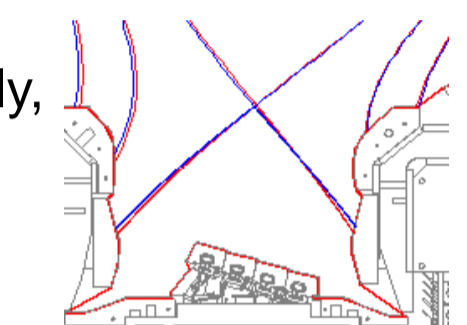
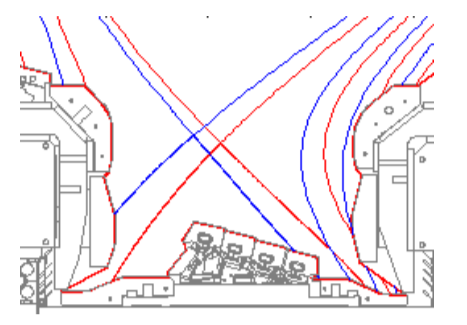


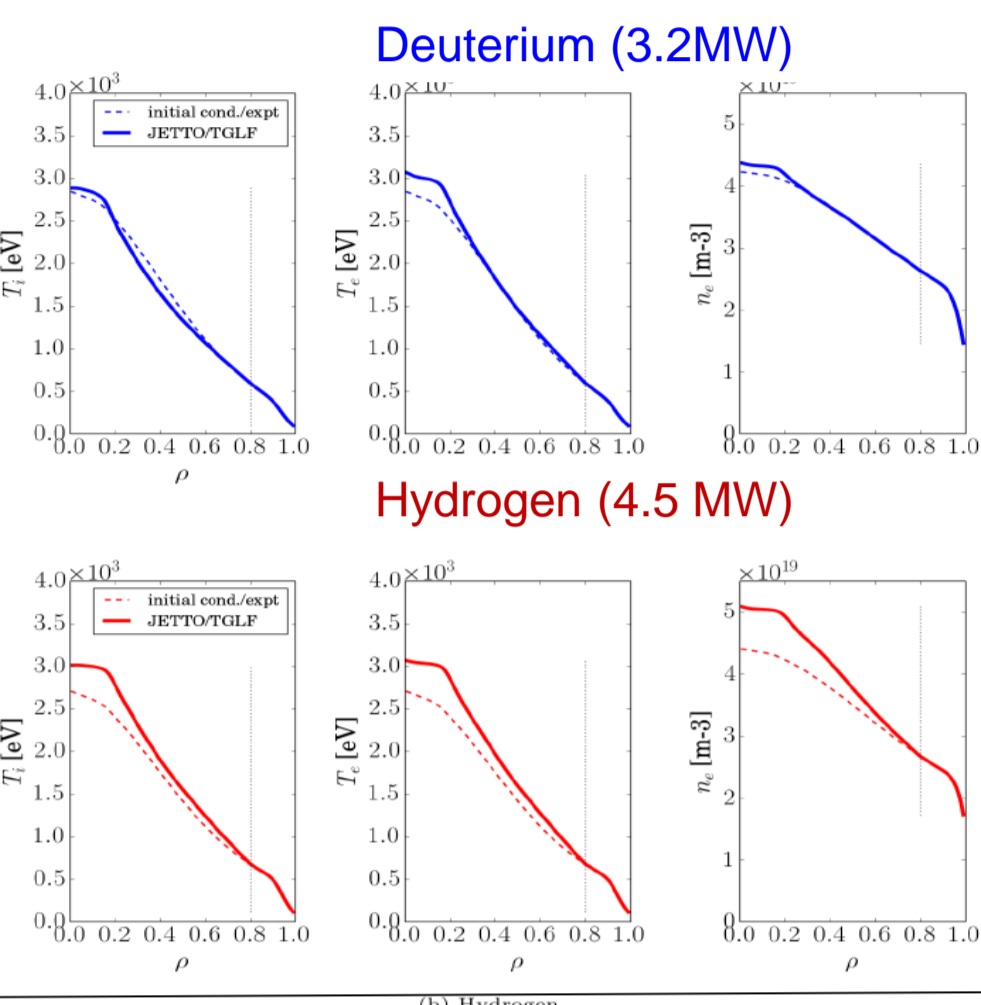
Datasets

- Deuterium & Hydrogen type I ELMy H-modes**, 171 samples (Maggi, PPCF 2018)
 - $B_T=1T$, $I_p=1MA$ ($q_{95} \approx 3$) and $B_T=1.7T$, $I_p=1.4MA$ ($q_{95} \approx 3.7$) (also 1.7T, 1.7MA in D only for dimensionless identity)
 - Mostly 'corner-corner (C/C)', configuration for best pumping and lowest P_{LH} power (some V/H too)
 - Gas scans and power scans:
 - Deuterium: $3.5MW \leq P_{NBI} \leq 17MW$, only NBI
 - Hydrogen: $5 MW \leq P_{NBI} \leq 10MW$, $0 \leq P_{ICRH} \leq 6.5MW$
 - NBI+ICRH required to achieve type I ELMy H-modes at 1.7T/1.4MA
- Deuterium & Hydrogen L-modes**, 20 samples
 - $B_T=2.9 T$, $I_p=2.5 MA$, $\langle n_e \rangle \geq 3.1 \times 10^{19} m^{-3}$ NBI power scan only, 20 samples, only NBI power scans $1.5MW \leq P_{NBI} \leq 9.5MW$
 - Divertor strike points on vertical tiles for highest P_{LH}
- Diagnostics**
 - Isotope ratio from high resolution Balmer- α spectroscopy and RGA
 - T_e and n_e from Thomson scattering (HRTS as used here and LIDAR)
 - T_i , ω_i and from CXRS (mostly Ne X line, only for 86 samples in H-mode)
 - n_i from n_e and Z_{eff} , assuming $Z_{eff}=1$ is due to Be only



