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Citation	Physical Review B (2018), 98(22)
Issue Date	2018-12-19
URL	http://hdl.handle.net/2433/241783
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Type	Journal Article
Textversion	publisher

Nematic transition and highly two-dimensional superconductivity in $\text{BaTi}_2\text{Bi}_2\text{O}$ revealed by ^{209}Bi -nuclear magnetic resonance/nuclear quadrupole resonance measurements

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(Received 20 August 2018; revised manuscript received 6 November 2018; published 19 December 2018)

In this Rapid Communication, a set of ^{209}Bi -nuclear magnetic resonance (NMR)/nuclear quadrupole resonance (NQR) measurements has been performed to investigate the physical properties of superconducting (SC) $\text{BaTi}_2\text{Bi}_2\text{O}$ from a microscopic point of view. The NMR and NQR spectra at 5 K can be reproduced with a nonzero in-plane anisotropic parameter η , indicating the breaking of the in-plane fourfold symmetry at the Bi site without any magnetic order, i.e., “the electronic nematic state.” In the SC state, the nuclear spin-lattice relaxation rate divided by temperature, $1/T_1T$, does not change even below T_c , while a clear SC transition was observed with a diamagnetic signal. This observation can be attributed to the strong two dimensionality in $\text{BaTi}_2\text{Bi}_2\text{O}$. Comparing the NMR/NQR results among $\text{BaTi}_2Pn_2\text{O}$ ($Pn = \text{As}, \text{Sb}, \text{and Bi}$), it was found that the normal and SC properties of $\text{BaTi}_2\text{Bi}_2\text{O}$ were considerably different from those of $\text{BaTi}_2\text{Sb}_2\text{O}$ and $\text{BaTi}_2\text{As}_2\text{O}$, which might explain the two-dome structure of T_c in this system.

DOI: [10.1103/PhysRevB.98.220507](https://doi.org/10.1103/PhysRevB.98.220507)

An electronic nematic transition, which is characterized by spontaneous rotational symmetry breaking in the electronic system, is often observed in strongly coupled superconductors [1–5]. The nematic phase extends beyond the superconducting (SC) phase in the phase diagram of most superconductors. The relationship between superconductivity and the nematic state, and the origin of the nematic transition, especially in iron-based superconductors, are phenomena of intensive debate in the superconductivity community. However, the relationship and the origin remain elusive. Therefore, to understand these phenomena, it is important to investigate and compare various superconductors that exhibit a nematic transition.

Recently, superconductivity in the vicinity of a charge density wave (CDW) phase was discovered in $\text{BaTi}_2Pn_2\text{O}$ ($Pn = \text{As}, \text{Sb}, \text{and Bi}$) [6–9], which possesses a two-dimensional layered structure as shown in Fig. 1(a). $\text{BaTi}_2Pn_2\text{O}$ crystallizes into a tetragonal structure of space group $P4/mmm$ (No. 123, D_{4h}^1) with alternately stacked Ti_2Pn_2O layers and Ba atoms along the c axis. The Ti_2Pn_2O layers contain a Ti_2O square net, which is an anticonfiguration to the CuO_2 square net. The edge-shared TiO_2Pn_4 octahedra form a square lattice, and the electronic state of Ti^{3+} is in the $3d^1$ state, which is regarded as an electron-hole symmetric state of the $3d^9$ state in Cu^{2+} [6]. $\text{BaTi}_2\text{As}_2\text{O}$ shows an anomaly in the nematic state at $T_A = 200$ K, which is ascribed to a CDW transition [10]. This anomaly was suppressed by the substitution of Sb for As and T_A becomes ~ 40 K in the end member $\text{BaTi}_2\text{Sb}_2\text{O}$ [2]. In $\text{BaTi}_2\text{Sb}_2\text{O}$, an SC transition was also observed at an SC transition temperature of $T_c = 1.2$ K. $^{121/123}\text{Sb}$ -nuclear magnetic resonance (NMR)/nuclear quadrupole resonance (NQR)

measurements in $\text{BaTi}_2\text{Sb}_2\text{O}$ revealed the breaking of in-plane fourfold symmetry at the Sb site below T_A without an internal field appearing at the Sb site [2], which indicated an electronic nematic transition at T_A . The muon spin rotation (μSR) measurements also indicated no internal magnetic field below T_A [12], and a long-range structural phase transition was found in the neutron diffraction measurements [13]. These results can be understood with commensurate nematic CDW ordering [14] or p - d bond order [15] below T_A . In both cases, the nematic state is realized in the CDW phase. Moreover, $^{121/123}\text{Sb}$ -NMR/NQR measurements strongly suggest that SC symmetry is a conventional s wave in $\text{BaTi}_2\text{Sb}_2\text{O}$ [2], which is in sharp contrast with those in the cuprate and iron-based superconductors.

