yielding a total of four or more components that could be further broadened by the proton disorder induced broadening effect of hexagonal, cubic, or amorphous ice; and complicated by Fermi resonance between  $2\nu_2$  and a smeared  $\nu_1$  vibrational state. Since intermolecular coupling increases with decreasing temperature, and the ice OH and OD bandwidths remain essentially unchanged with temperature, Scherer claimed that disorder must play a role in the OH or OD stretch width at higher temperatures. In conclusion Scherer (18) stated:

- (1) Fermi resonance between OH stretching and  $2\nu_2$  states has been observed.

- (2) Band fitting procedures that approximate the Evans hole with two bands are invalid.
- (3) Intermolecular coupling between adjacent water molecules is important in ice, amorphous ice, and liquid water.
- (4) Isobestic regions in the Raman spectra of water can not be interpreted as evidence for different molecular species.
- (5) Disorder is responsible for the bandwidth in liquid water and amorphous ice.
- (6) Liquid water at low temperatures is similar to ice near the melting point or amorphous ice at  $-180^{\circ}\text{C}$ .
- (7) There is evidence for non-tetrahedral HOH valence angles in ice, with some orientational preference for the c-axis hydrogen atoms to form more linear hydrogen bonds. (p. 211).

Rhim, Burum, and Elleman (19) studied the proton anisotropic shift in single crystal hexagonal ice and confirmed the Pauling model of ice but could not confirm or reject any possible proton ordering of ice. No hydrogen motion was observed at the liquid nitrogen temperatures used.

Recently Hagen, Tielens, and Greenberg (20) performed an extensive temperature study on the ice Ia to Ic transition between 10 and 140 K. They found the transition to occur between 130 and 140 K; with a 45 minute duration at 140 K, and 90 minute duration at 135 K. By monitoring the interference of a laser beam through the deposit Hagen et al. were able to measure the thickness of the sample. They calculated a  $dv/dR_{\text{o-o}} = 6700 \text{ cm}^{-1}$  (where  $R_{\text{o-o}}$  is the oxygen-oxygen lattice separation) and noted that the amorphous to cubic ice transformation was dependent on the thermal history of the sample.

## Matrix Isolated Decoupled Ice Spectra

One of the first matrix isolated systems to be studied was water isolated in solid nitrogen by Thiel, Becker, and Pimentel (21). By obtaining infrared spectra at various water to nitrogen ratios ( $M/A = 1000$  to  $10$ ; where  $M =$  moles matrix gas and  $A =$  moles water) Pimentel et al. distinguished between monomer, dimer, and possibly trimer bands of water. The bands at  $3725\text{ cm}^{-1}$  and  $3627\text{ cm}^{-1}$  were assigned to the  $\nu_3$  and  $\nu_1$  stretch of monomeric water and the  $3691\text{ cm}^{-1}$  and  $3546\text{ cm}^{-1}$  OH stretches to the dimer species. Because of the low number of dimer bands observed Pimentel et al. claimed the water dimer had a cyclic structure. They also noted a tenfold increase in the absorption coefficient of polymeric water relative to monomeric water.

In 1960 Haas and Hornig (22) published a benchmark paper in the study of ice. The spectra of HOD mixed with either a  $\text{H}_2\text{O}$  or  $\text{D}_2\text{O}$  matrix at  $-80^\circ\text{C}$  or  $-190^\circ\text{C}$  were recorded. The coupled  $\text{H}_2\text{O}$  and  $\text{D}_2\text{O}$  fundamentals and combinations were assigned, and interestingly, they were able to observe and assign the corresponding isotopically decoupled HOD modes. The HOD OD stretch was observed at  $2416\text{ cm}^{-1}$  in a predominately  $\text{H}_2\text{O}$  matrix, and the HOD OH stretch at  $3275\text{ cm}^{-1}$  in a predominately  $\text{D}_2\text{O}$  matrix. An unusual facet of the HOD OD stretch was the appearance of two distinct sidebands at  $2393\text{ cm}^{-1}$  and  $2442\text{ cm}^{-1}$  symmetrically placed about the  $2416\text{ cm}^{-1}$  OD stretch. The bands start appearing at HOD concentrations above 10 percent HOD in the ice matrix. Haas and Hornig attributed these sidebands to the coupling of adjacent or near neighbor OD oscillators producing an in and out-of-phase OD coupling at  $2393\text{ cm}^{-1}$  and  $2442\text{ cm}^{-1}$  respectively. The

narrow bandwidth of the HOD stretches in the presence of strong hydrogen bonding suggested that the large breadth of OH and OD stretches in ice spectra were due not to hydrogen bond variation but intermolecular coupling. The overtone  $2\nu_{\text{OH}}$  was assigned to  $6300 \text{ cm}^{-1}$  with a remarkably wide  $60 \text{ cm}^{-1}$  bandwidth even at low HOD concentrations where intermolecular coupling is absent. The large  $2\nu_{\text{OH}}$  width along with the greater  $\nu_{\text{OH}}$  width compared with  $\nu_{\text{OD}}$  suggested that a mass dependent phenomena such as tunneling was responsible for the  $2\nu$  width. They postulated that since the hydrogen can occupy two symmetric sites on a given oxygen-oxygen bond that a given proton could tunnel from one side of the oxygen-oxygen double well potential to the other. The presence of the  $2\nu_{\text{OH}}$  band also suggested that the thermal activation energy must be above 18 kcal but below 27.5 kcal if the brief Haas and Hornig calculations are accepted.

Tursi and Nixon (23) studied the water dimer in solid nitrogen. The  $\text{H}_2\text{O}$ ,  $\text{D}_2\text{O}$ , and HDO isotopes were isolated in  $\text{N}_2$  matrices at 20 K and 4 K. The sample to matrix ratios were varied and M/A plots as in Pimentel's paper were constructed. They observed a lack of the OH stretch in the HOD dimer implying that the deuterium hydrogen bond was stronger than the H hydrogen bond; thus producing a lack of these species in the matrix. Murby and Pullin (28), Bentwood, Barnes, and Orville-Thomas (29), and Barnes and Suzuki (30) have recently reinvestigated the self-association of matrix isolated water and concurred with Tursi and Nixon (23) that there seemed to be a preference for HOD to form OD hydrogen bonds rather than OH hydrogen bonds.

The decoupled frequencies of  $\text{OH}_2$  and  $\text{OD}_2$  oscillators in ice were

highly desired as appropriate starting points from which to build a coupled model of ice vibrations. In a series of papers (24, 25, 26) Devlin et al. elucidated these frequencies and assigned the vibrational frequency pattern of the isotopic ices. Ritzhaupt and Devlin (24) published the  $\nu_3-\nu_1$  OD stretch region of  $D_2O$  isolated and decoupled in a glassy  $H_2O$  matrix at 90 K. By simultaneously depositing a small amount of  $D_2O$  vapor with a large quantity of  $H_2O$  vapor they were able to obtain the sample with virtually no HOD contamination in the OD stretch region of isolated  $D_2O$  and with little isotopic scrambling at 90 K.

In a latter paper Ritzhaupt, Thornton, and Devlin (25) developed the simultaneous deposition technique so that  $D_2O$  spectra decoupled and isolated in cubic  $H_2O$  could be realized. They deposited a 0.1 micron film of  $H_2O$  (Ic) at 180 K, prior to the codeposition of dense  $H_2O$  and dilute  $D_2O$  vapors on a substrate held at 135 K. The epitaxial layer insured the ice deposit would be cubic, and the 135 K temperature that no deuterium or proton exchange occurred in the sample preparation period. The ice Ic  $D_2O$  stretching frequencies were identified as  $\nu_1 = 2367 \text{ cm}^{-1}$  and  $\nu_3 = 2444 \text{ cm}^{-1}$  for  $D_2O$  in  $H_2O$ . The bending modes ( $\nu_2$ ) were  $1220 \text{ cm}^{-1}$  for  $D_2O$  in cubic ice and  $1230 \text{ cm}^{-1}$  for amorphous or glassy ice. Both of these bending mode frequencies were high enough to discount any significant Fermi resonance interaction between  $2\nu_2$  and  $\nu_1$  in  $D_2O$  decoupled ice spectra. By raising the sample temperature to 140 to 160 K they were able to initiate proton-deuterium exchange as evidenced by the growth of the HOD band and decay of the  $D_2O$  bands. The ability to measure the exchange rates at different temperatures presented the opportunity of measuring the activation energy of the

exchange process. This they found to be approximately 10 kcal; well within the range suggested by Haas and Hornig (22).

By 1980 Ritzhaupt, Collier, Thornton, and Devlin (26) were using the epitaxial method of cubic ice deposition to unravel the uncoupled  $\text{H}_2\text{O}$  ice  $\nu_3$  frequencies. About 5 percent  $\text{H}_2\text{O}$  was codeposited with a  $\text{D}_2\text{O}$  matrix at 135 K. By allowing the sample to warm to 155 K the growth and decay of the appropriate isotopic exchange bands were observed. Both Raman spectra with polarization measurements and FTIR spectra were obtained. The use of digitized FTIR spectra allowed for the subtraction and scaling of different isotopic exchange spectra such that the change and identification of fundamental  $\text{H}_2\text{O}$  modes was enhanced. Ritzhaupt et al. placed the decoupled  $\text{H}_2\text{O}$  cubic ice  $\nu_3$  at  $3270 \text{ cm}^{-1}$ ,  $\nu_1$  at  $3200 \text{ cm}^{-1}$ , and  $\nu_2$  at  $1732 \text{ cm}^{-1}$ . Likewise the HOD  $\nu_{\text{OH}}$  was assigned to  $3270 \text{ cm}^{-1}$ ,  $\nu_{\text{OD}}$  to  $2418 \text{ cm}^{-1}$ , and  $\nu_2$  to  $1510 \text{ cm}^{-1}$  for a  $\text{D}_2\text{O}$  matrix and  $1465 \text{ cm}^{-1}$  for a  $\text{H}_2\text{O}$  matrix. Sceats, Stavola, and Rice (27) from their theoretical calculations suggested that Fermi resonance of  $2\nu_2$  with  $\nu_1$  actually pushed  $\nu_1$  and  $\nu_3$  bands into coincidence with the  $\nu_{\text{OH}}$  band yielding a single absorption. Ritzhaupt et al. showed that the coincidence of the cubic ice  $\nu_3$  and  $\nu_{\text{OH}}$  was correct but the Fermi resonance of  $2\nu_2$  with  $\nu_1$  actually increased the  $\nu_{\text{OH}} - \nu_1$  separation. With all the observed decoupled isotopic cubic ice frequencies available, Ritzhaupt et al. were able to perform overlay normal mode calculations for the calculated decoupled frequencies and quadratic force constants.

#### Theory of Ice Spectra

Much of the recent theoretical work done on coupled and decoupled

ice spectra has been done by Rice and Sceats who developed a method for predicting vibrational modes in highly intermolecularly coupled systems that is called the reduced vibrational model or RVM that greatly simplifies calculations.

McGraw, Madden, Bergren, Rice, and Sceats (31) used the RVM model to build a 64 molecule ice Ih cell in which the symmetry, ice rules, and zero net dipole moment were met. Like Whalley et al. (10) and Haas and Hornig (22), they contended that strong intermolecular coupling had to be accounted for before any realistic assignments could be made about the ice I stretching region. In examining their results they found that long range coupling between molecules not directly hydrogen bonded served to make features more diffuse but not to greatly affect the spectral distribution. In addition, the intermolecular coupling was predicted to be of the same magnitude as the intramolecular coupling. Rice et al. found that the modes predicted by the 64 molecule cell were complex mixtures of molecular motions such that identification of regions of the ice Ih stretching IR and Raman spectra with isolated molecular modes was not useful, but the predicted spectra agreed well with the observed spectra.

In a follow-up paper Madden, Bergren, McGraw, Rice, and Sceats (32) extended the theory and model to the OH stretching region of amorphous ice. Interestingly they found that the random orientation required of the new model cell had little affect on the predicted spectra but changes in the oxygen-oxygen lattice distances had a sizable one. They concluded that a distribution of force constants such as would result for oxygen-oxygen lattice changes rather than a change from the regular crystalline topology of ice Ih was



responsible for the infrared spectral changes between hexagonal and amorphous ice. This could lead to some interesting conclusions regarding the spectra of liquid water.

Sceats and Rice (33) utilized the overtone spectrum of ice I to develop an intramolecular potential. They assumed the intramolecular potential could be accurately described by a hydrogen bond dependence of the diagonal harmonic stretching force constant and the stretch-stretch interaction force constant. All other harmonic and diagonal anharmonic force constants were presumed to remain independent of hydrogen bond perturbations and hence could be derived from gas phase values or other sources. The results of their potential analysis of the ice overtone region indicated that the breadth of the  $2\nu_{OH}$  overtone may be due to intermolecular coupling rather than tunneling as suggested by Haas and Hornig (22).

The experimental work of Ritzhaupt et al. (24, 25) prompted Sceats, Stavola, and Rice (34) to extend their models to reproduce the stretching spectrum of  $D_2O$  in  $H_2O$  and  $H_2O$  in  $D_2O$  in both amorphous and crystalline ice I. They interpreted the amorphous  $D_2O$  decoupled in  $H_2O$  such that Fermi resonance of a near coincident  $2\nu_2$  with  $\nu_1$  pushes  $\nu_1$  further away from  $\nu_3$  thus increasing what should have been a smaller  $\nu_1-\nu_2$  splitting and increasing the width of  $\nu_1$  with the shifted  $2\nu_2$  state.

In cubic ice the situation is similar but the coincidence of  $2\nu_2$  and  $\nu_1$  is not as great and hence the resultant Fermi splitting gives a component on the high frequency side of  $\nu_3$ . For  $H_2O$  decoupled in  $D_2O$  they calculated a spectra with a very small  $\nu_1-\nu_3$  splitting such that Fermi resonance of  $2\nu_2$  with  $\nu_1$  produces a single peak centered

around the HDO OH decoupled stretching frequency. In a later paper by Ritzhaupt et al. (26) this was confirmed except that the  $\nu_1$  frequency was assigned to a much lower frequency rather than coincident with  $\nu_3$  and  $\nu_{OH}$  as predicted by Sceats, Stavola, and Rice.

Morse and Rice (35) compared different pair potentials suggested by researchers for computer simulations of ice and water. Since the structures of ice I, II, III, IV, V, VI, and VII are known, the accuracy with which these structures could be predicted by the pair potentials would be a sensitive test of their validity. They tested several potentials in predicting a proton ordered ice Ih structure and found the ab initio pair potential calculated by Matsuoka, Clementi, and Yoshimine (36) succeeded best.

#### Proton Transfer in Ice

Proton transfer has been studied in many systems by widely varying techniques. Much of the earlier work on proton transfer in ice and liquid water was accomplished through conductance studies. One noteworthy theoretical and experimental study was that of Conway, Bockris, and Linton (37). They developed a theory of proton conduction in liquid water that broke the conduction into two processes, one quantum mechanical tunneling of a hydrogen from one molecule to another and then, two, rotation of the molecule breaking the surrounding solvent cage bonds to transfer the hydrogen or proton to the other side of the cage. By such a series of transfers and rotations a proton could be transported through the liquid. Conway et al. claimed that the tunneling transfers occurred rapidly and the classical rotations much slower and thus the rotational step was rate determining. They

predicted from Morse potential curves that the classical proton transfer activation energy in liquid water was 3.85 kcal for hydrogen and 5.81 kcal for deuterium and the predicted proton to deuterium transfer rate ratio was 6:1. The rotational activation energy was found to be quite costly energetically and dependent on the mode of rotation. They calculated that the proton spends 1 percent of its time undergoing quantum mechanical tunneling transfers. In a note Conway and Bockris (38) suggested that if their theory was applied to ice the concentration of defect formed protons might be low enough that water molecule rotation may no longer be the rate determining step and hence the tunneling process could become rate determining and hence the mobility of the proton would increase until it was limited by the now faster rate determining tunneling transfer step.

Kim and Schmidt (39) developed a theory of proton transport in ice that suggested the mobility is limited not by barrier hopping but lattice scattering. One of the implications of this particular theory is that the mobility is predicted to increase with decreasing temperature.

Conductivity measurements of ice have been studied over several decades with widely varying results and much controversy. However by the mid 1960's much of the discrepancy was resolved and the problem traced to the use of two terminal conductivity cells which gave conductivities and resulting activation energies that were a mixture of bulk and surface ice conductivities. The use of guarded or three terminal conductivity cells or holders lead to more consistent measurements among researchers. Durand, Deleplanque, and Kahane (40) utilized ion exchange membranes for electrodes in conductivity measurements. The

temperature of the ice was kept between  $-20^{\circ}\text{C}$  and  $-100^{\circ}\text{C}$  to prevent surface conductivity effects thus negating the need for a third guard electrode. They measured a conductivity thermal activation energy of about 9 kcal/mole between  $-30^{\circ}\text{C}$  and  $-70^{\circ}\text{C}$  and a 13 kcal/mole activation energy between  $-90^{\circ}\text{C}$  and  $-100^{\circ}\text{C}$ .

In another similar conductivity measurement Bullemer, Eisel, Engelhardt, Riehl, and Seige (41) determined the bulk and surface conductivity of ice at  $-10^{\circ}\text{C}$ . They found an activation energy of  $8.5 \pm 0.5$  kcal/mole for the bulk conductivity and 26-30 kcal/mole for surface conductivity using a three terminal guarded potential probe method.

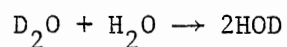
An alternate theoretical view of proton conduction in ice has been advanced by A. von Hippel et al. (42) in a series of papers describing their experimental and theoretical technique. A. von Hippel postulated that the Bjerrum D and L defects were formed in a single asymmetric stretch and bend combination photon excitation that did not require a molecular rotation that broke three hydrogen bonds. The calculated activation energy was about 13 kcal. Furthermore he predicted that Bjerrum defects would be the current carriers in the free volume of the crystal with ionic defects only playing a role near the electrodes. The movement of defects in the free volume was expected to be in polaronic fashion, and in rapidly expanding drift ellipsoids when under the influence of an externally applied electrical field. This is in contrast to the theory of Jaccard (43) who claimed that D, L Bjerrum defects were the majority current carriers but ionic defects also played a role and had a 400 times higher mobility which was explained by proton tunneling. Experimentally A. von Hippel et al.

measured a surface conduction activation energy of about 33 kcal/mole between 0°C and 30 °C which could be reduced to 20 kcal/mole by enclosing the sample in a vacuum. The bulk conductivity activation energy was about 12 kcal/mole with a three terminal holder. Below -35°C both two and three terminal holders gave consistent results of about 12 kcal/mole. Previously measured zero activation energies were claimed to result from measurements made in the initial millisecond time range where a stable conductive state had not yet been reached. Ice crystal age was also found to affect conductivity measurements. A. von Hippel did not feel the need to invoke proton tunneling to explain proton mobility in ice.

Hubmann (44) extended Jaccard's theory to accept the complementary action of ionic defects and Bjerrum D, L defects in proton migration in ice.

Evidence for proton tunneling and c-axis order in ice Ih was found by Lumpkin and Dixon (45). They measured the  $^{17}\text{O}$  quadrupole resonance of ice Ih at 4.2 and 1.3 K using a double resonance technique and found only two resonances consistent with the explanation offered.

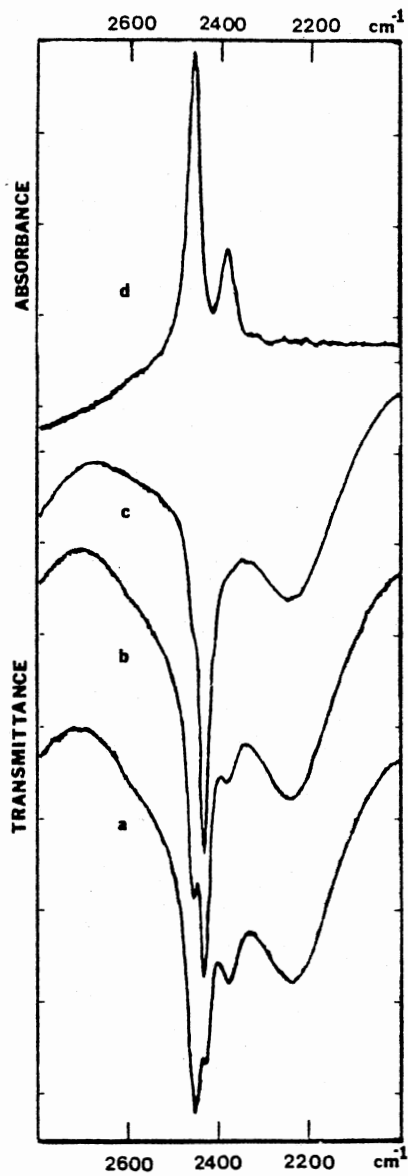
Ritzhaupt and Devlin (46) took the thin layer techniques developed in references 25 and 26 and isolated  $\text{D}_2\text{O}$  in cubic  $\text{H}_2\text{O}$  at 130 K. Two bands at  $2444\text{ cm}^{-1}$  and  $2367\text{ cm}^{-1}$  were assigned to the  $\text{D}_2\text{O}$   $\nu_3$  and  $\nu_1$  respectively. Upon warming the sample to 165-180 K the  $\text{D}_2\text{O}$   $\nu_3$  and  $\nu_1$  bands lost intensity while a band at  $2418\text{ cm}^{-1}$  attributed to the OD stretch of HOD grew in and gained intensity representative of the following exchange reaction.



Ritzhaupt and Devlin obtained FTIR spectra at several time

intervals during the exchange reaction at elevated temperatures. By subtracting the appropriate scaled pure HOD  $2418 \text{ cm}^{-1}$  component band from the  $\text{D}_2\text{O}$ -HOD exchange band complex the time dependence of  $\text{D}_2\text{O}$  loss could be followed by the  $\text{D}_2\text{O}$  band intensities after HOD subtraction or by the scaling factors used to scale the original  $2418 \text{ cm}^{-1}$  HOD band used in the subtraction. Figure 6 illustrates the spectral changes accompanying deuterium-hydrogen exchange. Over the temperature range of 165 K to 181 K the activation energy was found to be 9.3 kcal. They noted that this value was close to half the formation energy of ion pair defects (18.6 kcal) found by Bullemer et al. (41). This was considered indicative of the thermally activated proton tunneling mechanism suggested by Haas and Hornig (22) where a thermally activated  $\text{H}_3\text{O}^+$  ion defect is mobilized by proton tunneling. Such a process should reflect an Arrhenius type behavior. It was also suggested that such thermal activation would involve a ground to excited state stretching mode vibration whose energy of about 9.2 kcal/mode is amazingly close to the measured activation energy.

In a subsequent paper Ritzhaupt and Devlin (47) discovered that the previous samples of  $\text{D}_2\text{O}$  isolated in cubic  $\text{H}_2\text{O}$  discussed in reference 46 had been contaminated with minute quantities of aromatic compounds desorbing from the cell walls. As the contamination was slowly cleared up in subsequent experimental runs it was found that the rate of exchange reaction increased considerably with a new transitory peak at  $2396 \text{ cm}^{-1}$  slowly appearing as the reaction proceeded and then losing intensity as the exchange neared completion. Figure 7 shows the  $\text{D}_2\text{O}$   $2444 \nu_3 \text{ cm}^{-1}$ ,  $2367 \nu_1 \text{ cm}^{-1}$ , and  $2396 \text{ cm}^{-1}$  intermediary band spectral changes with time after interfering  $\text{H}_2\text{O}$  and HOD bands have been



Infrared curves for  $\approx 2\%$  D<sub>2</sub>O in H<sub>2</sub>O ice I<sub>c</sub>: (a) original deposit warmed to 165 K; (b) same sample after two hours at 165 K; (c) same sample (165 K) after seven hours at 165 K and following brief annealing at 185 K and (d) D<sub>2</sub>O absorbance spectrum from curve (b) using curve (c) to remove HOD absorbance.

Figure 6. D<sub>2</sub>O Spectral Changes with Isotopic Exchange

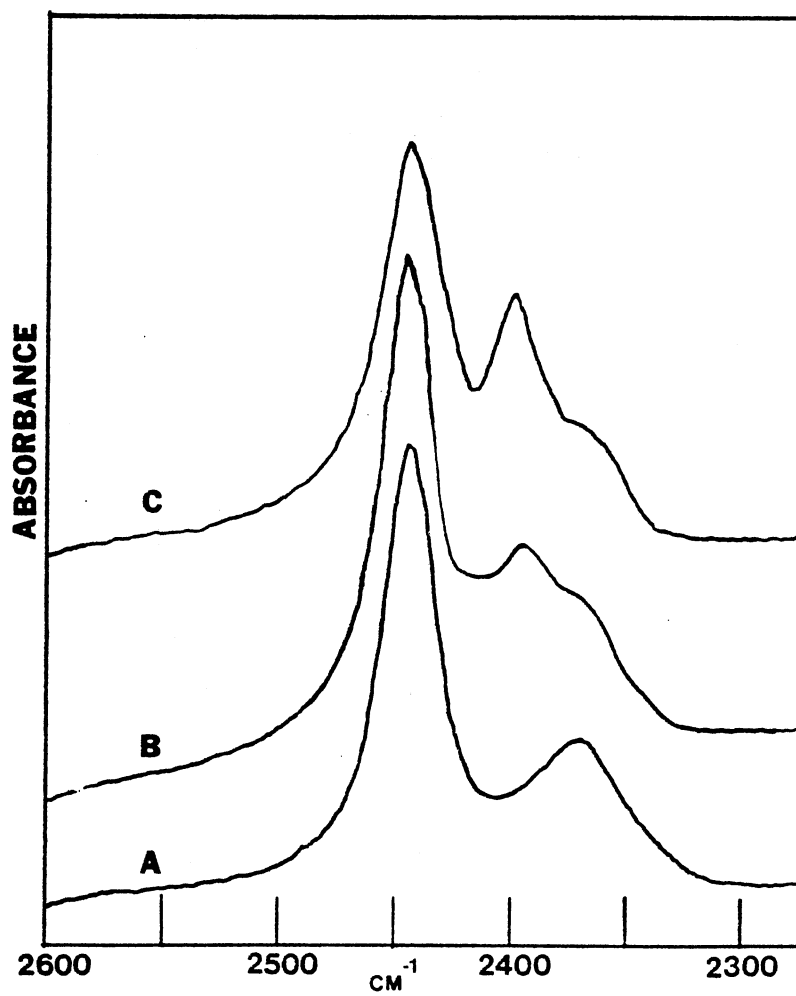
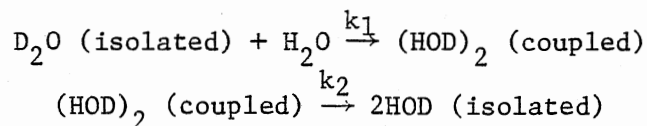


Figure 7.  $(\text{HOD})_2$  Spectral Changes with Isotopic Exchange



digitally subtracted. The  $2396\text{ cm}^{-1}$  band could clearly be identified with the  $2393\text{ cm}^{-1}$  dimer HOD coupled in-phase OD stretch assigned by Haas and Hornig (22). The  $2442\text{ cm}^{-1}$  out-of-phase coupled OD oscillator stretch should be obscured beneath the more intense  $\text{D}_2\text{O } \nu_3$  vibration at  $2444\text{ cm}^{-1}$ . This suggested the following exchange reactions.



Devlin et al. attributed the first step to a proton transfer which would produce a coupled HOD unit and the second step to a Bjerrum defect migration by a molecular rotation which decouples or destroys the coupled HOD producing two isolated HOD subunits. They explained the behavior of the doped and pure ice systems as follows. The doped system inhibited the transfer process enough that  $k_2 \gg k_1$  and hence the  $(\text{HOD})_2$  concentration was a very small steady state value and not spectrally observed. Hence the overall rate was slower and higher temperatures were needed to initiate the exchange. As the dopant was removed and  $k_1 \approx k_2$  the steady state situation was removed with a corresponding increase in  $(\text{HOD})_2$  concentration and greater overall exchange rates. The dopant appeared to have no effect on the rotational step. Rough calculations on the  $k_1$  rate or proton transfer process indicated an activation energy of 8.5 to 10.5 kcal. This suggested the dopants affected the proton transfer rate by trapping ion pair defects rather than altering the activation energy for defect formation. Devlin et al. suggested the possibility existed for elucidating the activation energy for both the ion pair defect formation and Bjerrum rotational defect migration.

In a recent note Kunst and Warman (48) used a microwave conduc-

tivity pulse radiolysis technique to find the mobility of protons and deuterons in ice. The proton virtual mobility in  $\text{H}_2\text{O}$  ice was determined to be  $.0064 \text{ cm}^2/\text{V}\cdot\text{s}$  at  $-5^\circ\text{C}$  and the deuteron virtual mobility in  $\text{D}_2\text{O}$  ice to be  $.0024 \text{ cm}^2/\text{V}\cdot\text{s}$  at  $-5^\circ\text{C}$  with a deuteron mobility to proton mobility ratio of .375.

#### Presentation of Problem

Ritzhaupt and Devlin (47) presented their analysis of the infrared  $\text{D}_2\text{O}$  stretching region of  $\text{D}_2\text{O}$  isolated and decoupled in cubic  $\text{H}_2\text{O}$  ice at 130-180 K. From the spectral changes observed it was concluded that proton or deuteron migration in ice followed a two fold process. The first step involved an ionic defect formation which mobilized a proton via a tunneling mechanism that allowed the defect to migrate throughout the crystal once formed. This proton transfer encountering a  $\text{D}_2\text{O}$  molecular subunit in an  $\text{H}_2\text{O}$  ice crystal would initiate isotopic scrambling by proton transfer; i.e. exchanging and transferring a proton for a deuteron leaving a vibrationally coupled  $(\text{HOD})_2$  dimer. The  $(\text{HOD})_2$  dimer would further scramble to two vibrationally decoupled and more spatially removed HOD molecular subunits in the  $\text{H}_2\text{O}$  ice crystal lattice by a molecular rotation or Bjerrum D, L defect migration.

The goal of this study is to find the best quantitative estimate of the ionic defect formation and Bjerrum D, L defect migration activation energies and examine their possible implications for proton transport in hydrogen bonded molecular networks. The study of these energies could yield possible insights into the role proton transport could have in alternate electrical control devices, and energy conversion and storage in biological systems.

Accomplishing this study will involve developing experimental techniques to implement the following procedures.

- (1) Collection time resolved IR spectra of  $D_2O-H_2O$  exchange processes over a suitable range of sampling times, time durations, samples, and experimental temperatures to perform kinetic analyses of the data obtained.
- (2) Assure sample purity such that the rates obtained reflect a reasonably pure ice situation.
- (3) Resolve the  $D_2O-HOD$  band complex obtained in each sample collection into its individual spectral components.
- (4) Relate the spectral intensity of the spectral components to each component concentration.
- (5) Develop a kinetic model for obtaining the proton transfer rate via ionic defect migration and Bjerrum D, L defect rotation rate from the spectral component concentrations.
- (6) Find the proton transfer and rotational rates over as wide a temperature range as experimentally feasible and find their respective activation energies.
- (7) Analyse the results for implications.

## CHAPTER II

### EXPERIMENTAL PROCEDURES AND RESULTS

#### Matrix Isolation Methods

Matrix isolation techniques are usually used to isolate a particular molecular specie in an inert matrix. This is accomplished by directing the gaseous mixture of the molecular specie and an inert gas matrix onto a cryogenic surface where the frozen gas mixture may be examined by optical spectroscopy or other means.

For this work matrix isolation techniques were used but the matrix gas of  $\text{H}_2\text{O}$  interacted significantly with the guest molecule  $\text{D}_2\text{O}$ . Though the electronic interaction between the matrix and guest were large, the different isotopic masses caused a decoupling of specific normal modes of vibration between the matrix and guest molecules. So the present work is better described as matrix decoupled spectroscopy.

All of the experimental work was accomplished by depositing vapors of  $\text{H}_2\text{O}$  and  $\text{D}_2\text{O}$  on a nitrogen gas cooled or mechanically cooled cesium iodide substrate. Two types of vapor deposition cells were used.

#### Metal Cell Technique

The first phase of the experimental work consisted of creating crystalline deposits of ice  $\text{I}_h$  with a 5% to 95% ratio of  $\text{D}_2\text{O}$  to  $\text{H}_2\text{O}$ . Figures 8 and 9 show the experimental apparatus used for this. The heart of the cryostat is an Air Products and Chemicals,

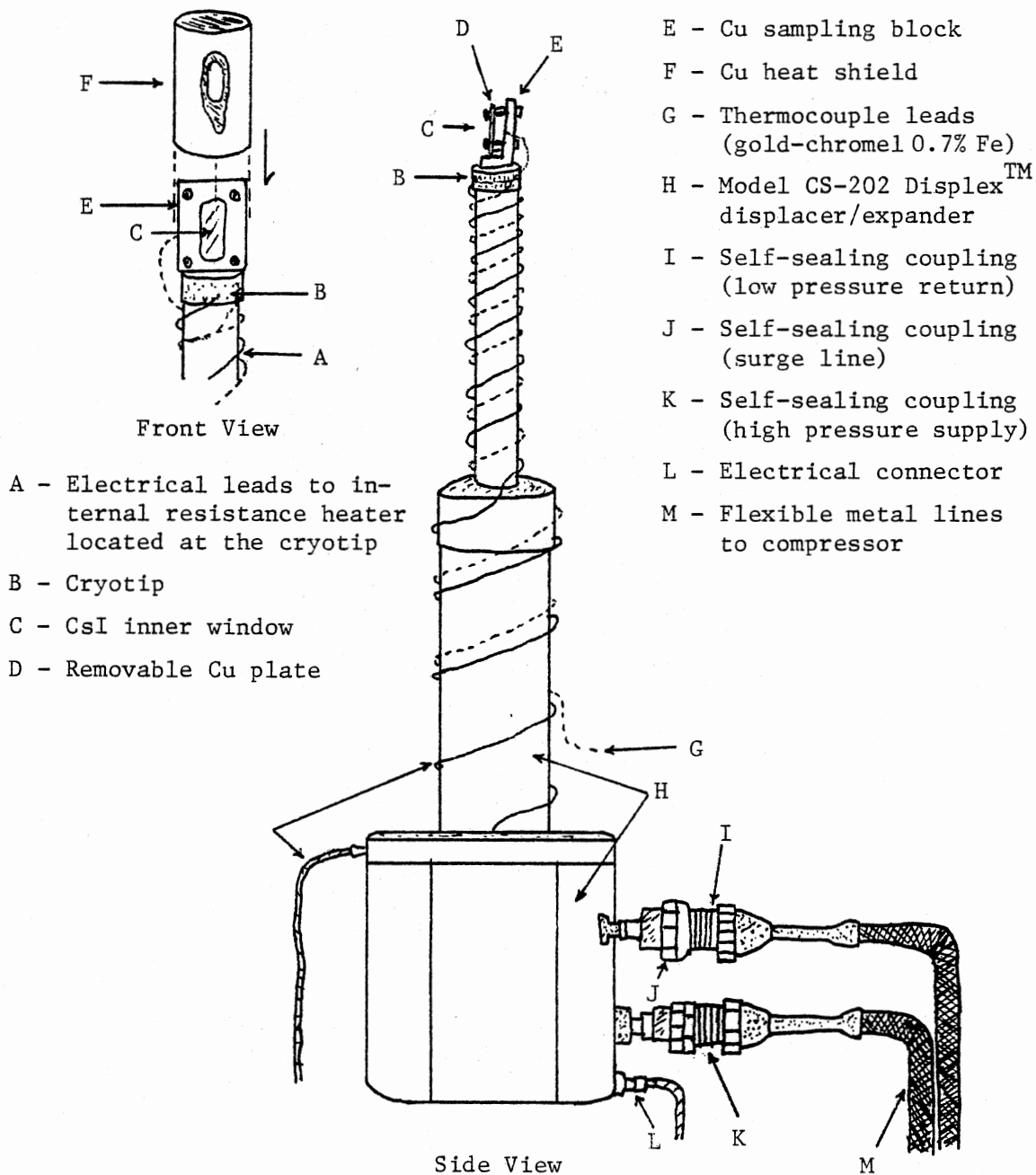


Figure 8. Air Products CS-202 Displacer/Expander Module with Tip and Sample Deposition Window

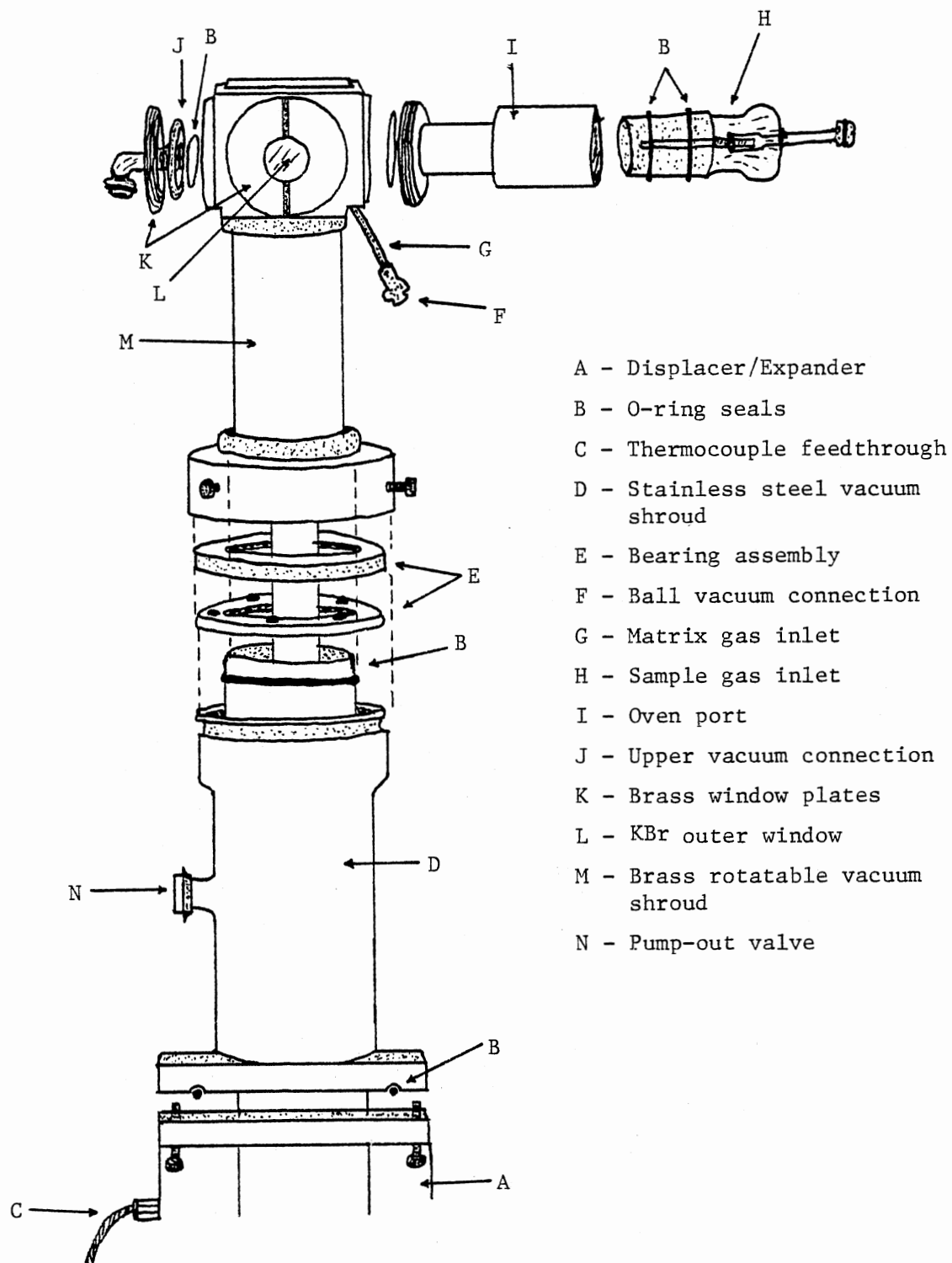


Figure 9. Brass Vacuum Shroud for Displacer/Expander Module

Inc. Model CS-202 Displex closed cycle helium refrigerator. Included in the system is a displacer/expander module connected by flexible lines to a compressor and control panel. The displacer/expander module has a refrigerated tip on which the cesium iodide substrate window is mounted. The module may assume any orientation and has a capability of temperature to 10 K. A small resistance heater mounted on the refrigerator tip allows temperature control from 10<sup>o</sup>K to 300<sup>o</sup>K with an accuracy of  $\pm .5^{\circ}\text{K}$ . If an insulating vacuum of 10<sup>-4</sup> torr is maintained around the cryotip, the cool-down time from room temperature to 10<sup>o</sup>K is approximately one hour. On the end of the cryotip is a copper block for mounting a cesium iodide substrate window for infrared transmission studies. Four screws and a copper backing plate holds the window to the salt plate holder. Thermal contact is assured by indium wire gaskets between the window and backing plate and copper block. Two gold (0.7% Fe) chromel thermocoupler measure the cryotip and CsI window temperature. One is attached directly to the refrigerator tip with indium solder for the internal heater control; and the other is sandwiched between the CsI and copper block in the indium gasket to monitor the CsI window surface temperature on the deposition side.

To maintain the necessary vacuum for vapor deposition and cryogenic temperatures a brass vacuum shroud was available that would cover and seal the displacer/expander module allowing vacuum operation and injection of sample and matrix gasses through two injection ports (Figure 9). Also as shown in Figure 9 the shroud has the opposing potassium bromide windows and a vacuum connection. The shroud was mounted on bearings to allow rotation of the shroud while under vacuum and was

equipped with vacuum valves so that once a good vacuum was reached, the displacer/expander module could be sealed off and removed from the vacuum pumps and removed to a FTIR infrared spectrophotometer for infrared studies of the deposit made on the cryotip cesium iodide window. The brass shroud could be rotated to bring the two shroud potassium bromide windows in line with the inner cesium iodide window for easy optical examination.

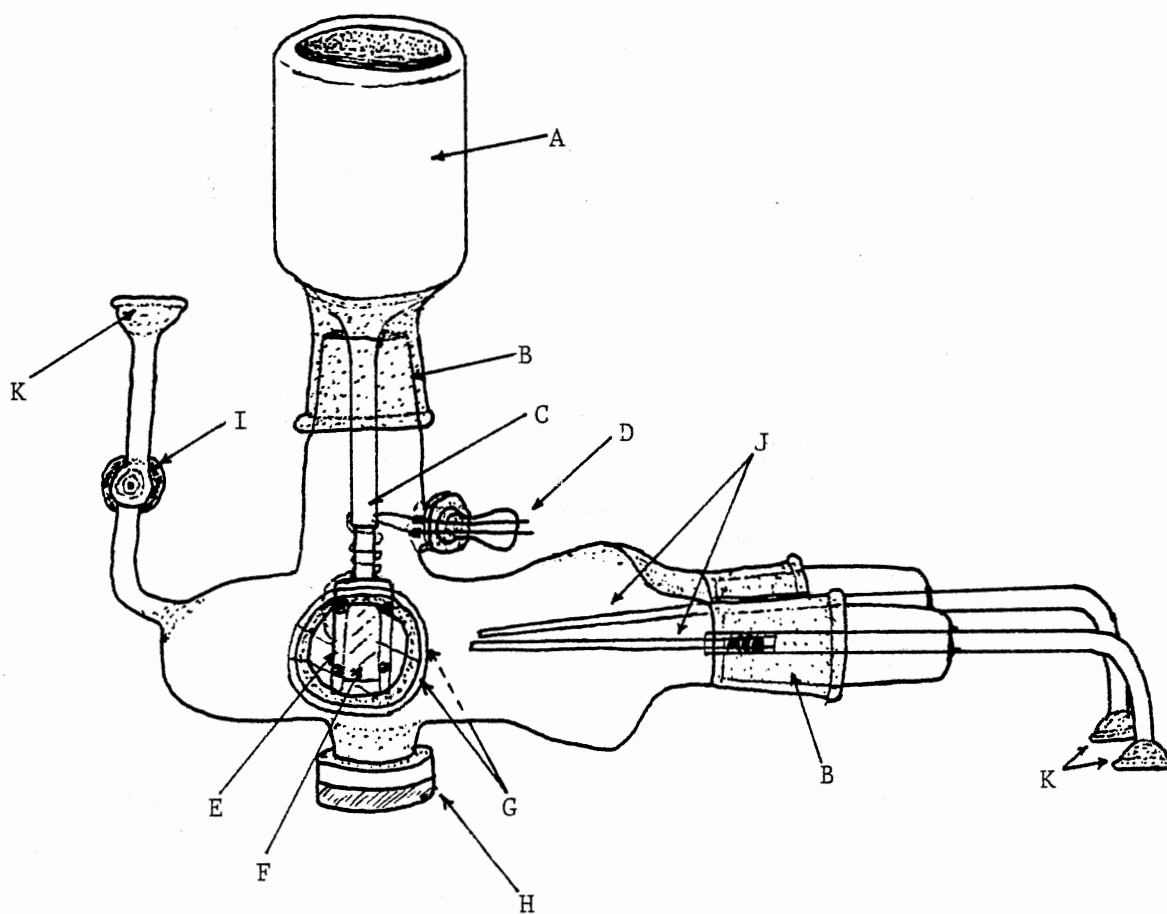
The matrix gas  $\text{H}_2\text{O}$  is introduced through G and the sample gas  $\text{D}_2\text{O}$  through H.  $\text{D}_2\text{O}$  vapors are sent to the cryostat from a vacuum valve sealed glass reservoir connected to a micro-flow needle valve that is connected to the sample line at C. The vapor pressure of  $\text{D}_2\text{O}$  at  $0^\circ\text{C}$  provided the necessary  $\text{D}_2\text{O}$  gas for deposition.  $\text{H}_2\text{O}$  vapors were likewise introduced into the cryostat. The  $\text{H}_2\text{O}$  vapor flow rate is controlled by a teflon needle valve above the  $\text{H}_2\text{O}$  water reservoir and monitored by a Fisher-Porter 1/16 glass flowmeter inserted into the  $\text{H}_2\text{O}$  sample line.

All vacuum seals in the system are made with lightly greased viton O-rings except for the  $\text{D}_2\text{O}$  and  $\text{H}_2\text{O}$  sample and matrix gas lines where the vacuum grease was omitted to prevent any possible grease contaminations in the deposits.

#### Glass Cell Apparatus

During the course of the investigation it was found necessary to make  $\text{D}_2\text{O}$ - $\text{H}_2\text{O}$  deposits doped with 7-azaindole. Rather than risk contamination of the metal cell where pure  $\text{D}_2\text{O}$ - $\text{H}_2\text{O}$  deposits were being made, these samples were prepared in an all glass vapor deposition cell as shown in Figure 10. This cell being very similar in function





- |                                       |                                    |
|---------------------------------------|------------------------------------|
| A - Liquid nitrogen dewar             | G - Outer KBr windows              |
| B - Vacuum greased taper joint        | H - Glass window                   |
| C - Dewar stem glass-metal connection | I - Teflon vacuum valve            |
| D - Thermocouple wire vacuum seal     | J - Sample and matrix gas inlet    |
| E - Cu sample block                   | K - Female ball vacuum connections |
| F - CsI window                        |                                    |

Figure 10. Glass Cryostat Cell for Vapor Deposition

to the closed-cycle helium refrigerator apparatus is comprised of a glass container with provisions for introducing the matrix and sample gasses through two sample lines terminating in nozzles aimed at a cesium iodide substrate window. The cesium iodide substrate window is attached to a copper holder with indium gaskets for thermal contact. The copper substrate holder is connected by a Kovar glass to metal seal to the bottom of a liquid nitrogen dewar for cooling. The dewar is attached to the glass cell by a greased glass taper joint so that the cesium iodide substrate window may be turned facing the sample gas inlet nozzles for deposition and then rotated  $90^\circ$  for optical spectroscopy through two potassium bromide windows fastened on opposing sides of the glass vapor deposition cell. The temperature was monitored by an iron-constantan thermocouple inserted into the indium wire gasket between the copper holder and deposition side of the inner cesium window. To obtain reproducible and consistent temperature measurements and correspondence between the metal cell and glass cell deposit temperature it was found necessary to temper the glass cell thermocouple leads by stripping off the outer insulation jacket of the thermocouple wire and wrapping the two individually insulated leads at least twenty times around the dewar stem attached to the copper holder before running the thermocouple leads outside the cell for voltage measurement. A Honeywell model no. 2745 potentiometer was used for all thermocouple temperature measurements. To check the temperature correspondence between the metal and glass cell measurements a series of spectra of  $D_2O$  decoupled in  $H_2O$  at various temperatures were made. The  $\nu_3$  asymmetric stretch bond of  $D_2O$  was expanded so that  $2^\circ K$  temperature shift effects could be detected in

the bandshapes as shown in Figure 11. By comparing spectra collected in the metal or glass cells with these spectra plots any temperature measurement inaccuracies could readily be detected. Both cells with different composition thermocouples were found to yield identical temperature shifted  $D_2O \nu_3$  bands at a given temperature setting over the 135 to 160 K range to at least a  $\pm 1$  K accuracy.

The matrix and sample gasses were introduced from  $D_2O$  and  $H_2O$  liquid reservoirs as in the metal cell apparatus. The 7-azaindole dopant was introduced by lightly dusting a glass wool plug with small crystals of 7-azaindole and placing it in the  $H_2O$  vapor sample line so that small quantities of sublimed 7-azaindole would be swept by the  $H_2O$  matrix vapors into the deposit. Temperature control with the glass cell was considerably more difficult. To circumvent the invariability of the liquid nitrogen boiling point for temperature control, a pressurized cool nitrogen gas cooler was designed as shown in Figure 12. A 50 liter liquid nitrogen dewar was mounted on a movable cart with a small glassware rack. The dewar was sealed off and a needle-valve controlled insulated tygon tubing line was used to direct the vented nitrogen gas to the glass cell dewar stem for cooling. Pressurization of the 50 liter dewar was provided by inserting a nichrome wire resistance heater into the liquid nitrogen of the 50 liter dewar and running the heater leads through a pressure seal in the 50 liter dewar top to a 110 AC variac. By controlling the heater voltage and adjusting the needle valve controlling the cool nitrogen gas flow, sufficient liquid nitrogen could be vaporized and blown into the glass cell dewar stem so that temperatures between 100 K and room temperature could be maintained with an accuracy of  $\pm .5$  K. By mounting

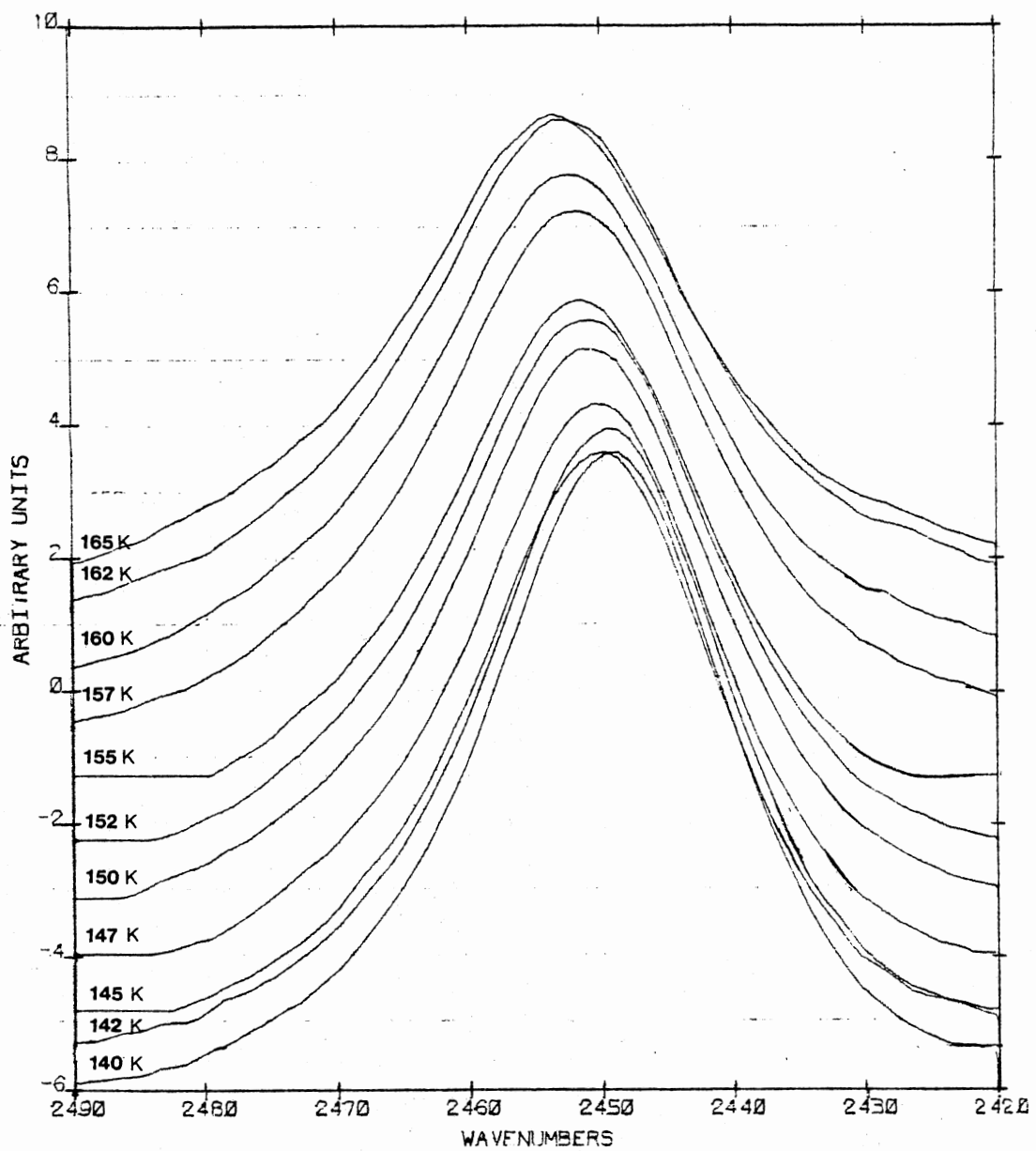
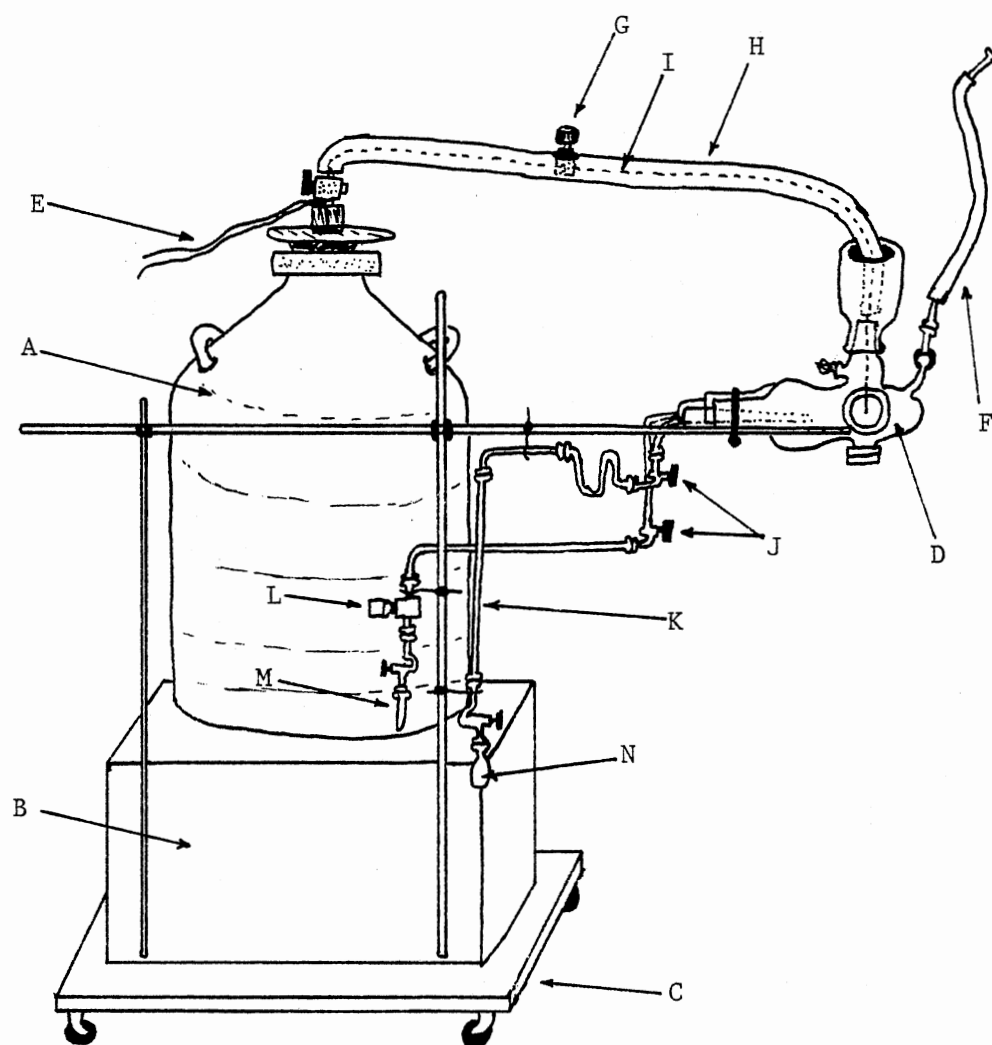


Figure 11. Expanded D<sub>2</sub>O Decoupled in H<sub>2</sub>O Ice  $\nu_3$  Bandshapes



- |                                     |                                |
|-------------------------------------|--------------------------------|
| A - Liquid N <sub>2</sub> dewar     | H - Foam insulation            |
| B - Dewar and rack support          | I - Tygon tubing               |
| C - Rollable cart                   | J - Vacuum valves              |
| D - Glass cryostat                  | K - Fisher-Porter flowmeter    |
| E - Resistance heater leads         | L - Micro-needle valve         |
| F - Flexible vacuum line            | M - D <sub>2</sub> O reservoir |
| G - N <sub>2</sub> gas needle valve | N - H <sub>2</sub> O reservoir |

Figure 12. Nitrogen Gas Cooler System for Glass Cryostat Cell

the glass dewar cell onto the 50 liter glassware rack the whole unit could be rolled to the vacuum rack for evacuation and deposition, then sealed off and disconnected from the vacuum rack to be moved to the FTIR for infrared transmission studies.

