

Particle precipitation by bipolar corona discharge ion winds

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Abstract: The paper reports the development of a particle precipitation based aerosol sampler using bipolar corona discharge ion winds with collected particles of minimized net charge. For the new approach, neutralized particles move towards a sampler under the effect of electric field and dual ion winds. Since there is no electrode or sampling chip installed inside the air-flow channel, impediments to airborne particle flow or ion winds are removed along the flow direction. In addition, the isolation of ion winds, which generate circuit, allows using various materials for the sampling chip including non-conductors and also protecting collected particles from any discharge ignition on the chip. The device mechanism is numerically simulated in OpenFOAM to study the electrofluidodynamic interaction of charged particles and bipolar ion winds. The efficiency of the new approach has been investigated by experiment with a maximum efficiency of 94%. The effects of flow rate, discharge voltage and electrode distances on the method are also evaluated.

Keywords: Charged particles, ionic wind, particle simulation, corona discharge, OpenFOAM

1. Introduction

Research in the healthcare shows that the particulate matter (PM) causes up to 30% of the total burden of diseases [1]. The inhalation of PM could yield numerous lung diseases such as pulmonary fibrosis [2], pulmonary inflammation, pleural effusion, granuloma [3] and cancer risk [4]. Thus, the particle sampling in bio-technology has attracted significant effort to develop new techniques in various application areas, especially in the biomedical engineering. The main objective of the sampling is to collect a sample of airborne particles and then transport them to a detection unit. The active sampling devices are designed based on three major collection methods: the impaction, impingement and filtration [5]. For the impaction approach, the sampling is carried out by the collision of particles to a wall and called as the inertial impaction. Meanwhile in the impingement scheme, a flow of particles is channelled to a collection chamber through nozzles of an air jet. The number of collected particles depends on the air jet characteristics such as the nozzle geometry and the particles' diameter [6–8]. Several commercial devices using this technique include Coriolis® μ air sampler [9] and Bio-sampler [10]. For the filtration method, particles are collected by driving air through a membrane of high density [11–14]. Similar to the impaction technique, devices using the filtration scheme would be overloaded when working in highly contaminated environments [15]. Among sampling devices, the electrostatic precipitator (EP) is evaluated as an efficient technique to filter nanoparticles at low pressure drop [16–19]. This technique is usually applied in bioengineering because it provides softer collection for sensitive micro-organisms which request a slow landing velocity on a substrate [20]. For this approach, a flow of particles is ionized at the upstream by an ionizer [21] before going through an electric field created by a pair of electrodes placed at the downstream of channel as described in Mainelis [22]. Most commonly, an electrostatic precipitator is developed for airborne nanoparticles

1 using a needle tip to generate a strong field towards a ground plate [23]. Another automated electrostatic sampler
2 for bio-aerosol was proposed [24] using a sphere ground electrode installed at the entrance of aerosol and a
3 sampling high voltage electrode placed in the middle of air channel. Several modifications of this approach were
4 carried out as follows. Instead of sphere ground electrode, a pair of pins is placed in parallel with the aerosol
5 direction and surrounded by a sheath flow while the sampling electrode is installed in the middle of the flow
6 direction and connected to ground [21]. Lee et al. utilized the point-to-plane configuration to create an electric
7 field across the aerosol direction drifting the aerosol to a plane of liquid [25]. Park and his co-workers [26] also
8 used the same configuration but with a liquid covered counter electrode drained out to an adenosine triphosphate
9 bioluminescence array in order to detect the increase of bioluminescent shortly. Another type of electrostatic
10 generated in axisymmetric configuration is wire-to-cylinder with a high voltage electrode applied on the outer
11 wall and the inner wall acts as the sampling stag [27–29]. In spite of many significant progresses, there are several
12 concerns for the mentioned techniques using the conventional corona discharge with single polarity (uni-polarity).
13 Indeed, although more investigations are requested, several recent publications [30,31] showed that electrostatic
14 field would damage the culturability of collected bio-particles (bacteria, fungi and other cellular derivatives) due
15 to the excessively charged stress. Furthermore, the particle collection by unipolar charge systems causes a
16 significant difference compared with real results in monitoring a collection due to the intrinsic charge of bio-
17 particles. It is up to the charge polarities of particles and generator, the error could be up to 50% as represented
18 in [30].

19 Recently, we have developed a novel method to generate ion winds using two pin electrodes which are placed in
20 parallel with a designed inter-electrode gap [32]. Such configuration allows the ion winds of opposite charge from
21 two electrodes moving forward to the space in front of the electrodes' tips [32–38]. In this work, a new
22 electrostatic aerosol sampling method using bipolar corona dual ion winds is developed. For the present approach,
23 particles collected at a collection chamber are almost neutralized owing to dual ion winds generated by an electric
24 field. Besides, the new configuration also possesses several advantages such as (i) the removal of undesired
25 impediments to airborne particle flow and ion wind along a channel and (ii) the flexibility in choosing various types
26 of material for the sampling chip including non-conductive ones. The performance of the present method is
27 investigated with both experimental work and numerical simulation using OpenFOAM finite volume method.

28 **2. Mechanism of the present aerosol electrostatic sampler**

29 As summarized above, traditional aerosol electrostatic precipitator is configured by one discharge electrode and
30 one counter electrode acting as the sampling electrode for the particle collection [23–25] (see Fig. 1a). Thus, the
31 sampling electrode must be conductive and therefore it is not suitable with several biocompatible materials, such
32 as polydimethylsiloxane, polymethylmethacrylate or silicone gel. Moreover, since the conventional EP creates a
33 considerable amount of charge, it is not appropriate for bio-sampling in several cases as mentioned in section 1.
34 In addition, it proved to be challenging to segregate charged droplets from a plate due to the presence of Van-de-
35 Wall force. This issue has recently attracted to numerous researches on bio-aerosol sampling using a solvent on
36 the lab-on-chip interface termed as the aerosol-to-hydrosol transfer [39–42].