With the substitution of Bi for Sb, T_A was suppressed, T_c shows a dome shape with maximum $T_c = 3.5$ K at $x = 0.2$, and superconductivity terminates at $x = 0.4$ in $\text{BaTi}_2(\text{Sb}_{1-x}\text{Bi}_x)_2\text{O}$ [7]. Interestingly, superconductivity reappears upon further substitution, and a two-dome structure in T_c was observed as shown in Fig. 1(b). In the end member $\text{BaTi}_2\text{Bi}_2\text{O}$, an SC transition was observed at 4.6 K without any trace of the CDW/spin density wave (SDW) transition. Since there is a possibility that the SC properties of $\text{BaTi}_2\text{Bi}_2\text{O}$ are different from those of $\text{BaTi}_2\text{Sb}_2\text{O}$, it is important to investigate $\text{BaTi}_2\text{Bi}_2\text{O}$ to understand the normal and SC properties of the $\text{BaTi}_2Pn_2\text{O}$ system.

In this Rapid Communication, ^{209}Bi -NMR/NQR measurements have been performed to investigate the physical properties of $\text{BaTi}_2\text{Bi}_2\text{O}$ from a microscopic point of view. From the temperature evolution of the NQR spectra, an electronic nematic transition at ~ 45 K was discovered. In the SC state, the nuclear spin-lattice relaxation rate divided by temperature, $1/T_1T$, does not change even below T_c , while a clear SC

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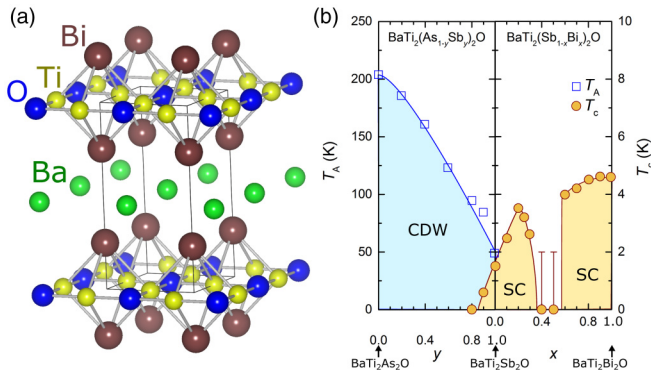


FIG. 1. (a) Crystal structure of $\text{BaTi}_2\text{Pn}_2\text{O}$. The figure of the crystal structure was modeled on the three-dimensional visualization program VESTA [11]. (b) Electronic phase diagram of $\text{BaTi}_2(\text{As}_{1-y}\text{Sb}_y)_2\text{O}$ and $\text{BaTi}_2(\text{Sb}_{1-x}\text{Bi}_x)_2\text{O}$ [7]. SC denotes the superconducting phases. The open squares represent the CDW transition temperature T_A . The solid circles represent the superconducting temperature T_c determined by magnetization and electrical resistivity measurements.

transition was observed by an ac susceptibility measurement. This is ascribed to the strong two dimensionality in $\text{BaTi}_2\text{Bi}_2\text{O}$. Although the NQR spectra in both $\text{BaTi}_2\text{Sb}_2\text{O}$ and $\text{BaTi}_2\text{Bi}_2\text{O}$ indicate the breaking of fourfold symmetry at the Sb/Bi site below ~ 40 K, the normal and SC properties of $\text{BaTi}_2\text{Bi}_2\text{O}$ are considerably different from those of $\text{BaTi}_2\text{Sb}_2\text{O}$.

