

## A NEW PROOF FOR THE LÉVY CONSTRUCTION OF SECOND KIND FOR STABLE LAWS

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We give a direct proof for the “Lévy construction of second kind” for stable laws on the real line without referring to the construction of “first kind.”

### 1. Introduction

Let  $X$  be a real-valued non-gaussian  $\alpha$ -stable random variable. It is well known that this is the case iff the Fourier transform (characteristic function) of  $X$  has the form

$$\begin{aligned} \varphi_X(u) = & \exp \left( iu\gamma + c_- \int_{-\infty}^0 \left( e^{iux} - 1 - \frac{iux}{1+x^2} \right) |x|^{-(1+\alpha)} dx + \right. \\ & \left. + c_+ \int_0^{\infty} \left( e^{iux} - 1 - \frac{iux}{1+x^2} \right) x^{-(1+\alpha)} dx \right), \quad u \in \mathbb{R}, \end{aligned}$$

with  $0 < \alpha < 2$ ,  $\gamma \in \mathbb{R}$ ,  $c_-, c_+ \geq 0$ ,  $c_- + c_+ > 0$ .

Possible “constructions” of  $X$  are the so-called Lévy constructions of “first” and “second kind.” These are the following.

Assume  $0 < \alpha < 2$ . Let  $\{N_t\}_{t \geq 0}$  be a Poisson process with parameter  $\lambda > 0$  and suppose  $\Gamma_j$  is the time of the  $j$ th jump of  $\{N_t\}_{t \geq 0}$ . Suppose  $\{Y_j\}_{j \in \mathbb{N}}$  is a sequence of i.i.d.  $\{-1, 1\}$ -valued random variables that is independent of the process  $\{N_t\}_{t \geq 0}$  and such that  $P(Y_j = 1) = p$ . Put  $a_j := 0$  for  $j \leq 1/\alpha$  and  $a_j := E(Y_j)E(\Gamma_j^{-1/\alpha})$  for  $j > 1/\alpha$ . Set

$$S_n(\alpha, \lambda, p, \gamma) := \gamma + \sum_{j=1}^n (\Gamma_j^{-1/\alpha} Y_j - a_j).$$

**Theorem 1** (Lévy construction of second kind). *The sum  $S_n(\alpha, \lambda, p, \gamma)$  converges to some  $S(\alpha, \lambda, p, \gamma)$  a.s. as  $n \rightarrow \infty$ , and  $S(\alpha, \lambda, p, \gamma)$  exhausts all (nondegenerate)  $\alpha$ -stable laws as  $(\lambda, p, \gamma) \in ]0, \infty[ \times ]0, 1] \times \mathbb{R}$ .*

For the “Lévy construction of first kind,” one just uses that

$$\mathcal{L}((\Gamma_1, \Gamma_2, \dots, \Gamma_n) \mid \Gamma_{n+1} = t) = \mathcal{L}((U_{[n:1]}, U_{[n:2]}, \dots, U_{[n:n]})), \quad (1)$$

where  $U_{[n:1]} < U_{[n:2]} < \dots < U_{[n:n]}$  denotes the increasing order statistics of independent random variables  $U_1, U_2, \dots, U_n$  distributed uniformly on  $[0, t]$ . Write

$$F_t(\alpha, \lambda, p, \gamma) := \gamma + \sum_{j=1}^{N_t} (U_{[N_t:j]}^{-1/\alpha} Y_j - a_j).$$

Then the “Lévy construction of first kind” is the following:

**Theorem 2** (Lévy construction of first kind). *The sum  $F_t(\alpha, \lambda, p, \gamma)$  converges weakly to some  $F(\alpha, \lambda, p, \gamma)$  as  $t \rightarrow \infty$ , and  $F(\alpha, \lambda, p, \gamma)$  exhausts all (nondegenerate)  $\alpha$ -stable laws as  $(\lambda, p, \gamma) \in ]0, \infty[ \times ]0, 1] \times \mathbb{R}$ .*

The classical proof of Theorem 1 proceeds in the manner that one first verifies Theorem 2 by calculating the Fourier transform of  $F_t(\alpha, \lambda, p, \gamma)$ , then uses (1), and at the end takes the limit as  $t \rightarrow \infty$ . In other words, the Lévy construction of second kind is deduced from that of the first kind by the equivalence (1). From the pedagogical

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point of view, this approach has one disadvantage: Although (1) seems to be quite intuitive, a formally absolute correct proof is quite cumbersome to write down. Often, textbooks just give a “proof” using manipulations with differentials (as, e.g., [1]). That is why in this note, we would like to show how a direct approach to the Lévy construction of second kind is possible without using (1). See, e.g., [3–7] for further information and generalizations of the Lévy construction.