37 For our research, the new technique to collect airborne particles of zero net charge is developed using a dual ion
38 wind based generator (see Fig. 1b). **The mechanism of the scheme is described as follows. With a high voltage**
39 **applied between two parallel pin electrodes, a curved electric field is generated and expands outward the space**
40 **between the two electrodes and then induces two ion winds simultaneously as shown in Fig. 1b. The two ion**

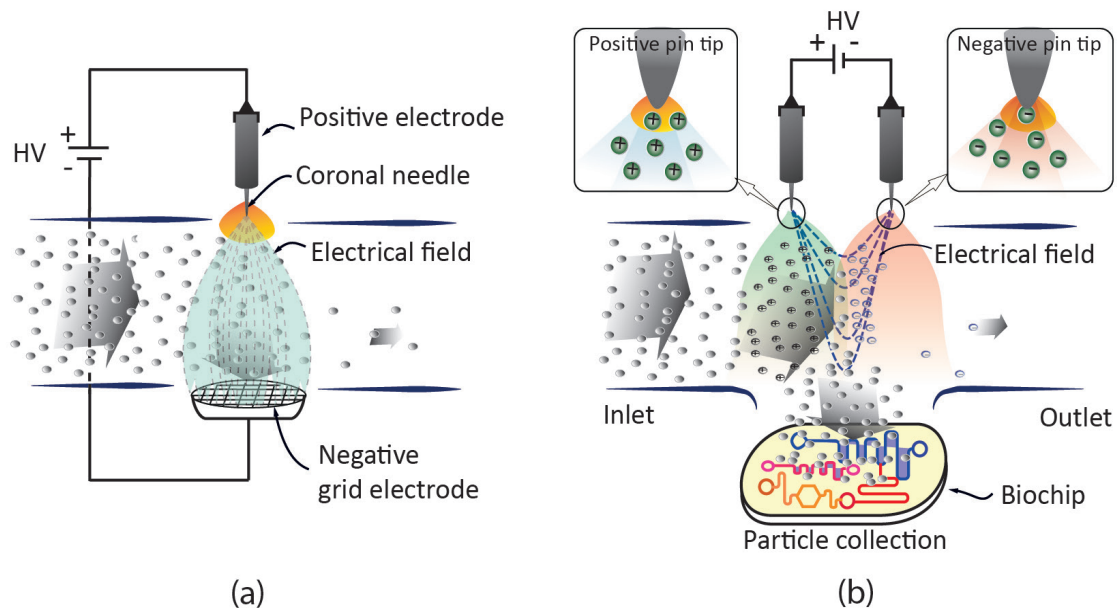


Figure 1. Configuration and mechanism of a bioaerosol electrostatic samplers. (a) Traditional configuration: High voltage is applied between discharged electrode and the collector to generate a high electric field, the collector is a part of high voltage circuit. (b) The present configuration: High voltage is separated from the collector; Electric field between two discharge electrodes creates ion winds. Particles are neutralized and directed towards the collector under the effect of the electric field and ion winds.

1 winds have the same but opposite charges. Regardless of the initial charge of the entrance particles, they are
 2 charged by the electric field while moving through the channel. Under the interaction with the ion wind at the
 3 upstream positive electrode, the aerodynamic force by ion wind and the electrostatic force drift charged particles
 4 toward a particle sampling chip (collector) installed in the opposite side of the electrodes. On the way to the
 5 collector, charged particles move downstream due to its initial airflow momentum and the curved electric field,
 6 and continue interacting with the negative ion wind from the downstream electrode. The negative ion wind
 7 further pushes particles toward the opposite side of the channel and neutralizes their charge as shown in Fig. 1b.
 8 As a result, the net charge of airborne particles is alleviated or even neutralized when they reach the collector due
 9 to the charge balance from two ion winds [37].

10 As presented, since neither the electrodes nor the sampling chip is installed inside the channel, there is not any
 11 impediment on the way of the airborne particles and the ion winds. Due to the insulation of sampling chip from
 12 the ion winds generating circuit, various materials can be used for the collector including glass, gel, metal or
 13 semiconductor. As the airborne particles sampled with very low net charge, the present scheme should be suitable
 14 for developing bio-aerosol electrostatic precipitator.

3. Experiment setup

3.1 Sampler prototype

A schema of the present aerosol sampler prototype is shown in Fig. 2. The system consists of a channel made of polypropylene for the aerosol, an ion wind generator assisting the sampling process and a sampling chamber for particle collection. The circular shaped inlet and outlet parts are connected smoothly with a body of 15 x 15 mm² rectangular cross-section and 70 mm length to avoid possible abrupts of the aerosol flow. As proposed by Asbach et al., [43] Tygon is selected as a material for two connecting tubes at the upstream and downstream of the channel to minimize the particle loss during the experiment. The ion wind generator is installed in the upper side of the channel centre while the sampling chamber is in the opposite side with an insulated tray of acrylic glass (see Fig. 2).

The ion wind generator consists of two stainless steel SUS304 electrodes placed in parallel with each other, both with 8 mm length and 0.5 mm diameter. The two electrodes with tip size of 80 μm are fixed inside a holder insulated from surroundings. The tips of electrodes are placed in a hole, 2 mm away from the wall in order not to cause an obstacle along the flow of particles as well as to protect the flow from undesired contaminants by electrodes. A battery-operated high voltage generator capable of generating up to 10 kV DC voltage is connected to the electrodes. In order to prepare the experiment, the air velocity is determined using a conventional air flow meter (AWM5102VN) to establish the velocity profile inside the channel. The current and voltage (I-V)

