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Ellipsometric characterization of amorphous and polycrystalline silicon films deposited using a single wafer reactor

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DIGITAL OBJECT IDENTIFIER
<http://dx.doi.org/10.1063/PT.3.1619>
The article by Thorne Lay and Hiroo Kanamori (10.1063/PT.3.1619) is an excellent review of the seismic energy released by the 11-March 2011 earthquake in Japan. The authors state that the seismic energy released was approximately five times as much energy as the 100-megaton atomic bombing of Nagasaki in 1945. While that of a 100-megaton atomic bombing is approximately five times as much energy as the 100-megaton atmospheric nuclear detonation event, the authors state that the 11-March 2011 earthquake had still more energy by a factor of about 3, or 15 times as much energy as the 100-megaton atomic bombing. The 1964 Chilean earthquake had still more energy by a factor of about 3, or 15 times as much energy as the 100-megaton atomic bombing. I believe the authors used the relation for seismic energy release rather than total strain energy release. The seismic energy underestimates the total strain energy release by a variable that depends on the fault plane. Accounting for total strain energy release would increase the earthquake energy number by orders of magnitude. Despite the catastrophic damage potential of nuclear bombs, the forces of nature occasionally unleash much larger energy releases. Although the nuclear bombs are under our control, earthquakes, volcanic eruptions, and extreme weather events are not. However, by judicious preparation and avoidance measures, humans can significantly diminish the damage of natural events. This article does not have any references.

Comment on this article
By the act of hitting a ball with a bat, one calculates the force energy to deliver the ball to its new location, but one must also take into account that the ball extended its energy release to that location which became struck by the ball as its momentum ceased and passed energy to the struck team. Therefore the parameters of the damage extend into the future when the received energy to that pushed upon, later becomes released in a new event. Perhaps calculations of one added that in while another's calculations did not. E.M.C.
Written by Edgar Mocarill, 14 July 2012 19:59

Ellipsometric characterization of amorphous and polycrystalline silicon films deposited using a single wafer reactor

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The optical functions of amorphous and polycrystalline silicon thin films deposited on single oxidized silicon substrates by chemical vapor deposition in a wide range of deposition temperatures have been determined using spectroscopic ellipsometry. The data analysis is performed by direct inversion of the experimental spectra, therefore, obtaining results independent of any film modeling. The optical results indicate that the film structure changes as the deposition temperature increases from amorphous to polycrystalline with different grain size and distribution. © 1997 American Institute of Physics. [S0003-6951(97)02203-1]

Polycrystalline silicon (*p*-Si) is certainly acknowledged as a key material in electronic device manufacturing. Recently, a new generation of deposition equipment has been developed, where *p*-Si films are deposited by single wafer chemical vapor deposition (CVD). These systems offer peculiar properties of uniformity on wafer and wafer-to-wafer reproducibility in spite of the fact that deposition occurs at a very high rate (about 50–100 nm/min). This technology is properly extendible to the application in the Gbit generation devices. A deep characterization of the films deposited under different conditions is necessary for a complete knowledge of the *p*-Si properties and of the deposition reactors and for the tuning of reliable processes for the different applications. Among the different characterization techniques, the optical ones are the most interesting, due to their nondestructiveness, high sensitivity, and *in situ* applicability without any influence on the process conditions.

In this letter, a spectroscopic ellipsometry study of phosphorus-doped *p*-Si films deposited by single wafer CVD in a wide range of deposition temperatures is reported. The complex dielectric function $\tilde{\epsilon}$ of the different films is obtained and a close correlation between $\tilde{\epsilon}$ and the structural properties of the samples is found.

Twelve *P*-doped (2×10^{20} *P* atoms/cm³) *p*-Si films were deposited from diluted SiH₄ and PH₃ using an apparatus HT Centura by Applied Materials on a 80-nm-thick thermal oxide film on (100)-oriented Czochralski silicon substrates. Each film was deposited in 150 s. After the deposition, the samples were stored in air, therefore, allowing the growth of a few-nm-thick native SiO₂ layers to occur. This series of samples has been deposited varying the substrate temperature from 605 to 660 °C in 5 °C steps (samples

1–12). Ellipsometry measurements were carried out in the spectral range from 1.5 to 5 eV (0.25–0.85 μm) using a rotating polarizer instrument SOPRA ES 4G, with 2 nm spectral resolution. The ellipsometric experimental data are the tan ψ and cos Δ functions,¹ as defined from

$$\tilde{\rho} = \frac{\tilde{r}_p}{\tilde{r}_s} = \tan \psi \exp i\Delta, \quad (1)$$

where \tilde{r}_p and \tilde{r}_s are the Fresnel reflection coefficients for *p*-polarized and *s*-polarized light, respectively.

