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The Effect of Mo Doping on The Charge Separation Dynamics and Photocurrent Performance of BiVO₄ Photoanodes

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Abstract

Doping with electron-rich elements in BiVO₄ photoanodes has been demonstrated as a desirable approach for improving their carrier mobility and charge separation efficiency. However, the effect of doping and dopant concentration on the carrier dynamics and photoelectrochemical performance remains unclear. In this work, we examined the effects of Mo doping on the charge separation dynamics and photocurrent performance in BiVO₄ photoanodes. We show that the photocurrent of BiVO₄ photoanodes increases with increasing concentration of the Mo dopant, which can be attributed to both the improved carrier mobility resulting from

increased electron density and charge separation efficiency due to the diminishing of trap states upon Mo doping. The effect of doping on the electronic structure, carrier dynamics and photocurrent performance of BiVO₄ photoanodes resulting from W and Mo dopants was also compared and discussed in this study. The knowledge gained from this work will provide important insights into the optimization of the carrier mobility and charge separation efficiency of BiVO₄ photoanodes by controlling the dopants and their concentrations.

Introduction

The development of efficient photoelectrode materials for the production of carbon-neutral and sustainable fuel sources by solar-driven water splitting is an urgent need to meet long-term global energy demands.¹⁴⁵ Despite the tremendous progress achieved towards the development of such photoelectrodes, it still remains a challenging task to develop a material that can meet the multiple resource intensive requirements of the solar-driven photocatalytic reactions.⁶⁴⁹ This particular challenge results from the oxidative half reaction, *i.e.* water oxidation to form protons and oxygen (OER), which requires a high overpotential and is considered as the bottleneck for artificial photosynthetic water splitting. Extensive efforts have thus been devoted to developing efficient photoanode materials for OER in the past decade.¹⁰⁻¹⁷

BiVO₄, owing to its moderate band gap (2.4 eV),¹³¹² capability to absorb visible light,²⁰²¹ high stability in aqueous solution,¹²²² as well as its significantly more positive valence band edge than the OER potential,²²²⁴ has become the most promising photoanode material. Regardless of its multifold benefits as a photoanode material, the performance of bare BiVO₄ for OER is not impressive. The key factors that lead to low OER efficiency have been ascribed to its poor carrier mobility, low charge separation efficiency, and weak water oxidation kinetics at the electrode surface.²⁵⁻²⁷ Doping with electron-rich elements has been shown as an attractive approach to improve the overall OER performance. For example, several studies have demonstrated improved electron mobility and photoelectrochemical (PEC) performance by W- and Mo-doping, which has been attributed to the increased electron density.¹³²²⁻³¹ In contrast, other studies suggested that doping can introduce trap states and serve as recombination centres that enhance the electron–hole recombination, which may decompensate the improved carrier mobility.²²²³³ Our recent studies on W-doped BiVO₄ photoanodes, however, suggested that the increased photocurrent upon W doping is largely attributed to the significant reduction of hole trap states, which inhibit the electron–hole pair recombination, while poor carrier mobility still remains as a limiting factor.³⁴ These previous studies suggest that doping may exert multiple influences on the electronic properties and charge separation efficiency despite its promise for overall OER improvement.

In order to gain a deep understanding of the doping mechanism in BiVO₄ photoanodes, it is necessary to examine the effect of doping on the carrier dynamics and OER performance using different dopants other than W. In this work, we investigate the effect of Mo doping on the morphology, carrier dynamics, and photocurrent performance of BiVO₄ photoanodes, as well as the dependence on Mo concentration. The results from these fundamental studies were also evaluated in comparison with our previous findings on W-doped BiVO₄ photoanodes.