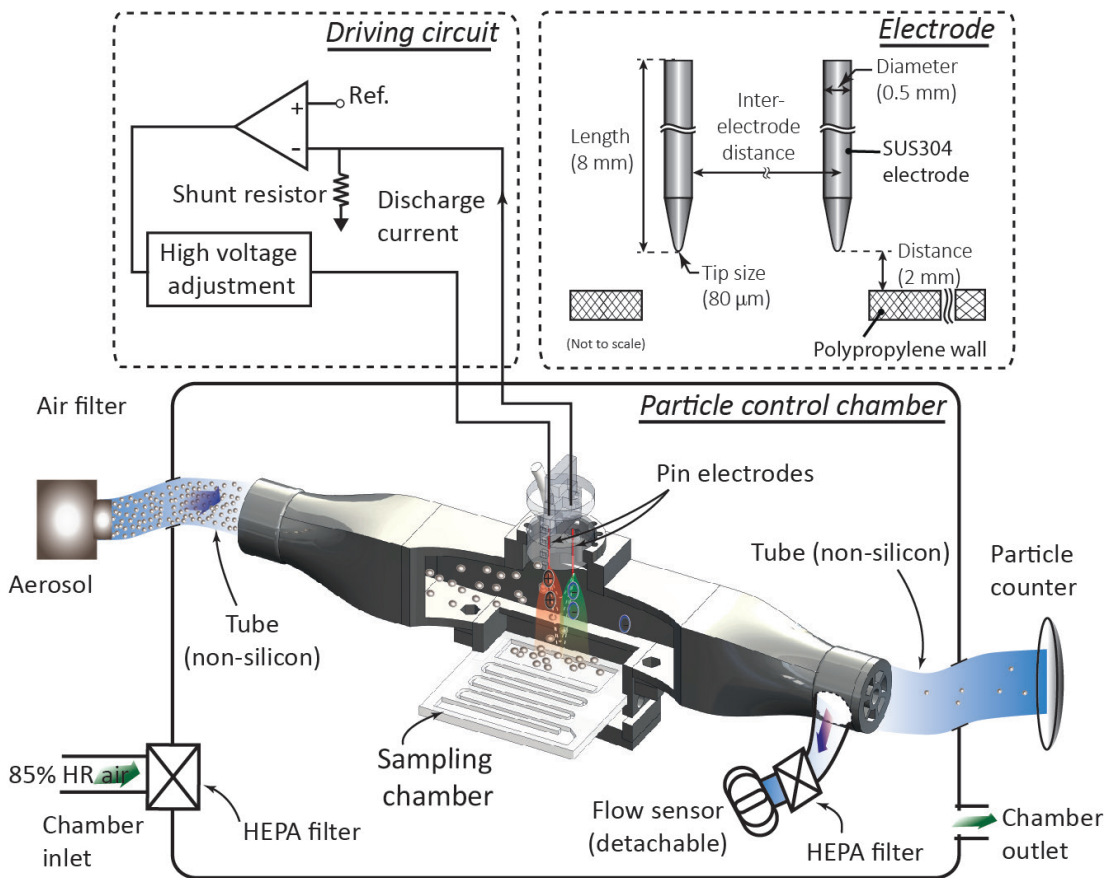


Figure 2. The present aerosol electrostatic sampler (AES): A schema of the AES together with its mechanism and configuration.

1 characteristic of ion wind is recorded for an entire range of air flow rates throughout the experiment by an
2 oscilloscope (HAMEG-R&S). The experiment was carried out with the ambient temperature of 24 - 25 °C, relative
3 humidity of 55% – 65% and the atmospheric pressure in a control box of the size of 25 cm × 40 cm × 65 cm, where
4 clean air is supplied through a filter (HEPA).

5 3.2. Measurement protocol

6 The efficiency of the present particle sampler is investigated based on the ratio of number of particles counted
7 at the outlet of channel in the two cases with and without the effect of ion winds for a given flow rate as follows
8 [24].

$$9 \quad \eta = 1 - \frac{N_{outlet-On}}{N_{outlet-Off}}, \quad (1)$$

10 where $N_{outlet-on}$ and $N_{outlet-off}$ are the particle concentrations measured at the air flow channel outlet with
11 and without the corona discharge, respectively.

12 For experiment, particles collected on the tray are followed up by an optical microscope and counted by a particle
13 counter. A sample of polystyrene particles with diameter of 1.0 μm is diluted with deionized water, filtered by 0.2
14 μm nano-porosity membrane (Merck Millipore Ltd), and stored in room condition for 30 minutes. The particles
15 are then stirred up by a mild ultrasonic mixer to prevent their agglomeration and then introduced into an aerosol
16 generator placed at the inlet of channel. Although the system can be operated with a large range of particle sizes
17 to simulate a natural human breath, at this step the trace by a sample of only one particle size of 1 μm allows to
18 investigate easily but efficiently the performance of the device.

19 The influence of corona parameters on the efficiency of the present device can be investigated using the Deutsch-
20 Anderson equation [44] as follows.

$$21 \quad \eta = 1 - e^{-\frac{n_p e C_C E}{3\pi\mu d_p} \times \frac{A}{Q}}, \quad (2)$$

22 where A is the collecting area; n_p the particle charge; e the elementary charge; C_C the slip correction factor; E the
23 electric field; μ the viscosity; Q the flow rate of aerosol and d_p the particle diameter. Eq. (2) depicts that collection
24 efficiency increases with the increase of particle charge (n_p) but with the decrease of flow rate.

25 The measurement protocol proceeds as follows. One minute after turning on the aerosol generator, a high
26 voltage is applied on the two electrodes to induce a corona discharge current. Until the current reaches a stable
27 state, an aerosol meter (TSI 9306) is switched on to start measuring the particle concentration every 20 seconds.
28 This measurement process is repeated five times and the experiment is conducted with a range of flow rates from
29 the generator. The process is controlled by a compressor and monitored at the channel outlet using a flow sensor
30 AWM5102VN (Honeywell, USA). Experimental results are presented in the next section.

31 3.3. Experimental results and discussions

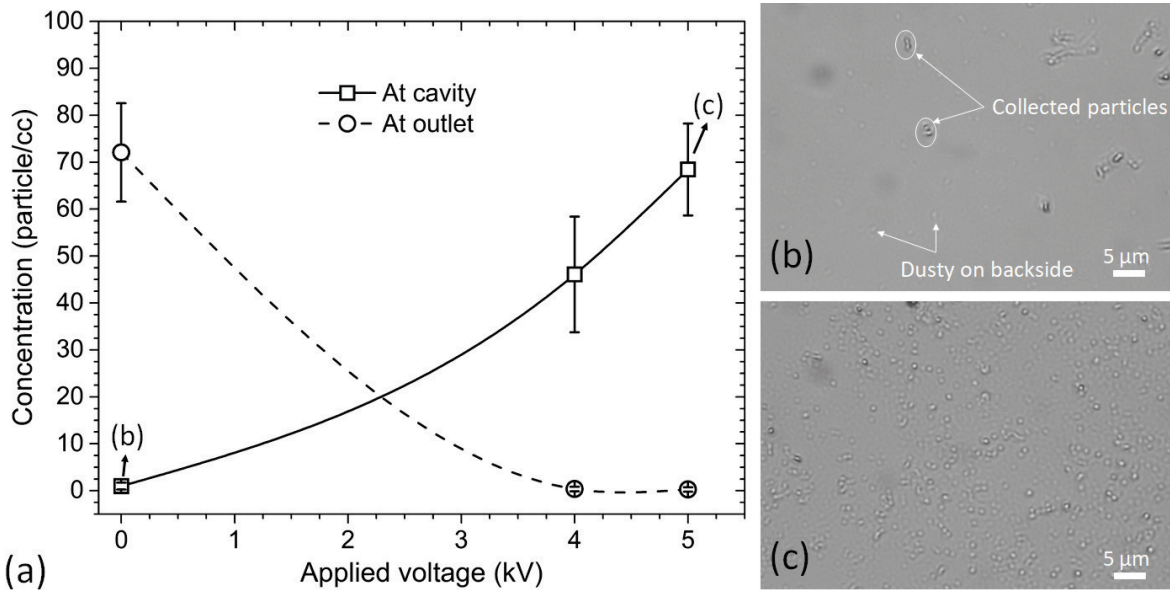
32 A range of flow rates of particles (2.0 lpm to 6.5 lpm) is used to investigate the efficiency of the present device
33 with different applied voltages. **In this range of flow rate, the initial charge of the aerosol is confirmed by an
34 electrometer (3068 TSI) to be less than -50 fA.** The densities of particles measured in the sampling chamber and
35 at the outlet with and without the effect of ion winds are plotted in Fig. 3 where the aerosol includes polymer
36 particles of 1.0 μm.