TGLF modelling of H & D L-modes with same W_{th}

- Predictive JETTO-TGLF (Staebler 2007, Romanelli 2014) modelling finds
 - ~identical confinement in D and H due to temperature profile shapes (stiffness)
 - Largest D/H differences in particle channel (not stiff)
 - Weak "anti-GB" scaling in experiment, missing in model
 - Including ExB effects does not improve confinement D relative to H in this case



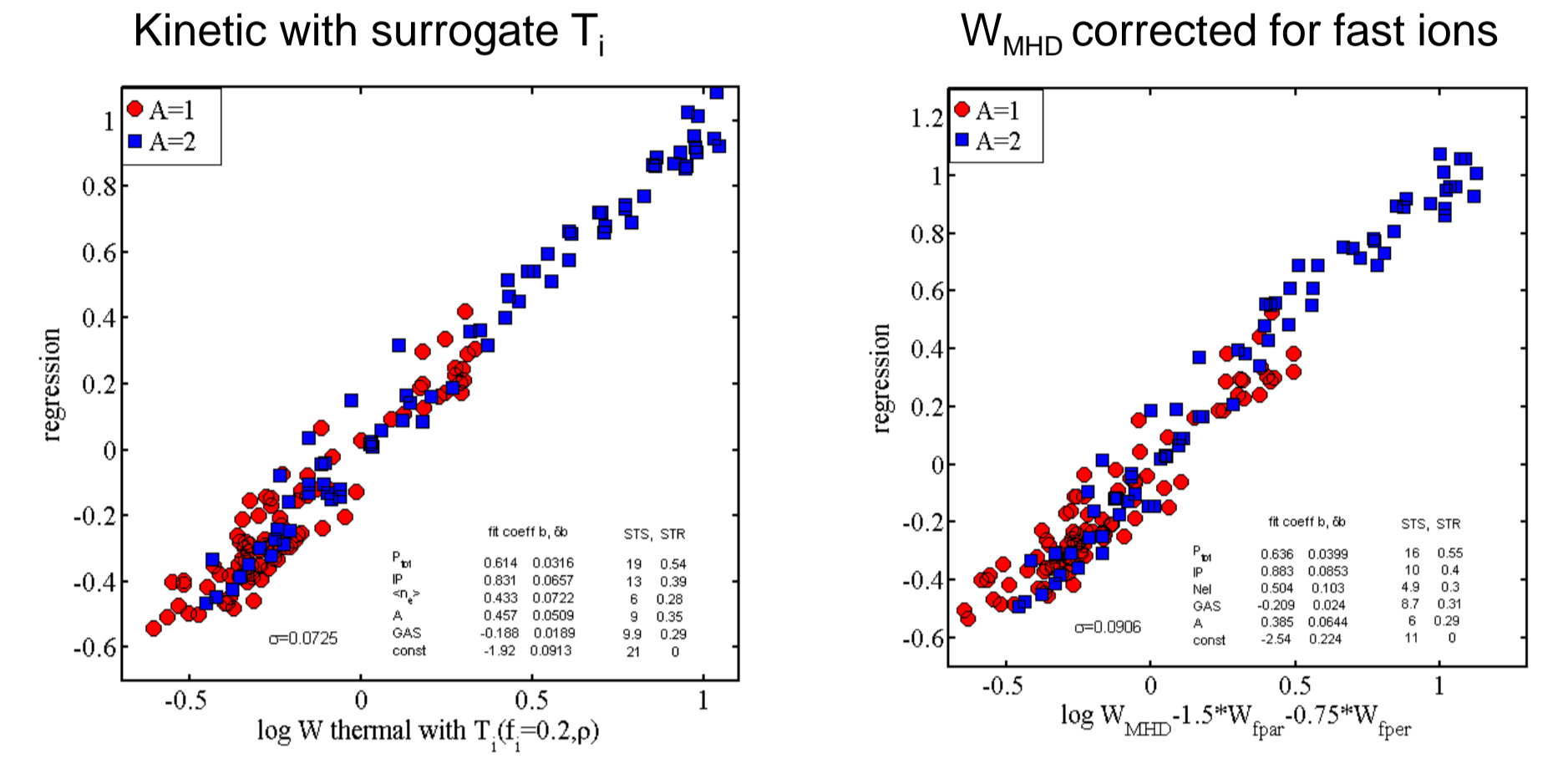
- Discussion:**
- Expectation of GB mass scaling of global confinement is naïve!
 - Assumes ion driven transport (ITG), neglects subdominant TEM, collisions
 - Assumes that fluxes are proportional to gradients
 - No complex non-linear / zonal flow / multi-scale effects
 - Assumes all else equal when isotope is changed (e.g. ExB shear in GB units, fast ion density, edge conditions...)
 - Assumes local property maps to a global confinement.
 - Boundary conditions can overrule local scaling with stiff profiles as recognised by Bateman et al, 1999

Type I ELMy H-mode: Global thermal energy confinement depends strongly on ion mass

Kinetic: $W_{th} \propto A^{0.46 \pm 0.06} p^{0.61 \pm 0.03} I_p^{0.83 \pm 0.07} \langle n_e \rangle^{0.43 \pm 0.08} G^{-0.19 \pm 0.02}$