In Fig. 1 typical experimental spectra of tan ψ and cos Δ for a *p*-Si sample (sample 9 deposited at 645 °C) are shown. Considering the penetration depth of the visible light in amorphous silicon (*a*-Si) and in single-crystal silicon and assuming that for all the *p*-Si samples the penetration depth lies between those of such two reference materials, in the spectra of Fig. 1 two regions can be distinguished.² In the high-energy region, i.e., from about 3 to 5 eV, the light penetration depth is about 10 nm therefore, only the surface native oxide covering the *p*-Si film and its interface with *p*-Si are investigated. In this region, the fitting procedures (discussed below) used to interpret the bare spectra give satisfactory results for samples 1–9 (dots in Fig. 1 are the best fit data of sample 9), suggesting a negligible effect of the roughness of the interface between the native SiO₂ and the *p*-Si layer, possibly present. On the contrary, the experimental data for the high-temperature samples (from 10–12) are affected by the surface roughness, which increases as the deposition temperature increases, as evidenced by the results of the fits. A quantitative evaluation of such a roughness obtained by atomic force microscopy (AFM) measurements gives a root-mean-square value varying from about 1 to about 5 nm from sample 1 to 12. In the spectral range from 1.5 to about 3 eV the light penetrates the *p*-Si and the thermal oxide layers, reaching the interface with the silicon sub-

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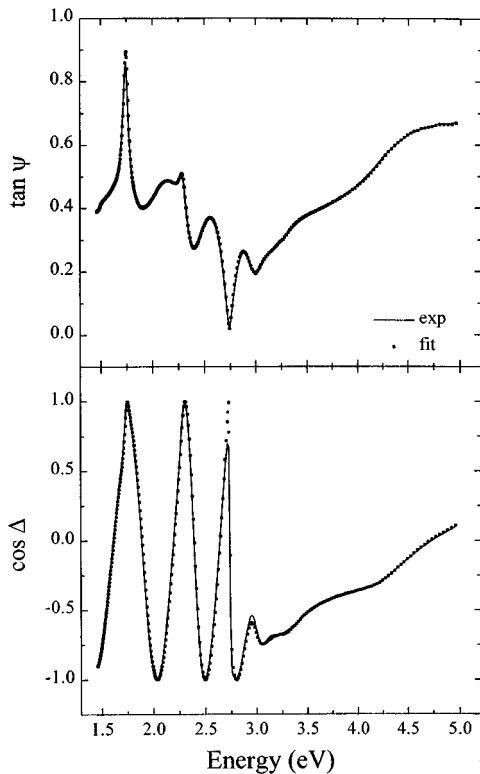


FIG. 1. Experimental $\tan \Psi$ and $\cos \Delta$ spectra of sample 9 (full lines) and best fit (dots).

strate. In this region, the experimental spectra of all samples are very similar to each other, showing strong interference fringes related to both the p -Si and the thermal oxide layers. Around 3 eV the fitting of experimental data is not very good due to the well-known poor quality of the p -Si/thermal oxide interface for deposition temperatures higher than about 615 °C.³ For lower energy the fitting of the experimental data is much easier, probably because of good and constant SiO_2 stoichiometry and structural order.

All the samples were modeled using the same three-layer structure, i.e. (from surface to substrate), a native oxide layer, the p -Si film of unknown thickness and optical properties, a thermal SiO_2 layer of 80 nm, and the single-crystal silicon substrate. The possibility of describing p -Si by a mixture of a -Si and single-crystal silicon, as far as it regards optical properties, is debatable and this point is discussed in the literature.^{4,5} In our opinion, as a first approach aimed at obtaining qualitative information on p -Si film crystallinity, a regression of the experimental $\tan \psi$ and $\cos \Delta$ considering the p -Si layer as a Bruggeman effective medium⁶ of the literature a -Si and single-crystal silicon⁷ can fruitfully be performed. In particular, for all the samples the thickness of the native oxide layer was determined by fitting the experimental curves in the high-energy region, while those of the p -Si and of the thermal oxide layers were determined in the low-energy region. Then, the fitting procedure over the whole spectral range was iterated several times till convergence was achieved. Both a low standard deviation⁸ σ and low correlation coefficients between the fitting parameters were used as criteria for selecting the best fit.

