EPJ Web of Conferences 93,01014 (2015) DOI: 10.1051/epjconf/20159301014 © Owned by the authors, published by EDP Sciences, 2015

The (n,γ) campaigns at EXILL

J. Jolie^{1, a}, J.-M. Régis¹, D. Wilmsen¹, S. Ahmed¹, M. Pfeiffer¹, N. Saed-Samii¹, N. Warr¹, A. Blanc², M. Jentschel², U. Köster², P. Mutti², T. Soldner², G. Simpson^{3,9}, G. De France⁴, W. Urban⁵, F. Drouet³, A. Vancraeyenest³, P. Baczyk⁵, M. Czerwinski⁵, A. Korgul⁵, C. Mazzocchi⁵, T. Rzaca-Urban⁵, A. Bruce⁶, O.J. Roberts⁶, L.M. Fraile⁷, H. Mach⁷, V. Paziy⁷, A. Ignatov⁸, S. Ilieva⁸, Th. Kröll⁸, M. Scheck^{8,9}, M. Thürauf⁸, D. Ivanova¹⁰, S. Kisyov^{10,23}, S. Lalkovski^{10,11}, Zs. Podolyák¹¹, P.H. Regan¹¹, W. Korten¹², M. Zielinska¹², M.D. Salsac¹², D. Habs¹³, P.G. Thirolf¹³, C. A. Ur¹⁴, C. Bernards¹⁵, R.F. Casten¹⁵, N. Cooper¹⁵, V. Werner^{15,8}, R.B. Cakirli¹⁶, S. Leoni¹⁷, G. Benzoni¹⁷, G. Bocchi¹⁷, S. Bottoni¹⁷, F.C.L. Crespi¹⁷, B. Fornal¹⁸, N. Cieplicka¹⁸, B. Szpak¹⁸, C.M. Petrache¹⁹, R. Leguillon^{19,20}, R. John²¹, C. Lorenz²¹, R. Massarczyk²¹, R. Schwengner²¹, D. Curien²², R. Lozeva²², L. Sengele²², N. Marginean²³, R. Lica²³

Abstract. At the PF1B cold neutron beam line at the Institut Laue Langevin, the EXILL array consisting of EXOGAM, GASP and ILL-Clover detectors was used to perform (n,γ) measurements at very high coincidence rates. About ten different reactions were measured in autumn 2012 using a highly collimated cold neutron beam. In spring 2013, the EXOGAM array was combined with 16 LaBr₃(Ce) scintillators in the EXILL&FATIMA campaign for the measurement of lifetimes using the generalised centroid difference method. We report on the properties of the set-ups and present first results from both campaigns.

1 Introduction

During autumn 2012 and spring 2013 the EXOGAM spectrometer [1] and additional detectors were installed on the high intensity cold neutron guide PF1B [2] of the Institut Laue Langevin (ILL) for the EXILL campaign. EXILL is partially a follow-up of a previous campaign

using 8 EUROBALL capsule Ge detectors and a highly collimated cold neutron beam for (n,γ) measurements [3]. The EXILL campaigns, which took two reactor cycles of 49 days each, used beside the (n,γ) also the (n,fission) reaction on ²³⁵U and ²⁴¹Pu targets. An overview of the (n,γ) experiments performed as well as the first results from the (n,γ) reaction on ¹⁹⁴Pt and ¹⁹⁵Pt is presented here.

¹IKP, University of Cologne, Zülpicher Str. 77, D-50937 Köln, Germany,

²ILL, 71 Avenue des Martyrs, 38042 Grenoble CEDEX 9, France

³LPSC, 53 Avenue des Martyrs, 38026 Grenoble, France

⁴GANIL, BP 55027, 14076 Caen CEDEX 5, France

⁵Faculty of Physics, University of Warsaw, ul. Hoza 69, PL-00-681 Warsaw, Poland

⁶SCEM, University of Brighton, Lewes Road, Brighton BN2 4GJ, UK

⁷Gruppo de Física Nuclear, Facultad de CC. Físicas, Universidad Complutese, 28040 Madrid, Spain

⁸Institut für Kernphysik, TU Darmstadt, Schlossgartenstrasse 9 64289 Darmstadt, Germany

