

# Investigation on the hydrodynamic scaling effect of an OWC type wave energy device using experiment and CFD simulation

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## Abstract

This paper presents a study of the effect of model scale on the performance of a fixed Oscillating Water Column (OWC) type Wave Energy Converter (WEC). Tank tests at two different scales, including the effect of scaling of the test tanks to minimise the bias introduced by different wave blockage effects. CFD simulations based on Reynolds Average Navier Stokes (RANS) method were then carried out for both scaled OWCs to investigate whether CFD simulation is able to reproduce the scale effect. Comparison between the tank test results and the CFD simulation results suggests that CFD simulation is capable of reproducing the hydrodynamic scaling effect with a good accuracy. Results also suggest that the hydrodynamic scaling effect is mainly introduced by the Reynolds number effect for cases investigated in the current study.

*Keywords:* Tank test, CFD simulation, Scale effect, Wave energy, Oscillating Water Column;

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## 1. Introduction

Being one of the promising renewable energy technologies, WECs have attracted worldwide attention during the last few decades as one of the more promising marine renewable energy technologies. Detailed reviews of wave

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5 energy technologies can be found in several studies, for instance, [1],[2] and  
6 [3]. Among all the proposed WEC technologies, the Oscillating Water Column  
7 (OWC) type WEC is probably one of the most extensively studied technologies  
8 due to its simple working principle [2].

9       Along with tank test, numerical studies of OWC type WECs have played an  
10 important role in accelerating the evolution of OWC technology. For example,  
11 Evans [4] derived the theoretical maximum efficiency for a 2D fixed symmetric  
12 OWC device by assuming that the OWC surface moves as a weightless rigid  
13 piston. Later Sarmiento and Falcão [5] improved the theory by allowing OWC  
14 surface variations using a surface pressure method and wave flume experiments  
15 were carried out to validate the surface pressure theory [6]. With the help  
16 of rapid development of computer technology, researchers started to simulate  
17 the OWC problem with more advanced methods, such as BEM [7] and CFD  
18 [8]. Different aspects of the OWC technology have been extensively studied by  
19 several researchers, such as hydrodynamic performance [9], optimisation of the  
20 OWC geometry [10] and optimisation of turbine-induced damping [11] etc.

21       Although significant progress on the development and understanding of the  
22 OWC technology has been made recently, there are still several challenges to  
23 overcome in performance prediction. The effect of model scale is probably one  
24 of the critical issues in the early development stage, since the assessment of  
25 the full-scale device performance is normally extrapolated from a model scale  
26 experiment or simulation result at the early stage. To fill the gaps in theory  
27 and guidelines for the requirements of scale testing of a WEC, Sheng et al. [12]  
28 presented a theoretical analysis on the scaling of physical modelling and power  
29 take-off system. In order to minimise viscous effects, it was recommended that  
30 a physical model test shall be carried out with critical Reynolds number above  
31 about  $10^5$ . This requirement, however, can not always be fulfilled especially for  
32 tests in relatively small wave tanks since the scale (hence Reynolds number) of  
33 the test is normally constrained by the tank size. In contrast, numerical simula-  
34 tion methods such as CFD do not have the same limitations of scale. Recently,  
35 Elhanafi et al. used an experiment-validated CFD model to investigate the air

36 compressibility effect at full scale [13]. Although those simulation works are val-  
37 idated against scaled tank test and excellent agreement between the simulation  
38 results and experiment results is achieved, there are, however, few published  
39 multi-scale tank test data which can validate the capability of simulation tools'  
40 to reproduce the aerodynamic and hydrodynamic scale effects.

41 Recently, Viviano et al. [14] tested a generalized small scale OWC and  
42 the results were compared with a similar large scale model to investigate the  
43 scale effect. In their study, OWC devices have the same width as the tank  
44 width. Therefore, 3D radiation and diffraction effect was excluded. This paper  
45 investigates two different scale tank tests of an idealized 3D OWC device. Cor-  
46 responding CFD simulations are then performed to investigate whether CFD  
47 simulation is capable of capturing the hydrodynamic scale effect. This work  
48 is structured as below: Section 2 describes the experimental work including a  
49 discussion of the uncertainty and error source. Corresponding CFD simulations  
50 are described in Section 3. Section 4 compares the results obtained from the  
51 tank tests and CFD. Conclusion and future works are summarized in Section 5.