The results of this first regression, reported in Table I as a percentage of a -Si in the effective medium approximation

TABLE I. Amorphous silicon content in the p -Si layers resulting from the first regression of experimental ellipsometric data, and layer thickness t and standard deviation σ obtained after regression refinement and iteration.

Sample	a -Si (%)	t (nm)	$\sigma \times 10^2$
1 (605 °C)	100.0	65.8	... ^a
2 (610 °C)	100.0	77.5	0.62
3 (615 °C)	99.0	89.0	0.63
4 (620 °C)	94.8	101.1	1.00
5 (625 °C)	81.6	113.5	2.86
6 (630 °C)	28.2	130.7	3.23
7 (635 °C)	16.3	147.9	1.36
8 (640 °C)	13.2	170.8	0.99
9 (645 °C)	16.9	200.8	... ^a
10 (650 °C)	14.2	238.1	1.62
11 (655 °C)	16.8	282.3	1.70
12 (660 °C)	15.9	328.7	2.19

^aReference sample.

(EMA) mixture, clearly show that sample 1 can be considered as fully amorphous and that in samples 8–12 the percentage of a -Si is the same. This suggests the possibility of employing the optical functions of the p -Si layer of sample 1 (amorphous) and of the p -Si layer of a sample among 9–12 as internal references to repeat the regression obtaining more reliable thickness values. As reference polycrystalline material, we chose p -Si of sample 9 because atomic force microscopy measurements show that it is affected by the lowest surface roughness among samples 9–12. As a starting step, we supposed that sample 1 and sample 9 are composed of a -Si and p -Si, respectively, which we expect to be similar to the ones reported in Ref. 3, deposited by conventional CVD. Therefore, a second regression only on these samples was performed using the same EMA description as before, but considering a mixture of these two literature⁴ materials and single-crystal silicon, obtaining reliable thickness values. These data, as well as the thickness of the thermal SiO_2 layers, are in agreement with transmission electron microscopy (TEM) results within 3%. We then abandoned the EMA description of the p -Si layer because, once the layer thickness is known, a direct numerical inversion of the ellipsometric data is possible and gives more reliable spectra of the complex dielectric function $\tilde{\epsilon} = \epsilon_1 + i\epsilon_2$, being free of any film modeling.

The optical functions ϵ_1 and ϵ_2 obtained for the amorphous p -Si layer in sample 1 are very similar to those of a -Si deposited by conventional CVD, while the functions of the p -Si layer in sample 9 are strongly different from those of p -Si deposited by the conventional CVD techniques.

To evaluate the complex dielectric function of all the other p -Si layers, a procedure for data analysis similar to that used for samples 1 and 9 was adopted, considering as reference data the $\tilde{\epsilon}$ of these two samples. First, data regression has been refined describing the p -Si layer as an EMA mixture to obtain reliable thickness values. The agreement found with TEM results is again very good (within 3%–4%). Then, the ϵ_1 and ϵ_2 functions were calculated by numerical inversion of the ellipsometric data. In particular, to obtain the film thickness the p -Si layer in sample 2 was modeled by a mixture of the p -Si of sample 1 and crystalline silicon, the p -Si layer in sample 3 was modeled by a mixture of p -Si of