⁹School of engineering, University of the West of Scotland, Paisley, PA1 2BE, UK

¹⁰Faculty Physics, University of Sofia, Bulgaria

¹¹Dep. Of Physics, Univ. of Surrey, Guildford GU2 7XH & Nat. Phys. Lab., Teddington, UK

¹²CEA, Centre de Saclay, IRFU, F-91191 Gif-sur-Yvette, France

¹³Fakultät f. Physik, LMU Münich, 85748 Garching, Germany

¹⁴INFN Sezione di Padova, 35131 Padova, Italy

¹⁵WNSL, Yale University, New Haven (CT), USA

¹⁶Istanbul University, Istanbul, Turkey

¹⁷University of Milano and INFN sez. Milano, Italy

¹⁸The Niewodniczanski Inst. of Nucl. Physics, Polish Academy of Sciences, Krakow, Poland

¹⁹CSNSM, Université Paris Sud and CNRS/IN2P3, 91405 Orsay, France

²⁰Advanced Science Research Center, Japan Atomic Energy Agency, 2-4 Shirakata Shirane, Tokai, Ibaraki 319-1195 Japan

²¹Helmholtz-Zentrum Dresden-Rossendorf, 01328 Dresden, Germany

²²IPHC, IN2P3/CNRS and the University of Strasbourg, 67037 Strasbourg, France

²³Horia Hulubei NIPNE, 77125 Bucharest, Romania

^a Corresponding author: jolie@ikp.uni-koeln.de

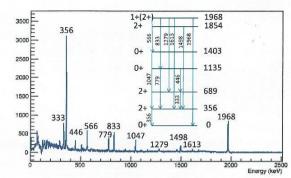


Figure 1. Coincident gamma rays in ¹⁹⁶Pt gated by the primary 5953 keV transition showing the deexcitation of the 1968 keV level

The main advantages of the (n,γ) reaction with thermal or cold neutrons are listed here. Because there is no Coulomb barrier the meV neutrons end up in capture states situated at the neutron separation energies. There is no orbital angular momentum transfer and the capture states have spin $J^{+1/2}^{\pi}$ and (or) $J^{-1/2}^{\pi}$ with J^{π} the spin and parity of the ground state of the target nucleus. They generally decay by primary (E1) transitions of several MeV populating, in a statistical way, excited states. By setting gates on these primary gamma rays a given level is selected and all coincident gamma-rays are emitted in the decay of the selected level to the ground state (see Figure 1). Finally, measurements of angular correlations allow the determination of spins and multipolarities.

2 The set-ups used

The collimator described in Ref [3] provided, at the target position, a 12-mm diameter cold neutron beam of ~10⁸n/(s.cm²). Perpendicular to the beam, 8 EXOGAM Clover detectors, with their BGO shields, were mounted in a ring while the additional detectors were mounted in the forward and backward directions at 45° (see figure 2). During the first reactor cycle and part of the second one, 6 unshielded GASP Ge detectors and 2 unshielded ILL-Clover Ge detectors were mounted in addition to EXOGAM. A digital data acquisition allowed event rates up to 0.84 MHz to be handled [4].

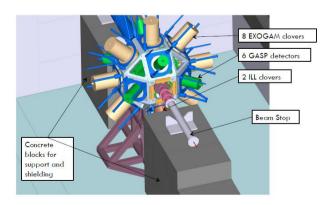


Figure 2. Set-up used for γ - γ and γ - γ - γ coincidences at the PF1B cold neutron beam.

With this set-up the following (n,γ) experiments were done:

- 'High precision measurement of multipole-mixing ratios to test the newly proposed candidate for a low-lying octupole-isovector state in ¹⁴⁴Nd.' M. Scheck (UWS) and T. Kroell (TU-Darmstadt).
- 'A γ - γ angular correlation measurement of ¹⁶⁸Er after (n, γ).' J. Jolie (Cologne)
- 'Low-spin doorway states feeding the 13/2 highspin isomer of ¹⁹⁵Pt following ¹⁹⁴Pt(n,γ).' D. Habs (LMU München)
- Properties of excited states of nuclei close to N=40 produced in (n,γ) reaction. A. Korgul (Warsaw)
- 'Search for particle-phonon coupled states in ⁴⁹Ca by the reaction ⁴⁸Ca (n,γ) ⁴⁹Ca.' S. Leoni (Milano)
- 'Studying three-step gamma cascades following neutron capture in ¹⁶¹Dy.' M. Krticka (Prague)
- 'Deformation of ground and beta-bands: the 2⁺ beta to 0⁺ beta decay in ¹⁵⁸Gd.' Ch. Bernards and V. Werner (Yale)
- 'Measurement of gamma-ray cascades in ⁷⁸Se-new information on the Pygmy dipole resonance.' R. Schwengner (Rossendorf)
- Identification of the nature of 3 levels in 96 Mo exploiting the 95 M(n, γ) reaction. M. Scheck (UWS)
- 'Low-spin structure and particle-phonon couplings in 210 Bi studied by using 209 Bi(n, γ) 210 Bi.' B. Fornal (Warsaw)
- 'The negative parity bands and octupole collectivity: Search for tetrahedral symmetry in 156Gd' D. Curien (Strasbourg)