Experimental

Materials

Bismuth(III) nitrate (Bi(NO₃)₃·5H₂O, Ward's Science, Rochester, NY), vanadyl(IV) acetylacetonate (C₁₀H₁₄O₅V, Acros Organics, Geel, Belgium), ammonium molybdate tetrahydrate ((NH₄)₆Mo₇O₂₄·4H₂O, Electron Microscopy Science, Hatfield, PA), acetic acid (\geq 99.7%, Sigma Aldrich, Milwaukee, WI), and acetylacetone (\geq 99.0%, TCI, Portland, OR) were used for film synthesis. Sodium sulfate (Na₂SO₄, Ward's Science) and highly-polished DI H₂O showing >16 Megaohm cm resistivity were used for electrolyte solution. FTO glass (2 mm thick, Solaronix, Aubonne, Switzerland) was cut to 1.25 × 2.50 cm dimensions for electrode fabrication. Nitric acid (Sigma Aldrich) and hydrogen peroxide (30%, J. T. Baker, Center Valley, PA) were used to make the piranha glass-etching solution.

Synthesis of Mo/BiVO₄

BiVO₄ films were prepared as in our previous work by a solution-based method.²⁶³² The Bi and V precursors (1 : 1 molar ratio) were dissolved in acetic acid (4 equivalents) and acetylacetone (1 equivalent), respectively, by sonication. Then, the two solutions were mixed and sonicated to make a 0.05 M BiVO₄ stock solution with respect to Bi and V. For undoped films, the stock solution was diluted to 0.04 M. For Mo-doped films, a 0.04 mM solution of Mo precursor in acetic acid was prepared and then added to the BiVO₄ stock solution in varying amounts (to obtain 0.2–1.8 mol% with respect to BiVO₄) along with any necessary amount of acetic acid to bring the final concentration of Mo/BiVO₄ to 0.04 M. Films were prepared by drop-coating onto a piranha-etched glass substrate or cleaned FTO. The films were dried in air for 30 minutes and then calcined at 450 °C for 90 minutes in air. After allowing to cool to room temperature slowly, films were then utilized for characterization.

Characterization

Steady-state UV-visible measurements were performed with an HP Agilent 8453 spectrometer. Scanning electron microscopy was performed with a JEOL JSM-6510LV operating in the secondary electron mode. X-ray diffraction (XRD) was performed using a Rigaku Miniflex II XRD diffractometer (Cu Kα). XRD samples were prepared by removing samples from the glass substrate and grinding in methanol for 10 minutes, followed by drying in air.

Transient absorption spectroscopy (TA)

The TA spectrometer is based on a regenerative amplified Ti–Sapphire laser system (Solstice, 800 nm, <100 fs FWHM, 3.5 mJ per pulse, 1 kHz repetition rate). The tunable 235–1150 nm pump is generated in TOPAS and the UV-visible probe is generated in a translated CaF₂ window by white light generation. 315 nm pump pulse (0.35μ J) was used for sample excitation. A Helios ultrafast spectrometer (Ultrafast Systems, LLC) was used for femtosecond TA measurements. The film sample was continuously translated to avoid sample degradation.

Photoelectrochemical cell (PEC) measurements

A custom three-electrode cell with a quartz window for electrode illumination was used for PEC measurements. A Ag/AgCl (3 M NaCl) reference electrode and a Pt wire counter electrode were used with a doped or un-doped BiVO₄– FTO working electrode in 0.1 M Na₂SO₄ (pH 7) electrolyte. A 300 W xenon lamp with an IR-filter (one foot long water

filter and an IR filter) and a UV-filter (>400 nm long-pass filter) output was used to illuminate the samples. The power density on the sample is 310 mW cm⁻². The current density–voltage (J-V) curve was recorded at a 25 mV s⁻¹ scan rate. The working electrode was illuminated from either front or back to observe photocurrent differences due to carrier transport.

Results and discussion

Although Mo-doped BiVO₄ films with Mo concentrations varying from 0.2 to 5% were all synthesized, this work will focus only on the films with 0.2%, 1.0%, and 1.8% Mo concentrations. This is because the films with Mo concentration above 2% show severe surface heterogeneity (Fig. S1a–c, ESI⁺), preventing our carrier dynamics studies using transient absorption spectroscopy, which is one of the core studies in this work.