## 52 **2. Physical experiments**

53 Offshore structures (such as offshore platforms) are generally designed in  
54 such a way so that the interaction with waves is small. Guidance for these  
55 structures on wave blockage may therefore not be well-suited to WECs which  
56 are designed to have maximum wave structure interaction. Wave blockage in this  
57 context refers to all hydrodynamic effects related to the transverse constraints  
58 of the tank walls on the hydrodynamic response – including wave reflection from  
59 the tank walls and local variation in flow velocity caused by reduced cross section  
60 area due to the presence of the model. Therefore, the impact of wave blockage  
61 on results should be carefully considered [15], especially when comparing the  
62 performance of two different scale devices, since the effect of wall reflections  
63 and flow variations may be confused with the scale effect.

64 *2.1. Facilities*

65 In order to minimise the bias from different wave blockage introduced by  
66 different tank widths, experiments were carried out in the Kelvin Hydrodynamic  
67 laboratory and the Henry Dyer Hydrodynamic Laboratory of the University of  
68 Strathclyde as shown in Figure 1. The Kelvin Hydrodynamic Laboratory has a  
69 dimension of  $76\text{ m} \times 4.6\text{ m} \times 2.5\text{ m}$  with water depth of 2.1 m and the Henry Dyer  
70 Hydrodynamic Laboratory has a dimension of  $21.6\text{ m} \times 1.53\text{ m} \times 1\text{ m}$  with water  
71 depth set to 0.7 m. Since the cross section dimension govern wave blockage, it is  
72 anticipated that these two tanks will provide similar wave blockage effect when  
73 the two models have a scale ratio of 3:1. Both tanks are equipped with flap  
74 type wave makers and a wave absorbing beach. For convenience, Kelvin tank  
75 is denoted as the large tank and Henry Dyer tank is denoted as the small tank  
76 hereafter.

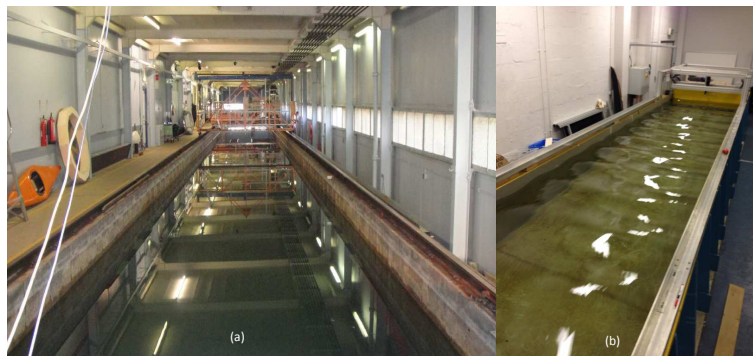


Figure 1: (a) Kelvin Hydrodynamic laboratory. (b) Henry Dyer Hydrodynamic laboratory.

77 *2.2. OWC device*

78 Simple acrylic hollow cylinders were selected to model an idealized OWC  
79 device for further investigation. Such a simple geometry allows easy scaling  
80 of the air compressibility by simply keeping the height of the air chamber the  
81 same for both scales [12]. A smooth plastic ring collar was fitted to the bottom  
82 of the device in order to have a better control of the sharp corners during the  
83 geometry scaling process and at the same time. (see Figure 2 for detail.). The

84 Power Take Off (PTO) system was modelled using an orifice plate to simulate  
 85 an idealised impulse turbine, because it has approximately quadratic pressure-  
 86 flow rate characteristics. This method of modelling the PTO has been used by  
 87 several researchers, for example, [13] and [16]). Instead of manufacturing several  
 88 different size orifice plates, 8 equal size and equally spaced circular openings were  
 89 drilled into the covering lid. By choosing different number of orifices open to  
 90 the air, different levels of damping could be achieved. More detailed geometry  
 91 information can be found in Table 1.