Here we report on results from the 194 Pt(n, γ) 195 Pt reaction.

During 30 days of the second cycle 16 LaBr₃(Ce) scintillators from the fast-timing array (FATIMA) collaboration were combined with the 8 EXOGAM detectors for fast-timing measurements. With this set-up the following (n,γ) experiments were done:

- 'Gamma-ray spectroscopy of 47 Ca by the reaction 46 Ca(n, γ) 47 Ca.' S. Leoni (Milano)
- 'Low-spin structure and particle-phonon couplings in ²¹⁰Bi studied by using ²⁰⁹Bi(n,γ) ²¹⁰Bi.' B. Fornal (Warsaw)
- Test of the O(6) selection rule in ¹⁹⁶Pt.' J. Jolie (Cologne)

From this campaign we report here on results of the $^{195}Pt(n,\gamma)^{196}Pt$ reaction.

3 First results for the 194 Pt(n, γ) 195 Pt reaction

Recently, an alternative production of the platinum isomer 195m Pt by resonant (γ,γ') photo excitation induced by intense monochromatic gamma beams was proposed

[5]. This method relies on a suitable 'doorway state' linking the low-spin ground state with the isomer. The selective population of this isomer is of special interest, as it shows unique properties for diagnostic and therapeutic purposes in nuclear medicine, especially in the treatment of cancer. The 13/2+ 195mPt isomer has a half-life of 4 days and it intensively emits X-rays (59% K_{α} at 66 keV, K_{β} at 76 keV) and gamma rays (11.7% at 99 keV). Thus, attaching the highly activated 195mPt to a suitable biomolecule, it can be used as a tracer for imaging with single-photon emission computed tomography (SPECT) or gamma cameras. Another benefit for the use of ^{195m}Pt refers to the frequently used platinum based chemotherapeutics like cis-platin or carbo-platin. Here, ^{195m}Pt could offer a useful tool to assess the tumor uptake during the chemotherapeutical treatment. In addition, ^{195m}Pt also emits an abundant amount of Auger and conversion electrons, which can be used therapeutically as their energy is deposited in a short range, maximizing the tumor to normal tissue dose ratio. Experimentally, a sample of highly enriched ¹⁹⁴Pt was placed in the very intense cold neutron beam in order to investigate the nuclear structure of ¹⁹⁵Pt via the reaction 194 Pt $(n,\gamma)^{195}$ Pt and search for such 'doorway states'.

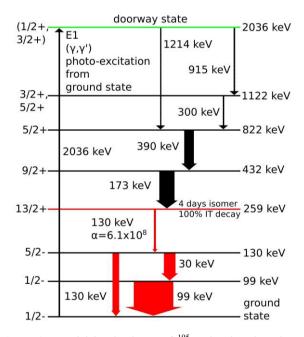


Figure 3. Partial level scheme of ¹⁹⁵Pt showing the relevant isomeric transitions (red), the isomer populating transitions (black) and the idea of (γ,γ') induced activation of the isomer ^{195m}Pt from the ground state of ¹⁹⁵Pt via the doorway state (green) (color online).

In a still ongoing analysis the level scheme is being investigated using coincidence measurements and angular correlations with special regard to the isomer populating pathways. Therefore, high energy primary transitions decaying from the neutron capture state at 6.105 MeV, especially to the strong isomer populating starting at the 1122-keV level with a cascade of 300 keV, 390 keV, 173 keV transitions, which first has been observed by D. D. Warner *et al.* in 1982 [6], were investigated. A 'doorway

state' in ¹⁹⁵Pt with an energy of 2.036 MeV was found, which is populating the isomer ^{195m}Pt and, in addition, directly decays into the ground state fulfilling the requirements of gamma ray induced activation of ^{195m}Pt (see figure 3, 4).