All infrared spectra were collected on a Digilab FTS-20C Fourier transform infrared spectrophotometer. This is a single beam instrument which collects digitalized interferograms via a Michelson interferometer for storage or immediate computer transform analysis to a real-time frequency infrared spectrum. All spectra can be ratioed against an earlier or later run reference spectra for background correction. The digitalized nature of the collected interferograms or computed spectra, dedicated mini-computer, and additional disk storage, allow for easy manipulation, expansion, addition and subtraction of infrared spectra. The faster collection speed, much improved accuracy, and superior quantitative features of FTIR spectroscopy were essential requirements for the completion of this work.

All spectra were collected with a triangular apodization function, a gain ranging factor (FTS-20C parameter GRR) of 40, a zero filling factor (ZFF) of 2, and detector signal sensitivities of 1 to 4 depending on the samples used and the particular experimental situation. The triangular apodization function weighs the interferogram to minimize tail fringes that may accompany Fourier transforms of infrared bands. The gain ranging factor determines the detector signal amplification changes over different parts of the interferogram to obtain the best signal to noise ratio. The zero filling factor adds additional zeroes to the end of the collected interferogram so that additional space filling points are integrated into the computed spectra. While not providing

additional resolution it smoothes out the digitalized nature of FTIR spectra. The overall detector sensitivity was adjusted to provide maximum signal gain without overflowing the analog to digital signal converter. The number of scans that were coadded to form the final interferogram varied from 10 to 500 scans depending on the time available for the scans, and signal to noise requirements.

#### Procedure

A typical  $D_2O/H_2O$  deposit and experimental kinetic analysis for the mechanically cooled metal deposition cell proceeded as follows.

If the cryostat's optical windows needed cleaning the cell assembly was partially dismantled, cleaned, and oven-dried to remove moisture, while the KBr and CsI windows were removed and repolished. Normally this was not required between runs, and the whole system was kept connected to the vacuum pump under a constant  $5 \times 10^{-3}$  torr vacuum to minimize contamination from outside air and remove volatile contaminations from inside the vacuum systems. Several hours before a run the diffusion pump was turned on, the liquid nitrogen cold trap filled, and the system allowed to evacuate to approximately a  $3 \times 10^{-4}$  torr vacuum. Next the cryostat was cooled to 170 K for an initial 1.5 minute deposit of  $H_2O$  vapor at a flow reading of 3.0 to 3.5 on the Fisher-Porter flowmeter scale. Subsequently the cryostat was cooled to 130 K for the main deposit. Since the main deposit was made at the amorphous to cubic ice transition temperature, an initial deposit of cubic ice at 170 K was needed to induce the main deposit to form cubic rather than amorphous ice. The crystallinity of the main  $D_2O/H_2O$  deposit was checked by examining the infrared spectrum in the 2800-2000

$\text{cm}^{-1}$  OD stretch region for features indicative of a crystalline or amorphous sample as shown in Figure 13. The main deposit of the  $\text{D}_2\text{O}/\text{H}_2\text{O}$  vapor mixture was made at 130 K for 20 minutes with the  $\text{H}_2\text{O}$  and  $\text{D}_2\text{O}$  flow rates adjusted to a 5% to 95% ratio of  $\text{D}_2\text{O}$  to  $\text{H}_2\text{O}$  vapors. The flow rate of  $\text{H}_2\text{O}$  on the Fisher-Porter flowmeter scale was held at approximately 7.4 while the  $\text{D}_2\text{O}$  flow rate was adjusted by trial and error to give the proper ratio using a micro needle valve. Both the  $\text{D}_2\text{O}$  and  $\text{H}_2\text{O}$  water reservoirs were immersed in a mixture of ice and water for temperature control.

Once the deposit was completed the cryostat was sealed off and removed from the vacuum line. Next it was moved to the FTIR spectrophotometer and suspended in the FTIR sample cavity so that the IR beam would pass through the sample but yet was mechanically isolated from the interferometer to minimize vibrations from the operating displacer/expander module. Then a  $2 \text{ cm}^{-1}$  resolution spectrum from  $4000\text{--}500 \text{ cm}^{-1}$  was collected to examine the nature of the deposit. Figure 14 shows a typical  $\text{D}_2\text{O}/\text{H}_2\text{O}$  ice FTIR spectrum. In all the deposits some HOD was found initially. The HOD OD stretch would fall between the  $\nu_3$  and  $\nu_1$  of  $\text{D}_2\text{O}$ . If the HOD to  $\text{D}_2\text{O}$  ratio was sufficiently small, the  $\text{D}_2\text{O}$  to  $\text{H}_2\text{O}$  ratio sufficiently close to an isolation ratio with little sign of dimer formation, and the sample crystalline, then a kinetic analysis was made. If not, the sample was prepared again.

To perform a kinetic analysis the cryostat was quickly warmed to the desired temperature (135, 140, 145, or 150 K) to initiate isotopic exchange, and then maintained there throughout the kinetic run with an accuracy of  $\pm .5 \text{ K}$ . Interferograms were then collected and filed at timed intervals at a  $2 \text{ cm}^{-1}$  resolution until exchange was complete or



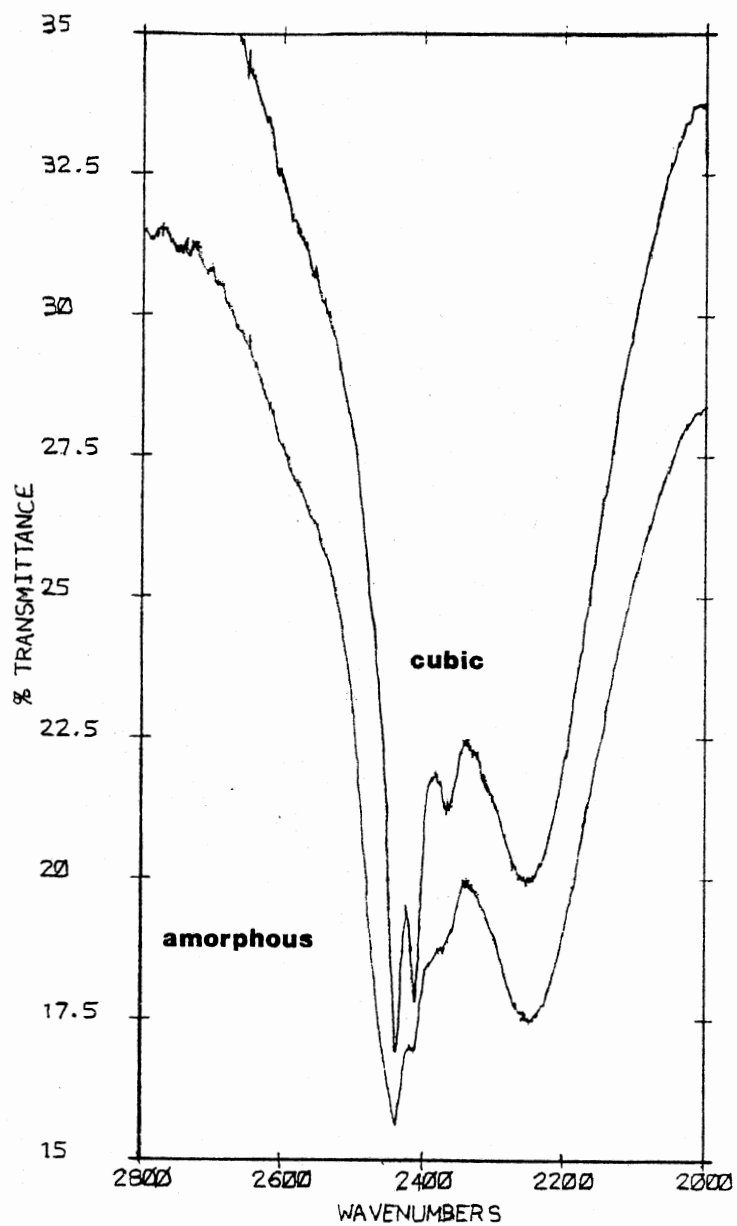


Figure 13. Spectral Changes Occurring between Amorphous and Cubic Decoupled D<sub>2</sub>O Ice

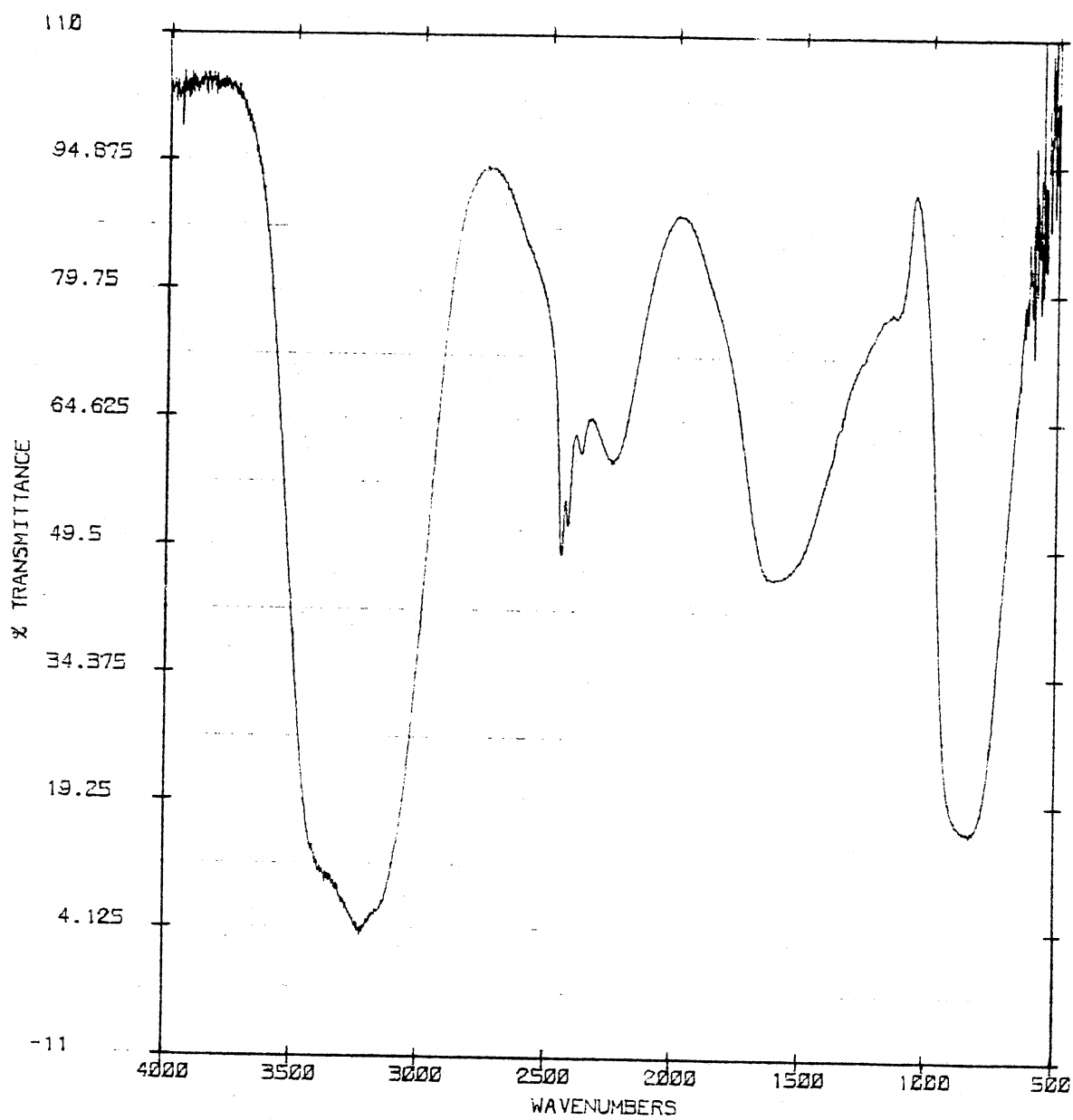


Figure 14. D<sub>2</sub>O Decoupled in H<sub>2</sub>O Cubic Ice Infrared Spectrum

sufficiently slow to warrant stopping. The collected interferograms were stored for later transformation to frequency spectra. The sample was then warmed to approximately 170 K for 10-20 minutes to completely convert the sample to HOD decoupled in H<sub>2</sub>O. Next the sample was cooled to the kinetic exchange temperature and a spectrum of HOD in H<sub>2</sub>O collected and stored. Lastly, the cryostat was reconnected to the vacuum and allowed to warm up to room temperature to vaporize and pump off the ice deposit. Thus with each sample a 130 K pre-kinetic run FTIR spectrum was made, a series of time correlated D<sub>2</sub>O/H<sub>2</sub>O exchange spectra at a chosen temperature, and a final converted HOD/H<sub>2</sub>O product spectra at the chosen reaction temperature.

The glass-cell 7-azaindole doped D<sub>2</sub>O in H<sub>2</sub>O deposits were prepared in a similar manner except that the deposit temperature was raised to 155 K to ensure crystallinity. This modification was used because doping inhibited any significant exchange at 155 K and affected the formation of a crystalline deposit at 130 K.

The collected glass-cell spectra were used to produce totally resolved D<sub>2</sub>O  $\nu_1$  and  $\nu_3$  bandshapes with no coupled HOD or isolated HOD band interference. These D<sub>2</sub>O bandshapes were made at 2 and 3 K intervals between 135 and 165 K. Subsequently they were used to resolve the D<sub>2</sub>O/HOD/H<sub>2</sub>O spectra collected with the metal cell.

Table I shows the collection parameters used for obtaining the kinetic data from the metal cell. The sampling time shows how long it took to collect the data for an interferogram, depending on the number of scans collected. Single beam spectra of the glass and metal cells cooled to approximately 150 K were used as background

TABLE I  
COLLECTION PARAMETERS USED FOR COLLECTING KINETIC DATA

Sample	Main Deposit Temp. ( $^{\circ}\text{K}$ )	Kinetic Temp. ( $^{\circ}\text{K}$ )	Resolution ( $\text{cm}^{-1}$ )	Scans Coadded	Sampling Time (sec)
2/2/82	130	150	2	20	50
6/7/82	130	150	2	10-20	25-50
6/9/82	125	150	2	10-20	25-50
1/29/82	130	145	2	20	50
5/10/82	130	145	2	20	50
5/24/82	130	145	2	10-20	25-50
6/3/82	130	145	2	20	50
5/11/82	130	140	2	20	50
5/25/82	130	140	2	20-50	50-120
5/15/82	130	135	2	20	50
5/26/82	130	135	2	20-50-100	50-120-240
6/11/82	125	135	2	20	50

references. Also spectra of these cells at 150 K with 20 minute  $H_2O$  deposits were used later in data manipulation.

The  $H_2O$  used was triply distilled and degassed under vacuum with several freeze-thaw cycles. The  $D_2O$  utilized was Merck and Co., Inc. deuterium oxide "100%" with a minimum isotopic purity of 99.96% deuterium. It was degassed under vacuum like  $H_2O$ . All chemical transfers were made with oven baked glassware under a dry nitrogen atmosphere. The 7-azaindole was obtained from Aldrich with a melting point of  $105-107^\circ$ , and IR and elemental analysis checks.

## CHAPTER III

### DATA TREATMENT AND THEORY DEVELOPMENT

#### Resolution of the Kinetic Spectral Data

The initial spectroscopic work on proton transfer in cubic ice deposits was reported by Devlin et al. (46, 47). As a result identification of the various vibrational bands in a typical  $D_2O/HOD/H_2O$  deposit spectrum was readily made. Figure 15 shows the 2600 to 2250  $cm^{-1}$  region of Figure 14 converted to absorbance units and expanded. Band for four distinct isotopic species are identifiable in the multiplet shown. The bands at 2444  $cm^{-1}$  and 2367  $cm^{-1}$  belong to  $\nu_3$  and  $\nu_1$  of  $D_2O$ . The strong band at 2418  $cm^{-1}$  is assigned to the OD stretch of uncoupled HOD. The weak shoulder at 2396  $cm^{-1}$  mentioned earlier is assigned to the in-phase coupled OD oscillator of coupled dimer HOD. Its companion out-of-phase OD stretch at 2442  $cm^{-1}$  is obscured beneath the  $D_2O$   $\nu_3$  band at 2444  $cm^{-1}$ . The baseline slant is due to the shoulder of the 2250  $cm^{-1}$   $H_2O$  combination band.

If the  $D_2O/HOD/H_2O$  deposit is warmed up and allowed to exchange the changes shown in Figure 16 occur. There the 2444 and 2367  $cm^{-1}$   $D_2O$  bands decrease in intensity and finally disappear, the 2396  $cm^{-1}$  coupled dimer HOD band remains steady or increases to subsequently disappear, and the decoupled or isolated HOD band at 2418  $cm^{-1}$  grows in intensity until eventually it becomes the sole spectral feature.

To follow the kinetics of the exchange process the multiplet

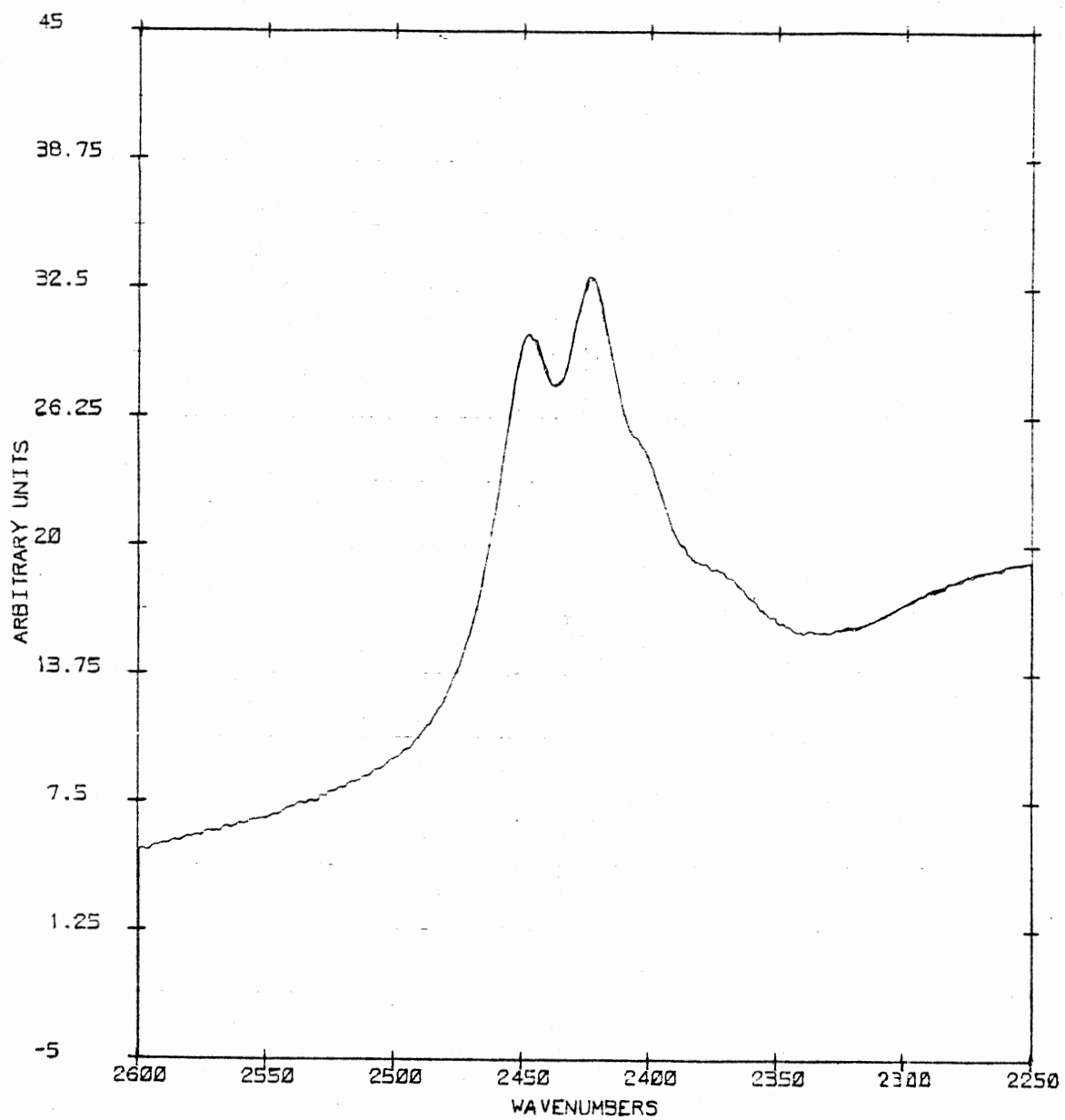


Figure 15. OD Stretch Multiplet of D<sub>2</sub>O Decoupled in H<sub>2</sub>O Ice

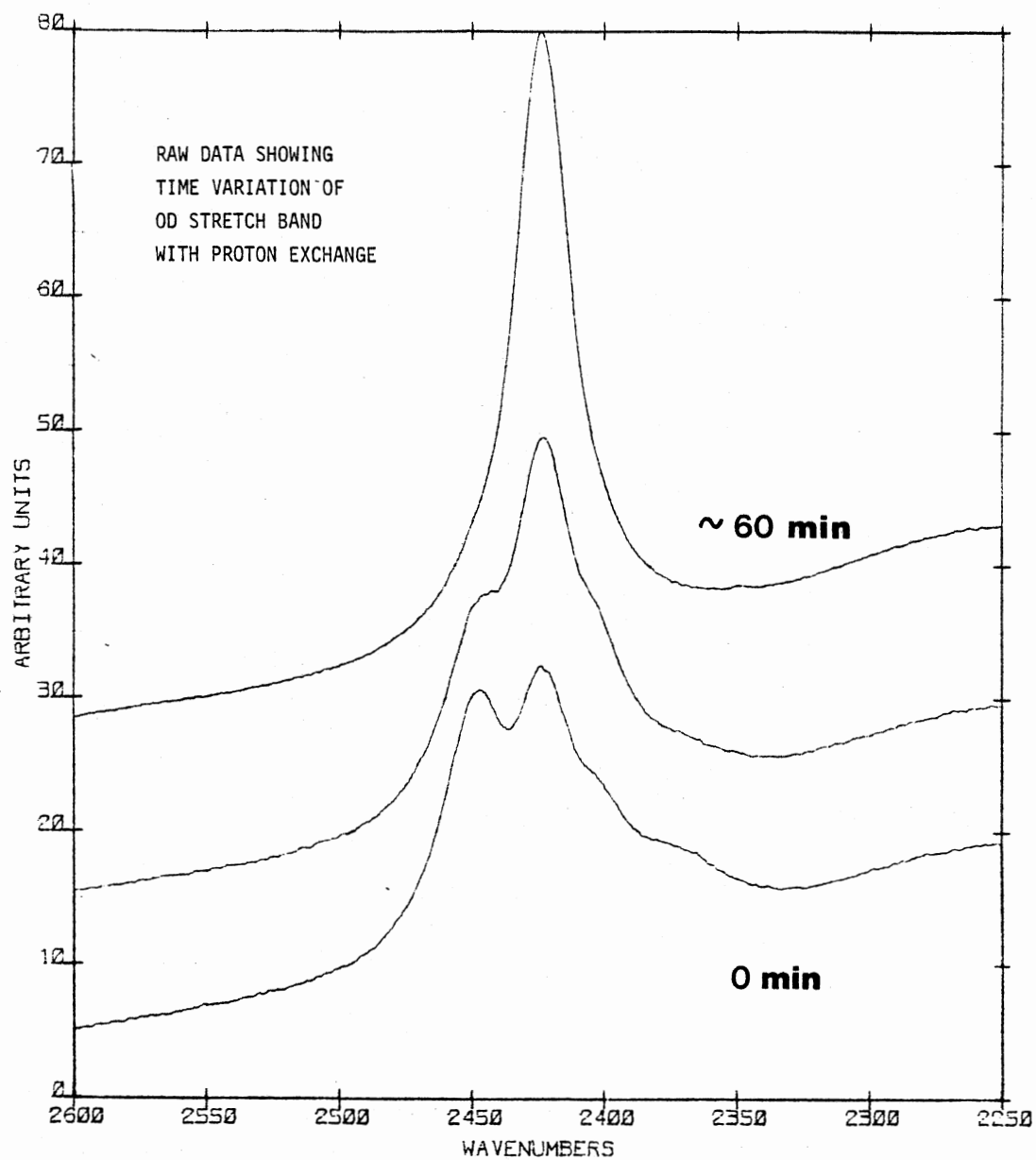


Figure 16. Time Variation of OD Stretch Multiplet with Proton Exchange



must be resolved into the bands for the individual kinetic species for absorbance measurement. For a mixed multiplet of four distinct species this can be done if three pure component species spectra are available. There each component spectra may be scaled appropriately and subtracted from the multiplet to leave a multiplet with one less component. The process may be repeated until only the unknown component is left.

If all four pure component spectra are available, a standard least square type fitting program can be used to determine the best fit scaling constants for each component spectrum. However, in resolving the OD stretch multiplet only three pure component spectra were available. The best route to extracting the component absorbances was to extract the absorbance for three available pure components and check to see if the remaining pure component bandshape agreed with predictions. Fortunately for the  $D_2O/HOD/H_2O$  deposit spectra, pure component spectra are obtainable for all except the coupled dimer HOD spectrum. In addition, from Haas and Hornig (22) some information is known about the coupled HOD bandshape. There should be two bands at about  $2393\text{ cm}^{-1}$  (in-phase stretch) and  $2442\text{ cm}^{-1}$  (out-of-phase stretch) with approximately a 2 to 1 (in-phase to out-of-phase) intensity ratio.

The  $2250\text{ cm}^{-1}$   $H_2O$  combination band pure component spectrum was obtained by making a deposit of  $H_2O$  only. The  $2418\text{ cm}^{-1}$  uncoupled HOD band was obtained by allowing the  $D_2O/HOD/H_2O$  deposit to convert completely to uncoupled HOD at the end of each kinetic run. Obtaining pure component  $D_2O$   $\nu_3$  and  $\nu_1$  bands at each of the temperatures used for kinetics was considerably more involved.

From Devlin et al. (46), it is known that by introducing a dopant

such as 7-azaindole into the ice samples the transfer rate  $k_1$  is slowed significantly relative to the  $k_2$  rotational rate, thus causing a very small steady state concentration of coupled HOD that is spectroscopically unobservable. Using the glass deposition cell with the glass cell procedure outlined in Chapter II, 7-azaindole doped  $D_2O/H_2O$  ice deposits were made. Curve A Figure 17 shows a typical spectrum of the  $2600-2250\text{ cm}^{-1}$  OD stretch region. Except for the absence of the  $2396$  and  $2442\text{ cm}^{-1}$  coupled HOD bands the features are identical to Figure 15. By warming the deposit sufficiently, the sample could be made to convert enough to uncoupled HOD, that (see Figure 17 curve B) the converted HOD spectrum could be scaled appropriately and subtracted from the  $D_2O/HOD$  multiplet shown in curve A. This left the pure component  $D_2O\ \nu_3$  and  $\nu_1$  bands at the appropriate temperature (curve C). Several doped  $D_2O/H_2O$  deposit spectra were collected at 2 and 3 K intervals from 135 to 160 K, then warmed and converted to uncoupled HOD. After cooling the resulting spectra were collected at the same temperature as the pre-converted  $D_2O/H_2O$  spectra. Thus at each temperature a  $D_2O/H_2O$  multiplet and uncoupled HOD spectra were available for producing a pure component  $D_2O\ \nu_3$  and  $\nu_1$  spectrum. Figure 18 shows these pure  $D_2O\ \nu_3$  and  $\nu_1$  spectra.

Now with three pure component spectra at the same temperature as the metal cell OD stretch multiplet it is possible to subtract out each component and find its respective absorbance until only the pure coupled HOD spectrum remained. The following procedure was used.

The  $2250\text{ cm}^{-1}$  combination band (Figure 19 curve B) is scaled and subtracted from the metal cell ice OD stretch multiplet (curve A) to

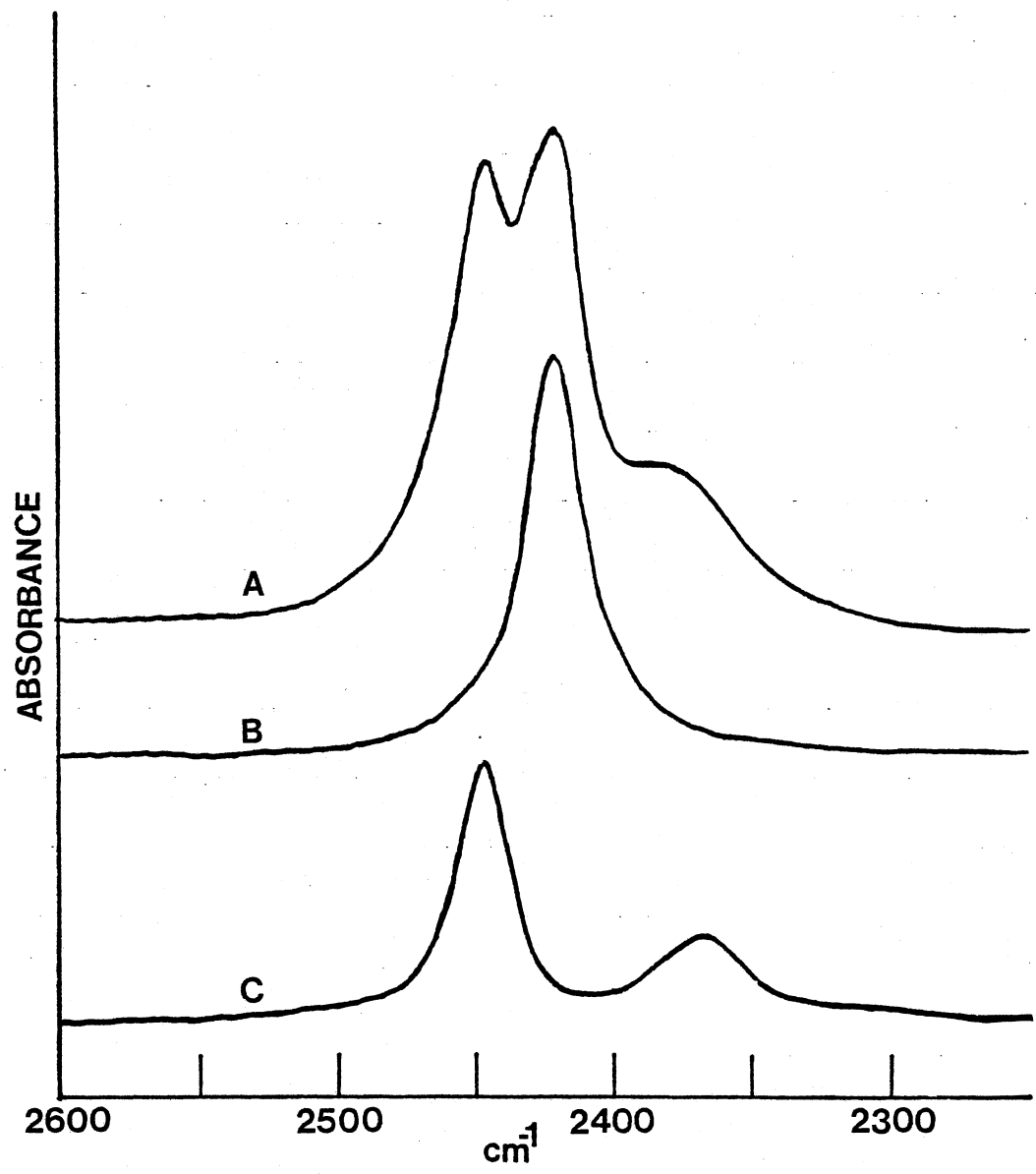


Figure 17. 7-Azaindole Doped  $\text{D}_2\text{O}/\text{H}_2\text{O}$  Ice Deposit Spectral Resolution

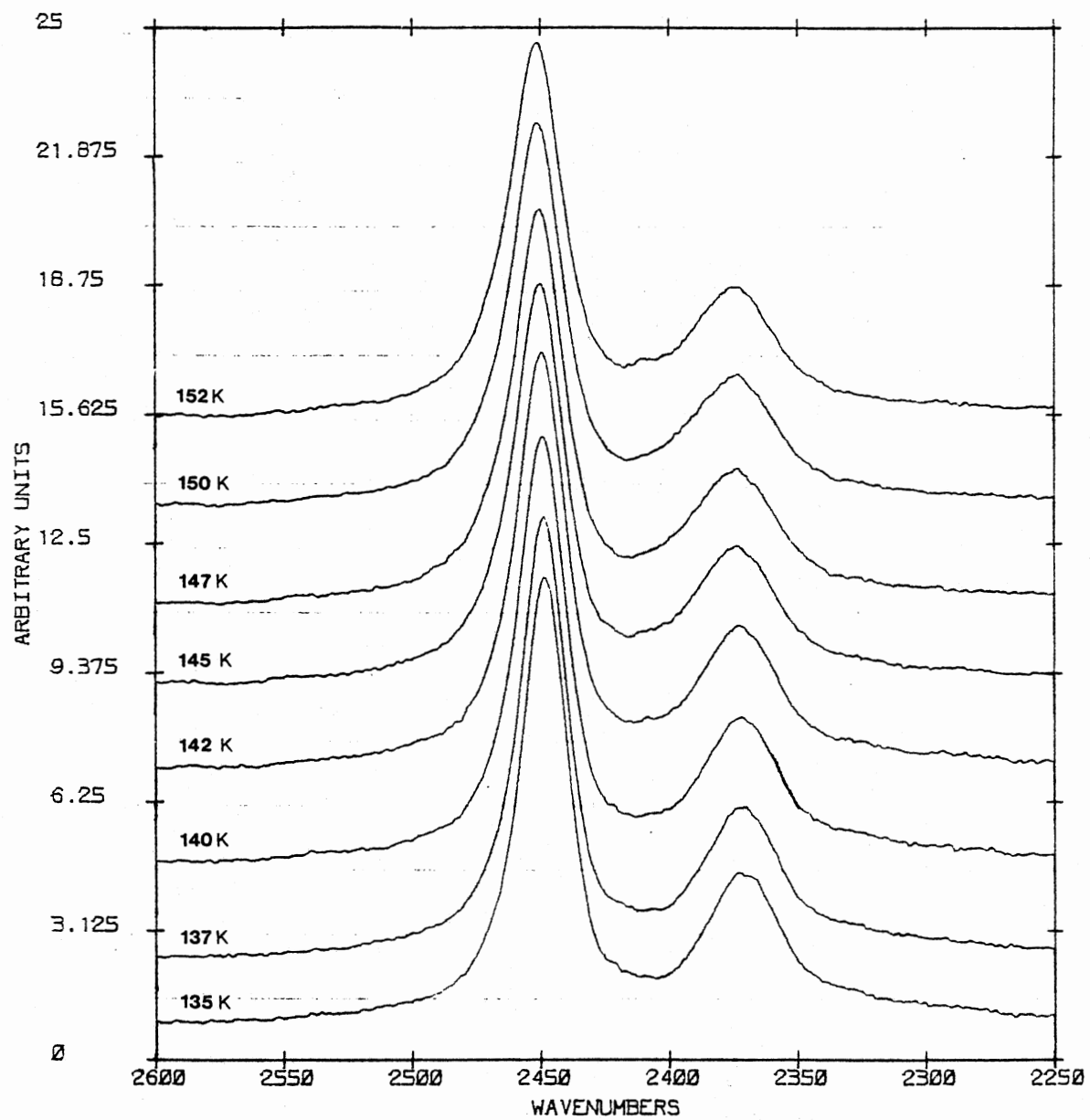


Figure 18. Pure Component  $D_2O$   $\nu_3$  Bandshapes from Glass Cell

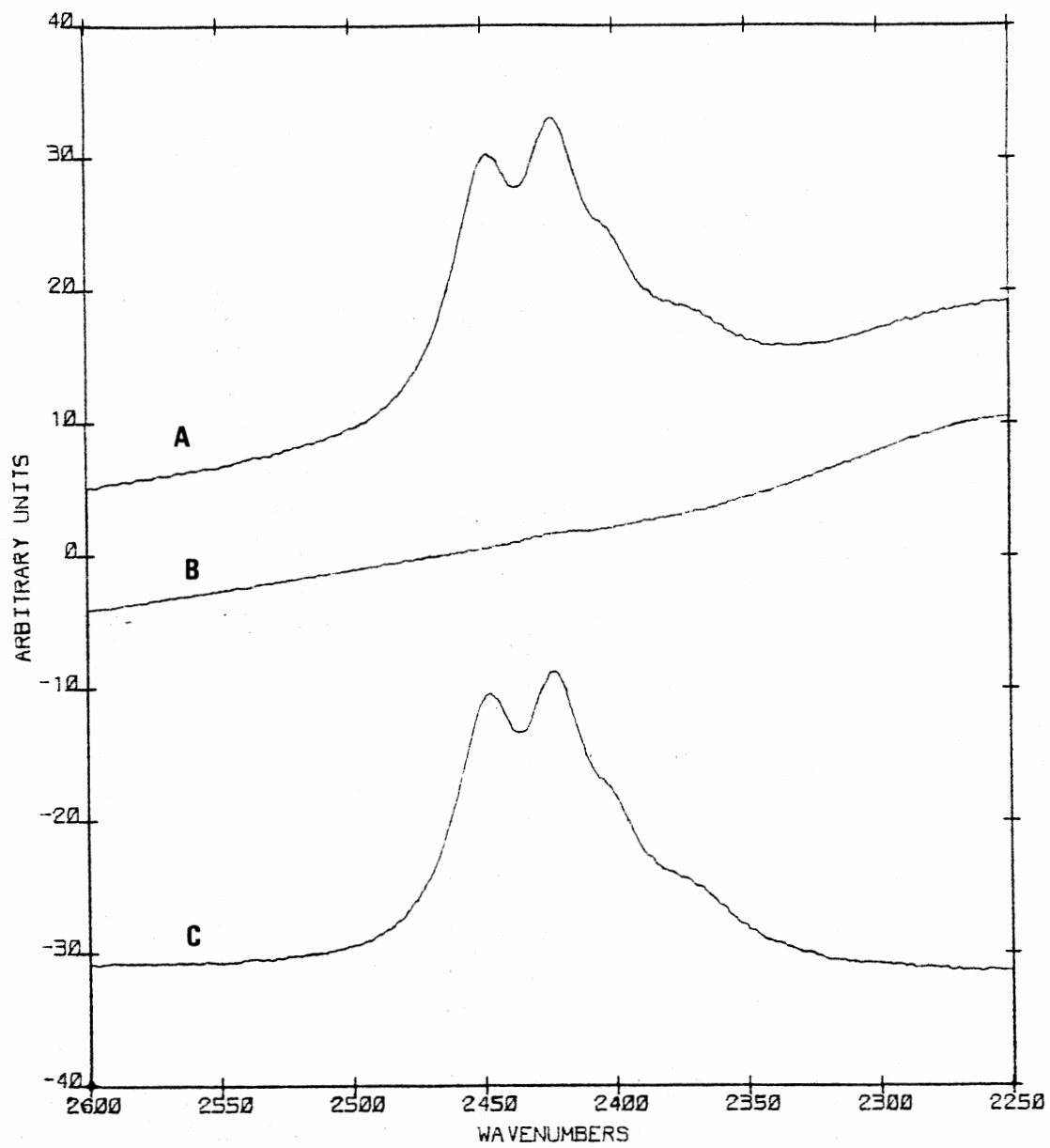


Figure 19. Subtraction of H<sub>2</sub>O Combination Shoulder from OD Multiplet

yield the multiplet on a flat background (curve C).

Next the pure  $D_2O$  spectrum (Figure 20 curve B), collected by the glass cell procedure at the same temperature as the OD stretch multiplet, was scaled and subtracted from the OD stretch multiplet (curve A) to yield the coupled and uncoupled HOD spectrum (curve C). Any significant over or under subtraction of the pure component  $D_2O$  (curve B) spectrum left small humps or dips at 2444 and 2367  $cm^{-1}$ . By multiplying the particular pure component  $D_2O$  absorbance by the scaling factor used, the absorbance of the  $D_2O$  concentration in the ice deposition sample was found. The absorbance of the pure component  $D_2O$  band was taken from the more intense  $\nu_3$  band.

The uncoupled 2418  $cm^{-1}$  HOD band (Figure 21 curve B) obtained at the end of each kinetic run was then scaled and subtracted from the previously found uncoupled and coupled HOD spectrum (Figure 21 curve A) to yield the pure component coupled HOD spectrum (curve C). The scaling factor used in the subtraction times the absorbance of the uncoupled HOD spectrum used (curve B) yielded the uncoupled HOD absorbance for that particular OD stretch multiplet. The absorbance of the coupled HOD was extracted from direct measurement of the in-phase 2396  $cm^{-1}$  band. Any over subtraction of uncoupled HOD (curve B) would reveal itself as too steep and negative a dip between the coupled HOD in and out-of-phase bands. Too little subtraction would show as insufficient dip or even a small residual hump between the two coupled HOD bands.

As a further check, the coupled and uncoupled spectrum (Figure 22 curve B) can be subtracted from the starting OD stretch multiplet (curve A) to yield a replica of the pure  $D_2O$  component bands used in the Figure 20

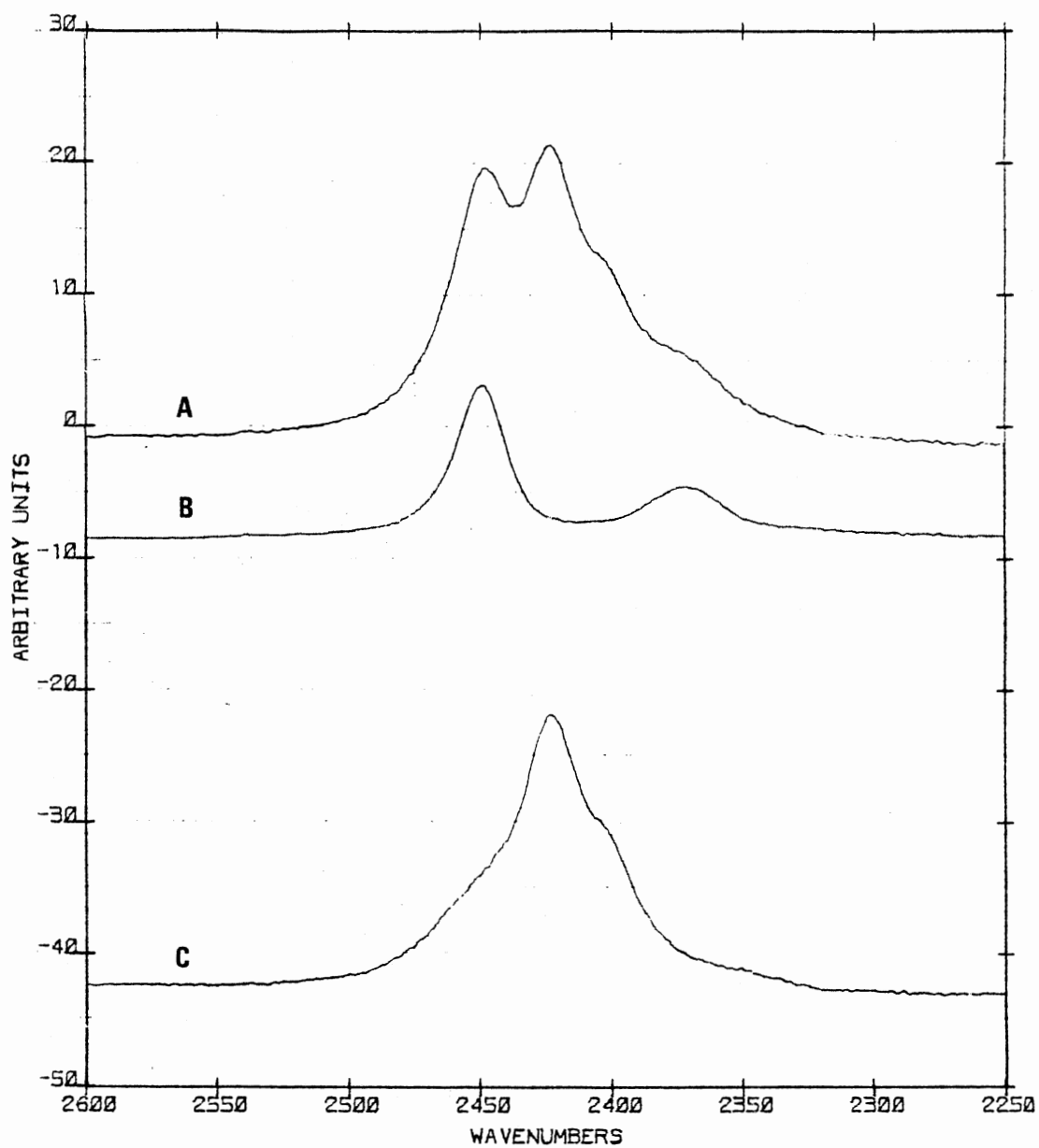


Figure 20. Subtraction of D<sub>2</sub>O Component from OD Multiplet

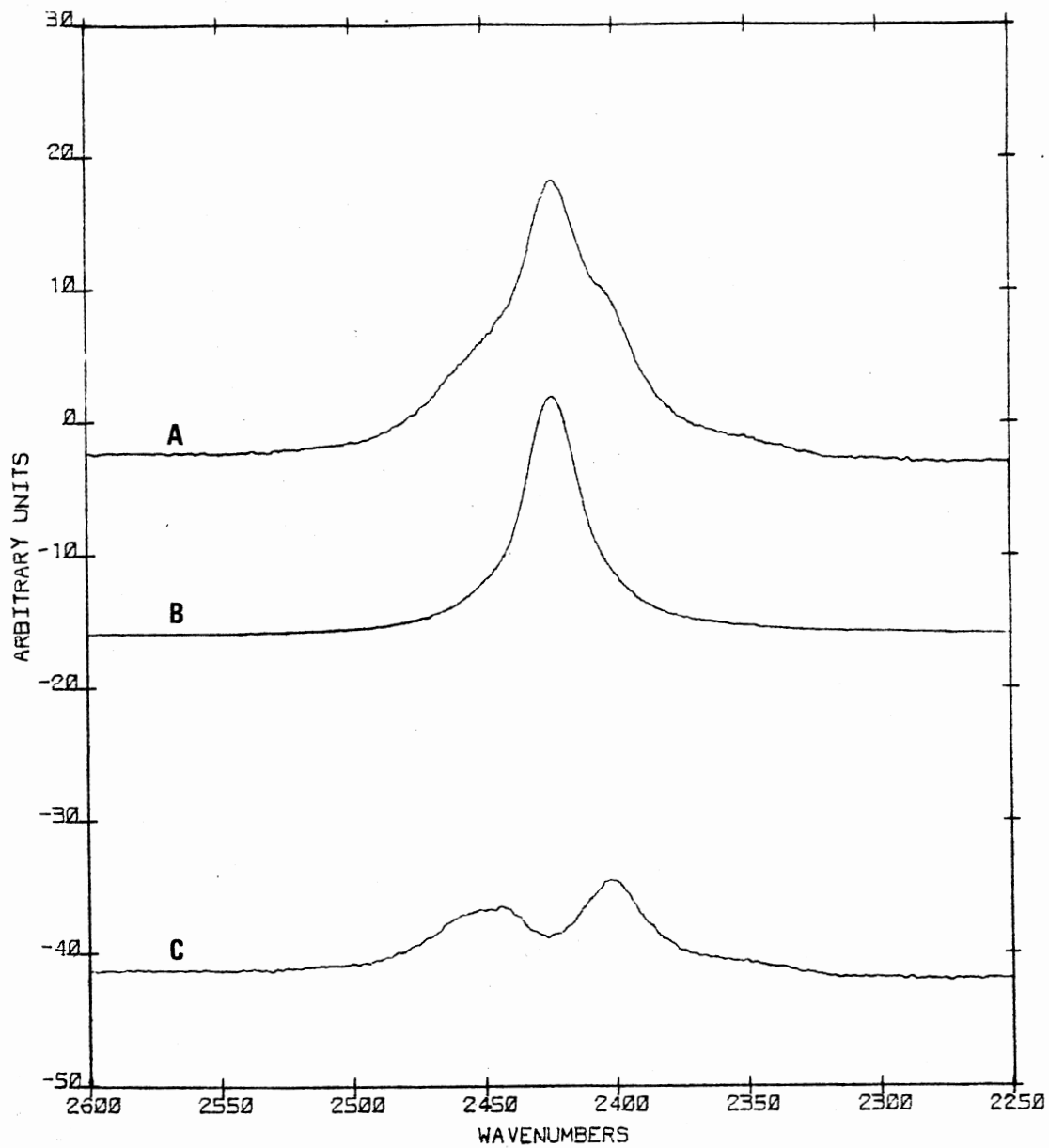


Figure 21. Subtraction of Uncoupled HOD Component from HOD Multiplet



subtraction, except that it has been scaled by the subtraction scaling factor used in the Figure 20 subtraction. This process was repeated for all multiplet resolutions at all data collection temperatures. Usually two or three adjustments of the subtraction scaling factors were necessary before the best resolution of the pure spectral components was achieved.

Since the  $D_2O$   $\nu_3$  band at  $2444\text{ cm}^{-1}$  and the coupled HOD out-of-phase band at  $2442\text{ cm}^{-1}$  overlap so completely the  $D_2O$  component scaling factor was not as correctly determined as the uncoupled HOD scaling factor. Figure 23 illustrates the effects of deviations of the  $D_2O$  pure component scaling factor from the correct scaling factor, on the coupled HOD bandshapes. Additionally, when the concentration and hence intensity of the coupled HOD was small, these deviations were difficult to detect. To circumvent this problem the following reasoning was used.

