

## Heat flux calculation using embedded thermocouple in W-coated CFC tiles in the JET tokamak

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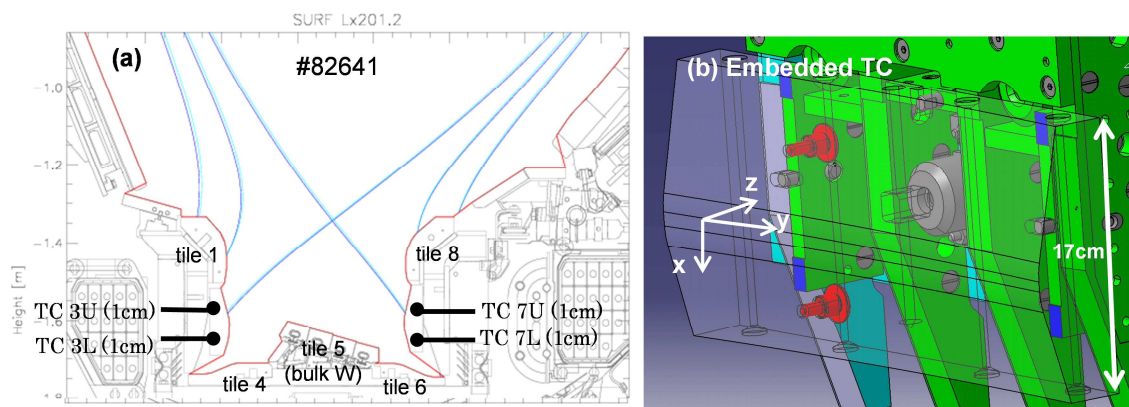
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**1. Introduction.** Surface temperature and heat flux measurements are important issues in high power fusion devices to guarantee safe plasma operation. In the JET tokamak, few embedded thermocouples (TC) located 1 cm below the tile surface are used to measure the bulk temperatures of the CFC tiles (coated with about 20  $\mu\text{m}$  of tungsten with the ITER-like wall). We propose here to use an inverse thermal calculation based on Quadruple thermal method to locally deduce the surface temperature (thus on the W coating) and deposited heat flux. The calculation requires the position of the peak (provided by offline magnetic reconstruction) and the qualitative 1D-shape of the heat flux deposited on the target (determined with an empirical expression of the target heat load profile found in ASDEX Upgrade and JET).

**2. Description of the JET experiment and diagnostic set-up.** The inverse heat flux calculation has been applied to 2.2T/2MA ELMy H-mode discharge #82641 performed with NBI heating (10MW during 4s) and with the high-field and low-field sides strike points on the vertical targets (low triangularity vertical tile configuration, see figure 1-a). The diagnostic set-up is made of infrared (IR) thermography system (wide angle equatorial viewing with spatial resolution of  $\sim 16\text{mm}$  in the divertor region) to measure the surface temperature and four embedded type-K thermocouples (TC) 10 mm below the tile surface to measure the bulk temperatures of the inner and outer vertical target plates (two TCs by tile as shown in figure 1-b, the upper TC is located at  $Z_{\text{TC}}^{\text{U}} = -1.557$  m and the lower one at  $Z_{\text{TC}}^{\text{L}} = -1.65$  m from the equatorial mid-plane). For #82641 the inner and outer strike point (ISP/OSP) positions are given by the magnetic equilibrium with typical uncertainties of  $\pm 0.5\text{cm}$  on the vertical tiles  $Z^{\text{ISP}} = -1.568\text{m}$  ( $\sim 1.2$  cm bellow TC3U) and  $Z^{\text{OSP}} = -1.565\text{m}$  ( $\sim 0.8$  cm bellow TC7U).