**2. Alternative proof of Theorem 1**

The convergence result  $S_n(\alpha, \lambda, p, \gamma) \xrightarrow{a.s.} S(\alpha, \lambda, p, \gamma)$  follows from the Three Series Theorem by observing that the sequence  $\{\Gamma_j^{-1/\alpha}\}_{j \geq 1}$  behaves as  $\{j/\lambda\}_{j \geq 1}$  and the conditional variance of  $\Gamma_j^{-1/\alpha} Y_j$  given  $Y_j$  is of the form  $m \Gamma_j^{-2/\alpha} \sim m_1 j^{-2/\alpha}$  as  $j \rightarrow \infty$  (cf. [2]). In order to verify the stability of  $S(\alpha, \lambda, p, \gamma)$ , observe that the addition of  $n$  independent copies of  $S(\alpha, \lambda, p, \gamma)$  corresponds to a superposition of  $n$  independent sequences  $\{\Gamma_j\}_{j \geq 1}$ , i.e., to the addition of  $n$  independent copies of the Poisson process  $\{N_t\}_{t \geq 0}$ , which is equivalent to the multiplication of the intensity parameter  $\lambda$  by  $n$ . In the sequence of jump times  $\{\Gamma_j\}_{j \geq 1}$  this corresponds to a division by  $n$ , hence in the sequence  $\{\Gamma_j^{-1/\alpha}\}_{j \geq 1}$  to a multiplication with  $n^{1/\alpha}$ . More precisely: Let, for  $1 \leq k \leq n$ , processes  $\{N_t^{(k)}\}_{t \geq 0}$  and  $\{Y_j^{(k)}\}_{j \geq 1}$  be given as above such that the processes  $D^{(k)} := \{(N_t^{(k)}, Y_j^{(k)})\}_{t \geq 0, j \geq 1}$  are i.i.d.,  $\gamma^{(k)} = \gamma \in \mathbb{R}$ ,  $S^{(k)}(\alpha, \lambda, p, \gamma)$  as above. Then

$$\begin{aligned} \mathcal{L} \left( \sum_{k=1}^n S^{(k)}(\alpha, \lambda, p, \gamma) \right) &= \mathcal{L}(n\gamma + \sum_{k=1}^n \sum_{j=1}^{\infty} ((\Gamma_j^{(k)})^{-1/\alpha} Y_j^{(k)} - a_j)) = \\ &= \mathcal{L} \left( \tilde{\gamma} + \sum_{j=1}^{\infty} \tilde{\Gamma}_j^{-1/\alpha} (\tilde{Y}_j - \tilde{a}_j) \right), \end{aligned}$$

where  $\{\tilde{\Gamma}_j\}_{j \geq 0}$  ( $\tilde{\Gamma}_0 := 0$ ) is defined as a process with independent increments and

$$\mathcal{L}(\tilde{\Gamma}_{j+1} - \tilde{\Gamma}_j) = \mathcal{L}(\tilde{\Gamma}_1) = \mathcal{L}(\min_{1 \leq k \leq n} \Gamma_1^{(k)}), \tag{2}$$

$\tilde{Y}_j, \tilde{a}_j$  by analogy as above ( $\tilde{\Gamma}_j$  is the time of the  $j$ th jump of the superposition of the processes  $\{N_t^{(k)}\}_{t \geq 0}$ ,  $k = 1, 2, \dots, n$ ; the property that the increments are i.i.d. follows from the fact that the processes  $\{N_t^{(k)}\}_{t \geq 0}$  are themselves independent processes with i.i.d. increments). Now

$$\begin{aligned} P(\min_{1 \leq k \leq n} \Gamma_1^{(k)} \geq x) &= \prod_{k=1}^n P(\Gamma_1^{(k)} \geq x) = \prod_{k=1}^n e^{-\lambda x} = e^{-n\lambda x} = \\ &= P(\Gamma_1^{(1)} \geq nx) = P(\Gamma_1^{(1)}/n \geq x), \end{aligned}$$

i.e.,  $\mathcal{L}(\min_{1 \leq k \leq n} \Gamma_1^{(k)}) = \mathcal{L}(\Gamma_1^{(1)}/n)$ , hence  $\mathcal{L}(\tilde{\Gamma}_1^{-1/\alpha}) = \mathcal{L}(n^{1/\alpha}(\Gamma_1^{(1)})^{-1/\alpha})$ . Thus

$$\begin{aligned} \mathcal{L} \left( \sum_{k=1}^n S^{(k)}(\alpha, \lambda, p, \gamma) \right) &= \mathcal{L}(\tilde{\gamma} + n^{1/\alpha} S^{(1)}(\alpha, \lambda, p, 0)) = \\ &= \mathcal{L}(\tilde{\gamma} + n^{1/\alpha} S(\alpha, \lambda, p, 0)). \end{aligned}$$