Both the in and out-of-phase bands arise from the same isotopic chemical species. Thus assuming Beer's law of absorption holds over the intensity range in question, there should be a constant ratio between the in-phase band and out-of-phase band irrespective of concentration or band intensity. To experimentally find this ratio, a series of the best most intense pure coupled HOD spectra were measured to find the best absorbance ratio. The best average ratio was 0.70 for the in-phase absorbance to out-of-phase absorbance. Multiplying the better determined in-phase coupled HOD band absorbance at  $2396\text{ cm}^{-1}$  by 0.70 should yield a better out-of-phase  $2442\text{ cm}^{-1}$  band absorbance. By subtracting this ratio determined absorbance from the spectrally measured coupled HOD out-of-phase absorbance the amount of over or

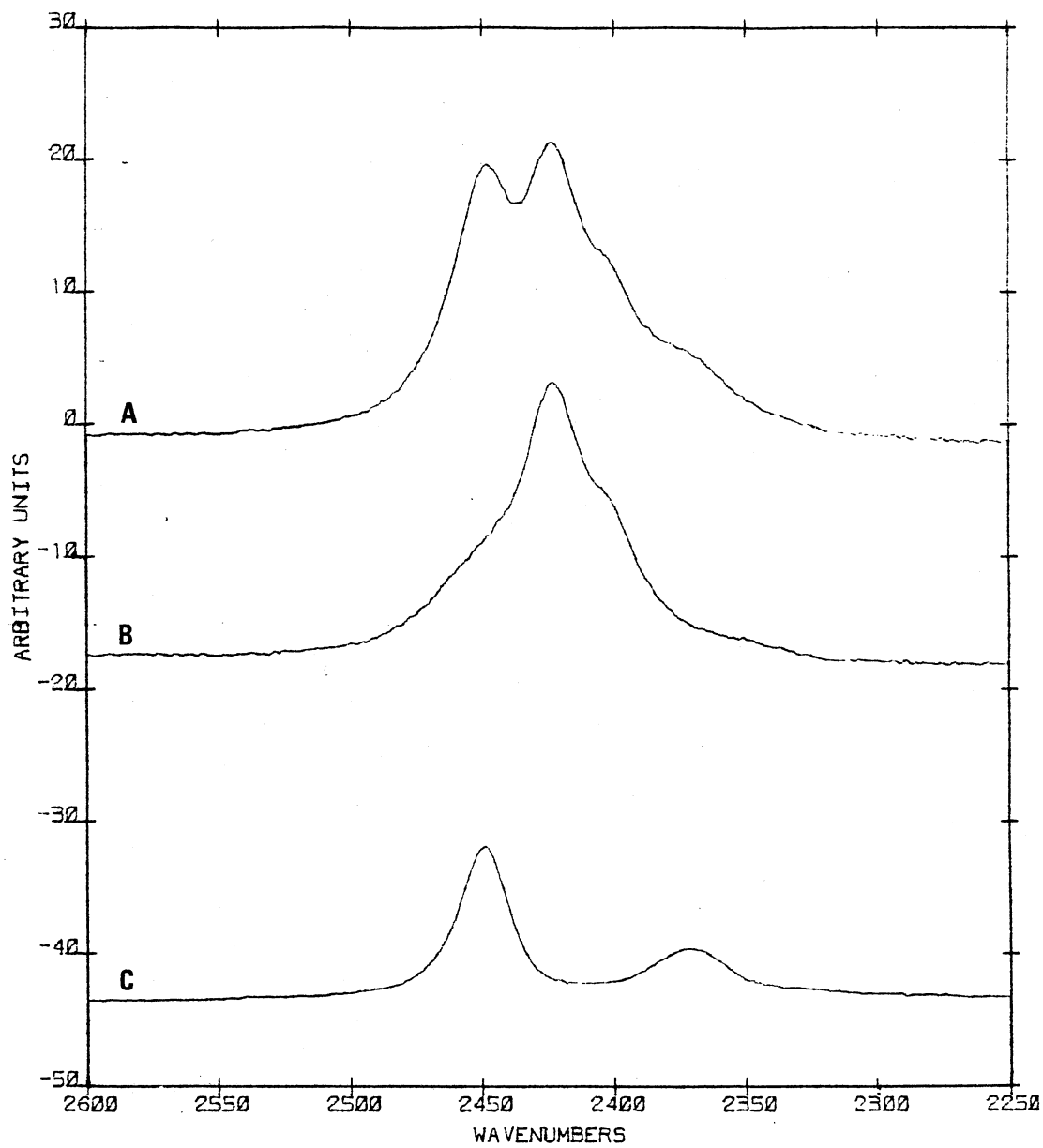


Figure 22. Subtraction of HOD Multiplet from OD Multiplet

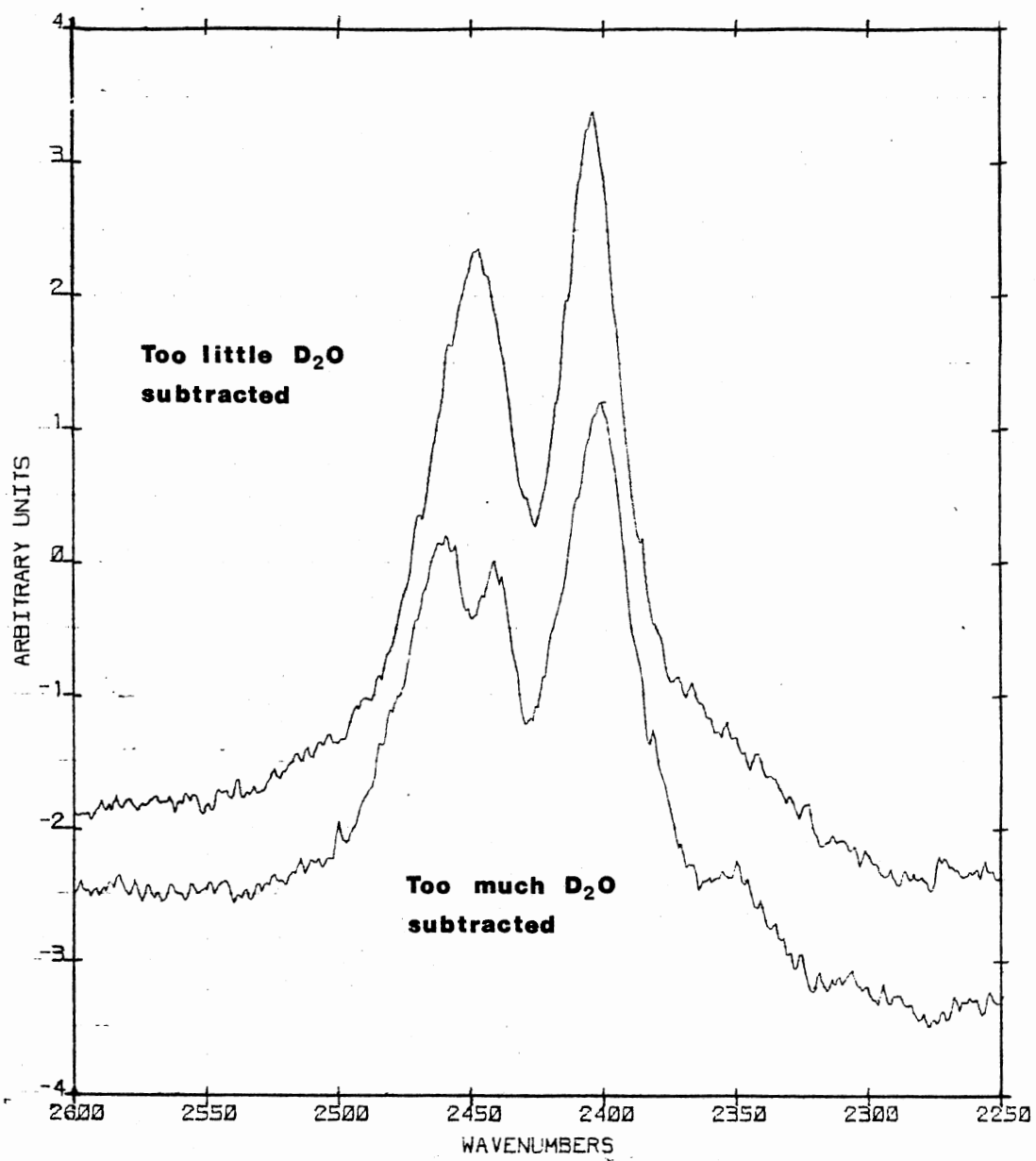


Figure 23. D<sub>2</sub>O Scaling Factor Deviation Effects on Coupled HOD Bandshapes

under subtraction of the pure component  $D_2O$  band from the OD stretch multiplet could be determined. This positive or negative deviation was added to the previously determined  $D_2O$  component absorbance for the OD stretch multiplet being resolved. The resulting  $D_2O$  absorbances are called corrected  $D_2O$  absorbances. Most deviations calculated were small and served to smooth out the data of the corrected  $D_2O$  absorbance values as compared to the uncorrected  $D_2O$  absorbance values. Table II lists the scaling factors, corrected  $D_2O$  absorbance, coupled HOD absorbance, uncoupled HOD absorbances for the kinetic data found at 135, 140, 145, and 150 K. Also listed are the unscaled absorbances of the fully converted uncoupled HOD bands used in subtracting out the uncoupled HOD (see Figure 21) and glass cell doped ice deposit determined pure  $D_2O$   $\nu_3$  bands (see Figure 20). All absorbance measurements determined from spectral plots were found by subtracting from the band peak absorbance the peak frequency absorbance of a baseline drawn tangest to the background absorption. All absorbance measurements are in Digilab's arbitrary absorbance units.

Table I shows the sampling times for the Table II data. The sampling or collection time for the interferogram is dependent on the resolution required, and number of interferometer scans coadded to form the final interferogram. For most data points the period of collection is small compared to the time interval between data points. The exceptions are the initial 150 K collected data. If the 2 to 4 minute interval of the 6/7/82 dataset is examined, it is found that the greatest change occurring with the  $D_2O$  absorbances are 12.639 at 2 minutes and 10.155 at 4 minutes. The Digilab FTIR takes 25 seconds to collect a 10 scan  $2\text{ cm}^{-1}$  resolution interferogram. Thus one quarter

TABLE II

ABSORBANCES AND SCALING FACTORS FOR RESOLVED OD STRETCH MULTIPLETS

Dataset	Time (minutes)	D <sub>2</sub> O		(HOD) <sub>2</sub>	HOD		
		Scaling Factor	Absorbance <sup>a,b</sup>	Absorbance <sup>a</sup>	Scaling Factor	Absorbance <sup>a</sup>	
		<u>150 K</u>					
2/2/82		Doped D <sub>2</sub> O absorbance = 9.20 <sup>c</sup>		Fully converted HOD absorbance = 36.60 <sup>d</sup>			
	0.0	1.27	11.869 <sup>e</sup>	3.56 <sup>e</sup>	0.31	11.346 <sup>e</sup>	
	2.0	0.98	9.281	5.05	0.35	12.81	
	4.0	0.80	7.24	5.6	0.42	15.372	
	6.0	0.65	5.665	5.45	0.49	17.934	
	9.5	0.45	3.885	4.65	0.60	21.96	
	12.5	0.33	3.041	3.85	0.67	24.522	
	15.5	0.27	2.379	3.15	0.73	26.718	
	20.5	0.19	1.873	2.25	0.79	28.914	
6/7/82		Doped D <sub>2</sub> O absorbance = 9.20		Fully converted HOD absorbance = 38.30			
	0.0	1.5	14.295	3.15	0.29	11.107	
	2.0	1.35	12.639	3.83	0.335	12.831	
	4.0	1.1	10.155	5.15	0.37	14.171	
	6.0	0.88	8.131	5.55	0.43	16.469	
	8.0	0.72	6.646	5.74	0.49	18.767	
	10.0	0.62	5.809	5.35	0.54	20.682	
	12.0	0.56	5.121	4.83	0.59	22.597	
	15.0	0.49	4.253	4.15	0.655	25.087	
	20.0	0.40	3.415	3.45	0.72	27.576	
	25.0	0.32	2.714	2.9	0.72	29.491	



TABLE II (Continued)

Dataset	Time (minutes)	D <sub>2</sub> O		(HOD) <sub>2</sub>	HOD	
		Scaling Factor	Absorbance <sup>a,b</sup>	Absorbance <sup>a</sup>	Scaling Factor	Absorbance <sup>a</sup>
	43	0.30	3.07	3.9	0.69	26.22
	52	0.26	2.59	3.4	0.73	27.74
	62	0.22	2.23	2.8	0.77	29.26
	72	0.18	2.125	2.55	0.79	30.02
5/10/82		Doped D <sub>2</sub> O absorbance = 9.50		Fully converted HOD absorbance = 35.80		
	0.0	1.1	10.367	3.69	0.36	12.888
	3.0	1.0	9.537	4.09	0.38	13.604
	6.0	0.93	8.805	4.4	0.40	14.32
	9.0	0.85	7.95	4.65	0.43	15.394
	12.0	0.79	7.675	4.5	0.445	15.931
	17.0	0.70	6.609	4.63	0.485	17.363
	22.0	0.63	5.947	4.54	0.52	18.616
	27.0	0.58	5.435	4.25	0.555	19.869
	32.0	0.54	4.821	4.07	0.59	21.122
	47.0	0.41	3.561	3.32	0.67	23.986
	62.0	0.29	2.895	2.8	0.72	25.776
	77.0	0.22	2.245	2.35	0.77	27.566
5/24/82		Doped D <sub>2</sub> O absorbance = 9.5		Fully converted HOD absorbance = 33.25		
	0.0	1.0	9.25	3.5	0.32	10.64
	3.0	0.88	8.117	3.89	0.34	11.305
	6.0	0.76	6.771	4.57	0.37	12.303
	9.0	0.64	5.925	4.65	0.40	13.3
	12.0	0.55	5.038	4.81	0.44	14.63
	17.0	0.40	3.809	4.83	0.50	10.625

TABLE II (Continued)

Dataset	Time (minutes)	D <sub>2</sub> O		(HOD) <sub>2</sub>	HOD	
		Scaling Factor	Absorbance <sup>a,b</sup>	Absorbance <sup>a</sup>	Scaling Factor	Absorbance <sup>a</sup>
	22.0	0.32	3.141	4.57	0.54	17.955
	27.0	0.29	2.997	4.14	0.57	18.953
	37.0	0.25	2.276	3.47	0.64	21.28
	47.0	0.22	1.935	3.05	0.685	22.776
6/3/82	Doped D <sub>2</sub> O absorbance = 9.50      Fully converted HOD absorbance = 34.00					
	0.0	1.25	12.13	1.85	0.365	12.41
	3.0	1.1	10.481	2.67	0.385	13.09
	6.0	1.0	9.456	3.52	0.405	13.77
	9.0	0.915	8.737	3.68	0.425	14.45
	12.0	0.85	8.139	3.98	0.445	15.13
	17.0	0.75	7.325	4.1	0.475	16.15
	22.0	0.68	6.646	4.02	0.51	17.34
	27.0	0.63	5.764	4.03	0.55	18.7
	32.0	0.59	5.315	3.7	0.585	19.89
	42.0	0.50	4.459	3.43	0.64	21.76
	52.0	0.41	3.755	3.1	0.69	23.46
	62.0	0.32	3.18	2.8	0.72	24.48
			<u>140 K</u>			
5/11/82	Doped D <sub>2</sub> O absorbance = 10.20      Fully converted HOD absorbance = 33.50					
	0.0	1.2	11.788	2.26	0.37	12.395
	3.0	1.05	10.843	2.81	0.39	13.065
	6.0	0.97	10.195	3.17	0.40	13.4
	9.0	0.93	9.636	3.5	0.405	13.568



TABLE II (Continued)

Dataset	Time (minutes)	D <sub>2</sub> O		(HOD) <sub>2</sub>	HOD	
		Scaling Factor	Absorbance <sup>a,b</sup>	Absorbance <sup>a</sup>	Scaling Factor	Absorbance <sup>a</sup>
	12.0	0.90	9.44	3.7	0.41	13.735
	17.0	0.85	8.921	3.97	0.42	14.07
	22.0	0.81	8.396	4.28	0.43	14.405
	27.0	0.775	7.969	4.38	0.445	14.908
	32.0	0.745	7.756	4.39	0.455	15.243
	42.0	0.69	7.008	4.4	0.48	16.08
	60.0	0.59	6.025	4.49	0.53	17.755
	80.0	0.50	5.205	4.05	0.58	19.43
	100.0	0.42	4.554	3.8	0.62	20.77
	119.8	0.35	4.03	3.6	0.65	21.775
5/25/82		Doped D <sub>2</sub> O absorbance = 10.20		Fully converted HOD absorbance = 37.60		
	0.0	1.41	14.384	2.54	0.31	11.656
	3.0	1.37	14.324	2.50	0.33	12.408
	6.0	1.33	13.951	2.85	0.33	12.408
	9.0	1.29	13.288	3.4	0.33	12.408
	12.0	1.25	12.937	3.59	0.34	12.784
	17.0	1.17	12.269	3.95	0.355	13.348
	27.0	1.1	11.306	4.52	0.37	13.912
	37.0	1.0	10.311	4.77	0.40	15.04
	47.0	0.92	9.797	4.91	0.42	15.792
	62.0	0.80	8.72	5.1	0.45	16.92
	82.0	0.71	7.587	5.15	0.49	18.424
	102.25	0.63	6.826	5.0	0.53	19.928
	132.0	0.54	5.839	4.77	0.57	21.432

TABLE II (Continued)

Dataset	Time (minutes)	D <sub>2</sub> O		(HOD) <sub>2</sub>	HOD	
		Scaling Factor	Absorbance <sup>a,b</sup>	Absorbance <sup>a</sup>	Scaling Factor	Absorbance <sup>a</sup>
<u>135 K</u>						
5/15/82		Doped D <sub>2</sub> O absorbance = 10.70		Fully converted HOD absorbance = 38.00		
	0.1	1.29	14.148	2.65	0.327	12.426
	3.0	1.27	13.806	2.79	0.329	12.502
	6.0	1.25	13.681	2.92	0.331	12.578
	9.0	1.23	13.356	3.15	0.733	12.654
	14.3	1.20	13.152	3.34	0.335	12.73
	24.0	1.15	12.635	3.7	0.34	12.92
	34.0	1.09	12.093	4.1	0.345	13.11
	49.1	1.04	11.571	4.51	0.35	13.3
	64.0	0.98	10.886	4.8	0.36	13.68
	84.0	0.91	10.397	5.1	0.37	14.06
	104.0	0.88	9.906	5.3	0.38	14.44
	124.0	0.85	9.573	5.46	0.39	14.82
	144.0	0.82	9.134	5.7	0.40	15.2
	174.0	0.78	8.721	5.65	0.42	15.96
5/26/82		Doped D <sub>2</sub> O absorbance = 10.70		Fully converted HOD absorbance = 23.65		
	0.0	0.74	8.155	1.49	0.356	8.419
	6.0	0.725	7.918	1.6	0.359	8.490
	11.0	0.71	7.714	1.69	0.362	8.561
	16.0	0.695	7.471	1.88	0.365	8.632
	25.0	0.68	7.265	2.13	0.368	8.703
	35.0	0.65	7.049	2.18	0.375	8.869
	55.0	0.60	6.536	2.42	0.39	9.224

TABLE II (Continued)

Dataset	Time (minutes)	D <sub>2</sub> O		(HOD) <sub>2</sub>	HOD	
		Scaling Factor	Absorbance <sup>a,b</sup>	Absorbance <sup>a</sup>	Scaling Factor	Absorbance <sup>a</sup>
	85.0	0.55	6.061	2.72	0.405	9.578
	115.0	0.51	5.68	2.91	0.42	9.933
	145.0	0.48	5.391	3.05	0.435	10.288
	185.0	0.45	4.981	3.12	0.45	10.643
	225.0	0.435	4.685	3.2	0.465	10.997
	285.0	0.405	4.208	3.18	0.49	11.589
6/11/82		Doped D <sub>2</sub> O absorbance = 10.7		Fully converted HOD absorbance = 35.50		
	0.0	1.28	14.053	2.09	0.320	11.36
	5.0	1.22	13.556	2.14	0.324	11.502
	10.0	1.20	12.99	2.5	0.326	11.573
	20.0	1.15	12.675	2.8	0.328	11.644
	30.0	1.10	12.368	3.26	0.33	11.715
	40.0	1.03	11.781	3.6	0.335	11.893
	60.0	0.98	11.166	4.0	0.35	12.425
	80.0	0.95	10.785	4.3	0.365	12.958
	100.0	0.93	10.266	4.55	0.37	13.135
	130.0	0.89	9.528	4.85	0.39	13.845
	160.0	0.85	9.011	5.12	0.41	14.555
	190.0	0.80	8.265	5.15	0.43	15.265
	220.0	0.75	7.815	5.3	0.45	15.975

<sup>a</sup> Digilab arbitrary absorbance units.

<sup>b</sup> Corrected D<sub>2</sub>O concentrations as described in text.

<sup>c</sup> D<sub>2</sub>O spectra obtained from glass cell ice deposits.

<sup>d</sup> Final converted HOD spectrum obtained at the end of each kinetic run.

<sup>e</sup> Data is accurate to 2 to 3 significant figures.

of the time interval was spent collecting data during which the  $D_2O$  absorbance was decreasing slightly. Since the ten scans are averaged, the interferometer would obtain a measured absorbance slightly lower for the 2 minute mark where the data collection began, than the true absorbance one would obtain for an instantaneous measurement at the 2 minute mark. If a linear decrease for the absorbance is assumed over the 25 second time interval, then an approximate average interferometer measured absorbance would be measured at about 13 seconds into the time interval, or at 2.23 minutes instead of 2 minutes. Since the absorbance changes by 2.484 (Digilab arbitrary absorbance units) over the 2 minute interval, then over 13 seconds the absorbance should change by approximately 0.31 absorbance units. This gives rise to approximately a 2.5% relative error in the two minute absorbance measurement of 12.639. In view of this small error, that decreases dramatically as the time intervals increase and rate of absorbance change decreases, the sampling error was ignored.

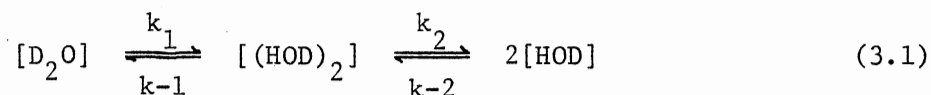
#### Determination of Relative Molar Absorptivities

Kinetic analysis is based on concentrations changes of chemical species. Therefore some method must be found for converting the absorbance measurements into concentration units before a kinetic mechanism may be proposed for isotopic exchange in cubic ice. To accomplish this, Beer's law is assumed to hold over the dynamic range of all kinetic absorbance measurements. All absorbance measurements, though recorded in arbitrary absorbance units, fell between 0.10 and 0.5 absorbance on a true absorbance scale; a range where Beer's law deviations are not

severe.

Due to the thin layer nature of the ice deposits and the method used in their preparation it is very difficult if not impossible to determine thickness, thickness uniformity, and absolute concentrations of the various isotopic species. Thus there was no means available for preparing standard Beer's law plots of each isotopic specie isolated in H<sub>2</sub>O cubic ice. Neither does the chemical literature contain reports of absorptivities of D<sub>2</sub>O, HOD, or (HOD)<sub>2</sub> isolated in a cubic ice lattice at the necessary frequencies. But fortunately since first order rate constants are dependent on the change in concentrations rather than absolute concentrations, relative concentrations and hence relative molar absorptivities are sufficient. If the relative magnitudes of the absorptivities are found, then the relative concentrations can be found from the absorbance data.

The kinetic expression of the isotopic exchange is



where [D<sub>2</sub>O] is the D<sub>2</sub>O concentration in moles/liter of the ice deposit, [(HOD)<sub>2</sub>] the coupled HOD dimer concentration, and [HOD] the uncoupled HOD concentration. From experimental data, (46, 47) it is known that the uncoupled HOD units in the cubic ice lattice eventually migrate apart. Therefore the equilibrium situation proposed in equation (3.1) is constantly perturbed until after an infinite period of time, essentially all the D<sub>2</sub>O and (HOD)<sub>2</sub> have converted to HOD. Thus at any time a material balance yields

$$[D_2O] + [(HOD)_2] = \frac{1}{2}[HOD_\infty] - \frac{1}{2}[HOD] \quad (3.2)$$

where  $[HOD_\infty]$  is the uncoupled HOD concentration after an infinite amount of time and the other concentrations at some lesser time.

Rearranging equation (3.2) gives

$$2[D_2O] + 2[(HOD)_2] + [HOD] - [HOD_\infty] = 0 \quad (3.3)$$

assuming Beer's law holds where  $[X] = \text{peak absorbance of } X/\epsilon_x b$ ,  $\epsilon_x$  is the molar absorptivity of component  $X$  and  $b$  the pathlength traversed in penetrating the sample medium then

$$\frac{2A_{D_2O}}{\epsilon_{D_2O} b} + \frac{2A_{(HOD)_2}}{\epsilon_{(HOD)_2} b} + \frac{A_{HOD}}{\epsilon_{HOD} b} - \frac{A_{HOD_\infty}}{\epsilon_{HOD} b} = 0 \quad (3.4)$$

$$\frac{2A_{D_2O}}{\epsilon_{D_2O}} + \frac{2A_{(HOD)_2}}{\epsilon_{(HOD)_2}} + \frac{A_{HOD} - A_{HOD_\infty}}{\epsilon_{HOD}} = 0 \quad (3.5)$$

By inserting absorbance data from three different time intervals three different equations may be constructed to solve for the three molar absorptivities. However these equations would form a homogeneous set of equations for which no unique solution exists; only a relational one. Thus to find the relational solution it is assumed that  $\epsilon_{HOD} = 1.0$ , and  $a = \frac{2}{\epsilon_{D_2O}}$ , and  $b = \frac{2}{\epsilon_{(HOD)_2}}$ . Rewriting, we obtain

$$aA_{D_2O} + bA_{(HOD)_2} = A_{HOD_\infty} - A_{HOD} \quad (3.6)$$

Then with absorbance data from two different time intervals of a particular kinetic run the relative molar absorptivities of  $D_2O$ ,  $(HOD)_2$ , and HOD may be found. Alternatively  $\epsilon_{(HOD)_2}$  or  $\epsilon_{D_2O}$  could be set equal to some constant and the relational solution then found.

REMOLE is a fortran program written to ascertain the relative molar absorptivities of cubic ice-isolated  $D_2O$ ,  $(HOD)_2$ , and HOD. The program uses equation 3.6 and the equivalent equation found by setting  $\epsilon_{HOD} =$

1.0 plus inputted absorbance data for  $D_2O$ ,  $(HOD)_2$ , and HOD to determine the best average relative molar absorptivities for each. A program listing with sample data and output is included in Appendix A.

Figure 24 shows a 150 K set of kinetic observations. REMOLE required that the observations of a dataset be numbered consecutively with integers starting from one. For example, observation 1 would contain the absorbances (Digilab arbitrary absorbance units) of  $D_2O$ ,  $(HOD)_2$ , and HOD at time = 0 and observation 2 the absorbances at time = 2 minutes. Then by inputting the starting and ending observation numbers REMOLE would calculate the average relative molar absorptivities for all unique binary combinations of observations between and including the starting and ending observations using simultaneous equations from equation (3.6). As a check, equation (3.6) was solved with both  $\epsilon_{HOD} = 1.0$ ; and its appropriate variant equation where  $\epsilon_{(HOD)_2} = 1.0$ . The relative molar absorptivities obtained from setting  $\epsilon_{(HOD)_2} = 1.0$  could be multiplied by a scaling factor for direct comparison with the molar absorptivities obtained from equation (3.6) with  $\epsilon_{HOD} = 1.0$ . As illustrated in Figure 24 observation 10 is at the end of the dataset where the arbitrary absorbance values of HOD are high and the values of  $D_2O$  and  $(HOD)_2$  quite low. Thus a 0.5 arbitrary absorbance error would be about a 15-20% error at observation 10 for  $D_2O$  and  $(HOD)_2$  but only about a 5% error for  $D_2O$  at lower observations. Therefore it is expected that binary combinations of lower numbered observations should give more reliable consistent absorptivities. REMOLE prints out the relative absorptivities for each binary combination of absorbances used in the simultaneous equation solution so that visual

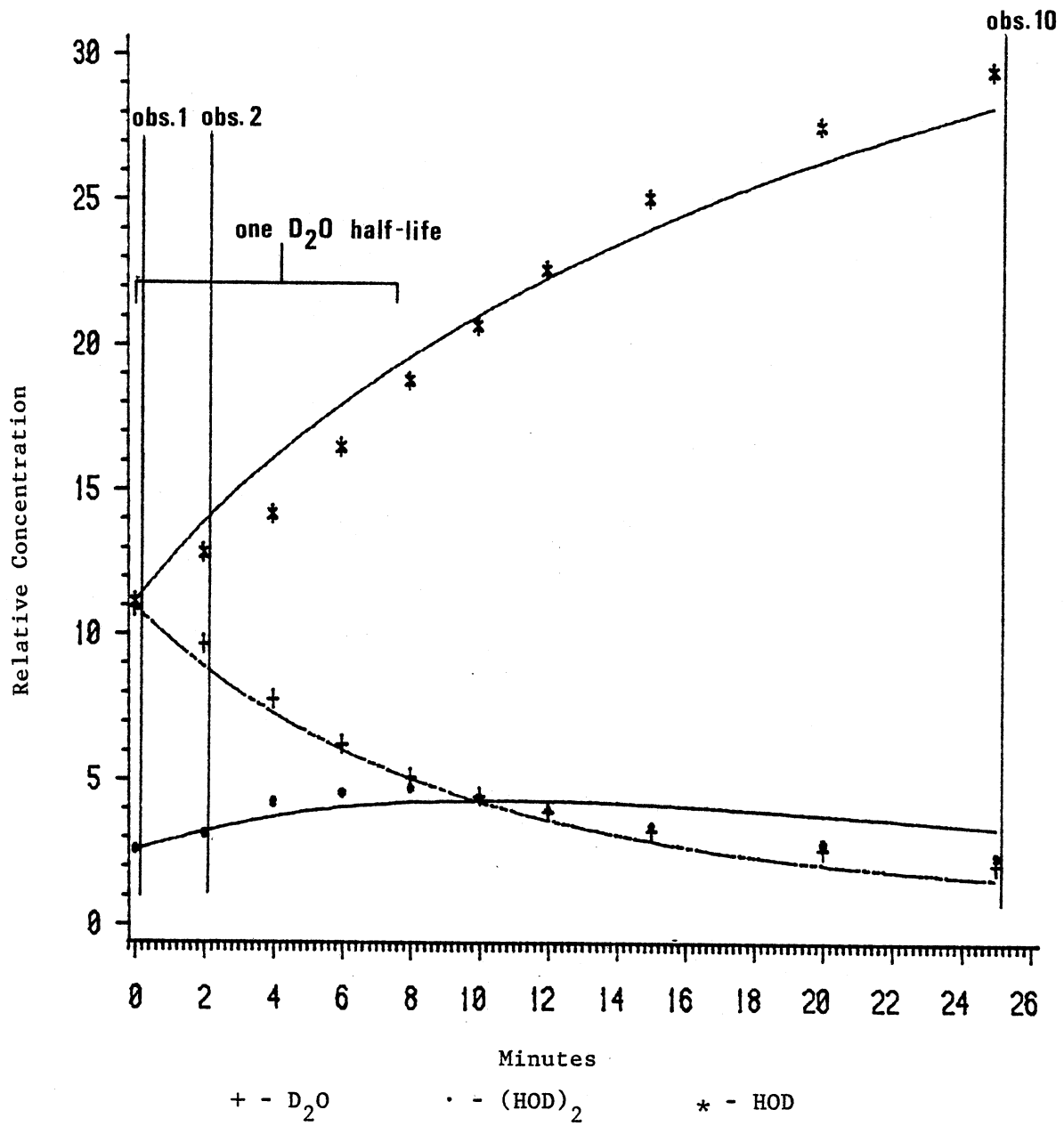


Figure 24. REMOLE Dataset Coverage and Analysis



inspection of the individual computed absorptivities was possible. An additional feature of REMOLE was to decrement the stopping observation by one and then recompute the average relative molar absorptivities. This process was repeated until only the starting and starting plus one observation absorbances were left for finding the relative molar absorptivities. By appropriate selection of starting and stopping observation numbers it is possible to confine the calculation to the most reliable part of the dataset. From calculations using REMOLE the best absorbances to use for relative-molar-absorptivity calculations were those within the first 1.5 half-lives of the  $D_2O$  absorbance.

Table III shows the results of these calculations for each dataset at each temperature. Included are the relative molar absorptivities for each dataset for calculations with  $\epsilon_{HOD}$  set equal to 1.0 and also for  $\epsilon'_{(HOD)_2}$  set equal to 1.0. Unless the  $\epsilon'$  values are multiplied by a  $\epsilon_{HOD}$  scaling factor, the two calculations cannot be directly compared. The average relative molar absorptivities for datasets and temperatures with their associated sigmas is shown at the bottom of Table III. The primed average absorptivities have been scaled by  $\epsilon_{(HOD)_2} = 1.222$  so that the different calculations may be compared for the extent of data induced error. No noticeable temperature dependent trend is observed in the relative molar absorptivities. This should be no problem since the calculated first order rate constants are not dependent on the magnitude of the concentrations just their rate of change. Table IV lists the relative concentration values calculated for each temperature and run-time using Beer's law and the final calculated average relative molar absorptivities of Table III.

TABLE III  
RELATIVE MOLAR ABSORPTIVITIES

Dataset	Relative Absorptivities $\epsilon_{\text{HOD}} = 1.0$			Relative Absorptivities $\epsilon_{(\text{HOD})_2} = 1.0$		
	$\epsilon_{\text{D}_2\text{O}}$	$\epsilon_{(\text{HOD})_2}$	$\epsilon_{\text{HOD}}$	$\epsilon'_{\text{D}_2\text{O}}$	$\epsilon'_{(\text{HOD})_2}$	$\epsilon'_{\text{HOD}}$
<u>150 K</u>						
2/2/82	1.292	1.113	1.0	1.172	1.0	0.9027
6/7/82	1.289	1.247	1.0	1.042	1.0	0.8063
6/9/82	1.422	1.338	1.0	1.148	1.0	0.7942
<u>145 K</u>						
1/29/82	1.327	1.171	1.0	1.152	1.0	0.8603
5/10/82	1.214	1.141	1.0	.9669	1.0	1.101
5/24/82	1.277	0.9053	1.0	1.434	1.0	1.111
6/3/82	1.277	1.295	1.0	1.003	1.0	0.7817
<u>140 K</u>						
5/11/82	1.377	1.303	1.0	1.097	1.0	0.7803
5/25/82	-----omitted due to inconsistent first data points-----					
<u>135 K</u>						
5/15/82	1.307	1.285	1.0	1.032	1.0	0.7863
5/26/82	1.284	1.202	1.0	1.099	1.0	0.8481
6/11/82	1.311	1.441	1.0	1.021	1.0	0.7469
Average Value	1.307	1.222	1.0	1.106	1.0	0.8653
$\sigma$	0.05228	0.1347	0.0	0.1216	0.0	0.1209
90% Confidence Interval	$\pm 0.026$	$\pm 0.067$	—	$\pm 0.060$	—	$\pm 0.060$
	$\epsilon'_{\text{D}_2\text{O}} \cdot \epsilon_{(\text{HOD})_2} = 1.352$			$\epsilon_{\text{D}_2\text{O}} = 1.307$		
	$\epsilon'_{(\text{HOD})_2} \cdot \epsilon_{(\text{HOD})_2} = 1.222$			$\epsilon_{(\text{HOD})_2} = 1.222$		
	$\epsilon'_{\text{HOD}} \cdot \epsilon_{(\text{HOD})_2} = 1.057$			$\epsilon_{\text{HOD}} = 1.000$		

TABLE IV  
RELATIVE CONCENTRATIONS DETERMINED FROM ABSORBANCE DATA

Dataset	Time (minutes)	[D <sub>2</sub> O]	[(HOD) <sub>2</sub> ]	[2HOD] <sup>a</sup>
<u>150 K</u>				
2/2/82	0.0	9.08	2.82	5.67
	2.0	7.10	4.13	6.41
	4.0	5.54	4.58	7.69
	6.0	4.33	4.46	8.97
	9.5	2.97	3.81	10.98
	12.5	2.33	3.15	12.26
	15.5	1.82	2.58	13.36
	20.5	1.43	1.84	14.46
6/7/82	0.0	10.94	2.58	5.55
	2.0	9.67	3.13	6.42
	4.0	7.77	4.21	7.09
	6.0	6.22	4.54	8.24
	8.0	5.09	4.70	9.38
	10.0	4.45	4.38	10.34
	12.0	3.92	3.95	11.30
	15.0	3.25	3.40	12.54
	20.0	2.61	2.82	13.79
	25.0	2.08	2.37	14.75
6/9/82	0.0	10.47	1.72	6.03
	1.0	9.82	1.80	6.42
	2.0	9.49	2.13	6.51
	3.0	9.02	2.29	6.70
	4.0	8.47	2.62	6.89
	5.0	7.71	2.99	7.47
	6.0	7.18	3.36	7.85
	8.0	6.57	3.36	8.62
	10.0	5.91	3.27	9.38
	12.0	5.23	3.27	10.15
	15.0	4.53	2.78	11.59
20.0	3.61	2.41	12.93	
<u>145 K</u>				
1/29/82	0.0	10.50	2.54	5.70
	3.0	8.84	3.68	6.27
	6.0	7.29	4.42	7.03
	9.0	6.44	4.87	7.60
	12.0	5.67	5.03	8.17
	17.0	4.71	4.82	9.12
	21.7	3.95	4.54	10.26
	27.0	3.35	4.21	11.21
	32.0	2.92	3.85	11.97

TABLE IV (Continued)

Dataset	Time (minutes)	[D <sub>2</sub> O]	[(HOD) <sub>2</sub> ]	[2HOD] <sup>a</sup>
	42.0	2.35	3.19	13.11
	52.0	1.98	2.78	13.87
	62.0	1.71	2.29	14.63
	72.0	1.63	2.09	15.01
5/10/82	0.0	7.93	3.02	6.44
	3.0	7.30	3.35	6.80
	6.0	6.74	3.60	7.16
	9.0	6.08	3.81	7.70
	12.0	5.87	3.68	7.97
	17.0	5.05	3.79	8.68
	22.0	4.55	3.71	9.31
	27.0	4.16	3.48	9.93
	32.0	3.69	3.33	10.56
	47.0	2.73	2.72	11.99
	62.0	2.22	2.29	12.89
	77.0	1.72	1.92	13.78
5/24/82	0.0	7.08	2.86	5.32
	3.0	6.21	3.18	5.65
	6.0	5.18	3.74	6.15
	9.0	4.53	3.81	6.65
	12.0	3.86	3.94	7.32
	17.0	2.91	3.95	8.31
	22.0	2.40	3.74	8.98
	27.0	2.29	3.39	9.48
	37.0	1.74	2.84	10.64
	47.0	1.48	2.50	11.39
6/3/82	0.0	9.28	1.51	6.21
	3.0	8.02	2.19	6.55
	6.0	7.24	2.88	6.89
	9.0	6.69	3.01	7.23
	12.0	6.23	3.26	7.57
	17.0	5.60	3.36	8.08
	22.0	5.09	3.29	8.67
	27.0	4.41	3.30	9.35
	32.0	4.07	3.93	9.95
	37.0	7.93	3.90	7.52
	47.0	7.50	4.02	7.90
	62.0	6.67	4.17	8.46
	82.0	5.81	4.21	9.21
	102.25	5.22	4.09	9.96
	132.0	4.47	3.90	10.72
<u>135 K</u>				
5/15/82	0.1	10.82	2.17	6.21
	3.0	10.56	2.28	6.25

TABLE IV (Continued)

Dataset	Time (minutes)	[D <sub>2</sub> O]	[(HOD) <sub>2</sub> ]	[2HOD] <sup>a</sup>
	6.0	10.47	2.39	6.29
	9.0	10.22	2.58	6.33
	14.3	10.06	2.73	6.37
	24.0	9.67	3.03	6.46
	34.0	9.25	3.36	6.56
	49.1	8.55	3.69	6.65
	64.0	8.33	3.93	6.84
	84.0	7.96	4.17	7.03
	104.0	7.58	4.34	7.22
	124.0	7.32	4.47	7.41
	144.0	6.99	4.66	7.60
	174.0	6.67	4.62	7.98
5/26/82	0.0	6.24	1.22	4.21
	6.0	6.06	1.31	4.25
	11.0	5.90	1.38	4.28
	16.0	5.72	1.54	4.32
	25.0	5.56	1.74	4.35
	35.0	5.39	1.78	4.43
	55.0	5.00	1.98	4.61
	85.0	4.64	2.23	4.79
	115.0	4.35	2.38	4.97
	145.0	4.13	2.50	5.14
	185.0	3.81	2.55	5.32
	225.0	3.59	2.62	5.50
	285.0	3.22	2.60	5.79
6/11/82	0.0	10.75	1.71	5.68
	5.0	10.37	1.75	5.75
	10.0	9.94	2.05	5.79
	20.0	9.70	2.29	5.82
	30.0	9.46	2.67	5.86
	40.0	9.01	2.95	5.95
	60.0	8.54	3.27	6.21
	80.0	8.25	3.52	6.48
	100.0	7.86	3.72	6.57
	130.0	7.29	3.97	6.92
	160.0	6.89	4.19	7.28
	190.0	6.32	4.21	7.63
	220.0	5.98	4.34	7.99

<sup>a</sup>Concentration of two HOD units, thus  $[2\text{HOD}] = \frac{1}{2} \cdot [\text{HOD}]$ .

### Model Development

One goal of this study was to find the proton transfer and rotational rates and their respective activation energies. To accomplish this goal a model must be developed that relates  $k_t$ , the rate constant for transfer of protons via an ionic defect migration, and  $k_r$ , the rate constant for migration via a D or L Bjerrum defect migration, to the kinetics spectroscopically observed.

Since all the oxygens in a cubic ice lattice are tetrahedrally linked one to another then any migration path of deuterons along these linkages can be represented as a chain of OH groups as shown in Figure 25a.

If the deuterons are examined in Figure 25a it can be seen that with proton transfer alone, only 3 unique positions of the deuterons in the OH chain are possible as shown and labeled in Figure 25b. If a Bjerrum D or L defect migrates through or along the chain then a proton or deuterium rotation may follow that leaves the two deuterons separated by two oxygen atoms instead of one. However, this separation of the deuterons will affect the positions of the deuterons relative to each other only if they are in positions 2 or 3 (i.e. after a transfer), and the defect approaches in such a manner that both defects are not rotated together thus preserving their orientation and translating that particular deuteron configuration through the crystal lattice. So if the number of oxygens between the two deuterons are counted and their position index as defined in Figure 25b noted, any particular configuration of deuterons with respect to successive rotations and transfers can be represented by using two subscripts

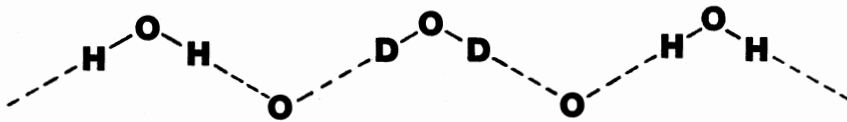


Figure 25a. Hydrogen Bonded Chain in Cubic Ice

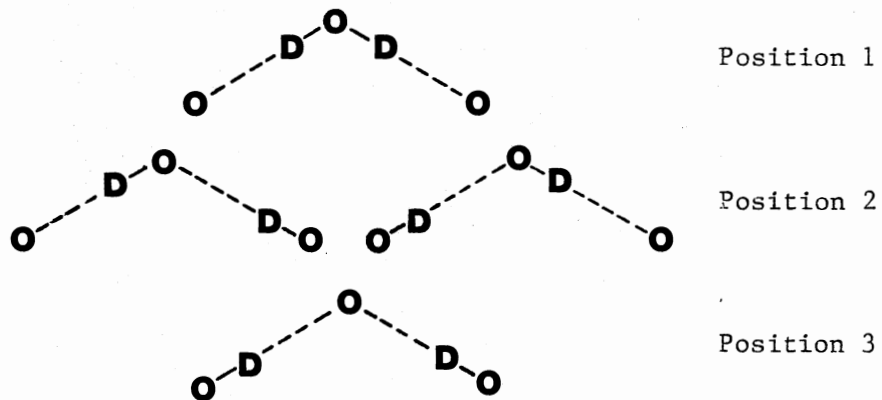
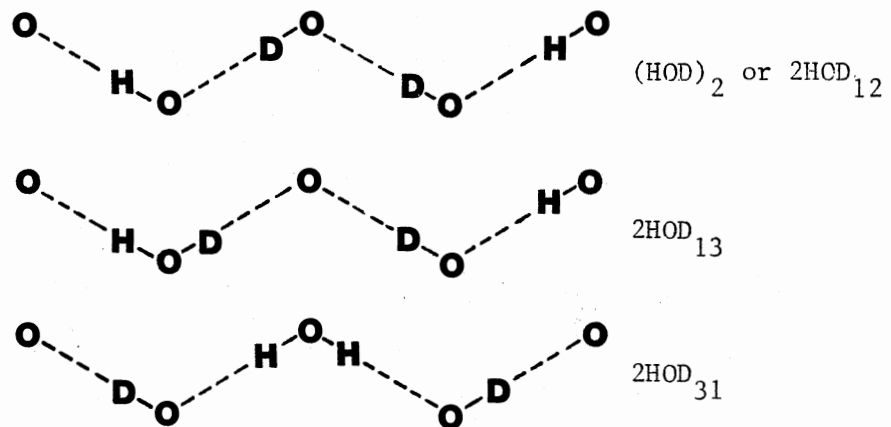


Figure 25b. Possible Proton Transfer Positions

Figure 25c. Example 2HOD<sub>nm</sub> Configurations

as shown in Figure 25c.

The next question presented is how does a  $D_2O$  unit in a crystal-line ice lattice convert to two HODs which then migrate away from each other. A  $D_2O$  unit can be represented by  $D_2O_{11}$  as shown in Figure 26. If an ion defect encounters  $D_2O_{11}$  then a deuteron transfer is possible leaving behind "coupled HOD" or  $2HOD_{12}$ . At this point another ion pair defect could migrate through the  $2HOD_{12}$  site transferring or shifting the other deuteron to form  $2HOD_{13}$  or "pseudo-coupled HOD". Alternately a D or L Bjerrum defect could migrate through the coupled HOD, or  $2HOD_{12}$  site leaving behind a  $2HOD_{21}$  or "separated HOD". At this point "pseudo-coupled HOD", or  $2HOD_{13}$ , must meet a D or L defect to proceed further to  $2HOD_{22}$  or "fully separated HOD". Likewise separated HOD or  $2HOD_{21}$  cannot migrate further unless an ion pair defect is encountered to leave  $2HOD_{22}$  or fully separated HOD behind. Thus the two divergent migration paths reunite back at  $2HOD_{22}$  where either a transfer or rotation can occur to the  $2HOD_{31}$  "position 1 form" or  $2HOD_{23}$  "position 3 form". These forms then migrate to a "position 2 form" and the sequence can repeat continually. But if deuterons can migrate apart from each other via a tunneling transfer or thermally activated rotation they can also migrate together. Figure 26 illustrates this aspect with double arrows and uses t or r to indicate whether the step is a transfer or rotation.  $D_2O_{11}$  and  $2HOD_{12}$  are spectroscopically distinguishable. However according to the experimental evidence to date, the other  $2HOD_{nm}$  are perceived spectroscopically as uncoupled HOD.

To relate this model to the measured  $D_2O$ ,  $(HOD)_2$ , and HOD absorbances some approximation must be incorporated to terminate the



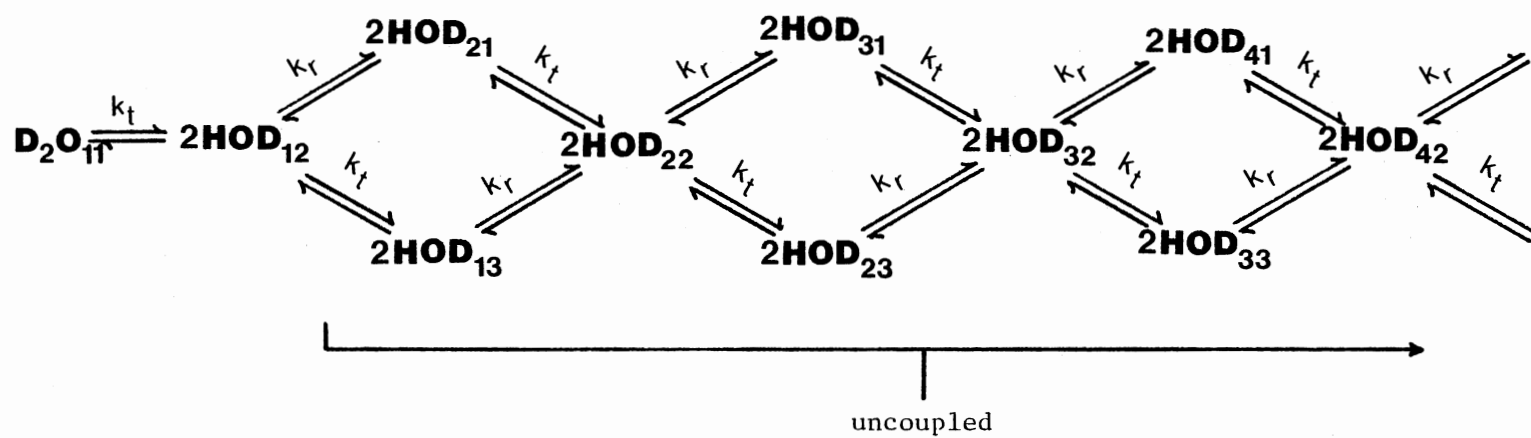
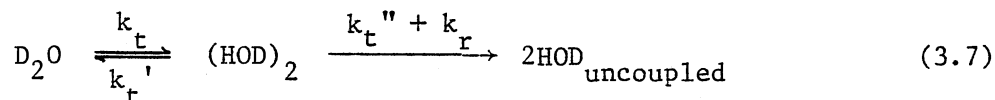
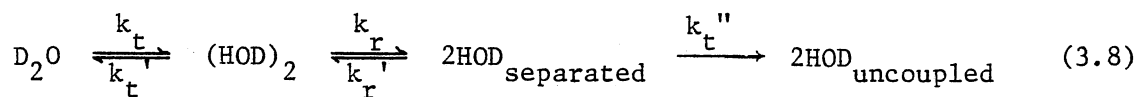


Figure 26.  $D_2O_{11}$  to  $2HOD_{nm}$  Migration Mechanism

number of steps and make any least square fitting technique practical. The first approximation suggested is that the exchange sequence terminates at  $2\text{HOD}_{21}$  and  $2\text{HOD}_{13}$ , the pseudo-coupled and separated forms of HOD. In addition there is no back reaction of  $2\text{HOD}_{21}$  or  $2\text{HOD}_{13}$  to  $2\text{HOD}_{12}$ . This would give



The lack of a back reaction from  $2\text{HOD}_{\text{uncoupled}}$  to  $(\text{HOD})_2$  could be justified by the probable small concentrations of  $2\text{HOD}_{21}$  and  $2\text{HOD}_{13}$  relative to  $2\text{HOD}_{12}$  or  $(\text{HOD})_2$ , due to the forward reactions of  $2\text{HOD}_{21}$  and  $2\text{HOD}_{13}$  that would tend to distribute the newly separated deuterons further apart in the ice lattice. This single equilibrium approximation can be expanded to a double equilibrium model provided an additional unknown parameter is introduced into the model, the  $2\text{HOD}_{21}$  concentration. If the  $2\text{HOD}_{12}$  to  $2\text{HOD}_{13}$  to  $2\text{HOD}_{22}$  forward and back pathways are ignored on the basis of rate constant magnitudes, (to be discussed later), then the following model approximation may be written.



Here the back reaction of  $2\text{HOD}_{\text{uncoupled}}$  or  $2\text{HOD}_{22}$  to  $2\text{HOD}_{\text{separated}}$  or  $2\text{HOD}_{21}$  is ignored on the basis of the small concentration of  $2\text{HOD}_{22}$  due to the  $2\text{HOD}_{22}$  to  $2\text{HOD}_{31}$  or  $2\text{HOD}_{23}$  forward reactions and the position of  $2\text{HOD}_{22}$  configuration in the migration pathway. The latter is suggested since the further two deuteriums migrate apart, the more options there are for different migration paths leading to different configurations. Hence the further two deuterons move apart the less

chance there is for the migration path to reverse to the original starting  $D_2O$  unit. The only remaining hurdle is to find the relationship between the various  $k_t$ s and  $k_r$ s before the models may be used.

#### Determining Rate Constant Relationships

In the gas or liquid phase, reaction rates are determined from collision frequencies, orientation factors, reaction cross sections, energy states populations, transition probabilities, and other factors. Mono and bi-molecular reactions can be expressed with first and second order kinetic rate equations. But reactions in the solid phase can be more complex and indeterminate. Orientations may be fixed, and two body collisions are almost nonexistent.

For a transfer or rotational movement of a deuteron to occur, a migrating ice defect must encounter a  $D_2O$  or HOD configuration. This suggests a second order process. But at a given temperature the concentration and rate of migration of the defects is constant. Furthermore, the defect is not consumed in the reaction so that its concentration still stays constant. Thus all the transfer rate constants ( $k_t$ ) and rotational rate constants ( $k_r$ ) expressed in Equation (3.7), (3.8), and Figure 26 are first order with respect to the  $D_2O$  or  $2HOD_{nm}$  configuration concentrations. The transfer rates between configurations should be equal except for the effect that the different number of approaches and subsequent conversion possibilities arising from different configurations should produce on the transfer rates. A similar situation exists for the rotational rates. The next logical step is to determine what different defect approaches and reacting pathways are possible for different deuteron configurations.