1 The results in Fig. 3 show that polymer particles almost move to the collecting tray under the effect of ion winds.
2 Indeed, the concentration of particle flow measured at the outlet without the ion wind is almost identical with
3 the concentration of particles measured in the sampling chamber when turning on the ion wind.

4 Theoretically, the particle charge increases with the increase of ion wind density and thus with the increase of
5 voltage [44]. In other words, the collection efficiency increases with the increase of voltage applied on the two
6 electrodes. Experimental results presented in Figs. 4 and 5 calculated from Eq. (1) are in good agreement with the
7 statements above and by Eq. (2) on the relationships of the collection efficiency of the present approach to the
8 flow rate of particles and the applied voltage, respectively.

9 Indeed, Fig. 4 shows that the collection efficiency of the system decreases with the increase of the aerosol flow
10 rate using an applied voltage of 5.5 kV. In addition, after reaching a maximum value of 94%, the efficiency starts
11 decreasing since the flow rate of 4 lpm. Meanwhile, the efficiency increases with the increase of applied voltage
12 with any flow rate of aerosol as shown in Fig. 5. It can be seen that with a relevant discharge voltage, the new
13 configuration is preferable to several ones published [24] whose efficiency is from 40% to 90% compared with
14 from 87% to 94% of the present system.

15 The research is extended with several samples included different particle sizes with a range of diameters from
16 300 nm to 1000 nm. For such experiments, aerosols include polymer particles of 1.0 μm and water particles
17 created by water atomizing using an aerosol generator (ATM 226). The distribution of different particle sizes at
18 the sampling chamber when turning on and off ion wind presented in the inset of Fig. 6 depicts that almost
19 particles of different sizes are drifted by the wind generated at a voltage 5.5 kV toward the sampling chamber,
20 meanwhile a small percentage of particles was found at the collector with the absence of ion wind.



21
22 *Figure 3. Experimental results of the present aerosol electrostatic sampler: Density of particles collected at the outlet of*
23 *the channel and the sampling chamber using polymer particles with diameter of 1.0 μm (Fig.3a). Images of collected*
24 *particles (small cycles) in the sampling chamber without the use of ion wind (Fig. 3b) and with ion wind (Fig. 3c) which*
25 *are observed using an optical microscopic.*

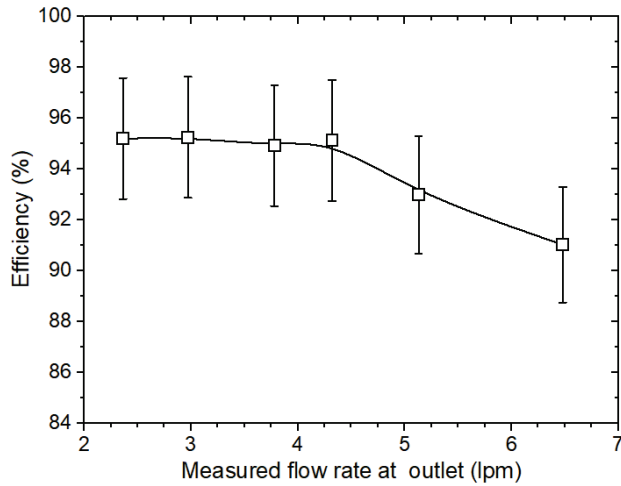


Figure 4. Experimental results of the present aerosol electrostatic sampler: The efficiency of the present device plotted versus the flow rate of particles with polymer diameter of $1.0\ \mu\text{m}$ using an applied voltage of $5.5\ \text{kV}$ and inter-electrode distance of $8.0\ \text{mm}$.

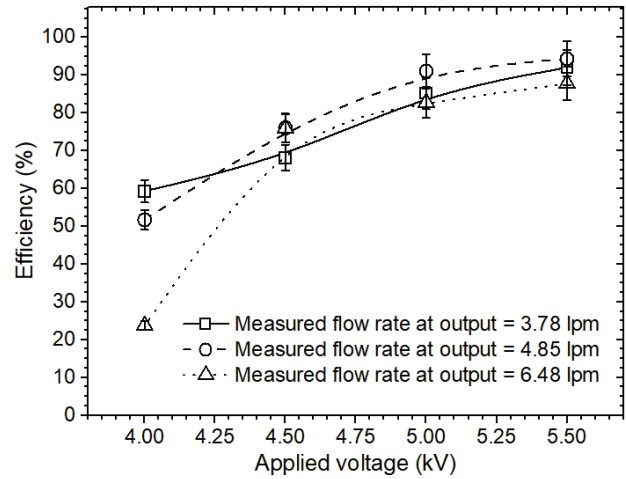


Figure 5. Experimental results of the present aerosol electrostatic sampler: The efficiency of the present device plotted versus the applied voltage for several different flow rates (3.78 , 4.85 and 6.48) lpm of particles with diameter of $1.0\ \mu\text{m}$ and using inter-electrode distance of $8.0\ \text{mm}$.

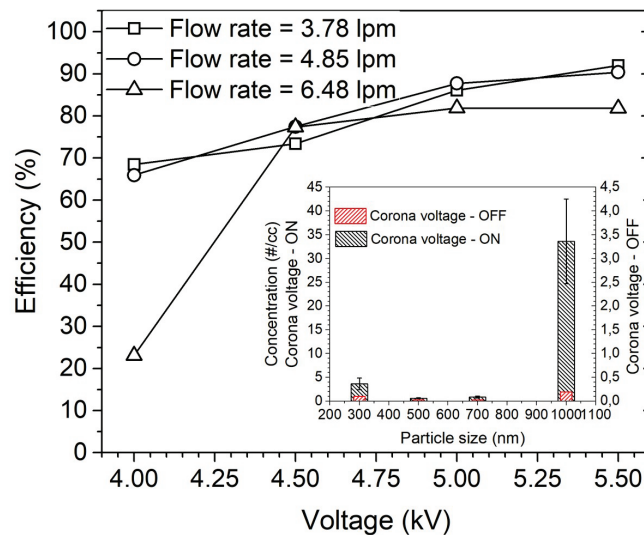


Figure 6. The efficiency of the present device plotted versus the applied voltage with several different flow rates (3.78 , 4.85 and 6.48) lpm of aerosol with particles of different diameters using inter-electrode gap of $8.0\ \text{mm}$.

