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Rapid destruction of protoplanetary discs due to external photoevaporation in star-forming regions

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ABSTRACT

We analyse N -body simulations of star-forming regions to investigate the effects of external far- and extreme-ultraviolet photoevaporation from massive stars on protoplanetary discs. By varying the initial conditions of simulated star-forming regions, such as the spatial distribution, net bulk motion (virial ratio), and density, we investigate which parameters most affect the rate at which discs are dispersed due to external photoevaporation. We find that disc dispersal due to external photoevaporation is faster in highly substructured star-forming regions than in smooth and centrally concentrated regions. Subvirial star-forming regions undergoing collapse also show higher rates of disc dispersal than regions that are in virial equilibrium or are expanding. In moderately dense ($\sim 100 M_{\odot} \text{pc}^{-3}$) regions, half of all protoplanetary discs with radii ≥ 100 au are photoevaporated within 1 Myr, three times faster than is currently suggested by observational studies. Discs in lower density star-forming regions ($\sim 10 M_{\odot} \text{pc}^{-3}$) survive for longer, but half are still dispersed on short time-scales (~ 2 Myr). This demonstrates that the initial conditions of the star-forming regions will greatly impact the evolution and lifetime of protoplanetary discs. These results also imply that either gas giant planet formation is extremely rapid and occurs before the gas component of discs is evaporated, or gas giants only form in low-density star-forming regions where no massive stars are present to photoevaporate gas from protoplanetary discs.

Key words: methods: numerical – protoplanetary discs – photodissociation region (PDR) – open clusters and associations: general.

1 INTRODUCTION

Protoplanetary discs (‘Proplyds’) are thin Keplerian discs around pre-main-sequence stars (Shu et al. 1994) and are the birth places of planets. Proplyds form as a result of angular momentum conservation during the gravitational collapse of clouds when stars are forming. Within the discs, dust can coagulate to form a range of objects from pebbles to planets.

The evolution and dispersal of Proplyds controls the planet formation process, and observations suggest that disc lifetimes are between ≈ 3 and 5 Myr (e.g. Zuckerman, Forveille & Kastner 1995; Haisch, Lada & Lada 2001; Pascucci et al. 2006; Richert et al. 2018). Internal processes remove mass from the proplyd and, after

several Myr, disc accretion slows significantly to the point that these processes begin removing more mass than can be replaced, leading to very rapid disc dispersal (Clarke, Gendrin & Sotomayor 2001; Owen, Ercolano & Clarke 2011).

Proplyd host stars do not form in isolation, but rather in clusters and associations with stellar densities that exceed that of the Galactic field by a few orders of magnitude (Lada & Lada 2003; Bressert et al. 2010). Tens to thousands of stars can form in these regions that are a fraction of a parsec in size (Clarke, Bonnell & Hillenbrand 2000). Observations of young star-forming regions have revealed that stars form in filamentary structures (André et al. 2010; Wang et al. 2014), resulting in hierarchical spatial distributions. The net motion of stars within these regions indicates that the structures are often collapsing (i.e. subvirial; Peretto, André & Belloche 2006; Foster et al. 2015; Kuhn et al. 2018).

The initial densities of star-forming regions are difficult to determine, and span a wide range (Bressert et al. 2010; King et al.

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2012), but many are thought to be at least $\sim 100 M_{\odot} \text{pc}^{-3}$ at the epoch of star formation (Parker 2014). Parker (2014) shows that two regions with similar present-day densities at present times may have originally had very different initial densities because initially dense regions expand much faster than lower density regions due to two-body relaxation. They compare the present-day stellar densities and amount of spatial substructure in seven star-forming regions, including the ONC and Upper Scorpius – in which both contain massive stars that could act as external photoionizing sources – to infer the likely range of initial stellar densities in each of these star-forming regions and all are consistent with having an initial density in the range $10\text{--}1000 M_{\odot} \text{pc}^3$.

External processes, such as close stellar interactions, can also cause proplyds to be truncated or destroyed, as well as disrupting the orbits of fledgling planets (Armitage 2000; Bonnell et al. 2001; Scally & Clarke 2001; Adams et al. 2006; Olczak, Pflanzner & Eckart 2008; Parker & Quanz 2012; Rosotti et al. 2014; Vincke, Breslau & Pflanzner 2015; Portegies Zwart 2016; Winter et al. 2018a). The density of the star-forming region will affect the rate of stellar interactions, with stars in low-density environments experiencing fewer dynamical interactions than higher density environments (Bressert et al. 2010; Wright et al. 2014). Furthermore, star-forming regions can contain massive stars ($> 15 M_{\odot}$), whose intense far-ultraviolet (FUV) and extreme-ultraviolet (EUV) radiation fields are significantly higher than those in the interstellar medium (Armitage 2000; Adams et al. 2004; Fatuzzo & Adams 2008). This high-energy radiation heats the gaseous material of the upper layers of the disc until the thermal energy of the heated layer exceeds the gravitational potential of the disc, causing it to escape as a photoevaporative wind (Hollenbach et al. 1994; Johnstone, Hollenbach & Bally 1998). This mass-loss will affect the evolution of proplyds, and reduce the reservoir of material available to form gas giant planets (Haworth et al. 2018a).

The effects of external photoevaporation appear to be observed in nearby star-forming regions, such as the Orion nebula cluster (ONC; McCaughrean & O’dell 1996; Eisner et al. 2016, 2018) and σ Orionis (Ansdell et al. 2017). The ONC has been preferentially observed due to its proximity to the Earth (~ 415 pc) and because the discs can be viewed in silhouette due to the bright nebulous background. Studies have also shown that 80–85 per cent of stars within the ONC host proplyds (Bally, O’Dell & McCaughrean 2000; Lada et al. 2000), making the ONC a favourable target for studying disc evolution and dispersal (though see Clarke 2007, for an alternative interpretation that posits that the ONC proplyds are merely ionization fronts of material left over from discs that are almost destroyed).

Due to two-body and violent relaxation, initially highly substructured star-forming regions can evolve to smooth and centrally concentrated clusters after only a few Myr (Allison et al. 2010; Parker et al. 2014). Furthermore, two clusters that presently have similar densities may have had very different initial densities because initially very dense clusters expand faster than lower density counterparts. As mentioned before, the initial density will affect the rate at which proplyds are disrupted and destroyed due to stellar interactions. However, how much the initial density and substructure of a star-forming region affect the rate of proplyd dispersal due to external photoevaporation is yet to be studied.

Previous studies into the effects of external photoevaporation on proplyds in star-forming regions have tended to calculate the background UV radiation without directly calculating the disc mass-loss (Armitage 2000; Adams et al. 2004). Scally &

Clarke (2001) did calculate mass-loss rates in simulations specifically tailored to match the ONC, but assumed rather low stellar densities ($\sim 40 M_{\odot} \text{pc}^{-3}$), whereas Parker (2014) suggests that the initial density of the ONC may have been much higher ($> 100 M_{\odot} \text{pc}^{-3}$).

These previous studies of external photoevaporation used spherically smooth spatial distributions with primordial mass segregation to model the environment of the ONC as observed today. However, observations of star-forming regions show that stars form in highly substructured filamentary environments, where the stars are moving with subvirial velocities. The initial net motion and spatial structure of a star-forming region will govern its future evolution, and by extension the degree to which planet formation is hindered. These initial conditions lead to dynamical mass segregation on time-scales of the age of the ONC (Allison et al. 2010; Parker et al. 2014), negating the requirement for primordial mass segregation (Bonnell & Davies 1998).

Here, we focus on initial conditions that more closely reflect observations of young star-forming regions (Cartwright & Whitworth 2004), and determine how much external photoevaporation affects the evolution of proplyds. Therefore, we do not centrally concentrate our massive stars, but randomly distribute them in our simulated star-forming regions. We run suites of simulations that cover a range of initial conditions, with varying initial density, spatial distribution, and net bulk motion (virial ratio). We then calculate and compare the mass-loss rates due to external photoevaporation for each set of initial conditions.

The paper is organized as follows. In Section 2, we describe our N -body simulations, proplyd assumptions and our external photoevaporation prescription; in Section 3 we present our results; we provide a discussion in Section 4, and we conclude in Section 5.

2 METHOD

In this section, we describe our method to select low-mass star-forming regions containing massive stars, before describing the subsequent N -body and stellar evolution of these regions.

2.1 Creating low-mass star-forming regions

We adopt two different masses (100 or $1000 M_{\odot}$) for our star-forming regions and populate these regions with stars drawn randomly from the initial mass function (IMF) parametrized in Maschberger (2013), which has a probability density function of the form

$$p(m) \propto \left(\frac{m}{\mu}\right)^{-\alpha} \left(1 + \left(\frac{m}{\mu}\right)^{1-\alpha}\right)^{-\beta}. \quad (1)$$

Here, $\mu = 0.2 M_{\odot}$ is the average stellar mass, $\alpha = 2.3$ is the Salpeter (1955) power-law exponent for higher mass stars, and $\beta = 1.4$ is used to describe the slope of the IMF for low-mass objects (which also deviates from the lognormal form; Bastian, Covey & Meyer 2010). Finally, we sample from this IMF within the mass range $m_{\text{low}} = 0.1 M_{\odot}$ to $m_{\text{up}} = 50 M_{\odot}$.