Figure 27 and Table V show a brief representation and resulting mathematical formulation for the  $D_2O_{11} \rightarrow 2HOD_{12}$  exchange reaction. Reaction I Figure 27 shows three sequential steps labeled the attack, transfer, and leaving factors. If the relative probability of occurrence for each step is determined for each reaction then a comparison between different  $k_t$ s may be made. For the attack step a hydronium ionic defect moves to the  $D_2O$  site. There are two different hydrogen bonds on which the attacking hydrogen can tunnel. Then the attack factor is written as  $2\nu_H$  where  $\nu_H$  is the tunneling frequency of the attacking proton (see Table V). On the transfer step either the proton can tunnel back to its original oxygen at a frequency of  $\nu_H$  or one of the two deuterons may tunnel at a frequency of  $\nu_D$  to another oxygen to yield  $2HOD_{12}$  or coupled HOD. The total tunneling frequency would be  $\nu_H + 2\nu_D$  and the tunneling frequency for the desired reaction would be  $2\nu_D$ . Then a probability of occurrence for the desired event could be defined as  $2\nu_D/(\nu_H + 2\nu_D)$ . To assure that the transfer step, which leaves a hydronium ion defect at the  $HOD_{12}$  site, does not back exchange to  $D_2O_{11}$  a leaving factor is included using the same logic as for the transfer step. Using these three factors a relative probability of reaction can be assigned to reaction I through VI by

$$P_{R_x} = (\text{attack factor } R_x)(\text{transfer factor } R_x)(\text{leaving factor } R_x) \quad (3.9)$$

The relationship between any given transfer reaction can be written as

$$k_t' = \frac{(\text{attack factor } k_t')(\text{transfer factor } k_t')(\text{leaving factor } k_t')}{(\text{attack factor } k_t)(\text{transfer factor } k_t)(\text{leaving factor } k_t)} k_t \quad (3.10)$$

Table V lists the transfer factors found for Figure 27 and 28 for the forward and back transfer exchange reactions between  $D_2O_{11}$ ,  $2HOD_{12}$ ,

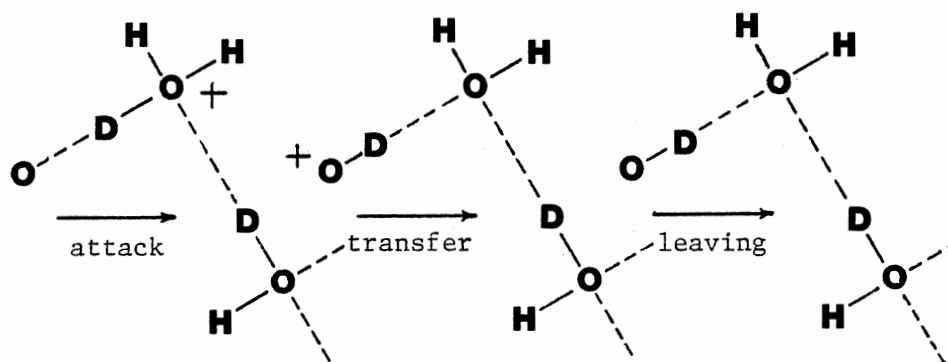
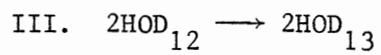
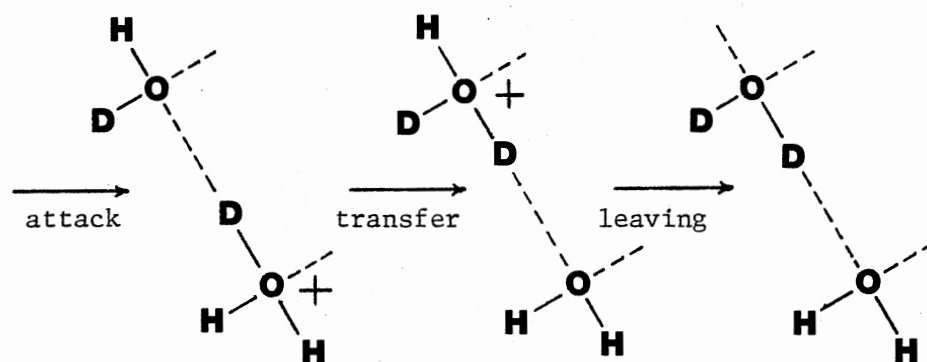
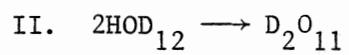
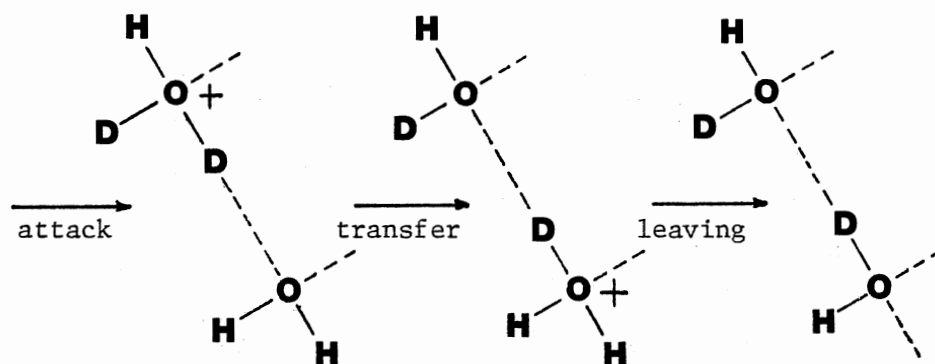
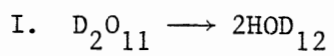


Figure 27. Proton Transfer Steps for Deuterium Migration

TABLE V  
RELATIVE PROBABILITY FACTORS FOR TRANSFER MIGRATION

Reaction	Attack Factor	Transfer Factor	Leaving Factor
$D_2O_{11} \rightarrow 2HOD_{12}$	$2v_H$	$\frac{2v_D}{v_H + v_D}$	$\frac{2v_H}{2v_H + v_D}$
$2HOD_{12} \rightarrow D_2O_{11}$	$2v_H$	$\frac{v_D}{2v_H + v_D}$	$\frac{2v_H}{v_H + 2v_D}$
$2HOD_{12} \rightarrow 2HOD_{13}$	$1v_H$	$\frac{v_D}{2v_H + v_D}$	$\frac{2v_H}{2v_H + v_D}$
$2HOD_{13} \rightarrow 2HOD_{12}$	$4v_H$	$\frac{v_D}{2v_H + v_D}$	$\frac{2v_H}{2v_H + v_D}$
$2HOD_{21} \rightarrow 2HOD_{22}$	$4v_H$	$\frac{v_D}{2v_H + v_D}$	$\frac{2v_H}{2v_H + v_D}$
$2HOD_{22} \rightarrow 2HOD_{21}$	$2v_H$	$\frac{v_D}{2v_H + v_D}$	$\frac{2v_H}{2v_H + v_D}$
$2HOD_{22} \rightarrow 2HOD_{23}$	$2v_H$	$\frac{v_D}{2v_H + v_D}$	$\frac{2v_H}{2v_H + v_D}$
$2HOD_{23} \rightarrow 2HOD_{22}$	$4v_H$	$\frac{v_D}{2v_H + v_D}$	$\frac{2v_H}{2v_H + v_D}$
For $n \geq 2$			
$2HOD_{n1} \rightarrow 2HOD_{n2}$	$4v_H$	$\frac{v_D}{2v_H + v_D}$	$\frac{2v_H}{2v_H + v_D}$
$2HOD_{n2} \rightarrow 2HOD_{n1}$	$2v_H$	$\frac{v_D}{2v_H + v_D}$	$\frac{2v_H}{2v_H + v_D}$

TABLE V (Continued)

Reaction	Attack Factor	Transfer Factor	Leaving Factor
$2\text{HOD}_{n2} \rightarrow 2\text{HOD}_{n3}$	$2v_H$	$\frac{v_D}{2v_H + v_D}$	$\frac{2v_H}{2v_H + v_D}$
$2\text{HOD}_{n3} \rightarrow 2\text{HOD}_{n2}$	$4v_H$	$\frac{v_D}{2v_H + v_D}$	$\frac{2v_H}{2v_H + v_D}$

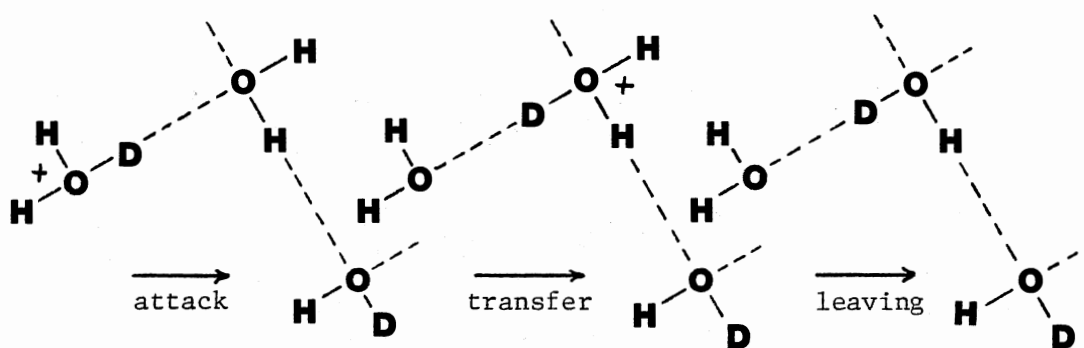
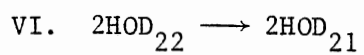
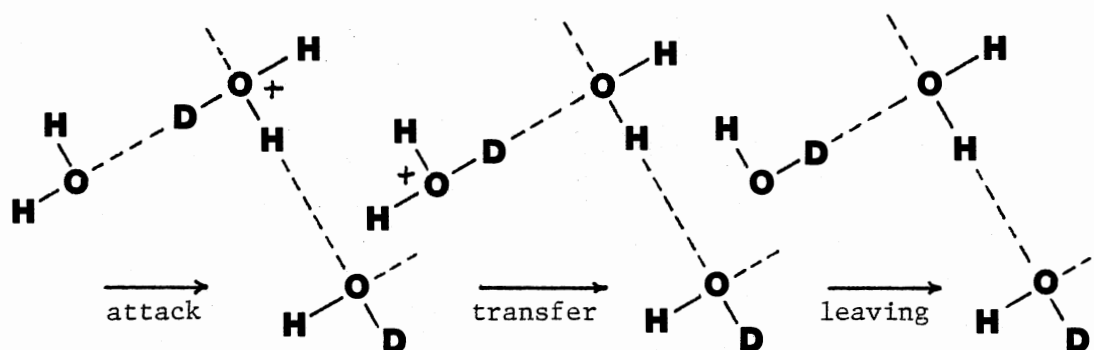
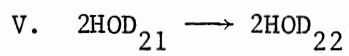
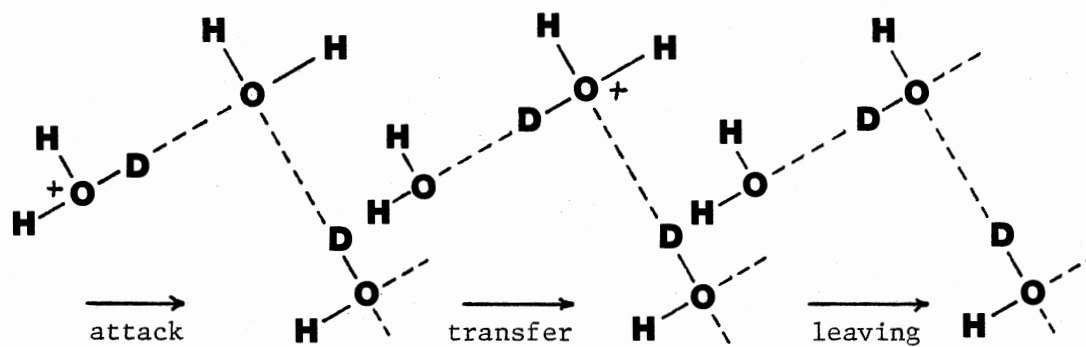
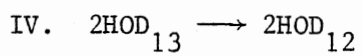


Figure 28. Proton Transfer Steps for Deuterium Migration



$2\text{HOD}_{13}$ ,  $2\text{HOD}_{21}$ , and  $2\text{HOD}_{22}$  the  $\text{D}_2\text{O}_1$  coupled HOD, pseudo HOD, separated HOD, and fully separated HOD molecular units of cubic ice. It should be noted that the hydroxide ionic defect does not appear to play a significant part in the  $\text{D}_2\text{O}$  isotopic exchange reaction (46, 47).

Examining Table VI and Figures 29 and 30 reveals a similar formulation for the rotational step or a Bjerrum D or L defect migrating through a  $2\text{HOD}_{nm}$  site. Reaction I the  $2\text{HOD}_{12} \rightarrow 2\text{HOD}_{21}$  reaction in Figure 29 shows how an L defect can migrate through the  $2\text{HOD}_{12}$  unit to change it to  $2\text{HOD}_{21}$ . The L defect may enter the  $2\text{HOD}_{23}$  along 3 hydrogen bonded paths, shown by dotted lines. However, the top hydrogen bond pathway does not change the configuration. Thus there are only two attack paths. This gives an attack factor of 2 for the L defect (see Table VI). In the transfer step the L defect moves to the oxygen-oxygen linkage between the two HODs as a result of the deuteron rotating down to the lower O-O linkage. In the bottom O-O linkage, there are four different protons that can rotate into the L defect, as will always be the case of an L defect migrating through a perfect Bernard-Fowler ice crystal. If the bottom protons rotates in, the L defect will move backwards thus not initiating the transfer step. If the deuteron rotates in, the L defect moves up the hydroxyl chain to the middle position shown in the transfer step. However, if the proton on the bottom HOD unit rotated in, the L defect would simply shift to another attack position identical to the position shown in the attack step. Thus of four possible proton or deuteron rotations only three change the L defect with respect to the reaction pathway, and only one of these three, namely the deuteron rotation, initiates the transfer step. Therefore the relative probability for the transfer

TABLE VI  
RELATIVE PROBABILITY FACTORS FOR ROTATIONAL MIGRATION

Reaction	Attach Factor	Transfer Factor	Leaving Factor	Total
<u>L Defect</u>				
$2\text{HOD}_{12} \rightarrow 2\text{HOD}_{21}$	2	$\frac{1}{3}$	$\frac{1}{2}$	$\frac{1}{3}$
$2\text{HOD}_{21} \rightarrow 2\text{HOD}_{12}$	2	$\frac{1}{3}$	$\frac{1}{2}$	$\frac{1}{3}$
$2\text{HOD}_{13} \rightarrow 2\text{HOD}_{22}$	4	$\frac{1}{3}$	$\frac{3}{4}$	1
Sequence A				
$2\text{HOD}_{22} \rightarrow 2\text{HOD}_{13}$	2	$\frac{1}{3}$	$\frac{1}{4}$	$\frac{1}{6}$
Sequence B				
$2\text{HOD}_{22} \rightarrow 2\text{HOD}_{13}$	1	$\frac{1}{3}$	$\frac{1}{4}$	$\frac{1}{12}$
Total				
$2\text{HOD}_{22} \rightarrow 2\text{HOD}_{13}$	$\frac{1}{3} (\text{Sequence A}) + \frac{2}{3} (\text{Sequence B}) = \frac{1}{9}$			
<u>D Defect</u>				
$2\text{HOD}_{12} \rightarrow 2\text{HOD}_{21}$	1	$\frac{1}{3}$	$\frac{1}{2}$	$\frac{1}{6}$
$2\text{HOD}_{21} \rightarrow 2\text{HOD}_{12}$	1	$\frac{1}{3}$	$\frac{1}{2}$	$\frac{1}{6}$
$2\text{HOD}_{13} \rightarrow 2\text{HOD}_{22}$	2	$\frac{1}{2}$	$\frac{2}{3}$	$\frac{2}{3}$
$2\text{HOD}_{22} \rightarrow 2\text{HOD}_{13}$	1	$\frac{1}{3}$	$\frac{1}{3}$	$\frac{1}{9}$

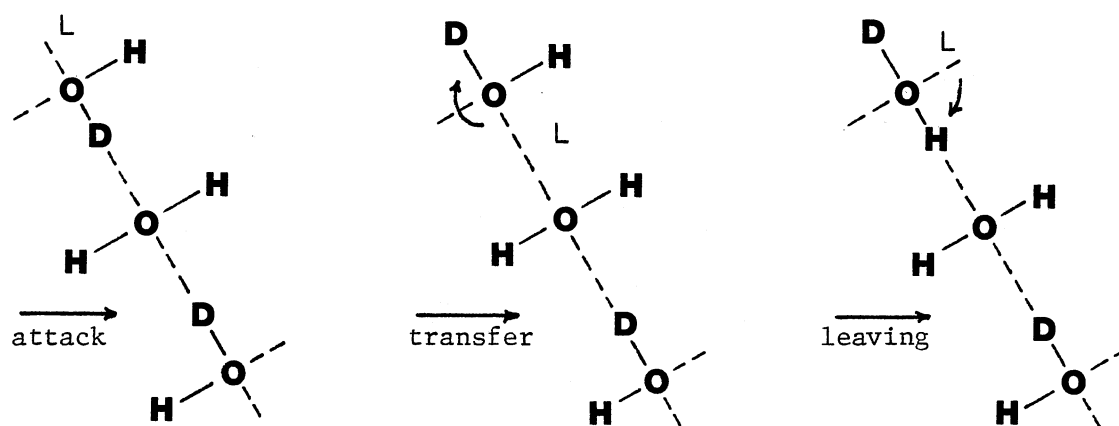
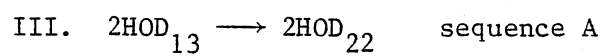
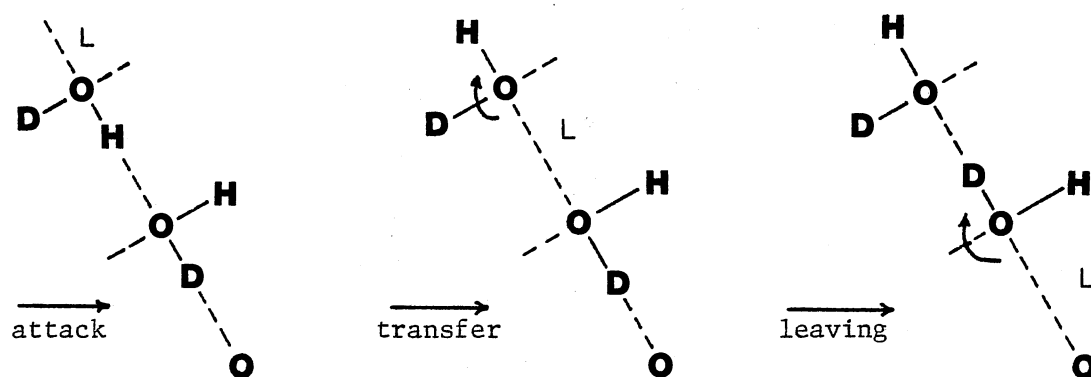
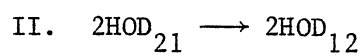
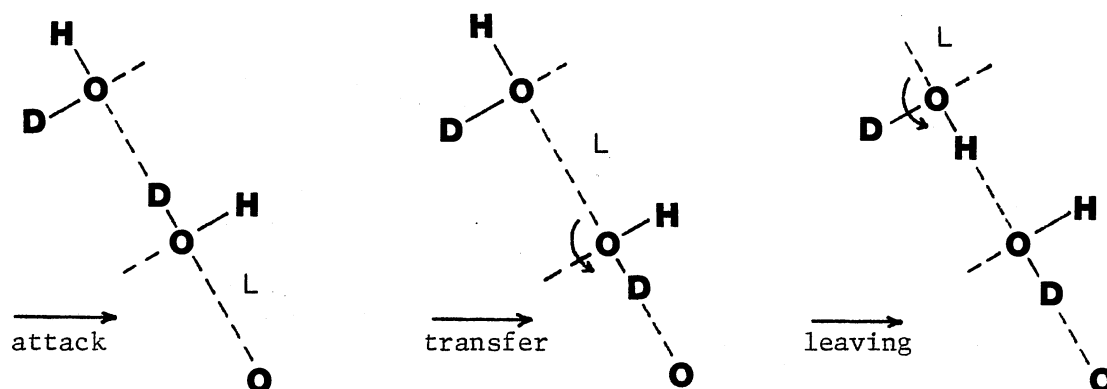
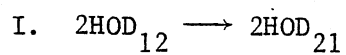
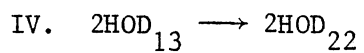
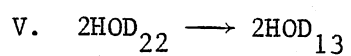
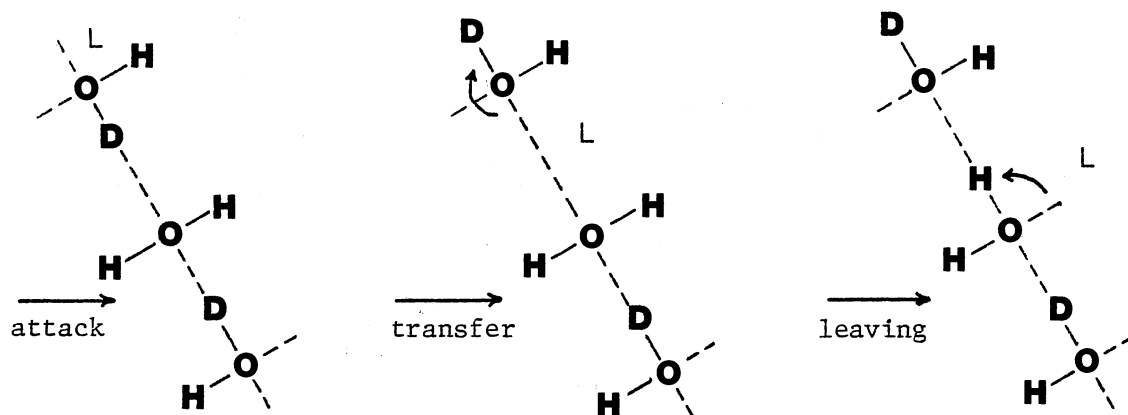


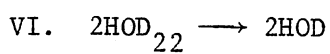
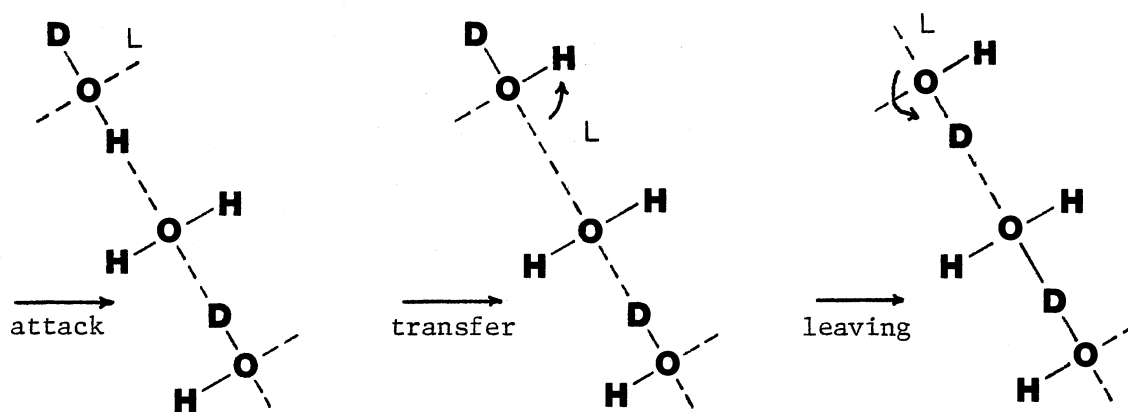
Figure 29. Proton Rotational Steps for Deuteron Migration



sequence B



sequence A result



sequence B result

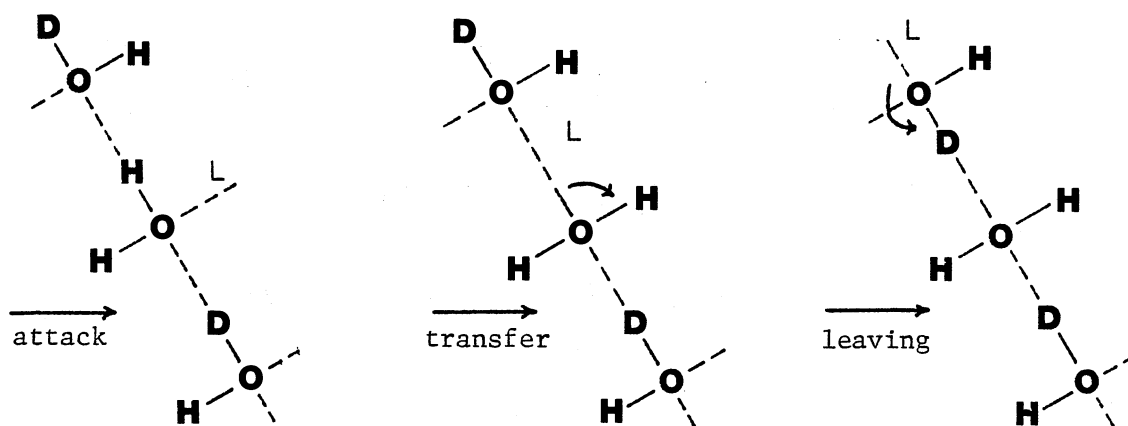


Figure 30. Proton Rotational Steps for Deuteron Migration

step is  $\frac{1}{3}$ . For the leaving step all four proton or deuterons will change the L defect with respect to the reaction pathway. Two of these (the protons rotating in) leave the  $2\text{HOD}_{21}$  intact as the L defect moves on. But if either of the two deuterons rotates in then the L defect movement leaves a  $2\text{HOD}_{12}$  configuration behind and a L defect position identical to the position shown in the attack step. Thus for the leaving step there is a  $\frac{2}{4}$  or  $\frac{1}{2}$  probability that the desired reaction or step occurs. The total relative probability is now found by multiplying these factors together to yield  $\frac{1}{3}$  as shown in Table VI. The relative probability for the back reaction  $2\text{HOD}_{21} \rightarrow 2\text{HOD}_{12}$  is likewise determined.

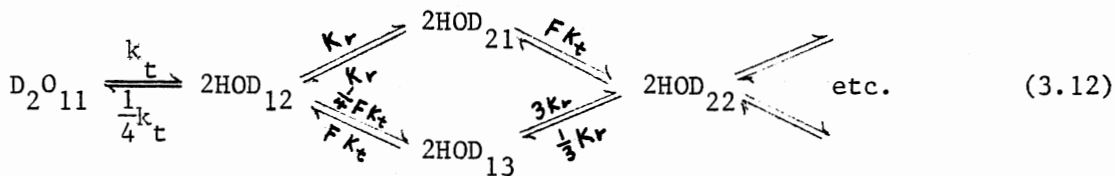
The  $2\text{HOD}_{13} \rightarrow 2\text{HOD}_{22}$  (pseudo to fully separated HOD) reaction is slightly more complex. As shown in reaction III Figure 29 and reaction IV Figure 30, there are two different sequences to convert  $2\text{HOD}_{13}$  to  $2\text{HOD}_{22}$ , each leaving a different configured  $2\text{HOD}_{22}$ . Sequence A or reaction III leaves a  $2\text{HOD}_{22}$  product with the end HOD group providing the rotated proton for the leaving step. Sequence B or reaction IV has the middle  $\text{H}_2\text{O}$  group providing the proton for the leaving step. Since the middle  $\text{H}_2\text{O}$  has two protons to rotate into the L defect but the end HOD group of sequence A has only one proton to rotate into the L defect on the transfer step then sequence B should occur twice as much as sequence A. For the forward  $2\text{HOD}_{13} \rightarrow 2\text{HOD}_{22}$  reaction this should present no problem. But for the reverse  $2\text{HOD}_{22} \rightarrow 2\text{HOD}_{13}$  reaction each sequence has different attack factors. So to account for this pathway splitting the sequence A total relative probability is weighted by  $\frac{1}{3}$  to account for its share of the probability of  $2\text{HOD}_{13} \rightarrow 2\text{HOD}_{22}$  occurring. Likewise sequence B is weighted by  $\frac{2}{3}$ .

The total reaction relative probability is given by the sum of the sequence relative probabilities.

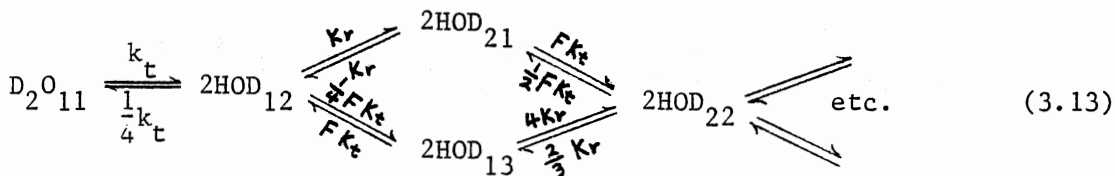
Similar relative probability calculations for a D defect migration through a  $D_2O_{11}$  or  $2HOD_{nm}$  site are also shown in Table VI. As with transfer rate constants the rotational rate constants can be related through

$$k_r' = \frac{(\text{total relative probability } k_r') k_r}{(\text{total relative probability } k_r)} \quad (3.11)$$

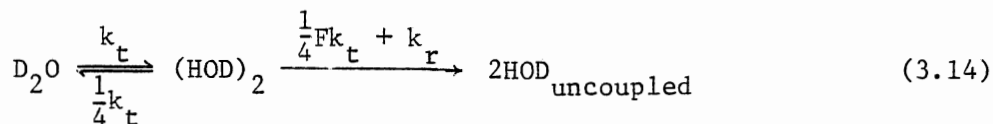
Using (3.10) and (3.11) the relationships between the various  $k_t$ 's and  $k_r$ 's (the transfer and rotational rate constants) of equation (3.7) and (3.8) are written in terms of one  $k_t$  and one  $k_r$ . Figure 26 may be rewritten in terms of the  $D_2O_{11} \rightarrow 2HOD_{12}$   $k_t$  and  $2HOD_{12} \rightarrow 2HOD_{21}$   $k_r$  as



where  $F = \left( \frac{v_H + 2v_D}{2v_H + v_D} \right)$  and the  $k_r$  relationships are calculated using the L defect factors. If the D defect factors are used for  $k_r$  then Figure 26 may be written

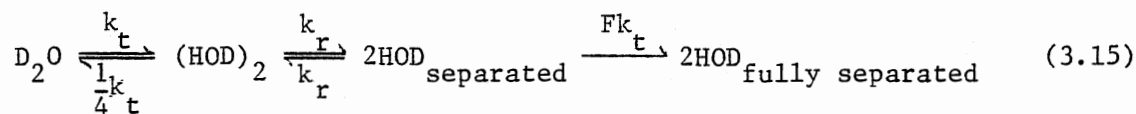


Equation (3.7) the single equilibrium model approximation may be written as



and the double equilibrium model given in Equation (3.7) using either D

or L rotational defects as



Since the rate constants depleting  $2HOD_{13}$  in equations (3.12, 3.13) are significantly higher than the feeding rate constants, it is expected that the concentration of  $2HOD_{13}$  in an exchange process would be small and steady state. Furthermore since the amount of  $2HOD_{12}$  or coupled HOD being converted to  $2HOD_{13}$  is small ( $\frac{1}{4}k_t$ ) compared to the other decoupling rate constants, it is assumed that the deuterons primarily migrate apart through the  $2HOD_{12}$ ,  $2HOD_{21}$ ,  $2HOD_{22}$  pathway, as suggested by the double equilibrium (3.15) model. The single equilibrium (3.14) model assumes that no back reaction of  $2HOD_{21}$  or  $2HOD_{13}$  to  $2HOD_{12}$  occurs. Perusal of these rate constants shows that they are equal or greater in magnitude than the forward rate constants and hence the assumption may not be valid. The double equilibrium (3.15) has this same problem with the  $2HOD_{21}$  to  $2HOD_{22}$  forward and back rates. Here, however, the back reaction rate is significantly smaller than the forward ( $Fk_t$  forward to  $\frac{1}{2}Fk_t$  back) and the concentration of  $2HOD_{22}$  should be considerably smaller than  $2HOD_{21}$ . In summation the single equilibrium model (3.14) has three unknown parameters  $k_t$ ,  $k_r$ , and  $F$ ; and includes no provision for any back exchange into the coupled HOD or  $2HOD_{12}$  configuration. The double equilibrium model (3.15) has four unknown parameters  $k_t$ ,  $k_r$ ,  $F$ , and an initial HOD separated or  $2HOD_{21}$  concentration. It provides for a  $2HOD_{21}$  back reaction to  $2HOD_{12}$  or coupled HOD, but includes a spectroscopically unobservable species,  $2HOD_{21}$ . Fitting these two models to the experimental data should provide a measure of the importance of the  $2HOD_{21}$  configuration in the

isotopic scrambling of  $D_2O$  in cubic ice.

Kunst and Warman (48) determined the virtual mobility of protons in ice at  $-5^\circ C$  as  $6.4 \times 10^{-3} \text{ cm}^{-2} \text{ V}^{-1} \text{ sec}^{-1}$  and deuterons in  $D_2O$  ice at  $-5^\circ C$  as  $2.4 \times 10^{-3} \text{ cm}^{-2} \text{ V}^{-1} \text{ sec}^{-1}$ . At  $-30^\circ C$  they reported that the ratio remained the same though the magnitudes increased. Using these values  $F$  is calculated as .736842. This value was used throughout the calculations and any effects resulting from changing this value are noted.

Determining an initial  $2HOD_{21}$  or separated HOD concentration was not straightforward. If an ice deposit of  $D_2O$  isolated in  $H_2O$  with no coupled or uncoupled HOD could be made, then the initial  $2HOD_{21}$  concentration could be set equal to zero. However, all ice deposits made contained small amounts of coupled HOD initially and some uncoupled HOD in addition. So as a reasonable approximation the initial  $2HOD_{21}$  concentration was taken as some function of the initial  $(HOD)_2$  or  $2HOD_{12}$  concentration. The idea being that whatever processes (limited pre-exchange or deposition contamination) created the initial  $2HOD_{12}$  concentrations would also create the initial  $2HOD_{21}$  concentration in the ice deposit. Fractional values of initial  $2HOD_{12}$  between 0.0 and 1.0 were tried.

Thus with an initial guessed  $2HOD_{21}$  concentration and the previously calculated  $F$  constant, both the single and double equilibrium models were used to fit the kinetic concentration data. These calculations served to determine the appropriate transfer ( $k_t$ ) and rotational ( $k_r$ ) rate constants, for isotopic exchange processes in cubic ice at temperatures ranging from 135-150 K.



### Solution and Data Analysis of the Models

The differential rate equations for the single equilibrium (3.14) model are

$$\frac{d[D_2O]}{dt} = \frac{1}{4}k_t [(HOD)_2] - k_t[D_2O] \quad (3.16)$$

$$\frac{d[(HOD)_2]}{dt} = k_t[D_2O] - \frac{1}{4}[(HOD)_2] - \left(\frac{1}{4}Fk_t + k_r\right)[(HOD)_2]$$

$$\frac{d[2HOD_{uncoupled}]}{dt} = \left(\frac{1}{4}Fk_t + k_r\right)[(HOD)_2]$$

Likewise for the two equilibrium model (3.15) the differential rate equations are

$$\frac{d[D_2O]}{dt} = \frac{1}{4}k_t [(HOD)_2] - k_t[D_2O] \quad (3.17)$$

$$\frac{d[D_2O]}{dt} = k_t[D_2O] - \frac{1}{4}k_t [(HOD)_2] - k_r[(HOD)_2] + k_r[2HOD_{separated}]$$

$$\frac{d[2HOD_{separated}]}{dt} = k_r[(HOD)_2] - k_r[2HOD_{separated}] - Fk_t[2HOD_{separated}]$$

$$\frac{d[2HOD_{fully\ separated}]}{dt} = Fk_t[2HOD_{fully\ separated}]$$

$$\text{and } [2HOD_{uncoupled}] = [2HOD_{separated}] + [2HOD_{fully\ separated}]$$

Before any model fitting is accomplished these differential equations must be solved. For complex mechanisms two basic approaches are possible. One, find a way to solve the equations analytically, or, two numerically solve them using computer differential equation solvers. Matsen and Franklin (49) have published the general solution to the four component coupled set of first order reactions. The technique sets up the differential equations as a secular determinant similar to a standard eigenvalue problem. Solutions were then found

from the resulting eigenvalues and eigenvectors. Lewis and Johnson (50) recast the equations into a matrix formulation that is solved by matrix algebra. Unfortunately, these analytical solutions contain many of the problems often encountered in numerical solutions. Therefore, since the differential equations are not stiff; and fast, accurate, easy to use differential equation solvers are available, the numerical solution method was chosen to solve the equations.

A variable order Adams predictor-corrector differential equation solver written by the International Mathematical and Statistical Library group was used to solve all the differential equations of the model simultaneously. Among the features of this routine are variable step sizes, abilities to help handle mildly stiff differential equations, built in user determined relative error bounds, and internal error checking.

The subroutine is named DGEAR and was called from a main Fortran program whenever a solution to the model differential equations was required. DGEAR required that the user supply the differential equations in another subroutine named FCN. FCN required the initial time,  $D_2O$ ,  $(HOD)_2$ ,  $HOD_{\text{separated}}$ ,  $HOD_{\text{fully separated}}$  concentrations, final time at which the concentration solutions are desired, and various control parameters. DGEAR then outputs the various concentrations at the desired time of reaction along with an error flag indicating the reliability of the answers. Initial step-sizes of 0.00001 and relative error bounds of 0.0001 was found to give the best accuracy and speed. A description of DGEAR along with sample input and output is included in Appendix B. Due to copyright restrictions, a program listing is not included.

The next problem was to elucidate some method of varying the rate constants until the calculated concentration output is as identical to the experimental concentrations as possible. The method of nonlinear least squares was chosen. Though the methods of nonlinear least squares are legion they all attempt to minimize the square of the error residuals to find the best value of the minimizing parameters. The Marquardt strategy (5]) of nonlinear least squares is one of the most successful and efficient nonlinear routines used. Accordingly, it was used as a subroutine to find the best transfer and rotational rate constants for each dataset at each temperature. NLLSQ1 is a modified Fortran nonlinear least squares program based on the Marquardt strategy. It is composed of a main calling program which reads the data to be fitted, an initial parameter guess, and various control cards. After placing the needed information into array form, NLLSQ1 calls the Marquardt strategy subroutine NLLSQ. The NLLSQ subroutine also required that two additional subroutines DCALC and YCALC be written by the user. YCALC contains the function to be fitted such that at a given call with a chosen set of independent function variables the resulting dependent variables may be calculated and returned to NLLSQ. DCALC may optionally contain the function derivatives or be used as a dummy routine while the subroutine NLLSQ numerically calculated the derivatives. The latter option was used in this work. Once the data is least square fitted, the final optimized parameters (transfer and rotational rate constants), calculated dependent variables ( $[D_2O]$ ,  $[(HOD)_2]$ ,  $[2HOD_{uncoupled}]$ ), and root mean square deviation of fit is printed.

Least squares analysis strives to fit an equation  $Y = f(X_1, X_2, X_3, \dots)$

to an observed set of  $Y, X_1, X_2, X_3 \dots$  data.  $Y$  is called the dependent variable, and  $X_i$  the independent variable. For the single and double equilibrium models the dependent variables are the concentrations of  $D_2O$ ,  $(HOD)_2$ , and uncoupled HOD. The uncoupled HOD is calculated from the sum of separated HOD and fully separated HOD for the double equilibrium models, and fully separated HOD for the single equilibrium model. The time of reaction  $t$  is the independent variable and  $k_t$  and  $k_r$  the minimizing parameters. The initial concentrations of  $D_2O$ ,  $(HOD)_2$ , separated HOD, and fully separated HOD were function constants that changed with each dataset analyzed. But since NLLSQ operates on functions with one dependent variable, only one species concentration could be fitted per run, unless the three dependent concentrations are transformed into a single dependent variable. Thus to obtain a consistent fit of the models to all observed concentrations the three concentrations were transformed into a single dependent variable. This was accomplished by including the  $(HOD)_2$  and uncoupled HOD values after the  $D_2O$  concentration values in the dependent variable array. Thus the first third of the array contained  $D_2O$  concentration values, the second third the  $(HOD)_2$  values, and the last third the uncoupled HOD concentrations values. For the independent or time variable array the reaction time in minutes was repeated three times in the array therefore maintaining a one to one correspondence between the  $i$ th observations of the time and concentration arrays.

When NLLSQ called the subroutine YCALC, an index of observation requested, present calculated values of  $k_t$  and  $k_r$ , and concentration and time value arrays were inputted as arguments. YCALC would initialize the needed information required by DGEAR for solving the

differential kinetic rate equations. Using the time value found at the inputted index of observation DGEAR would solve the rate equations for the concentrations of  $D_2O$ ,  $(HOD)_2$ , and uncoupled HOD for that observation time. YCALC then determined which third of the concentration and time arrays the observation index referred to and returned either the  $D_2O$ ,  $(HOD)_2$ , or uncoupled HOD calculated concentration value to the NLLSQ calling subroutine. By calling YCALC for every observation a calculated array of concentrations values was determined for subsequent refinement of the transfer and rotational rate constants.

NLLSQ solved for  $k_t$  and  $k_r$  by iterative improvements of the initial guessed  $k_t$  and  $k_r$  rate constants. When the iteration failed to yield any significant improvement of  $k_t$  and  $k_r$  and specific internal convergence criteria were met, the iteration was stopped and the final computed concentration and time arrays along with the final  $k_t$  and  $k_r$  values were printed. By examining the root mean square deviation of fit and comparing the calculated versus observed values a judgement was made on the success of the least square fit for the model being tested. Most of the datasets were fit within 6-8 iterative cycles with good convergence criteria using the double equilibrium model. The single equilibrium model was less successful. To avoid the possibility of having converged on a local minima different initial  $k_t$  and  $k_r$  values were tried to see if the final results differed. A copy of NLLSQ with YCALC programmed for the double equilibrium model is included in Appendix C. Also included is sample input, output, and a written description of the NLLSQ nonlinear least square subroutine.

To determine the activation energies associated with the transfer and rotational rate constants the Arrhenius expression for kinetic rate constant temperature dependence was used.

$$\ln K = \frac{E_a}{RT} + \ln A \quad (3.18)$$

Here, A, the pre-exponential factor, was assumed to be constant. SAS, the Statistical Analysis System, marketed by the SAS Institute is a system of computer software designed to perform complex or simple statistical analysis of user data. It features a programming language capability, several statistical procedures to perform linear least square analysis, correlation calculations, and line printer or plotter graphing capabilities.

An SAS program ACTIV was written to accomplish the least square analysis of the rate constants versus temperature using equation (3.18). Each least square analysis included the NLLSQ calculated rate constants for all 12 kinetic datasets. A correlation analysis was carried out on the calculated and observed values obtained from the Arrhenius plots. The procedure was then repeated over all 12 kinetic datasets every time the model or a model functional constant was changed. An examination of the Arrhenius plots and calculated correlation values permitted an evaluation on which models gave the best, most consistent activation energies, and what range of error is expected in the activation energy values. A listing of ACTIV is included in Appendix D.

## CHAPTER IV

### RESULTS AND CONCLUSIONS

#### Spectral and Concentration Data

One of the by-products of resolving the OD stretch multiplet at 2000-2800  $\text{cm}^{-1}$  was an uncluttered look at the coupled HOD in and out-of-phase modes. Figure 31 illustrates one of the better  $(\text{HOD})_2$  spectrum collected. This particular spectrum was obtained at 145 K at 17 min into the exchange reaction. The in-phase component is located at 2404  $\text{cm}^{-1}$  and the out-of-phase component at 2446  $\text{cm}^{-1}$ , thus yielding a 42  $\text{cm}^{-1}$  split. The peak absorbance ratio of the two peaks is .70 but the intensity or band areas are almost equal. This is due to the in-phase bandwidth measuring 32  $\text{cm}^{-1}$  while the out-of-phase component is 35  $\text{cm}^{-1}$ .

This is in excellent agreement with Haas and Hornig's (22) coupled HOD split of 49  $\text{cm}^{-1}$  with frequencies of 2442  $\text{cm}^{-1}$  and 2393  $\text{cm}^{-1}$ . However, Haas and Hornig's spectra were obtained at 190 K. Thus allowing for temperature shifts, the agreement is excellent indeed. The peak height ratios of the present work are comparable to those of Haas and Hornig but their predicted 2 to 1 intensity ratio does not agree with the present work. This is probably due to their inability to resolve the HOD couplet and obtain a more quantitative measurement of band areas.

One interesting feature of the  $(\text{HOD})_2$  doublet is the 35 and 32

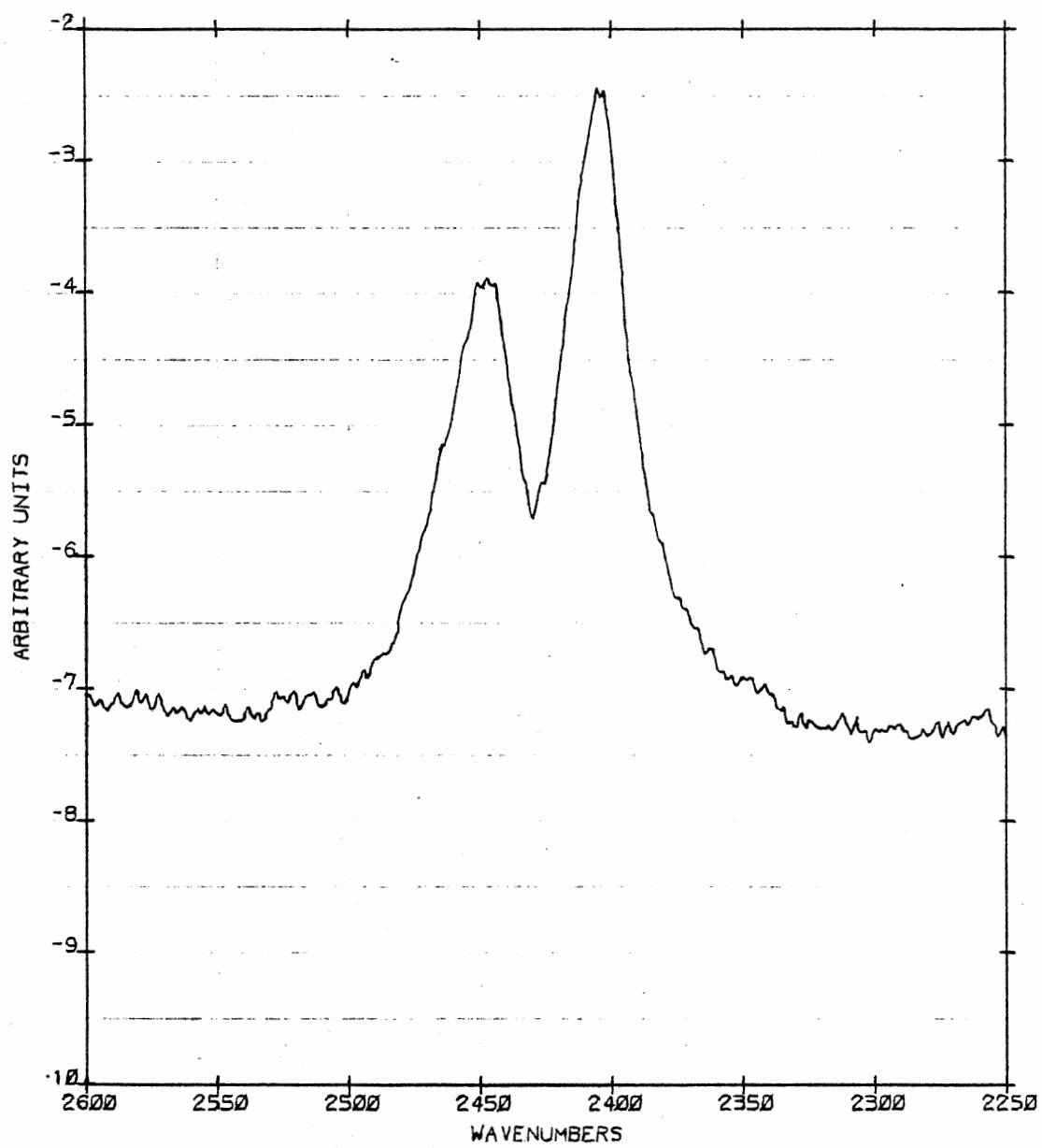


Figure 31. Coupled HOD Absorbance Spectrum of 145 K



$\text{cm}^{-1}$  bandwidths which are about 50% wider than the  $\text{D}_2\text{O}$  and uncoupled HOD bandwidths. Both the  $\text{D}_2\text{O}$  and uncoupled HOD oscillators have exactly one possible configuration if the surrounding  $\text{H}_2\text{O}$  network is ignored. However the  $(\text{HOD})_2$  oscillator has two possible configurations for the hydrogens valence bonded to the two coupled OD oscillators as shown in Figure 32a. This may explain the extra band breadth observed.

Occasionally in some  $(\text{HOD})_2$  spectra an extra wing or band was observed around  $2460 \text{ cm}^{-1}$ . The intensity would vary from a slight shoulder to an intense peak as shown in Figure 32b. Usually the shoulder would appear in initially collected spectra and quickly disappear as the exchange reaction occurred. A clue to the origin of this shoulder band was discovered when a few ice deposits made at 125 K instead of 130 K were more seriously affected by this band. A check of the lower 135 and 140 K runs revealed that when the band occurred at these temperatures it dissipated less rapidly than at the higher temperature kinetic runs. The  $\nu_3$  band of  $\text{D}_2\text{O}$  is the most intense feature of the initial ice deposit. If the deposit was slightly amorphous the  $\nu_3$  band would be slightly amorphous in character, and hence have a slightly greater bandwidth. But since the  $\text{D}_2\text{O}$  band being subtracted out was fully crystalline, in resolving the  $(\text{HOD})_2$  spectrum any residual amorphous character of  $\nu_3$  would not be subtracted and show up as a wing or band near the  $\nu_3$  frequency. As the deposit started to exchange at a higher temperature the ice deposit would quickly anneal and likewise any amorphous character of  $\nu_3$  would disappear. Therefore this wing was regarded as an amorphous residue and discounted as a spectrally significant feature when it occasionally

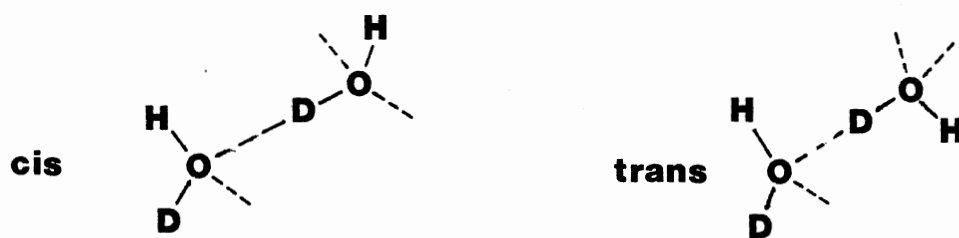


Figure 32a. Possible Coupled Configurations in Cubic Ice Lattice

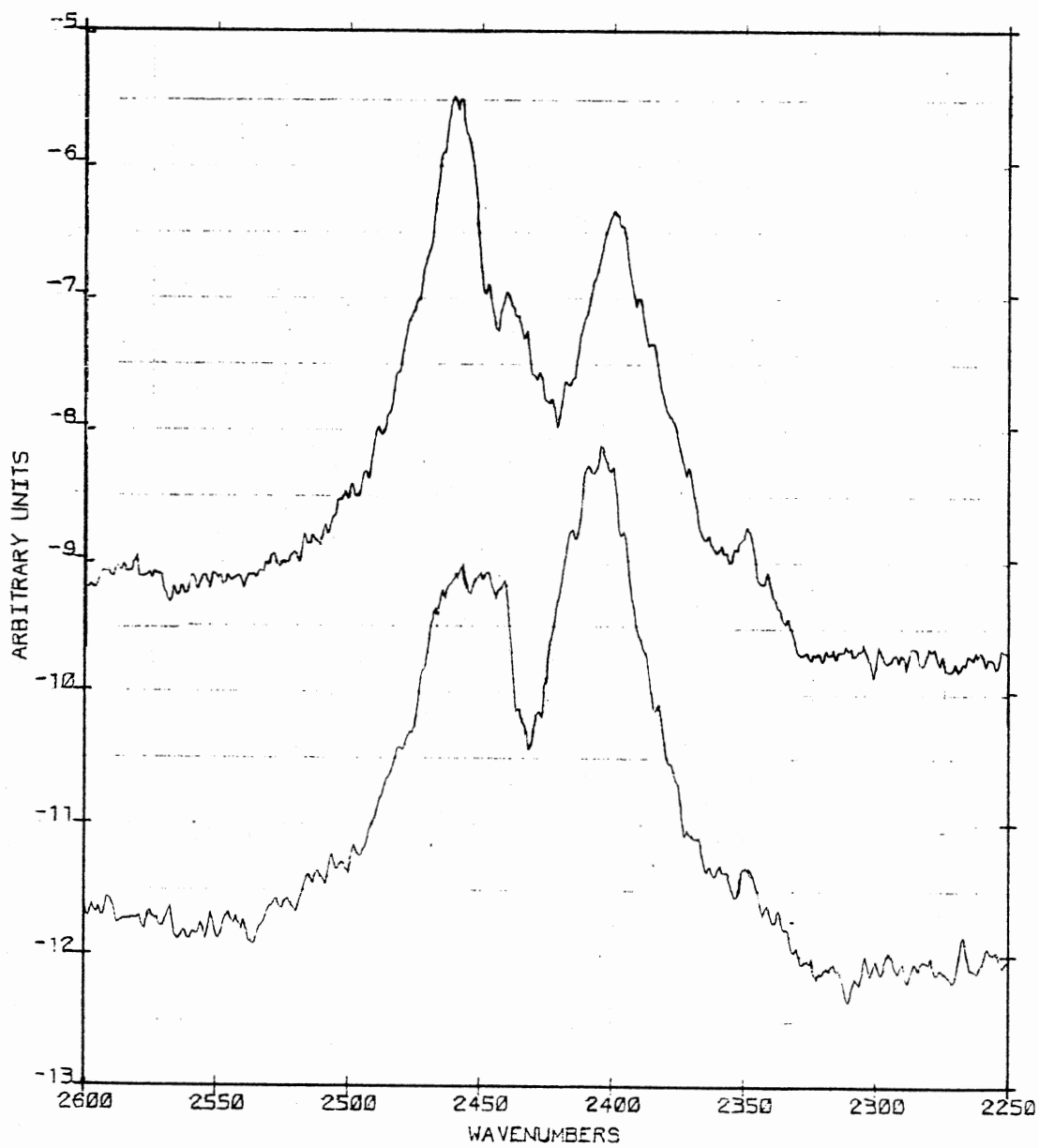


Figure 32b. Amorphous Induced Bands in Coupled HOD Bandshapes

appeared.

In Table IV the higher temperature (150 K, 145 K) deposits exchanged more completely. Likewise, the error for fitting the kinetic dataset was greatest for 150 K and 145 K and least for the 135 K temperature datasets. In the 150 K datasets and most of the 145 K datasets the reaction was followed through at least three half-lives of the  $D_2O$  concentrations. Thus they should provide a fairly rigorous test of the models proposed. The interesting facet of all the datasets is the extensive retardation of the reaction rates after about one to two  $D_2O$  half-lives. The reaction appears to follow an initial first order decay in  $D_2O$  concentration until the concentration level of all species is sufficiently low so that the feedback mechanisms present in the isotopic exchange mechanism becomes important. At that time large concentration changes appear to be stymied by the various back reactions.

#### Nonlinear Least Square Strategy

When fitting the one-equilibrium model to the kinetic data, the error became large near the completion of the data. Thus to improve the accuracy of the fit the latter data points of the 150, 145, and 140 K datasets were deleted. At least 1.5  $D_2O$  concentration half-lives worth of data were retained in each dataset. Though a more acceptable fit was obtained, the accuracy of the fit is questionable. Table VII lists the rate constants obtained from this fit and Figure 33 displays the calculated and experimental relative concentration values for the 150 K 6/7/82 dataset.

As observed in Figure 33 there is difficulty in fitting the

TABLE VII

LEAST SQUARE CALCULATED RATE CONSTANTS FOR ONE AND TWO EQUILIBRIUM MODELS

Model pa	One equil.		Two equil.		Two equil.		Two equil.	
	.73684		.73684		.73684		.73684	
Initial HOD <sub>21</sub> <sup>b</sup>	-		0.0		0.25		0.5	
data	shortened		full		full		full	
Dataset	K <sub>trans</sub>	K <sub>rot</sub>	K <sub>trans</sub>	K <sub>rot</sub>	K <sub>trans</sub>	K <sub>rot</sub>	K <sub>trans</sub>	K <sub>rot</sub>
<u>150 K</u>								
2/2/82	.21300	.21830	.17495	.28568	.17496	.36879	.17430	.52666
RMSD <sup>c</sup>	1.354		.5590		.5321		.5028	
6/7/82	.13660	.19843	.11100	.29613	.11115	.42103	.11099	.74646
RMSD	1.219		.5713		.5566		.5402	
6/9/82	.079542	.20734	.068384	.31804	.068654	.50200	.068748	1.4265
RMSD	.7247		.5357		.5345		.5310	
<u>145 K</u>								
1/29/82	.079151	.077733	.059852	.089019	.059440	.10920	.058936	.14205
RMSD	1.173		.3984		.38037		.3633	
5/10/82	.042268	.061057	.033598	.075077	.034035	.11292	.034340	.23085
RMSD	.7774		.4493		.4424		.4307	
5/24/82	.082043	.069061	.061907	.074030	.061634	.096254	.061115	.13880
RMSD	.7908		.2278		.2082		.1877	
6/3/83	.042662	.068473	.032988	.088568	.033011	.11470	.032995	.16417
RMSD	.6586		.4111		.4029		.3940	
<u>140 K</u>								
5/11/82	.018463	.023522	.014124	.023192	.014096	.028898	.014048	.038439
RMSD	.6210		.3752		.3678		.3605	
5/25/82	.011992	.017621	.0092350	.019506	.0092242	.025360	.0092230	.036234
RMSD	.6657		.2383		.23270		.2279	

TABLE VII (Continued)

Model F <sup>a</sup>	One equil. .73684		Two equil. .73684		Two equil. .73684		Two equil. .73684	
	shortened		full		full		full	
Initial HOD <sub>21</sub> <sup>b</sup> data Dataset	K <sub>trans</sub>	K <sub>rot</sub>	K <sub>trans</sub>	K <sub>rot</sub>	K <sub>trans</sub>	K <sub>rot</sub>	K <sub>trans</sub>	K <sub>rot</sub>
135 K								
5/15/82	.0042115	.0040228	.0037078	.0033044	.0037055	.0040094	.0037008	.0050681
RMSD	.3258		.1997		.1961		.1926	
5/26/82	.0037345	.0043476	.0031359	.0037833	.0031316	.0046006	.0031249	.0058465
RMSD	.3059		.1491		.1489		.1484	
6/11/82	.0038900	.0047298	.0033464	.0039339	.0033437	.0046827	.0033392	.0057573
RMSD	.4305		.2068		.2022		.1979	

<sup>a</sup>Frequency factor;  $F = \left( \frac{\nu_H + 2\nu_D}{\nu_D + 2\nu_H} \right)$ .

<sup>b</sup>Initial HOD<sub>21</sub> determined by  $[\text{HOD}_{21}]_{\text{initial}} = (\text{table value shown}) \cdot [\text{HOD}_{12}]_{\text{Initial}}$ .

<sup>c</sup>Root mean square deviation of nonlinear least square fit.