The inset shows a distribution of particles collected in the sampling chamber using aerosol of different particle sizes with and without the ion wind. A flow rate of $3.78\ \text{lpm}$ included polymer particles of $1.0\ \mu\text{m}$ ($1000\ \text{nm}$) and smaller water particles is used. Inter-electrode distance is $8.0\ \text{mm}$

- 1 Similar to the aerosol of particles with one unit size, the collection efficiency determined by Eq. (1) for the whole
- 2 measured particles, increases with the increase of applied voltage with any flow rate of aerosol as shown in Fig.

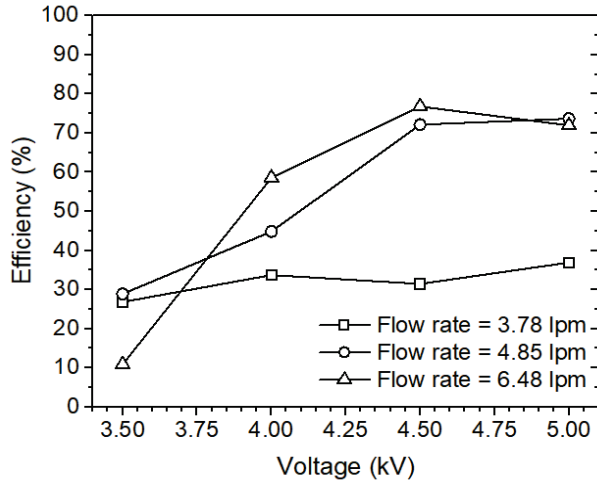


Figure 7a. The efficiency of the present device plotted versus the applied voltage with several different flow rates (3.78, 4.85 and 6.48) lpm of particles of $1.0\ \mu\text{m}$ using the inter-electrode distance of 5mm.

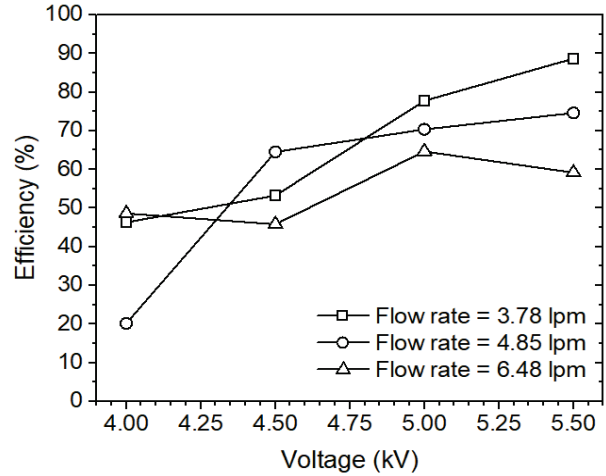


Figure 7b. The efficiency of the present device plotted versus the applied voltage with several different flow rates (3.78, 4.85 and 6.48) lpm of particles of $1.0\ \mu\text{m}$ using the inter-electrode distance of 10 mm.

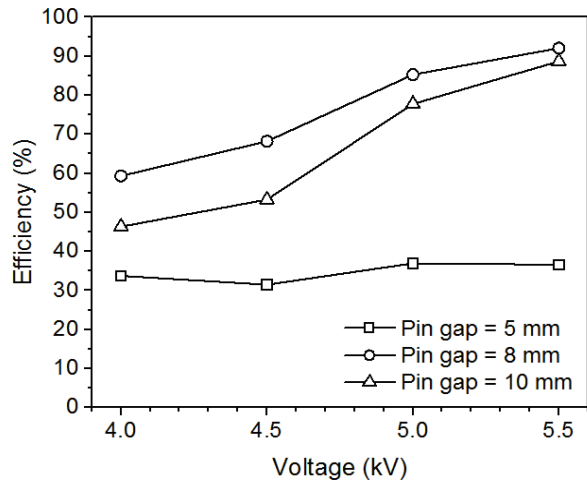


Figure 7c. The efficiency of the present device plotted versus the applied voltage with several different inter-electrode distances (5, 8, 10) mm and a flow rate of 3.78 lpm with particles of $1.0\ \mu\text{m}$.

1 6. In addition, the difference of collection efficiency is insignificant for different flow rates after the applied voltage
 2 overcame 5.0 kV and the efficiency reaches a maximum value at the applied voltage of around 5.5kV.

3 Besides applied voltage, the effect of inter-electrode distance on the efficiency of the present device is also
 4 considered for aerosol with polymer particles of $1.0\ \mu\text{m}$. Experimental results by Figs. 7a & b show that the effect
 5 of the inter-electrode distance on the efficiency of sampler depends on the flowrate of aerosol. For example, the
 6 efficiency is the highest with the inter-electrode distance of 5 mm (Fig. 7a) but the lowest with the distance of 10
 7 mm (Fig. 7b) using the same aerosol flowrate of 6.48 lpm. In other words, there is relevant flowrate for a designed
 8 inter-electrode distance of the present device. For example, the flow-rate of 3.78 lpm is relatively relevant for

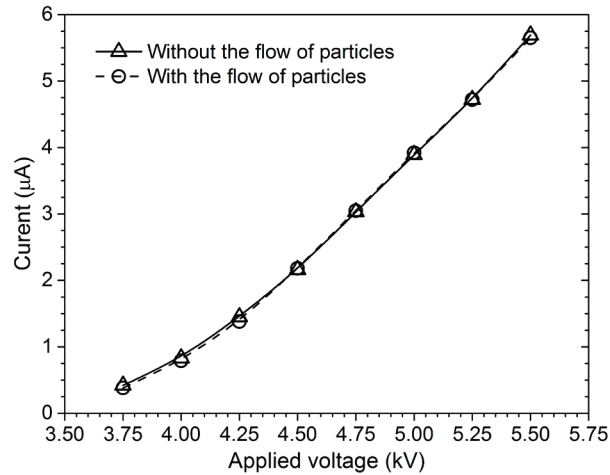


Figure 8. I-V characteristics of corona measured with and without the flow of particles during the discharge process.