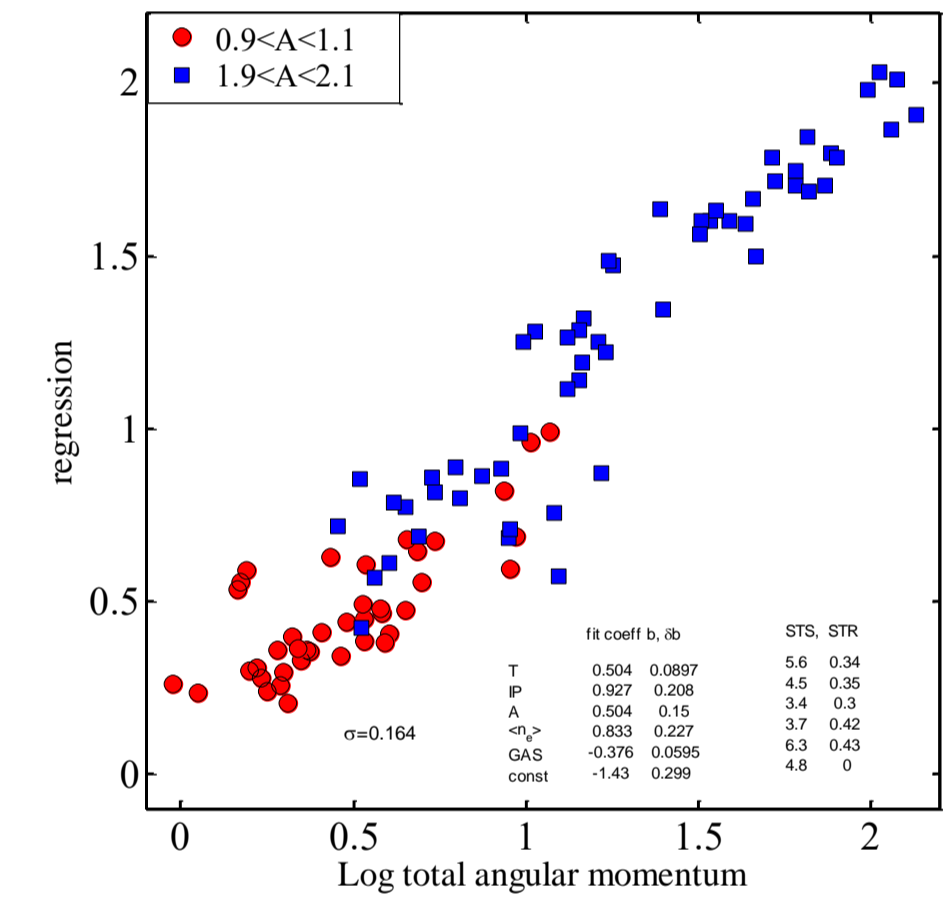
EFIT-based: $W_{th-EFIT} \propto A^{0.39 \pm 0.06} p^{0.64 \pm 0.04} I_p^{0.88 \pm 0.09} \langle n_e \rangle^{0.50 \pm 0.11} G^{-0.21 \pm 0.025}$

G is gas injection rate; is as statistically relevant and more significant than $\langle n_e \rangle$



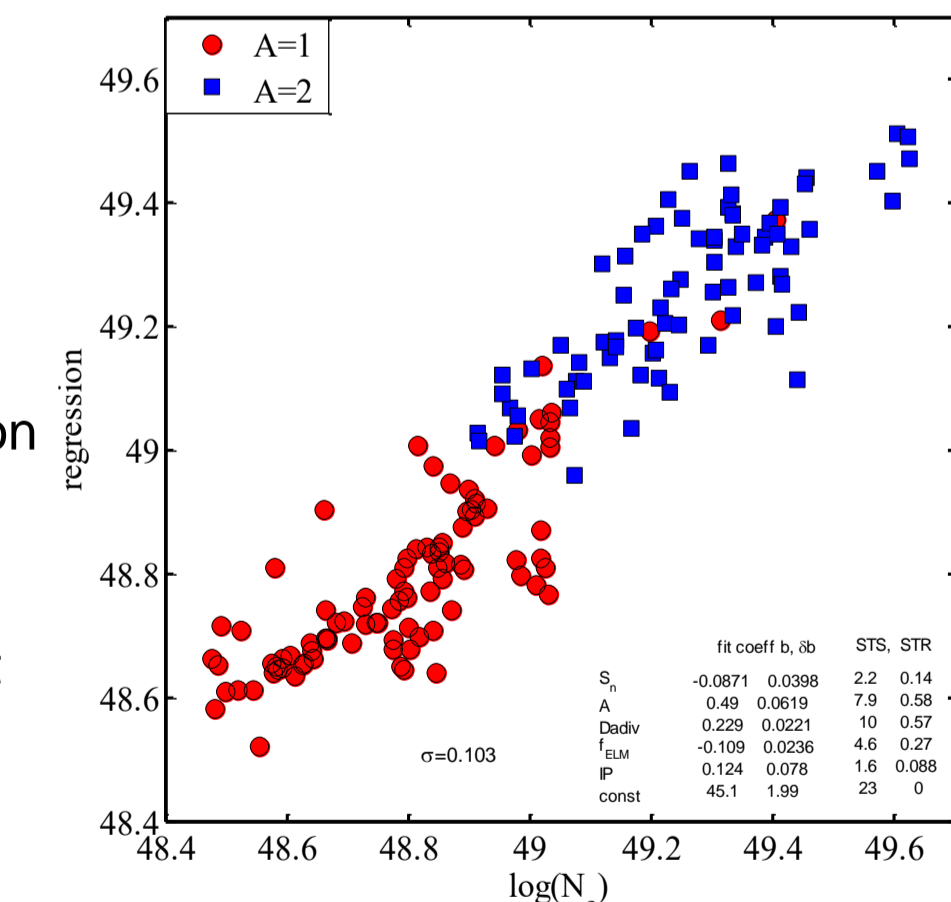
H-mode: Momentum dependence strong

- τ_E/τ_E in range 0.8-1.6
- Dependencies for angular momentum L are similar to total thermal energy
- $L \propto A^{0.5} T^{0.5} I_p^{0.93} \langle n_e \rangle^{0.83} G^{-0.38}$
- T is total NBI torque
- Significant, because momentum carried by ions only, no issue with equipartition with electrons
- Consistent with overall transport being dominated by ion channel if Prandtl number ≈ 1 , consistent with core GK modelling finding that ITG turbulence dominates in core