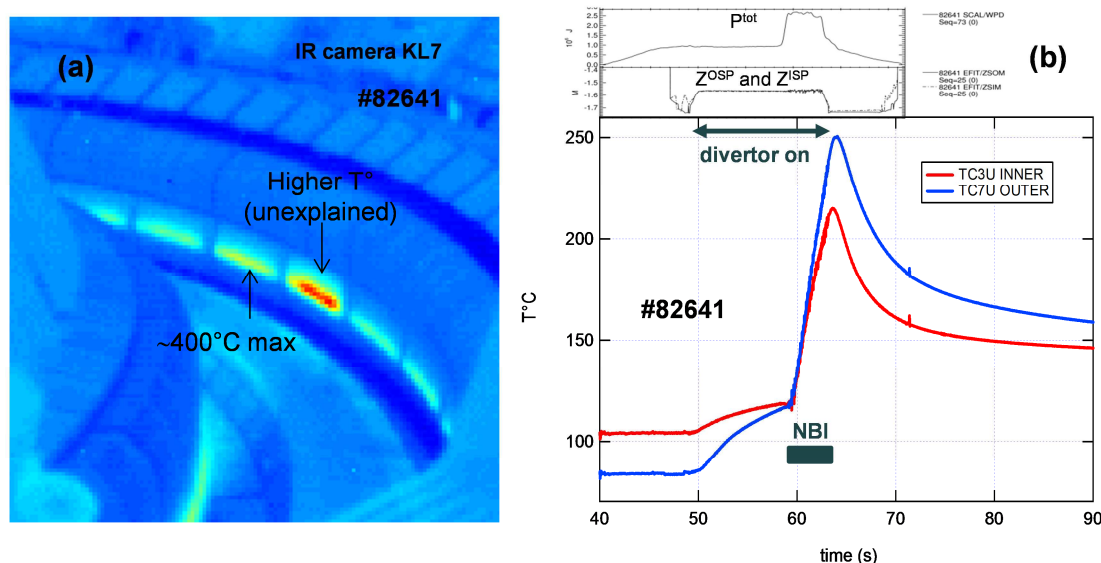
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<sup>\*</sup> See the Appendix of F. Romanelli et al., Proceedings of the 23rd IAEA Fusion Energy Conference 2010, Daejeon, Korea



**Figure 1:** (a) Poloidal cross-section of the divertor with the magnetic equilibrium and position of the TCs in the inner and outer targets. (b) View of the outer vertical tile with upper and lower embedded TCs.

The IR and TC data reported during #82641 are shown in figure 2-a and 2-b respectively. The brightness surface temperature measurement gives about 400°C (during the inter-ELM phase) on the outer tile shown in figure 2-a at the end of the 4s NBI heating phase assuming the emissivity of the W coated tile is constant and equal to 0.4 [1]. In the same time the bulk temperature increases by 100° and 140° on the inner and outer tile respectively (the total heating integrates inter-ELM and ELM phases).



**Figure 2:** (a) IR data (t=63s) taken from the wide angle viewing camera (KL7 IR camera) during pulse #82641. (b) Inner (TC3U, red curve) and outer (TC7U, blue curve) TC data during pulse #82641.

**3. Inverse heat flux calculation using TC data.** We propose to use an inverse thermal calculation to deduce the local surface temperature (thus on the W coating) and deposited heat flux using TC data taking into account the bi-dimensional effects due to the orthotropic material ( $k_y=k_z$  and  $k_x=k_z/4$ ) and the non-uniformity of the heat flux in the x-direction (poloidal). The method is the same as already applied with the MkiIGB CFC divertor: deconvolution of the thermocouple data with the 2D step response of the tile at the thermocouple location [2]. The 2D step response is computed with the thermal quadrupoles