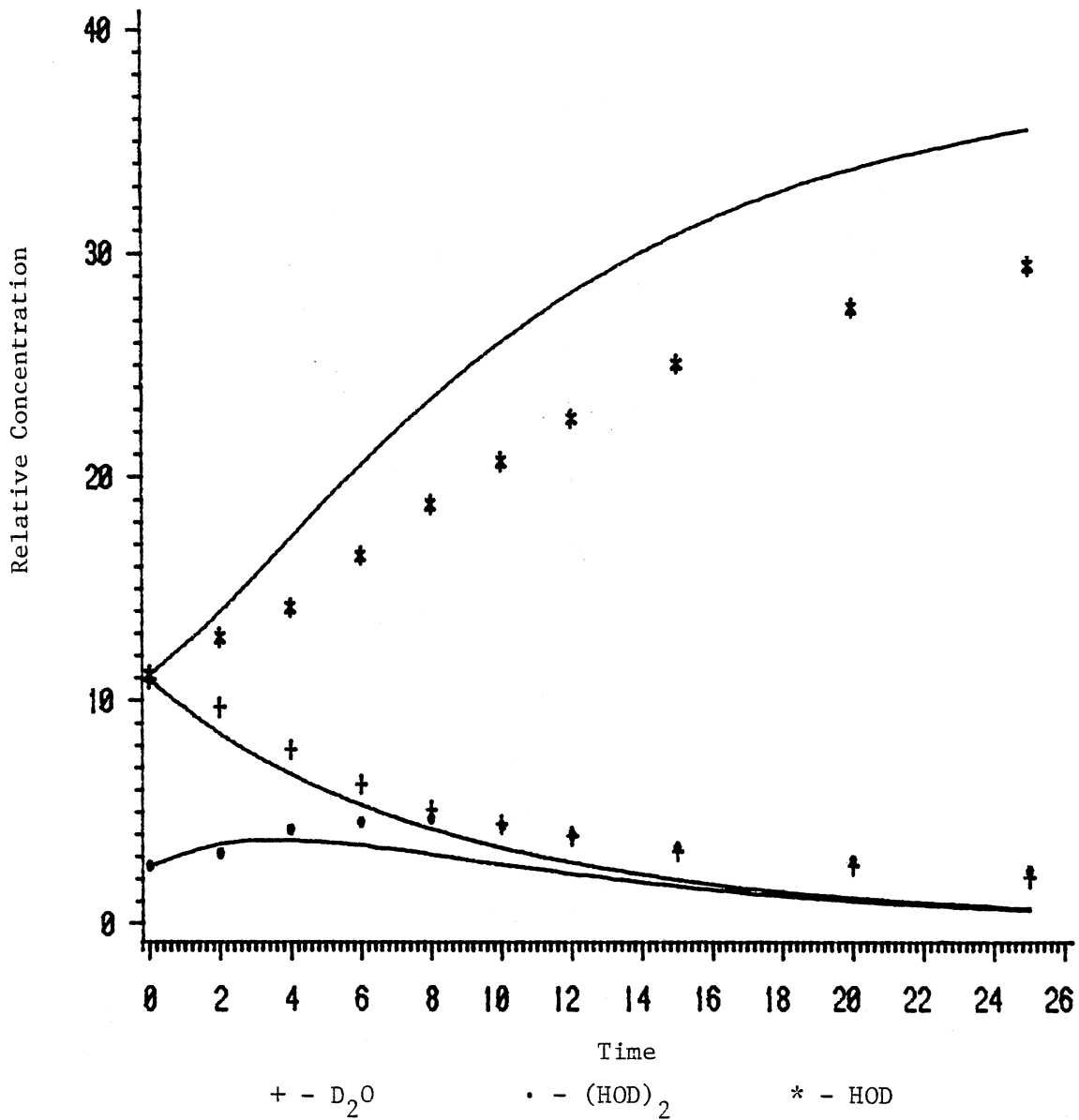


Figure 33. Experimental and One Equilibrium Model Calculated Relative Concentration Curves for 6/7/82 150 K Data

$(\text{HOD})_2$  concentration at the longer time intervals. The lack of back exchange to  $(\text{HOD})_2$  favors a more rapid depletion of  $(\text{HOD})_2$ , thus more quickly eliminating any feedback from  $(\text{HOD})_2$  to  $\text{D}_2\text{O}$ . As a result both the  $(\text{HOD})_2$  and  $\text{D}_2\text{O}$  concentrations fall off too rapidly.

For the two-equilibrium model both the above described shortened datasets and the full datasets were fitted. The full datasets yielded the most consistent results and hence were used for the remainder of the calculations. Several fits were made, each with different initial separated HOD concentrations. Table VII lists the transfer and rotational rate constants found for the three best runs. They are initial separated HOD values of 0.0,  $0.25 \times (\text{HOD})_2$  initial, and  $0.5 \times (\text{HOD})_2$  initial. The root mean square deviation of the fit for each dataset is included beneath each rate constant entry. The 150 K and 145 K determined transfer rate constants show the greatest error spread, while the rotational rate constants seem to be more consistent. This is not entirely unexpected. The concentration of hydronium ions in water at  $25^\circ\text{C}$  is about  $1.0 \times 10^{-7}$  moles/liter. So the concentration of ionic defects in ice at 145 K should be even less. This concentration of defects is less than two parts per billion. It is very difficult to obtain water of this purity. Accordingly the transfer rate constant variance at fixed temperatures is attributed to the ionic defect's sensitivity to small changes in the concentration of impurities in the ice such as arises from normal operation of a vacuum system. So the transfer rate constants listed are assumed to arise from a pure ice deposition situation perturbed by minute contaminations. The fact that the transfer rate constants are so

sensitive to small contamination changes provides evidence that the ice deposition is pure enough that contamination concentrations are of the same order as the hydronium ion defect concentration. If the contamination was greater, the trapping would be large enough to significantly lower the transfer rates and suppress  $(\text{HOD})_2$  formation as illustrated by Devlin et al. (46, 47).

Not surprisingly the rotational rates appear to be more consistent. This may reflect their insensitivity to impurity doping as suggested by Devlin et al. (47).

The transfer rate constants are largely insensitive to the initial separated HOD concentration and corresponding differing rotational rate constants. The rotational rates are affected by the initial separated HOD concentrations as Table VII illustrates. The  $.5 \times (\text{HOD})_2$  initial concentration data has the smallest RMSD of fit but yet this is misleading. The least square analysis showed the .5 calculation to be the least correlated with respect to activation energy plots. In addition for some rate constants the standard error associated with the rotational rate constant was 50% greater than for the 0.0, and .25 calculations.

Figures 34-45 show calculated and experimental concentration versus time curves for each dataset. All three two-equilibrium calculations gave practically identical calculated concentration values so the .25 initial  $\text{HOD}_{21}$ ,  $F = .7368$  rate constant calculated concentration values were used. For each graph the symbols represent the experimental data and the lines the corresponding calculated values. Representative estimated error bars are shown for Figure 34 but are omitted from the rest for clarity.



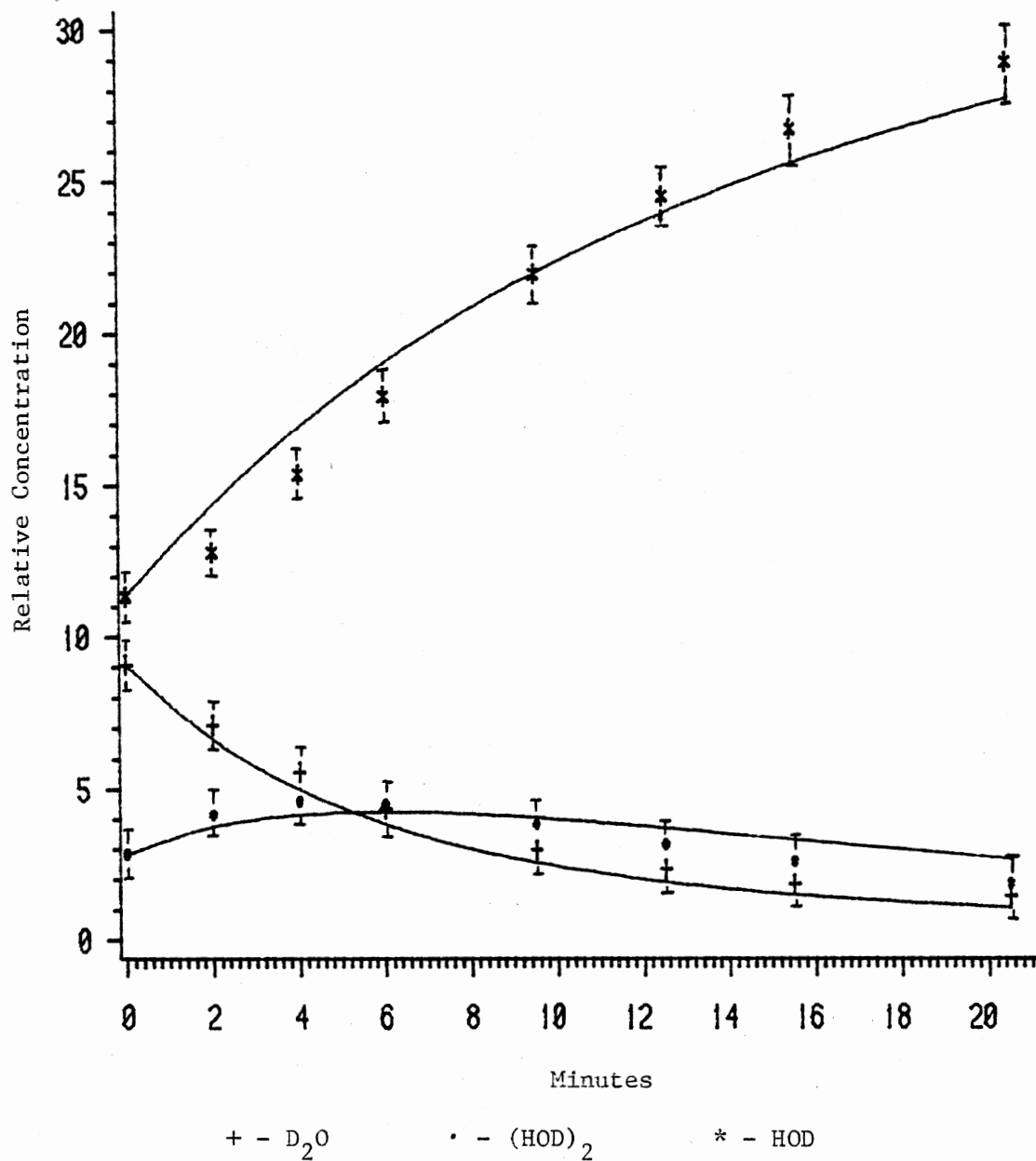


Figure 34. Experimental and Two Equilibrium Model Calculated Relative Concentration Curves for 2/2/82 150 K Data

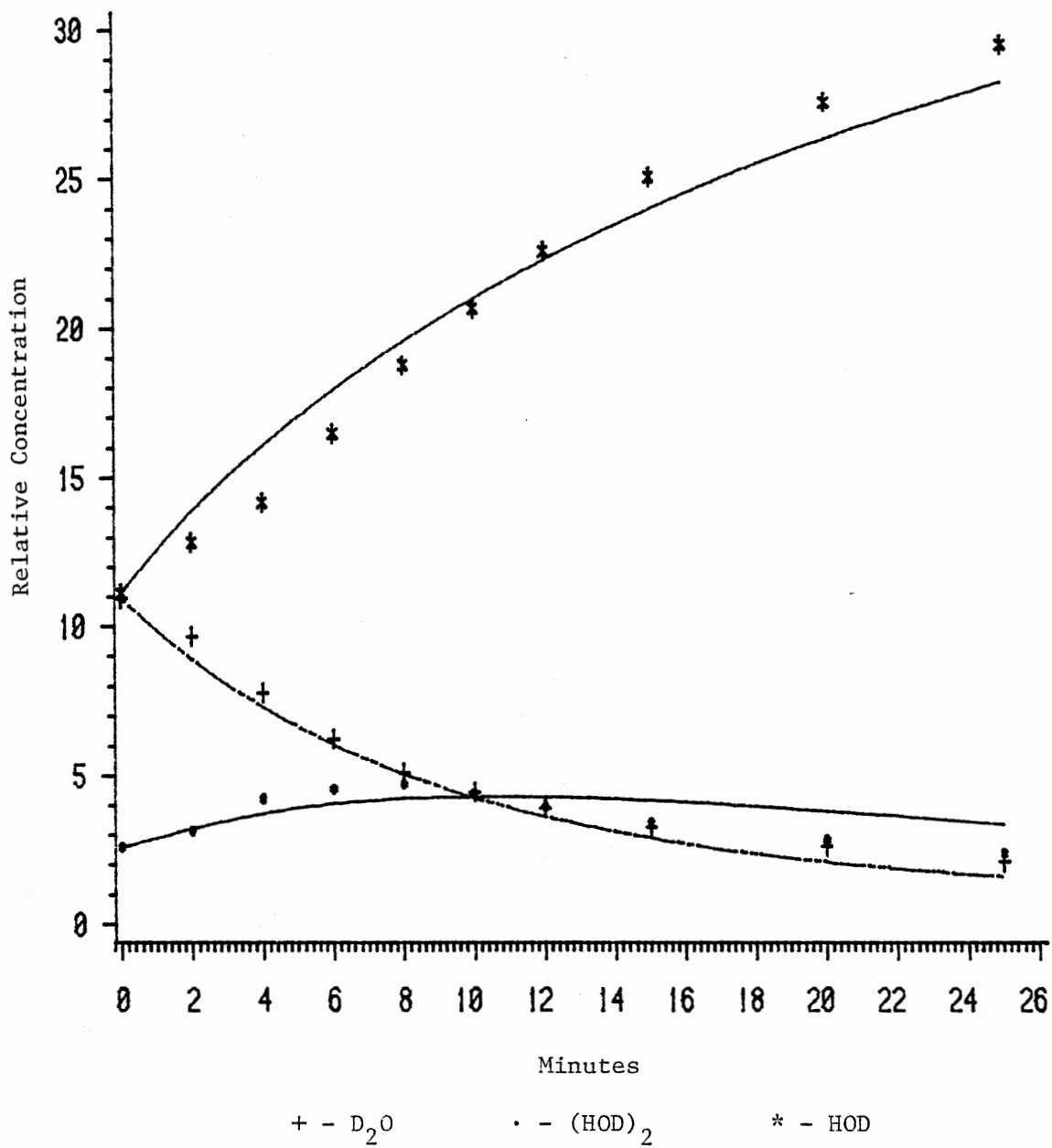


Figure 35. Experimental and Two Equilibrium Model Calculated Relative Concentration Curves for 6/7/82 150 K Data

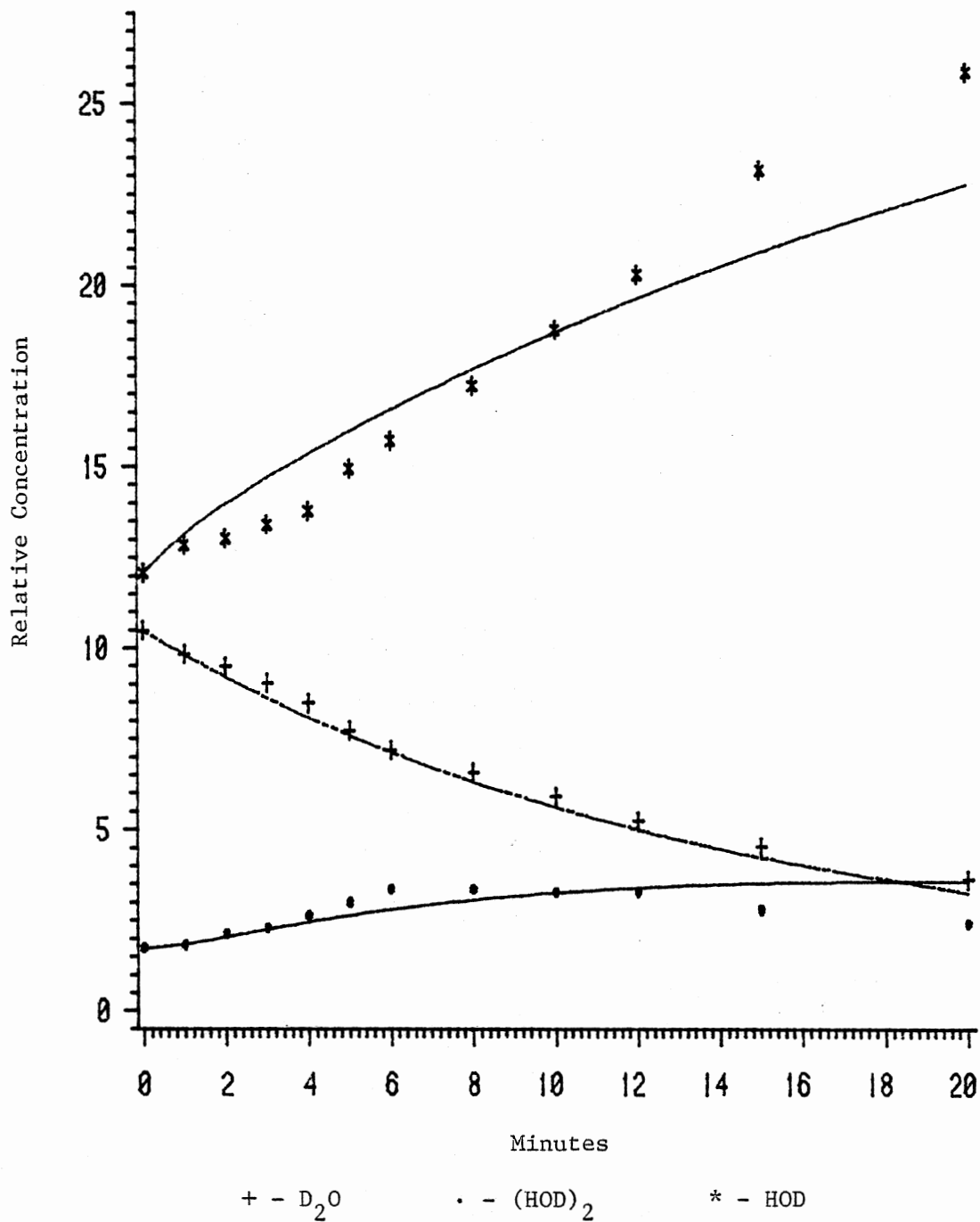


Figure 36. Experimental and Two Equilibrium Model Calculated Relative Concentration Curves for 6/9/82 150 K Data

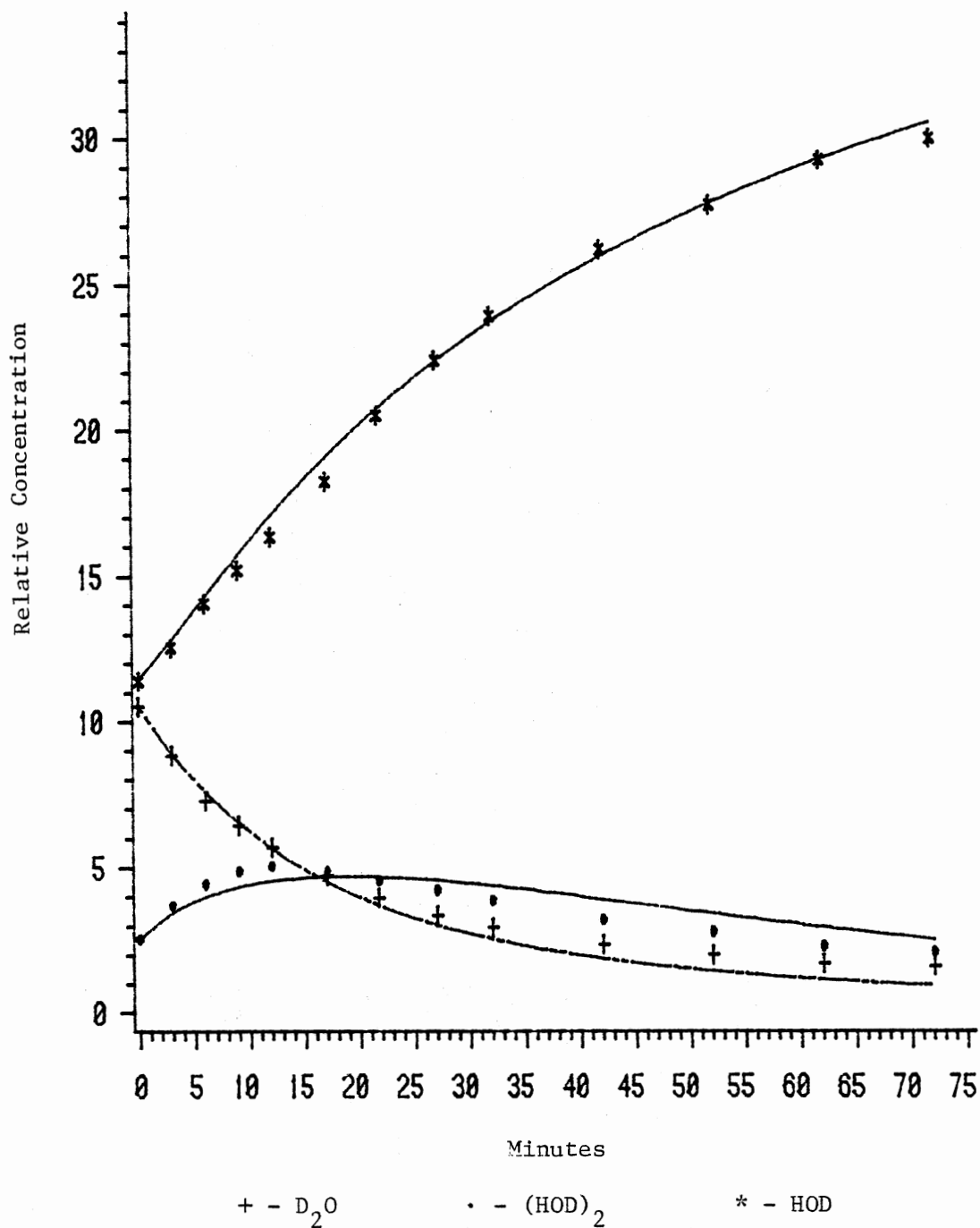


Figure 37. Experimental and Two Equilibrium Model Calculated Relative Concentration Curves for 1/29/82 145 K Data

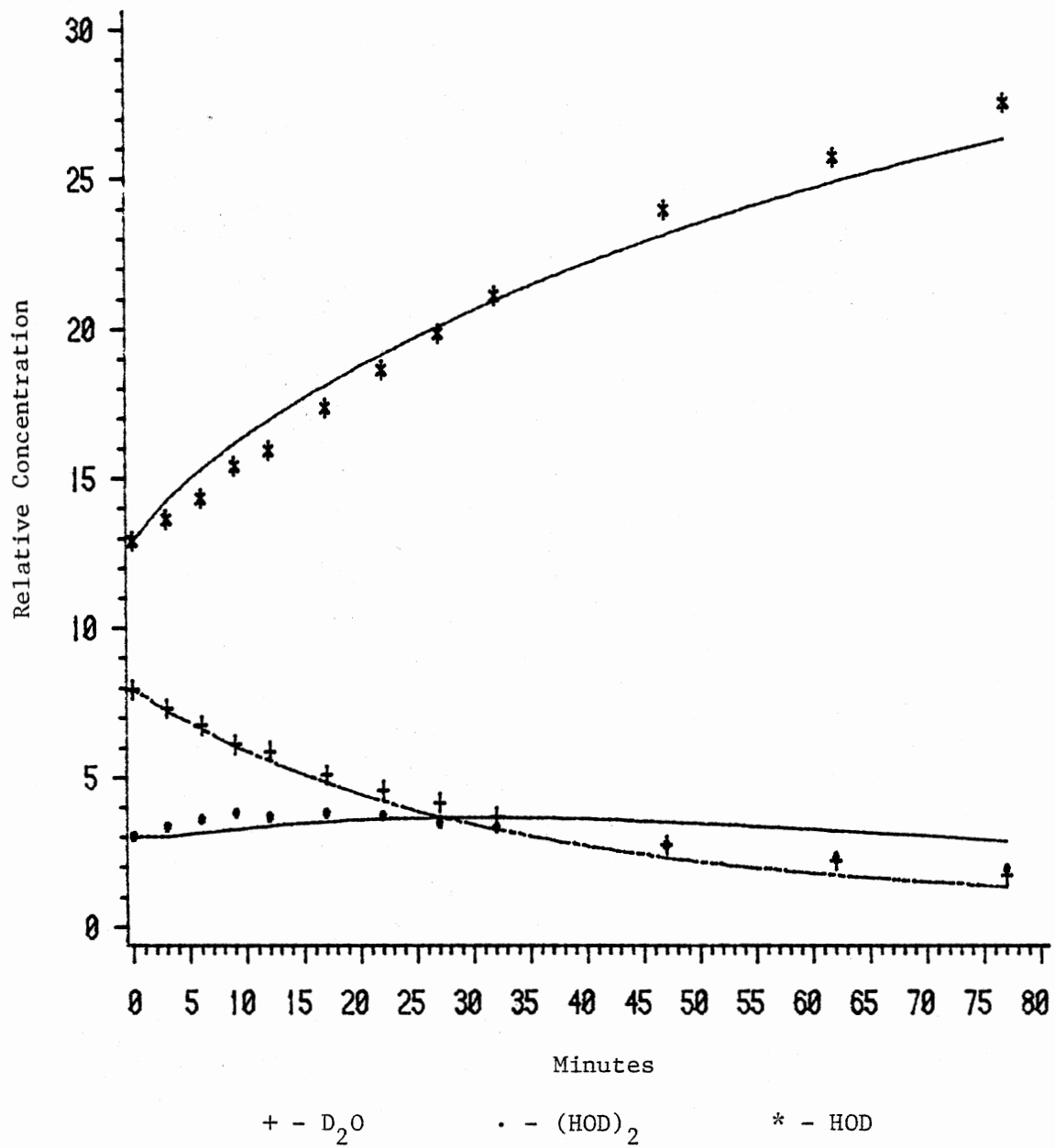


Figure 38. Experimental and Two Equilibrium Model Calculated Relative Concentration Curves for 5/10/82 145 K Data

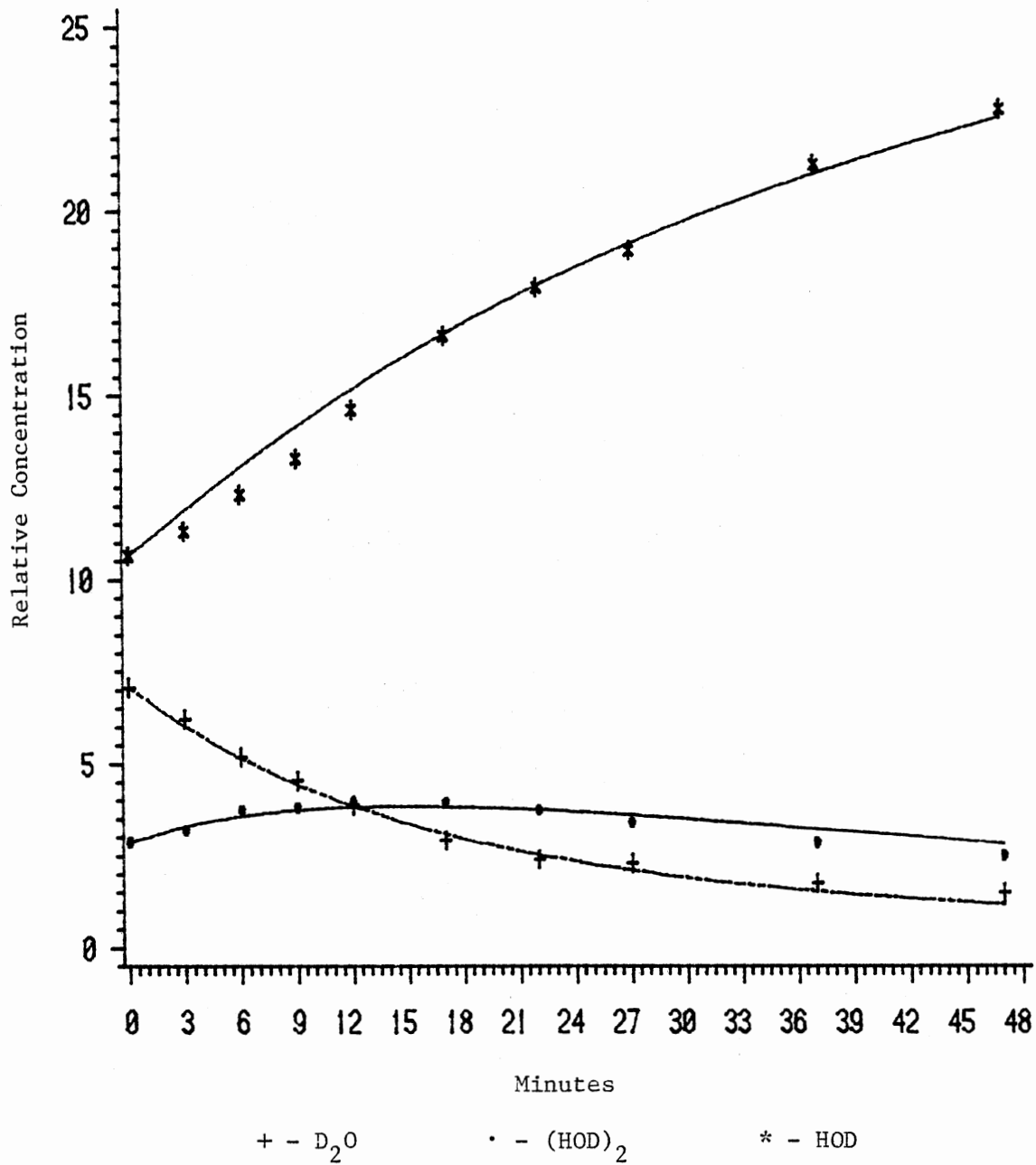


Figure 39. Experimental and Two Equilibrium Model Calculated Relative Concentration Curves for 5/24/82 145 K Data

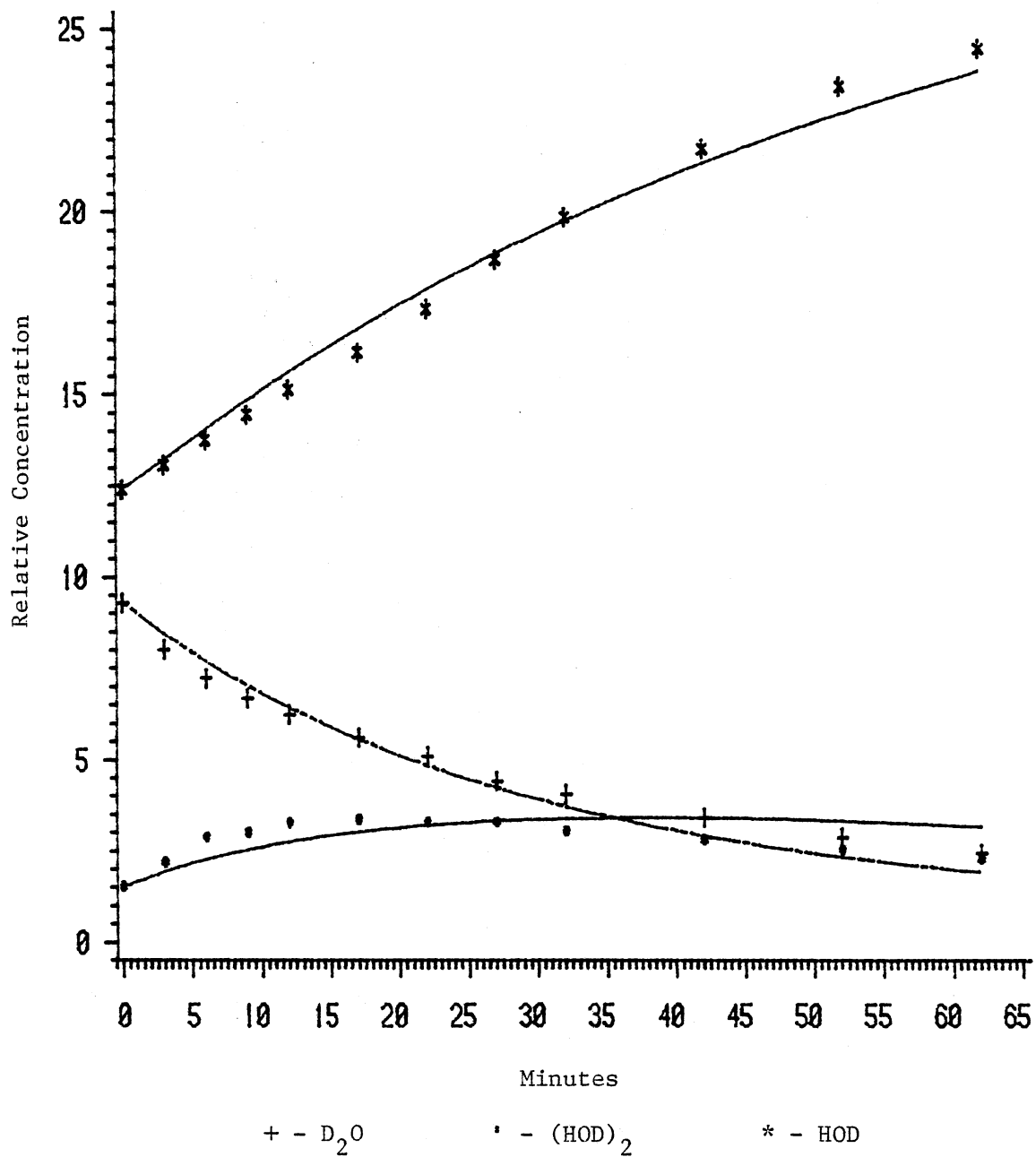


Figure 40. Experimental and Two Equilibrium Model Calculated Relative Concentration Curves for 6/3/82 145 K Data

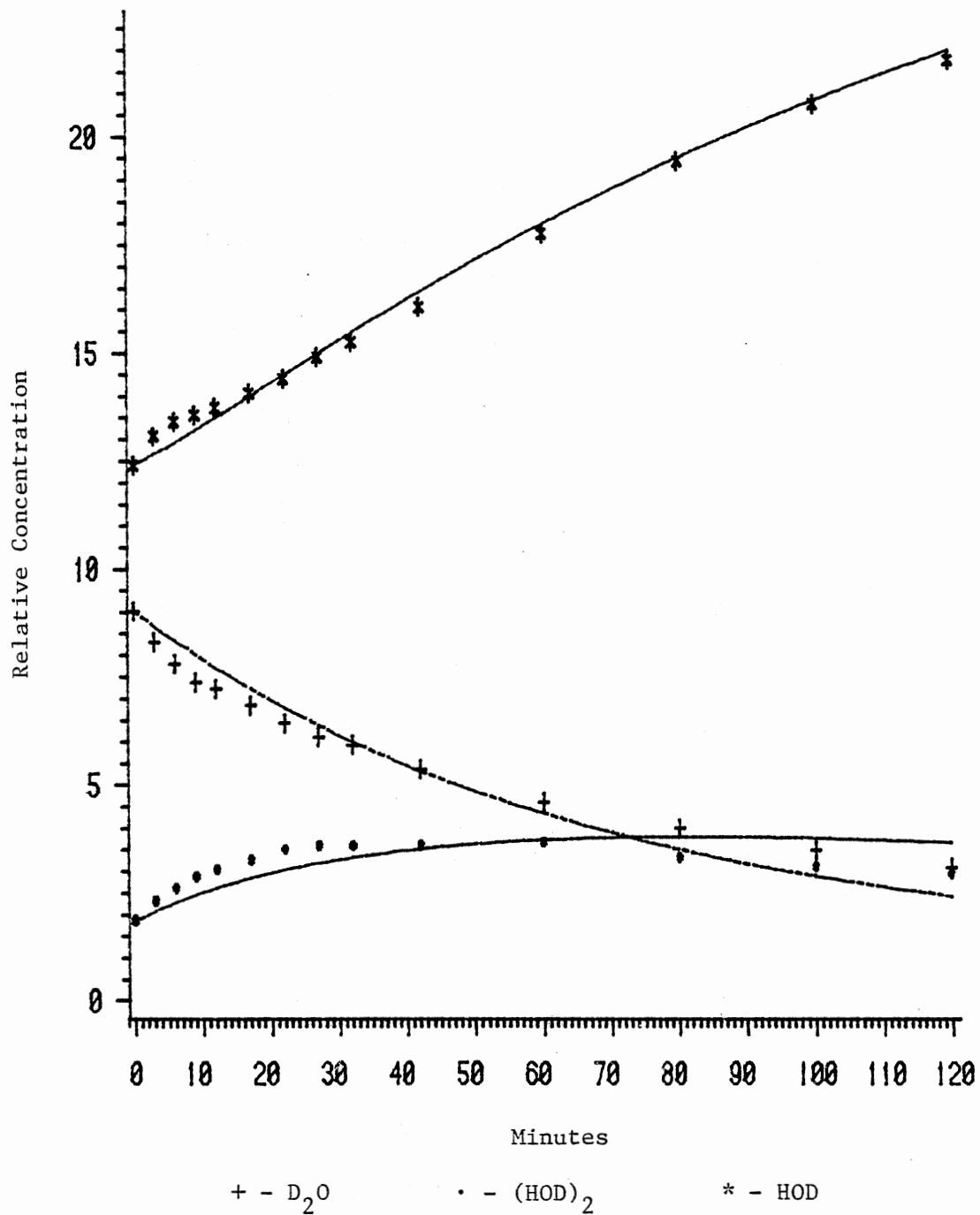


Figure 41. Experimental and Two Equilibrium Model Calculated Relative Concentration Curves for 5/4/82 140 K Data



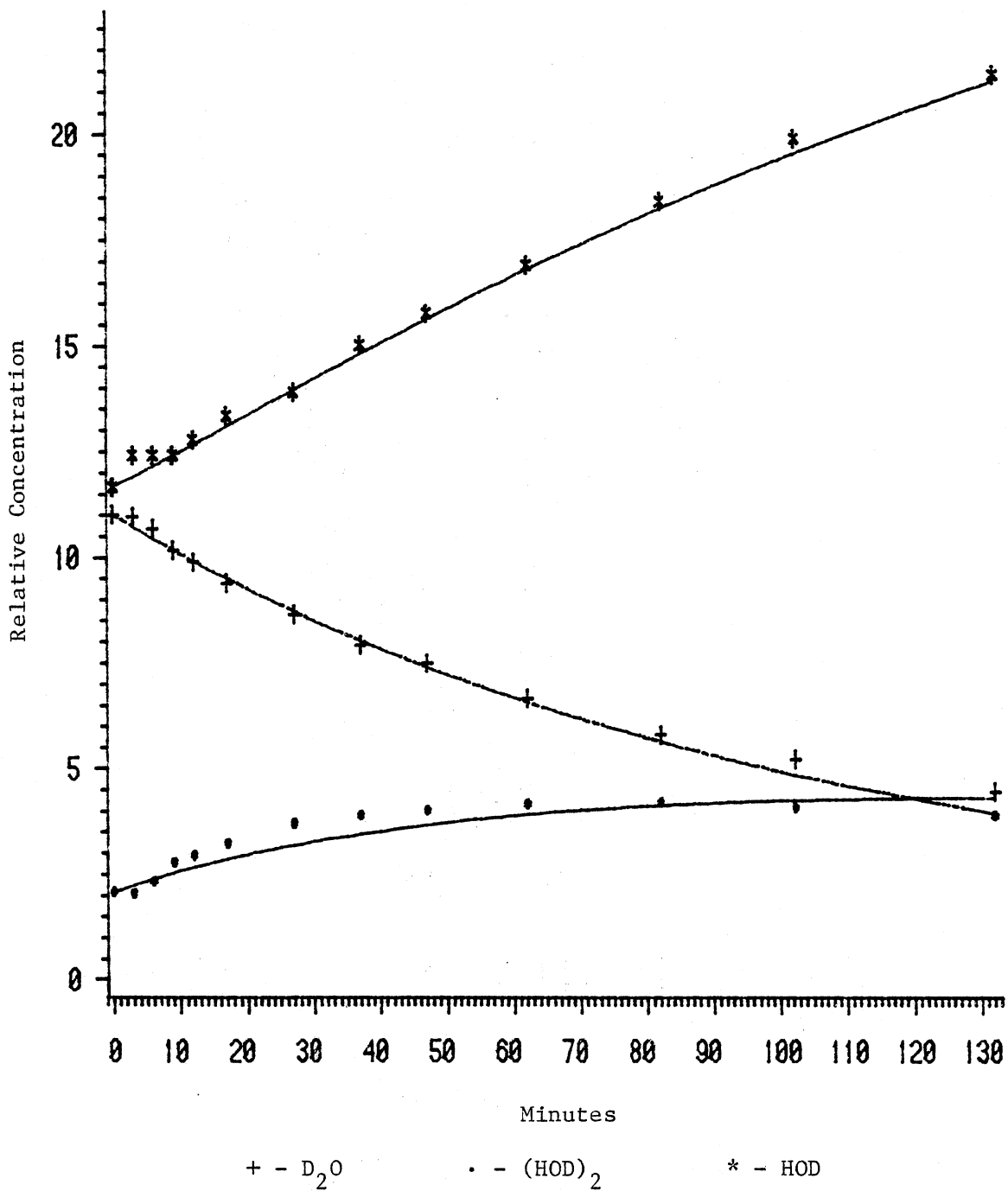


Figure 42. Experimental and Two Equilibrium Model Calculated Relative Concentration Curves for 5/25/82 140 K Data

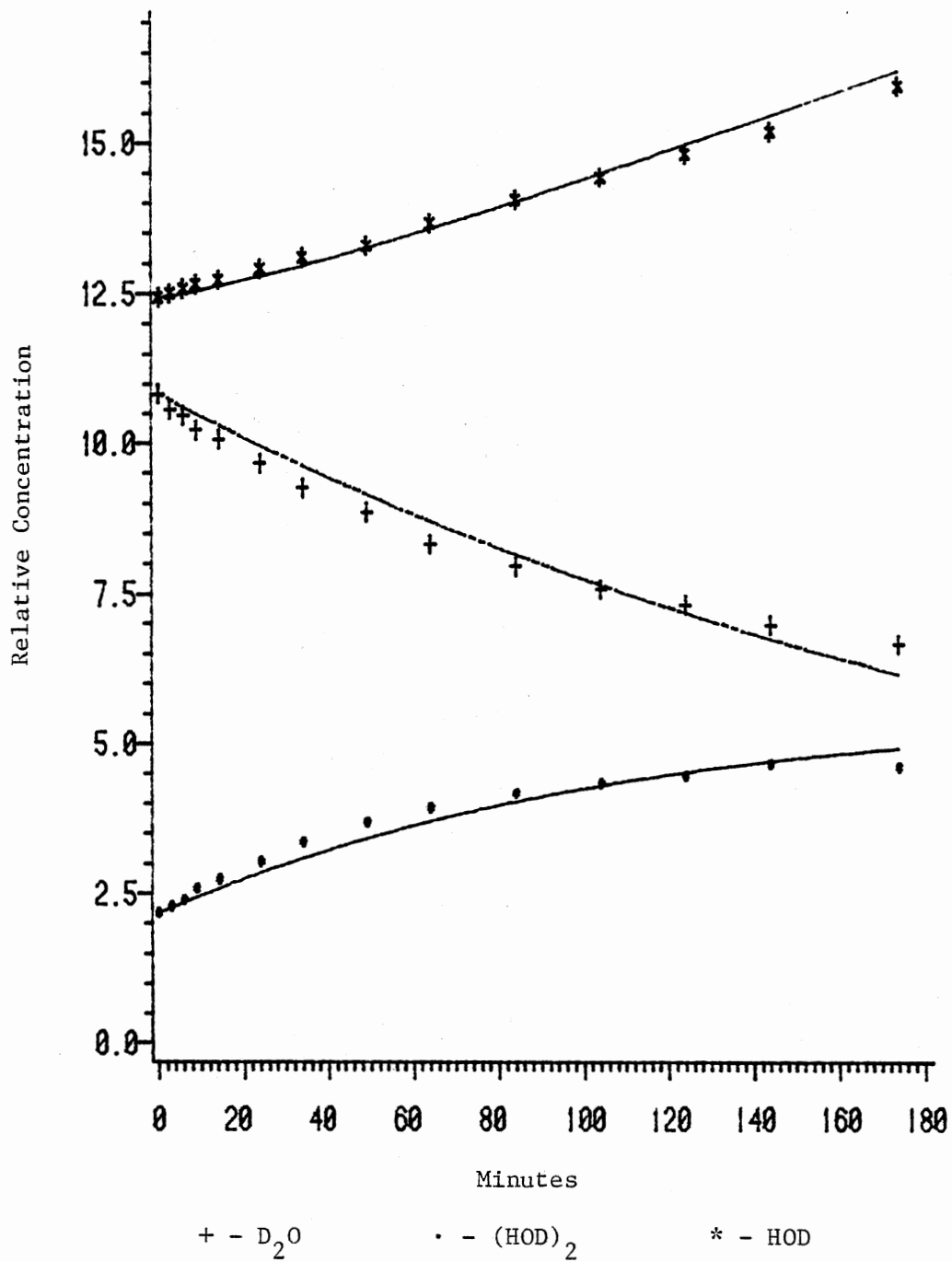


Figure 43. Experimental and Two Equilibrium Model Calculated Relative Concentration Curves for 5/15/83 135 K Data

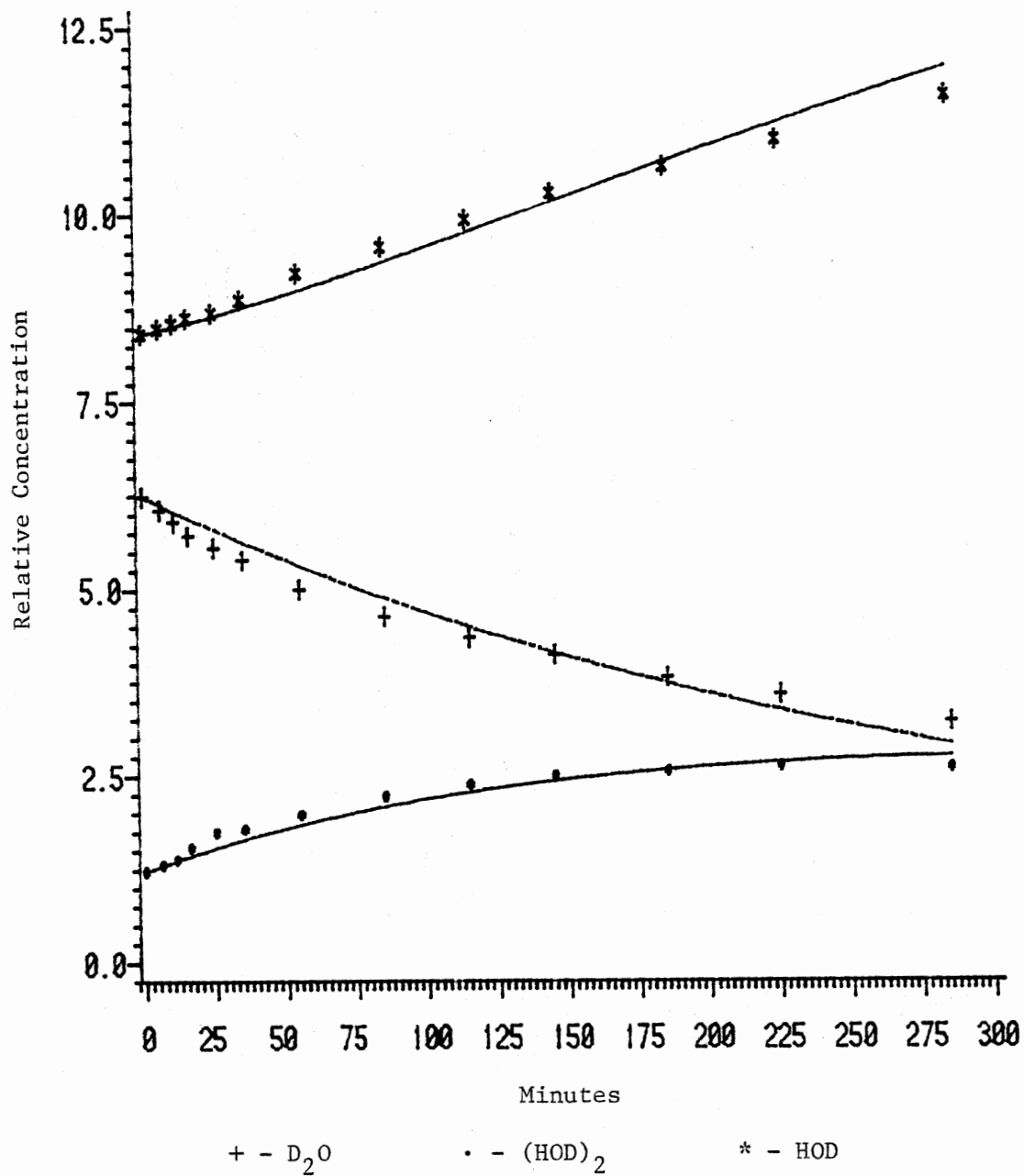


Figure 44. Experimental and Two Equilibrium Model Calculated Relative Concentration Curves for 5/26/82 135 K Data

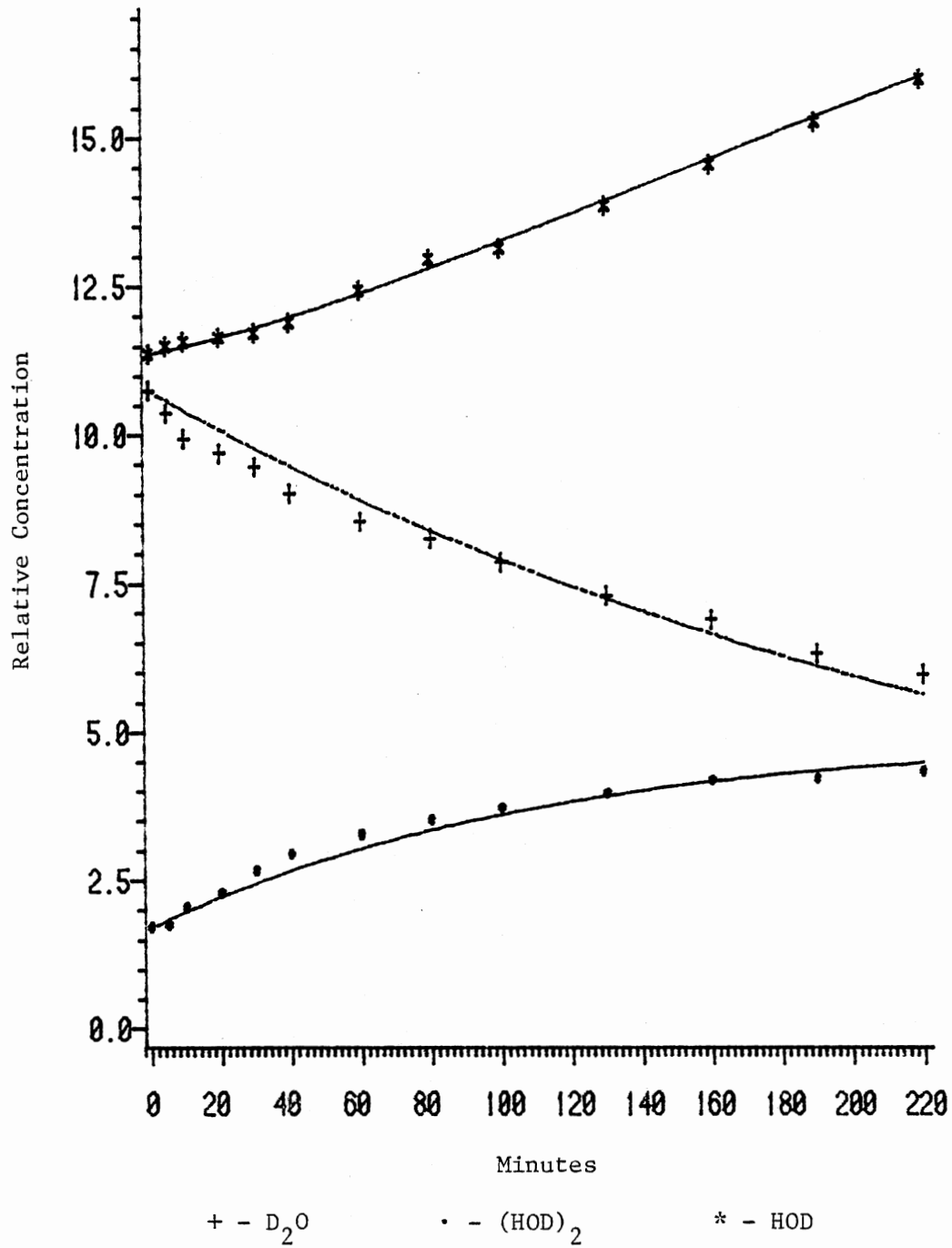


Figure 45. Experimental and Two Equilibrium Model Calculated Relative Concentration Curves for 6/11/82 135 K Data

Figures 46-57 show the calculated  $D_2O_{11}$  and  $2HOD_{nm}$  configuration concentrations determined from the two-equilibrium model used for Figures 34-45. The top decaying curve is the  $D_2O$  configuration relative concentration. The increasing curve just below  $D_2O_{11}$  corresponds to the  $2HOD_{22}$  and higher (fully separated) configuration concentration. Note that the concentration is less than one half the corresponding uncoupled HOD concentrations seen in Figures 34-45. This is because the  $2HOD_{22}$  concentration includes both HOD units rather than a single HOD unit. Also, the  $2HOD_{21}$  configuration concentration (separated HOD) is shown separately, rather than added into the  $2HOD_{22}$  concentration to form the uncoupled HOD shown in Figures 34-45. The lowest intermediary curve in Figures 46-57 is the  $2HOD_{21}$  configuration relative concentration (separated HOD). The intermediary curve just above it is the  $2HOD_{12}$  configuration concentration (coupled HOD). The initial value of  $2HOD_{21}$  is always one quarter the  $2HOD_{12}$  concentration as defined in the model. It is interesting to note that the  $2HOD_{21}$  concentration never exceeds the  $2HOD_{12}$  concentration in the data shown, but quickly establishes a slightly lower concentration. This is probably due to the equal magnitude of the forward and reverse rotational rate constants between these two configurations.