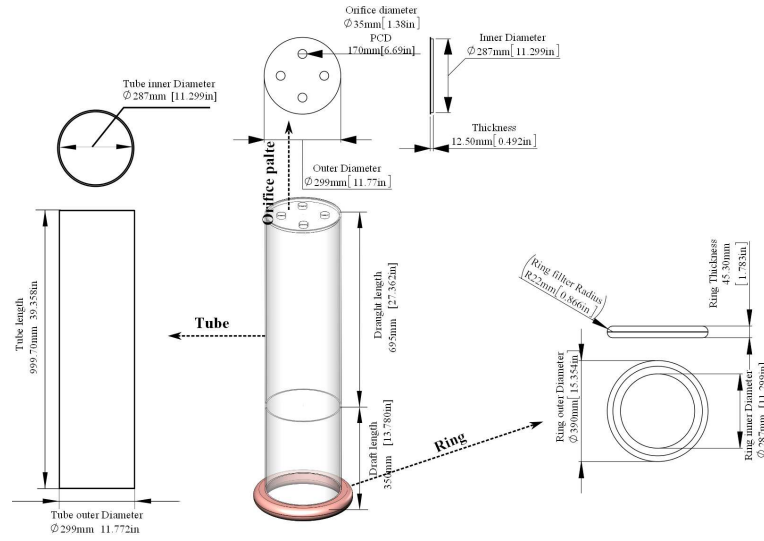


Figure 2: CAD illustration of the large scale device.

### 92 2.3. OWC performance and testing procedures

93 When assessing the performance of the OWC device, it is critical to assess  
 94 the available wave power from the incident wave. Conventionally, a reference  
 95 wave probe is located some distance in front of the device (i.e. between device  
 96 and wave maker) to measure the incident wave. That measured incident wave  
 97 information may be different from the wave arriving at the device due to spatial  
 98 variations of waves in the tank and effects of wave decay. Besides, the wave  
 99 measured by the reference probe may include the waves due to radiation (from

Table 1: Geometry details of the two OWC devices, scale ratio 1:3.

Component	Parameters (mm)	Large scale	Small scale
OWC model	total length	1045.0	808.5
	draft	350.0	116.7
Orifice plate	Plate Diameter	299.0	100.0
	Thickness	12.5	4.0
	Orifices Diameter	35.0	11.6
	Orifices position (PCD)	170	56.5
Tube	Outer Diameter	299.0	100.0
	Inner Diameter	287.0	96.0
Ring	Inner Diameter	299.5	100.5
	Outer Diameter	390.0	130
	Thickness	45.3	15.1
	Fillet radius	22.0	7.5

100 wave and OWC interaction)and scattering as well as the incident wave. There-  
 101 fore, in the present work, taking the advantage of the high level of repeatability  
 102 of the wave makers, the waves were first calibrated at the target location where  
 103 the devices would be deployed prior to installation of the model. The incident  
 104 average wave energy flux ( $P_{avail}$ ) can then be determined by the calibrated wave  
 105 information through

$$P_{avail} = \frac{1}{2} \rho g A^2 C_g \quad (1)$$

106

107 where  $\rho$  is the density of water,  $g$  is the gravitational acceleration,  $A$  is the  
 108 measured wave amplitude and  $C_g$  is the wave group velocity defined as

$$C_g = \frac{1}{2} \frac{\omega}{k} \left( 1 + \frac{2kh}{\sinh(2kh)} \right) \quad (2)$$

109 The  $\omega, k, h$  in Equation 2.3 are the circular wave frequency, the wave number  
 110 and the water depth, respectively.

111 Devices were then fixed in the center-line of each tank. Regular waves with  
 112 non-dimensional frequencies ( $Kh$ , here  $K = \omega^2/g$ ) from 2 to 8 with constant  
 113 wave height (0.06m for the large scale test and 0.02 for the small scale test.)  
 114 were then tested. The mean captured power by the OWC device is calculated  
 115 via

$$P = \frac{1}{T} \int_0^T \Delta p(t)q(t)dt \quad (3)$$

116 where  $T$  is the wave period and  $\Delta p(t)$  is the instantaneous pressure difference  
 117 across the orifice plate. This is measured by a differential pressure transducer  
 118 installed on the top of the orifice plate. A Honeywell 163PC0D75 ( $\pm 622.27$  Pa)  
 119 low pressure differential transducer was used for the large scale tests and a SEN-  
 120 SIRION SDP1000-L025 ( $\pm 62$  Pa) low differential pressure sensor was employed  
 121 to measure the pressure for the small scale tests,  $q(t)$  in Equation 2.3 is the  
 122 instantaneous volume flow rate driven by the water column and is defined by

$$q(t) = A_w \frac{\partial \eta}{\partial t} \quad (4)$$

123 where  $A_w$  is the cross section area of the OWC and  $\eta$  is the OWC elevation  
 124 measured by wave probes located in the middle of each device.