Since this is true for all  $n \geq 1$ , this means that  $S(\alpha, \lambda, p, \gamma)$  obeys an  $\alpha$ -stable law.

It remains to show that every  $\alpha$ -stable law is of the form  $\mathcal{L}(S(\alpha, \lambda, p, \gamma))$ . It holds that

$$S(\alpha, 1, 1, 0) \stackrel{a.s.}{=} \lim_{t \rightarrow \infty} S_t(\alpha, 1),$$

where

$$S_t(\alpha, \lambda) := \sum_{j=1}^{N_t} \Gamma_j^{-1/\alpha} - \sum_{j=1}^{\lambda t} a_j.$$

Here  $\{\Gamma_j\}_{j \geq 1}$  as above with parameter  $\lambda$ . Observe that  $N_t - \lambda t \stackrel{a.s.}{=} o(t^{1/\alpha})$ ,  $t \rightarrow \infty$ , by the Law of the Iterated Logarithm and thus  $\sum_{j=1}^{N_t} a_j - \sum_{j=1}^{\lambda t} a_j \stackrel{a.s.}{\rightarrow} 0$ ,  $t \rightarrow \infty$  (cf. [2]). For every  $n \geq 1$  we have that

$$\mathcal{L}(S_t(\alpha, 1)) = \mathcal{L}(S_t(\alpha, 1/n) + b_n)$$

for suitable  $b_n \in \mathbb{R}$ , i.e.  $S_t(\alpha, 1)$  is infinitely divisible. Since for all  $t \geq 0$  the Lévy measure in the Lévy-Hinčin formula of  $\mathcal{L}(S_t(\alpha, 1))$  is concentrated on  $[0, \infty[$ , the same must hold for the limit  $\mathcal{L}(S(\alpha, 1, 1, 0))$  (see e.g. [1, Theorem 9.22]), hence the Fourier transform of  $\mu^{(0)} := \mathcal{L}(S(\alpha, 1, 1, 0))$  is of the form

$$\hat{\mu}^{(0)}(u) = \exp \left( iu\gamma^{(0)} + c_+^{(0)} \int_0^\infty \left( e^{iux} - 1 - \frac{iux}{1+x^2} \right) x^{-1+\alpha} dx \right)$$

for some  $c_+^{(0)} > 0$ . Now take any (nondegenerate)  $\alpha$ -stable law  $\mu$  given by the Fourier transform

$$\hat{\mu}(u) = \exp \left( iu\gamma^{(0)} + c_- \int_{-\infty}^0 \left( e^{iux} - 1 - \frac{iux}{1+x^2} \right) |x|^{-1+\alpha} dx + c_+ \int_0^\infty \left( e^{iux} - 1 - \frac{iux}{1+x^2} \right) x^{-1+\alpha} dx \right)$$

( $c_- + c_+ > 0$ ). Then we have

$$\mu = \mathcal{L}(c'_+ S'(\alpha, 1, 1, 0) - c'_- S''(\alpha, 1, 1, 0) + \gamma')$$

with  $c'_+ := (c_+/c_+^{(0)})^{1/\alpha}$  and  $c'_- := (c_-/c_+^{(0)})^{1/\alpha}$ , where  $S'(\alpha, 1, 1, 0)$  and  $S''(\alpha, 1, 1, 0)$  are i.i.d. random variables obeying the law  $\mathcal{L}(S(\alpha, 1, 1, 0))$ . However,

$$\mathcal{L}(c'_+ S'(\alpha, 1, 1, 0) - c'_- S''(\alpha, 1, 1, 0) + \gamma') = \mathcal{L}(S(\alpha, (c_+ + c_-)/c_+^{(0)}, c_+/(c_+ + c_-), \gamma')),$$

i.e.,  $\mu$  has indeed a Lévy construction of the second kind.

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