Polycrystalline samples of $\text{BaTi}_2\text{Bi}_2\text{O}$ were synthesized via the conventional solid-state reaction [7]. Stoichiometric amounts of BaO, Ti, and Bi were mixed and pelletized. The pellet was then wrapped in a Ta foil and sealed in a quartz tube. The reaction temperature was 850°C . To prevent sample degradation by air and/or moisture, the polycrystalline samples were mixed with Araldite adhesive. The mixture was then solidified with random crystal orientation. All the procedures were performed in a glove box filled with Ar. The SC transition at 4.6 K was confirmed by a dc magnetization measurement with a commercial superconducting quantum interference device (SQUID) magnetometer (Quantum Design, MPMS3), and the ac susceptibility measurement was performed using an NMR coil after the mixing. A spin-echo technique was used for the NMR/NQR measurements. The ^{209}Bi -NMR spectra (nuclear spin $I = 9/2$, nuclear gyromagnetic ratio $^{209}\gamma/2\pi = 6.842$ MHz/T, and natural abundance 100%) were obtained as a function of magnetic field in a fixed frequency $f = 68.41$ MHz (~ 10 T). The ^{209}Bi nuclear spin-lattice relaxation rate $1/T_1$ was determined by fitting the time variation of the spin-echo intensity after the saturation of the nuclear magnetization to a theoretical function for $I = 9/2$.

The ^{209}Bi -NQR measurements indicate an electronic nematic state at low temperatures. Figure 2 shows the ^{209}Bi -NQR spectra that were obtained by the frequency-swept method at 5 K. When $I \geq 1$, a nucleus has an electric quadrupole moment eQ as well as a magnetic dipole moment; thus, the degeneracy of nuclear energy levels is lifted even at zero magnetic field due to the interaction between eQ and the

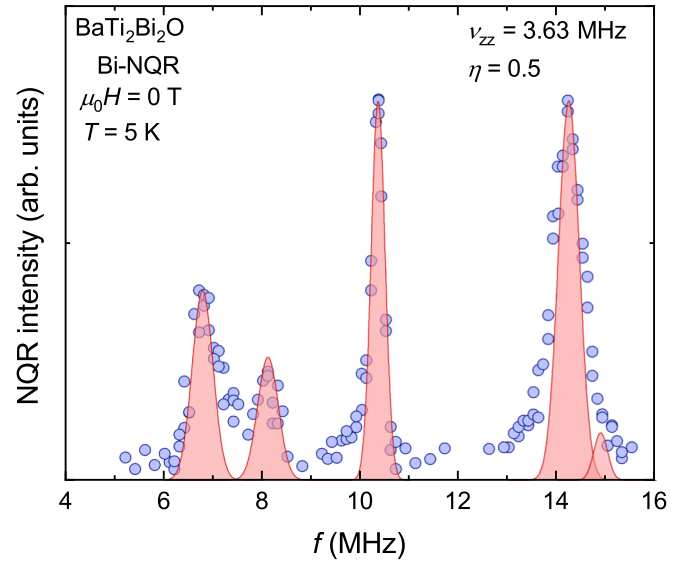


FIG. 2. ^{209}Bi -NQR spectra obtained by the frequency-swept method at 5 K. From the observed ^{209}Bi -NQR spectra, the quadrupole parameters for Bi nuclei are evaluated as shown in the figure. The solid red curves are the simulation of the NQR spectra using the estimated quadrupole parameters.

electric field gradient (EFG). This interaction is described as

$$\mathcal{H}_Q = \frac{h\nu_{zz}}{6} \left\{ (3I_z^2 - I^2) + \frac{1}{2}\eta(I_+^2 + I_-^2) \right\}, \quad (1)$$

where h is the Planck's constant, ν_{zz} is the quadrupole frequency along the principal axis (c axis) of the EFG and is defined as $\nu_{zz} \equiv 3e^2qQ/2I(2I-1)$ with $eq = V_{zz}$, and η is an asymmetry parameter of the EFG expressed as $(V_{xx} - V_{yy})/V_{zz}$ with $V_{\alpha\alpha}$, which is the second derivative of the electric potential V and the EFG along the α direction ($\alpha = x, y, z$). When ^{209}Bi is in the EFG, the degenerate ten nuclear-spin states are split into five energy levels, yielding four (or more) resonance frequencies as shown in Fig. 2. The NQR parameters $\nu_{zz} = 3.63$ MHz and $\eta = 0.5$ were obtained by comparing the observed ^{209}Bi -NQR spectra and calculated resonance frequencies obtained from the diagonalization of Eq. (1). A nonzero η implies the breaking of fourfold symmetry at the Bi site, which has a C_4 symmetry in the tetragonal structure (also called the “electronic nematic state”). It was observed that the low-intensity peak at the highest frequency corresponded to the $\nu_1 + \nu_2$, which is caused by the formation of hybrid states due to nonzero η [16].