125 To compare the performance of the two devices, the so called capture factor  
 126 (capture width ratio) is introduced, defined as

$$Cf = \frac{P}{P_{avail}D_{out}} \quad (5)$$

127 where  $D_{out}$  is the characteristic length of the WEC device. In this case,  $D_{out}$   
 128 is the outer diameter of the OWC device (tube).

#### 129 *2.4. Experiment uncertainty and error*

130 Uncertainty analysis was performed in line with International Towing Tank  
 131 recommendation and guidelines ([17],[18]) in this paper. The main uncertainty  
 132 source in the test comes from the instruments used for measurements. This kind  
 133 of uncertainty (Type B, or systematic uncertainty) can be quantified through

134 instrument calibration or stated by the manufacturer. Combined with Type A  
135 (random uncertainty) obtained from repeated tests, uncertainty in the physical  
136 quantity of interested (e.g. mean captured power) can be calculated using un-  
137 certainty propagation analysis. For example, the total uncertainty in the peak  
138 mean power captured (1.1 W) caused by the pressure and volume flow rate mea-  
139 surement is 0.047 W for the large scale test. Detailed information on uncertainty  
140 analysis can be found in the above references. In the present study, the results  
141 of uncertainty analysis will be presented via error bars showing 95% confidence  
142 intervals with testing results.

143       Apart from the uncertainties whose impact can be directly assessed in the  
144 form of physical quantities of interest, there exist some uncertainties that cannot  
145 be modelled explicitly by uncertainty propagation. For example, the uncertainty  
146 in the draft will lead to an uncertainty in the natural frequency of the OWC  
147 and hence, in turn uncertainty in the captured power. These kind of uncertain-  
148 ties cannot be directly related to the final power output through uncertainty  
149 propagation rules, and therefore can only be quantified separately. The draft  
150 was set by visual alignment of the water surface and the draft line; hence the  
151 effect of the meniscus may lead to a draft different from the target value. The  
152 uncertainty in the draft is estimated to be about 1-2 mm . Similar uncertain-  
153 ties includes the uncertainty in the orifice size measurement, roundness of the  
154 OWC tube and the non-horizontality of the water column surface. Although no  
155 transversal oscillations were observed during the tests, it should be noted that  
156 the non-horizontality of the water column surface may bias the volume flow rate  
157 determination since the cross-section area in equation 4 is assumed to be flat  
158 and horizontal.

159       In addition to the uncertainties, there are also some known and unavoidable  
160 scaling discrepancies between the two models to available materials and manu-  
161 facturing accuracy. For instance, the diameter of the tube and the thickness of  
162 the orifice plate are not scaled precisely, as shown by Table 1.



### 163 **3. CFD simulation**

164 To simulate the air-water two phase interaction problem, a Finite Volume  
165 Method (FVM) based software STAR-CCM+ is selected to simulate the two  
166 different scale OWC devices. This software has been widely used by several re-  
167 searchers to simulate the OWC problem, for example Lopez[11] and Elhanafi[13].

#### 168 *3.1. Governing equations and numerical solver settings*

169 Star-CMM+ uses a predictor-corrector method to link the continuity and  
170 momentum equations. The shear stress transport(SST)  $k - \omega$  model [19] is se-  
171 lected in current work to model the turbulence. The Volume of Fluid (VOF)[20]  
172 method along with high-resolution interface-capturing (HRIC) scheme [21] are  
173 employed to resolve the free surface. Simulations are carried out by using a  
174 segregated flow model and isothermal ideal gas is selected to account for the air  
175 compressibility. The isothermal law is selected in current study because of the  
176 fact that the air compressibility and temperature variation at such small scales  
177 are negligible. Assuming isothermal avoids solving an ordinary energy transport  
178 equation, and hence reduces computational time.

