

## Scaling of confinement in the ITPA L-mode database with dimensionless variables

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### Introduction

The International Tokamak Physics Activity (ITPA) L-mode database has been augmented with several new contributions since the previous release of an L-mode scaling expression [1]. Among the issues that this database allows us to address, the scaling of the energy confinement time with dimensionless parameters ( $\rho^*$ ,  $\beta$ ,  $\nu^*$ ) is of physical importance and of some interest for extrapolation to Next Step machines. Recently, this scaling has been revisited for the H-mode regime, following dimensionless scaling experiments on JET and DIII-D which showed a weak impact of  $\beta$  on the confinement, in contrast to the strong  $\beta$  degradation featured by the IPB98(y, 2) scaling expression [2]. It has been found that including the errors on the experimental measurement in the analysis could change the  $\beta$ -dependence of the scaling expression: the exponent of the  $\beta$ -dependence can be reduced down to a negligible value, while at the same time a degradation with the effective collisionality  $\nu^*$  appears [3]. In this work, the same statistical methods are applied to the new version of the L-mode database, in order to investigate the possible changes on the  $\beta$ -dependence of the L-mode scaling expression.

### Description of the L-mode data base

The L-mode database DB3v1 used in this paper features now about 3000 L-mode entries, including recent contributions from HL-1M, NSTX and JET. The different variables used in the analysis are:

$W_{th}$	The thermal stored energy (MJ)	$n$	the central line average density ( $10^{19} \text{ m}^{-3}$ )
$P$	the loss power (MW)	$I$	the current (MA)
$R$	the major radius (m)	$B$	the toroidal field (T)
$a$	the minor radius (m)	$M$	the isotope mass
$A$	the cross sectional area ( $\text{m}^2$ )	$\tau_{th}$	the thermal energy confinement time

The selection criteria are standard L-mode only (no improved confinement discharge), hydrogenic discharges only, semi-stationary,  $0.4 \leq T_{i0}/T_{e0} \leq 2.5$  and  $0.5 < li \leq 2$ . Then, the number of observations by tokamak and their relative weights, deduced by  $w_j \sim (2 + N^{1/2}/4)^{-1}$ , are:

	$N_i$	$w_j$		$N_i$	$w_j$		$N_i$	$w_j$
ASDEX	237	0.17	JFT2M	127	0.21	TdeV	7	0.38
CMOD	225	0.17	JT60	174	0.19	TEXTOR	81	0.24
D3D	124	0.21	NSTX	8	0.37	TFTR	400	0.14
FTU	6	0.38	PDX	32	0.29	TSUPRA	75	0.24
HL-1M	4	0.40	RTP	10	0.36			
JET	714	0.12	T10	28	0.30	Total	2252	

### Scaling laws of confinement

In this configuration, the standard log-linear method gives the scaling law of confinement below:

$$\tau_E = 0.016 I^{0.84} B^{0.12} R^{1.55} n^{0.46} a^{-0.83} A^{0.63} M^{0.21} P^{-0.76}$$

Then, for non-dimensional parameters, the law becomes:

$$\omega_c \tau_E \propto \beta^{-1.55} \rho^{*-1.84} \nu^{*0.38}$$

This scaling law is similar to the ITERL-97P one [1]:

$$\tau_E = 0.011 I^{0.96} B^{0.03} R^{1.77} n^{0.40} a^{-1.22} A^{0.64} M^{0.20} P^{-0.73}$$

$$\omega_c \tau_E \propto \beta^{-1.39} \rho^{*-1.98} \nu^{*0.19}$$

#### • Error in Variable

The error in variable technique (EIV), introduced to plasma physics in Kardaun et al. [4] and subsequently in Cordey et al. [3] weightens measurements with their average errors given in the above table:

	$\delta R$	$\delta a$	$\delta A$	$\delta B$	$\delta I$	$\delta n$	$\delta M$	$\delta P$	$\delta W_{th}$
<b>Average error</b>	1.3%	2.9%	4.7%	1.5%	1.3%	5.0%	8.4%	14.2%	14.1%

For the same assumptions on the data errorbars as used for the H-mode DB analysis, the original  $\beta$ -degradation of the energy confinement time is reduced to -0.77 while the  $\nu^*$ -dependence is almost null and the  $\rho^*$  one is nearly constant (Bohm scaling). This result is qualitatively similar to what has been obtained for the H mode DB, i.e. the EIV method provides a reduction of the  $\alpha_\beta$  exponent and simultaneously decreases the  $\nu^*$  exponent.