The nonzero η was also confirmed by the field-swept NMR spectrum as shown in Fig. 3. For the NMR measurements, although the nuclear energy levels were already split by the electric quadrupole interaction, magnetic fields were still applied to lift the degeneracy of the spin degrees of freedom. The total effective Hamiltonian could then be expressed as

$$\mathcal{H} = \mathcal{H}_Z + \mathcal{H}_Q = -\frac{\gamma}{2\pi} h(1+K)\mathbf{I} \cdot \mathbf{H} + \mathcal{H}_Q, \quad (2)$$

where K is the Knight shift and \mathbf{H} is an external field. As the NMR spectra are varied against the angle between the principal axis of the EFG and magnetic field direction, the sum of the spectrum for all the corresponding angles

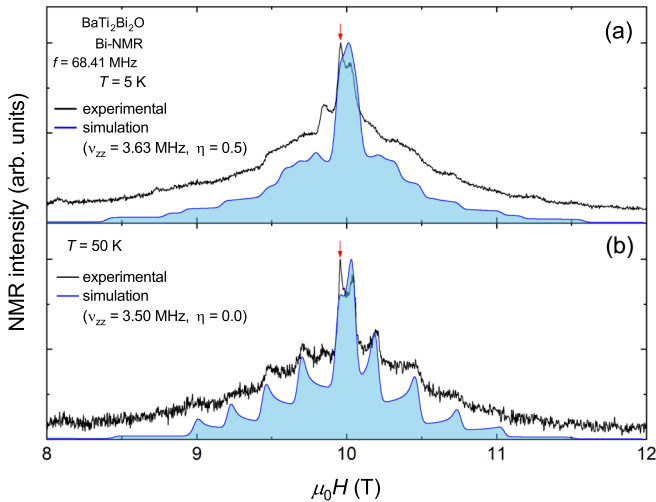


FIG. 3. Field-swept ^{209}Bi -NMR spectra at (a) 5 K and (b) 50 K at 68.41 MHz. In addition, simulations of the NMR spectra are shown. At 5 K, the same NQR parameters as those for the NQR measurements are used for the simulation. The arrows indicate the magnetic field at which T_1 is measured.

is observed in the case of the powder samples. The NMR spectrum at 5 K was consistently reproduced by the NQR parameters determined through the NQR measurements at 5 K. In contrast, the NMR spectrum at 50 K can be fitted by the simulation with $\nu_{zz} = 3.50$ MHz and $\eta = 0.0$. The small difference between experimental and simulated values might have originated from the impurity phase and/or the degree of orientation. It should be noted that the double-horn-shaped satellite signals are the characteristic feature of the $\eta = 0.0$. This reflects the preservation of the fourfold symmetry of the crystal structure. To estimate the transition temperature, the temperature evolution of the NQR spectra was measured. Figure 4 shows the temperature dependence

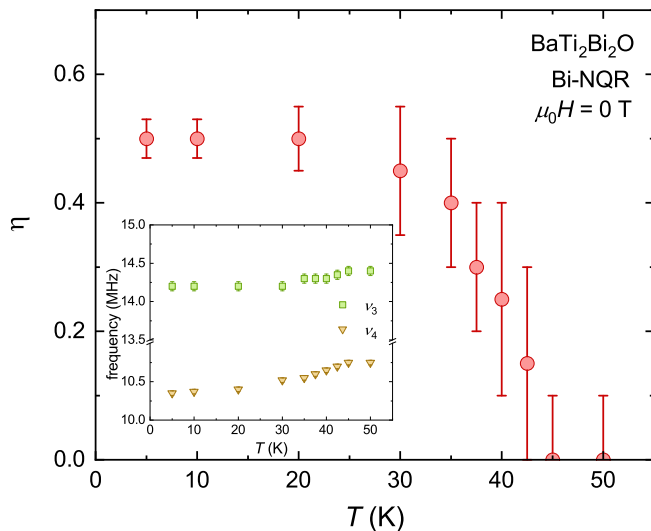


FIG. 4. Temperature dependence of η deduced from the first and second largest NQR peaks (~ 10 and 14 MHz). Inset: Temperature dependence of NQR frequency.