#### 179 *3.2. Numerical wave tank construction*

180 Wave generation is realised by specifying the time varying wave particle  
181 velocity and wave elevation (hence the phase volume fraction) at the inlet of  
182 the CFD domain (Figure 3). Fifth order Stokes wave theory is adopted to  
183 calculate the required velocity and wave elevation profile.

184 Wave damping at the end of the Numerical Wave Tank (NWT) is achieved  
185 by introducing a resistance to the vertical motion in the form of a momentum  
186 source in a pre-defined zone (for example, zone B in Figure 3) [22]. The length  
187 of the damping zone B is set to be two wavelength ( $\lambda$ ) for good absorption  
188 performance.

189 Wave reflection from the inlet boundary is absorbed using the Euler Overlay  
190 Method (EOM) [23]. This method computes the difference between the analyt-  
191 ical wave information and the actual wave information in the specified region

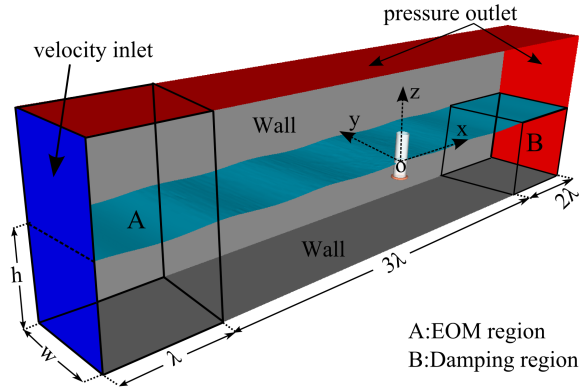


Figure 3: CFD domain and boundary conditions. Here  $h$  is the water depth (2.1 m for the large scale and 0.7m for the small scale.) and  $w$  is half tank width (2.3 m for the large scale and 0.765 m for the small scale simulation.)

192 (zone A in Figure 3, the actual wave is then forced to the analytical wave by  
 193 adding corresponding source or sink into the governing equations. The source  
 194 or sink term takes the following form

$$S(\phi) = -c(\phi - \phi^*) \quad (6)$$

195 Where  $S(\phi)$  is the source or sink corresponding to variable  $\phi$  (time varying wave  
 196 particle velocity distribution along the water depth direction and instantaneous  
 197 wave elevation .). In order to make a smooth transition between the computed  
 198 and the analytical results, a distance dependent weighting function  $c$  is intro-  
 199 duced into the source and sink term. The weighting function has the following  
 200 form

$$c = c_0 \cos^2(\pi x/2) \quad (7)$$

201 Where  $c_0$  is the maximum value of the forcing coefficient and  $x$  is the relative  
 202 distance within the EOM zone ( $x$  equals to zero at the beginning and 1 at the  
 203 velocity inlet, meaning the forcing takes no effect at the end of the EOM zone  
 204 and gives the maximum impact at the velocity inlet.). The choice of the value  
 205 of  $c_0$  is problem dependent [24], a value of 100 was found to be sufficient and  
 206 efficient for the present work.

207 When wave generation is considered, mesh topology normally has a signifi-  
 208 cant impact on the quality of the simulated wave due to numerical dissipation.  
 209 A denser mesh can normally reduce the numerical dissipation at the cost of  
 210 longer computation time. The mesh distribution around the free surface was  
 211 first investigated by performing simulations of a selected wave with different  
 mesh settings. Those simulations were executed in a pseudo-2D manner which

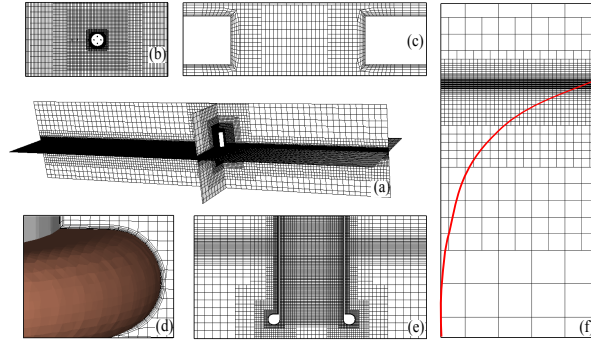


Figure 4: Mesh distribution: (a)overview of mesh distribution, (b) Mesh distribution at the free surface along tank width direction. (c) Mesh distribution around the orifice (sectional view). (d) Mesh distribution around the collar ring. (f) Free surface mesh distribution along water depth direction.