(analytical solution of the heat transfer equation solved in Laplace-Fourier space) therefore with no thermal properties variation with temperature (linear calculation). The geometry of the carbon tile is approximated by a rectangular tile with a total average thickness  $e$  of 4 cm and a width  $L$  of 17.2 cm (Fig. 1-b). Uncertainties attributed to the main simplification (geometry plus linearity) are numerically evaluated using 2D FEM modeling, the error bar is  $\sim 15\%$  in normal plasma condition (when  $T < 500^\circ\text{C}$ ). The calculation requires the position of the peak (given by the offline reconstruction of the magnetic equilibrium) and the 1D-shape of the heat flux deposited on the target (qualitative shape). The target heat load profile is described with an heuristic expression based on the convolution of the scrape-off layer (SOL) exponential profile with a Gaussian width  $S$  representing the diffusion of the heat flux toward the targets (resulting from the parallel and perpendicular transport in the divertor) [3]:

$$q(s) = \frac{Q_0}{2} \exp\left(\frac{s}{2\lambda_{qf_x}}\right)^2 \frac{\bar{s}}{\lambda_{qf_x}} \cdot \operatorname{erfc}\left(\frac{s}{2\lambda_{qf_x}} - \frac{\bar{s}}{S}\right) + Q_{BG} \text{ and } \bar{s} = s - s_0 \quad (1)$$

Where  $s$  is the target coordinate,  $s_0$  is the strike line position,  $f_x$  is the magnetic flux expansion (integrated from the outer midplane to the target),  $\lambda_q$  is the SOL power decay length and  $Q_0$  is the peak heat load on the strike line position (i.e. the parameter that we propose to deduce with the TC data). The power decay length  $\lambda_q$  is given by an empirical scaling based on IR thermography analysis (mainly performed during JET and AUG C-wall operation):

$$\lambda_q^{JET} = 0.7 \cdot B^{-0.84}(T) \cdot q^{1.23} \cdot P_{SOL}^{0.14}(MW) \cdot R_{geo}^{0.02}(m) \quad (2)$$

Where  $B$  is the magnetic field,  $q$  is the cylindrical safety factor and  $P_{SOL}$  is the power crossing the separatrix ( $P_{SOL} = P_{tot} - P_{rad}$ ). For #82641,  $B = 2.2$  T,  $q = 3.2$ ,  $P_{SOL} \sim 7$  MW, we find  $\lambda_q = 2$  mm. The vertical outer target heat load profile given by (1) is shown in figure 3 for #82641,  $f_x = 8$ ,  $\lambda_q = 2$  mm, full line  $S = 1$  mm ( $S/\lambda_q = 0.5$ , mean value of the JET data) and dashed line  $S = 2$  mm ( $S/\lambda_q = 1$ , maximum value of the JET data). The inverse heat flux calculation has been applied with the 1D-shape presented above ( $\lambda_q = 2$  mm,  $S = 1$  mm,  $f_x = 8$  and  $Q_{BG} = 100$  kW/m<sup>2</sup> attributed to plasma radiation). Peak heat loads are found to be  $Q^{ISP} = 3$  MW/m<sup>2</sup> and  $Q^{OSP} = 5$  MW/m<sup>2</sup> with  $\pm 15\%$  uncertainties at the inner (ISP) and outer strike point (OSP). The corresponding