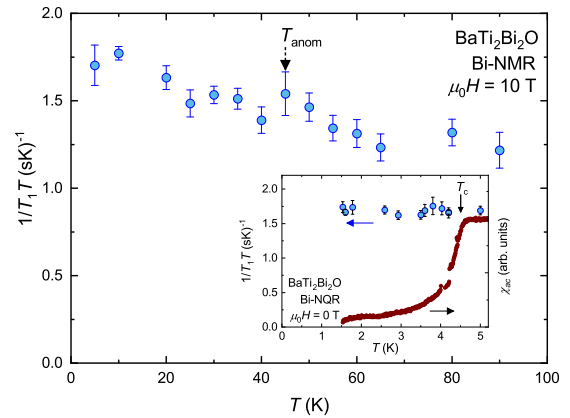


FIG. 5. Temperature dependence of $1/T_1T$ at ~ 10 T. The dashed arrow indicates the temperature at which the anomaly was observed in the NMR spectrum T_{anom} . Inset: Temperature dependence of $1/T_1T$ and ac susceptibility at 0 T. The solid arrow indicates the SC transition temperature T_c .

of η deduced from the first and second largest NQR peaks (~ 10 and 14 MHz). The figure indicates that η assumes a nonzero value below $T_{\text{anom}} \sim 45$ K. As such, a nematic phase transition was not reported in previous experiments [8]. The most plausible reason behind this is that the anomalous transition may have been so small that it could be detected only by a highly sensitive probe in an electric environment, such as the NQR measurement. A similar nematic transition was observed at ~ 40 K in $\text{BaTi}_2\text{Sb}_2\text{O}$. To reveal the origin of the nematic transitions and the mechanism by which the nematic phases cover the two SC domes in $\text{BaTi}_2(\text{Sb}_{1-x}\text{Bi}_x)_2\text{O}$, it is important to first understand the relationship between superconductivity and the electronic nematic state.

Figure 5 shows the temperature dependence of $1/T_1T$ at ~ 10 T. $1/T_1T$ was measured at the center peak of the Bi-NMR spectrum as indicated by arrows in Fig. 3. $1/T_1T$ slightly increases upon cooling and no measurable anomaly was observed around T_{anom} . To investigate the SC properties, the temperature dependence of $1/T_1T$ at the NQR signal (14.2 MHz peak) as shown in the inset of Fig. 5 was also measured. While a clear SC transition was observed at 4.6 K by the ac susceptibility measurement, $1/T_1T$ did not change even below T_c .

First, we discuss the unaltered behavior of $1/T_1T$ in the SC state. In conventional superconductors, $1/T_1T$ decreases at temperatures well below T_c due to the reduction of the quasiparticle density of states around the Fermi energy E_F by opening the SC energy gap. There are two possible reasons why $1/T_1T$ did not decrease in the SC state. One possibility is that spin diffusion due to low dimensionality prevented the reduction of $1/T_1T$ [17]. A similar absence of a clear reduction of $1/T_1T$ was observed in $(\text{La}_{0.87}\text{Ca}_{0.13})\text{FePO}$ [18] and certain cuprate superconductors [19,20]. This brings us to the other possibility that the coupling between the Ti_2O layer and Bi atom was so weak that the effect of the SC transition could not be detected even by ^{209}Bi -NQR measurements. Such a scenario was discovered at the Cu(1) site in $\text{YBa}_2\text{Cu}_3\text{O}_{7-\delta}$ and $\text{YBa}_2\text{Cu}_4\text{O}_{8+\delta}$ [21,22]. $1/T_1T$'s at the first and second largest NQR peaks (10.35 and 14.2 MHz)