212  
 213 employed only one cell in the tank width direction. The mesh distribution  
 214 around the free surface was designed in such a way that the size of the mesh  
 215 is controlled by the aspect ratio (defined as the ratio between the mesh size  
 216 in the wave height and the wavelength direction.). The mesh topology in the  
 217 water depth direction was decided based on the maximum water particle veloc-  
 218 ity profile. For example, as shown in Figure 4 (f), the mesh gets coarser with  
 219 increasing water depth as the water particle velocity (illustrated by the red solid  
 220 line) reduces. Mesh aspect ratio of 1/2, 1/4, 1/8 and 1/16 were investigated.  
 221 the resulting mesh resolution in the wave propagation direction varies from 60  
 222 cells in one wavelength to 140 cells in one wave length with a uniform step of  
 223 20 cells. The corresponding mesh resolution in the wave height direction is then  
 224 decided by multiplying the mesh size in wavelength direction with the defined

225 aspect ratio. The total number of cells division within the wave height varies  
226 from about 5 to about 35 depending on the aspect ratio and the mesh resolution  
227 in the wavelength direction.

228 As for the numerical dissipation caused by the mesh, inappropriate temporal  
229 discretization will also lead to dramatic numerical dissipation. A second-order  
230 time discretization method is selected to resolve time marching for better accu-  
231 racy. Time step size  $\Delta t$  is decided based on the Courant-Fridrichs-Lewy (CFL)  
232 number via

$$\Delta t = \frac{CFL \cdot \Delta x}{u} \quad (8)$$

233 Where  $\Delta x$  is the size of a single cell at the free surface in x direction (see Figure 3  
234 for coordinate system). The denominator  $u$  is the wave phase velocity in current  
235 study. A CFL number of 0.5 is normally enough to meet the requirement of  
236 a second order temporal discretization scheme; in the present study a value of  
237 0.25 is selected to give an extra safe margin.

238 A regular wave with wave height equal to 0.06m at  $Kh = 4.9$  is tested for  
239 those proposed mesh settings with correspondingly calculated time step size. It  
240 is found that a mesh aspect ratio of 1/8 provides the most economic result for  
241 the current study. Figure 5 demonstrates the spatial distribution of the wave  
242 elevation along the tank length after 40 seconds simulated physical time. As can  
243 be seen, the wave crest height increases with denser meshes in one wavelength  
244 indicating the numerical dissipation is relieved with denser mesh. Comparison of  
245 the wave height measured at 2 wavelength away from the wave generating inlet  
246 and the theoretical wave height suggests the maximum and minimum discrep-  
247 ancy is about 3.7% and 1.7%, respectively. With an improvement of only 2%,  
248 the simulation with 140 cells in the wavelength direction took about 17 hours  
249 on a desktop PC with 32G RAM and 4 core Intel I7-2600 processor (3.4 GHz)  
250 while the case with 60 cells only took about 2 hours with the same computer.  
251 The 80 cells simulation appears to be the most economic case with a discrepancy  
252 of about 2.9% and 3 hours running time. Therefore, this mesh setting for wave  
253 capturing is selected for the OWC simulations.

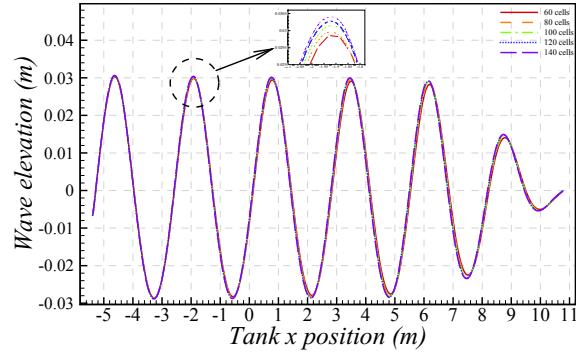


Figure 5: Mesh effect on the free surface elevation spatial distribution along the tank after 40 seconds simulated physical time. The legend states the number of cells in one wavelength.

254 Simulations using different time step sizes are executed to check the reliability of the proposed time step size determination method. The time step size  
 255 calculated, based on Equation 2.3, yields about 0.004 seconds. As suggested by  
 256 Figure6, simulation results converge when the time step size is smaller than 0.005  
 257 indicating the validity of the proposed time step size determination method.