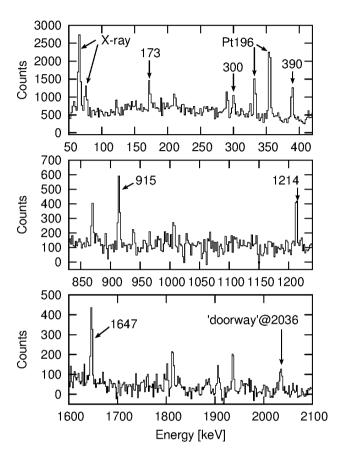


Figure 4. Background corrected coincidence spectrum of the relevant energy ranges in ¹⁹⁵Pt gated on the primary transition of 4.069 MeV leading to the 'doorway state' at 2.036 MeV. The upper figure shows the strong isomer populating transitions of the 'Warner's cascade'; the middle figure shows two transitions depopulating the 'doorway' and feeding the cascade; the lower figure shows the doorway state decaying directly into the ground state (cf. fig. 3).

Moreover other promising γ -transitional pathways leading to the isomer were observed.

4 First results for the 195 Pt(n, γ) 196 Pt reaction

Many years ago, the even-even nucleus ¹⁹⁶Pt was proposed as an example of the SO(6)-limit of the Interacting Boson Model (IBM) [7]. When the Hamiltonian for N s,d bosons is composed by the linear combination of second order Casimir operators of the group chain: SO(3) \subset SO(5) \subset SO(6) \subset U(6), the quantum numbers classifying the irreducible representation (irrep) of these groups: L, v, σ , N are valid and the Hamiltonian is analytically solvable. The lowest states have $\sigma = N$. States at higher energies can have

 $\sigma = N-2$, N-4,.... Recently, the need to test the SO(6) symmetry and not only the SO(5) properties in ¹⁹⁶Pt was stressed [8]. When testing the goodness of the description of an SO(6) nucleus one thus should look for properties involving states with different σ . One of these properties concerns the electric quadrupole transition operator. In the SO(6) limit, the E2 transition operator is a generator of SO(6). Now, generators cannot act outside their irrep and the strict selection rule $\Delta \sigma = \theta$ follows. In order to test this selection rule, absolute B(E2) values between $\sigma = N$ -2 and $\sigma = N$ states need to be measured. This is not an easy task as the first $\sigma = N-2$ states in ¹⁹⁶Pt are at high energy and have low spins. They are not populated by standard in-beam reactions with ions, but they are by neutron capture. It is, therefore, not an accident that Ref [7] relied on ILL data. In 1990 an attempt was made to measure the lifetime of the lowest $\sigma = N-2$ state at 1402.7 keV using the GRID method [9] at GAMS4 [10]. For this 0^{+}_{3} state a lower lifetime limit of $\tau > 1.8$ ps could be established showing that the B(E2) towards the 2⁺₁ was smaller than 0.034 e²b² and, as such, at least an order of magnitude smaller than allowed transitions between states with $\sigma = N$.

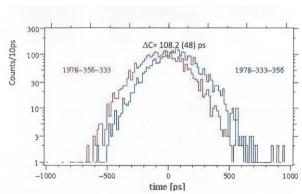


Figure 5. The binned time spectra using the 356 keV ground state transition as start for the TAC (left curve) and once as stop (right curve) (color online).

Using the FATIMA and EXILL set-ups triple coincidences can be used in which a Ge gate selects the cascade of gamma-rays to be measured by the LaBr₃(Ce) scintillators. These fast signals can then be used for the generalized mirror symmetric centroid difference method [11,12] to measure lifetimes. The prompt response difference curve (PRD) was obtained using an ¹⁵²Eu source and transitions in ⁴⁹Ti after neutron capture [13,14]. The method can be used to measure lifetimes down to about 10 ps.