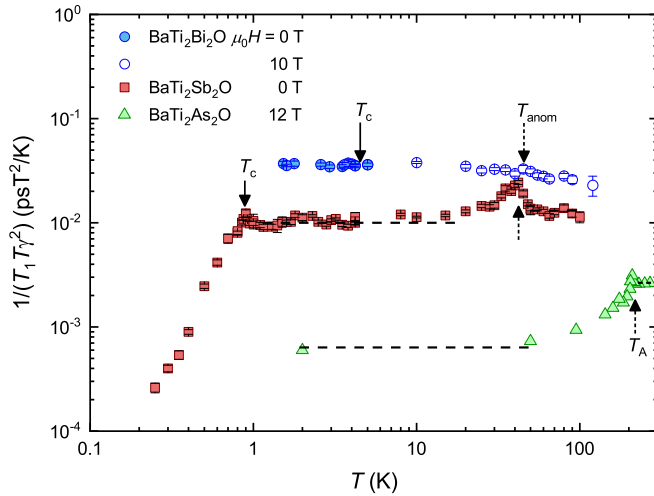


FIG. 6. Temperature dependence of $1/T_1 T 1/(T_1 T \gamma^2)$ in $\text{BaTi}_2\text{Bi}_2\text{O}$, $\text{BaTi}_2\text{Sb}_2\text{O}$ [2], and $\text{BaTi}_2\text{As}_2\text{O}$ [15]. For $\text{BaTi}_2\text{Bi}_2\text{O}$, we plot the data at 10 T above 5 K and at 0 T below 5 K. For $\text{BaTi}_2\text{As}_2\text{O}$, we use the angle average of $1/T_1 T$ for simplicity. The dashed arrows indicate T_{anom} or T_A . The solid arrows indicate T_c .

are not different, although the frequency dependency is expected in the spin diffusion scenario [17]. In both cases, low dimensionality is important for such an anomalous behavior to occur; therefore, it can be concluded that the NMR results indicate that $\text{BaTi}_2\text{Bi}_2\text{O}$ is a two-dimensional superconductor, and T_c in this system is enhanced by low dimensionality [23]. Such two dimensionality was not predicted by earlier band calculations [24,25], which imply three-dimensional Fermi surfaces. It is plausible that this discrepancy is related to the electronic nematic transition. The electronic state at low temperatures differs from that at room temperature due to the nematic transition at ~ 45 K.

Next, we compare the NMR/NQR results in $\text{BaTi}_2\text{Bi}_2\text{O}$ with those in $\text{BaTi}_2\text{Sb}_2\text{O}$ and $\text{BaTi}_2\text{As}_2\text{O}$. In general, $1/T_1 T$ can be described as

$$\frac{1}{T_1 T} = \frac{2\gamma_N^2 k_B}{(\gamma_e^2 \hbar)^2} \lim_{\omega \rightarrow 0} \sum_q A_q A_{-q} \frac{\chi''_{\perp}(\mathbf{q}, \omega_0)}{\omega_0}, \quad (3)$$

where γ_N (γ_e) is the nuclear (electronic) gyromagnetic ratio, k_B is the Boltzmann's constant, \hbar is the reduced Planck's constant, A_q is the Fourier transform of the hyperfine coupling, and $\chi''_{\perp}(\mathbf{q}, \omega)$ is the transverse component of the imaginary part of the dynamical susceptibility. Then, the value of $1/T_1 T$ in different nuclei may be compared after the normalization by γ_N^2 . Figure 6 shows the temperature dependence of $1/(T_1 T \gamma_N^2)$ in $\text{BaTi}_2\text{Bi}_2\text{O}$, $\text{BaTi}_2\text{Sb}_2\text{O}$, and $\text{BaTi}_2\text{As}_2\text{O}$ [15]. While the NQR/NMR spectra in those compounds indicate the breaking of the fourfold symmetry at the Pn site, the temperature dependences of $1/T_1 T$ are considerably different from each

