

# On the ratio $T_2/T_1$ for non-Ohmic spectral densities

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There has been some recent interest<sup>1-8</sup> in the ratio  $T_2/T_1$ , where  $T_2$  and  $T_1$  are the phase and population relaxation times, respectively, for a two-level system (TLS). Budimir and Skinner<sup>1</sup> showed that when the TLS is strongly coupled to a stochastic bath, under certain circumstances one finds the unusual result that  $T_2 > 2T_1$ . The validity of this result was later extended to finite temperatures by considering a TLS coupled to a quantum-mechanical harmonic bath.<sup>3,9</sup> In this calculation the spectral density of the bath was Ohmic, that is, it was proportional to frequency in the low-frequency limit. Given that for other types of spin-boson Hamiltonians (e.g., the tunneling problem<sup>10,11</sup> and the pure dephasing problem<sup>12</sup>) Ohmic and non-Ohmic spectral densities give qualitatively different results, it is natural to wonder whether the possibility that  $T_2 > 2T_1$  is specific to the Ohmic model, or is, in fact, more generally valid.

To address this question, in this Note we consider the super-Ohmic spectral density used recently by Suárez *et al.*<sup>13</sup>

$$\Gamma_1(\omega) = A(\omega/\omega_c)^3 e^{-\omega/\omega_c} \quad (\omega > 0), \quad (1)$$

which has the same low-frequency behavior as the Debye model together with the standard deformation potential approximation. (Herein we follow exactly the definitions and notation of Ref. 3.) For the "complex" coupling model [ $\Gamma_2(\omega) = 0$ ], it was found that<sup>3</sup>

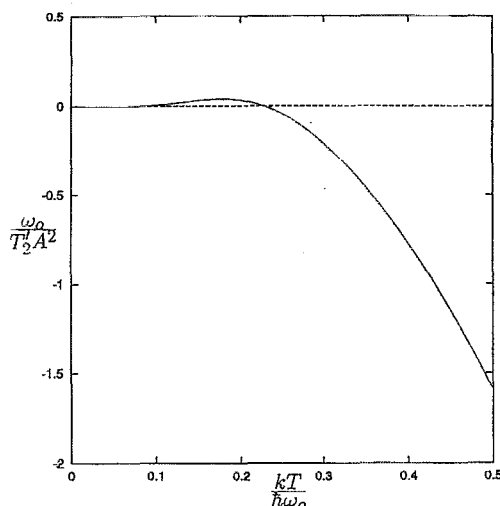


FIG. 1.  $\omega_c/T_2^2 A^2$  vs  $kT/\hbar\omega_0$  for  $\omega_c/\omega_0 = 0.4$ .

$$\frac{1}{T_2'} = \frac{4\delta^4}{\pi} P \int_{-\infty}^{\infty} \frac{d\omega \omega}{\omega^2 - \omega_0^2} \times \frac{\partial}{\partial \omega} \{ \Gamma_1(|\omega|)^2 n(|\omega|) [n(|\omega|) + 1] \} + O(\delta^6), \quad (2)$$

where  $1/T_2' \equiv 1/T_2 - 1/2T_1$ ,  $\delta$  is the dimensionless coupling constant,  $P$  is the Cauchy principal value,  $\omega_0$  is the TLS frequency, and  $n(\omega) = (e^{\hbar\omega/kT} - 1)^{-1}$ . In the high-temperature limit we can write (setting  $\delta = 1$ )  $1/T_2' = (AkT/\hbar)^2 \omega_0^{-3} f(\omega_c/\omega_0)$ , and the dimensionless function  $f$  can be evaluated numerically. We find that for  $0.281 < \omega_c/\omega_0 < 1.006$ ,  $f(\omega_c/\omega_0) < 0$ , which means that  $1/T_2' < 0$ , and that  $T_2 > 2T_1$ . Choosing now the value of  $\omega_c/\omega_0 = 0.4$ , we can calculate the temperature dependence of  $1/T_2'$ , which is shown in Fig. 1. We find that for  $kT/\hbar\omega_0 > 0.228$ ,  $T_2 > 2T_1$ .

Thus, we see that, similar to the Ohmic case, for a reasonably wide range of parameters  $T_2 > 2T_1$ . This shows that the breakdown of the usual inequality  $T_2 \leq 2T_1$  is not simply a peculiarity of the Ohmic model, but is indeed a more general result.

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