$$\tau_E = 0.0041 I^{0.73} B^{0.16} R^{1.97} n^{0.43} a^{-2.18} A^{1.14} M^{0.39} P^{-0.66}$$

$$\omega_c \tau_E \propto \beta^{-0.77} \rho^{*-2.07} \nu^{*0.08}$$

The results of the EIV method depend on the set of estimated errors, as shown in figure 1. To entirely eliminate the  $\beta$ -degradation, the estimated loss power error must be increased and/or the estimated stored energy error decreased. For  $\delta P = 28.8\%$  and  $\delta W_{th} = 11\%$ , we obtain:

$$\tau_E = 0.0020 I^{0.64} B^{0.07} R^{2.12} n^{0.29} a^{-3.00} A^{1.35} M^{0.62} P^{-0.41}$$

$$\omega_c \tau_E \propto \beta^{-0.00} \rho^{*-2.20} \nu^{*-0.87}$$

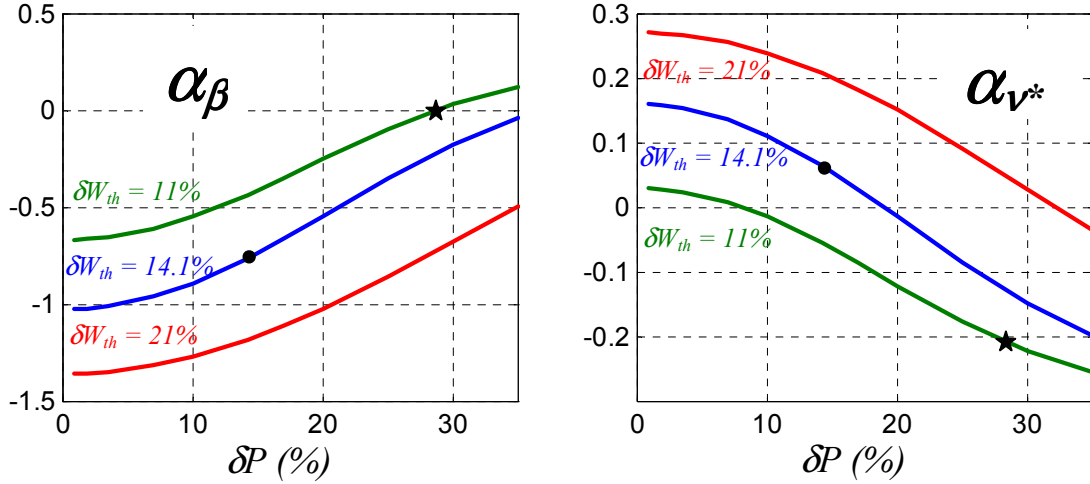


Fig. 1 – The exponents of  $\beta$  and  $v^*$  versus the estimated error on the loss power  $\delta P$  for different values of the uncertainty on the thermal energy  $\delta W_{th}$ . The stars indicate a possible couple  $(\delta P, \delta W_{th})$  that gives zero  $\beta$ -dependence; the dots correspond to the error dataset used for the H-mode DB analysis.

- **Sensitivity of the  $\beta$  exponent**

The form of the scaling is  $\tau_E = C I^{\alpha_I} B^{\alpha_B} R^{\alpha_R} n^{\alpha_n} a^{\alpha_a} A^{\alpha_A} M^{\alpha_M} P^{\alpha_P}$  and  $\omega_c \tau_E \propto \beta^{\alpha_\beta} \rho^{*\alpha_{\rho^*}} v^{*\alpha_{v^*}}$ ; the relations between the engineer exponents and the non-dimensional exponents can be written:

$$\alpha_\beta = \frac{7\alpha_p + \alpha_I + \alpha_B + 1 + 4\alpha_n}{4(1 + \alpha_p)} \quad \alpha_p = -\frac{9\alpha_p + 3(\alpha_I + \alpha_B + 1) + 4\alpha_n}{2(1 + \alpha_p)} \quad \alpha_v = -\frac{3\alpha_p + \alpha_I + \alpha_B + 1}{4(1 + \alpha_p)}$$

From these expressions, we can deduce that  $\alpha_\beta$  is very sensitive to  $\alpha_p$  especially when this one is close to  $-1$ . For exponents in the range of the L-mode scaling law ( $\alpha_p \sim -0.7$ ), a variation of 0.01 on  $\alpha_p$  induces an approximate variation of 0.1 on  $\alpha_\beta$ !

### Application to ITER configuration

Table beside shows the confinement time for different scaling laws when the loss power is equal to 28 MW (corresponding to an additional power of 32 MW which is just below the transition limit  $P_{L-H}$  calculated with  $B = 5.3\text{T}$ ,  $n = 5.10^{19}\text{m}^{-3}$ ,  $R = 6.2\text{m}$ ,  $a = 2\text{m}$ ,  $M = 2.5$   $I = 5\text{ MA}$ ,  $A = 22\text{ m}^2$ ).