The morphologies of un-doped and Mo-doped BiVO₄ films with Mo concentrations varying from 0.2%, 1.0%, to 1.8% were examined by SEM. As shown in Fig. 1, a porous nanostructure was present in all BiVO₄ films, while a rougher, more disordered structure was observed in the doped films with increasing Mo concentrations. The thickness of these films was controlled to be $\sim 2.8 \,\mu$ m (Fig. S1d, ESI⁺). Fig. 2a compares the XRD patterns of undoped and Mo-doped BiVO₄ films with different concentrations of Mo. With increasing Mo concentration, the splitting of the diffraction peaks at 35°, 46°, and 58° becomes less prominent. These results are consistent with previous literature data,³⁴³⁸ indicating that the structure of Mo/BiVO₄ films slightly changes from monoclinic to a mixture of monoclinic and tetragonal phases. Previous studies have shown that the bandgap energies for BiVO₄films with monoclinic to tetragonal phases are 2.4 eV and 2.9 eV, respectively.³⁴⁹ If Mo doping causes significant phase transfer from monoclinic to tetragonal structures, one would expect a peak shift in the UV-visible spectra of Mo-doped films compared to that of un-doped BiVO₄ films (Fig. 2b). These results, together with the less diffuse XRD patterns³⁴⁰ of Mo-doped films than those of the un-doped BiVO₄ film, suggest that Mo-doped BiVO₄ thin films mostly retain a monoclinic-scheelite structure.



Fig. 1 SEM images of the un-doped $BiVO_4$ film (a) and $Mo/BiVO_4$ films with 0.2% (b), 1.0% (c), and 1.8% (d) Mo concentration.





The effect of Mo concentration on carrier dynamics was examined using femtosecond transient absorption spectroscopy (TA). Fig. 3a–c show the TA spectra of Mo/BiVO₄ films with 0.2%, 1.0%, and 1.8% Mo concentrations, respectively. Similar to the spectral features of the un-doped BiVO₄ film published previously,^{3437–49} all Mo/BiVO₄ samples exhibit four main features, including two absorption bands centered at 370 nm and 470 nm, a broad featureless absorption >550 nm, and a negative feature that corresponds to the ground state (GS) bleach band. The three absorption bands exhibit a strong probe-wavelength-dependent feature (Fig. S2, ESI⁺), which is consistent with our previous report on the hole trapping process and can be assigned to the trapped hole absorption.³⁴Fig. 3d compares the initial TA spectra (at 0.5–1 ps delay time) of un-doped and Mo-doped BiVO₄ films. These spectra have been normalized at the maximum of the GS bleach band for better comparison. It is immediately noticeable that the initial intensities of these hole absorption bands in the Mo-doped films are significantly reduced, where more reduction occurs in the films with higher Mo concentrations. These results together suggest that the intrinsic hole traps in the un-doped BiVO₄ film are reduced due to Mo doping.



Fig. 3 Femtosecond TA spectra of 0.2% Mo/BiVO₄ (a), 1.0% Mo/BiVO₄ (b), and 1.8% Mo/BiVO₄ (c). (d) The comparison of TA spectra of un-doped and Mo-doped BiVO₄ films with different doping concentrations at 0.5-1 ps time delay. The comparison of TA kinetic traces at 472 nm (e) and 430 nm (f) for un-doped and Mo-doped BiVO₄ films with different Mo concentrations.