We use a ‘soft-sampling’ technique for sampling the IMF (Elmegreen 2006), which implies that the only formal limit on the most massive star that can form is that of the upper limit of the IMF (Parker & Goodwin 2007). From an ensemble of Monte Carlo simulations, we find that typically, a $1000 M_{\odot}$ star-forming region will contain five massive stars ($M_{\star} > 15 M_{\odot}$).

Low-mass star-forming regions rarely contain massive stars; however, if the only limit on the mass of the star that can form is the total mass of the star-forming region itself then occasionally we would expect a low-mass star-forming region to contain one or more massive stars and such regions are observed (e.g. γ^2 Vel, a low-mass region containing at least two massive stars, Jeffries et al. 2014). Note that we are not explicitly attempting to model the γ^2 Vel star-forming region, which harbours a dense (sub)cluster within a more diffuse region. Instead, we are pointing out that observational examples of low-mass star-forming regions such as this and σ Orionis occasionally contain massive stars and their photoevaporative effects on discs in these low-mass star-forming regions could be important.

In order to demonstrate the importance of low-mass star-forming regions, let us consider the demographics of star-forming regions from randomly sampling the underlying probability distributions. The observed mass function of star-forming regions follows a power law of the form

$$N \propto M_{\text{cl}}^{-2}, \quad (2)$$

where N is the number of star-forming regions with mass M_{cl} (Lada & Lada 2003), which implies that there many more low-mass star-forming regions compared to high-mass regions.

We follow the procedure in Parker & Goodwin (2007) and Nicholson & Parker (2017) and sample 1×10^6 star-forming regions in the mass range $50\text{--}10^5 M_{\odot}$. Of these, ~ 1200 lie in the mass range $1000 \pm 10 M_{\odot}$, and $\sim 15\,000$ lie in the mass range $100 \pm 10 M_{\odot}$.

We then randomly populate our star-forming regions with stars drawn from equation (1), until the total mass of stars equals or exceeds the chosen star-forming region mass from equation (2).

Nicholson & Parker (2017) found that ~ 10 per cent of low-mass star-forming regions contain at least one massive star ($> 15 M_{\odot}$) when using the ‘soft-sampling’ technique described above, and 1 per cent of low-mass regions contain two massive stars. Furthermore, when taking into account the decreasing probability of forming a high-mass star-forming region (equation 2), the number of low-mass ($M_{\text{cl}} = 100 \pm 1 M_{\odot}$) regions containing at least one massive star is ~ 3100 , which is actually greater than the total number of high-mass ($M_{\text{cl}} = 1000 \pm 10 M_{\odot}$) star-forming regions (1200). Of these 1200 high-mass regions, ~ 1000 contain at least one massive ($> 15 M_{\odot}$) star.

If we translate these numbers into the total number of stars that may be affected by photoevaporation, the average number of stars in our $M_{\text{cl}} = 100 \pm 1 M_{\odot}$ star-forming regions containing at least one massive star is ~ 110 , so in the 3100 low-mass regions that contain at least one massive star there are $\sim 341\,000$ stars in total. The average number of stars in our high-mass star-forming regions ($M_{\text{cl}} = 1000 \pm 10 M_{\odot}$) is ~ 1710 , and so the 1000 regions that contain at least one massive star host a total of $\sim 1\,700\,000$ stars that could be affected by photoevaporation. In short, the fraction of stars originating in low-mass star-forming regions containing at least one massive star is ~ 20 per cent of the total number of stars originating from high-mass regions containing at least one massive star. If we stipulate that the high-mass regions must contain three or more massive stars, only ~ 580 regions of 1200 fulfil this criteria and host a total of 986 000 stars. The fraction of stars originating in low-mass star-forming regions containing at least one massive star is ~ 35 per cent of the total number of stars originating from high-mass regions containing at least three massive stars.

Crucially, this makes no assumption about the disruption and dissolution of these star-forming regions, and how many stars from each type of region eventually enter the Galactic field. The Galactic

potential will influence the destruction of low-mass star-forming regions much more than high-mass regions (Binney & Tremaine 2008), which take longer to dissolve into the Galactic field (and some remain as long-lived open clusters). Therefore, the majority of planet-hosting field stars may come from lower mass regions.

Given their significant contribution to the integrated stellar mass function, we therefore also investigate low-mass star-forming regions ($100 M_{\odot}$) that contain either one or two massive stars – these represent an unusual sampling of the IMF but allow us to investigate the effects of photoevaporation in less populous star-forming regions.

Hence, we have three different star-forming region set-ups; a $100 M_{\odot}$ region with one massive star ($38 M_{\odot}$), a $100 M_{\odot}$ region with two massive stars (42 and $23 M_{\odot}$), and one $1000 M_{\odot}$ region with five massive stars (43, 33, 26, 17, and $17 M_{\odot}$). These regions were selected as the median outcomes of Monte Carlo sampling of 1×10^6 star-forming regions (Nicholson & Parker 2017), and then filled with stellar masses drawn from the IMF (Maschberger 2013). We then selected the median regions in terms of the total number of stars from within the mass ranges of 100 ± 1 and $1000 \pm 10 M_{\odot}$, with the stipulation that they had to contain massive stars. For the $100 M_{\odot}$ mass regions, we specifically selected the median region containing one and two massive stars. For the $1000 M_{\odot}$, we selected the average cluster that contained three or more massive stars.

The external photoevaporation prescriptions we will adopt in this work are those from Scally & Clarke (2001) that only weakly depend on the adopted stellar IMF, but in a future paper we will assign an FUV flux and an EUV flux to each intermediate/high-mass star based on its mass and then determine the respective fluxes incident on every low-mass star and use the recent FRIED grid of models (Haworth et al. 2018b) to determine mass-loss for individual discs.

2.2 N-body simulations

Our simulations are created with initial substructure by following the box-fractal method in Goodwin & Whitworth (2004). We use a range of fractal dimensions for varying amounts of substructure: $D = 1.6$ (highly substructured), $D = 2.0$ (moderately substructured), and $D = 3.0$ (smooth). The method also correlates stellar velocities on local scales so that nearby stars have similar velocities, but more distant stars can have a wide range of different velocities.

The velocities are then scaled to the virial ratio α_{vir} , where $\alpha_{\text{vir}} = T/\Omega$; T is the total kinetic energy and Ω is the total potential energy of the stars. A range of virial ratios are used: $\alpha_{\text{vir}} = 0.3$ (subvirial, or collapsing), $\alpha_{\text{vir}} = 0.5$ (virial equilibrium), and $\alpha_{\text{vir}} = 0.7$ (supervirial, or expanding). Note that this virial ratio determines the net bulk motion, i.e. whether the star-forming region will collapse or expand. The correlated velocities on local scales mean that the local velocity dispersion can be subvirial, facilitating a violent relaxation process as the star-forming region evolves. Such local subvirial velocity dispersions are observed in the earliest stages of star formation (Larson 1981; Peretto et al. 2006; Foster et al. 2015).

We create star-forming regions with stellar densities of 100 or $10 M_{\odot} \text{pc}^{-3}$ for the $1000 M_{\odot}$ regions; for the $100 M_{\odot}$ regions we set an initial density of $100 M_{\odot} \text{pc}^{-3}$. Such densities bracket the range observed in present-day star-forming regions (Bressert et al. 2010) as well as allowing for potentially higher *primordial* densities (Parker 2014).

Finally, we created a set of simulations with a Plummer sphere distribution (Plummer 1911) to facilitate comparisons with previous studies. We use the same IMF from our $1000 M_{\odot}$ simulations to

create two clusters with Plummer sphere distributions that have initial densities of 10 and 100 $M_{\odot} \text{pc}^{-3}$.

We evolve each of our star-forming regions for 10 Myr using the `kira` integrator within the `Starlab` environment (Portegies Zwart et al. 2001). Stellar evolution is implemented using the `SeBa` look-up tables (Portegies Zwart & Verbunt 1996). No binary or multiple stellar systems are included in these simulations. To gauge the amount of stochasticity in the disc photoevaporation, we run 20 realizations of the same initial conditions, identical apart from the random number seed used to assign the positions and velocities.

2.3 Protoplanetary discs and external photoevaporation

The mass-loss rate of discs at a certain distance from a neighbouring massive star is determined by the strength of the star's FUV ($h\nu < 13.6 \text{ eV}$) and EUV ($h\nu > 13.6 \text{ eV}$) fluxes at that distance. Mass-loss due to FUV photons is caused by heating the circumstellar disc, which creates an unbound neutral layer that can drift towards the ionization front, where it meets the EUV field. FUV is independent of the distance from the massive star because the only requirement is that the FUV flux is strong enough to heat the disc above its escape velocity. With EUV, the mass-loss rate depends on the EUV flux and so is directly dependent on the distance from the massive star(s).