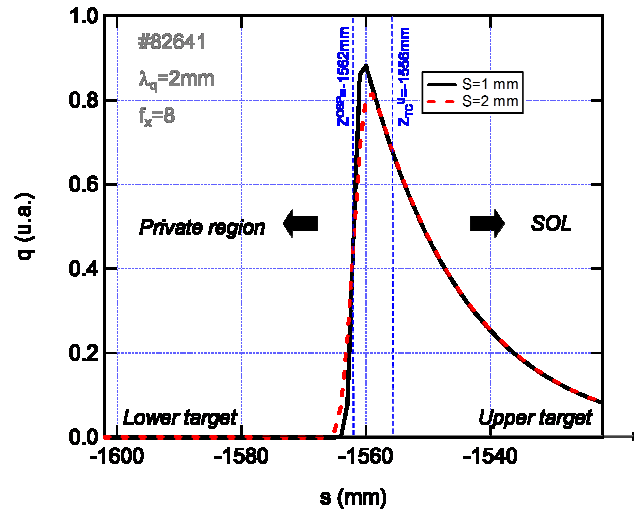
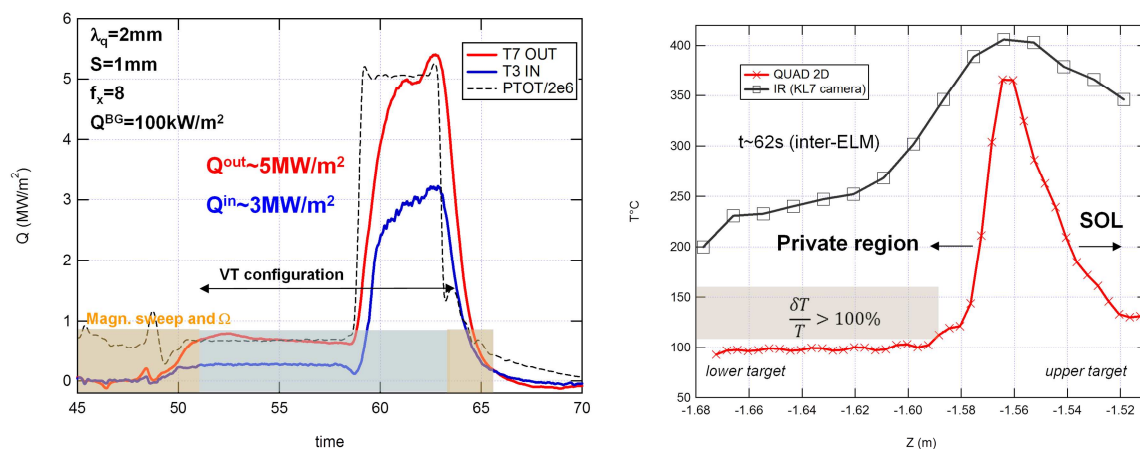


Figure 3: Normalized target heat load profile on the outer vertical tile for  $S = 1$  mm and  $S = 2$  mm.

surface temperatures and total accumulated energy (output power integrated over the pulse) are of 400°C/12MJ (OSP) and 280°C/5.6MJ (ISP) at the end of the heating phase (lasting 4s). The surface temperature target profiles derived from TC data (QUAD calculation) and IR images (using an emissivity of  $\epsilon$  of 0.4 [1]) are compared in figure 4-b. A relatively good agreement is obtained near the strike point (typically when  $T > 300^\circ\text{C}$ ) while a strong discrepancy is observed on the rest of the target (when temperature are below 300°). Further investigation would be required to understand the discrepancy at low temperatures (taking into account emissivity and reflection issues for IR data and strike point positioning and  $\lambda_q$  uncertainties with the inverse TC calculation).



**Figure 4:** (a) Inner / outer heat load computed with the 2D-QUAD method. (b) Surface temperature profiles along the OUTER vertical tile measured with TC (crosses) and wide angle IR (squares).

**4. Conclusions.** The combination of the inter-ELM empirical target heat load profile proposed in [3] (limited to attached plasmas) with an inverse heat flux calculation using embedded TC data provides a new way to compute heat load and surface temperature on targets (integrating inter-ELM and ELM phases). The first application to JET data with the ILW shows satisfactory results (temperature, heat load and energy). The main advantage of the method is that it is independent of surface layers, emissivity and reflection issues (associated to IR thermography in metal environment) while its main disadvantages are the poor spatial covering (due to the low number of TC present in the machine) in one side and the strike point position uncertainty (of the order of  $\text{cm} \gg \lambda_q$ ) in the other side. In metallic wall environment IR thermography and embedded TC are two complementary systems to guarantee the accuracy of the temperature and heat load measurements.

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[3] T. Eich et al., Phys. Rev. Lett. 107 (2011) 215001