1 the inter-electrode distance of 8 mm compared with the distances 5mm and 10mm because it yields a higher
 2 efficiency of the device as given in Fig. 7c.

3 Furthermore, the charge of the particle flow throughout the channel measured by an electrometer probe (3068
 4 TSI) installed in the device is approximately -10 fA, which is much lower than one of microamperes as presented
 5 recently in [37]. In other words, the airborne particles are nearly neutralised by the positive and negative ion
 6 winds.

7 For the present configuration of sampler, powers used for the corona discharge and the device operation are very
 8 low, around 20 mW and 70 mW, respectively. This allows the present device to be operated by a small battery,
 9 which is promising in the development of portable sampling devices. **Finally, the experiment also demonstrates**
 10 **that particles can be deposited on a sampling tray which is made entirely of glass. It means that insulating**
 11 **materials are applicable for collectors in the present system.**

12 4. Numerical simulation of the new configuration

13 The simulation of particle flow through the device under the effect of ion winds is presented in this section. The
 14 numerical results are then compared with experimental ones presented in section 3. A multi physic simulation is
 15 carried out to analyze the characteristics of the present configuration, in which several experimental results on
 16 corona discharge are introduced as the boundary conditions of the problem. In this work, we consider three issues:
 17 (i) the electro-hydrodynamics for the generation of ion winds; (ii) the charging process of particles caused by a
 18 strong electric field and (iii) the migration of charged particles drifted by the fluidic and electrostatic forces.

19 4.1. Governing equations and boundary conditions

20 For the electro-hydrodynamics, by neglecting the buoyancy force due to temperature variation, ion wind is
 21 considered as an incompressible turbulent flow of ions where the ionization region surrounding the electrode tips
 22 is modeled with a corresponding charge concentration. At the steady state, the migration of ions in the inter-
 23 electrode zone, the interaction of ions within the electric field and the charge consumption by the ion

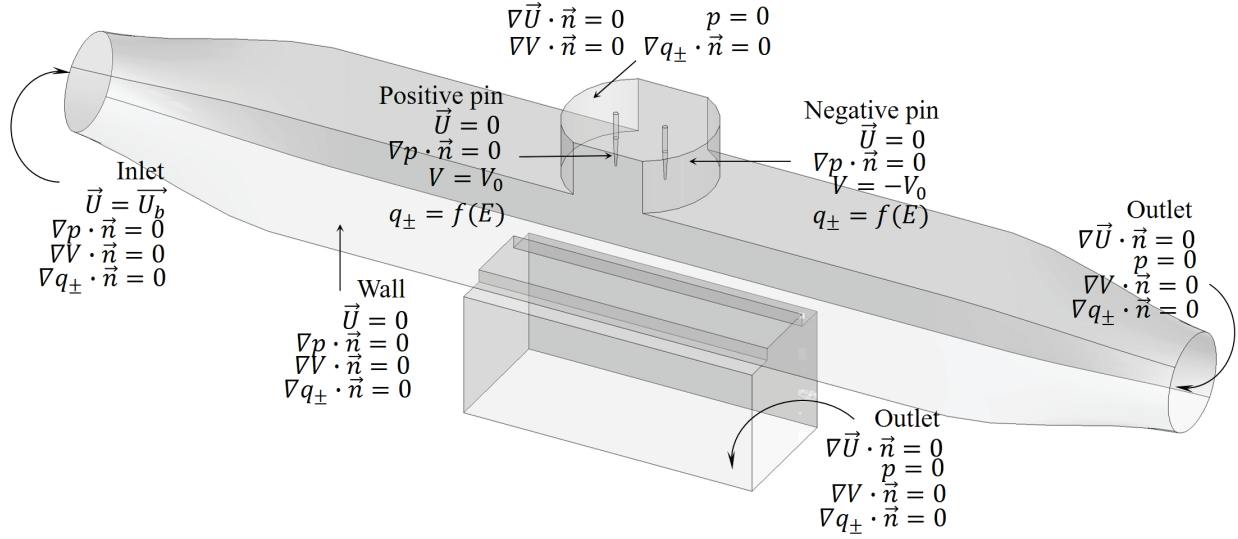


Figure 9. Simulation of the present sampler: considered domain and boundary conditions.

1 recombination process are governed by the Gauss' law of electrical field and the conservations of charge,
 2 momentum and mass as follows [45–47].

3
$$\vec{E} = \nabla\phi, \quad (3a)$$

4
$$\nabla \cdot (\pm\mu\vec{E}\rho_{\pm} \pm \vec{U}\rho_{\pm}) = -R_i\rho_+\rho_-/q_e, \quad (3b)$$

5
$$\nabla \cdot (\nabla\phi) = -(\rho_+ - \rho_-)/\epsilon_0, \quad (3c)$$

6
$$\nabla \cdot (\vec{U}\vec{U}) - \nabla \cdot (\nu\nabla\vec{U}) = -\nabla p + f_e/\rho_a, \quad (3d)$$

7
$$\nabla \cdot \vec{U} = 0, \quad (3e)$$

8
$$f_e = (\rho_+ - \rho_-)\vec{E}, \quad (3f)$$

9 where μ is the mobility of positive and negative charges ($\mu = 1.6 \times 10^{-4} m^2 \cdot V^{-1} \cdot s^{-1}$); $R_i = 10^{-13} m^3 \cdot s^{-1}$ the
 10 rate constant for ion-ion recombination; $q_e = 1.62 \times 10^{-19} C$ the charge of electrons; \vec{U} velocity of air drifted by the
 11 motion of charge; $\epsilon_0 = 8.854 \times 10^{-12} C \cdot V^{-1} \cdot m^{-1}$ the permittivity of the air; p the pressure; $\nu = 15.7 \times$
 12 $10^{-3} m^2 \cdot s^{-1}$ the kinematic viscosity, and $\rho_a = 1.2041 kg \cdot m^{-3}$ the air density.