We use the same prescriptions for FUV and EUV photoevaporation as Scally & Clarke (2001):

$$\dot{M}_{\text{FUV}} \approx 2 \times 10^{-9} r_d M_{\odot} \text{yr}^{-1}, \quad (3)$$

$$\dot{M}_{\text{EUV}} \approx 8 \times 10^{-12} r_d^{3/2} \frac{\Phi_i}{d^2} M_{\odot} \text{yr}^{-1}, \quad (4)$$

where r_d is the radius of the disc in astronomical units, Φ_i is the ionizing EUV photon luminosity from each massive star in units of 10^{49} s^{-1} , and d is the distance from the massive star in parsecs. The UV photon rate (Φ_i) for the massive stars ($> 15 M_{\odot}$) is dependent on stellar mass, and we use the values from Vacca, Garmany & Shull (1996) and Sternberg, Hoffmann & Pauldrach (2003).

These photoevaporation rates were derived assuming a disc density profile $\Sigma \propto r_d^{-2}$ (Hollenbach, Yorke & Johnstone 2000; Hartmann 2009); however, our analysis does not take into account the evolution of the surface density profile if the disc radius were to change significantly.

Observations of star-forming regions show that disc radii can extend to several 100s au (e.g. Ansdell et al. 2018). However, the typical initial radius of proplyds is still debated in the literature and therefore we sample a wide range of initial radii: 10, 50, 100, 200, and 1000 au. We adopt a single value for each disc radius, focusing primarily on 100 au discs, and then repeat the analysis for the five values. In reality, the radius of the disc will change due to internal processes such as viscous evolution, and due to internal and external photoevaporation, but we are unable to account for this (and the changing disc density profile) in our N -body simulations.

The initial disc masses are also debated, with theoretical constraints from the Minimum Mass Solar Nebula (MMSN, Weiden-schilling 1977; Hayashi 1981) and observations (Andrews et al. 2013; Ansdell et al. 2016) suggesting an upper limit of $M_{\text{disc}} = 0.01 M_{\star}$. Following the example in Scally & Clarke (2001), the initial disc masses in our simulation are 10 per cent of the host star mass ($M_{\text{disc}} = 0.1 M_{\star}$). Current observations suggest that disc masses are ~ 1 per cent of the host star mass. We select 10 per cent so that we are sampling the upper range of the disc masses. While

we do not account for accretion on to the proplyds, our discs are large enough in mass that we can neglect the accretion on to the disc as it will be minimal in comparison. For completeness, we ran a set of simulations where the disc masses were 1 per cent of the stellar host mass, which is more consistent with the MMSN estimates.

We subtract mass from our discs according to equations (2) and (3). The rate of photoevaporation due to EUV radiation is dependent on distance from the ionizing source, d , whereas the photoevaporation rate due to FUV is largely independent of distance from the source (Störzer & Hollenbach 1999). Following Störzer & Hollenbach (1999) and Scally & Clarke (2001), we apply equation (2) if the disc-hosting stars are within 0.3 pc of the ionizing source, noting that this distance is calibrated to models where $\theta^1 \text{C Ori}$ is the most massive star ($40 M_{\odot}$), which is commensurate with the most massive star in our simulations. However, we note that this is likely an underestimate of the amount of photoevaporation due to FUV fields in star-forming regions (Facchini, Clarke & Bisbas 2016; Haworth et al. 2018a).

3 RESULTS

We focus on 1000 M_{\odot} star-forming regions, which typically contain approximately five massive stars ($M_{\star} > 15 M_{\odot}$) that act as photoionizing sources (cf. Scally & Clarke 2001). Our 1000 M_{\odot} cluster contains five massive stars; 43.2, 32.7, 25.7, and two 17 M_{\odot} with Φ_i values of 1.1, 0.47, 0.19, and ~ 0.013 , respectively. We focus on the results for two different initial stellar densities, ~ 10 and $\sim 100 M_{\odot} \text{pc}^{-3}$ and, apart from the final section, the assumed initial mass for every disc is $M_{\text{disc}} = 0.1 M_{\star}$.

We present the results from varying different initial properties within the star-forming regions, focusing on the spatial distribution (fractal dimension, D) and net bulk motion (virial ratio, α_{vir}). We focus on proplyds that have a radius of 100 au, however, the effects of external photoevaporation on discs with different radii and mass are discussed later in Sections 3.3 and 3.4.

We compare our results to the observed disc fractions in both Haisch et al. (2001) and Richert et al. (2018) using ages from the models in Siess, Dufour & Forestini (2000). We discuss the caveats associated with these models in Section 4.

We later present two low-mass clusters (100 M_{\odot}) with an initial density of $\sim 100 M_{\odot} \text{pc}^{-3}$ that are subvirial ($\alpha_{\text{vir}} = 0.3$) and highly substructured ($D = 1.6$). Our clusters contain either one (38 M_{\odot}) or two (42 and 23 M_{\odot}) massive stars. The corresponding Φ values are $\Phi = 0.76$ and $\Phi = 1.01, 0.11$, respectively.

3.1 Substructure in star-forming regions

We first present the results from four simulations of star-forming regions, where in each simulation the star-forming region has a different initial spatial distribution; $D = 1.6$ (highly substructured), $D = 2.0$ (moderately substructured), $D = 3.0$ (smooth), and a Plummer sphere spatial distribution. In all simulations, the star-forming region is subvirial ($\alpha_{\text{vir}} = 0.3$).

In Fig. 1, we show the average fraction of stars that have retained their (100 au) discs from 20 runs of each simulation, where the initial substructure of the star-forming region is varied. We present the results for regions with two different initial stellar densities 10 and 100 $M_{\odot} \text{pc}^{-3}$, respectively.

Fig. 1(a) shows the results from a star-forming region with an initial density of 10 $M_{\odot} \text{pc}^{-3}$. Within highly substructured regions ($D = 1.6$), the time taken for half of the stars to lose their discs due to external photoevaporation is 2.12 Myr. In moderately

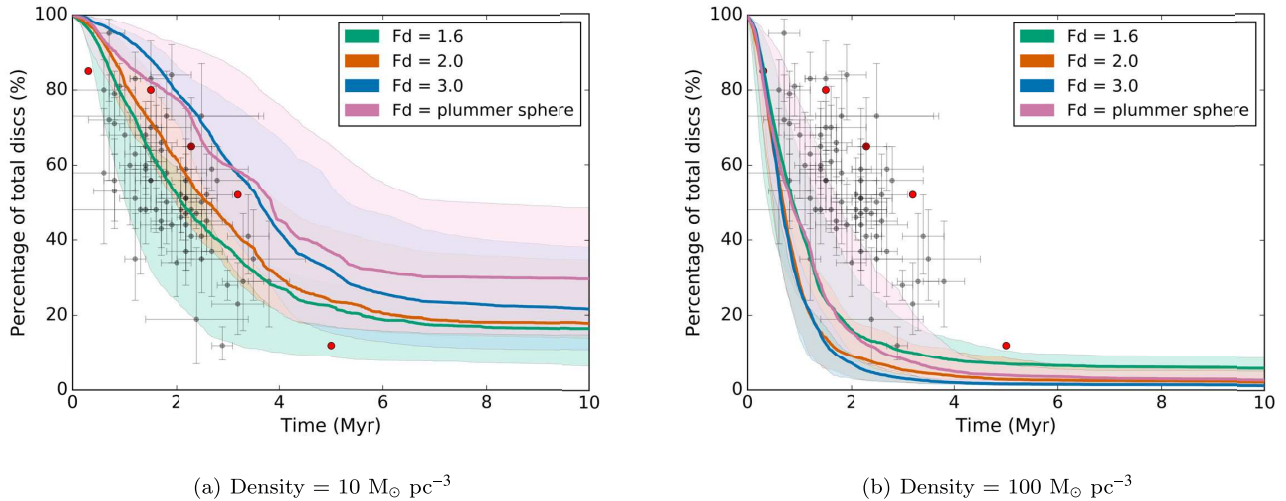


Figure 1. The average percentage of stars retaining their 100 au disc overtime within a subvirial ($\alpha_{\text{vir}} = 0.3$) cluster. The amount of substructure in the star-forming region is varied from highly substructured ($D = 1.6$) to smooth and centrally concentrated (Plummer sphere). Two different initial densities (10 and $100 M_{\odot} \text{ pc}^{-3}$) are considered. Each coloured line represents a different initial spatial distribution. The red data points are observational values from Haisch et al. (2001). The grey data points are from Richert et al. (2018) using stellar ages from the models in Siess et al. (2000). The coloured shaded regions show the complete range of values from the 20 runs for each set of initial conditions.

substructured regions ($D = 2.0$), this time increases to 2.60 Myr. However, the average percentage of remaining discs with time in both remain relatively similar throughout the 10 Myr. For regions with an initially smooth and spherical distribution ($D = 3.0$), the time taken for half of the discs to disperse is 3.62 Myr. Discs within Plummer spheres have the longest lifetimes (3.85 Myr), with an average of ~ 29.7 per cent of discs surviving for longer than 10 Myr.