The reduction of hole traps in Mo-doped BiVO₄ films is further supported by the diminishing rising component at 472 nm which has been assigned to the hole trapping process previously.^{37,39,40} As shown in Fig. 3e, a clear rising component was observed in the kinetic trace of the un-doped BiVO₄ film. In contrast, this rising component becomes slower in 0.2% Mo/BiVO₄ and 1.0% Mo/BiVO₄ films, and eventually disappears in the 1.8% Mo/BiVO₄ film. These kinetic traces can be adequately fitted using a multi-exponential function. The kinetic traces for un-doped BiVO₄, 0.2% Mo/BiVO₄, and 1.0% Mo/BiVO₄ films can be fitted by a single rise component and three decay components, and the kinetic trace of 1.8% Mo/BiVO₄ film can be fitted by three decay components. The fitting parameters and equations are summarized in Table 1. The obtained time constants for the hole trapping process are 2.1 ps, 4.2 ps, and 5.3 ps for un-doped BiVO₄, 0.2% Mo/BiVO₄, and 1.0% Mo/BiVO₄ films, respectively, suggesting a retarded hole trapping process with increasing Mo concentration. The remaining three decay components can be assigned to the charge recombination process between the trapped holes with electrons in the conduction band or electron trap states.^{37–40}

Probe	Sample (%)	τ 1, ps (A1, %)	τ 2, ps (A2, %)	τ ₃, ps (A₃, %)	τ 4, ps	s (A₄, %)
The 43	80 nm kinet	ic traces in witl	$\frac{\text{Fig. 3f}}{\sum_{i} A_{i} e^{-i/\tau_{i}}}$	itted with $-e^{-t/\tau_r}$. r i	$\sum_{i} A_{i} e^{i}$ ndicat	. The 472 nm kinetic traces in <u>Fig. 3e</u> are fitted ees the rising component.
430 nm	0	1.6 (64)	23 (14.5)	874 (15.1)	≫5 ns	(6.4)
	0.2	2.2 (58.7)	37.7 (17.0)	1130 (12.7)		(11.5)
	1.0	2.2 (45.5)	35.3 (22.8)	1130 (18.9)		(12.8)
	1.8	2.2 (37.2)	49.1 (24.0)	1130 (24.2)		(13.8)
472 nm	0	'2.1 (100)	16.5 (32.5)	970 (44.4)	≫5 ns	(23.1)
	0.2	r4.2 (100)	21.1 (39.6)	964 (40.9)		(19.5)
	1.0	'5.3 (100)	15.6 (39.7)	794 (39.7)		(20.6)
	1.8		5.0 (57.6)	628 (26.8)		(15.6)

Table 1 Fitting parameters for TA kinetics of un-doped BiVO4 and Mo-doped BiVO4 films

In addition to hole trapping dynamics, the effect of Mo doping on the electron dynamics was examined by probing the GS bleach kinetics. Fig. 3f compares the GS bleach recovery kinetics (430 nm) for un-doped BiVO₄ and Mo-doped BiVO₄ films. These kinetic traces can all be fit by four-exponential-decay functions with their parameters listed in Table 1. The fastest time component (τ_1 in Table 1), likely resulting from an electron trapping process, increases from 1.6 ps in the un-doped BiVO₄ film to 2.2 ps in Mo-doped BiVO₄ films. The increase of time constant is accompanied by decreasing amplitude, suggesting the reduction of the electron trapping process. In addition, the overall GS recovery kinetics becomes slower in Mo-doped BiVO₄ films, following the order of 1.8% Mo/BiVO₄ < 1.0% Mo/BiVO₄ < 0.2% Mo/BiVO₄ < un-doped BiVO₄, suggesting that the overall electron—hole recombination is inhibited with the increase of Mo concentration. These results imply that Mo doping altered the nature of trap states in BiVO₄ films through diminishing the hole and electron trapping processes, which together elongate the electron—hole pair lifetime. Due to the presence of the long decay component (\gg 5 ns) which is beyond the TA time window, the electron—hole lifetime cannot be quantitatively determined from these fitting parameters. Instead, the electron—hole lifetime was estimated

from the half-recovery time of GS bleach, the time when the GS bleach recovers to 50%. The half-recovery lifetimes were 2.9, 6.0, 17.2, and 48.5 ps for un-doped BiVO₄, 0.2% Mo/BiVO₄, 1.0% Mo/BiVO₄, and 1.8% Mo/BiVO₄ films, respectively.