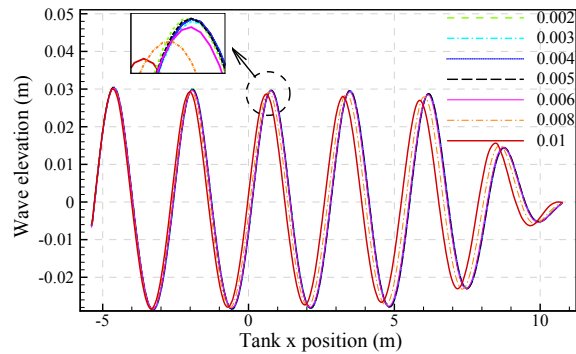


Figure 6: Time step size effect on the wave elevation after 40 seconds physical time simulation.

258

### 259 3.3. OWC simulation

260 The OWC device is then fixed in the middle of the NWT as shown in Fig-  
 261 ure 3. Boundary conditions are illustrated in Figure 3. It should be noted here  
 262 that simulations took advantage of the symmetry of the problem about the tank

263 centreline; hence, only half of the tank was simulated. The total length of the  
264 NWT is set to be  $6 \lambda$  which may or may not be sufficient for accurate simula-  
265 tion of the device's performance. This simplification is made due to the large  
266 demand of computational resource for the OWC simulations. The dense mesh  
267 inside the OWC device (Figure 4 (e)), the extra refined mesh around the orifice  
268 plate (Figure 4 (c)), the mesh refinement outside the OWC device at the free  
269 surface and the employment of modelling of the boundary layer at the device  
270 surface leads to a mesh with typically 2 million volume cells. The transients  
271 of the oscillation of the water column (especially at high frequencies) require  
272 long physical time to reach steady state. Along with the required small time  
273 step size, this places further demands on the computational resource required.  
274 A typical simulation (50 seconds simulated physical time) around the resonant  
275 point takes about 64 hours using 48 cores high performance computer (inter  
276 Xeon X5650 2.66 GHz CPU).

277 In the small scale simulation, the small scale OWC device was scaled directly  
278 from the large scale simulation according to proper scaling law, rather than  
279 reproduced according to the small scale device used in the small scale physical  
280 experiment. Hence, the CFD simulation does not have geometry scaling errors.  
281 In addition to waves, the mesh distribution and time step size settings for the  
282 small scale simulation are directly scaled according to Froude similarity rule as  
283 well.

## 284 **4. Results**

285 Study of the OWC device without orifice plate is first carried out to in-  
286 vestigate the scale effect on the hydrodynamics without PTO damping. These  
287 results are presented in section 4.1. Results of OWC with PTO damping are  
288 presented in Section 4.2.

### 289 *4.1. Tests with no PTO damping*

290 The Response Amplitude Operator, defined as the ratio between the mea-  
291 sured OWC motion amplitude and the incident wave amplitude are presented

292 in Figure 6 for both tank test and CFD simulations. The RAO obtained from  
 293 the CFD simulation has been corrected assuming an incident wave amplitude  
 294 reduction of 3% due to numerical dissipation based on the NWT study.