	$\tau_E$	$\beta_N$
<b>ITERL-97P</b>	2.82	0.44
<b>Log-linear</b>	2.42	0.39
<b>EIV</b>	3.24	0.51

### Discussion

The high sensitivity of the  $\alpha_\beta$  exponent to the value of  $\alpha_p$  explains why  $\alpha_\beta$  can be significantly reduced by applying the EIV method. Nevertheless, even with the EIV method, a quite large value of  $\delta P$  and simultaneously a somewhat reduced value of  $\delta W_{th}$  have to be chosen in order to eliminate completely the  $\beta$ -degradation. Therefore the standard L mode dataset suggests that a  $\beta$ -degradation has still to be expected in L mode.

This result is in contrast to the H mode database, where using the EIV method with a realistic set of assumed errors is able to eliminate entirely the  $\beta$ -degradation.

The large variation of  $\alpha_\beta$  and  $\alpha_v$  as a function of the error in  $P$  has only a limited influence on the prediction for ITER FEAT, see [3].

The  $\beta$ -degradation remains even if the high  $\beta$  shots are removed from the scaling. Therefore, the  $\beta$ -degradation of the global dataset is not due to a change of the physics at high  $\beta$ . Likely L-mode discharges do not have high enough pressure to get close to ballooning stability limits.

Since there seem to be significant a scatter of the predicted  $\tau_E$ , scaling with respect to the measured value at high  $\tau_E$  (in particular JET and TFTR data), we have applied log-linear regression independently to JET and to TFTR data. The two machines have enough entries in the database so that the regression is statistically relevant, once having removed the geometry scale parameters  $R$ ,  $a$  and  $A$ . The result of these specific device regressions gives also reduced  $\beta$ -degradation with respect to the one obtained for the whole dataset (-0.5 for JET and -0.8 for TFTR). These values are of the same order than what is given by the EIV method for the whole database using the standard error set.

Dedicated L-mode  $\beta$  scaling experiments on various tokamaks yielded quite contrasted results on the  $\beta$ -degradation: DIII-D found  $\alpha_\beta \sim 0$  [5], TFTR  $\sim -0.5$  [6], while Tore Supra indicated a much larger  $\beta$ -degradation  $\alpha_\beta \sim -1.3$  [7]. Further experiments are needed in the various tokamaks in order to clarify this issue of  $\beta$ -degradation. Indeed, a clear guidance from experiments is needed in order to optimize the statistical treatment to apply to the global database.

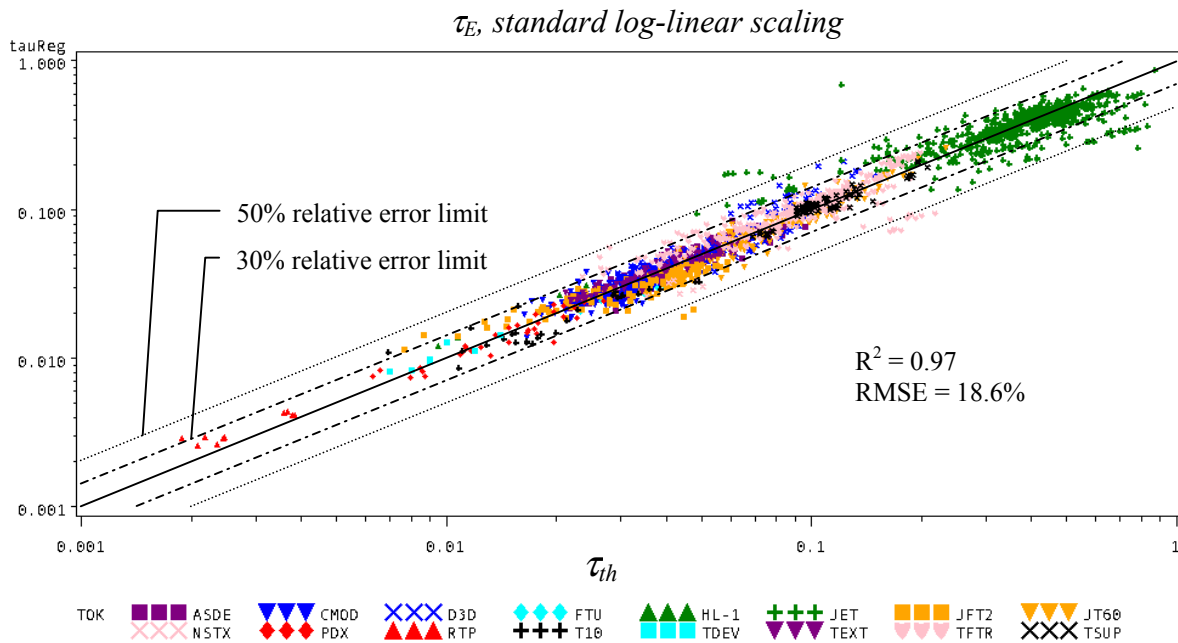


Fig. 2 -  $\tau_E$  scaling for log-linear regression vs. estimated  $\tau_{th}$ .

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