Plummer (1911) models (and other models that describe smooth star clusters such as a King 1966 profile or an Elson, Fall & Freeman 1987 profile) are intended to model dynamically relaxed systems, whereas young star-forming regions are yet to relax. Therefore, even a smooth box fractal ($D = 3.0$) contains kinematic substructure, which causes the dynamical evolution of such a region to be more violent than a smooth Plummer sphere. It is therefore unsurprising that fewer discs survive in kinematically substructured fractal regions than in Plummer spheres with a similar density. Fig. 1(b) shows the results for star-forming regions with an initial density of $100 M_{\odot} \text{ pc}^{-3}$. For discs in the highly substructured regions ($D = 1.6$), the time taken for half of the stars within the cluster to lose their proplyds is 0.87 Myr. The majority of discs within the highly substructured region ($D = 1.6$) are dispersed after 10 Myr, with ~ 6 per cent surviving for the length of the simulation. The majority of discs within smooth, spherical regions are also dispersed within a short time frame, with only ~ 3 per cent remaining after 10 Myr. In Table 1, we summarize the average time taken for half of the stars in a star-forming region to lose their discs for each spatial distribution and Table 2 summarizes the percentage of discs remaining at the end of the 10 Myr simulation.

In the low-density simulations, regions with more spatial substructure photoevaporate discs faster than smoother regions of a comparable density (Fig. 1a). The reason for this is that the more substructured regions are initially further from dynamical equilibrium than the smooth regions, and low-mass stars undergo more close interactions with the high-mass stars as the regions relax.

Interestingly, in the high-density simulations (Fig. 1b), whilst the fraction of discs remaining after 10 Myr is lower than in the low-density simulations, the initially more substructured star-forming

Table 1. The time taken for half of stars within the star-forming region to lose the gas from their 100 au proplyds in a $1000 M_{\odot}$ subvirial ($\alpha_{\text{vir}} = 0.3$) region with two different initial densities; 10 and $100 M_{\odot} \text{ pc}^{-3}$. Four different spatial distributions are analysed; $D = 1.6$ (highly substructured), $D = 2.0$ (moderately substructured), $D = 3.0$ (smooth), and a Plummer sphere distribution. The highest and lowest values from the 20 different runs are included.

Fractal dimension (D)	Half-life of Proplyds	
	$\rho = 10 M_{\odot} \text{ pc}^{-3}$	$\rho = 100 M_{\odot} \text{ pc}^{-3}$
1.6	$2.12 \pm_{1.11}^{0.51}$ Myr	$0.87 \pm_{0.49}^{0.50}$ Myr
2.0	$2.60 \pm_{0.62}^{1.36}$ Myr	$0.67 \pm_{0.22}^{0.21}$ Myr
3.0	$3.62 \pm_{0.89}^{1.68}$ Myr	$0.65 \pm_{0.16}^{0.10}$ Myr
Plummer sphere	$3.85 \pm_{1.34}^{3.70}$ Myr	$0.84 \pm_{0.29}^{0.90}$ Myr

Table 2. The average percentage of 100 au discs remaining after 10 Myr within a subvirial ($\alpha_{\text{vir}} = 0.3$) star-forming region from 20 realizations of each simulation. The amount of substructure is varied from highly substructured ($D = 1.6$) to a smooth and centrally concentrated Plummer sphere. The highest and lowest values from the 20 different runs are included. Two different initial densities (10 and $100 M_{\odot} \text{ pc}^{-3}$) are considered.

Fractal dimension (D)	Percentage of discs remaining after 10 Myr	
	$\rho = 10 M_{\odot} \text{ pc}^{-3}$	$\rho = 100 M_{\odot} \text{ pc}^{-3}$
1.6	$16.40 \pm_{9.8}^{4.58}$ per cent	$5.99 \pm_{2.67}^{2.88}$ per cent
2.0	$17.75 \pm_{4.04}^{17.03}$ per cent	$2.27 \pm_{0.92}^{4.04}$ per cent
3.0	$21.60 \pm_{10.88}^{16.63}$ per cent	$1.35 \pm_{0.42}^{0.63}$ per cent
Plummer Sphere	$29.70 \pm_{14.98}^{18.95}$ per cent	$2.81 \pm_{1.63}^{2.36}$ per cent

regions contain more discs than the smooth regions after 10 Myr (and their disc-destruction half-life is longer, see Table 1). We attribute this to the higher ejection rate of massive stars in dense, substructured star-forming regions (Parker et al. 2014), which

Table 3. The time taken for half of stars within the cluster to lose the gas within their 100 au proplyds in a $1000 M_{\odot}$, moderately substructured ($D = 2.0$) star-forming region for two different initial densities: 10 and $100 M_{\odot} \text{pc}^{-3}$. Three different virial ratios are analysed: $\alpha_{\text{vir}} = 0.3$ (subvirial, or collapsing), $\alpha_{\text{vir}} = 0.5$ (virial equilibrium), and $\alpha_{\text{vir}} = 0.7$ (supervirial, or expanding).

Virial ratio (α_{vir})	Half-life of proplyds	
	$\rho = 10 M_{\odot} \text{pc}^{-3}$	$\rho = 100 M_{\odot} \text{pc}^{-3}$
0.3	$2.60 \pm_{0.62}^{1.36}$ Myr	$0.67 \pm_{0.22}^{0.21}$ Myr
0.5	$3.10 \pm_{1.15}^{2.73}$ Myr	$0.68 \pm_{0.35}^{0.66}$ Myr
0.7	$3.53 \pm_{1.40}^{1.72}$ Myr	$0.63 \pm_{0.23}^{0.57}$ Myr

Table 4. The average percentage from 20 runs of simulations of 100 au discs remaining after 10 Myr within a moderately substructured ($D = 2.0$) cluster. The bulk motion (virial ratio) of the star-forming region is varied, from collapsing (subvirial, $\alpha_{\text{vir}} = 0.3$) to expanding (supervirial, $\alpha_{\text{vir}} = 0.7$). The highest and lowest values from the 20 different runs are included.

Virial Ratio (α_{vir})	Percentage of discs remaining after 10 Myr	
	$\rho = 10 M_{\odot} \text{pc}^{-3}$	$\rho = 100 M_{\odot} \text{pc}^{-3}$
0.3	$17.75 \pm_{4.04}^{17.03}$ per cent	$2.27 \pm_{0.92}^{4.04}$ per cent
0.5	$29.77 \pm_{17.95}^{15.73}$ per cent	$4.16 \pm_{2.06}^{15.14}$ per cent
0.7	$32.67 \pm_{16.14}^{10.9}$ per cent	$5.00 \pm_{3.23}^{8.04}$ per cent

means that some of the ionizing sources are no longer near the majority of the proplyds as early as 1 Myr after the start of dynamical evolution.

3.2 Virial ratio

We explore how changing the net bulk motion of the star-forming region affects the rate of disc dispersal due to external photoevaporation. We run simulations of our star-forming region with three different virial ratios; 0.3 (subvirial, or collapsing), 0.5 (virial equilibrium), and 0.7 (supervirial, or expanding). We keep the fractal dimension constant, adopting $D = 2.0$, and as before we analyse simulations with two different initial densities: 10 and $100 M_{\odot} \text{pc}^{-3}$.

In Table 3, we summarize the average time taken for half of the stars in each region to lose their (100 au) discs for a given bulk virial ratio. Table 4 shows the percentage of discs remaining at the end of the 10 Myr simulation. Fig. 2 shows the average fraction of stars that have retained their discs from the 20 runs of each simulation for a star-forming region where we vary the initial bulk motion (virial ratio).

Fig. 2(a) shows the average mass-loss rate in a star-forming region with an initial density of $10 M_{\odot} \text{pc}^{-3}$. The time taken for half of the stars within a collapsing (subvirial) star-forming region to lose their discs is 2.60 Myr. In regions that are expanding (supervirial), this time increases to 3.53 Myr. The percentage of discs within the subvirial region after 10 Myr is 17.75 per cent, in comparison to discs within an expanding region where the percentage rises to 32.67 per cent.

The initial net bulk motion of low-density star-forming regions affects the amount of discs that are photoevaporated due to external radiation, with subvirial regions evaporating more discs at a faster rate than either virialized or supervirial regions.

Fig. 2(b) shows the results for a star-forming region with an initial density of $100 M_{\odot} \text{pc}^{-3}$. The time taken for half the stars within a collapsing region to lose their discs is 0.67 Myr. This time is similar for regions in virial equilibrium and expanding regions (0.68 and 0.63 Myr, respectively). The lower disc half-life for the supervirial regions could again be due to massive stars being ejected in the (sub)virial regions. The percentage of discs remaining after 10 Myr in subvirial star-forming regions is 2.27 per cent, whereas in regions where the net motion is expansive, this is increased to 5.00 per cent.

3.3 Disc radii

Here, we present the rates of disc dispersal for different initial disc radii in a star-forming region with two different initial densities (10 and $100 M_{\odot} \text{pc}^{-3}$). The region has a fractal dimension of $D = 2.0$ (moderately substructured) and a virial ratio of $\alpha_{\text{vir}} = 0.3$ (subvirial).