13 The ion density of corona is assumed to be constant ($\rho_{\pm} = I/(\mu E_w a)$) on the face of electrodes, where a is the
 14 area of electrode tips, E_w the electric field at electrode tips determined by the Peek's law for air: $E_w =$
 15 $31(kV/cm)[1 + 0.308/(0.5R)^{1/2}]$, with R being the radius of electrode tips in cm and I the corona discharge
 16 current given by the I-V experimental characteristics in Fig. 8. **The current in Fig. 8 was measured at the negative**
 17 **electrode for cases with and without air flow and accordingly used as boundary conditions in simulation. Because**
 18 **both electrodes are connected to a single power source, which is battery operated, this current is considered to**
 19 **be equal at both positive and negative electrodes based on the simplified Kirchhoff's current law for charge**
 20 **conservation. In this work, the particle flow does not significantly affect the corona discharge because (i) there is**

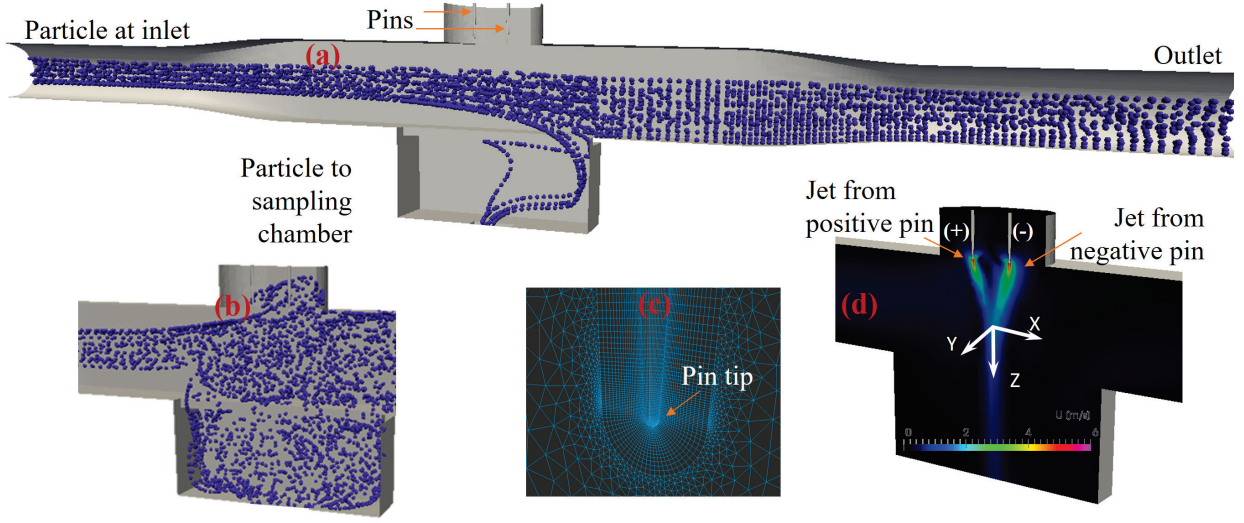


Figure 10. Simulation of the present device: Particle trajectories from the numerical simulation with ion wind is *off* (a), ion wind is *on* (b), meshing at the vicinity of pin tips (c) and ion wind from the dual electrodes (d).

1 no obstacle along the channel as presented in section 3 and (ii) the maximum velocity of particle flow is less than
2 1 m/s [48–51].

3 Neglecting the random motion [52], ions drift along the electric field lines to reach and charge particles. The
4 charging rate is determined by the following equation [53,54].

$$5 \quad \frac{dq}{dt} = \pi \epsilon_0 \rho_o \mu q_s \left(1 - \frac{q}{q_s}\right)^2, \quad (4)$$

6 where q_s the saturation charge of particles and given by $q_s = \left(1 + 2 \frac{K-1}{K+2}\right) r_p^2 E / \epsilon_0$, with r_p the radius of the
7 particle; K the dielectric constant of particles and ρ_o the ion concentration in the vicinity of particles.

8 Particles' locations and their velocities are tracked by the Newtonian second law as follows.

$$9 \quad m_p \frac{d\vec{U}_p}{dt} = \vec{f}_D + \vec{f}_E, \quad (5)$$

10 where m_p is the mass of a particle, f_D the drag force of air acting on the particle. f_E is the Coulomb's force of the
11 electric field since the particle is charged with q_p and given by

$$12 \quad \vec{f}_E = q_p \vec{E}_p, \quad (6)$$

13 with \vec{E}_p the electric field projected onto the particle position, meanwhile the drag force f_D is determined by

$$14 \quad \vec{f}_D = \frac{1}{2} C_D \rho_a A (\vec{U} - \vec{U}_p) |\vec{U} - \vec{U}_p|, \quad (7)$$

15 where C_D is the drag coefficient, ρ_a the air density, $\vec{U} - \vec{U}_p$ the difference between the air and particle velocities
16 at a position of the particle. In this simulation, particles are assumed to be sphere whose surface area is $A = \pi d^2$
17 with d the diameter of particle.

1 The drag coefficient (C_D) is then computed by the Shiller-Neumann correlation [55]

$$2 \quad C_D = \frac{24}{Re_p} (1 + 0.15 Re_p^{0.687}), \quad (8)$$

3 where Re_p is the Reynolds number $Re_p = \rho_a D |\vec{U} - \vec{U}_p| / \mu_g$ with ρ_a the viscosity of air.

4 At an instant, with an obtained velocity of particles, their position \vec{x}_p is determined straightforwardly by the
5 following kinematic equation

$$6 \quad \frac{d\vec{x}_p}{dt} = \vec{U}_p \quad (9)$$

7 Figure 9 represents the domain under consideration together with boundary conditions of the problem. A set of
8 400 particles with diameter of 1.0 μm is introduced into the system. The charge of particles is calculated at the
9 saturated state q_s . The velocity profile of particle flow, electric field lines, generated ion wind and particles'
10 trajectory are determined using OpenFOAM finite volume based solver [56,57].

11 4.2. Numerical results and discussion

12 With a meshing at the vicinity of electrode tips as described in Fig. 10c, the ion wind generated at the electrodes
13 is given in Fig. 10d. In addition, trajectories of particles in the device without and with the use of ion wind by 5.5
14 kV are represented in Figs. 10a&b, respectively. The numerical results show that there is approximately 62% of
15 particles reaching the sampling chamber with ion wind generated by a voltage of 5.5 kV compared with 13%
16 without ion wind.