In addition to the effect of Mo doping on the carrier dynamics of BiVO₄ films, the dependence of carrier mobility in BiVO₄ films on Mo concentration was also examined using linear sweep voltammetry performed under visible light illumination. It has been shown that the difference in photocurrents between back- and front-illumination is related to the electron mobility because the back-illumination generates excitons close to the FTO backcontact whereas front-illumination generates excitons primarily at the surface where the electrons need to reach the back contact of the electrode. (Fig. 4a) therefore, the higher photocurrent from back-illumination than from front-illumination suggests that slow electron transport is one of the limiting factors for photocurrent generation. For the un-doped BiVO₄ sample (Fig. 4a), there is a significant difference in photocurrent when illuminated from the front compared to from the back during photocatalysis, where much higher photocurrent under back-side illumination was observed, suggesting poor electron mobility in the un-doped BiVO₄ sample. It is noticeable that this difference becomes smaller with increasing Mo doping level (Fig. 4b–d), suggesting that electron mobility improves with increasing Mo doping. Indeed, little difference was observed between frontand back-illumination for the 1.8% Mo/BiVO₄ film (Fig. 4d), indicating that electron mobility is no longer a limiting factor for OER performance with the addition of 1.8% Mo dopant. These findings clearly demonstrate that electron mobility in BiVO₄ photoanodes is significantly improved upon Mo doping.



Fig. 4 The comparison of photocurrent from back-side- and front-side-illumination for un-doped $BiVO_4$ (a), 0.2% Mo/BiVO_4 (b), 1.0% Mo/BiVO_4 (c), and 1.8% Mo/BiVO_4 (d) thin films.

With the understanding of the effect of Mo doping on the carrier dynamics and electron mobility of $BiVO_4$ photoanodes, we now turn to evaluate the correlation of these properties with OER performance. Fig. <u>5</u> compares the OER photocurrents of the un-doped $BiVO_4$ film and Mo-doped $BiVO_4$ films with different Mo concentrations under back-side illumination. Although not shown, there was negligible current observed for all samples without light illumination. In contrast, under light illumination, it is immediately apparent that Modoped photoanodes produce a higher photocurrent than the un-doped BiVO₄ photoanode, beginning to diverge at overpotentials as low as 0.3 V vs. Ag/AgCl. Despite the little difference in photocurrent between 0.2% and 1.0% Mo/BiVO₄, the overall trend for photocurrent shows an increase with increasing doping concentration, where the sample with the highest doping level (1.8% Mo/BiVO₄) shows the maximum photocurrent.





Previous studies have shown that the PEC performance of BiVO₄ photoanodes is strongly dependent on their crystal structure. Kudo and coworkers have found that BiVO₄ photoanodes with a monoclinic scheelite structure demonstrated much higher photocatalytic activity than those with a tetragonal structure, which has been mainly attributed to the enhanced light absorption of the former with a bandgap of 2.4 eV, while the latter have a bandgap of 2.9 eV.¹⁰²² According to these findings, one would expect that the change in the BiVO₄ structure from monoclinic scheelite to tetragonal due to Mo doping would reduce its photocatalytic activity. However, our work along with some previous studies²⁰²⁵ showed improved PEC performance regardless of the change in the crystal structure. Given that other factors in addition to the crystal structure can also influence the PEC performance, and doping only causes small changes in the crystal structure, we believe that the much higher photocurrent in the Mo-doped films is mainly ascribed to the improved electron mobility due to increasing electron density and increased carrier lifetime due to diminishing electron and hole trap states, which likely compensates the negative effect resulting from crystal structure change.