H-mode: particle confinement dependence strong

- $N_e = \int n_e dV$ total electron content
- Divertor Balmer-alpha Γ_{div} and NBI source S_{NBI} as proxies for source
- Regression: $N_e \propto A^{0.49} \Gamma_{div}^{0.23} S_{NBI}^{-0.08} I_p^{0.12} f_{ELM}^{-0.11}$ or
- Strongest dependencies of N_e are on ion mass and divertor source
- Beam fuelling contribution S_N to N_e insignificant (STS-2) and irrelevant (STR-0.1, negative). Weak I_p dependence intriguing...
- Weak but significant negative dependence on ELM frequency f_{ELM}



Energy, momentum and particle confinement have isotope dependency $\sim A^{0.4}$ to $A^{0.5}$ in type I ELMy H-mode! Coincidence or clue?

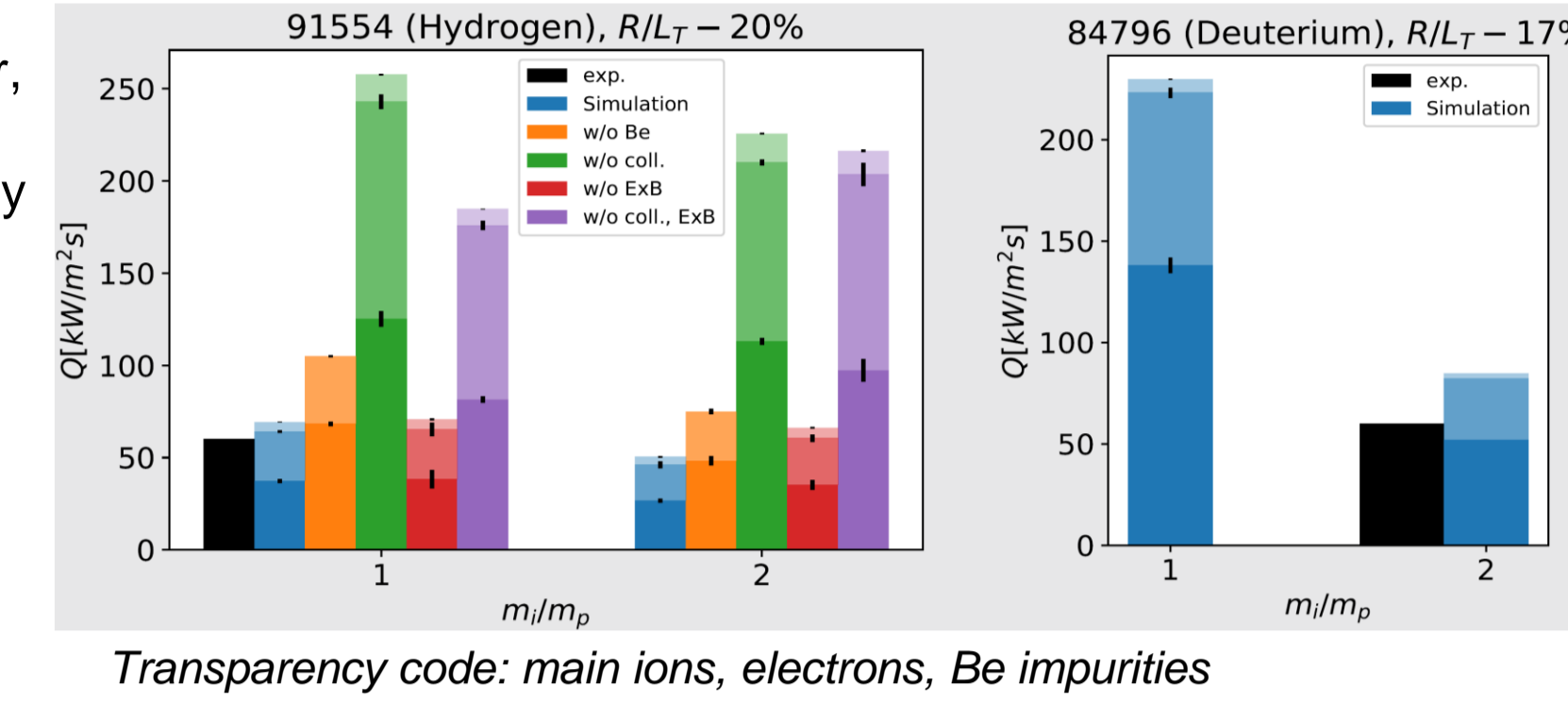
- Particle density n_e lower in hydrogen entails lower energy ($\propto n_e T_e$) confinement because absolute temperatures are similar in shape (stiff) and in absolute value