The experiment took place during 19 hours. Because the thermal cross section for 195 Pt(n, γ) 196 Pt is 28 barns and dominates over those of the other Pt isotopes, a nat Pt target of 143 mg was used. The Ge total count rate was 263 kHz and the LaBr₃ one 192 kHz. In total 3 x 10 10 events were measured. Figure 5 shows the time spectra of the 333-356 keV cascade gated by the 1978-keV transition in 196 Pt. The observed centroid difference between both spectra corresponds to $\Delta C = 108(5)$ ps. Using the PRD curve and $\Delta C = PRD(333)$ -PRD(356) + 2τ together with PRD(333)-PRD(356) = 8(10) ps [12] a

lifetime of $\tau(2^+_1) = 50(6)$ ps is deduced which agrees very well with the literature value of 49.2(2) ps. Being confident that the method works, the lifetime of the 1402-keV 0^+ state was then measured. Again the 356-keV transition provided the decay signal for the generalized centroid difference method.

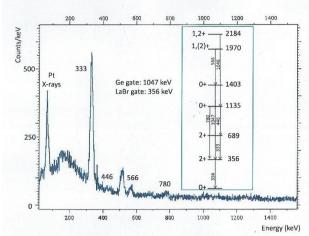


Figure 6. Relevant part of the level scheme of ¹⁹⁶Pt (insert) and the double-gated LaBr₃(Ce) spectrum.

The feeding signal was obtained from the 566-keV transition populating the 1402-keV state. The Ge gate was set on the 1048-keV transition of the decay towards the first excited state. Figure 6 shows the double-gated LaBr₃(Ce) spectrum. The result is shown in Figure 7. The effective lifetime deduced from the centroid shift $\Delta C = 53(11)$ ps using PRD(566)-PRD(356) = -54(10) then yields $\tau_{\text{eff}} = \tau(2^+_1) + \tau(0^+_3) = 54(7)$ ps leading to an upper limit for the lifetime of $\tau(0^{+}_{3})$ < 12 ps. Using also the lower limit from [10] we find that 0.75 W.u < $B(E2;0^{+}_{3}\rightarrow 2^{+}_{1}) < 5 \text{ W.u and } 0.06 \text{ W.u } < B(E2;0^{+}_{3}\rightarrow 2^{+}_{2})$ < 0.41 W.u. Clearly no collective B(E2) values are found between the states, confirming the validity of the SO(6) symmetry. A recent Gammasphere experiment using multiple Coulomb excitation confirmed our result with $0.62 \text{ W.u} < B(E2;0^{+}_{3} \rightarrow 2^{+}_{1}) < 4.9 \text{ W.u} [15].$

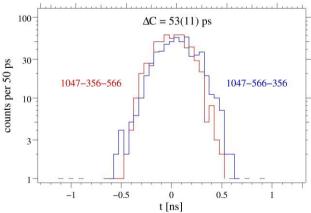


Figure 7. The binned time spectra using the 356 keV ground state transition as start for the TAC (left curve) and as stop (right curve) (color online).

5 Conclusion

The EXILL campaigns have provided a huge amount of new data using a high intensity cold neutron beam and the (n,γ) and (n,fission) reactions. Here we reported on the (n,γ) experiments from the EXILL and EXILL&FATIMA campaigns and showed promising first results on ¹⁹⁵Pt and ¹⁹⁶Pt. The success of the campaigns has led to the proposal of a dedicated new ILL instrument for prompt fission studies using gamma-ray spectroscopy called FIPPS [16]. The option of using also the (n,γ) reaction at FIPPS is important.

Acknowledgements

This work was supported by NUPNET by contract 05P12PKNUF, 05P12RDNUP and DNC7RP01/4, by the Science and Technology Facilities Councils (UK), by the Spanish MINECO via the FPA2010-17142 and PRI-PIMNUP-2011-1338 projects, and by the DFG by cluster of Excellence 'Origin and Structure of the Universe', grant KR 1796/2-1 and SFB634. This material is based in part upon work supported by the U.S.D.O.E., SC, Office of nuclear physics, under Award No. DE-FG02-91ER40609. The EXILL campaign would not have been possible without the support of several services at the ILL and the LPSC. We are grateful to the EXOGAM collaboration for the loan of the detectors, to GANIL for assistance during installation and dismantling, to the INFN Legnaro laboratory for the loan of the GASP detectors and to the FATIMA collaboration for the use of the LaBr₃(Ce) scintillators.

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