other. In $\text{BaTi}_2\text{Sb}_2\text{O}$ and $\text{BaTi}_2\text{As}_2\text{O}$, $1/T_1 T$ was clearly enhanced toward T_A . The ratio of $1/T_1$ between the two isotopes of the Sb nuclei is close to the ratio of the square of the nuclear gyromagnetic ratio γ_N , suggesting that this enhancement originates from the magnetic nature [2]. In addition, a reduction of the constant value of $1/T_1 T$ below T_A was observed, indicating a decrease in the density of states due to the nematic transition. The value of $1/(T_1 T \gamma_N^2)$ in $\text{BaTi}_2\text{As}_2\text{O}$ is one order of magnitude smaller than those in $\text{BaTi}_2\text{Sb}_2\text{O}$ and $\text{BaTi}_2\text{Bi}_2\text{O}$, possibly due to the small hyperfine coupling constant. In contrast, $1/T_1 T$ shows no measurable anomaly around T_{anom} in $\text{BaTi}_2\text{Bi}_2\text{O}$. While the system does exhibit strong two dimensionality, the value of $1/T_1 T$ in $\text{BaTi}_2\text{Bi}_2\text{O}$ is larger than those in $\text{BaTi}_2\text{Sb}_2\text{O}$ and $\text{BaTi}_2\text{As}_2\text{O}$. This observation can be attributed to the existence of strong coupling between the atoms of Bi and Ba and the lack thereof between the Bi and Ti_2O layers. From this, it can hence be inferred that the BaBi_2 layer in $\text{BaTi}_2\text{Bi}_2\text{O}$ plays the role of a block layer whose electronic state is different from that of the Ti_2O plane.

Furthermore, a clear coherence peak immediately below T_c and an exponential decay at low temperatures as evidence of an s -wave superconductivity were observed in $\text{BaTi}_2\text{Sb}_2\text{O}$, while no anomaly was observed in $\text{BaTi}_2\text{Bi}_2\text{O}$. These differences may have arisen from the difference of the dimensionality between two compounds. Since $\text{BaTi}_2\text{Sb}_2\text{O}$ is three dimensional, the anomaly in the nematic transition and superconductivity can be detected by $^{121/123}\text{Sb}$ -NMR/NQR. However, such an anomaly was not observed by ^{209}Bi -NMR/NQR in $\text{BaTi}_2\text{Bi}_2\text{O}$ because of its strong two dimensionality. It is apparent that the SC Ti_2O layer is sandwiched by the nonsuperconducting BaBi_2 block layer with the different electronic state; thus, $\text{BaTi}_2\text{Bi}_2\text{O}$ is regarded as a two-dimensional superconductor. Therefore, it is important that the validation of two-dimensional superconductivity is performed using $^{47/49}\text{Ti}$ or ^{17}O NMR measurements.

In conclusion, the ^{209}Bi -NMR/NQR measurements were performed to investigate the physical properties of superconducting $\text{BaTi}_2\text{Bi}_2\text{O}$ from a microscopic point of view. The temperature evolution of the NQR spectra indicates an electronic nematic order at $T_{\text{anom}} \sim 45$ K. Comparing the NMR/NQR results among $\text{BaTi}_2 Pn_2\text{O}$ ($Pn = \text{As}, \text{Sb}, \text{and Bi}$), it was found that the normal and SC properties of $\text{BaTi}_2\text{Bi}_2\text{O}$ were somewhat different from those of $\text{BaTi}_2\text{Sb}_2\text{O}$ and $\text{BaTi}_2\text{As}_2\text{O}$. The observed two-dome structure in T_c on $\text{BaTi}_2(\text{Sb}_{1-x}\text{Bi}_x)_2\text{O}$ may have originated from these differences.

The authors acknowledge S. Yonezawa, Y. Maeno, and Y. Matsuda for fruitful discussions. This work was partially supported by the Kyoto University LTM Center and Grant-in-Aids for Scientific Research (KAKENHI) (Grants No. JP15H05882, No. JP15H05884, No. JP15K21732, No. JP15H05745, No. JP15K17698, and No. JP17K14339).

[1] R. Okazaki, T. Shibauchi, H. J. Shi, Y. Haga, T. D. Matsuda, E. Yamamoto, Y. Onuki, H. Ikeda, and Y. Matsuda, *Science* **331**, 439 (2011).

[2] S. Kitagawa, K. Ishida, K. Nakano, T. Yajima, and H. Kageyama, *Phys. Rev. B* **87**, 060510(R) (2013).