\Rightarrow global confinement scaling with mass is BAKED INTO PEDESTAL

- This conclusion was already drawn by Bateman et al, 1999

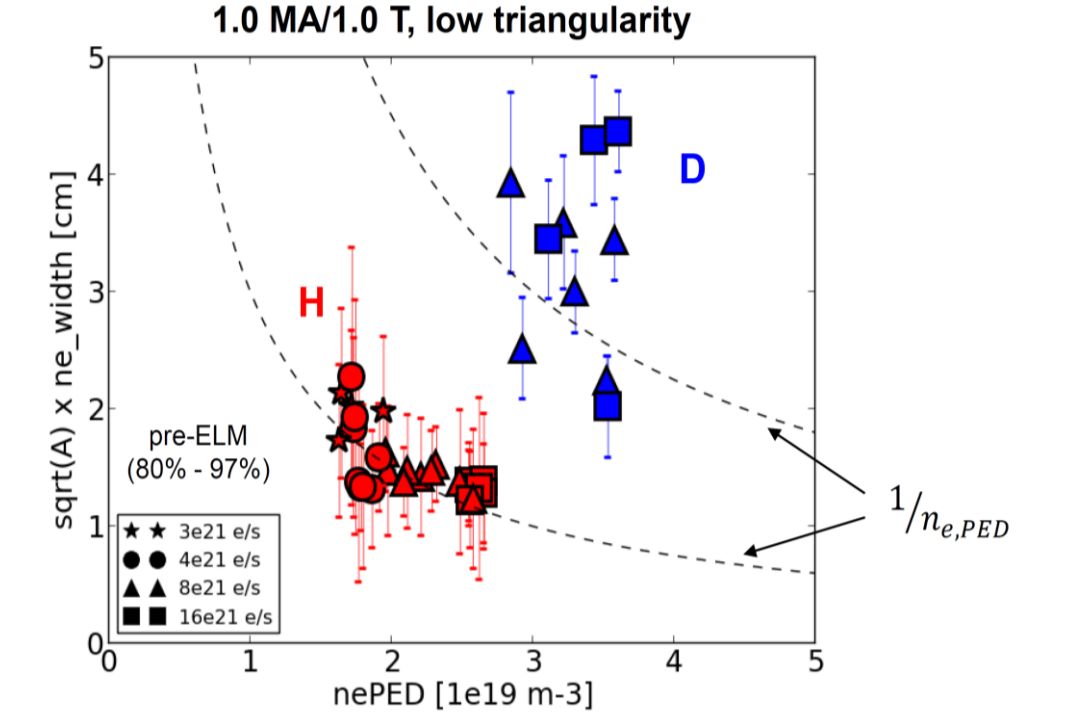
H-mode: Nonlinear GENE simulations reverse GyroBohm scaling

- Pair (H&D) with $P_{aux}=10MW$ simulated, non-linear, flux-tube, $\rho=0.5$, assuming $A=1$ & 2
- Absolute heat fluxes reproduced if ∇T_e reduced by ~20%, provided
 - collisions are included
 - dilution by Be impurities included
- Strong overprediction if collisions are neglected
- Note this is local - results may be different in future global simulations with imposed boundary conditions (M. Oberparleiter, 2019)