Fig. 3 shows the percentage of proplyds with initial radii ranging between 10–1000 au that have some remaining mass over 10 Myr in a $1000 M_{\odot}$ star-forming region with different initial stellar densities: $10 M_{\odot} \text{pc}^{-3}$ in Fig. 3(a) and $100 M_{\odot} \text{pc}^{-3}$ in Fig. 3(b).

In the lower density star-forming regions (Fig. 3a), the time taken for half of the 100 au discs to be completely photoevaporated is 2.60 Myr. Discs with radii of 10 au have much greater lifetimes, with an average of ~ 77 per cent of discs surviving the full length of the simulation. The majority of discs with very large radii (1000 au) are still depleted within very short time-scales. Disc depletion rates begin to switch off after ~ 4 Myr due to a combination of a large decrease in density of the star-forming region, which peaks at ~ 2 Myr, and the death of the most massive star at 4.33 Myr.

Fig. 3(b) shows that the majority (90 per cent or more) of discs with radii > 10 au are completely photoevaporated before the end of the 10 Myr simulation in moderately dense star-forming regions. The time taken for half of the stars in the region to lose their 100 au discs is 0.67 Myr. The vast majority of the largest discs (1000 au) are photoevaporated completely within 2 Myr, with half of the stars in the region losing their discs within < 0.1 Myr.

We also ran simulations for low-mass star-forming regions ($100 M_{\odot}$) with an initial density of $\sim 100 M_{\odot} \text{pc}^{-3}$. These low-mass regions contain two massive stars (42 and $23 M_{\odot}$), which represent an unusual sampling of the IMF (Parker & Goodwin 2007), but is observed in nature (e.g. γ^2 Vel, Jeffries et al. 2014). In these low-mass regions, half of discs with radii of 100 au dissipated in 0.51 Myr. This is comparable to the $1000 M_{\odot}$ regions with a similar density. We will further explore the effects of different stellar IMFs on disc dispersal in a future paper.

3.4 Disc masses

We have assumed that the disc masses are 10 per cent of the host star's mass, which is likely to be an overestimate and various studies suggest that the disc mass is as low as 1 per cent of the host star's mass (Weidenschilling 1977; Hayashi 1981; Andrews et al. 2013).

In Fig. 4, we show the results for a star-forming region with our two different initial densities (10 and $100 M_{\odot} \text{pc}^{-3}$, respectively), where the initial disc masses are set to $M_{\text{disc}} = 0.01 M_{\star}$.

Fig. 4(a) shows that on average the time taken for half of the stars within the low-density star-forming region to lose their 100 au disc is 0.71 Myr, less than half of the time taken for discs with 10 per cent of the mass of their stellar host. For discs with a radii of 10 au, the half-life is 3.31 Myr.

The time-scale for half of the 100 au discs to dissipate in the moderately dense ($100 M_{\odot} \text{pc}^{-3}$) star-forming region (see Fig. 4b)

Table 5. The time taken for half of stars in a star-forming region to lose the gas within their 100 au proplyds in a $1000 M_{\odot}$, moderately substructured ($D = 2.0$) region for two different initial densities: 10 and $100 M_{\odot} \text{pc}^{-3}$ and two different masses of disc, 10 per cent and 1 per cent. Three different virial ratios are analysed: $\alpha_{\text{vir}} = 0.3$ (subvirial, or collapsing), $\alpha_{\text{vir}} = 0.5$ (virial equilibrium), and $\alpha_{\text{vir}} = 0.7$ (super virial, or expanding). The highest and lowest values from the 20 different runs are included.

Disc radius (au)	Half-life of cluster proplyds			
	Disc mass = $0.1 M_{\star}$		Disc mass = $0.01 M_{\star}$	
	$\rho = 10 M_{\odot} \text{pc}^{-3}$	$\rho = 100 M_{\odot} \text{pc}^{-3}$	$\rho = 10 M_{\odot} \text{pc}^{-3}$	$\rho = 100 M_{\odot} \text{pc}^{-3}$
10	> 50 per cent remaining	$3.92 \pm_{1.87}^{5.91}$ Myr	$3.31 \pm_{0.9}^{5.16}$ Myr	$0.84 \pm_{0.27}^{0.22}$ Myr
50	$3.94 \pm_{0.82}^{5.96}$ Myr	$1.04 \pm_{0.24}^{0.40}$ Myr	$1.22 \pm_{0.35}^{1.18}$ Myr	$0.28 \pm_{0.09}^{0.07}$ Myr
100	$2.60 \pm_{0.62}^{1.36}$ Myr	$0.67 \pm_{0.22}^{0.21}$ Myr	$0.71 \pm_{0.21}^{0.96}$ Myr	$0.14 \pm_{0.05}^{0.05}$ Myr
200	$1.55 \pm_{0.44}^{1.34}$ Myr	$0.36 \pm_{0.11}^{0.09}$ Myr	$0.39 \pm_{0.12}^{0.58}$ Myr	$0.06 \pm_{0.02}^{0.03}$ Myr
1000	$0.37 \pm_{0.11}^{0.55}$ Myr	$0.06 \pm_{0.02}^{0.02}$ Myr	$0.15 \pm_{0.05}^{0.20}$ Myr	$0.02 \pm_{0.01}^{0.02}$ Myr

Table 6. The average percentage of 100 au discs remaining after 10 Myr within a moderately substructured ($D = 2.0$) star-forming region for two different initial densities, (10 and $100 M_{\odot} \text{pc}^{-3}$), with two different initial disc masses, 10 per cent and 1 per cent the mass of the host star. The bulk motion (virial ratio) of the star-forming region is varied, from collapsing (subvirial, $\alpha_{\text{vir}} = 0.3$) to expanding (supervirial, $\alpha_{\text{vir}} = 0.7$). The highest and lowest values from the 20 different runs are included.

Disc radius (au)	Percentage of discs remaining after 10 Myr			
	Disc mass = $0.1 M_{\star}$		Disc mass = $0.01 M_{\star}$	
	$\rho = 10 M_{\odot} \text{pc}^{-3}$	$\rho = 100 M_{\odot} \text{pc}^{-3}$	$\rho = 10 M_{\odot} \text{pc}^{-3}$	$\rho = 100 M_{\odot} \text{pc}^{-3}$
10	$77.29 \pm_{4.71}^{12.79}$ per cent	$39.49 \pm_{14.17}^{14.34}$ per cent	$28.81 \pm_{9.09}^{19.13}$ per cent	$3.99 \pm_{2.18}^{6.65}$ per cent
50	$34.44 \pm_{7.57}^{18.08}$ per cent	$6.73 \pm_{3.62}^{5.97}$ per cent	$4.79 \pm_{1.43}^{9.34}$ per cent	$0.40 \pm_{0.23}^{1.53}$ per cent
100	$17.75 \pm_{4.04}^{17.03}$ per cent	$2.27 \pm_{0.92}^{4.04}$ per cent	$2.27 \pm_{1.13}^{3.91}$ per cent	$0.08 \pm_{0.08}^{0.59}$ per cent
200	$7.12 \pm_{1.82}^{10.79}$ per cent	$0.74 \pm_{0.36}^{1.95}$ per cent	$1.09 \pm_{14.17}^{1.81}$ per cent	$0.00 \pm_{0.00}^{0.38}$ per cent
1000	$1.09 \pm_{0.9}^{1.73}$ per cent	$0.0 \pm_{0.00}^{0.38}$ per cent	$0.29 \pm_{0.29}^{0.51}$ per cent	$0.00 \pm_{0.00}^{0.04}$ per cent

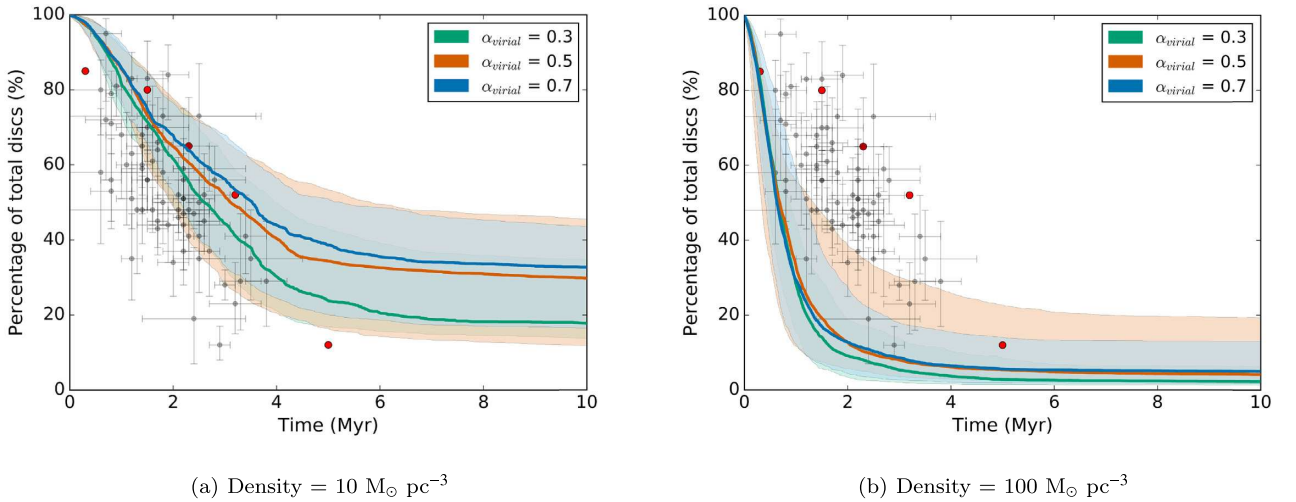


Figure 2. The average percentage of stars retaining their 100 au disc with time in a $1000 M_{\odot}$, moderately substructured ($D = 2.0$) star-forming region with an initial density of 10 and $100 M_{\odot} \text{pc}^{-3}$. Each coloured line represents a different virial ratio. The red data points are observational values from Haisch et al. (2001). The grey data points are from Richert et al. (2018) using ages from the stellar model in Siess et al. (2000). The shaded regions show all values between the maximum and minimum values from all 20 runs of the simulations.