Compared with the one-equilibrium model, the two-equilibrium model is a significantly better fit, especially at the completion of the exchange reaction. The fit of  $(HOD)_2$  is improved with deviations occurring at two to three  $D_2O$  half-lives instead of one half-life as in the one-equilibrium model. The two-equilibrium model tends to overestimate the  $(HOD)_2$  concentration and underestimate the

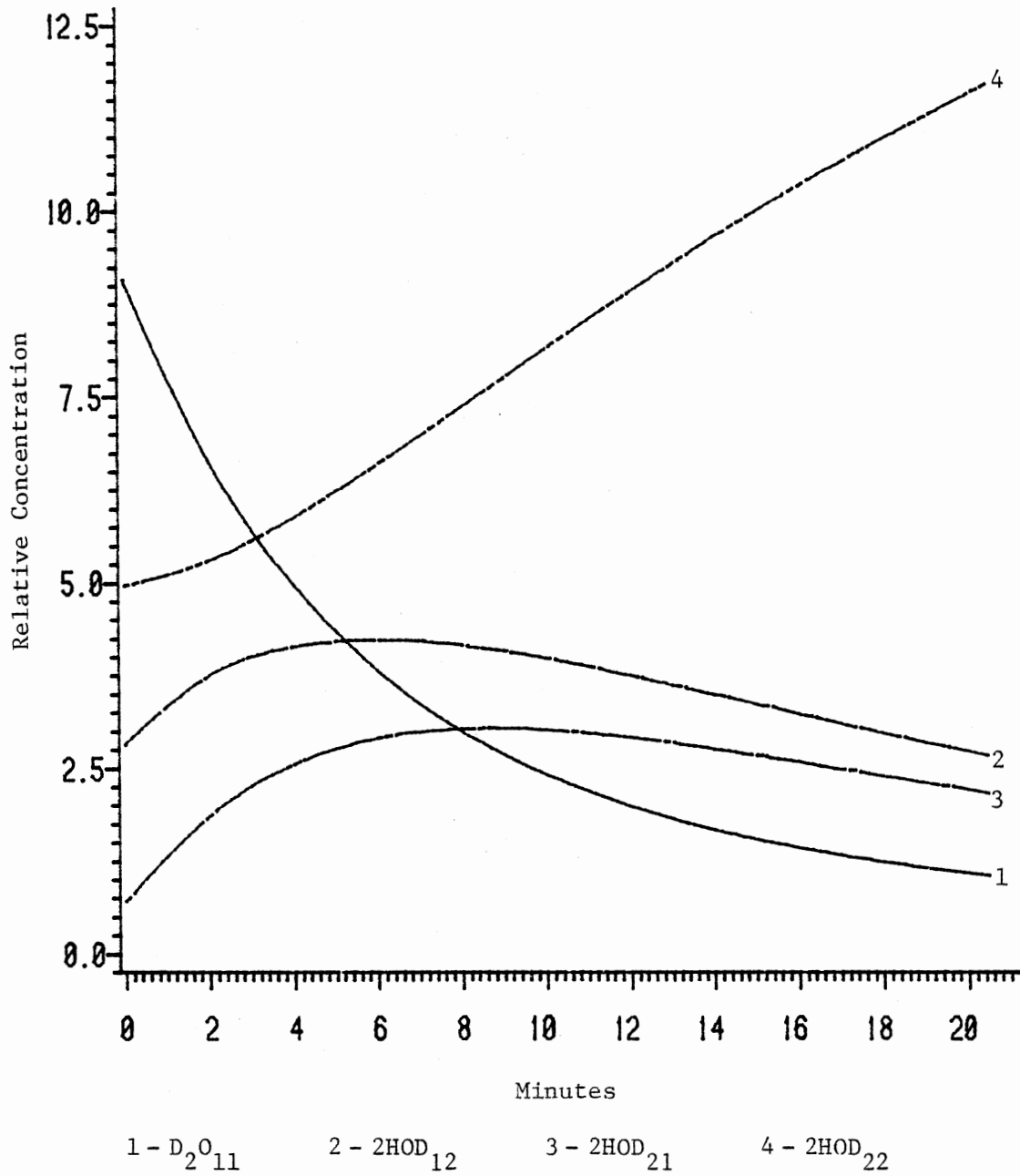


Figure 46. Calculated  $D_2O_{11}$  and  $2HOD_{nm}$  Configuration Concentration Curves for 2/2/82 150 K Data

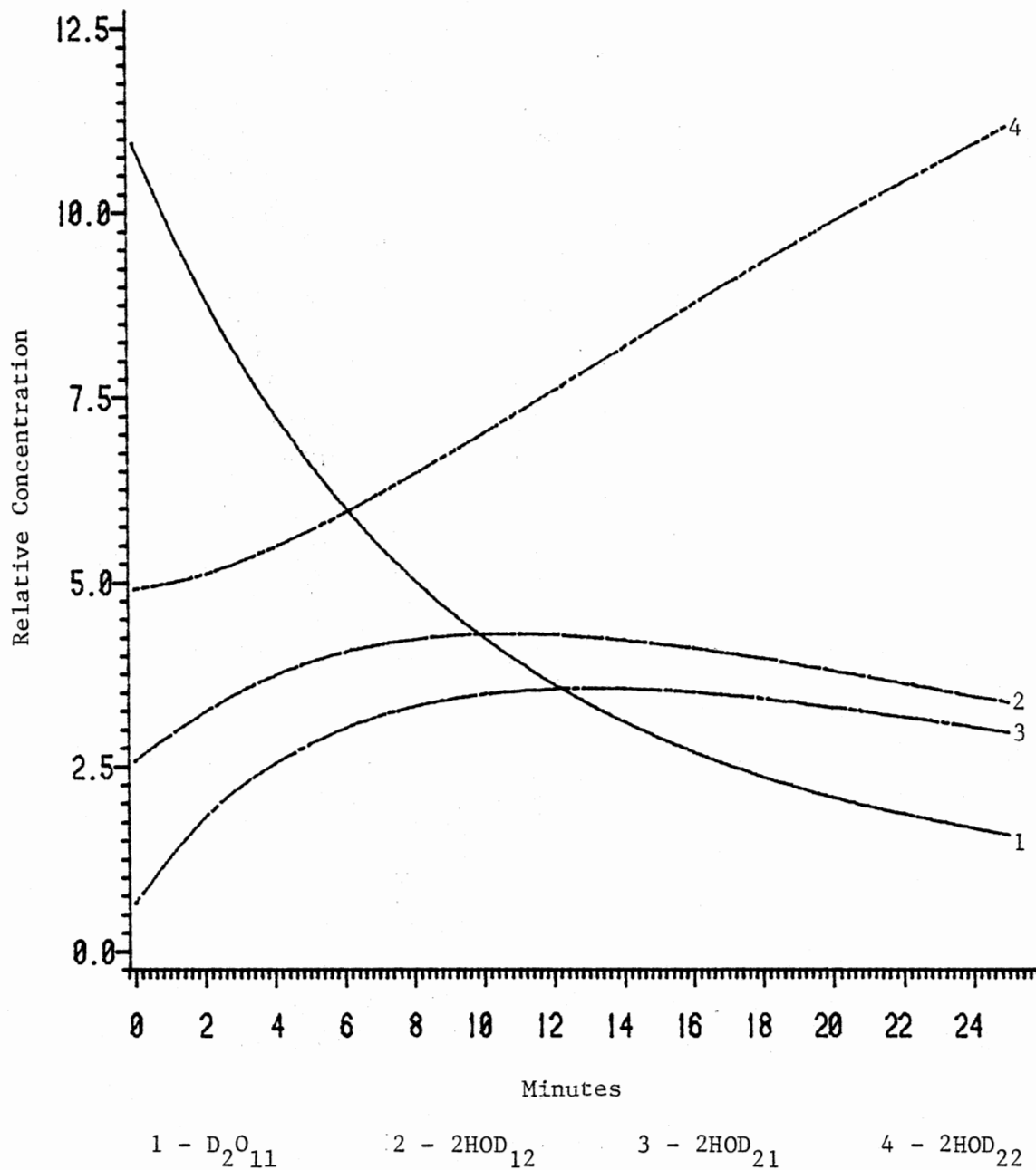


Figure 47. Calculated  $D_2O_{11}$  and  $2HOD_{nm}$  Configuration Concentration Curves for 6/7/82 150 K Data

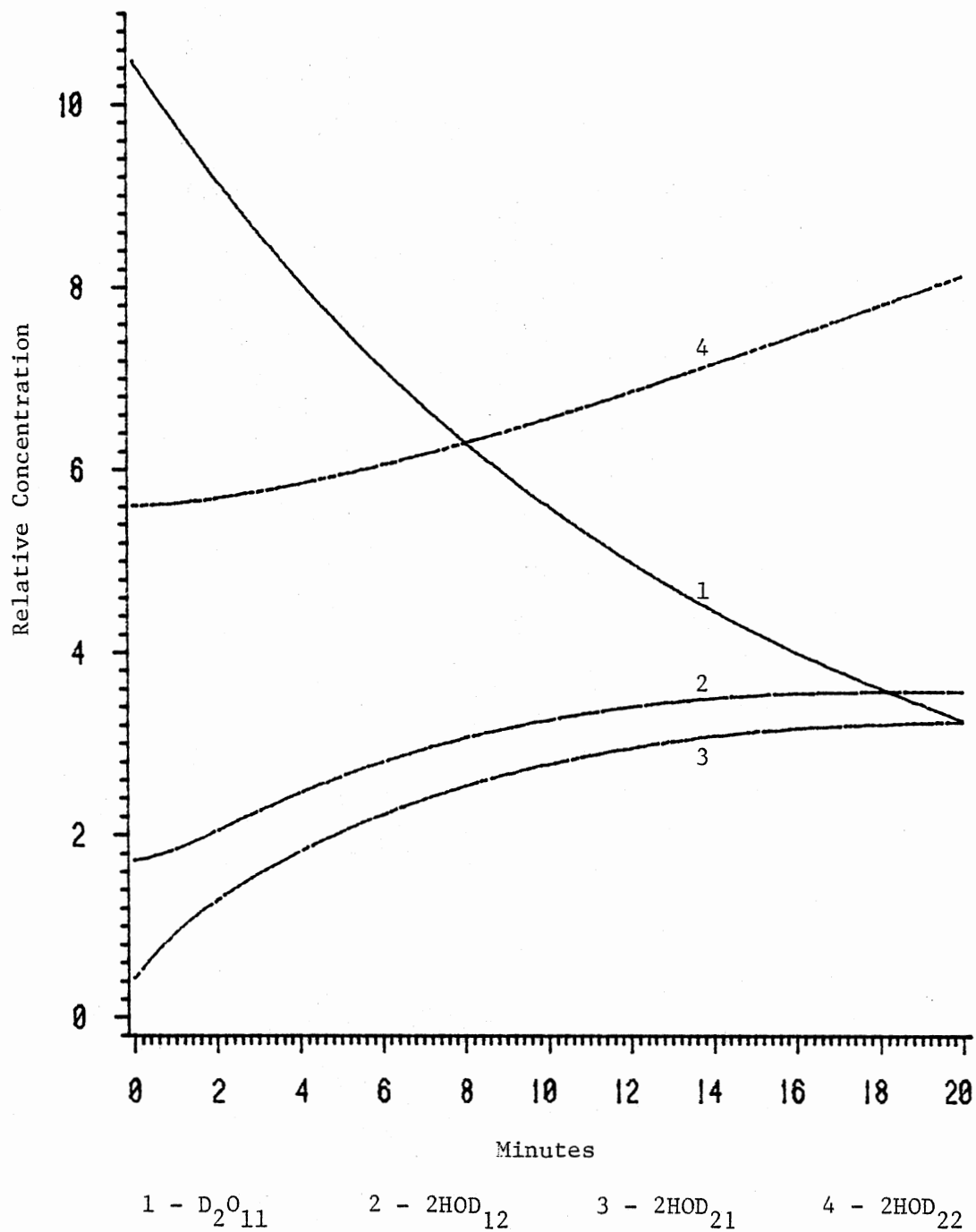


Figure 48. Calculated D<sub>2</sub>O<sub>11</sub> and 2HOD<sub>nm</sub> Configuration Concentration Curves for 6/9/82 150 K Data



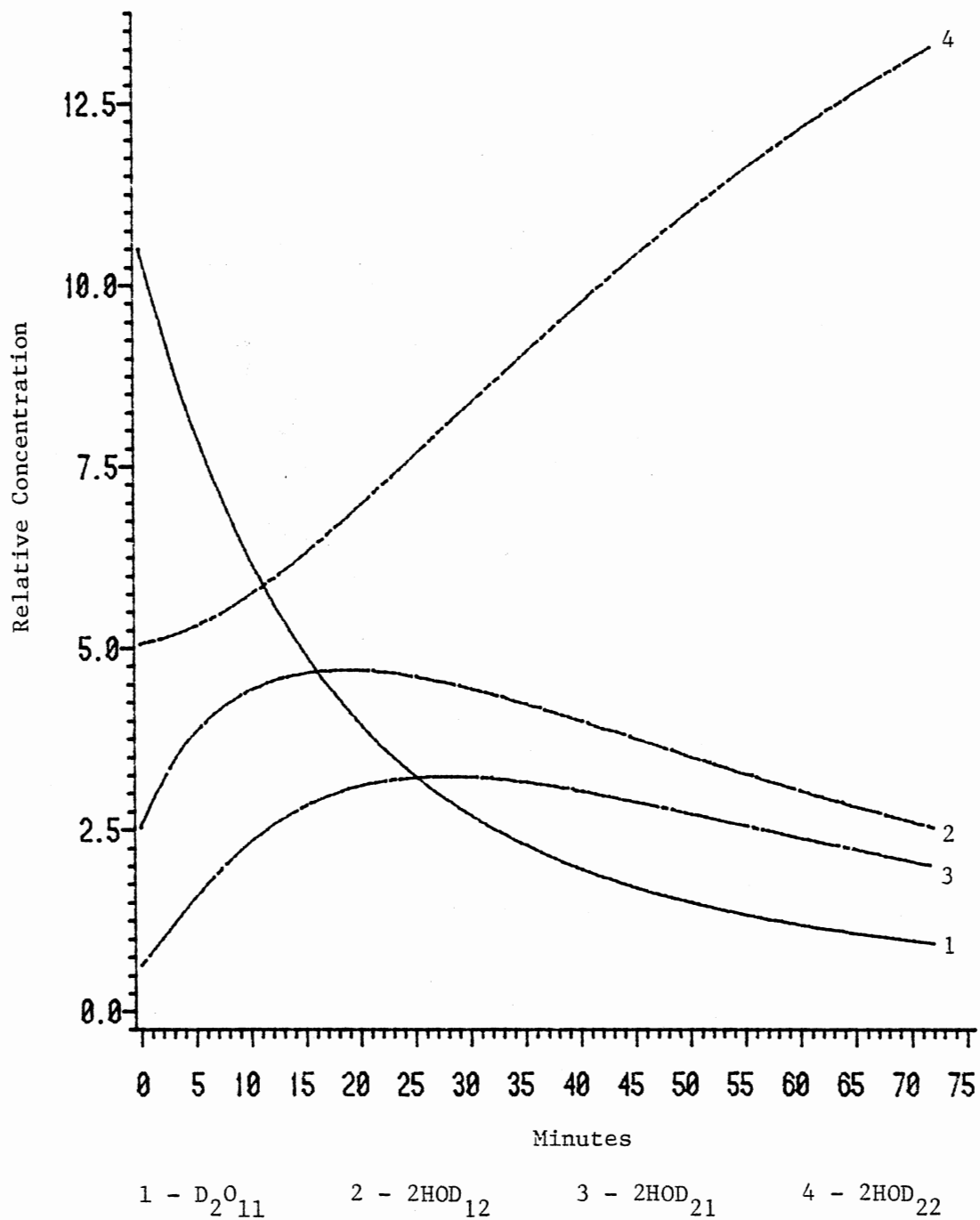


Figure 49. Calculated  $D_2O_{11}$  and  $2HOD_{nm}$  Configuration Concentration Curves for 1/29/82 145 K Data

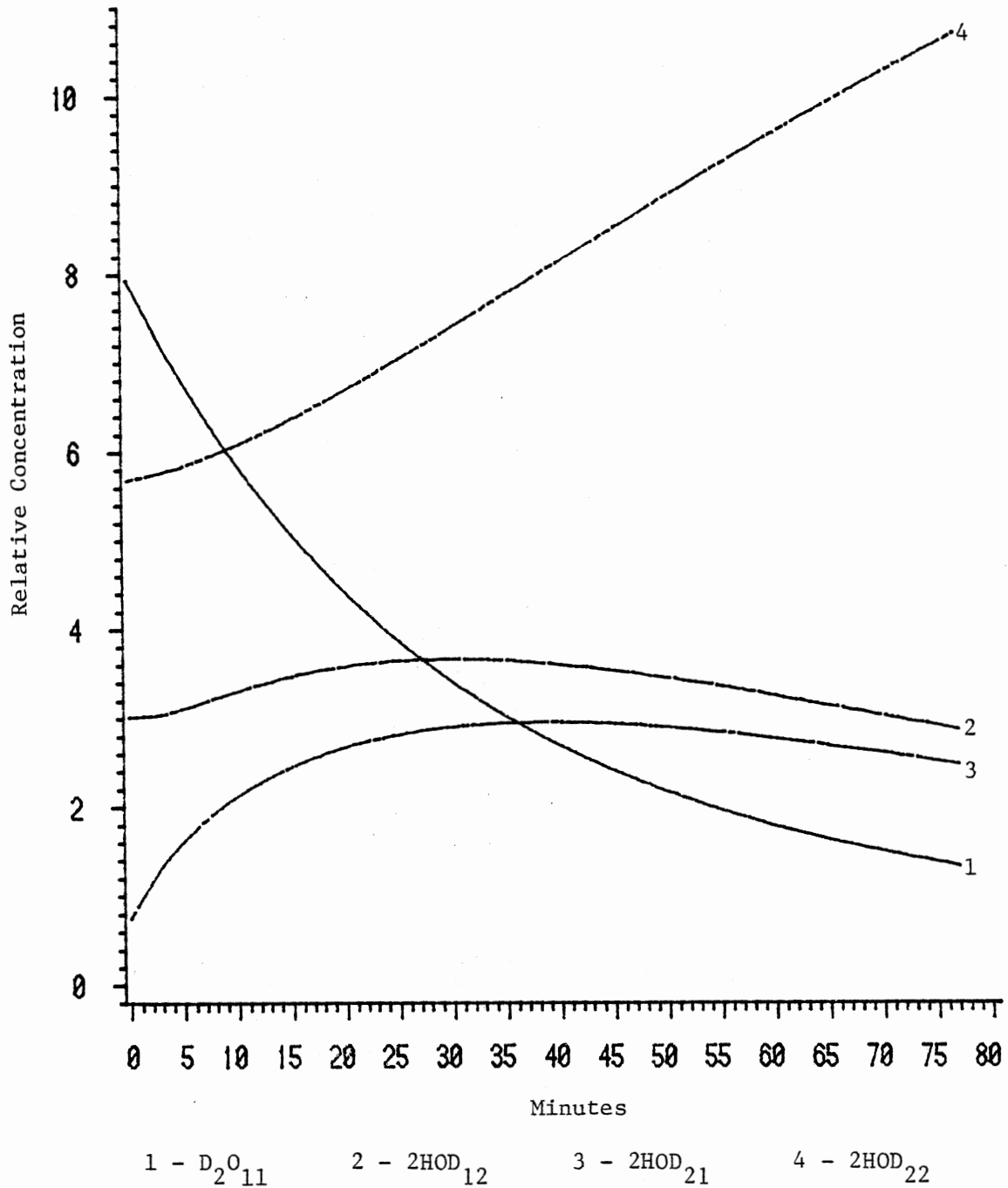


Figure 50. Calculated  $D_2O_{11}$  and  $2HOD_{nm}$  Configuration Concentration Curves for 5/10/82 145 K Data

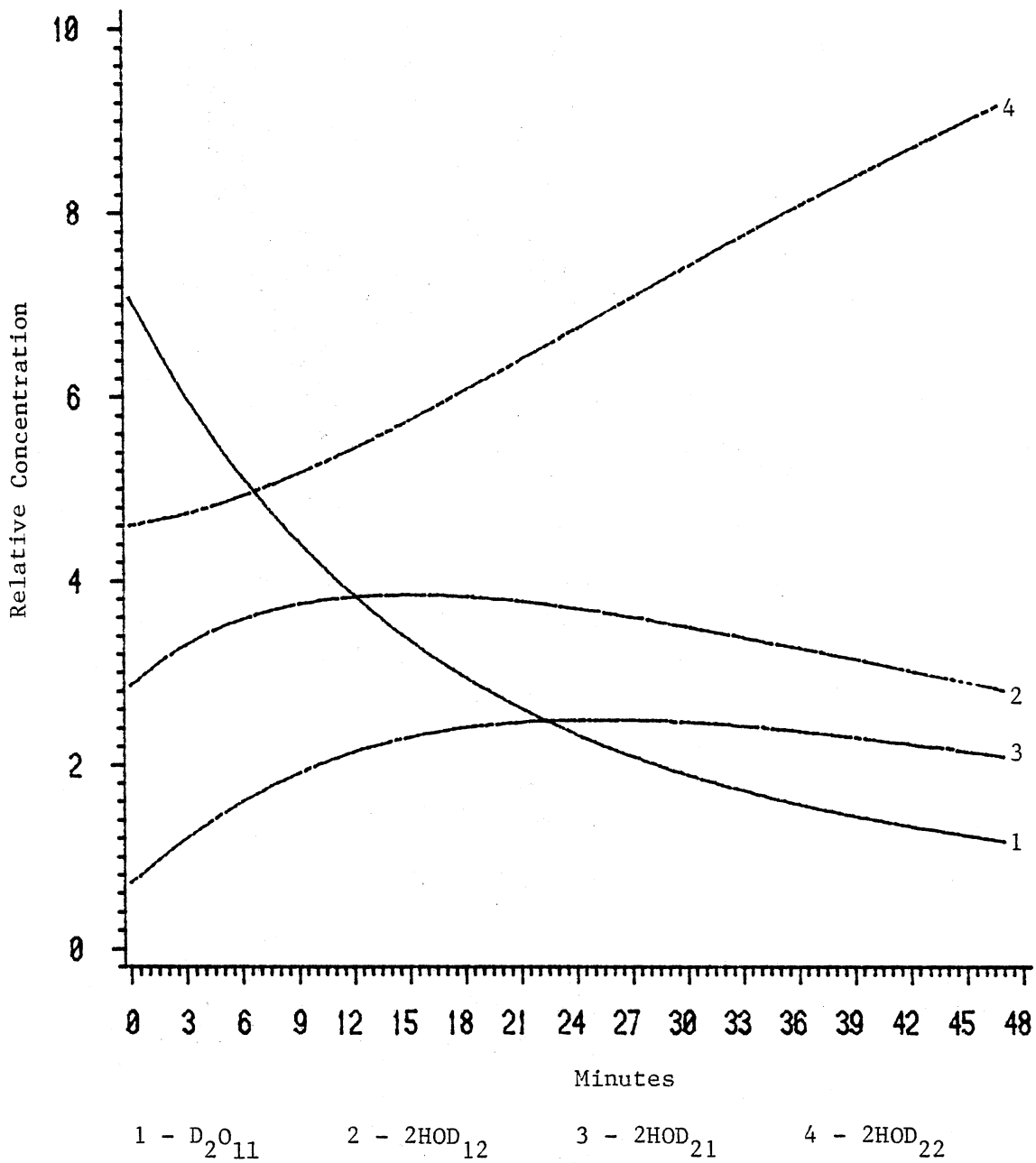


Figure 51. Calculated  $D_2O_{11}$  and  $2HOD_{nm}$  Configuration Concentration Curves for 5/24/82 145 K Data

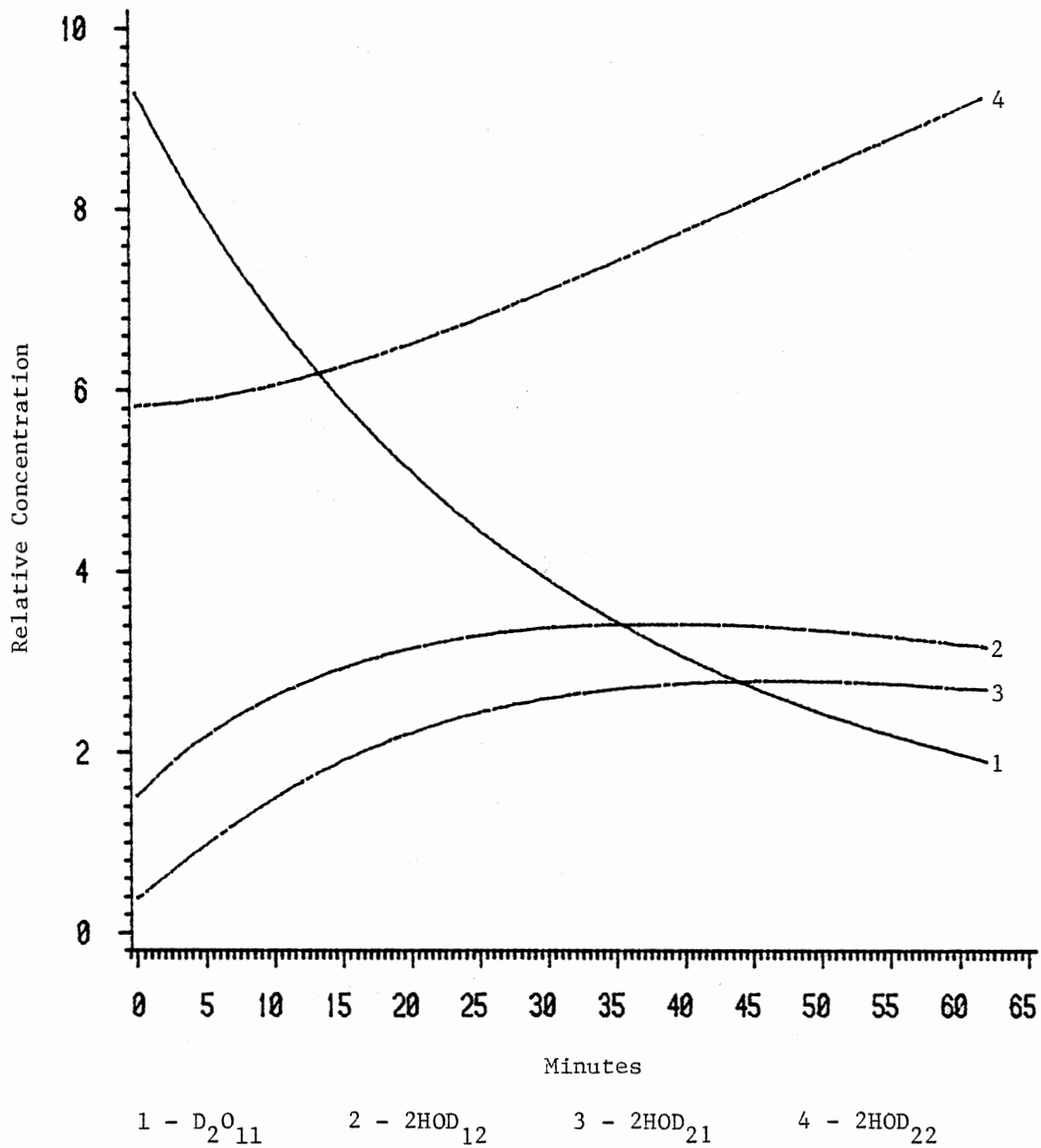


Figure 52. Calculated  $D_2O_{11}$  and  $2HOD_{nm}$  Configuration Concentration Curves for 6/3/82 145 K Data

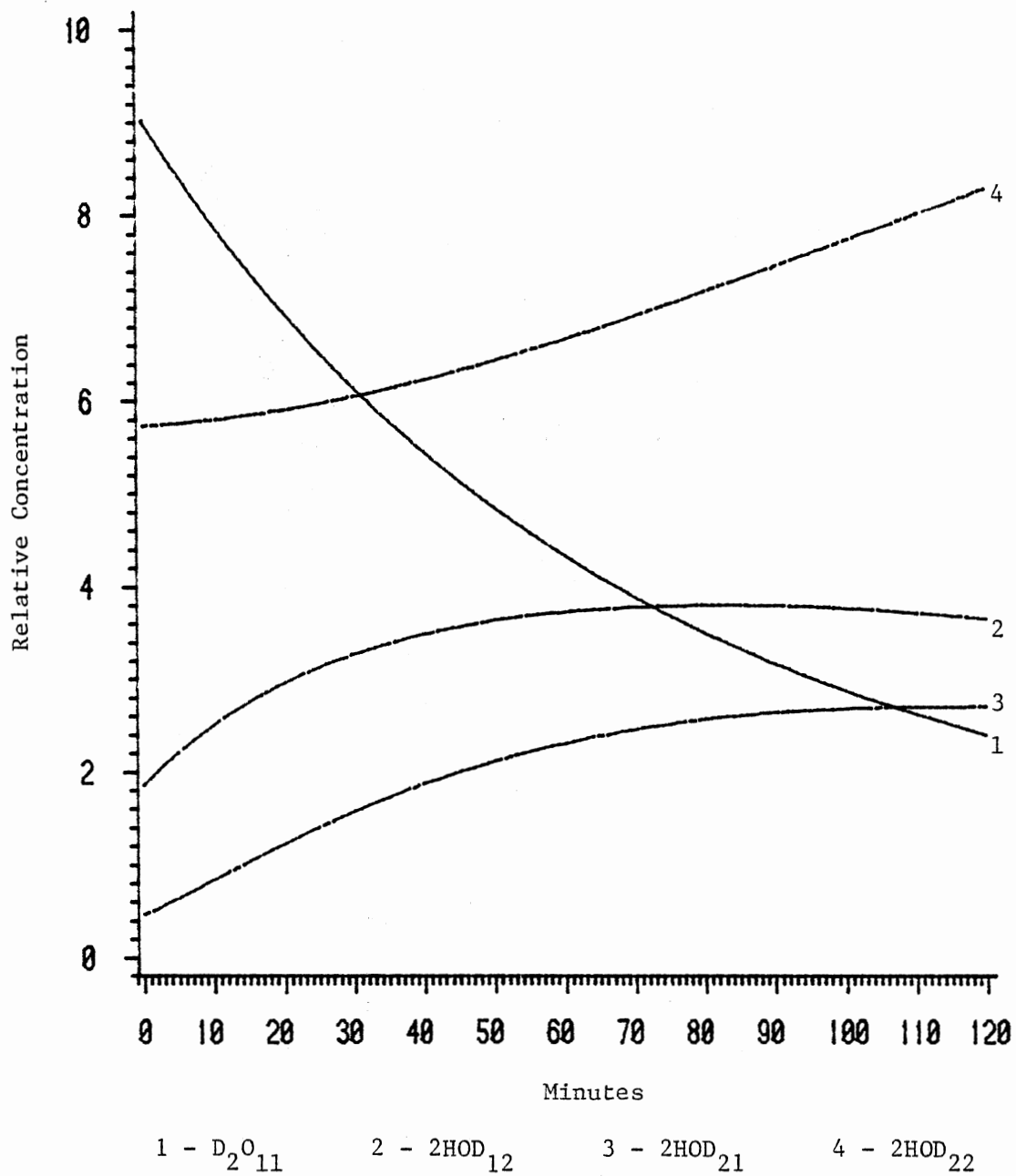


Figure 53. Calculated  $D_2O_{11}$  and  $2HOD_{nm}$  Configuration Concentration Curves for 5/11/82 140 K Data

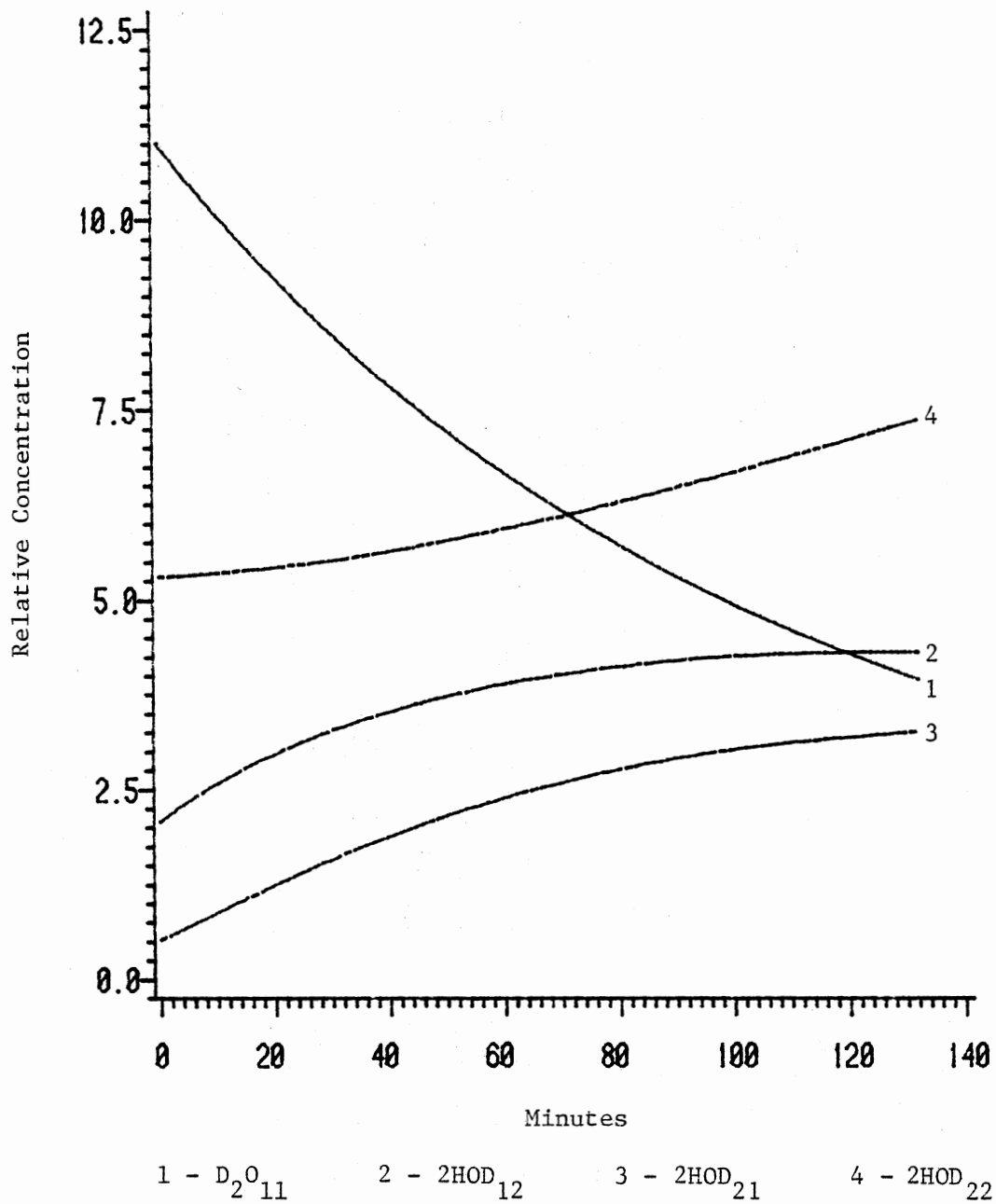


Figure 54. Calculated  $D_2O_{11}$  and  $2HOD_{nm}$  Configuration Concentration Curves for 5/25/82 140 K Data

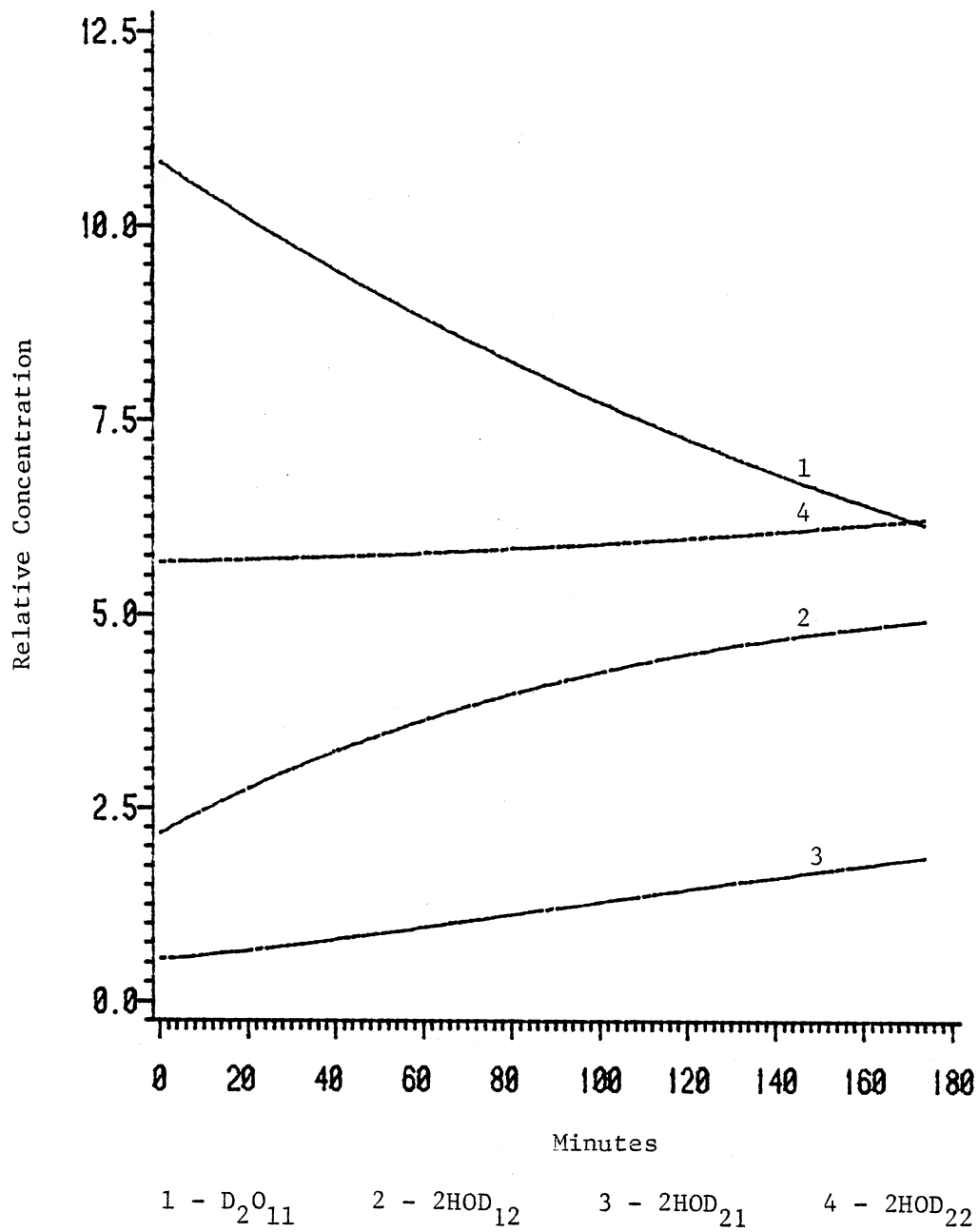


Figure 55. Calculated  $D_2O_{11}$  and  $2HOD_{nm}$  Configuration Concentration Curves for 5/15/82 135 K Data

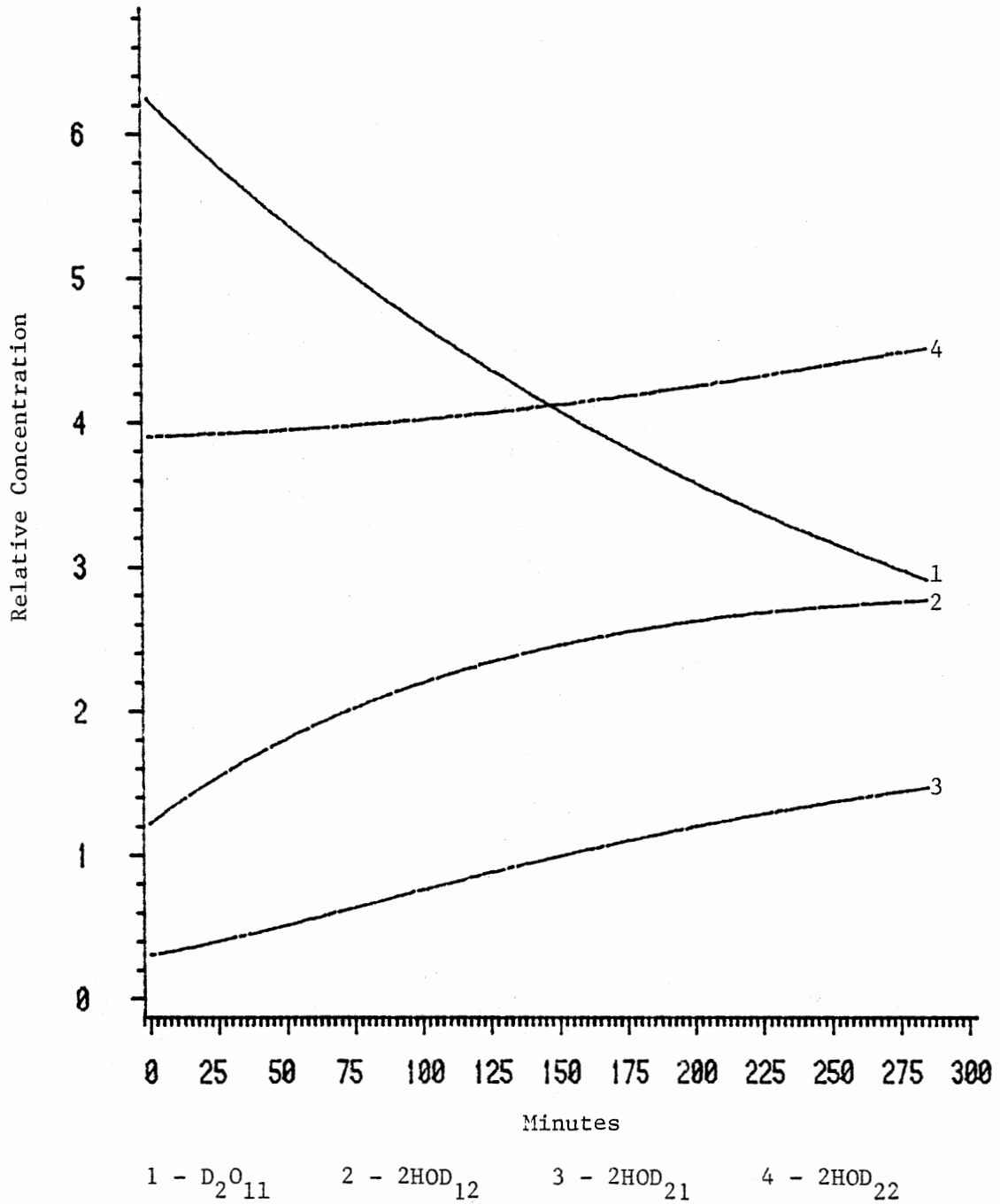


Figure 56. Calculated  $D_2O_{11}$  and  $2HOD_{nm}$  Configuration Concentration Curves for 5/28/82 135 K Data



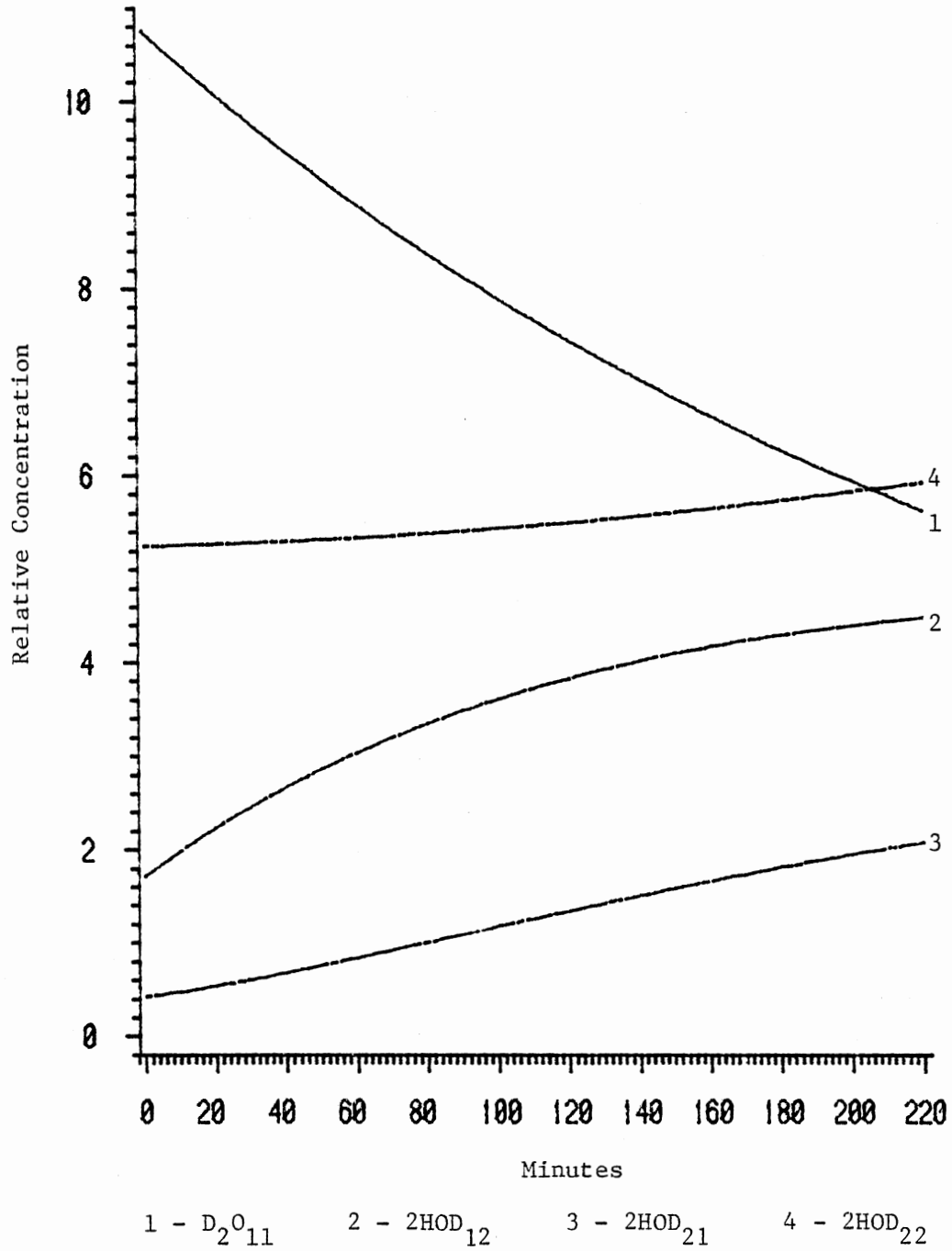


Figure 57. Calculated  $D_2O_{11}$  and  $2HOD_{nm}$  Configuration Concentration Curves for 6/11/82 135 K Data

$D_2O$  concentration after several half-lives of the  $D_2O$  concentration. This may be due to the slightly fast depletion of  $D_2O$  which builds up the  $(HOD)_2$  too fast. The exclusion of the  $(HOD)_2$  to pseudo-HOD to separated HOD migration path (see Figure 26,  $2HOD_{12} \leftrightarrow 2HOD_{13} \leftrightarrow 2HOD_{22}$ ) may account for some of the deviations. But attempts to include this pathway in the model failed due to too high a correlation in the rate constants and too many unknowns in the model equation.

#### Transfer and Rotational Activation Energies

The two-equilibrium model calculated rate constants shown in Table VII were least square fitted to standard Arrhenius plots to find the activation energies predicted for each calculation. Table VIII lists these results.

The correlation factors for the Arrhenius plots tend to support the contention that the 0.0 and .25 calculations are the best fit. For the transfer rate constants the .25 calculation has the best correlation and for the rotational rate constants the 0.0 calculation. In contrast the .5 calculation has the lowest correlation for both the transfer and rotational rate constant plots. A further indication of the poorer fit of the .5 calculation is the .561 standard error which is twice the standard error of the 0.0 and .25 calculations. A .7 calculation not shown gave an even poorer fit. So in summarizing, the 0.0 and .25 calculations gave the best activation energies for the rotational and transfer rate constants. Both calculations give essentially the same rate constants and hence the same activation energies. The upper plot of Figure 58 shows the Arrhenius plot for the transfer rate constants. From this a transfer

TABLE VIII  
TRANSFER AND ROTATIONAL ACTIVATION ENERGIES

Model initial HOD <sub>21</sub>	Two-equilibrium 0.0	Two-equilibrium 0.25	Two-equilibrium 0.5
Transfer Activation Energy (kcal/mole)	9.531	9.537	9.513
Standard Error	0.658	0.653	0.661
Least Square Correlation	0.97697	0.97736	0.97667
Rotational Activation Energy (kcal/mole)	11.767	12.199	13.296
Standard Error	0.228	0.244	0.561
Least Square Correlation	0.99813	0.99801	0.99123

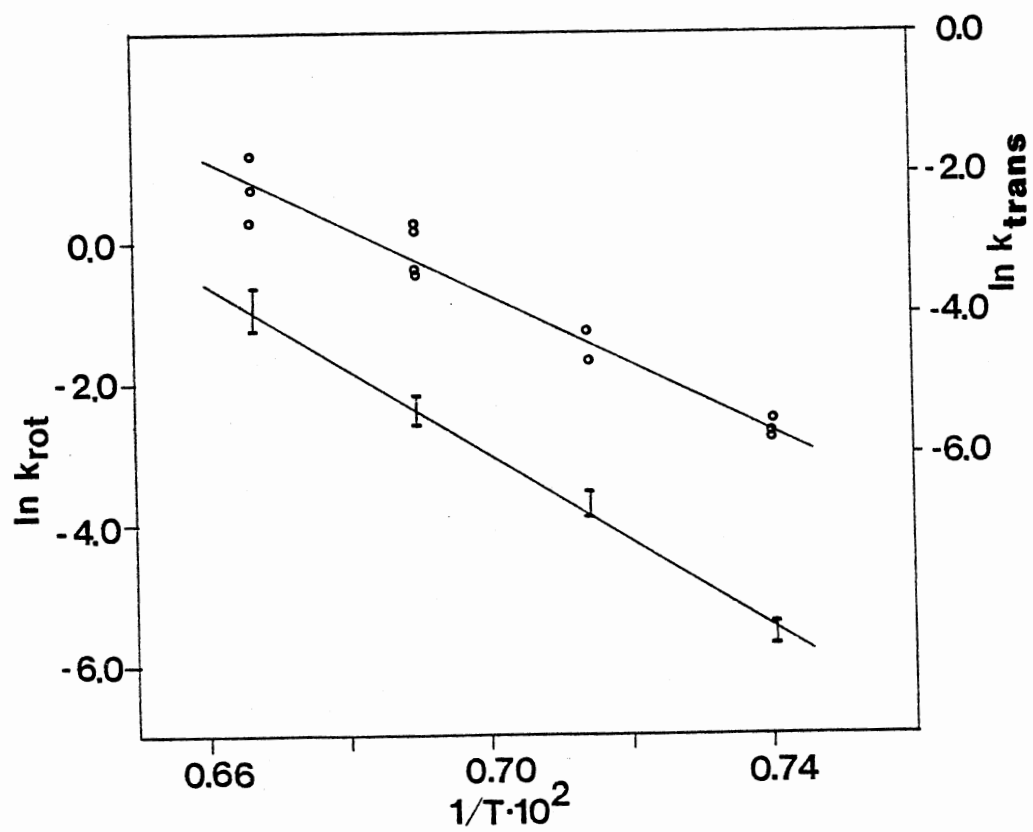


Figure 58. Transfer and Rotational Rate Constants Arrhenius Plots

rate constant activation energy of  $9.5 \pm .7$  kcal/mole with a pre-exponential factor of  $(9.4 \pm .7) \times 10^{12}$  is determined.

The rotational activation energy is not as well determined. Though the standard errors are smaller than those of the transfer rate constant activation energy the 0.0 and .25 calculations still give different activation energies. Unfortunately there is no suitable criteria to determine which value is correct. Thus since the initial separated HOD concentration cannot be determined closer than 0.0 to .25 times the initial  $(\text{HOD})_2$  concentration the 0.0 and .25 calculations should represent acceptable bounds to the rotational rate constant activation energy.

The lower Arrhenius plot of Figure 58 shows the error bars determined from the lowest and highest rate constant from either the 0.0 or .25 calculations for a given temperature. Thus the upper error bar of the 150 K data point was determined from the natural log of .5020 the 150 K 6/9/82 dataset .25 calculated rotational rate constant, and the lower error bar from the log of .2857 the 150 K 2/2/82 dataset 0.0 calculated rotational rate constant. Then taking the average of the 0.0 calculation 11.8 and .25 calculation 12.2 activation energies and drawing in the resulting straight line plot gives the lower plot shown in Figure 58. So in spite of the initial separated HOD indeterminacy, the rotational activation energy can still be error bounded to an accuracy equal or better than the transfer activation energy. Interestingly, unlike the transfer plot, the rotational plot errors are fairly consistent over the 150 to 135 K temperature range. This may reflect the possibility that the  $(\text{HOD})_2$  concentrations are less affected by sampling and spectral resolution

errors in obtaining the coupled HOD absorbances. From the lower plot a rotational rate constant activation energy of  $12.0 \pm .5$  kcal/mole with a pre-exponential factor of  $(1.1 \pm 1.1) \times 10^{17}$  is determined. The error limits were calculated by adding the calculated standard errors to one half the difference between the 0.0 and .25 least square determined activation energies and pre-exponential factors.

#### Discussion of Activation Energies

The calculated  $9.5 \pm .7$  kcal/mole transfer activation energy is quite close to the 9.3 kcal/mole value found by Ritzhaupt and Devlin (46) on their studies of doped ice systems. The transfer activation energy reflects an independence of the dopants that drastically affect the transfer rates. This gives further credence to the suggestion that the dopants affect the concentration of ionic defects and hence the rate. Accordingly the activation energy measured represents the ion pair defect formation energy. This ionization is dependent on the temperature, but independent of defect concentration or dopants present in the ice lattice.

A. von Hippel (42) suggested that D and L Bjerrum defects were the only charge carriers in bulk ice conductivity measurements. In view of the magnitudes of the transfer and rotational rate constants and the magnitudes of the  $D_2O$  and  $(HOD)_2$  concentrations, this view is hard to support at 135-150 K temperatures. Both the ionic and Bjerrum defects appear to be equally active in isotopic scrambling for pure ice systems. But the  $12.0 \pm .5$  kcal/mole rotational activation energy does imply that the rotational step may become more

predominant and the transfer step more rate determining at higher temperatures. Thus conductivity activation energies of pure ice in the 135 to 180 K region may be hard to relate to a particular defect activation energy. The 9 and 13 kcal/mole conductivity activation energies at 240-200 K and 180-170 K, reported by Durand et al. (40) may be a reflection of this problem.

Nagle and Tristram-Nagle (52) have reviewed various models for proton transport in biological membranes. In particular the hydrogen bonded chain mechanisms in which various amino acid hydroxyl and amine groups hydrogen bond to form a chain capable of transmitting protons via a hop/turn mechanism. For a particular hydrogen bonded chain (HBC) to transmit a proton more than once each hopping or ionic defect migration must be followed by a turn or D,L defect migration through the chain. The model presented in Figure 26 presents a similar situation. Here a particular deuteron cannot move through the ice lattice unless an alternating sequence of transfer and rotational steps is followed. Thus if  $D_2O$  completely scrambles to uncoupled HOD both transfer and rotational defects should be present. But if proton transfer could be induced without a rotational decoupling then a statistical 1:2:1 ratio of  $D_2O$ , coupled HOD, and uncoupled HOD should be produced. This should be identifiable spectroscopically provided the  $2HOD_{13}$  configuration contributes to the uncoupled HOD band intensity. Laser photoexcitation of the  $D_2O$  OD stretch modes and overtones in a  $D_2O/H_2O$  deposit cooled below 120 K may be one method of accomplishing this via a photoionization pathway.

### Summary

The primary goal of this study was to elucidate the kinetics and resulting activation energies for isotopic scrambling in pure cubic ice. Pertinent to this goal was the identification and quantification of mixed isotope vapor deposited ice spectra and development of an experimentally consistent kinetic model to explain the observed spectral changes.

Twelve  $D_2O/H_2O$  vapor deposited ice samples were made at 130 K and 125 K where no isotopic exchange would occur. After warming these samples to 135, 140, 145, or 150 K to initiate exchange, FTIR spectra were collected at various time intervals to monitor the respective spectral changes. To help resolve these spectra into their various components for absorbance and concentration measurements, additional  $D_2O/H_2O$  ice samples lightly doped with 7-azaindole were made. The IR spectra collected from these ice samples were used to produce pure component  $D_2O$  spectra at the required kinetic temperatures. Using the pure component uncoupled HOD spectra collected at the completion of isotopic exchange, and previously determined pure component  $D_2O$  spectra, the FTIR spectra collected for each ice sample could be resolved into the individual  $D_2O$ , and uncoupled HOD spectra. Accordingly absorbances were extracted from these spectra.

Relative concentrations were ascertained from calculated relative molar absorptivities. Using a material balance equation a system of simultaneous equations utilizing the kinetic absorbance data was solved repetitively for the relative molar absorptivities of  $D_2O$ , coupled HOD, and uncoupled HOD.



With the kinetic concentration data in hand, a model was developed relating the various component concentration changes to the migration of ionic and Bjerrum D,L defects. By breaking the interaction of an ionic or Bjerrum D,L defect with a  $D_2O$  or HOD unit in an ice crystal into three steps, the attack factor, the transfer factor, and the leaving factor, it was possible to relate the individual transfer and rotational steps necessary in isotopic scrambling to a common set of transfer and rotation rate constants reflecting the rate of ionic defect migration and D,L defect migration. Approximations yielded both a one- or two-equilibrium model suitable for nonlinear least square fit in the experimental concentration data.

Using an iterative nonlinear least square routine both the one- and two-equilibrium models were fit to the data. The two-equilibrium model with feedback from separated HOD to coupled HOD gave a superior fit. From the rate constants determined with the two equilibrium model Arrhenius plots were constructed and analyzed for the respective transfer and rotational rate activation energies.

As with most investigations, more questions are raised than are answered. One of the interesting prospects raised in this study is what would be the effect of  $NH_3$  or HF doping on the transfer or rotation rates? If rate suppression or enhancement is observed, then possibly the degree of suppression or enhancement could be used as a probe of the proton donicity or affinity of the dopant. A modification of the simultaneous deposit of sample and dopant would be a layered deposit where each layer is a different isotopic or chemical species. Two isotopes could be separated by a third sandwiched

layer which would have to transport protons for isotopic scrambling to occur. Proton transport across a specifically prepared molecular film could be monitored by this method. Thus there are numerous prospects for further study of proton transport in vapor deposited ice samples.