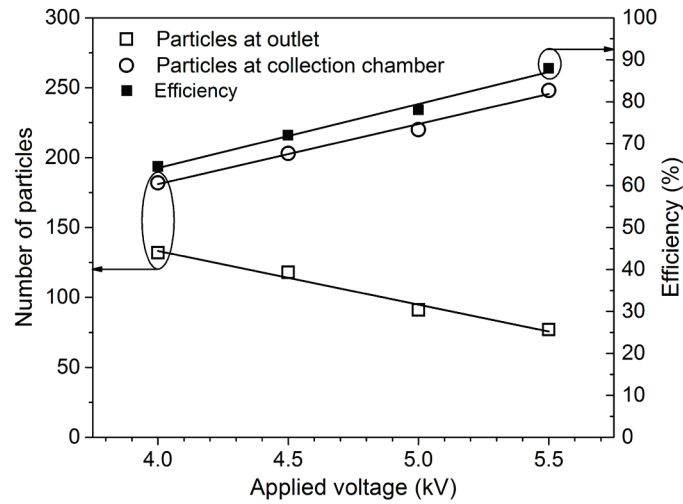
17 The simulation is repeated with a range of applied voltages $V = (4.0 - 5.5)$ kV using a flowrate of 5.0 lpm.
18 Numerical experiments show that the number of particles reaching the sampling chamber and thus, the efficiency
19 of the sampler increases with the increase of applied voltage as shown in Fig. 11. These simulated results are in
20 agreement with ones by the experiments using the same parameters except a small difference of the flowrate of
21 particle flow (4.86 lpm vs. 5.0 lpm) as presented in Fig. 5.

22 The linear relation between the efficiency and applied voltage with the flowrate of 5lpm observed by both
23 experimental and simulation works (Figs. 5 and 11) allows to predict the characteristics of the present aerosol
24 sampler. In addition, the percentage of lost particles by the present approach is significantly improved, it is 18%
25 compared with 26% by a point-to-plane system reported in [58]. **Nevertheless, the particle loss depends on the
26 design of an aerosol sampling system and hence, it is necessary to further consider for specific applications which
27 are the object of our ongoing work.**

28 Finally, the simulation is also extended with aerosol of particles with different sizes ranging 200 to 2000
29 nanometres using a flowrate of 3.78 lpm, applied voltage of 4.5 kV and inter-electrode distance of 8.0 mm. **At the
30 inlet, 400 particles are introduced into the system through a circular area whose diameter is 0.75 of that of the
31 inlet. The particles' size satisfies the Gauss distribution with the mean diameter of 1 μm and a standard deviation
32 3σ covering the range 200 nm – 2000nm. Since particles with different quantities are randomly released, while
33 moving inside the channel, random correlations of particle sizes and initial positions are determined. Thus, the
34 number of particles of a size is possibly lost more than others, and this yields a variation of the collection efficiency
35 for particles with different sizes.** Fig. 12 represents the particle distribution in the device, especially at the
36 collecting chamber with and without the use of ion winds. The simulation results show that the efficiency of the
37 collection vibrates from 50% to 73% for the particle sizes smaller than 1 μm . The efficiency then gradually
38 decreases with particle size bigger than 1 μm . This can be explained as the inertia effect on the electrostatic force

1 for particles having a significant mass. In addition, the numerical results by Fig. 12c show an agreement with the
 2 experimental one in Fig. 6 where the effect of ion wind was observed for a range of particle sizes

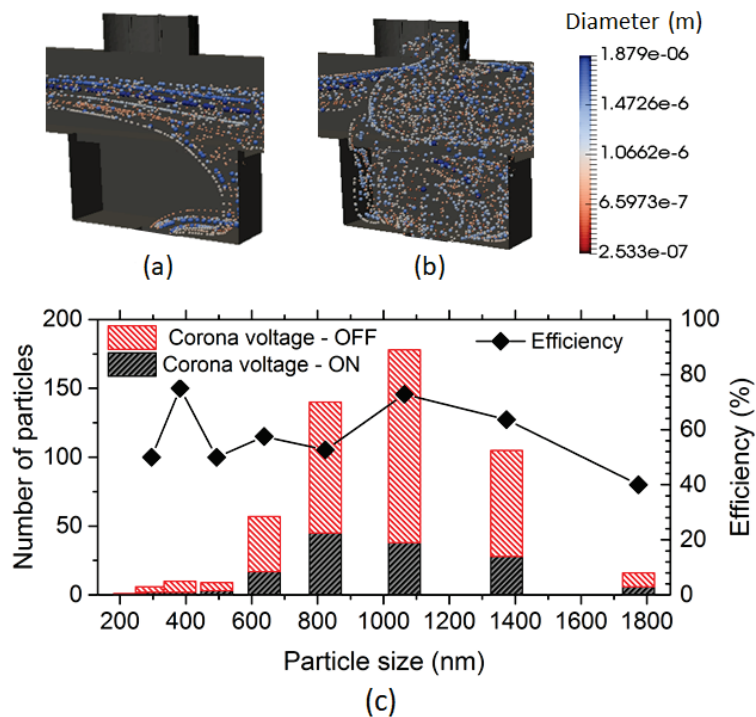
3



4

5 *Figure 11. Simulation of the sampler: Efficiency of the present sampler at different applied voltages with the flow rate*
 6 *of 5 lpm and the number of particles released at inlet = 400*

7



8

9

10 *Figure 12. Numerical simulation of the present device using aerosol of different particle sizes: Particle trajectories in the*
 11 *sampling chamber by the simulation without ion wind (a) and with ion wind (b); Particle distribution and efficiency using*
 12 *ion wind generated by an applied voltage of 4.5kV (c).*

1 Although the numerical simulation reasonably reflected the characteristics and behaviours of the device
2 compared with experimental results, it is necessary to consider and investigate several problems in our future
3 work, including the difference between negative and positive ions; the introduction of the diffusive charging of
4 particles which dominates the charging system for nano-particles; the coalescence and effect of shape for
5 submicron particles (less than 0.1 μ m) [59,60]; and the role of charging by electrostatic force, drifting by
6 aerodynamic force and the electro-coalescence in a simulation model [61,62] .

7 **5. Conclusion**

8 A new electrostatic sampling method based on dual ion winds has been developed to collect particles without
9 net charge. For the approach, a flow of airborne particles moves through a channel under the effect of dual ion
10 winds of opposing charges generated simultaneously by electrodes arranged in parallel. As results of the new
11 configuration, airborne particles can move smoothly through the device to reach the collection chamber because
12 there is not any impediment along the channel. In addition, particles reaching to sampler are almost neutralized
13 due to corona dual ion winds. Another advantage of the present approach is that the material of sampling chip
14 can be chosen arbitrarily including glass, gel, metal or semiconductor owing to the isolation of sampling chip from
15 ion winds generating circuit. The present aerosol sampling method is investigated by experimental work for a
16 range of discharge voltages (4.0 kV to 5.5 kV) and flow rates of airborne particles (2.0 to 6.5 lpm). A numerical
17 simulation is also carried out and numerical results support those by experiments. Although further investigation
18 is required, the present approach would yield a promising platform for a system of automated bio-aerosol
19 sampling.

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