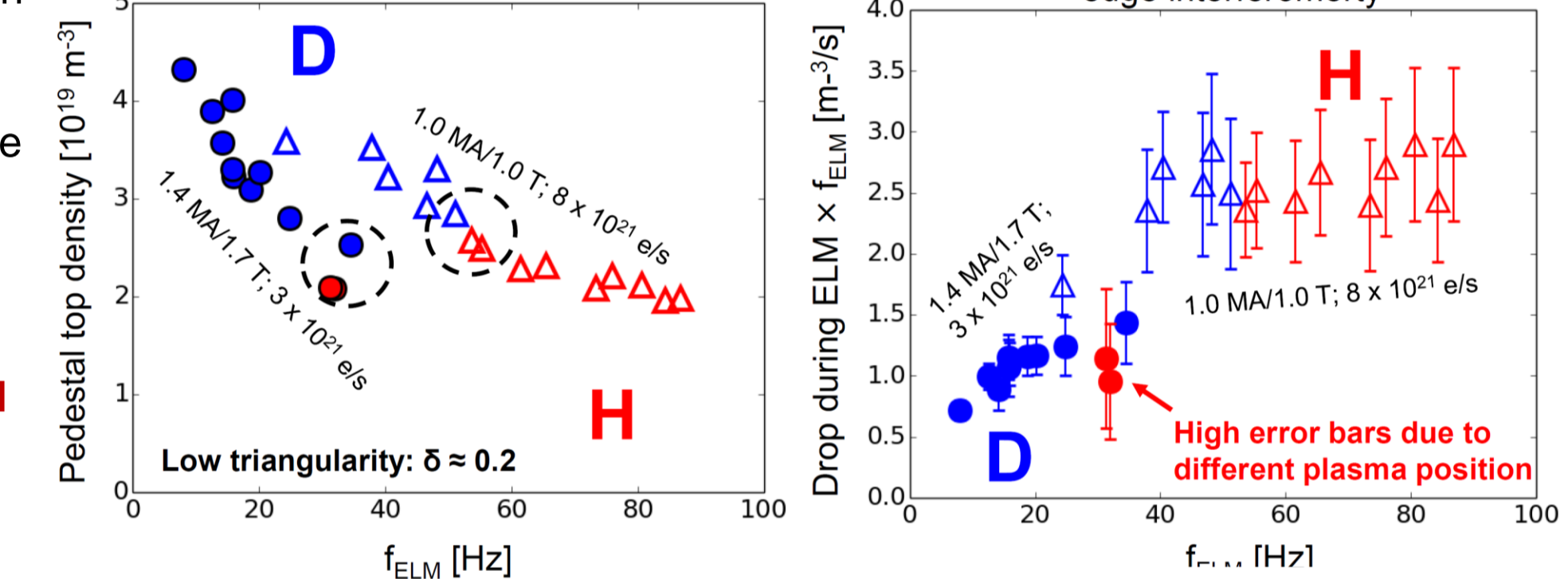
Type I ELMy H-mode: role of pedestal

- The lower particle confinement in hydrogen is add odds with idea that the higher thermal velocity should make fuelling easier
- \Rightarrow Transport more than overrides fuelling by neutrals
- Pedestal width model based on neutral penetration (Groebner 2002): $\Delta_{nc} \propto A^{-1/2} (T_{ped}/T_{sep})^{1/2} n_{ped}^{-1}$
- Scaling is not followed in the dataset, even reversed at 1MA, 1T!



\Rightarrow Transport processes that override neutral penetration differences are at work in pedestal (Horvath, NF to be submitted)

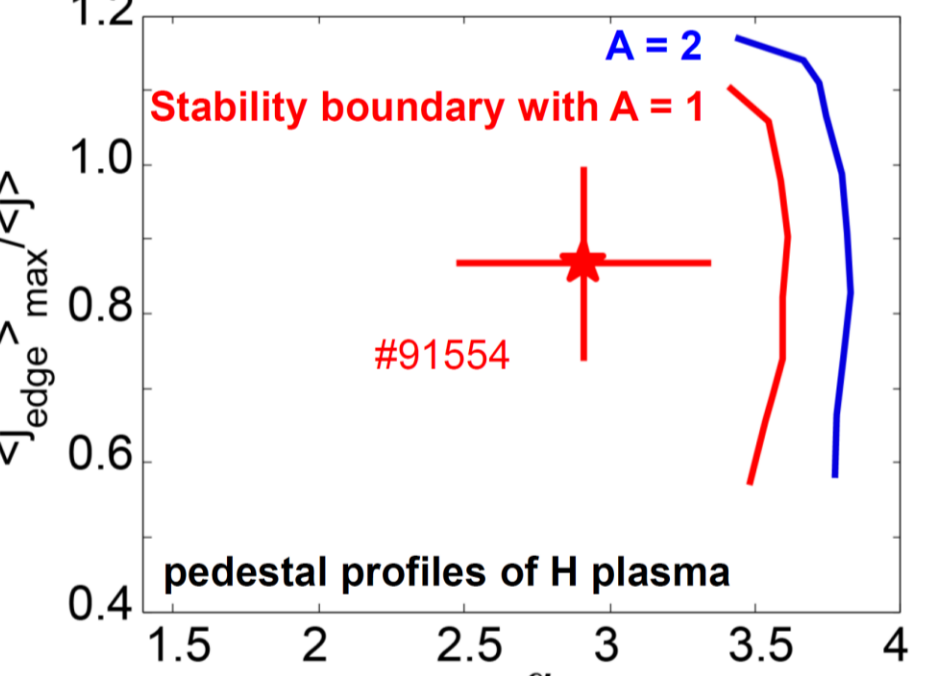
- Type I ELMs more frequent in H than D from same gas rate & power
- Pedestal density decrease with f_{ELM}
- However for $f_{ELM} > 40Hz$, ELM particle loss/ELM decreases and time average losses $\delta n \times f_{ELM}$ saturate



\Rightarrow ELMs alone cannot explain differences in density between H and D (Horvath, NF to be submitted)

Differences in pedestal stability may play a role:

- Operating point for H well inside s- α stability boundary (D is ~at boundary)
- Small (5%) reduction of stability boundary expected from peeling-ballooning stability criterion $\gamma > \omega_{dia}/2$ because $\gamma \propto A^{-1/2}$
- A possibly larger effect (up to 15%) would be expected if T_{ersep} is higher in H than in D ($T_{ersep} \approx 100eV$), as suggested by EDGE2D-EIRENE simulations ($T_{ersep} \leq 160eV$ in H).
- Strong ∇p region would be shifted outwards \Rightarrow P-B stability reduced, boundary shrinks
- Experimental data validation undergoing, but challenging



Summary & discussion

Type I ELMy H-mode:

- Strong isotope dependence in all transport channels: $\sim A^{0.4}$ to $\sim A^{0.5}$
- Gyrokinetic GENE analysis shows ITG is dominant
- GENE reverses GB scaling thanks to collisions and impurities
- Low particle confinement in hydrogen likely due to pedestal and edge transport processes
- Low particle confinement in hydrogen, leads to lower n_e , entailing lower energy and lower momentum confinement than in D

L-mode:

- Weak dependence of global energy / particle confinement on isotope ($\sim A^{0.14}$)
- Stiffness in TGLF QL modelling overcomes intrinsic GB dependence in local QL models, leading to ~no isotope scaling, but NOT to observed anti-GB scaling

Take home:

- Key to understanding and prediction remains edge/pedestal physics

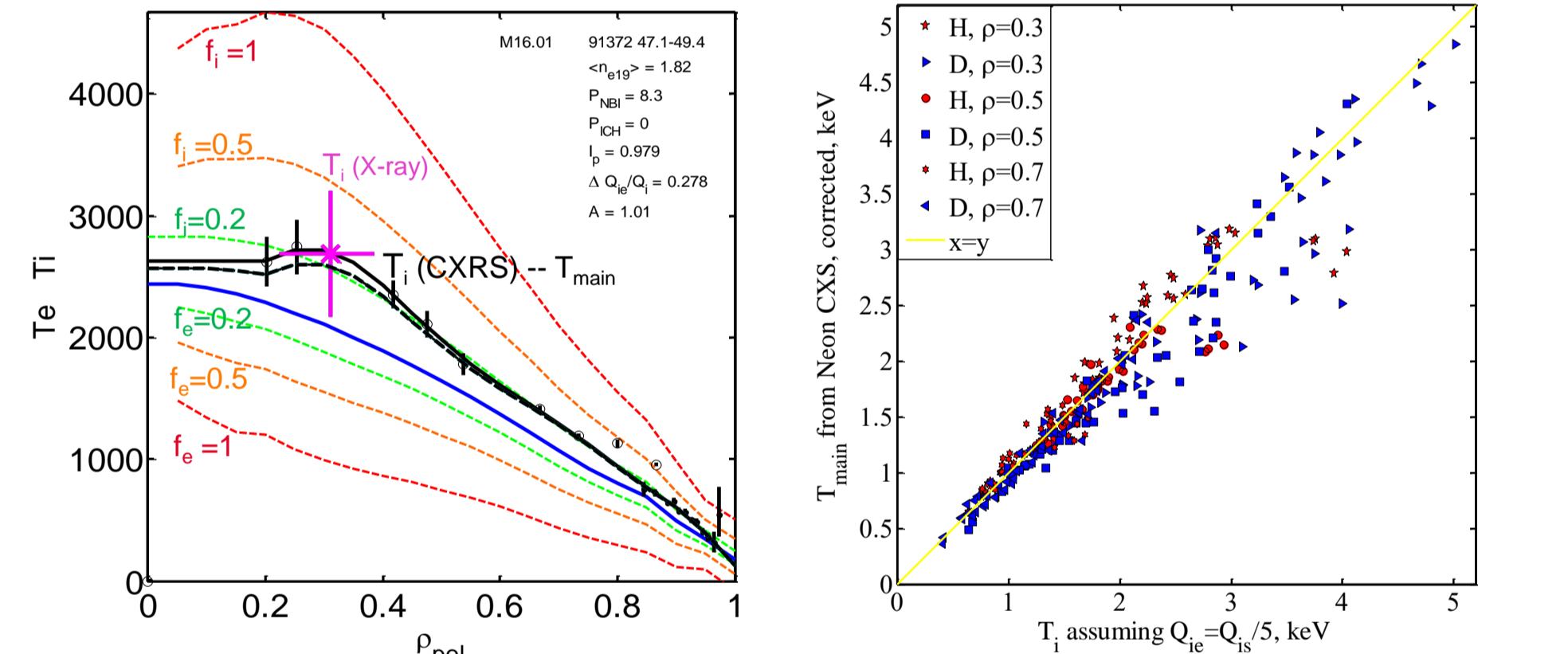
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 M. Oberparleiter, to be submitted
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Main ion temperature constrained by power balance

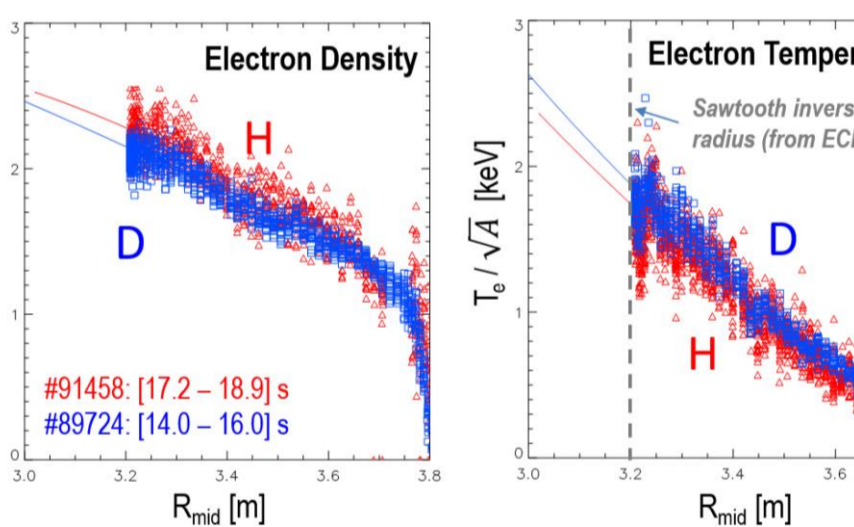
- Previous analyses (Maggi PPCF 2018) assumed $T_i=T_e$. Since then, CXRS data (C, Ne) became available, but only for half the dataset (and still noisy).
- Detailed analysis (H Weisen 2017) extends T_{imp} data to whole dataset for the main ions T_{main}
- A family of assumed $T_i(f_i, p)$ can be defined, assuming a certain fraction f_i of Q_{is} is transferred to electrons by equipartition.
- Shown: $f_i=0.2, 0.5$ & 1 (& if $Q_{ie} < 0$, $f_e = Q_{ie}/Q_{es}$)
- Multi-ion power balance shows T_{main} in range (0.95-1) $\times T_{imp}$ in core Most $T_{main}(p)$ close to $T_i(f_i=0.2, p)$, i.e. net ion heat flux is ~80% of deposited
- $T_i(p) = T_i(f_i=0.2, p)$ assumed in this study to extend T_{main} data to entire database



L-mode isotope identity experiment satisfies scale invariance

- An H/D dimensionless L-mode identity pair in ρ^* , β , v^* and q was successfully created by scaling the dimensional parameters as follows:
 - $I_p, B_T \propto A^{3/4}$, $n \propto A, T \propto A^{1/2}$
 - Scale invariance was achieved, i.e. the pair had identical normalised confinement time
 - $\omega_{ci} \tau_{EiH} \propto B_T \tau_{EiH} / A$
 - This is consistent with ion scale transport depending on ρ^* , β , v^* and, within errors, no additional isotope dependence
 - compared to the dimensional set still calls for understanding (Maggi, NF, to be submitted)