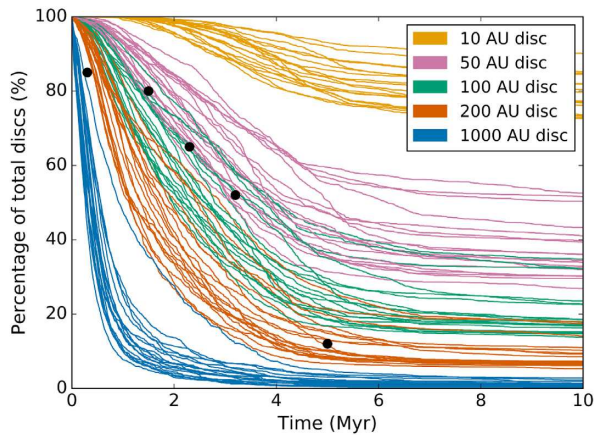
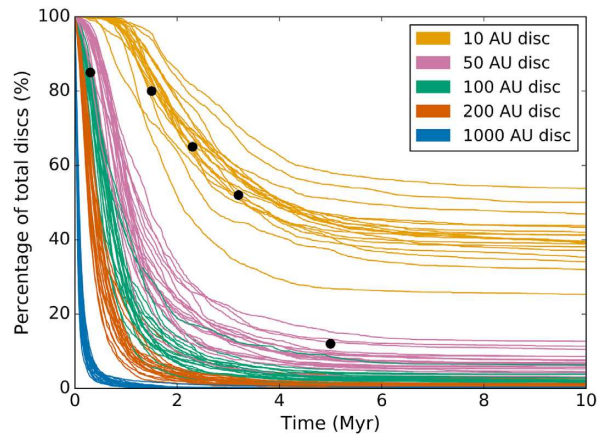
(a) Density = $10 M_{\odot} \text{ pc}^{-3}$ (b) Density = $100 M_{\odot} \text{ pc}^{-3}$

Figure 3. The percentage of total remaining discs overtime for a $1000 M_{\odot}$ star-forming region with an initial density of ~ 10 and $100 M_{\odot} \text{ pc}^{-3}$ (panels a and b, respectively). The cluster is moderately substructured ($D = 2.0$) and is subvirial ($\alpha_{\text{vir}} = 0.3$), and each colour represents a different initial disc radius. The disc masses are 10 per cent of the host star mass. The multiple coloured lines are each a single run of the 20 simulation runs. The black data points are observational values from Haisch et al. (2001).

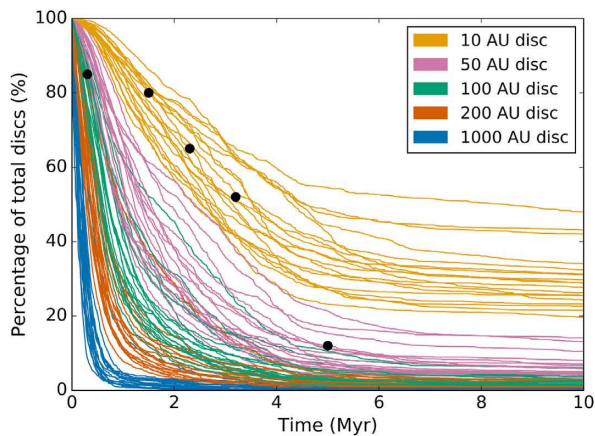
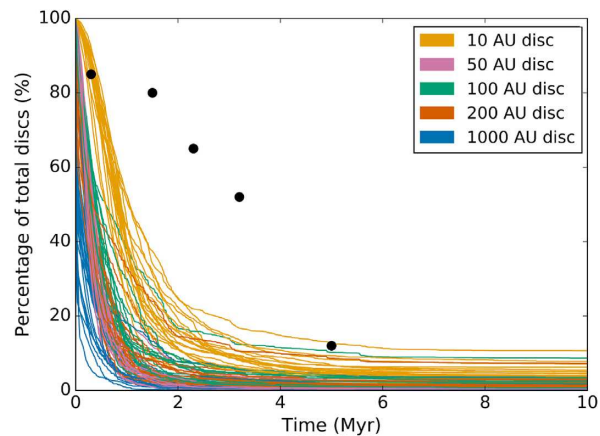
(a) Density = $10 M_{\odot} \text{ pc}^{-3}$ (b) Density = $100 M_{\odot} \text{ pc}^{-3}$

Figure 4. The percentage of total remaining discs overtime for a star-forming region of $1000 M_{\odot}$ with an initial density of ~ 10 and $100 M_{\odot} \text{ pc}^{-3}$, respectively, a fractal dimension of $D = 2.0$ and a virial ratio of $\alpha_{\text{vir}} = 0.3$. The initial disc masses are 1 per cent of the host star mass. Each colour represents a different disc radius. The multiple coloured lines are each a single run of the 20 simulation runs. The black data points are from observational values from Haisch et al. (2001).

is ~ 0.14 Myr. For discs with a radius of 10 au, the half-life is ~ 0.84 Myr. Less than 5 per cent of 10 au discs survive for more than 3 Myr.

3.5 Mass of star-forming regions

We also ran simulations for two different low-mass star-forming regions ($100 M_{\odot}$) with an initial density of $\sim 100 M_{\odot} \text{ pc}^{-3}$, which were subvirial ($\alpha_{\text{vir}} = 0.3$) and substructured ($D = 1.6$). These low-mass regions contain one ($38 M_{\odot}$) or two (42 and $23 M_{\odot}$) massive stars, which represent an unusual sampling of the IMF (Parker & Goodwin 2007), but is observed in nature (e.g. γ^2 Vel; Jeffries et al. 2014). Our expectation from randomly sampling the IMF is that 10 per cent of all star-forming regions can host a massive star, and 1 per cent of regions will host two massive stars. We note that the

lack of massive star(s) in *any* star-forming region would preclude disc destruction from photoevaporation, though as discussed in Section 2 it is unclear which type of star-forming region (in terms of total mass, M_{cl}) contributes the most (planet hosting) stars to the Galactic field.

In both of these low-mass regions, half of discs with radii of 100 au dissipated before ~ 1 Myr (Fig. 5). The time taken for half of the discs to be destroyed in a region with one massive star is 0.95 Myr. This time is reduced to 0.37 Myr for the cluster with two massive stars.

At the end of the 10 Myr simulation, 15.5 per cent of discs within the region with one massive star are surviving. Within the region containing two massive stars, less than 5 per cent of discs are remaining, double the number of discs remaining in higher mass regions with the same initial conditions.

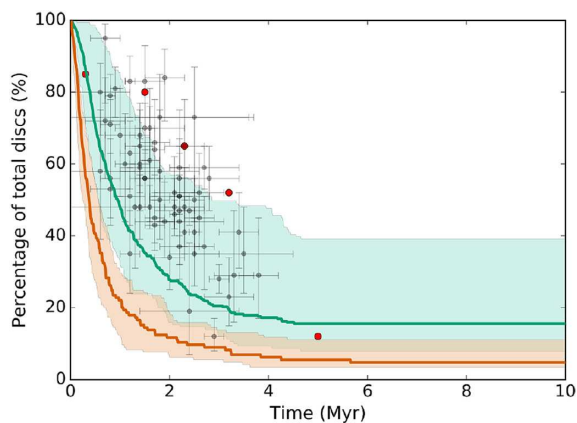


Figure 5. The median percentage of proplyds (100 au) remaining with time for two $100 M_{\odot}$ clusters with initial densities of $100 M_{\odot} \text{pc}^{-3}$ but different numbers of massive stars. The green line shows values for a cluster with one massive star ($>15 M_{\odot}$) and the orange line a cluster with two massive stars. The red data points are observational values from Haisch et al. (2001). The grey data points are from Richert et al. (2018) using ages from the stellar model in Siess et al. (2000). The coloured shaded regions show the complete range of values from the 20 runs for each set of the different clusters.

4 DISCUSSION

The initial conditions of a star-forming region will affect the rate at which proplyds are photoevaporated due to the radiation from nearby massive stars. The initial substructure and net bulk motion of a star-forming region impacts the rate of disc dispersal.