Very recently, we have reported the origin of improved charge separation and OER performance of BiVO₄ photoanodes to be W doping.²⁴ The results showed that the improved photocurrent of BiVO₄ is mainly attributed to the inhibited electron—hole recombination due to the reduction of trap states, while the carrier mobility has little contribution. In contrast, the current study on Mo-doping effect demonstrates that the significant improvement in photocurrent can be ascribed to both the increased carrier mobility and charge separation efficiency. In order to better understand the doping mechanism in BiVO₄ photoanodes and identify the appropriate dopants for optimum OER performance, it is necessary to compare the origin of the doping effect on the electronic structure, carrier dynamics and photocurrent performance of BiVO₄ photoanodes resulting from Mo and W dopants.

Fig. 6a compares the XRD patterns of 1.8% W/BiVO₄ and 1.8% Mo/BiVO₄ films. Slightly sharper diffraction peaks at 35°, 46°, and 58° were observed for the W/BiVO₄ film than for the Mo/BiVO₄ film, indicating that Mo/BiVO₄ more likely retains the monoclinic-scheelite phase than the W/BiVO₄ film. This finding was further confirmed by their UV-visible spectra, where the spectrum of the W/BiVO₄ film is found to be more blue-shifted compared to the spectrum of Mo/BiVO₄ (Fig. 6b). The possibility of changes in the crystal phase with doping necessitates the structural analysis by X-ray absorption spectroscopy, in which the local structure of individual atoms can be probed and determined. As shown in Fig. S3–S5 (ESI[±]), the local structure at both Bi and V centers, probed by X-ray absorption spectroscopy (XAS), shows a negligible change in both the oxidation state and the local geometry between W/BiVO₄ and Mo/BiVO₄, although their local structures are different from that of the un-doped BiVO₄ film.²⁴ These results imply that the change in the ground state electronic structure of BiVO₄ films than in W/BiVO₄ films.



Fig. 6 The comparison of XRD patterns (a), UV-visible absorption spectra (b), ground state bleach kinetics at 430 nm (c), and photocurrents (d) of 1.8% Mo/BiVO₄ and 1.8% W/BiVO₄ films.

In addition to their ground state electronic structure, the carrier dynamics of 1.8% W/BiVO₄ and 1.8% Mo/BiVO₄ films were also compared. As indicated by the noteworthy reduction of the initial amplitude of TA spectra of both doped BiVO₄ films compared to the un-doped one (Fig. S6, ESI⁺), doping caused a substantial reduction of hole traps in the doped films. However, the GS bleach recovery in 1.8% Mo/BiVO₄ appears to be faster than that in 1.8% W/BiVO₄ (Fig. 6c), suggesting enhanced electron–hole recombination in the former. This result conflicts with the photocurrent measurement, where the photocurrent generated by 1.8% Mo/BiVO₄ photoanode is much higher than that of 1.8% W/BiVO₄ photoanode (Fig. 6d). Given that Mo doping caused significantly increased electron mobility in 1.8% Mo/BiVO₄ photoanode, we therefore attribute the much higher

photocurrent in the former to the increased electron mobility upon Mo doping, which compensates its less efficient charge separation efficiency.

Conclusions

In conclusion, the effect of Mo doping on the carrier dynamics and OER performance of BiVO₄ photoanodes was investigated. Using transient absorption spectroscopy, we show that both electron and hole traps are reduced upon Mo doping, inhibiting the electron—hole recombination. The studies of photocurrent measurement using linear sweep voltammetry reveal an improved photocurrent with increasing Mo concentration, which is accompanied by the significantly enhanced electron mobility. Compared to W-doped BiVO₄ photoanodes, despite the enhanced electron—hole recombination in Mo/BiVO₄ photoanodes, the OER photocurrent is much higher in the latter, which can be attributed to the significantly improved electron mobility in Mo/BiVO₄ photoanodes, compensating the less efficient charge separation efficiency. These findings provide further understanding of the doping physics in BiVO₄ photoanodes and will provide important insights into the optimization of OER performance by identifying appropriate dopants and dopant concentrations.