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APPENDIX A

FORTRAN PROGRAM REMOLE

C OSU BILL COLLIER CHEMISTRY PS203 4/27/82  
 C MAIN PROGRAM TO FIND D2O, (HOD)2, HOD RELATIVE MOLAR ABSORBTIVITIES  
 C FROM ABSORBANCE DATA AND MATERIAL BALANCE  
 C DOUBLE PRECISION USED

C  
 C INITIALIZE PROGRAM AND READ CONTROL  
 C CARDS. TIME DATA LISTED NEXT  
 C WITH TWO CARDS FOR GRAPHING  
 C ROUTINE LABELS IN FRONT OF EACH  
 C DATA SET

0001  
 0002 IMPLICIT REAL\*8 (A-H,O-Z)  
 DIMENSION ICARD(100),T(50),A(50),B(50),C(50),ITITLE(144),  
 0003 1IMAG4(5151),RANGE(4),ICHR(10),OUT(50,3),AA(2,2),BB(2),WKAREA(20)  
 0004 DATA IN/5/,IOUT/6/,RANGE/4\*0.0000/  
 DATA ICHAR/1HA,1HB,1HC,1H ,1H ,1H ,1H ,1H ,1H ,1H /

C  
 C START DATA SETS LOOP, READ GRAPH  
 C TITLE, READ HOD AT T=INFINITY  
 C AB UNITS (CINF), ISTART, ISTOP  
 C LIMITS AND THEN EACH DATA SET

0005 I = 0  
 0006 10 CONTINUE  
 0007 READ (IN,515) ISTART,ISTOP,NNCARD,CINF  
 0008 IF (ISTART.EQ.9999) STOP  
 0009 IF (ISTOP.LE.ISTART) GOTO 50  
 0010 IF (ISTART.GT.0) GOTO 15  
 0011 ISTART = IABS(ISTART)  
 0012 GOTO 35  
 0013 15 READ (IN,510) (ITITLE(M),M=1,144)  
 0014 I = I+1  
 0015 WRITE (IOUT,600)  
 0016 WRITE (IOUT,605) I  
 0017 DO 20 J=1,NNCARD  
 0018 READ (IN,520) T(J),A(J),B(J),C(J)  
 0019 C(J) = C(J)\*CINF  
 0020 WRITE (IOUT,610) T(J),A(J),B(J),C(J)  
 0021 20 CONTINUE

C  
 C PLOT INPUTTED DATA ON LINE  
 C PRINTER ROUTINE. ALIGN DATA FILES

0022 DO 30 J=1,NNCARD  
 0023 OUT(J,1) = A(J)  
 0024 OUT(J,2) = B(J)  
 0025 OUT(J,3) = C(J)  
 0026 30 CONTINUE  
 0027 INC = 1  
 0028 NP = NNCARD  
 0029 M = 3  
 0030 IOPT = 1  
 0031 IY = 50  
 0032 CALL USPLTD (T,OUT,IY,NP,M,INC,ITITLE,RANGE,ICHR,IOPT,IMAG4,  
 1 IER)

C  
 C SOLVE FOR RELATIVE MOLAR ABSORBTION  
 C COEFFICIENTS FOR ALL UNIQUE BINARY

COMBINATIONS OF OBSERVATIONS (TIMES)  
 ISTART TO ISTOP

C  
 C  
 C

```

0033      35  WRITE (IOUT,615)
0034          SUME1 = 0.0D00
0035          SUME2 = 0.0D00
0036          SUME3 = 0.0D00
0037          SUMEE1 = 0.0D00
0038          SUMEE2 = 0.0D00
0039          SUMEE3 = 0.0D00
0040          NMIN1 = ISTOP-1
0041          ICOUNT = 0
0042          DO 40 J=ISTART,NMIN1
0043              NPLUS1 = J+1
0044              DO 45 K=NPLUS1,ISTOP
    
```

C  
 C  
 C  
 C  
 C  
 C  
 C

SET UP AA MATRIX AND BB VECTOR  
 FOR E1, E2, E3 SOLUTIONS, SET E3  
 EQUAL TO 1.0 AND SOLVE TWO MATERIAL  
 BALANCE EQUATIONS FOR E1 AND E2 WHERE  
 $2A/E1+2B/E2 = CINF-C$

```

0045          AA(1,1) = A(J)
0046          AA(1,2) = B(J)
0047          AA(2,1) = A(K)
0048          AA(2,2) = B(K)
0049          BB(1) = CINF-C(J)
0050          BB(2) = CINF-C(K)
0051          CALL LEQT1F (AA,1,2,2, BB,2, WKAREA, IER)
0052          E1 = 2.0D00/BB(1)
0053          E2 = 2.0D00/BB(2)
0054          E3 = 1.0D00
0055          SUME1 = SUME1+E1
0056          SUME2 = SUME2+E2
0057          SUME3 = SUME3+E3
    
```

C  
 C  
 C  
 C  
 C  
 C  
 C

SET UP AA MATRIX AND BB VECTOR  
 FOR E1, E2, E3 SOLUTIONS, SET E2  
 EQUAL TO 1.0 AND SOLVE TWO MATERIAL  
 BALANCE EQUATIONS FOR E1 AND E3 WHERE  
 $CINF-C/2E3-A/E1 = B$

```

0058          AA(1,1) = CINF-C(J)
0059          AA(1,2) = -A(J)
0060          AA(2,1) = CINF-C(K)
0061          AA(2,2) = -A(K)
0062          BB(1) = B(J)
0063          BB(2) = B(K)
0064          CALL LEQT1F (AA,1,2,2, BB,2, WKAREA, IER)
0065          EE1 = 1.0D00/BB(2)
0066          EE2 = 1.0D00
0067          EE3 = 0.5D00/BB(1)
0068          SUMEE1 = SUMEE1+EE1
0069          SUMEE2 = SUMEE2+EE2
0070          SUMEE3 = SUMEE3+EE3
0071          ICOUNT = ICOUNT+1
0072          WRITE (IOUT,620) J,K,T(J),T(K),E1,E2,E3,EE1,EE2,EE3
0073
    
```

45 CONTINUE





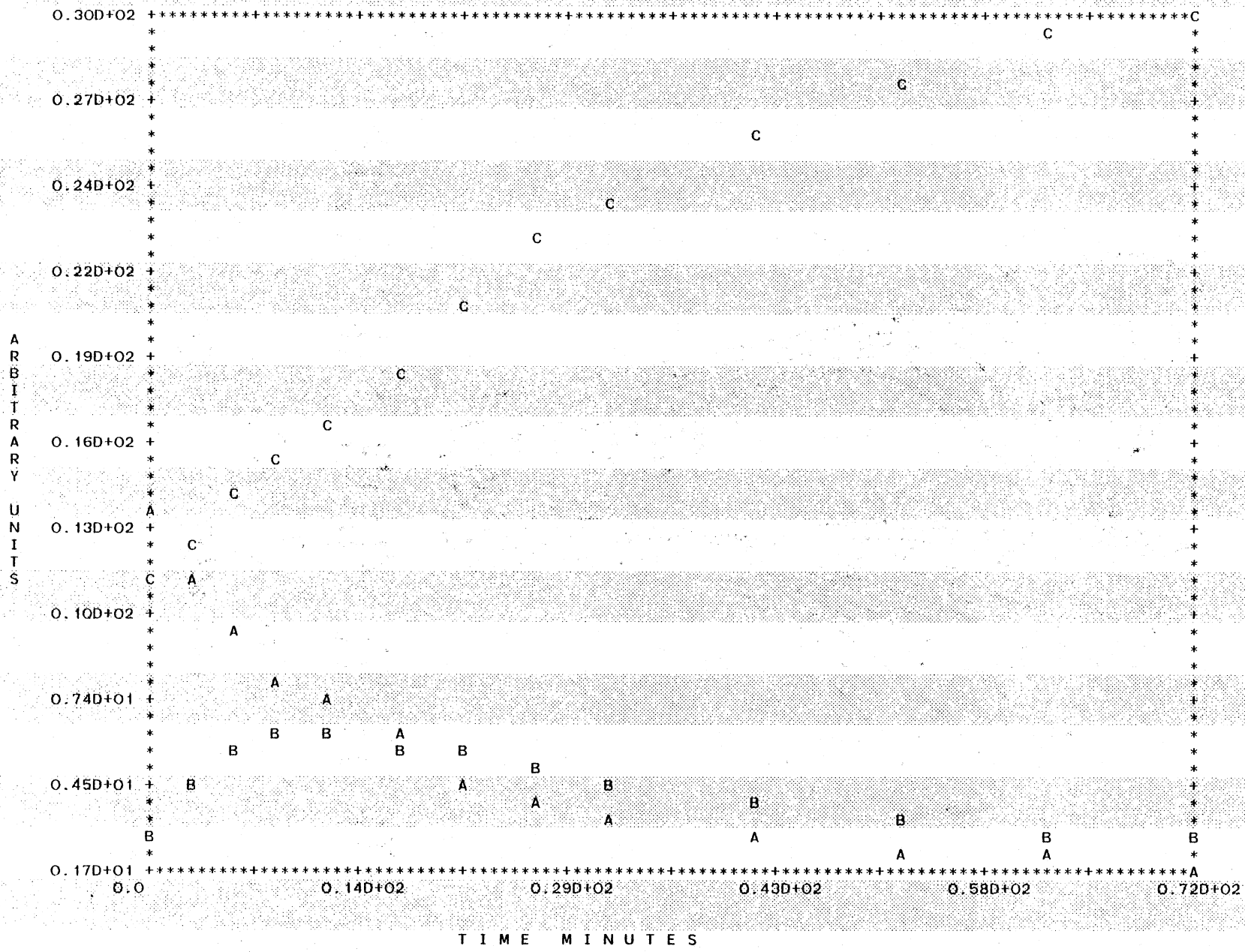
PS203 BILL COLLIER CHEMISTRY

RELATIVE MOLAR ABSORPTIVITIES BY SIMULTANEOUS EQUATIONS

DATASET # 1

MIN	D2O	(HOD)2	HOD
0.0	13.5000	3.1000	11.4000
3.0000	11.4000	4.5000	12.5400
6.0000	9.4000	5.4000	14.0600
9.0000	8.1500	5.9500	15.2000
12.0000	7.2500	6.1500	16.3400
17.0000	6.1500	5.9000	18.2400
21.7000	4.8000	5.5500	20.5200
27.0000	4.1000	5.1500	22.4200
32.0000	3.5000	4.7000	23.9400
42.0000	2.8500	3.9000	26.2200
52.0000	2.5000	3.4000	27.7400
62.0000	2.1000	2.8000	29.2600
72.0000	1.7000	2.5500	30.0200

D2O ISOLATED IN H2O 1/29/82 RUN 145K UNCORRECTED D2O VALUES



CALCULATED MOLAR ABSORBTIVITIES AT EACH STEP

OBS. #	OBS. #	TIME	TIME	E1	E2	E3	EE1	EE2	EE3
1	2	0.0	3.0	0.12464D+01	0.12557D+01	0.10000D+01	0.99254D+00	0.10000D+01	0.79634D+00
1	3	0.0	6.0	0.12606D+01	0.11964D+01	0.10000D+01	0.10536D+01	0.10000D+01	0.83581D+00
1	4	0.0	9.0	0.12572D+01	0.12100D+01	0.10000D+01	0.10390D+01	0.10000D+01	0.82646D+00
1	5	0.0	12.0	0.12557D+01	0.12164D+01	0.10000D+01	0.10323D+01	0.10000D+01	0.82213D+00
1	6	0.0	17.0	0.12664D+01	0.11745D+01	0.10000D+01	0.10782D+01	0.10000D+01	0.85145D+00
1	7	0.0	21.7	0.12852D+01	0.11089D+01	0.10000D+01	0.11590D+01	0.10000D+01	0.90182D+00
1	8	0.0	27.0	0.12812D+01	0.11220D+01	0.10000D+01	0.11418D+01	0.10000D+01	0.89123D+00
1	9	0.0	32.0	0.12918D+01	0.10878D+01	0.10000D+01	0.11876D+01	0.10000D+01	0.91930D+00
1	10	0.0	42.0	0.13036D+01	0.10530D+01	0.10000D+01	0.12380D+01	0.10000D+01	0.94967D+00
1	11	0.0	52.0	0.13013D+01	0.10596D+01	0.10000D+01	0.12281D+01	0.10000D+01	0.94377D+00
1	12	0.0	62.0	0.13206D+01	0.10072D+01	0.10000D+01	0.13111D+01	0.10000D+01	0.99281D+00
1	13	0.0	72.0	0.13532D+01	0.93281D+00	0.10000D+01	0.14506D+01	0.10000D+01	0.10720D+01
2	3	3.0	6.0	0.12946D+01	0.11467D+01	0.10000D+01	0.11290D+01	0.10000D+01	0.87207D+00
2	4	3.0	9.0	0.12746D+01	0.11886D+01	0.10000D+01	0.10723D+01	0.10000D+01	0.84129D+00
2	5	3.0	12.0	0.12683D+01	0.12026D+01	0.10000D+01	0.10546D+01	0.10000D+01	0.83152D+00
2	6	3.0	17.0	0.12916D+01	0.11527D+01	0.10000D+01	0.11206D+01	0.10000D+01	0.86756D+00
2	7	3.0	21.7	0.13304D+01	0.10814D+01	0.10000D+01	0.12302D+01	0.10000D+01	0.92469D+00
2	8	3.0	27.0	0.13198D+01	0.10996D+01	0.10000D+01	0.12002D+01	0.10000D+01	0.90941D+00
2	9	3.0	32.0	0.13417D+01	0.10630D+01	0.10000D+01	0.12621D+01	0.10000D+01	0.94071D+00
2	10	3.0	42.0	0.13670D+01	0.10249D+01	0.10000D+01	0.13337D+01	0.10000D+01	0.97568D+00
2	11	3.0	52.0	0.13621D+01	0.10320D+01	0.10000D+01	0.13198D+01	0.10000D+01	0.96899D+00
2	12	3.0	62.0	0.14062D+01	0.97336D+00	0.10000D+01	0.14447D+01	0.10000D+01	0.10274D+01
2	13	3.0	72.0	0.14766D+01	0.89830D+00	0.10000D+01	0.16437D+01	0.10000D+01	0.11132D+01
3	4	6.0	9.0	0.12338D+01	0.12411D+01	0.10000D+01	0.99410D+00	0.10000D+01	0.80575D+00
3	5	6.0	12.0	0.12330D+01	0.12424D+01	0.10000D+01	0.99247D+00	0.10000D+01	0.80490D+00
3	6	6.0	17.0	0.12883D+01	0.11555D+01	0.10000D+01	0.11150D+01	0.10000D+01	0.86546D+00
3	7	6.0	21.7	0.13645D+01	0.10628D+01	0.10000D+01	0.12840D+01	0.10000D+01	0.94095D+00
3	8	6.0	27.0	0.13417D+01	0.10878D+01	0.10000D+01	0.12334D+01	0.10000D+01	0.91926D+00
3	9	6.0	32.0	0.13816D+01	0.10452D+01	0.10000D+01	0.13219D+01	0.10000D+01	0.95676D+00
3	10	6.0	42.0	0.14297D+01	0.10009D+01	0.10000D+01	0.14285D+01	0.10000D+01	0.99913D+00
3	11	6.0	52.0	0.14204D+01	0.10089D+01	0.10000D+01	0.14079D+01	0.10000D+01	0.99117D+00
3	12	6.0	62.0	0.15104D+01	0.93972D+00	0.10000D+01	0.16073D+01	0.10000D+01	0.10642D+01
3	13	6.0	72.0	0.16475D+01	0.86204D+00	0.10000D+01	0.19111D+01	0.10000D+01	0.11600D+01
4	5	9.0	12.0	0.12316D+01	0.12441D+01	0.10000D+01	0.98995D+00	0.10000D+01	0.80379D+00
4	6	9.0	17.0	0.13562D+01	0.11038D+01	0.10000D+01	0.12287D+01	0.10000D+01	0.90598D+00
4	7	9.0	21.7	0.14798D+01	0.10098D+01	0.10000D+01	0.14654D+01	0.10000D+01	0.99031D+00
4	8	9.0	27.0	0.14222D+01	0.10495D+01	0.10000D+01	0.13551D+01	0.10000D+01	0.95284D+00
4	9	9.0	32.0	0.14875D+01	0.10049D+01	0.10000D+01	0.14802D+01	0.10000D+01	0.99511D+00
4	10	9.0	42.0	0.15750D+01	0.95578D+00	0.10000D+01	0.16478D+01	0.10000D+01	0.10463D+01
4	11	9.0	52.0	0.15583D+01	0.96435D+00	0.10000D+01	0.16159D+01	0.10000D+01	0.10370D+01
4	12	9.0	62.0	0.17445D+01	0.88433D+00	0.10000D+01	0.19727D+01	0.10000D+01	0.11308D+01
4	13	9.0	72.0	0.20016D+01	0.81193D+00	0.10000D+01	0.24652D+01	0.10000D+01	0.12316D+01
5	6	12.0	17.0	0.15797D+01	0.98547D+00	0.10000D+01	0.16030D+01	0.10000D+01	0.10147D+01
5	7	12.0	21.7	0.16863D+01	0.94170D+00	0.10000D+01	0.17907D+01	0.10000D+01	0.10619D+01
5	8	12.0	27.0	0.15411D+01	0.10040D+01	0.10000D+01	0.15350D+01	0.10000D+01	0.99604D+00
5	9	12.0	32.0	0.16370D+01	0.96077D+00	0.10000D+01	0.17038D+01	0.10000D+01	0.10408D+01
5	10	12.0	42.0	0.17872D+01	0.90796D+00	0.10000D+01	0.19684D+01	0.10000D+01	0.11014D+01
5	11	12.0	52.0	0.17591D+01	0.91673D+00	0.10000D+01	0.19189D+01	0.10000D+01	0.10908D+01
5	12	12.0	62.0	0.21415D+01	0.82611D+00	0.10000D+01	0.25923D+01	0.10000D+01	0.12105D+01
5	13	12.0	72.0	0.26096D+01	0.76380D+00	0.10000D+01	0.34167D+01	0.10000D+01	0.13092D+01
6	7	17.0	21.7	0.17786D+01	0.91868D+00	0.10000D+01	0.19360D+01	0.10000D+01	0.10885D+01
6	8	17.0	27.0	0.15205D+01	0.10111D+01	0.10000D+01	0.15039D+01	0.10000D+01	0.98904D+00
6	9	17.0	32.0	0.16647D+01	0.95384D+00	0.10000D+01	0.17452D+01	0.10000D+01	0.10484D+01
6	10	17.0	42.0	0.18963D+01	0.88897D+00	0.10000D+01	0.21332D+01	0.10000D+01	0.11249D+01
6	11	17.0	52.0	0.18526D+01	0.89934D+00	0.10000D+01	0.20600D+01	0.10000D+01	0.11119D+01
6	12	17.0	62.0	0.25678D+01	0.78825D+00	0.10000D+01	0.32576D+01	0.10000D+01	0.12686D+01
6	13	17.0	72.0	0.34195D+01	0.73006D+00	0.10000D+01	0.46839D+01	0.10000D+01	0.13697D+01
7	8	21.7	27.0	0.11061D+01	0.12612D+01	0.10000D+01	0.87701D+00	0.10000D+01	0.79288D+00
7	9	21.7	32.0	0.15000D+01	0.90000D+00	0.10000D+01	0.15000D+01	0.10000D+01	0.10000D+01

7	10	21.7	42.0	0.20784D+01	0.86307D+00	0.10000D+01	0.24082D+01	0.10000D+01	0.11387D+01
7	11	21.7	52.0	0.19646D+01	0.88140D+00	0.10000D+01	0.22290D+01	0.10000D+01	0.11346D+01
7	12	21.7	62.0	0.81693D+01	0.68078D+00	0.10000D+01	0.12000D+02	0.10000D+01	0.14689D+01
7	13	21.7	72.0	0.19684D+02	0.65324D+00	0.10000D+01	0.30133D+02	0.10000D+01	0.15308D+01
8	9	27.0	32.0	0.30477D+01	0.79910D+00	0.10000D+01	0.38140D+01	0.10000D+01	0.12514D+01
8	10	27.0	42.0	0.27632D+02	0.67394D+00	0.10000D+01	0.41000D+02	0.10000D+01	0.14838D+01
8	11	27.0	52.0	0.16015D+02	0.68357D+00	0.10000D+01	0.23429D+02	0.10000D+01	0.14629D+01
8	12	27.0	62.0	-0.95890D+00	0.42683D+00	0.10000D+01	-0.22466D+01	0.10000D+01	0.23429D+01
8	13	27.0	72.0	-0.24854D+01	0.54557D+00	0.10000D+01	-0.45556D+01	0.10000D+01	0.18329D+01
9	10	32.0	42.0	-0.95865D+00	0.44003D+00	0.10000D+01	-0.21786D+01	0.10000D+01	0.22725D+01
9	11	32.0	52.0	-0.71770D+00	0.39474D+00	0.10000D+01	-0.18182D+01	0.10000D+01	0.25333D+01
9	12	32.0	62.0	0.81871D-01	-0.13158D+00	0.10000D+01	-0.62222D+00	0.10000D+01	-0.76000D+01
9	13	32.0	72.0	-0.11313D+01	0.46425D+00	0.10000D+01	-0.24368D+01	0.10000D+01	0.21540D+01
10	11	42.0	52.0	-0.31579D+01	0.57416D+00	0.10000D+01	-0.55000D+01	0.10000D+01	0.17417D+01
10	12	42.0	62.0	0.38113D+00	-0.24561D+01	0.10000D+01	-0.15517D+00	0.10000D+01	-0.40714D+00
10	13	42.0	72.0	-0.11773D+01	0.46927D+00	0.10000D+01	-0.25088D+01	0.10000D+01	0.21310D+01
11	12	52.0	62.0	0.28340D+00	-0.92105D+00	0.10000D+01	-0.30769D+00	0.10000D+01	-0.10857D+01
11	13	52.0	72.0	-0.12281D+01	0.47448D+00	0.10000D+01	-0.25882D+01	0.10000D+01	0.21076D+01
12	13	62.0	72.0	-0.20877D+02	0.62632D+00	0.10000D+01	-0.33333D+02	0.10000D+01	0.15966D+01

AVERAGE E1    AVERAGE E2    AVERAGE E3    AVERAGE EE1    AVERAGE EE2    AVERAGE EE3

0.17510D+01    0.85713D+00    0.10000D+01    0.19007D+01    0.10000D+01    0.99247D+00



CALCULATED MOLAR ABSORPTIVITIES AT EACH STEP

OBS. #	OBS. #	TIME	TIME	E1	E2	E3	EE1	EE2	EE3
1	2	0.0	3.0	0.12464D+01	0.12557D+01	0.10000D+01	0.99254D+00	0.10000D+01	0.79634D+00
1	3	0.0	6.0	0.12606D+01	0.11964D+01	0.10000D+01	0.10536D+01	0.10000D+01	0.83581D+00
1	4	0.0	9.0	0.12572D+01	0.12100D+01	0.10000D+01	0.10390D+01	0.10000D+01	0.82646D+00
1	5	0.0	12.0	0.12557D+01	0.12164D+01	0.10000D+01	0.10323D+01	0.10000D+01	0.82213D+00
1	6	0.0	17.0	0.12664D+01	0.11745D+01	0.10000D+01	0.10782D+01	0.10000D+01	0.85145D+00
1	7	0.0	21.7	0.12852D+01	0.11089D+01	0.10000D+01	0.11590D+01	0.10000D+01	0.90182D+00
1	8	0.0	27.0	0.12812D+01	0.11220D+01	0.10000D+01	0.11418D+01	0.10000D+01	0.89123D+00
1	9	0.0	32.0	0.12918D+01	0.10878D+01	0.10000D+01	0.11876D+01	0.10000D+01	0.91930D+00
1	10	0.0	42.0	0.13036D+01	0.10530D+01	0.10000D+01	0.12380D+01	0.10000D+01	0.94967D+00
1	11	0.0	52.0	0.13013D+01	0.10596D+01	0.10000D+01	0.12281D+01	0.10000D+01	0.94377D+00
2	3	3.0	6.0	0.12946D+01	0.11467D+01	0.10000D+01	0.11290D+01	0.10000D+01	0.87207D+00
2	4	3.0	9.0	0.12746D+01	0.11886D+01	0.10000D+01	0.10723D+01	0.10000D+01	0.84129D+00
2	5	3.0	12.0	0.12683D+01	0.12026D+01	0.10000D+01	0.10546D+01	0.10000D+01	0.83152D+00
2	6	3.0	17.0	0.12916D+01	0.11527D+01	0.10000D+01	0.11206D+01	0.10000D+01	0.86756D+00
2	7	3.0	21.7	0.13304D+01	0.10814D+01	0.10000D+01	0.12302D+01	0.10000D+01	0.92469D+00
2	8	3.0	27.0	0.13198D+01	0.10996D+01	0.10000D+01	0.12002D+01	0.10000D+01	0.90941D+00
2	9	3.0	32.0	0.13417D+01	0.10630D+01	0.10000D+01	0.12621D+01	0.10000D+01	0.94071D+00
2	10	3.0	42.0	0.13670D+01	0.10249D+01	0.10000D+01	0.13337D+01	0.10000D+01	0.97568D+00
2	11	3.0	52.0	0.13621D+01	0.10320D+01	0.10000D+01	0.13198D+01	0.10000D+01	0.96899D+00
3	4	6.0	9.0	0.12338D+01	0.12411D+01	0.10000D+01	0.99410D+00	0.10000D+01	0.80575D+00
3	5	6.0	12.0	0.12330D+01	0.12424D+01	0.10000D+01	0.99247D+00	0.10000D+01	0.80490D+00
3	6	6.0	17.0	0.12883D+01	0.11555D+01	0.10000D+01	0.11150D+01	0.10000D+01	0.86546D+00
3	7	6.0	21.7	0.13645D+01	0.10628D+01	0.10000D+01	0.12840D+01	0.10000D+01	0.94095D+00
3	8	6.0	27.0	0.13417D+01	0.10878D+01	0.10000D+01	0.12334D+01	0.10000D+01	0.91926D+00
3	9	6.0	32.0	0.13816D+01	0.10452D+01	0.10000D+01	0.13219D+01	0.10000D+01	0.95676D+00
3	10	6.0	42.0	0.14297D+01	0.10009D+01	0.10000D+01	0.14285D+01	0.10000D+01	0.99913D+00
3	11	6.0	52.0	0.14204D+01	0.10089D+01	0.10000D+01	0.14079D+01	0.10000D+01	0.99117D+00
4	5	9.0	12.0	0.12316D+01	0.12441D+01	0.10000D+01	0.98995D+00	0.10000D+01	0.80379D+00
4	6	9.0	17.0	0.13562D+01	0.11038D+01	0.10000D+01	0.12287D+01	0.10000D+01	0.90598D+00
4	7	9.0	21.7	0.14798D+01	0.10098D+01	0.10000D+01	0.14654D+01	0.10000D+01	0.99031D+00
4	8	9.0	27.0	0.14222D+01	0.10495D+01	0.10000D+01	0.13551D+01	0.10000D+01	0.95284D+00
4	9	9.0	32.0	0.14875D+01	0.10049D+01	0.10000D+01	0.14802D+01	0.10000D+01	0.99511D+00
4	10	9.0	42.0	0.15750D+01	0.95578D+00	0.10000D+01	0.16478D+01	0.10000D+01	0.10463D+01
4	11	9.0	52.0	0.15583D+01	0.96435D+00	0.10000D+01	0.16159D+01	0.10000D+01	0.10370D+01
5	6	12.0	17.0	0.15797D+01	0.98547D+00	0.10000D+01	0.16030D+01	0.10000D+01	0.10147D+01
5	7	12.0	21.7	0.16863D+01	0.94170D+00	0.10000D+01	0.17907D+01	0.10000D+01	0.10619D+01
5	8	12.0	27.0	0.15411D+01	0.10040D+01	0.10000D+01	0.15350D+01	0.10000D+01	0.99604D+00
5	9	12.0	32.0	0.16370D+01	0.96077D+00	0.10000D+01	0.17038D+01	0.10000D+01	0.10408D+01
5	10	12.0	42.0	0.17872D+01	0.90796D+00	0.10000D+01	0.19684D+01	0.10000D+01	0.11014D+01
5	11	12.0	52.0	0.17591D+01	0.91673D+00	0.10000D+01	0.19189D+01	0.10000D+01	0.10908D+01
6	7	17.0	21.7	0.17786D+01	0.91868D+00	0.10000D+01	0.19360D+01	0.10000D+01	0.10885D+01
6	8	17.0	27.0	0.15205D+01	0.10111D+01	0.10000D+01	0.15039D+01	0.10000D+01	0.98904D+00
6	9	17.0	32.0	0.16647D+01	0.95384D+00	0.10000D+01	0.17452D+01	0.10000D+01	0.10484D+01
6	10	17.0	42.0	0.18963D+01	0.88897D+00	0.10000D+01	0.21332D+01	0.10000D+01	0.11249D+01
6	11	17.0	52.0	0.18526D+01	0.89934D+00	0.10000D+01	0.20600D+01	0.10000D+01	0.11119D+01
7	8	21.7	27.0	0.11061D+01	0.12612D+01	0.10000D+01	0.87701D+00	0.10000D+01	0.79288D+00
7	9	21.7	32.0	0.15207D+01	0.99398D+00	0.10000D+01	0.15300D+01	0.10000D+01	0.10061D+01
7	10	21.7	42.0	0.20784D+01	0.86307D+00	0.10000D+01	0.24082D+01	0.10000D+01	0.11587D+01
7	11	21.7	52.0	0.19646D+01	0.88140D+00	0.10000D+01	0.22290D+01	0.10000D+01	0.11346D+01
8	9	27.0	32.0	0.30477D+01	0.79910D+00	0.10000D+01	0.38140D+01	0.10000D+01	0.12514D+01
8	10	27.0	42.0	0.27632D+02	0.67394D+00	0.10000D+01	0.41000D+02	0.10000D+01	0.14838D+01
8	11	27.0	52.0	0.16015D+02	0.68357D+00	0.10000D+01	0.23429D+02	0.10000D+01	0.14629D+01
9	10	32.0	42.0	-0.95865D+00	0.44003D+00	0.10000D+01	-0.21786D+01	0.10000D+01	0.22725D+01
9	11	32.0	52.0	-0.71770D+00	0.39474D+00	0.10000D+01	-0.18182D+01	0.10000D+01	0.25333D+01
10	11	42.0	52.0	-0.31579D+01	0.57416D+00	0.10000D+01	-0.55000D+01	0.10000D+01	0.17417D+01

AVERAGE E1    AVERAGE E2    AVERAGE E3    AVERAGE EE1    AVERAGE EE2    AVERAGE EE3

CALCULATED MOLAR ABSORPTIVITIES AT EACH STEP

OBS. #	OBS. #	TIME	TIME	E1	E2	E3	EE1	EE2	EE3
1	2	0.0	3.0	0.12464D+01	0.12557D+01	0.10000D+01	0.99254D+00	0.10000D+01	0.79634D+00
1	3	0.0	6.0	0.12606D+01	0.11964D+01	0.10000D+01	0.10536D+01	0.10000D+01	0.83581D+00
1	4	0.0	9.0	0.12572D+01	0.12100D+01	0.10000D+01	0.10390D+01	0.10000D+01	0.82646D+00
1	5	0.0	12.0	0.12557D+01	0.12164D+01	0.10000D+01	0.10323D+01	0.10000D+01	0.82213D+00
1	6	0.0	17.0	0.12664D+01	0.11745D+01	0.10000D+01	0.10782D+01	0.10000D+01	0.85145D+00
1	7	0.0	21.7	0.12852D+01	0.11089D+01	0.10000D+01	0.11590D+01	0.10000D+01	0.90182D+00
1	8	0.0	27.0	0.12812D+01	0.11220D+01	0.10000D+01	0.11418D+01	0.10000D+01	0.89123D+00
2	3	3.0	6.0	0.12946D+01	0.11467D+01	0.10000D+01	0.11290D+01	0.10000D+01	0.87207D+00
2	4	3.0	9.0	0.12746D+01	0.11886D+01	0.10000D+01	0.10723D+01	0.10000D+01	0.84129D+00
2	5	3.0	12.0	0.12683D+01	0.12026D+01	0.10000D+01	0.10546D+01	0.10000D+01	0.83152D+00
2	6	3.0	17.0	0.12916D+01	0.11527D+01	0.10000D+01	0.11206D+01	0.10000D+01	0.86756D+00
2	7	3.0	21.7	0.13304D+01	0.10814D+01	0.10000D+01	0.12302D+01	0.10000D+01	0.92469D+00
2	8	3.0	27.0	0.13198D+01	0.10996D+01	0.10000D+01	0.12002D+01	0.10000D+01	0.90941D+00
3	4	6.0	9.0	0.12338D+01	0.12411D+01	0.10000D+01	0.99410D+00	0.10000D+01	0.80575D+00
3	5	6.0	12.0	0.12330D+01	0.12424D+01	0.10000D+01	0.99247D+00	0.10000D+01	0.80490D+00
3	6	6.0	17.0	0.12883D+01	0.11555D+01	0.10000D+01	0.11150D+01	0.10000D+01	0.86546D+00
3	7	6.0	21.7	0.13645D+01	0.10628D+01	0.10000D+01	0.12840D+01	0.10000D+01	0.94095D+00
3	8	6.0	27.0	0.13417D+01	0.10878D+01	0.10000D+01	0.12334D+01	0.10000D+01	0.91926D+00
4	5	9.0	12.0	0.12316D+01	0.12441D+01	0.10000D+01	0.98995D+00	0.10000D+01	0.80379D+00
4	6	9.0	17.0	0.13562D+01	0.11038D+01	0.10000D+01	0.12287D+01	0.10000D+01	0.90598D+00
4	7	9.0	21.7	0.14798D+01	0.10098D+01	0.10000D+01	0.14654D+01	0.10000D+01	0.99031D+00
4	8	9.0	27.0	0.14222D+01	0.10495D+01	0.10000D+01	0.13551D+01	0.10000D+01	0.95284D+00
5	6	12.0	17.0	0.15797D+01	0.98547D+00	0.10000D+01	0.16030D+01	0.10000D+01	0.10147D+01
5	7	12.0	21.7	0.16863D+01	0.94170D+00	0.10000D+01	0.17907D+01	0.10000D+01	0.10619D+01
5	8	12.0	27.0	0.15411D+01	0.10040D+01	0.10000D+01	0.15350D+01	0.10000D+01	0.99604D+00
6	7	17.0	21.7	0.17786D+01	0.91868D+00	0.10000D+01	0.19360D+01	0.10000D+01	0.10885D+01
6	8	17.0	27.0	0.15205D+01	0.10111D+01	0.10000D+01	0.15039D+01	0.10000D+01	0.98904D+00
7	8	21.7	27.0	0.11061D+01	0.12612D+01	0.10000D+01	0.87701D+00	0.10000D+01	0.79288D+00

AVERAGE E1    AVERAGE E2    AVERAGE E3    AVERAGE EE1    AVERAGE EE2    AVERAGE EE3

0.13498D+01    0.11241D+01    0.10000D+01    0.12217D+01    0.10000D+01    0.89658D+00

CALCULATED MOLAR ABSORBTIVITIES AT EACH STEP

OBS. #	OBS. #	TIME	TIME	E1	E2	E3	EE1	EE2	EE3
1	2	0.0	3.0	0.12464D+01	0.12557D+01	0.10000D+01	0.99254D+00	0.10000D+01	0.79634D+00
1	3	0.0	6.0	0.12606D+01	0.11964D+01	0.10000D+01	0.10536D+01	0.10000D+01	0.83581D+00
1	4	0.0	9.0	0.12572D+01	0.12100D+01	0.10000D+01	0.10390D+01	0.10000D+01	0.82646D+00
1	5	0.0	12.0	0.12557D+01	0.12164D+01	0.10000D+01	0.10323D+01	0.10000D+01	0.82213D+00
1	6	0.0	17.0	0.12664D+01	0.11745D+01	0.10000D+01	0.10782D+01	0.10000D+01	0.85145D+00
1	7	0.0	21.7	0.12852D+01	0.11089D+01	0.10000D+01	0.11590D+01	0.10000D+01	0.90182D+00
2	3	3.0	6.0	0.12946D+01	0.11467D+01	0.10000D+01	0.11290D+01	0.10000D+01	0.87207D+00
2	4	3.0	9.0	0.12746D+01	0.11886D+01	0.10000D+01	0.10723D+01	0.10000D+01	0.84129D+00
2	5	3.0	12.0	0.12683D+01	0.12026D+01	0.10000D+01	0.10546D+01	0.10000D+01	0.83152D+00
2	6	3.0	17.0	0.12916D+01	0.11527D+01	0.10000D+01	0.11206D+01	0.10000D+01	0.86756D+00
2	7	3.0	21.7	0.13304D+01	0.10814D+01	0.10000D+01	0.12302D+01	0.10000D+01	0.92469D+00
3	4	6.0	9.0	0.12338D+01	0.12411D+01	0.10000D+01	0.99410D+00	0.10000D+01	0.80575D+00
3	5	6.0	12.0	0.12330D+01	0.12424D+01	0.10000D+01	0.99247D+00	0.10000D+01	0.80490D+00
3	6	6.0	17.0	0.12883D+01	0.11555D+01	0.10000D+01	0.11150D+01	0.10000D+01	0.86546D+00
3	7	6.0	21.7	0.13645D+01	0.10628D+01	0.10000D+01	0.12840D+01	0.10000D+01	0.94095D+00
4	5	9.0	12.0	0.12316D+01	0.12441D+01	0.10000D+01	0.98995D+00	0.10000D+01	0.80379D+00
4	6	9.0	17.0	0.13562D+01	0.11038D+01	0.10000D+01	0.12287D+01	0.10000D+01	0.90598D+00
4	7	9.0	21.7	0.14798D+01	0.10098D+01	0.10000D+01	0.14654D+01	0.10000D+01	0.99031D+00
5	6	12.0	17.0	0.15797D+01	0.98547D+00	0.10000D+01	0.16030D+01	0.10000D+01	0.10147D+01
5	7	12.0	21.7	0.16863D+01	0.94170D+00	0.10000D+01	0.17907D+01	0.10000D+01	0.10619D+01
6	7	17.0	21.7	0.17786D+01	0.91868D+00	0.10000D+01	0.19360D+01	0.10000D+01	0.10885D+01

AVERAGE E1    AVERAGE E2    AVERAGE E3    AVERAGE EE1    AVERAGE EE2    AVERAGE EE3

0.13458D+01    0.11352D+01    0.10000D+01    0.12077D+01    0.10000D+01    0.88826D+00



CALCULATED MOLAR ABSORBTIVITIES AT EACH STEP

OBS. #	OBS. #	TIME	TIME	E1	E2	E3	EE1	EE2	EE3
1	2	0.0	3.0	0.12464D+01	0.12557D+01	0.10000D+01	0.99254D+00	0.10000D+01	0.79634D+00
1	3	0.0	6.0	0.12606D+01	0.11964D+01	0.10000D+01	0.10536D+01	0.10000D+01	0.83581D+00
1	4	0.0	9.0	0.12572D+01	0.12100D+01	0.10000D+01	0.10390D+01	0.10000D+01	0.82646D+00
1	5	0.0	12.0	0.12557D+01	0.12164D+01	0.10000D+01	0.10323D+01	0.10000D+01	0.82213D+00
1	6	0.0	17.0	0.12664D+01	0.11745D+01	0.10000D+01	0.10782D+01	0.10000D+01	0.85145D+00
2	3	3.0	6.0	0.12946D+01	0.11467D+01	0.10000D+01	0.11290D+01	0.10000D+01	0.87207D+00
2	4	3.0	9.0	0.12746D+01	0.11886D+01	0.10000D+01	0.10723D+01	0.10000D+01	0.84129D+00
2	5	3.0	12.0	0.12683D+01	0.12026D+01	0.10000D+01	0.10546D+01	0.10000D+01	0.83152D+00
2	6	3.0	17.0	0.12916D+01	0.11527D+01	0.10000D+01	0.11206D+01	0.10000D+01	0.86756D+00
3	4	6.0	9.0	0.12338D+01	0.12411D+01	0.10000D+01	0.99410D+00	0.10000D+01	0.80575D+00
3	5	6.0	12.0	0.12330D+01	0.12424D+01	0.10000D+01	0.99247D+00	0.10000D+01	0.80490D+00
3	6	6.0	17.0	0.12883D+01	0.11555D+01	0.10000D+01	0.11150D+01	0.10000D+01	0.86546D+00
4	5	9.0	12.0	0.12316D+01	0.12441D+01	0.10000D+01	0.98995D+00	0.10000D+01	0.80379D+00
4	6	9.0	17.0	0.13562D+01	0.11038D+01	0.10000D+01	0.12287D+01	0.10000D+01	0.90598D+00
5	6	12.0	17.0	0.15797D+01	0.98547D+00	0.10000D+01	0.16030D+01	0.10000D+01	0.10147D+01

AVERAGE E1    AVERAGE E2    AVERAGE E3    AVERAGE EE1    AVERAGE EE2    AVERAGE EE3

0.12892D+01    0.11811D+01    0.10000D+01    0.10997D+01    0.10000D+01    0.84968D+00

APPENDIX B

INPUT INSTRUCTIONS TO FORTRAN SUBROUTINE DGEAR

PURPOSE - DIFFERENTIAL EQUATION SOLVER - VARIABLE ORDER  
 ADAMS PREDICTOR CORRECTOR METHOD OR  
 GEARS METHOD

USAGE - CALL DGEAR (N,FCN,FCNJ,X,H,Y,XEND,TOL,METH,  
 MITER,INDEX,IWK,WK,IER)

ARGUMENTS N - INPUT NUMBER OF FIRST-ORDER DIFFERENTIAL  
 EQUATIONS. N CAN BE REDUCED, BUT NEVER  
 INCREASED, DURING A PROBLEM. N MUST BE  
 GREATER THAN ZERO.

FCN - NAME OF SUBROUTINE FOR EVALUATING FUNCTIONS.  
 (INPUT)  
 THE SUBROUTINE ITSELF MUST ALSO BE PROVIDED  
 BY THE USER AND IT SHOULD BE OF THE  
 FOLLOWING FORM  
 SUBROUTINE FCN (N,X,Y,YPRIME)  
 REAL Y(N),YPRIME(N)

FCN SHOULD EVALUATE YPRIME(1),...,YPRIME(N)  
 GIVEN N,X, AND Y(1),...,Y(N). YPRIME(I)  
 IS THE FIRST DERIVATIVE OF Y(I) WITH  
 RESPECT TO X.

FCN MUST APPEAR IN AN EXTERNAL STATEMENT IN  
 THE CALLING PROGRAM AND N,X,Y(1),...,Y(N)  
 MUST NOT BE ALTERED BY FCN.

FCNJ - NAME OF THE SUBROUTINE FOR COMPUTING THE N BY  
 N JACOBIAN MATRIX OF PARTIAL DERIVATIVES.  
 (INPUT)  
 THE SUBROUTINE ITSELF MUST ALSO BE PROVIDED  
 BY THE USER AND IT SHOULD BE OF THE  
 FOLLOWING FORM  
 SUBROUTINE FCNJ (N,X,Y,PD)  
 REAL Y(N),PD(N,N)

FCNJ SHOULD COMPUTE PD(I,J) WHICH IS THE  
 PARTIAL DERIVATIVE OF YPRIME(I) WITH  
 RESPECT TO Y(J).

FCNJ MUST APPEAR IN AN EXTERNAL STATEMENT IN  
 THE CALLING PROGRAM AND N,X,Y(1),...,Y(N)  
 MUST NOT BE ALTERED BY FCNJ.

FCNJ IS CALLED ONLY IF MITER IS EQUAL TO  
 ONE. OTHERWISE A DUMMY ROUTINE CAN BE  
 SUBSTITUTED. SEE REMARKS.

X - INDEPENDENT VARIABLE. (INPUT AND OUTPUT)  
 ON INPUT, X SUPPLIES THE INITIAL VALUE AND  
 AND IS USED ONLY ON THE FIRST CALL.  
 ON OUTPUT, X IS REPLACED WITH THE CURRENT  
 VALUE OF THE INDEPENDENT VARIABLE AT WHICH  
 INTEGRATION HAS BEEN COMPLETED.

- H - INPUT/OUTPUT.  
 ON INPUT, H CONTAINS THE NEXT STEP SIZE IN X. H IS USED ONLY ON THE FIRST CALL.  
 ON OUTPUT, H CONTAINS THE STEP SIZE USED LAST, WHETHER SUCCESSFULLY OR NOT.
- Y - DEPENDENT VARIABLES, VECTOR OF LENGTH N.  
 (INPUT AND OUTPUT)  
 ON INPUT, Y(1),...,Y(N) SUPPLY INITIAL VALUES.  
 ON OUTPUT, Y(1),...,Y(N) ARE REPLACED WITH A COMPUTED VALUE AT XEND.
- XEND - INPUT VALUE OF X AT WHICH SOLUTION IS DESIRED NEXT. INTEGRATION WILL NORMALLY GO SLIGHTLY BEYOND XEND AND THE ROUTINE WILL INTERPOLATE TO X = XEND.  
 NOTE THAT (X-XEND)\*H MUST BE LESS THAN ZERO (X AND H AS SPECIFIED ON INPUT).
- TOL - INPUT RELATIVE ERROR BOUND. TOL MUST BE GREATER THAN ZERO. TOL IS USED ONLY ON THE FIRST CALL UNLESS INDEX IS EQUAL TO -1. TOL SHOULD BE AT LEAST AN ORDER OF MAGNITUDE LARGER THAN THE UNIT ROUND OFF BUT GENERALLY NOT LARGER THAN .001. SINGLE STEP ERROR ESTIMATES DIVIDED BY YMAX(I) WILL BE KEPT LESS THAN TOL IN ROOT-MEAN-SQUARE NORM (EUCLIDEAN NORM DIVIDED BY SQRT(N)). THE VECTOR YMAX OF WEIGHTS IS COMPUTED INTERNALLY AND STORED IN WORK VECTOR WK. INITIALLY YMAX(I) IS THE ABSOLUTE VALUE OF Y(I), WITH A DEFAULT VALUE OF ONE IF Y(I) IS EQUAL TO ZERO. THEREAFTER, YMAX(I) IS THE LARGEST VALUE OF THE ABSOLUTE VALUE OF Y(I) SEEN SO FAR, OR THE INITIAL VALUE OF YMAX(I) IF THAT IS LARGER.
- METH - INPUT BASIC METHOD INDICATOR.  
 USED ONLY ON THE FIRST CALL UNLESS INDEX IS EQUAL TO -1.  
 METH = 1, IMPLIES THAT THE ADAMS METHOD IS TO BE USED.  
 METH = 2, IMPLIES THAT THE STIFF METHODS OF GEAR, OR THE BACKWARD DIFFERENTIATION FORMULAE ARE TO BE USED.
- MITER - INPUT ITERATION METHOD INDICATOR.  
 MITER = 0, IMPLIES THAT FUNCTIONAL ITERATION IS USED. NO PARTIAL DERIVATIVES ARE NEEDED. A DUMMY FCNJ CAN BE USED.  
 MITER = 1, IMPLIES THAT THE CHORD METHOD IS USED WITH AN ANALYTIC JACOBIAN. FOR THIS METHOD, THE USER SUPPLIES SUBROUTINE FCNJ.  
 MITER = 2, IMPLIES THAT THE CHORD METHOD IS USED WITH THE JACOBIAN CALCULATED INTERNALLY BY FINITE DIFFERENCES. A DUMMY FCNJ CAN BE USED.

MITER = 3, IMPLIES THAT THE CHORD METHOD  
IS USED WITH THE JACOBIAN REPLACED BY 168  
A DIAGONAL APPROXIMATION BASED ON A  
DIRECTIONAL DERIVATIVE.  
A DUMMY FCNJ CAN BE USED.

- INDEX - INPUT AND OUTPUT PARAMETER USED TO INDICATE  
THE TYPE OF CALL TO THE SUBROUTINE.  
ON INPUT, INDEX = 1, IMPLIES THAT THIS IS THE  
FIRST CALL FOR THIS PROBLEM. ON OUTPUT INDEX  
IS RESET TO 0 IF INTEGRATION WAS SUCCESSFUL.  
OTHERWISE, THE VALUE OF INDEX IS UNCHANGED.  
ON INPUT, INDEX = 0, IMPLIES THAT THIS IS NOT  
THE FIRST CALL FOR THIS PROBLEM. INDEX IS  
UNCHANGED ON OUTPUT.  
ON INPUT, INDEX = -1, IMPLIES THAT THIS IS NOT  
THE FIRST CALL FOR THIS PROBLEM, AND THE  
USER HAS RESET TOL. ON OUTPUT INDEX IS RESET  
TO 0 IF INTEGRATION WAS SUCCESSFUL.  
OTHERWISE THE VALUE OF INDEX IS UNCHANGED.  
ON INPUT, INDEX = 2, IMPLIES THAT THIS IS NOT  
THE FIRST CALL FOR THIS PROBLEM. INTEGRATION  
IS TO CONTINUE AND XEND IS TO BE HIT EXACTLY  
(NO INTERPOLATION IS DONE). THIS VALUE OF  
INDEX ASSUMES THAT XEND IS BEYOND THE  
CURRENT VALUE OF X. INDEX IS UNCHANGED ON  
OUTPUT.  
ON INPUT, INDEX = 3, IMPLIES THAT THIS IS NOT  
THE FIRST CALL FOR THIS PROBLEM. INTEGRATION  
IS TO CONTINUE AND CONTROL IS TO BE RETURNED  
TO THE CALLING PROGRAM AFTER ONE STEP. XEND  
IS IGNORED. INDEX IS UNCHANGED ON OUTPUT.
- IWK - INTEGER WORK VECTOR OF LENGTH N. USED ONLY IF  
MITER IS EQUAL TO 1 OR 2.
- WK - REAL WORK VECTOR OF LENGTH  $4*N+NMETH+NMITER$ .  
THE VALUE OF NMETH DEPENDS ON THE VALUE OF  
METH.  
IF METH IS EQUAL TO 1,  
NMETH IS EQUAL TO  $N*13$ .  
IF METH IS EQUAL TO 2,  
NMETH IS EQUAL TO  $N*6$ .  
THE VALUE OF NMITER DEPENDS ON THE VALUE OF  
MITER.  
IF MITER IS EQUAL TO 1 OR 2,  
NMITER IS EQUAL TO  $N*(N+1)$ .  
IF MITER IS EQUAL TO 3,  
NMITER IS EQUAL TO N.  
IF MITER IS EQUAL TO 0,  
NMITER IS EQUAL TO 1.  
WK MUST REMAIN UNCHANGED BETWEEN SUCCESSIVE  
CALLS DURING INTEGRATION.
- IER - ERROR PARAMETER. (OUTPUT)  
WARNING ERROR  
IER = 33, IMPLIES THAT  $X+H$  WILL EQUAL X ON  
THE NEXT STEP. THIS CONDITION DOES NOT  
FORCE THE ROUTINE TO HALT. HOWEVER, IT  
DOES INDICATE ONE OF TWO CONDITIONS.

THE USER MIGHT BE REQUIRING TOO MUCH ACCURACY VIA THE INPUT PARAMETER TOL. 169  
IN THIS CASE THE USER SHOULD CONSIDER INCREASING THE VALUE OF TOL. THE OTHER CONDITION WHICH MIGHT GIVE RISE TO THIS ERROR MESSAGE IS THAT THE SYSTEM OF DIFFERENTIAL EQUATIONS BEING SOLVED IS STIFF (EITHER IN GENERAL OR OVER THE SUBINTERVAL OF THE PROBLEM BEING SOLVED AT THE TIME OF THE ERROR). IN THIS CASE THE USER SHOULD CONSIDER USING A NONZERO VALUE FOR THE INPUT PARAMETER MITER.

WARNING WITH FIX ERROR

IER = 66, IMPLIES THAT THE ERROR TEST FAILED. H WAS REDUCED BY .1 ONE OR MORE TIMES AND THE STEP WAS TRIED AGAIN SUCCESSFULLY.

IER = 67, IMPLIES THAT CORRECTOR CONVERGENCE COULD NOT BE ACHIEVED. H WAS REDUCED BY .1 ONE OR MORE TIMES AND THE STEP WAS TRIED AGAIN SUCCESSFULLY.

TERMINAL ERROR

IER = 132, IMPLIES THE INTEGRATION WAS HALTED AFTER FAILING TO PASS THE ERROR TEST EVEN AFTER REDUCING H BY A FACTOR OF 1.0E10 FROM ITS INITIAL VALUE. SEE REMARKS.

IER = 133, IMPLIES THE INTEGRATION WAS HALTED AFTER FAILING TO ACHIEVE CORRECTOR CONVERGENCE EVEN AFTER REDUCING H BY A FACTOR OF 1.0E10 FROM ITS INITIAL VALUE. SEE REMARKS.

IER = 134, IMPLIES THAT AFTER SOME INITIAL SUCCESS, THE INTEGRATION WAS HALTED EITHER BY REPEATED ERROR TEST FAILURES OR BY A TEST ON TOL. SEE REMARKS.

IER = 135, IMPLIES THAT ONE OF THE INPUT PARAMETERS N,X,H,XEND,TOL,METH,MITER, OR INDEX WAS SPECIFIED INCORRECTLY.

IER = 136, IMPLIES THAT INDEX HAD A VALUE OF -1 ON INPUT, BUT THE DESIRED CHANGES OF PARAMETERS WERE NOT IMPLEMENTED BECAUSE XEND WAS NOT BEYOND X. INTERPOLATION TO X = XEND WAS PERFORMED. TO TRY AGAIN, SIMPLY CALL AGAIN WITH INDEX EQUAL TO -1 AND A NEW VALUE FOR XEND.