4.1 Changing the initial conditions of star-forming regions

In our low-density simulations, highly substructured ($D = 1.6$) regions disperse half of the proplyds within 1.51 Myr, more than twice as fast as smooth ($D = 3.0$) regions. In simulations with a Plummer sphere distribution, more than 30 per cent of discs remain at the end of the 10 Myr simulation, almost double that of discs within highly substructured clusters.

At these low densities, the degree of substructure matters because a more substructured star-forming region is further from dynamical equilibrium than a smooth region. When this occurs, a low-mass star is likely to have more close encounters with a massive ionizing star than in a smooth region.

In moderately dense initial conditions ($100 M_{\odot} \text{pc}^{-3}$), the difference in the fraction of discs that are photoevaporated between different initial spatial distributions decreases greatly, although regions with a Plummer sphere distribution retain more of their discs than regions with initial substructure. However, the average of all runs indicates that the amount of initial substructure has little effect on the survival rates of discs at these densities and fewer than 50 per cent of discs remain after 1 Myr.

The effect of changing the net bulk motion of the star-forming region has a similar impact on the rate of disc dispersal as the initial substructure has. For low-density regions ($10 M_{\odot} \text{pc}^{-3}$), the difference between the amount of discs surviving within a collapsing and an expanding star-forming region is ~ 15 per cent, with the collapsing regions enabling more photoevaporation than in expanding regions. Again, approximately double the number of discs remain in expanding clusters than in collapsing clusters. In moderately dense clusters it is similar, with the difference being ~ 3 per cent.

For low-mass star-forming regions ($100 M_{\odot}$), disc dispersal rates are similar to those of in higher mass regions. Whilst the UV field strength can vary due to different realizations of the IMF (Fatuzzo & Adams 2008), these low-mass regions show that the mere presence of a high-mass star ($>15 M_{\odot}$) will cause disc lifetimes to be shortened dramatically.

Our simulations are set up to mimic the observations of star formation in filaments, where the pre-stellar cores have subvirial motion (Larson 1981; Foster et al. 2015). The local velocity dispersion is therefore always subvirial to some degree and because mass-loss due to photoevaporation is so fast (equations 2 and 3), most of the photoevaporation occurs during the substructured phase of a star-forming region.

Previous studies investigating the effects of external photoevaporation on disc dispersal rates assumed smooth and centrally concentrated spatial distributions (Sclally & Clarke 2001; Adams et al. 2004, 2006; Winter et al. 2018b), replicating environments such as the present-day conditions of the ONC. However, using the present-day spatial and kinematic distributions to model star clusters may not accurately replicate the dynamical history of the star-forming region from which the cluster formed (Parker et al. 2014).

We cannot provide a direct comparison with the work by Fatuzzo & Adams (2008) and Winter et al. (2018b) because these authors assume initially smooth initial conditions, and we have shown that the severity of photoevaporation depends on the degree of initial substructure, as well as the initial positions of the most massive stars. We have distributed the massive stars randomly within the substructure, and after dynamical evolution these massive stars migrate towards the centre of the star cluster as it forms (Allison et al. 2010; Parker et al. 2014, 2016).

Furthermore, we have used the photoevaporation prescription from Sclally & Clarke (2001), rather than determine the photoevaporation rate from the EUV/FUV fluxes as a function of the flux in the interstellar medium (the so-called G_0 value, $1.6 \times 10^{-3} \text{erg s}^{-1} \text{cm}^{-2}$; Habing 1968). In comparison with Sclally & Clarke (2001), we find that discs are destroyed earlier in ONC-type regions because the initial densities are higher (in line with current observations; Parker 2014), and the star-forming regions are substructured (Cartwright & Whitworth 2004).

An initially highly substructured star-forming region can become smooth and centrally concentrated within a few Myr due to a combination of violent and two-body relaxation. Proplyds in these highly substructured environments will be photoevaporated at faster rates than discs within initially smooth regions. Even though they will both appear smooth within a few Myr, the percentage of discs remaining, and possibly the population of planets within the regions, will vary greatly.

The initial density of the cluster has the largest effect on the disc dispersal rate due to external photoevaporation. The ‘moderately dense’ clusters reflect the likely initial densities of many star-forming regions (Parker 2014). However, we find that these ‘moderately dense’ environments are very destructive for proplyds and evaporate discs at rates faster than suggested by observations (compare the black points in Figs 1 and 2 with our simulated data). Our results suggest that proplyds (or at least their gas content) would always be significantly depleted in moderately dense ($100 M_{\odot} \text{pc}^{-3}$) star-forming regions, if those regions contain massive stars.

Haisch et al. (2001) finds that the fraction of disc-hosting stars in young star-forming regions falls to 50 per cent after ~ 3 Myr, whereas Richert et al. (2018) find that after only ~ 2 Myr half of

the discs remain in their observed regions. In comparison, more than half of the discs in our simulations that are within dense environments are destroyed within ~ 1 Myr. One interesting data point in the Haisch et al. (2001) sample is the ONC. With an age of ~ 1 Myr, the centre of the ONC contains four massive stars and a density of $\sim 400 M_{\odot} \text{pc}^{-3}$ (Hillenbrand & Hartmann 1998). Observations of this part of the ONC suggest that ~ 80 – 85 per cent of stars within the cluster are surrounded by bright ionization fronts, interpreted to be discs, with radii of ~ 1000 au (Bally et al. 2000; Lada et al. 2000).

The age of our simulated regions, where 80 per cent of stars still possess a 100 au disc with some mass, is 0.48 Myr – likely to be less than half the age of the ONC. This suggests that the massive stars within the ONC should have destroyed the majority of 100 au discs. From similar arguments, Clarke (2007) concluded that the possible discs in the ONC with radii > 10 au are likely to be merely ionization fronts, containing little mass. Our simulations with different initial disc radii show that the radius of the disc will greatly affect the rate at which it is photoevaporated (see also Clarke 2007) due to the dependence on disc radii within the FUV and EUV photoevaporation prescriptions.

Recent surveys suggest that most stars in the Galactic field host planets, and many of these are gas or ice giants (Mayor et al. 2011). This implies that the majority of planet-forming discs were able to survive a significant amount of time in their birth environment. Our simulations suggest that this is only possible in low-density regions that contain no photoionizing sources (i.e. massive stars). Therefore, (giant) planet formation must occur on very rapid time-scales (< 1 Myr), or stars that host giant planets must have formed in very benign environments.

Many observed proplyds are located in low-mass, low-density star-forming regions (Andrews et al. 2013; Ansdell et al. 2018) and would be unaffected by external photoevaporation. However, many star-forming regions are typically moderate-density ($\sim 100 M_{\odot} \text{pc}^{-3}$) environments (Parker 2014), and our results suggest that the majority of Proplyds in star-forming regions with these densities do not survive for long enough periods of time to form giant planets.

4.2 Caveats

There are several caveats to our results, which we discuss next.

The effects of external EUV radiation on proplyds can be reduced when thick winds are present, caused by FUV heating of the disc (Alexander et al. 2014). However, the majority of the disc mass-loss occurs due to FUV radiation. We repeated our analysis without EUV photoevaporation and find the disc dispersal rates to be similar.

It is possible that we are overestimating the amount of photoevaporation from massive stars. However, recent research suggests that the prescriptions used here are actually underestimating the amount of FUV radiation that discs receive (Facchini et al. 2016; Haworth et al. 2018a). As FUV is the dominant source of external photoevaporation, the proplyds in our simulations could dissipate on even shorter time-scales.

Star formation is an inherently inefficient process, with typically only ~ 30 per cent of the mass of a giant molecular cloud converted into stars. Young star-forming regions are observed to contain a large amount of dust and gas, which could shield the proplyds from significant photoevaporation. At these early stages, the stellar density within the substructure is highest, and is therefore when the largest percentage of stars are in closest proximity to the massive stars. However, hydrodynamical simulations of star-forming regions

show that massive stars blow large ($\sim \text{pc}$ -scale) cavities within the gas on short time-scales (Dale, Ercolano & Bonnell 2013), and so low-mass disc-hosting stars that would be affected by EUV/FUV radiation will likely reside in the cavities blown out by the massive stars. If the gas and dust could shield the disc, this would protect them for a very short period of time (Gorti, Hollenbach & Dullemond 2015). Whether this grace period would be long enough to allow gas-rich giant planets to form is uncertain.

Given that most star-forming regions have stellar densities above a few $M_{\odot} \text{pc}^{-3}$ (Bressert et al. 2010), external photoevaporation will detrimentally affect proplyds in any star-forming region that contains massive stars. This implies that star-forming regions that do not contain massive stars are more likely to form giant planets, but we note that massive stars appear necessary in order to deliver short-lived radioisotopes to the young Solar system (Lugaro, Ott & Kereszturi 2018). The number of massive stars in a star-forming region appears to only be limited by the mass of the star-forming cloud (Parker & Goodwin 2007), but this also means that low-mass star-forming regions ($< 10^4 M_{\odot}$) stochastically sample the IMF, meaning that our simulations cannot be described as ‘typical’ star-forming regions.