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References

- 1. S. Y. Reece, J. A. Hamel, K. Sung, T. D. Jarvi, A. J. Esswein, J. J. H. Pijpers and D. G. Nocera, *Science*, 2011, **334**, 645–648.
- 2. M. Gratzel, Nature, 2001, 414, 338-344
- 3. M. G. Walter, E. L. Warren, J. R. McKone, S. W. Boettcher, Q. X. Mi, E. A. Santori and N. S. Lewis, *Chem. Rev.*, 2010, **110**, 6446–6473
- 4. W. J. Youngblood, S. H. A. Lee, Y. Kobayashi, E. A. Hernandez-Pagan, P. G. Hoertz, T. A. Moore, A. L. Moore, D. Gust and T. E. Mallouk, *J. Am. Chem. Soc.*, 2009, **131**, 926–927 .
- 5. D. Gust, T. A. Moore and A. L. Moore, Acc. Chem. Res., 2009, 42, 1890–1898
- 6. A. Fujishima and K. Honda, Nature, 1972, 238, 2
- 7. M. W. Kanan and D. G. Nocera, Science, 2008, **321**, 1072–1075.
- 8. F. D. Lin and S. W. Boettcher, Nat. Mater., 2014, 13, 81–86.
- 9. Y. J. Lin, S. Zhou, X. H. Liu, S. Sheehan and D. W. Wang, J. Am. Chem. Soc., 2009, 131, 2772–2773.
- 10. Y. Tian and T. Tatsuma, J. Am. Chem. Soc., 2005, 127, 7632–7637.
- 11. M. Ni, M. K. H. Leung, D. Y. C. Leung and K. Sumathy, Renewable Sustainable Energy Rev., 2007, 11, 401–425.
- 12. I. S. Cho, M. Logar, C. H. Lee, L. L. Cai, F. B. Prinz and X. L. Zheng, *Nano Lett.*, 2014, **14**, 24–31.
- 13. H. Dotan, K. Sivula, M. Gratzel, A. Rothschild and S. C. Warren, *Energy Environ. Sci.*, 2011, 4, 958–964.
- 14. M. Bar, L. Weinhardt, B. Marsen, B. Cole, N. Gaillard, E. Miller and C. Heske, Appl. Phys. Lett., 2010, 96, 032107
- 15. Q. X. Mi, Y. Ping, Y. Li, B. F. Cao, B. S. Brunschwig, P. G. Khalifah, G. A. Galli, H. B. Gray and N. S. Lewis, *J. Am. Chem. Soc.*, 2012, **134**, 18318–18324 .
- 16. J. W. Jang, C. Du, Y. Ye, Y. Lin, X. Yao, J. Thorne, E. Liu and G. McMahon, Nat. Commun., 2015, 6, 7447–7451
- 17. Y. Park, K. J. McDonald and K. S. Choi, Chem. Soc. Rev., 2013, 42, 2321–2337
- 18. A. Kudo, K. Omori and H. Kato, J. Am. Chem. Soc., 1999, 121, 11459–11467.
- 19. F. F. Abdi, L. H. Han, A. H. M. Smets, M. Zeman, B. Dam and R. van de Krol, Nat. Commun., 2013, 4, 2195.