PRECISION/HARDWARE - SINGLE AND DOUBLE/H32  
- SINGLE/H36,H48,H60

REQD. IMSL ROUTINES - DGRCS,DGRIN,DGRPS,DGRST,LUDATF,LUELMF,UERTST,  
UGETIO

- REMARKS
1. THE EXTERNAL SUBROUTINE FCNJ IS USED ONLY WHEN INPUT PARAMETER MITER IS EQUAL TO 1. OTHERWISE, A DUMMY FUNCTION CAN BE USED. THE DUMMY SUBROUTINE SHOULD BE OF THE FOLLOWING FORM
 

```
SUBROUTINE FCNJ (N,X,Y,PD)
  INTEGER N
  REAL Y(N),PD(N,N),X
  RETURN
  END
```
  2. AFTER THE INITIAL CALL, IF A NORMAL RETURN OCCURRED (IER=0) AND A NORMAL CONTINUATION IS DESIRED, SIMPLY RESET XEND AND CALL DGEAR AGAIN. ALL OTHER PARAMETERS WILL BE READY FOR THE NEXT CALL. A CHANGE OF PARAMETERS WITH INDEX EQUAL TO -1 CAN BE MADE AFTER EITHER A SUCCESSFUL OR AN UNSUCCESSFUL RETURN.
  3. THE COMMON BLOCK /GEAR/ USED IN DGEAR NEEDS TO BE PRESERVED BETWEEN CALLS TO DGEAR. IF IT IS NECESSARY FOR THE COMMON BLOCK TO EXIST IN THE CALLING PROGRAM THE FOLLOWING STATEMENT SHOULD BE INCLUDED
 

```
COMMON /GEAR/ DUMMY(48),SDUMMY(4),IDUMMY(38)
```

 WHERE DUMMY, SDUMMY, AND IDUMMY ARE VARIABLE NAMES NOT USED ELSEWHERE IN THE CALLING PROGRAM. (FOR DOUBLE PRECISION DUMMY IS TYPE DOUBLE AND SDUMMY IS TYPE REAL)
  4. THE CHOICE OF VALUES FOR METH AND MITER MAY REQUIRE SOME EXPERIMENTATION, AND ALSO SOME CONSIDERATION OF THE NATURE OF THE PROBLEM AND OF STORAGE REQUIREMENTS. THE PRIME CONSIDERATION IS STIFFNESS. IF THE PROBLEM IS NOT STIFF, THE BEST CHOICE IS PROBABLY METH = 1 WITH MITER = 0. IF THE PROBLEM IS STIFF TO A SIGNIFICANT DEGREE, THEN METH SHOULD BE 2 AND MITER SHOULD BE 1, 2, OR 3. IF THE USER HAS NO KNOWLEDGE OF THE INHERENT TIME CONSTANTS OF THE PROBLEM, WITH WHICH TO PREDICT ITS STIFFNESS, ONE WAY TO DETERMINE THIS IS TO TRY METH = 1 AND MITER = 0 FIRST, AND LOOK AT THE BEHAVIOR OF THE SOLUTION COMPUTED AND THE STEP SIZES USED. IF THE TYPICAL VALUES OF H ARE MUCH SMALLER THAN THE SOLUTION BEHAVIOR WOULD SEEM TO REQUIRE (THAT IS, MORE THAN 100 STEPS ARE TAKEN OVER AN INTERVAL IN WHICH THE SOLUTIONS CHANGE BY LESS THAN ONE PERCENT), THEN THE PROBLEM IS PROBABLY STIFF AND THE DEGREE OF STIFFNESS CAN BE ESTIMATED FROM THE VALUES OF H USED AND THE SMOOTHNESS OF THE SOLUTION. IF THE DEGREE OF STIFFNESS IS ONLY SLIGHT, IT MAY BE THAT METH = 1 IS MORE EFFICIENT THAN METH = 2. EXPERIMENTATION WOULD BE REQUIRED TO DETERMINE THIS. REGARDLESS OF METH, THE LEAST EFFECTIVE VALUE OF MITER IS 0, AND THE MOST EFFECTIVE IS 1 OR 2. MITER = 3 IS GENERALLY SOMEWHERE IN BETWEEN. SINCE THE STORAGE REQUIREMENTS GO UP IN THE SAME ORDER AS EFFECTIVENESS, TRADE-OFF CONSIDERATIONS ARE NECESSARY. FOR REASONS OF ACCURACY AND SPEED, THE CHOICE OF MITER = 1 IS GENERALLY PREFERRED TO MITER = 2, UNLESS THE SYSTEM IS FAIRLY COMPLICATED

(AND FCNJ IS THUS NOT FEASIBLE TO CODE). THE ACCURACY OF THE FCNJ CALCULATION CAN BE CHECKED BY COMPARISON OF THE JACOBIAN WITH THAT GENERATED WITH MITER = 2. IF THE JACOBIAN MATRIX IS SIGNIFICANTLY DIAGONALLY DOMINANT, THEN THE OPTION MITER = 3 IS LIKELY TO BE NEARLY AS EFFECTIVE AS MITER = 1 OR 2, AND WILL SAVE CONSIDERABLE STORAGE AND RUN TIME. IT IS POSSIBLE, AND POTENTIALLY QUITE DESIRABLE, TO USE DIFFERENT VALUES OF METH AND MITER IN DIFFERENT SUBINTERVALS OF THE PROBLEM. FOR EXAMPLE, IF THE PROBLEM IS NON-STIFF INITIALLY AND STIFF LATER, METH = 1 AND MITER = 0 MIGHT BE SET INITIALLY, AND METH = 2 AND MITER = 1 LATER.

5. THE INITIAL VALUE OF THE STEP SIZE, H, SHOULD BE CHOSEN CONSIDERABLY SMALLER THAN THE AVERAGE VALUE EXPECTED FOR THE PROBLEM, AS THE FIRST-ORDER METHOD WITH WHICH DGEAR BEGINS IS NOT GENERALLY THE MOST EFFICIENT ONE. HOWEVER, FOR THE FIRST STEP, AS FOR EVERY STEP, DGEAR TESTS FOR THE POSSIBILITY THAT THE STEP SIZE WAS TOO LARGE TO PASS THE ERROR TEST (BASED ON TOL), AND IF SO ADJUSTS THE STEP SIZE DOWN AUTOMATICALLY. THIS DOWNWARD ADJUSTMENT, IF ANY, IS NOTED BY IER HAVING THE VALUES 66 OR 67, AND SUBSEQUENT RUNS ON THE SAME OR SIMILAR PROBLEM SHOULD BE STARTED WITH AN APPROPRIATELY SMALLER VALUE OF H.
6. SOME OF THE VALUES OF INTEREST LOCATED IN THE COMMON BLOCK /GEAR/ ARE
  - A. HUSED, THE STEP SIZE H LAST USED SUCCESSFULLY (DUMMY(8))
  - B. NQUSED, THE ORDER LAST USED SUCCESSFULLY (IDUMMY(6))
  - C. NSTEP, THE CUMULATIVE NUMBER OF STEPS TAKEN (IDUMMY(7))
  - D. NFE, THE CUMULATIVE NUMBER OF FCN EVALUATIONS (IDUMMY(8))
  - E. NJE, THE CUMULATIVE NUMBER OF JACOBIAN EVALUATIONS, AND HENCE ALSO OF MATRIX LU DECOMPOSITIONS (IDUMMY(9))
7. THE NORMAL USAGE OF DGEAR MAY BE SUMMARIZED AS FOLLOWS
  - A. SET THE INITIAL VALUES IN Y.
  - B. SET N, X, H, TOL, METH, AND MITER.
  - C. SET XEND TO THE FIRST OUTPUT POINT, AND INDEX TO 1.
  - D. CALL DGEAR
  - E. EXIT IF IER IS GREATER THAN 128.
  - F. OTHERWISE, DO DESIRED OUTPUT OF Y.
  - G. EXIT IF THE PROBLEM IS FINISHED.
  - H. OTHERWISE, RESET XEND TO THE NEXT OUTPUT POINT, AND RETURN TO STEP D.
8. THE ERROR WHICH IS CONTROLLED BY WAY OF THE PARAMETER TOL IS AN ESTIMATE OF THE LOCAL TRUNCATION ERROR, THAT IS, THE ERROR COMMITTED ON TAKING A SINGLE STEP WITH THE METHOD, STARTING WITH DATA REGARDED AS EXACT. THIS IS TO BE DISTINGUISHED FROM THE GLOBAL TRUNCATION ERROR, WHICH IS THE ERROR IN ANY GIVEN COMPUTED VALUE OF Y(X) AS A RESULT OF THE LOCAL TRUNCATION ERRORS FROM ALL STEPS TAKEN TO OBTAIN Y(X). THE LATTER ERROR



ACCUMULATES IN A NON-TRIVIAL WAY FROM THE LOCAL ERRORS, AND IS NEITHER ESTIMATED NOR CONTROLLED BY 172 THE ROUTINE. SINCE IT IS USUALLY THE GLOBAL ERROR THAT A USER WANTS TO HAVE UNDER CONTROL, SOME EXPERIMENTATION MAY BE NECESSARY TO GET THE RIGHT VALUE OF TOL TO ACHIEVE THE USERS NEEDS. IF THE PROBLEM IS MATHEMATICALLY STABLE, AND THE METHOD USED IS APPROPRIATELY STABLE, THEN THE GLOBAL ERROR AT A GIVEN X SHOULD VARY SMOOTHLY WITH TOL IN A MONOTONE INCREASING MANNER.

9. IF THE ROUTINE RETURNS WITH IER VALUES OF 132, 133, OR 134, THE USER SHOULD CHECK TO SEE IF TOO MUCH ACCURACY IS BEING REQUIRED. THE USER MAY WISH TO SET TOL TO A LARGER VALUE AND CONTINUE. ANOTHER POSSIBLE CAUSE OF THESE ERROR CONDITIONS IS AN ERROR IN THE CODING OF THE EXTERNAL FUNCTIONS FCN OR FCNJ. IF NO ERRORS ARE FOUND, IT MAY BE NECESSARY TO MONITOR INTERMEDIATE QUANTITIES GENERATED BY THE ROUTINE. THESE QUANTITIES ARE STORED IN THE WORK VECTOR WK AND INDEXED BY SPECIFIC ELEMENTS IN THE COMMON BLOCK /GEAR/. IF IER IS 132 OR 134, THE COMPONENTS CAUSING THE ERROR TEST FAILURE CAN BE IDENTIFIED FROM LARGE VALUES OF THE QUANTITY

$WK(IDUMMY(11)+I)/WK(I)$ , FOR  $I=1, \dots, N$ .

ONE CAUSE OF THIS MAY BE A VERY SMALL BUT NONZERO INITIAL VALUE OF  $ABS(Y(I))$ .

IF IER IS 133, SEVERAL POSSIBILITIES EXIST.

IT MAY BE INSTRUCTIVE TO TRY DIFFERENT VALUES OF MITER.

ALTERNATIVELY, THE USER MIGHT MONITOR SUCCESSIVE CORRECTOR ITERATES CONTAINED IN  $WK(IDUMMY(12)+I)$ , FOR  $I=1, \dots, N$ . ANOTHER POSSIBILITY MIGHT BE TO MONITOR THE JACOBIAN MATRIX, IF ONE IS USED, STORED, BY COLUMN, IN  $WK(IDUMMY(10)+I)$ , FOR  $I=1, \dots, N*N$  IF MITER IS EQUAL TO 1 OR 2, OR FOR  $I=1, \dots, N$  IF MITER IS EQUAL TO 3.

### Algorithm

DGEAR finds approximations to the solution of a system of first order ordinary differential equations of the form  $y'=f(x,y)$  with initial conditions. The basic methods used for the solution are of implicit linear multistep type. There are two classes of such methods available to the user. The first is the implicit Adams methods (up to order twelve), and the second is the backward differentiation formula (BDF) methods (up to order five), also called Gear's stiff methods. In either case the implicitness of the basic formula requires that an algebraic system of equations be solved at each step. A variety of corrector iteration methods is available for this.

DGEAR and the associated nuclei are adaptations of a package designed by A. C. Hindmarsh based on C. W. Gear's subroutine DIFSUB.

See references:

1. Hindmarsh, A.C., "GEAR: Ordinary Differential Equation System Solver", Lawrence Livermore Laboratory, Report UCID-30001, Revision 3, December, 1974..

2. Gear, C.W., Numerical Initial Value Problems in Ordinary Differential Equations, Prentice-Hall, Englewood Cliffs, New Jersey, 1971.

### Example

This example illustrates the basic usage of DGEAR. A table of solution values for  $x = 1.0, 2.0, \dots, 10.0$  is obtained for the predator-prey problem:

$$\begin{aligned} y_1' &= 2y_1(1-y_2) & y_1 &= 1 \\ y_2' &= y_2(y_1-1) & y_2 &= 3 \end{aligned} \quad \text{at } x = 0$$

Input:

```

INTEGER    N,METH,MITER,INDEX,IWK(2),IER,K
REAL      Y(2),WK(35),X,TOL,XEND,H
EXTERNAL  FCN,FCNJ
N         = 2
X         = 0.0
Y(1)     = 1.0
Y(2)     = 3.0
TOL      = .00001
H        = .00001
METH     = 1
MITER    = 0
INDEX    = 1
DO 10 K=1,10
  XEND=FLOAT(K)
  CALL DGEAR (N,FCN,FCNJ,X,H,Y,XEND,TOL,METH,MITER,INDEX,IWK,WK,IER)
  IF(IER.GT.128) GO TO 20
C          Y(1) and Y(2) are current solution values at X.
C          Insert write statement here.
10 CONTINUE
  STOP
20 CONTINUE
C          Handle IER.GT.128
C          Items that may help diagnose the problem should be
C          output here.
C          TOL,N,Y(1),...,Y(N),XEND,H,X,METH,MITER, AND INDEX.
STOP
END
SUBROUTINE FCN(N,X,Y,YPRIME)
INTEGER    N
REAL      Y(N),YPRIME(N),X
YPRIME(1) = 2.0*Y(1)*(1.0-Y(2))
YPRIME(2) = Y(2)*(Y(1)-1.0)
RETURN
END

SUBROUTINE FCNJ(N,X,Y,PD)
INTEGER    N
REAL      Y(N),PD(N,N),X
RETURN
END

```

Output:

174

IER = 0

X	Y(1)	Y(2)
1.	0.08	1.46
2.	0.085	0.58
3.	0.29	.25
4.	1.45	0.19
5.	4.05	1.44
6.	0.18	2.26
7.	0.07	0.91
8.	0.15	0.37
9.	0.65	0.19
10.	3.15	0.35

APPENDIX C

FORTRAN PROGRAM NLLSQ1

REQUESTED OPTIONS (EXECUTE): NODECK,NOLIST,OPT(O)

OPTIONS IN EFFECT: NOLIST NOMAP NOXREF NOGOSTMT NODECK SOURCE TERM OBJECT FIXED NOTEST NOTRMFLG SRCFLG NOSYM  
 OPT(O) LANGLVL(77) NOFIPS FLAG(I) NAME(MAIN ) LINECOUNT(60) CHARLEN(500) SDUMP

	*	1	2	3	4	5	6	7	8
ISN	1								
		C							IMPLICIT REAL *8 (A-H,O-Z)
		C							SAMPLE MAINLINE FOR NLLSQ
									UNIVERSITY OF OKLAHOMA 1975 E.L.ENWALL
ISN	2								DIMENSION X(100,10),Y(100), B(10), IB(10), TITLE(10), JJ(10)
ISN	3								COMMON /ABYC/ NNOB
ISN	4								COMMON /CDYC/ XSTRT,YASTRT,YBSTRT,YCSTRT
ISN	5		1						CONTINUE
		C							BLANK IB
ISN	6								DO 2 K=1,10
ISN	7		2						IB(K)=0
ISN	8								READ (5,12,END=9) TITLE
ISN	9								WRITE (6,13) TITLE
		C							READ INTEGERS
ISN	10								READ (5,10) NO,NP,NOM,M,ID,IP
ISN	11								WRITE (6,15) NO,NP,NOM,M,ID,IP
ISN	12								IF (NO.EQ.O) STOP
ISN	14								NNP=IABS(NP)
ISN	15								NNO=IABS(NO)
ISN	16								MM=IABS(M)
		C							IF NOM.NE.O READ AND SET FLAGS FOR OMITTED PARAMETERS
ISN	17								IF (NOM.LE.O) GO TO 4
ISN	18								READ (5,10) (JJ(K),K=1,NOM)
ISN	19								WRITE (6,14) (JJ(K),K=1,NOM)
ISN	20								DO 3 K=1,NOM
ISN	21								KK=JJ(K)
ISN	22		3						IB(KK)=1
		C							IF IP.EQ.O READ NEW PARAMETERS
ISN	23		4						IF (IP.EQ.O) READ (5,11) (B(K),K=1,NNP)
ISN	25								IF (IP.EQ.O) WRITE (6,16) (B(K),K=1,NNP)
		C							IF NO.LT.O DO NOT READ NEW DATA
ISN	27								IF (NO.LT.O) GO TO 8
ISN	28								READ (5,11) XSTRT,YASTRT,YBSTRT,YCSTRT
ISN	29								DO 5 I=1,NNO
ISN	30								READ (5,11) Y(I),(X(I,K),K=1,MM)
ISN	31		5						CONTINUE
		C							MESSAGE RAW DATA-----USER ROUTINE
ISN	32								DO 50 II=1,NNO
ISN	33								TEMP = Y(II)
ISN	34								Y(II) = X(II,1)/1.307DOO
ISN	35								X(II,1) = TEMP
ISN	36								Y(NNO+II) = X(II,2)/1.222DOO
ISN	37								X(NNO+II,1) = X(II,1)
ISN	38								X(NNO+II,2) = X(II,2)
ISN	39								X(NNO+II,3) = X(II,3)
ISN	40								Y(2*NNO+II) = X(II,3)/2.0DOO
ISN	41								X(2*NNO+II,1) = X(II,1)
ISN	42								X(2*NNO+II,2) = X(II,2)
ISN	43								X(2*NNO+II,3) = X(II,3)
ISN	44		50						CONTINUE
		C							END MESSAGE
ISN	45		8						CONTINUE

\* . \* . 1 . . . 2 . . . 3 . . . 4 . . . 5 . . . 6 . . . 7 . \* . . . 8

```

ISN      46      NNOB= NNO                      00000520
ISN      47      NNO = 3*NNO                    00000530
ISN      48      NCYCLS=20                      00000540
ISN      49      CALL NLLSQ (X,Y,B,IB,NP,NNO,M,IRET,NCYCLS) 00000550
ISN      50      IF (IRET.NE.5) WRITE (6,17) IRET 00000560
ISN      52      GO TO 1                        00000570
ISN      53      9 STOP                          00000580
          C
          C
ISN      54      10 FORMAT (25I3)                00000610
ISN      55      11 FORMAT (8F10.0)             00000620
ISN      56      12 FORMAT (10A8)              00000630
ISN      57      13 FORMAT ('1',10A8)          00000640
ISN      58      14 FORMAT (' INVARIANT PARAMETER SUBSCRIPTS ',10I3) 00000650
ISN      59      15 FORMAT (' NO=',I3,' NP=',I3,' NOM=',I3,' M=',I3,' D=',I3,' IP=',I3,1) 00000660
ISN      60      16 FORMAT (' INITIAL PARAMETER VALUES',/,2(5X,5(1PD15.7))) 00000680
ISN      61      17 FORMAT ('O','PROBLEM IN SUBROUTINE NLLSQ CALCULATIONS IRET= ',I5) 00000690
ISN      62      END                            00000700
    
```

\*\*\* VS FORTRAN ERROR MESSAGES \*\*\*

IFX0030I CNTL O(I) TRMFLG HAS BEEN SPECIFIED BUT SYSTEM IS NOT A TERMINAL. TRMFLG CANCELED.

\*STATISTICS\* SOURCE STATEMENTS = 58, PROGRAM SIZE = 13062 BYTES, PROGRAM NAME = MAIN PAGE: 1.

\*STATISTICS\* 1 DIAGNOSTIC GENERATED. SEVERITY CODE IS 0.

\*\*\*\*\* END OF COMPILATION 1 \*\*\*\*\*

OPTIONS IN EFFECT: NDLIST NUMAP NOXREF NOGOSTMT NODECK SOURCE TERM OBJECT FIXED NOTEST NOTRMFLG SRCFLG NOSYM  
OPT(O) LANGLVL(77) NOFIPS FLAG(I) NAME(MAIN ) LINECOUNT(60) CHARLEN(500) SDUMP

\*.....1.....2.....3.....4.....5.....6.....7.\*.....8

ISN	1	SUBROUTINE DCALC (X,Y,B,D,NP,I)	00000710
ISN	2	IMPLICIT REAL *8 (A-H,O-Z)	00000720
ISN	3	DIMENSION X(100,10),Y(100),B(10),D(10)	00000730
	C	D CALC ROUTINE SPECIFIC TO RUN OR DUMMY	00000740
ISN	4	RETURN	00000750
ISN	5	END	00000760

\*STATISTICS\* SOURCE STATEMENTS = 5, PROGRAM SIZE = 432 BYTES, PROGRAM NAME = DCALC PAGE: 3.

\*STATISTICS\* NO DIAGNOSTICS GENERATED.

\*\*\*\*\* END OF COMPILATION 2 \*\*\*\*\*

OPTIONS IN EFFECT: NOLIST NOMAP NOXREF NOGOSTMT NODECK SOURCE TERM OBJECT FIXED NOTEST NOTRMFLG SRCFLG NOSYM  
 OPT(O) LANGLVL(77) NOFIPS FLAG(I) NAME(MAIN ) LINECOUNT(60) CHARLEN(500) SDUMP

\*.....1.....2.....3.....4.....5.....6.....7.\*.....8

```

ISN      1      SUBROUTINE YCALC (X,Y,B,YC,WDEL,I)      00000770
ISN      2      IMPLICIT REAL*8 (A-H,O-Z)            00000780
ISN      3      DIMENSION X(100,10),Y(100),B(10),YY(4),WK(400) 00000790
ISN      4      COMMON XKT,XKR,F                      00000800
ISN      5      COMMON /ABYC/ NNOB                    00000810
ISN      6      COMMON /CDYC/ XSTRT,YASTRT,YBSTRT,YCSTRT 00000820
ISN      7      EXTERNAL FCN,FCNJ                    00000830
ISN      8      N = 4                                  00000840
ISN      9      METH = 1                              00000850
ISN     10      MITER = 0                             00000860
ISN     11      INDEX = 1                             00000870
ISN     12      F = 0.5D00                            00000880
ISN     13      XKT = B(1)                             00000890
ISN     14      XKR = B(2)                             00000900
ISN     15      XX = XSTRT                             00000910
ISN     16      YY(1) = YASTRT/1.307D00               00000920
ISN     17      YY(2) = YBSTRT/1.22D00                00000930
ISN     18      YY(3) = YBSTRT*0.0D00                 00000940
ISN     19      YY(4) = YCSTRT/2.0D00-YY(3)           00000945
ISN     20      H = 1.0D-05                           00000950
ISN     21      TOL = 1.0D-04                         00000960
ISN     22      XEND = X(I,1)                          00000970
ISN     23      IF (I.GT.(2*NNOB)) GOTO 20             00000980
ISN     24      IF (I.GT.NNOB) GOTO 10                 00000990
ISN     25      CALL DGEAR (N,FCN,FCNJ,XX,H,YY,XEND,TOL,METH,MITER,INDEX,IWK,WK, 00001000
ISN          1IER)                                     00001010
ISN     26      YC = YY(1)                             00001020
ISN     27      X(I,4) = YY(1)                         00001030
ISN     28      X(I,5) = YY(3)                         00001040
ISN     29      WDEL = Y(I)-YC                         00001050
ISN     30      RETURN                                  00001060
ISN     31      10 CALL DGEAR (N,FCN,FCNJ,XX,H,YY,XEND,TOL,METH,MITER,INDEX,IWK,WK, 00001070
ISN          1IER)                                     00001080
ISN     32      YC = YY(2)                             00001090
ISN     33      X(I,4) = YY(2)                         00001100
ISN     34      X(I,5) = YY(3)                         00001110
ISN     35      WDEL = Y(I)-YC                         00001120
ISN     36      RETURN                                  00001130
ISN     37      20 CALL DGEAR (N,FCN,FCNJ,XX,H,YY,XEND,TOL,METH,MITER,INDEX,IWK,WK, 00001140
ISN          1IER)                                     00001150
ISN     38      YC = YY(3)+YY(4)                       00001160
ISN     39      X(I,4) = YY(1)                         00001170
ISN     40      X(I,5) = YY(4)                         00001180
ISN     41      WDEL = Y(I)-YC                         00001190
ISN     42      RETURN                                  00001200
ISN     43      END                                    00001210
    
```

\*STATISTICS\* SOURCE STATEMENTS = 43, PROGRAM SIZE = 5752 BYTES, PROGRAM NAME = YCALC PAGE: 4.

\*STATISTICS\* NO DIAGNOSTICS GENERATED.

\*\*\*\*\* END OF COMPILATION 3 \*\*\*\*\*



OPTIONS IN EFFECT: NOLIST NOMAP NOXREF NOGOSTMT NODECK SOURCE TERM OBJECT FIXED NOTEST NOTRMFLG SRCFLG NOSYM  
OPT(O) LANGLVL(77) NOFIPS FLAG(I) NAME(MAIN ) LINECOUNT(60) CHARLEN(500) SDUMP

\*.....\*...1.....2.....3.....4.....5.....6.....7.\*.....8

ISN	1	SUBROUTINE FCN (N,X,Y,YPRIME)	00001220
ISN	2	IMPLICIT REAL*8 (A-H,O-Z)	00001230
ISN	3	DIMENSION Y(N),YPRIME(N),B(10)	00001240
ISN	4	COMMON XKI,XKR,F	00001250
ISN	5	YPRIME(1) = 0.25000*XKT+Y(2)-XKT*Y(1)	00001260
ISN	6	YPRIME(2) = XKT*Y(1)-0.25000*XKT*Y(2)-XKR*Y(2)+XKR*Y(3)	00001270
ISN	7	YPRIME(3) = XKR*Y(2)-XKR*Y(3)-F*XKT*Y(3)	00001280
ISN	8	YPRIME(4) = F*XKT*Y(3)	00001285
ISN	9	RETURN	00001290
ISN	10	END	00001300

\*STATISTICS\* SOURCE STATEMENTS = 10, PROGRAM SIZE = 1038 BYTES, PROGRAM NAME = FCN PAGE: 5.

\*STATISTICS\* NO DIAGNOSTICS GENERATED.

\*\*\*\*\* END OF COMPILATION 4 \*\*\*\*\*

OPTIONS IN EFFECT: NOLIST NOMAP NOXREF NOGOSTMT NODECK SOURCE TERM OBJECT FIXED NOTEST NOTRMFLG SRCFLG NOSYM  
OPT(O) LANGLVL(77) NOFIPS FLAG(I) NAME(MAIN ) LINECOUNT(60) CHARLEN(500) SDUMP

\*...\*...1.....2.....3.....4.....5.....6.....7.\*.....8

ISN	1	SUBROUTINE FCNJ (N,X,Y,PD)	00001310
ISN	2	IMPLICIT REAL*8 (A-H,O-Z)	00001320
ISN	3	DIMENSION Y(N),PD(N,N)	00001330
ISN	4	RETURN	00001340
ISN	5	END	00001350

\*STATISTICS\* SOURCE STATEMENTS = 5, PROGRAM SIZE = 556 BYTES, PROGRAM NAME = FCNJ PAGE: 6.

\*STATISTICS\* NO DIAGNOSTICS GENERATED.

\*\*\*\*\* END OF COMPILATION 5 \*\*\*\*\*

OPTIONS IN EFFECT: NOLIST NOMAP NOXREF NOGOSTMT NODECK SOURCE TERM OBJECT FIXED NOTEST NOTRMFLG SRCFLG NOSYM  
 OPT(O) LANGLVL(77) NOFIPS FLAG(I) NAME(MAIN ) LINECOUNT(60) CHARLEN(500) SDUMP

```

      *.....*.....1.....2.....3.....4.....5.....6.....7.*.....8
ISN      1      SUBROUTINE NLLSQ (X,Y,B,IB,NP,NO,MM,IRET,NCYCLS)      00001360
      C      NON LINEAR LEAST SQUARES ROUTINE USING MARQUARDT CONVERG STRATEGY 00001370
      C      J.SOC.INDUST.APPL.MATH.VOL.11.NO.2.JUNE 1963--D.W.MARQUARDT      00001380
      C      DIMENSIONED FOR 10 VARIABLES, 100 OBSERVATIONS, 5 IND. VARIABLES 00001390
      C      NP NEGATIVE CALCULATES ESTIMATED DERIVATIVES      00001400
      C      MM NEGATIVE PRINTS B VALUES ON EVERY CYCLE      00001410
      C      WRITTEN U OF OKLA MARCH 1975 E.L.ENWALL      00001420
ISN      2      IMPLICIT REAL *8 (A-H,O-Z)      00001430
ISN      3      REAL*8 PHIR,PHI,PHIL10,PHILB10,DEN,DARCOS,SUMA      00001440
ISN      4      DIMENSION X(100,10), Y(100), B(10), D(10), IB(10)      00001450
ISN      5      DIMENSION A(10,10), AS(10,10), S(100), BS(10), DB(10), DDB(10), DG      00001460
      1(10), V(10), AAS(10,10), AM(100)      00001470
ISN      6      REAL GAMZRO/4,BE01/,DDEL/5.E-5/,EPS/2.5E-09/,FEPS/1.E-6/,LAMBDA      00001480
ISN      7      COMMON XK1,XK4,XK2      00001490
ISN      8      DATA 101/'B('/,102/'')='/,NIP/O/      00001500
***ERROR 1755(W)***      A CHARACTER CONSTANT HAS BEEN ASSIGNED TO A NON-CHARACTER VARIABLE OR ARRAY. INITIALIZE WITH CONSTANT OF
      LIKE TYPE TO AVOID ERROR MESSAGE.
***ERROR 1755(W)***      A CHARACTER CONSTANT HAS BEEN ASSIGNED TO A NON-CHARACTER VARIABLE OR ARRAY. INITIALIZE WITH CONSTANT OF
      LIKE TYPE TO AVOID ERROR MESSAGE.
ISN      9      IF (NCYCLS.LT.1.OR.NCYCLS.GT.100) NCYCLS=20      00001510
ISN     11      LAMBDA=0.01000001      00001520
ISN     12      ND=1      00001530
ISN     13      NPRT=1      00001540
ISN     14      IRET=0      00001550
ISN     15      IF (NP.LE.0) ND=-1      00001560
ISN     17      IF (MM.LE.0) NPRT=-1      00001570
ISN     19      NP=IABS(NP)      00001580
ISN     20      IF (NP.NE.0) GO TO 101      00001590
ISN     21      WRITE (6,189)      00001600
ISN     22      RETURN      00001610
ISN     23      101 MM=IABS(MM)      00001620
ISN     24      WRITE (6,176)      00001630
ISN     25      NCYC=0      00001640
ISN     26      102 CONTINUE      00001650
      C      TEST FOR TOO MANY CYCLES      00001660
ISN     27      IF (NCYC.GT.NCYCLS) GO TO 147      00001670
ISN     28      NCYC=NCYC+1      00001680
      C      BUILD MATRIX      00001690
ISN     29      NIP=0      00001700
ISN     30      DO 103 K=1,NP      00001710
ISN     31      IF (IB(K).NE.0) NIP=NIP+1      00001720
ISN     33      D(K)=0.0      00001730
ISN     34      V(K)=0.0      00001740
ISN     35      DO 103 L=1,NP      00001750
ISN     36      A(K,L)=0.0      00001760
ISN     37      103 CONTINUE      00001770
ISN     38      PIH=0.0      00001780
ISN     39      DO 109 I=1,NO      00001790
ISN     40      CALL YCALC (X,Y,B,YC,WDEL,I)      00001800
ISN     41      IF (ND.GT.0) GO TO 105      00001810
ISN     42      SDEL=WDEL      00001820
ISN     43      DO 104 K=1,NP      00001830
ISN     44      IF (IB(K).NE.0) GO TO 104      00001840

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      * . . . 1 . . . 2 . . . 3 . . . 4 . . . 5 . . . 6 . . . 7 . . . * . . . 8
ISN   45      BZ=B(K)                                00001850
ISN   46      DD=BZ*DDEL                             00001860
ISN   47      B(K)=BZ+DD                             00001870
ISN   48      CALL YCALC (X,Y,B,YC,WDEL,I)           00001880
ISN   49      D(K)=(SDEL-WDEL)/DD                    00001890
ISN   50      B(K)=BZ                                00001900
ISN   51      104 CONTINUE                            00001910
ISN   52      WDEL=SDEL                              00001920
ISN   53      GO TO 106                               00001930
ISN   54      105 CALL DCALC (X,Y,B,D,NP,I)           00001940
ISN   55      106 PHI=PHI+WDEL*WDEL                  00001950
ISN   56      DO 108 K=1,NP                           00001960
ISN   57      IF (IB(K).NE.O) GO TO 108               00001970
ISN   58      V(K)=V(K)+D(K)*WDEL                     00001980
ISN   59      DO 107 L=1,NP                           00001990
ISN   60      IF (IB(L).NE.O) GO TO 107               00002000
ISN   61      A(K,L)=A(K,L)+D(K)*D(L)                 00002010
ISN   62      107 CONTINUE                            00002020
ISN   63      108 CONTINUE                            00002030
ISN   64      109 CONTINUE                            00002040
ISN   65      PHIR=PHI                                00002050
ISN   66      C   CALC RMSD                            00002060
ISN   67      RMSD=(PHIR/(NO-NP+NIP))**.5             00002070
ISN   68      IF (NPRT.LE.O.OR.NCYC.EQ.1) WRITE (6,182) NCYC,RMSD,LAMBDA 00002080
ISN   69      IF (NPRT.LE.O.OR.NCYC.EQ.1) WRITE (6,187) (IO1,K,IO2,B(K),K=1,NP) 00002090
ISN   70      M=O                                     00002100
ISN   71      C   SAVE DIAGONAL ELEMENTS               00002110
ISN   72      DO 110 K=1,NP                           00002120
ISN   73      IF (IB(K).NE.O) A(K,K)=1.O              00002130
ISN   74      IF (A(K,K).LE.O.O) GO TO 169             00002140
ISN   75      110 DG(K)=A(K,K)**.5                     00002150
ISN   76      DO 111 K=1,NP                           00002160
ISN   77      V(K)=V(K)/DG(K)                         00002170
ISN   78      DO 111 L=1,NP                           00002180
ISN   79      C   SAVE ORIGINAL MATRIX                 00002190
ISN   80      AAS(K,L)=A(K,L)                         00002200
ISN   81      C   SCALE MATRIX                         00002210
ISN   82      111 A(K,L)=A(K,L)/(DG(K)*DG(L))          00002220
ISN   83      C   SAVE MATRIX, ORIG PARAMS             00002230
ISN   84      DO 112 K=1,NP                           00002240
ISN   85      BS(K)=B(K)                              00002250
ISN   86      DO 112 L=1,NP                           00002260
ISN   87      AS(K,L)=A(K,L)                          00002270
ISN   88      112 IF (K.EQ.L) A(K,L)=A(K,L)+LAMBDA    00002280
ISN   89      IRET=1                                   00002290
ISN   90      GO TO 149                                00002300
ISN   91      113 CONTINUE                            00002310
ISN   92      DO 114 K=1,NP                           00002320
ISN   93      114 DDB(K)=DB(K)                        00002330
ISN   94      C   CHECK FOR CONVERGENCE                00002340
ISN   95      DO 115 K=1,NP                           00002350
ISN   96      Z=DDB(K)**2/(FEPS+B(K)**2)              00002360
ISN   97      IF (Z.GT.EPS) GO TO 116                 00002370
ISN   98      115 CONTINUE                            00002380
ISN   99      WRITE (6,179)                           00002390
ISN  100      GO TO 134                               00002400

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* . . . 1 . . . 2 . . . 3 . . . 4 . . . 5 . . . 6 . . . 7 . . . 8
C      CALC PHI (LAMBDA)                                00002410
ISN    99      116 PHIL=0.0                               00002420
ISN   100      DO 117 I=1,NO                               00002430
ISN   101      CALL YCALC (X,Y,B,YC,WDEL,I)               00002440
ISN   102      117 PHIL=PHIL+WDEL*WDEL                   00002450
C      TEST PHIL                                         00002460
ISN   103      IF (PHIL.LE.PHIR) GO TO 128                00002470
ISN   104      118 CONTINUE                               00002480
ISN   105      IF (LAMBDA.GE.1.0) GO TO 123                00002490
C      DO FOR LAMBDA*10                                   00002500
ISN   106      LAMBDA=LAMBDA*10                           00002510
ISN   107      DO 119 K=1,NP                               00002520
ISN   108      DO 119 I=1,NP                               00002530
ISN   109      A(K,L)=AS(K,L)                             00002540
ISN   110      119 IF (K.EQ.L) A(K,L)=A(K,L)+LAMBDA      00002550
ISN   112      IRET=2                                     00002560
ISN   113      GO TO 149                                   00002570
ISN   114      120 CONTINUE                               00002580
C      CALC PHI (LAMBDA*10)                              00002590
ISN   115      121 PHIL10=0.0                             00002600
ISN   116      DO 122 I=1,NO                               00002610
ISN   117      CALL YCALC (X,Y,B,YC,WDEL,I)               00002620
ISN   118      122 PHIL10=PHIL10+WDEL*WDEL               00002630
ISN   119      IF (PHIL10.GE.PHIR.AND.LAMBDA.GE.1.0) GO TO 123 00002640
ISN   120      IF (PHIL10.GE.PHIR) GO TO 118              00002650
C      NEW LAMBDA, NEW CYCLE                              00002660
ISN   121      GO TO 102                                  00002670
C      LARGE LAMBDA, GAMMA TEST                          00002680
ISN   122      123 CONTINUE                               00002690
ISN   123      DTG=0.0                                     00002700
ISN   124      GTG=0.0                                     00002710
ISN   125      BTB=0.0                                     00002720
ISN   126      DO 124 K=1,NP                               00002730
ISN   127      DTG=DTG+DB(K)*DG(K)                        00002740
ISN   128      GTG=GTG+DG(K)*DG(K)                       00002750
ISN   129      124 BTB=BTB+DB(K)*DB(K)                   00002760
ISN   130      DEN=DTG/((BTB*GTG)**.5)                   00002770
ISN   131      GAM=(180./3.141592)*DARCOS(DEN)            00002780
ISN   132      IF (NPRT.LE.0.OR.NCYC.EQ.1) WRITE (6,177) GAM,LAMBDA 00002790
ISN   134      IF (GAM.GE.90.0) GO TO 148                 00002800
ISN   135      IF (GAM.GE.GAMZRO) GO TO 102               00002810
ISN   136      DO 125 K=1,NP                               00002820
ISN   137      DB(K)=DB(K)/2.                             00002830
ISN   138      125 B(K)=BS(K)+DB(K)                       00002840
ISN   139      DO 126 K=1,NP                               00002850
ISN   140      Z=DB(K)**2/(FEPS+B(K)**2)                 00002860
ISN   141      IF (Z.GT.EPS) GO TO 127                   00002870
ISN   142      126 CONTINUE                               00002880
ISN   143      GO TO 146                                   00002890
ISN   144      127 GO TO 121                               00002900
C      CALC PHI (LAMBDA/10)                              00002910
ISN   145      128 LAMBDA=LAMBDA/10.                     00002920
ISN   146      PHLB10=0.0                                 00002930
ISN   147      DO 129 K=1,NP                               00002940
ISN   148      DO 129 L=1,NP                               00002950
ISN   149      A(K,L)=AS(K,L)                             00002960

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*.....1.....2.....3.....4.....5.....6.....7.*.....8
ISN 150 129 IF (K.EQ.L) A(K,L)=A(K,L)+LAMBDA 00002970
ISN 152 IRET=3 00002980
ISN 153 GO TO 149 00002990
ISN 154 130 CONTINUE 00003000
ISN 155 DO 131 I=1,NO 00003010
ISN 156 CALL YCALC (X,Y,B,YC,WDEL,I) 00003020
ISN 157 131 PHLB10=PHLB10+WDEL*WDEL 00003030
ISN 158 IF (PHLB10.GT.PHIR) GO TO 132 00003040
ISN 159 GO TO 102 00003050
C RESTORE STARTING VALUE OF LAMBDA 00003060
ISN 160 132 LAMBDA=LAMBDA*10. 00003070
ISN 161 DO 133 K=1,NP 00003080
ISN 162 133 B(K)=BS(K)+DDB(K) 00003090
ISN 163 GO TO 102 00003100
C CONVERGENCE 00003110
ISN 164 134 CONTINUE 00003120
ISN 165 WRITE (6,178) NCGC,PHIR,RMSD 00003130
C SET UP FOR FINAL MATRIX INVERSION, CALC EST CONDITION NUMBER 00003140
C SAVE SCALED MATRIX IN AS 00003150
ISN 166 AMIN=1.E25 00003160
ISN 167 AMAX=0.0 00003170
ISN 168 DO 135 K=1,NP 00003180
ISN 169 IF (AAS(K,K).GT.AMAX) AMAX=AAS(K,K) 00003190
ISN 171 IF (AAS(K,K).LT.AMIN) AMIN=AAS(K,K) 00003200
ISN 173 DO 135 L=1,NP 00003210
ISN 174 A(K,L)=AAS(K,L)/(DG(K)*DG(L)) 00003220
ISN 175 135 AS(K,L)=A(K,L) 00003230
ISN 176 CONDNO=AMAX/AMIN 00003240
ISN 177 WRITE (6,185) CONDNO 00003250
ISN 178 IRET=4 00003260
ISN 179 GO TO 149 00003270
ISN 180 136 CONTINUE 00003280
C SCALED INVERSE IN A 00003290
ISN 181 DO 137 K=1,NP 00003300
ISN 182 DO 137 L=1,NP 00003310
ISN 183 137 AAS(K,L)=A(K,L)/(DG(K)*DG(L)) 00003320
C TRUE INVERSE IN AAS 00003330
ISN 184 DO 138 K=1,NP 00003340
ISN 185 138 B(K)=BS(K) 00003350
ISN 186 DO 139 K=1,NP 00003360
ISN 187 139 DG(K)=AAS(K,K)**.5 00003370
ISN 188 DO 140 K=1,NP 00003380
ISN 189 DO 140 L=1,NP 00003390
ISN 190 140 AAS(K,L)=AAS(K,L)/(DG(K)*DG(L)) 00003400
C AAS NOW CORRELATION MATRIX 00003410
C CALC STD ERRORS IN EACH PARAMETER 00003420
ISN 191 DO 141 K=1,NP 00003430
ISN 192 DB(K)=DG(K)*RMSD 00003440
ISN 193 IF (IB(K).NE.O) DB(K)=O.O 00003450
ISN 195 141 CONTINUE 00003460
ISN 196 WRITE (6,170) 00003470
ISN 197 WRITE (6,171) (K,BS(K),DB(K),K=1,NP) 00003480
ISN 198 WRITE (6,173) 00003490
ISN 199 DO 142 I=1,NO 00003500
ISN 200 CALL YCALC (X,Y,B,YC,WDEL,I) 00003510
ISN 201 WRITE (6,172) Y(I),YC,WDEL,(X(I,K),K=1,MM) 00003520

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*.....*.....1.....2.....3.....4.....5.....6.....7.....*.....8
ISN      202      142 CONTINUE                                00003530
ISN      203      WRITE (6,175)                            00003540
ISN      204      DO 143 K=1,NP                             00003550
ISN      205      143 WRITE (6,174) K,(AAS(K,L),L=1,NP)    00003560
ISN      206      DO 144 K=1,NP                             00003570
ISN      207      DO 144 L=1,NP                             00003580
ISN      208      AAS(K,L)=0.0                              00003590
ISN      209      DO 144 J=1,NP                             00003600
ISN      210      144 AAS(K,L)=AAS(K,L)+A(K,J)*AS(J,L)      00003610
ISN      211      C WRITE OUT PRODUCT OF INVERSE AND STARTING MATRICES
ISN      212      WRITE (6,184)                              00003620
ISN      213      DO 145 K=1,NP                             00003630
ISN      214      145 WRITE (6,183) K,(AAS(K,L),L=1,NP)    00003640
ISN      215      IRET=5                                     00003650
ISN      216      RETURN                                     00003660
ISN      217      C G L TEST                                 00003670
ISN      218      146 WRITE (6,180)                          00003680
ISN      219      GO TO 134                                  00003690
ISN      220      C TOO MANY CYCLES                          00003700
ISN      221      147 CONTINUE                                00003710
ISN      222      WRITE (6,181)                              00003720
ISN      223      GO TO 134                                  00003730
ISN      224      148 WRITE (6,186) GAM                      00003740
ISN      225      GO TO 134                                  00003750
ISN      226      149 CONTINUE                                00003760
ISN      227      C INTERNAL SUBROUTINE TO INVERT MATRIX CHOLESKI METHOD
ISN      228      M=1                                        00003770
ISN      229      DO 150 K=1,NP                             00003780
ISN      230      DO 150 L=K,NP                             00003790
ISN      231      AM(M)=A(K,L)                              00003800
ISN      232      150 M=M+1                                  00003810
ISN      233      K=1                                        00003820
ISN      234      IF (NP.GT.1) GO TO 151                    00003830
ISN      235      AM(1)=1.0/AM(1)                            00003840
ISN      236      GO TO 165                                  00003850
ISN      237      151 DO 158 M=1,NP                          00003860
ISN      238      IMAX=M-1                                   00003870
ISN      239      DO 157 L=M,NP                              00003880
ISN      240      SUMA=0.0                                   00003890
ISN      241      KLI=L                                     00003900
ISN      242      KMI=M                                     00003910
ISN      243      IF (IMAX) 154,154,152                    00003920
ISN      244      152 DO 153 I=1,IMAX                        00003930
ISN      245      SUMA=SUMA+AM(KLI)*AM(KMI)                 00003940
ISN      246      J=NP-I                                    00003950
ISN      247      KLI=KLI+J                                  00003960
ISN      248      153 KMI=KMI+J                              00003970
ISN      249      154 TERM=AM(K)-SUMA                       00003980
ISN      250      IF (L-M) 155,155,156                     00003990
ISN      251      155 IF (TERM.LE.0.0) GO TO 169            00004000
ISN      252      DENOM=TERM**0.5                           00004010
ISN      253      AM(K)=DENOM                                00004020
ISN      254      GO TO 157                                  00004030
ISN      255      156 AM(K)=TERM/DENOM                      00004040
ISN      256      157 K=K+1                                  00004050
ISN      257      158 CONTINUE                               00004060

```



```

* . . . * . . . 1 . . . 2 . . . 3 . . . 4 . . . 5 . . . 6 . . . 7 . . . * . . . 8
ISN   254      AM(1)=1.0/AM(1)                                00004090
ISN   255      KDM=1                                          00004100
ISN   256      DO 161 L=2,NP                                  00004110
ISN   257      KDM=KDM+NP-L+2                                  00004120
ISN   258      TERM=1.0/AM(KDM)                              00004130
ISN   259      AM(KDM)=TERM                                    00004140
ISN   260      KMI=0                                          00004150
ISN   261      KLI=L                                          00004160
ISN   262      IMAX=L-1                                        00004170
ISN   263      DO 160 M=1,IMAX                                00004180
ISN   264      K=KLI                                          00004190
ISN   265      SUMA=0.0                                        00004200
ISN   266      DO 159 I=M,IMAX                                00004210
ISN   267      II=KMI+I                                       00004220
ISN   268      SUMA=SUMA-AM(KLI)*AM(II)                       00004230
ISN   269      159 KLI=KLI+NP-I                                00004240
ISN   270      AM(K)=SUMA*TERM                                00004250
ISN   271      J=NP-M                                          00004260
ISN   272      KLI=K+J                                         00004270
ISN   273      160 KMI=KMI+J                                    00004280
ISN   274      161 CONTINUE                                    00004290
ISN   275      K=1                                             00004300
ISN   276      DO 164 M=1,NP                                  00004310
ISN   277      KLI=K                                           00004320
ISN   278      DO 163 L=M,NP                                  00004330
ISN   279      KMI=K                                           00004340
ISN   280      IMAX=NP-L+1                                     00004350
ISN   281      SUMA=0.0                                        00004360
ISN   282      DO 162 I=1,IMAX                                00004370
ISN   283      SUMA=SUMA+AM(KLI)*AM(KMI)                     00004380
ISN   284      KLI=KLI+1                                       00004390
ISN   285      162 KMI=KMI+1                                    00004400
ISN   286      AM(K)=SUMA                                       00004410
ISN   287      163 K=K+1                                       00004420
ISN   288      164 CONTINUE                                    00004430
ISN   289      165 CONTINUE                                    00004440
ISN   290      M=1                                             00004450
ISN   291      DO 166 K=1,NP                                  00004460
ISN   292      DO 166 L=K,NP                                  00004470
ISN   293      A(K,L)=AM(M)                                    00004480
ISN   294      A(L,K)=AM(M)                                    00004490
ISN   295      166 M=M+1                                       00004500
ISN   296      DO 168 K=1,NP                                  00004510
ISN   297      DB(K)=0.0                                       00004520
ISN   298      IF (IB(K).NE.0) GO TO 168                      00004530
ISN   299      DO 167 L=1,NP                                  00004540
ISN   300      167 DB(K)=DB(K)+A(K,L)*V(L)                    00004550
ISN   301      DB(K)=DB(K)/DG(K)                                00004560
ISN   302      168 B(K)=BS(K)+DB(K)                             00004570
ISN   303      GO TO (113,120,130,136), IRET                  00004580
ISN   304      169 WRITE (6,188) M,NCYC,IRET                  00004590
ISN   305      IRET=IRET+1000*NCYC+M*1000**2                 00004600
ISN   306      RETURN                                          00004610
C
C
ISN   307      170 FORMAT ('OPARAMETERS AND SIGMAS')          00004620
C
C
ISN   307      170 FORMAT ('OPARAMETERS AND SIGMAS')          00004630
C
C
ISN   307      170 FORMAT ('OPARAMETERS AND SIGMAS')          00004640

```



\*.....1.....2.....3.....4.....5.....6.....7.\*.....8

```

ISN      308      171 FORMAT ('OB(',I2,')=' ,IPE13.5, ' +- ',E13.5)                00004650
ISN      309      172 FORMAT (1X,1P3E11.3,(T35,5E10.3))                          00004660
***ERROR 1264(W)***      SLASH, COLON, OR COMMA EXPECTED AFTER DELIMITER GROUP "2".  COMMA ASSUMED.
ISN      310      173 FORMAT ('O',5X,'Y',10X,'YC',10X,'DEL',8X,'X-1-M')          00004670
ISN      311      174 FORMAT (15,5X,10F7.3)                                     00004680
ISN      312      175 FORMAT (' CORRELATION MATRIX')                            00004690
ISN      313      176 FORMAT ('ONON LINEAR LSQ ROUTINE WITH MARQUARDT STRATEGY')    00004700
ISN      314      177 FORMAT (' GAMMA LAMBDA TEST GAMMA=',F4.0, '; LAMBDA=',IPE7.0) 00004710
ISN      315      178 FORMAT (' IN',I3, ' CYCLES; SUM WDEL SQ=',IPE12.5, '; RMSD=',E12.6)00004720
ISN      316      179 FORMAT (' CONVERGENCE')                                    00004730
ISN      317      180 FORMAT (' GAMMA EPSILON TEST -- CONVERGENCE')              00004740
ISN      318      181 FORMAT (' NCYCLES EXCEEDED')                              00004750
ISN      319      182 FORMAT (' CYCLE',I2, '; RMSD=',IPE12.5, '; LAMBDA=',E7.1)    00004760
ISN      320      183 FORMAT (15,5X,5(IPE13.4),/,12X,5E13.4)                   00004770
ISN      321      184 FORMAT ('      B*B INVERSE')                              00004780
ISN      322      185 FORMAT (' MATRIX CONDITION NUMBER=',IPE7.1)              00004790
ISN      323      186 FORMAT (' GAMMA TEST NO CONVERGENCE GAMMA=',F4.0)         00004800
ISN      324      187 FORMAT (5X,5(A3,I2,A2,IPE10.3))                          00004810
ISN      325      188 FORMAT (' MATRIX INVERSION ERROR M=',I5, ' NCYC=',I5, ' IRET=',I5) 00004820
ISN      326      189 FORMAT (' NUM OF PARAMETERS = 0')                        00004830
ISN      327      END                                                            00004840

```

\*\*\* VS FORTRAN ERROR MESSAGES \*\*\*

```

IFX1755I  DATA    4(W)    8    A CHARACTER CONSTANT HAS BEEN ASSIGNED TO A NON-CHARACTER VARIABLE OR ARRAY. INITIALIZE WITH
CONSTANT OF LIKE TYPE TO AVOID ERROR MESSAGE.

IFX1755I  DATA    4(W)    8    A CHARACTER CONSTANT HAS BEEN ASSIGNED TO A NON-CHARACTER VARIABLE OR ARRAY. INITIALIZE WITH
CONSTANT OF LIKE TYPE TO AVOID ERROR MESSAGE.

IFX1264I  FORM     4(W)   309    SLASH, COLON, OR COMMA EXPECTED AFTER DELIMITER GROUP "2".  COMMA ASSUMED.

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\*STATISTICS\* SOURCE STATEMENTS = 313, PROGRAM SIZE = 22746 BYTES, PROGRAM NAME = NLLSQ PAGE: 7.

\*STATISTICS\* 3 DIAGNOSTICS GENERATED. HIGHEST SEVERITY CODE IS 4.

\*\*\*\*\* END OF COMPILATION 6 \*\*\*\*\*

APPENDIX D

SAS PROGRAM ACTIV

SAS PROGRAM ACTIV TO FIND ACTIVATION ENERGIES  
FROM RATE CONSTANT DATA

```

//ACTIV    JOB (13530,240-94-9716), 'COLLIER', TIME=(00,40), CLASS=A,
//  MSGCLASS=A
// *ROUTE PRINT LOCAL
// *PASSWORD ?
//  EXEC SAS
// *SYSIN DD *
MACRO SOLVE
PROC GLM DATA=SUBDATA;
  TITLE HEADING;
  MODEL LKT=RETEMP;
  OUTPUT OUT=SUBDATA P=LPKT;
PROC CORR DATA=SUBDATA;
  VAR LPKT LKT;
PROC PLOT DATA=SUBDATA;
  PLOT LKT*RETEMP LPKT*RETEMP='P' / OVERLAY;
PROC PRINT DATA=SUBDATA;
  VAR TEMP KT RETEMP LKT LPKT;
PROC GLM DATA=SUBDATA;
  MODEL LKR=RETEMP;
  OUTPUT OUT=SUBDATA P=LPKR;
PROC CORR DATA=SUBDATA;
  VAR LPKR LKR;
PROC PLOT DATA=SUBDATA;
  PLOT LKR*RETEMP LPKR*RETEMP='P' / OVERLAY;
PROC PRINT DATA=SUBDATA;
  VAR TEMP KR RETEMP LKR LPKR%
MACRO INDATA
DATA SUBDATA;
  INPUT TEMP KT KR;
  LKT=LOG(KT);
  LKR=LOG(KR);
  RETEMP=1.0/TEMP;
  CARD%
MACRO HEADING LEAST SQUARES ANALYSIS DATASET HOD=0.0 AND F=.5%
MACRO SUBDATA HODOF5%
INDATA;
150.      .1558      .2250
150.      .1046      .2079
150.      .06363     .1974
145.      .06018     .08179
145.      .03102     .05578
145.      .06154     .06897
145.      .03442     .06602
140.      .01501     .02201
140.      .009688    .01621
135.      .003708    .003328
135.      .003138    .003862
135.      .003346    .003985
SOLVE;

```

PAGE 2

MACRO HEADING LEAST SQUARES ANALYSIS DATASET HOD=0.0 AND F=.7368%  
 MACRO SUBDATA HOD0F73%

INDATA;

150.	.1556	.2143
150.	.1046	.1974
150.	.06366	.1924
145.	.05999	.07713
145.	.03105	.05444
145.	.06144	.06632
145.	.03442	.06388
140.	.01500	.02135
140.	.009688	.01578
135.	.003707	.003304
135.	.003136	.003783
135.	.003346	.003934

SOLVE;

MACRO HEADING LEAST SQUARES ANALYSIS DATASET HOD=0.0 AND F=1.0%  
 MACRO SUBDATA HOD0F1%

INDATA;

150.	.1551	.2041
150.	.1045	.1878
150.	.06367	.1873
145.	.05977	.07297
145.	.03105	.05307
145.	.06131	.06386
145.	.03441	.06178
140.	.01498	.02068
140.	.009684	.01535
135.	.003708	.003280
135.	.003134	.003706
135.	.003346	.003882

SOLVE;

MACRO HEADING LEAST SQUARES ANALYSIS DATASET HOD=.25 AND F=.5%  
 MACRO SUBDATA HOD25F5%

INDATA;

150.	.1567	.3059
150.	.1047	.2927
150.	.06367	.2841
145.	.06024	.1050
145.	.03121	.08420
145.	.06168	.09504
145.	.03442	.08425
140.	.01501	.02809
140.	.009691	.02076
135.	.003706	.004098
135.	.003134	.004810
135.	.003344	.004823

SOLVE;

MACRO HEADING LEAST SQUARES ANALYSIS DATASET HOD=.25 AND F=.7368%  
 MACRO SUBDATA HOD25F73%

INDATA;

150.	.1559	.2760
150.	.1046	.2617
150.	.06367	.2674

VITA <sup>2</sup>

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Candidate for the Degree of  
Doctor of Philosophy

Thesis: SPECTROSCOPICALLY EVALUATED RATES AND ENERGIES FOR PROTON  
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