- [3] S.-H. Baek, D. V. Efremov, J. M. Ok, J. S. Kim, J. van den Brink, and B. Büchner, *Nat. Mater.* **14**, 210 (2014).
- [4] S. Hosoi, K. Matsuura, K. Ishida, H. Wang, Y. Mizukami, T. Watashige, S. Kasahara, Y. Matsuda, and T. Shibauchi, *Proc. Natl. Acad. Sci. U.S.A.* **113**, 8139 (2016).
- [5] Y. Sato, S. Kasahara, H. Murayama, Y. Kasahara, E.-G. Moon, T. Nishizaki, T. Loew, B. K. J. Porras, T. Shibauchi, and Y. Matsuda, *Nat. Phys.* **13**, 1074 (2017).
- [6] T. Yajima, K. Nakano, F. Takeiri, T. Ono, Y. Hosokoshi, Y. Matsushita, J. Hester, Y. Kobayashi, and H. Kageyama, *J. Phys. Soc. Jpn.* **81**, 103706 (2012).
- [7] T. Yajima, K. Nakano, F. Takeiri, Y. Nozaki, Y. Kobayashi, and H. Kageyama, *J. Phys. Soc. Jpn.* **82**, 033705 (2013).
- [8] T. Yajima, K. Nakano, F. Takeiri, J. Hester, T. Yamamoto, Y. Kobayashi, N. Tsuji, J. Kim, A. Fujiwara, and H. Kageyama, *J. Phys. Soc. Jpn.* **82**, 013703 (2013).
- [9] T. Yajima, *Condens. Matter* **2**, 4 (2017).
- [10] X. F. Wang, Y. J. Yan, J. J. Ying, Q. J. Li, M. Zhang, N. Xu, and X. H. Chen, *J. Phys.: Condens. Matter* **22**, 075702 (2010).
- [11] K. Momma and F. Izumi, *J. Appl. Crystallogr.* **44**, 1272 (2011).
- [12] Y. Nozaki, K. Nakano, T. Yajima, H. Kageyama, B. Frandsen, L. Liu, S. Cheung, T. Goko, Y. J. Uemura, T. S. J. Munsie, T. Medina, G. M. Luke, J. Munevar, D. Nishio-Hamane, and C. M. Brown, *Phys. Rev. B* **88**, 214506 (2013).
- [13] B. A. Frandsen, E. S. Bozin, H. Hu, Y. Zhu, Y. Nozaki, H. Kageyama, Y. J. Uemura, W.-G. Yin, and S. J. L. Billinge, *Nat. Commun.* **5**, 5761 (2014).
- [14] H. Nakaoka, Y. Yamakawa, and H. Kontani, *Phys. Rev. B* **93**, 245122 (2016).
- [15] D. W. Song, J. Li, D. Zhao, L. K. Ma, L. X. Zheng, S. J. Li, L. P. Nie, X. G. Luo, Z. P. Yin, T. Wu, and X. H. Chen, *arXiv:1806.11272*.
- [16] K. Karube, T. Hattori, Y. Ihara, Y. Nakai, K. Ishida, N. Tamura, K. Deguchi, N. K. Sato, and H. Harima, *J. Phys. Soc. Jpn.* **80**, 064711 (2011).
- [17] M.-H. Julien, *J. Phys. Soc. Jpn.* **77**, 125002 (2008).
- [18] Y. Nakai, K. Ishida, Y. Kamihara, M. Hirano, and H. Hosono, *Phys. Rev. Lett.* **101**, 077006 (2008).
- [19] S. Kambe, H. Yasuoka, A. Hayashi, and Y. Ueda, *Phys. Rev. Lett.* **73**, 197 (1994).
- [20] M.-H. Julien, *Physica B* **329-333**, 693 (2003).
- [21] M. Mali, D. Brinkmann, L. Pauli, J. Roos, H. Zimmermann, and J. Hulliger, *Phys. Lett. A* **124**, 112 (1987).
- [22] H. Zimmermann, M. Mali, D. Brinkmann, J. Karpinski, E. Kaldis, and S. Rusiecki, *Physica C* **159**, 681 (1989).
- [23] Y. Yanase, T. Jujo, T. Nomura, H. Ikeda, T. Hotta, and K. Yamada, *Phys. Rep.* **387**, 1 (2003).
- [24] D. V. Suetin and A. L. Ivanovskii, *JETP Lett.* **97**, 220 (2013).
- [25] K. Nakano, K. Hongo, and R. Maezono, *Sci. Rep.* **6**, 29661 (2016).