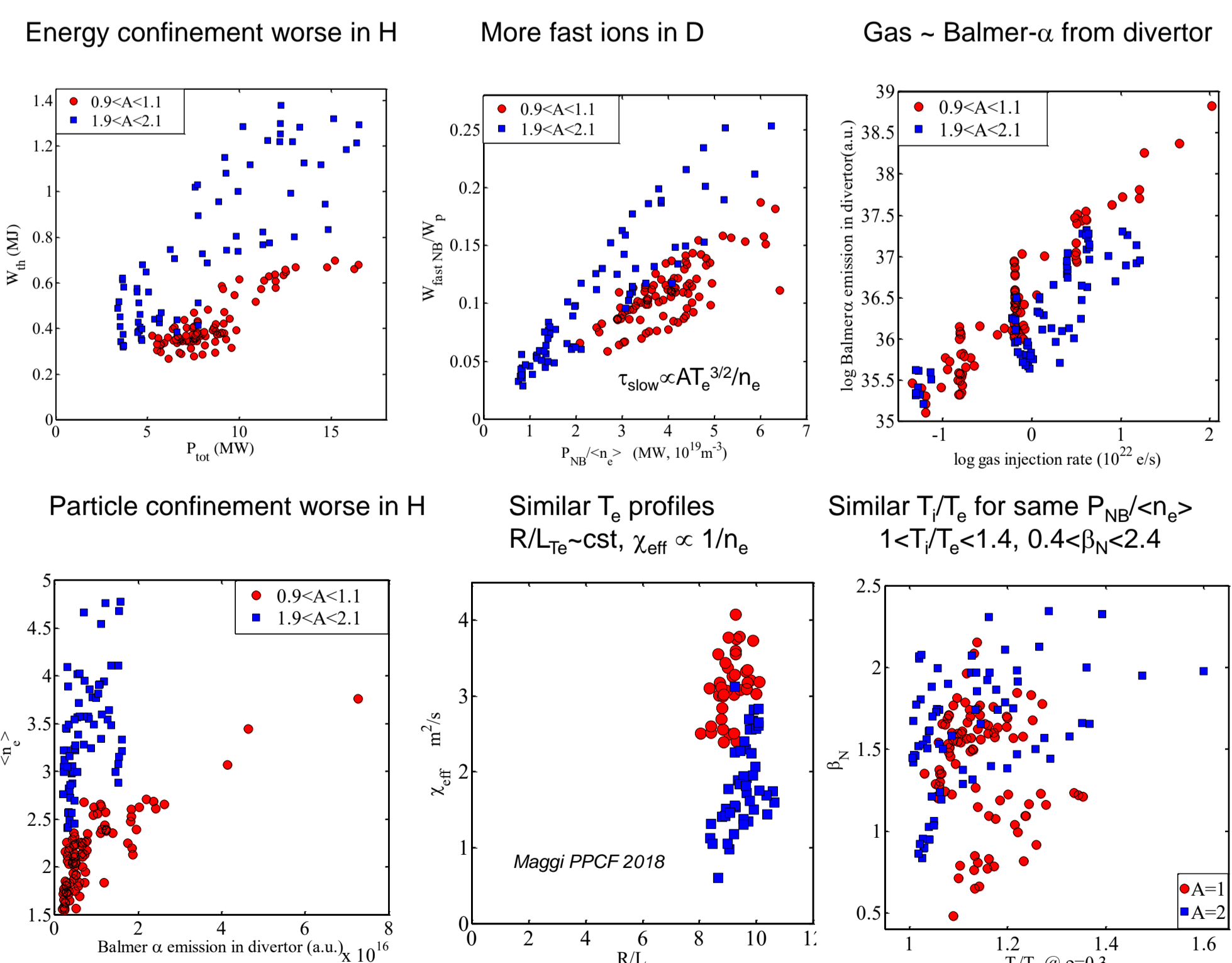
Pulse #	#91458	#89724
Isotope	H	D
Time interval [s]	17.2 - 18.9	14.0 - 16.0
B_T [T]	1.74	2.95
I_p [MA]	1.44	2.46
P_aux [MW] ($\pm 10\%$)	2.56	6.24
P_aux/B_T^3 [MW/T^3]	1.02	1.03
Z_eff ($\pm 10\%$)	1.4	1.35
τ_{EiH} [s] ($\pm 10\%$)	0.155	0.19
B_T τ_{EiH} / A [T/s]	0.27	0.28



L-modes: isotope dependence is weak

- NBI power scans in D and H at $n_e \approx \text{const}$
- Stiff temperature profiles: $R/L_{Te} \approx 8 m^{-1}$, $T_i/T_e \approx 1$
- GK analysis shows dominant mode is ITG in core (Maggi EPS 2018)
- Robust regressions for thermal stored energy W_{th} without and with n_e
 - $W_{th} \propto A^{0.15} p^{0.37}$ or $W_{th} \propto A^{0.14} p^{0.35} (n_e)^{0.62}$ (figure)
- Small ($\pm 10\%$, unintended) variations in density show global particle confinement also weakly dependent on isotope:
 - $N_e \propto A^{0.12} \Gamma_{main}^{0.27}$
- with $N_e = \int n_e dV$ total electron content and Γ_{main} the main chamber Balmer-alpha emission (a.u.) from a horizontal midplane viewing line, taken as a proxy for the particle source

Type I ELMy H-mode: overview



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