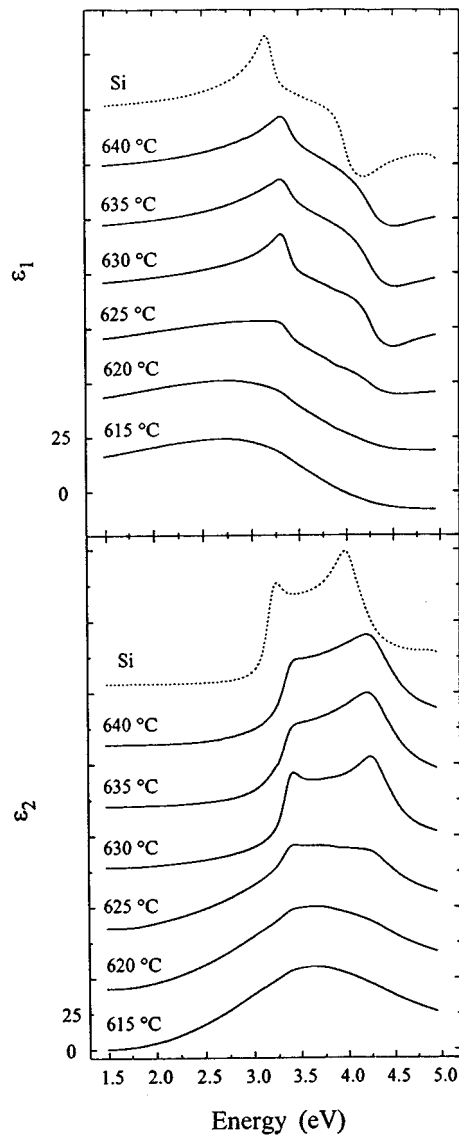


FIG. 2. Real and imaginary parts ϵ_1 and ϵ_2 of the complex dielectric functions of samples from 3 to 8 (deposition temperature from 615 to 640 °C, full lines) and of literature single-crystal silicon (dashed lines). The vertical scales refer to the 615 °C spectra; each of the other ϵ_1 and ϵ_2 spectra has been shifted by 30 and 40 units, respectively, with respect to the preceding one.

sample 2 and crystalline silicon, and so on for sample 4. A similar procedure for the analysis of the experimental data of samples 8 to 5 was adopted, considering the optical functions of the p -Si layer in sample 9. More precisely, such optical functions together with a -Si and crystalline silicon functions were used for the EMA modeling of the p -Si layer in sample 8 to get its thickness, and so on down to sample 5, obtaining very good results. The $\tilde{\epsilon}$ function for the high-temperature p -Si films (samples 10–12) were also derived in the same way, starting from the reference $\tilde{\epsilon}$ of p -Si in sample 9. The layer thickness t and the parameter σ in Table I are the p -Si thickness and standard deviation resulting from such analysis.

A low standard deviation σ is reached for samples 1–3 and 7–11, deposited at the lowest and the highest temperatures, respectively, while σ is less satisfactory for samples 5, 6, and 7. As for the p -Si layers in the samples, TEM images

show that they are dispersions of isolated crystalline grains in an amorphous matrix. Such grains are larger as the deposition temperature is higher, i.e., from sample 5–7, with dimensions of the order of the wavelength. Therefore, for such samples the EMA is very poor and this justifies the high standard deviation reported in Table I. The high value of σ for sample 12 has to be attributed to the mentioned surface roughness.

Figure 2 shows the spectra of ϵ_1 and ϵ_2 for samples 3–8, deposited at temperatures from 615 to 640 °C (full lines), obtained from the numerical inversion of the ellipsometric data, and for literature crystalline silicon (dashed lines). The results for the fully amorphous samples (1 and 2) and for the polycrystalline ones (9–12) are not reported, being exactly the same as those for samples 3 and 8, respectively.

The shape of the ϵ_1 and ϵ_2 curves changes from that of a -Si to very close to that of crystalline Si as the deposition temperature increases from 615 to 630 °C. The optical behavior of the p -Si deposited at this last temperature is the most similar to that of single-crystal silicon. Indeed, the intensity of the peaks in ϵ_1 and ϵ_2 at about 3.3 and 4.5 eV for the samples deposited at higher temperatures is lower. These peaks are indicative of the crystallinity of the samples, being related to critical points in the joint density of states of single-crystal silicon,^{9,10} and are very sensitive to the size of the crystals. Therefore, the height of these peaks in the spectra of the sample deposited at 630 °C indicate that it contains the largest crystalline grains.

In conclusion, the optical functions of P -doped polycrystalline silicon films deposited by CVD were obtained for a wide range of deposition temperatures from spectroscopic ellipsometry measurements by direct inversion of the experimental data, therefore, independently of any modeling of the film. The low-temperature films are amorphous and very similar to conventional CVD a -Si, while the high-temperature ones are polycrystalline with a strongly different optical behavior with respect to conventional CVD p -Si films. A transition temperature was observed in the range from 625 to 640 °C, where the p -Si structure shifts from that of a crystalline grain dispersion in an a -Si matrix to that of a true polycrystal, with closely packed grains.

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