Quantifying disc dispersal is further complicated by how difficult it is to determine the ages of young stars, especially before 1 Myr (Siess et al. 2000). We use the stellar ages from the Siess et al. (2000) model. However, models of pre-main sequence stellar evolution calculate different ages depending on the physics that is implemented. Of the three models presented in Richert et al. (2018), we use the ages from Siess et al. (2000) so that we are comparing the lower end of cluster ages to our simulations. The average stellar age calculated for the clusters in Richert et al. (2018) is significantly shorter than in more recent models from Feiden (2016). Using these lower age limits, we more than halve the possible average lifetimes of the discs within the observed clusters.

Furthermore, recent work by Bell et al. (2013) suggests that the ages of pre-main sequence stars may be underestimated by a factor of 2, meaning that the observed discs (e.g. Haisch et al. 2001) could be a factor of 2 older. This would make it even more surprising that discs would remain around low-mass stars, if those stars form in regions containing massive stars.

There is also the question of how quickly the photoionizing massive stars form. In the competitive accretion models (Bonnell et al. 2001), massive stars gradually gain in mass over ~ 1 Myr (Wang et al. 2010), suggesting high-mass stars form later than low-mass stars (Tan et al. 2014), which would in turn decrease the amount of time low-mass stars spend near the photoionizing sources (Dale, Ercolano & Bonnell 2012; Dale et al. 2014). In our simulations all stars form simultaneously, and therefore the disc-hosting low-mass stars do not have this grace period, which would increase disc lifetimes.

The growth of planetesimals into planets can be greatly accelerated by the accretion of cm-scale pebbles. Johansen & Lambrechts (2017) show that once a $10^{-2} M_{\oplus}$ planetesimal has formed it can grow to Jupiter mass in 1 Myr when starting as far out as about 15 au. An initial phase of accreting pebbles forms a $10 M_{\oplus}$ core in about 0.8 Myr, which then undergoes runaway gas accretion to reach Jupiter mass. Such processes potentially allow close-in giant planets to be formed even in the relatively hostile conditions that we consider here.

However, photoevaporation by the central star can cause large amounts of mass-loss in the inner disc, potentially affecting giant planet formation (Alexander et al. 2014). Grain size also has a significant effect on disc dispersal rates. Mass-loss occurs much

more quickly when grain growth has occurred because the FUV radiation can penetrate deeper into the disc (Facchini et al. 2016).

Discs that can survive in moderately dense environments have small radii (10–50 au). This is because of the disc radius dependence in the external photoevaporation prescriptions. Internal UV radiation can cause significant mass-loss and erosion of the disc within short time-scales (1 Myr, Gorti & Hollenbach 2008). The time-scale for internal disc dispersal is very short (10^5 yr), with a UV switch being triggered due to the slowing of accretion on to the inner 10 au of the disc (Clarke et al. 2001), also calling into question the survivability of small discs.

Our disc radii are fixed, but in reality disc radii change with time, often in an inside-out fashion where the initial radius is small (and not as susceptible to photoevaporation) compared to later in the disc’s life. We include several different disc radii to help visualize what happens for different disc initial conditions, but we cannot model the full viscous evolution in our post-processing analysis.

In our simulations, we have elected to keep the stellar IMF constant across different realizations of the spatial and kinematic initial conditions of our star-forming regions. The reasons for this are two-fold. First, we wish to isolate the possible effects of stochastic dynamical evolution (Allison et al. 2010; Parker & Goodwin 2012; Parker et al. 2014), which could lead to different photoevaporation rates even if the ionizing flux from massive stars were kept constant. The uncertainties shown by the shaded regions in Figs 1, 2, and 5 show this stochasticity for the same initial conditions. Secondly, the photoevaporation prescriptions we adopt (following Scally & Clarke 2001) are actually quite insensitive to the mass of the most massive stars (but rather depend on whether the massive stars are present or not).

However, Fatuzzo & Adams (2008) show that the FUV and EUV fluxes can vary if the stellar IMF is extremely top-heavy and contains more massive stars than expected on average. In a forthcoming paper, we will calculate the EUV and FUV fluxes in our substructured star-forming regions and use the recent FRIED models of disc photoevaporation from Haworth et al. (2018b) to determine mass-loss based on these models, and whether it depends strongly on stochastic sampling of the stellar IMF.

Similarly, in some of our simulations the massive stars are ejected early on, which is a common occurrence in simulations of dense star-forming regions (Allison et al. 2010; Oh & Kroupa 2016, Schoettler et al., in preparation). We will also quantify the effects of these ejections on the fraction of surviving proplyds in an upcoming paper.

The majority of discs observed with ALMA have been located in low-mass, low-density star-forming regions. Current observations suggest that the majority of stars form in moderately dense ($\sim 100 M_{\odot} \text{pc}^{-3}$) environments (Parker 2014). However, the majority of proplyds in clusters with these densities do not survive for long enough periods of time to form planets, as planet formation is thought to take place over a few million years (Pollack et al. 1996). The fact that the majority of stars have planetary systems around them poses important questions as a result of the discrepancies that seemingly arise. This may indicate that the majority of stars form in low-mass clusters where there are few to no high-mass stars.

We adopt initial disc masses that are 10 per cent of the mass of the host star, which is likely to be a large overestimate. When looking at more realistic values (1 per cent), discs are destroyed on even shorter time-scales. However, it should be noted that accretion and internal photoevaporation will have much larger effects on disc mass evolution for these lower mass discs.

5 CONCLUSIONS

We have calculated the mass-loss due to external photoevaporation of proplyds in N -body simulations of the evolution of star-forming regions. We ran a suite of simulations where we vary the initial spatial structure, bulk motion, and initial density of the regions. We compared our simulations that more closely represent observed star-forming regions (subvirial, substructured) with those of primordially mass segregated, spherical clusters, similar to those used in previous studies of external photoevaporation.

The parameter that most affects rates of disc dispersal is the initial density of the star-forming region. The majority of proplyds within simulated regions that mimic the conditions in nearby star-forming regions are dispersed due to external photoevaporation within very short time-scales. In moderately dense ($\sim 100 M_{\odot} \text{pc}^{-3}$) star-forming regions that have moderate levels of substructure ($D = 2.0$) and are collapsing ($\alpha_{\text{vir}} = 0.3$), we find the time taken for half of 100 au discs to dissipate is significantly shorter (three times less) than suggested in observational studies (Haisch et al. 2001). Lower density clusters ($\sim 10 M_{\odot} \text{pc}^{-3}$) allow discs to survive long enough to match observations of disc lifetimes, although the half-life of 100 au discs is still less than that found by Haisch et al. (2001).

The initial spatial distribution of the star-forming region also affects the rate of proplyd dispersal due to external photoevaporation. The degree to which initial substructure affects disc dispersal rates depends on the initial density. In moderately dense ($\sim 100 M_{\odot} \text{pc}^{-3}$) regions, the effects are washed out, but in lower density regions ($\sim 10 M_{\odot} \text{pc}^{-3}$) we find that the more fractal and clumpy a star-forming region is, the higher the rate of disc dispersal. This is due to violent relaxation and the rapid increase in density (sometimes up to an order of a magnitude) of the star-forming region within a short amount of time. As most star-forming regions appear to have a high degree of substructure, it is important for future studies of disc dispersal to take the initial conditions into consideration due to external photoevaporation in dense environments.

The virial ratio of the star-forming region affects the rate of disc dispersal in a similar way to substructure. Regions that have a low initial density and are collapsing photoevaporate more discs on average than clusters that are expanding. The effects of varying the initial net bulk motion in moderately density clusters is negligible.

The majority of observed stars in the Galactic field host planetary systems, implying their proplyds survived long enough for formation to take place. There are three possible scenarios to resolve this apparent tension between observations and our simulations:

(i) The majority of planets may not form in moderately dense star-forming regions ($\sim 100 M_{\odot} \text{pc}^{-3}$); rather, they would form in low-density regions with no photoionizing massive stars present. Many of the proplyds have been observed in these low-density ambient environments (Ansdell et al. 2018), but significant numbers of proplyds (or at least their remnants) have been observed in dense, hostile regions such as the ONC (McCaughrean & O’dell 1996).

(ii) If some planets do form in dense, clustered environments containing massive stars (such as the ONC), then this suggests that giant planet formation must happen on very short time-scales (less than 1–2 Myr), or be confined to discs with radii significantly smaller than the orbit of Neptune in our Solar system. Johansen & Lambrechts (2017) show that giant planet formation can occur on these time-scales once large enough planetesimals have formed. However, internal photoevaporation processes can deplete the inner disc and set limits on the formation time of giant planets (Alexander et al. 2014).

(iii) The current calculations of mass-loss in discs due to external photoevaporation are severely overestimating the detrimental effects of EUV and FUV radiation. However, recent research (Facchini et al. 2016; Haworth et al. 2018a) suggests that photoevaporative mass-loss rates caused by FUV radiation may be underestimated, and our calculations also underestimated the effects as we adopt conservatively high initial disc masses.

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