- 20. K. Sayama, A. Nomura, T. Arai, T. Sugita, R. Abe, M. Yanagida, T. Oi, Y. Iwasaki, Y. Abe and H. Sugihara, *J. Phys. Chem. B*, 2006, **110**, 11352–11360.
- 21. H. M. Luo, A. H. Mueller, T. M. McCleskey, A. K. Burrell, E. Bauer and Q. X. Jia, *J. Phys. Chem. C*, 2008, **112**, 6099–6102.
- 22. A. Iwase and A. Kudo, J. Mater. Chem., 2010, 20, 7536–7542
- 23. A. Walsh, Y. Yan, M. N. Huda, M. M. Al-Jassim and S. H. Wei, Chem. Mater., 2009, 21, 547–551.
- 24. S. P. Berglund, D. W. Flaherty, N. T. Hahn, A. J. Bard and C. B. Mullins, J. Phys. Chem. C, 2011, 115, 3794–3802 .
- 25. H. Ye, J. Lee, J. S. Jang and A. J. Bard, J. Phys. Chem. C, 2010, 114, 13322–13328.
- 26. P. M. Rao, L. L. Cai, C. Liu, I. S. Cho, C. H. Lee, J. M. Weisse, P. D. Yang and X. L. Zheng, *Nano Lett.*, 2014, **14**, 1099–1105 .
- 27. F. F. Abdi, T. J. Savenije, M. M. May, B. Dam and R. van de Krol, J. Phys. Chem. Lett., 2013, 4, 2752–2757.
- W. J. Luo, Z. S. Yang, Z. S. Li, J. Y. Zhang, J. G. Liu, Z. Y. Zhao, Z. Q. Wang, S. C. Yan, T. Yu and Z. G. Zou, *Energy Environ. Sci.*, 2011, 4, 4046–4051.
- 29. H. S. Park, K. E. Kweon, H. Ye, E. Paek, G. S. Hwang and A. J. Bard, J. Phys. Chem. C, 2011, 115, 17870–17879.
- 30. S. P. Berglund, A. J. E. Rettie, S. Hoang and C. B. Mullins, Phys. Chem. Chem. Phys., 2012, 14, 7065–7075.
- 31. K. P. S. Parmar, H. J. Kang, A. Bist, P. Dua, J. S. Jang and J. S. Lee, ChemSusChem, 2012, 5, 1926–1934.
- 32. W. J. Luo, Z. S. Li, T. Yu and Z. G. Zou, J. Phys. Chem. C, 2012, 116, 5076-5081.
- 33. K. E. Karakitsou and X. E. Verykios, J. Phys. Chem., 1993, 97, 1184–1189.
- 34. B. Pattengale, J. Ludwig and J. Huang, J. Phys. Chem. C, 2016, 120, 1421–1427.
- A. J. E. Rettie, H. C. Lee, L. G. Marshall, J. F. Lin, C. Capan, J. Lindemuth, J. S. McCloy, J. S. Zhou, A. J. Bard and C. B. Mullins, *J. Am. Chem. Soc.*, 2013, **135**, 11389–11396.
- 36. S. M. Thalluri, S. Hernandez, S. Bensaid, G. Saracco and N. Russo, Appl. Catal., B, 2016, 180, 630–636.
- 37. N. Aiga, Q. X. Jia, K. Watanabe, A. Kudo, T. Sugimoto and Y. Matsumoto, *J. Phys. Chem. C*, 2013, **117**, 9881–9886.
- 38. Y. M. Ma, S. R. Pendlebury, A. Reynal, F. Le Formal and J. R. Durrant, Chem. Sci., 2014, 5, 2964–2973.
- 39. J. Ravensbergen, F. F. Abdi, J. H. van Santen, R. N. Frese, B. Dam, R. van de Krol and J. T. M. Kennis, *J. Phys. Chem. C*, 2014, **118**, 27793–27800.
- 40. I. Grigioni, K. G. Stamplecoskie, E. Selli and P. V. Kamat, J. Phys. Chem. C, 2015, 119, 20792–20800.
- 41. D. K. Zhong, S. Choi and D. R. Gamelin, J. Am. Chem. Soc., 2011, 133, 18370–18377.
- 42. F. F. Abdi and R. van de Krol, J. Phys. Chem. C, 2012, 116, 9398-9404.
- 43. J. A. Seabold and K. S. Choi, J. Am. Chem. Soc., 2012, 134, 2186–2192.