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# ANALYZING POWER MEASUREMENTS FOR THE EXCITATION OF STATES IN ${ }^{28}$ Si and ${ }^{24} \mathrm{Mg}$ BY INELASTIC SCATTERING OF POLARIZED PROTONS 

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Analyzing powers $A_{y}(\theta)$ for the excitation of states in ${ }^{28} \mathrm{Si}$ and ${ }^{24} \mathrm{Mg}$ with excitation energies up to 16 MeV have been measured with a $135-\mathrm{MeV}$ polarized proton beam. The scattered protons were detected with the QDDM magnetic spectrograph at angles between $25^{\circ}$ and $65^{\circ}$ with an overall resolution of about 70 keV . Results for the $6^{-}, \mathrm{T}=1(14.35 \mathrm{MeV}), 6^{-}, \mathrm{T}=0$ (11.58 $\mathrm{MeV})$, and $5^{-}, \mathrm{T}=0(9.70 \mathrm{MeV})$ states in ${ }^{28} \mathrm{Si}^{1}$, whose predominant configurations are all $\left(d_{5 / 2}\right)^{-1}\left(f_{7 / 2}\right)$, are shown in Fig. 1, where they are compared with the results of DWIA calculations using the t-matrix effective interaction derived by Love from the free nucleon-nucleon scattering data. ${ }^{2}$ The cross section for the $6^{-}$, $T=1$, state is due mainly to the tensor direct term in the interaction, while that for the $6^{-}$, $\mathrm{T}=0$, state $i s$ due mainly to tensor and spin-orbit exchange terms, and that for the $5^{-}, \mathrm{T}=0$, state is due mainly to spin-orbit and central interaction terms. The $A_{y}(\theta)$ results for the $6^{-}$states are sensitive to interference both between the central and spin-orbit parts and between the spin-orbit and tensor parts of


Figure 1. Analyzing powers, $A_{y}(\theta)$, for the $135-\mathrm{MeV}$ ( $\vec{p}, p^{\prime}$ ) excitation of ( $a$ ) the $6^{-}, T=0$, state at 11.58 MeV , (b) the $6^{-}, T=1$, state at 14.35 MeV , and (c) the $5^{-}, T=0$, state at 9.70 MeV . The experimental data are compared with results of DWIA calculations using the Love t-matrix.
the $t$-matrix. For the $5^{-}$state, on the other hand, the calculated result for $A_{y}$ is predominantly sensitive to central-spin-orbit interference, and the change of sign of $A_{y}$ near $40^{\circ}$ is rather well correlated with the change of sign in the central part of the $t$-matrix at the corresponding momentum transfer, as given by the Love interaction.

The differential cross sections and analyzing powers for the lower-lying states of ${ }^{28} \mathrm{Si}$ and ${ }^{24} \mathrm{Mg}$ have also been measured. In ${ }^{24} \mathrm{Mg}$, for example, the excitation of states in the $K=0$ and $K=2$ bands can be compared to calculations using the Chalk River pro-jected-Hartree-Fock wave functions. The remarkable agreement between theory and experiment for inelastic
electron scattering ${ }^{3}$ allows a detailed study of the proton scattering mechanism. The totally anomalous shape of the electromagnetic form factor for the $4_{1}{ }^{+}(\mathrm{K}=0)$ state in ${ }^{24} \mathrm{Mg}$ is well reproduced in the differential cross section for ( $p, p^{\prime}$ ).

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> SPIN-ORBIT EFFECTS IN THE EXCITATION OF PROTON AND NEUTRON STATES IN THE (p,p') REACTION AT $160 \mathrm{MeV}, 120 \mathrm{MeV}$, AND 95 MeV $\begin{gathered}\text { Alan Scott, F. Todd Baker, M.A. Grimm, J.H. Johnson, V. Penumetcha, } \\ \text { R.C. Styles, J.A. Mowrey, and W.G. Love } \\ \text { University of Georgia, Athens, Georgia } 30602\end{gathered}$ W.P. Jones and J.D. Wiggins

Large differences in the shapes of measured differential cross sections were found earlier ${ }^{1}$ for the excitations of the $4_{1}^{+}$proton state in ${ }^{90} \mathrm{Zr}$ and the $4_{1}^{+}$ neutron state in ${ }^{92} \mathrm{Zr}$. To obtain a good fit to the data for this proton state in ${ }^{90} \mathrm{Zr}$ with purely collective calculations, an enhanced spin-orbit contribution $\left(B_{4}^{S O} / B_{4}=1.25\right)$ was required, but no satisfactory fits were found for this neutron state in ${ }^{92} \mathbf{Z r}$. Collective fits to the data for the $2_{1}^{+}, 4_{1}^{+}, 6_{1}^{+}, 8_{1}^{+}$proton states in ${ }^{90} \mathrm{Zr}$ showed ${ }^{2}$ the increasing dominance of the spinorbit contribution as the multipolarity increased. Recent calculations show the cross sections for the $2_{1}^{+}$, $4_{1}^{+}, 6_{1}^{+}$, and $8_{1}^{+}$states in ${ }^{90} \mathrm{Zr}$ to be underpredicted in the DWIA by factors of $30,10,3$, and 2 respectively when only the $\left(g_{g / 2}\right)^{2}$ valence terms are included for
the central, spin-orbit and tensor amplitudes, ${ }^{3}$ suggesting the need for core polarization amplitudes similar in magnitude to those required at lower energies. ${ }^{4,5}$ These DWIA calculations showed the relative importance of the spin-orbit part of the t-matrix increased in this sequence as the multipolarity increased. ${ }^{3}$ The dominance of the spin-orbit contributions for the $8_{1}^{+}$ state in ${ }^{90} \mathrm{Zr}$ is shown in the DWIA and collective calculations of Figures $1(a)$ and $1(b)$ respectively.

Large spin-orbit effects at this energy ( $E_{p}=160 \mathrm{MeV}$ ) clearly suggested the need for ( $p, p^{\prime}$ ) asymmetry measurements. Analyzing power data have been obtained at 14 angles from $16^{\circ}$ to $44^{\circ}$ for the $2_{1}^{+}, 4_{1}^{+}$, $5_{1}^{-}, 3_{1}^{-}$, and $2_{2}^{+}$states in ${ }^{90} \mathrm{Zr}$ and for the $2_{1}^{+}, 4_{1}^{+}, 3_{1}^{-}$, $2_{2}^{+}, 2_{3}^{+}$, and $5_{1}^{-}$states in ${ }^{92} \mathrm{Zr}$, and at 8 angles from $26^{\circ}$