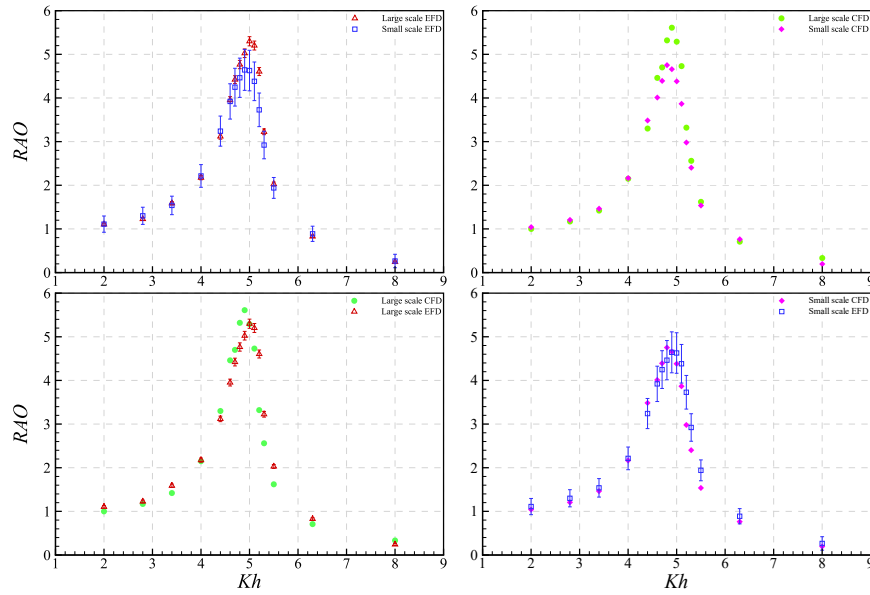


Figure 7: Comparison of RAO obtained from tank test(EFD) and CFD simulation for both large and small scales.

295 As suggested by Figure 7, the uncertainty in the small scale experiment is  
 296 much higher than that of the large scale test. This large uncertainty comes  
 297 from the uncertainty in the wave probes which were used to measure the OWC  
 298 oscillation and incident wave. The absolute Type B (systematic) uncertainty  
 299 introduced by those wave probes used in the small scale test is in fact similar to  
 300 that of the large scale test. However, due to the scaled incident wave amplitude  
 301 (and hence also the OWC response) the uncertainty in the RAO calculation  
 302 increased dramatically through uncertainty propagation. In spite of the large  
 303 uncertainty, it is clear that the response of the small scale test is smaller than  
 304 that of the large scale test around the peak response frequencies even with such a  
 305 large uncertainty. The comparison of the large and small scale CFD simulations

306 clearly confirmed this observation. The smaller RAO obtained from the small  
 307 scale model can be explained by the dissimilarity of Reynolds number, that in  
 308 the small scale OWC is much smaller meaning higher relative viscous losses.

309 The comparison of the CFD simulation and tank test result suggests the  
 310 CFD slightly over predicts the peak RAO obtained from tank test for both  
 311 scales. A frequency shift in the peak response frequency is also observed for  
 312 both scales. The uncertainty in the draft in the tank test mentioned previously  
 313 may contribute to the different in natural frequency. This difference in the  
 314 natural frequency will in turn contribute to the difference in the peak response.  
 315 Nevertheless, the difference in the peak response period is small. Table 2 lists  
 316 the peak response period for CFD simulation and tank test.

Table 2: Peak response period obtained by CFD and tank test

	Large scale	Small scale
CFD	1.313 (s)	0.758 (s)
EFD	1.300 (s)	0.751 (s)

317 The maximum and minimum possible peak RAO based on the 95% uncer-  
 318 tainty values for the large and the small scale tests are listed in Table 3. The  
 319 corresponding possible maximum and minimum difference in the peak RAO be-  
 320 tween the large scale and the small scale tests are 1.23 and 0.08, respectively.  
 321 The possible relative difference thus varies from 1.5% to 22.7% (the minimum  
 322 relative difference is defined as the ratio of the minimum difference to the large  
 323 scale minimum possible RAO value, and the maximum relative difference is de-  
 324 fined as the ratio of the maximum difference to the large scale maximum possible  
 325 RAO value.). Calculations based on the measured peak RAO suggest that the  
 326 small scale test under-predicts the large scale by about 12.4%. Results from the  
 327 CFD simulation suggest an absolute difference in the peak RAO of 0.86 yielding  
 328 a relative difference about 15.3%.



Table 3: Tank test peak RAO and possible minimum and maximum value for the large and small scale OWC.

	Mearsured peak RAO	Uncertainty at peak RAO	Maximum possible RAO	Minimum possible RAO
Large scale	5.300	0.105	5.405	5.195
Small scale	4.644	0.469	5.113	4.175

329 *4.2. Tests with PTO damping*

330 Results are presented here for the case with 4 orifices open to the atmo-  
 331 sphere. This configuration gave the maximum power output compared with  
 332 other conditions (e.g. 8 orifices open to the atmosphere.).

333 The RAO, the pressure amplitude and the capture factor for OWC de-  
 334 vices with modelled PTO damping are presented in Figure 8, Figure 9 and Fig-  
 335 ure 10, respectively.

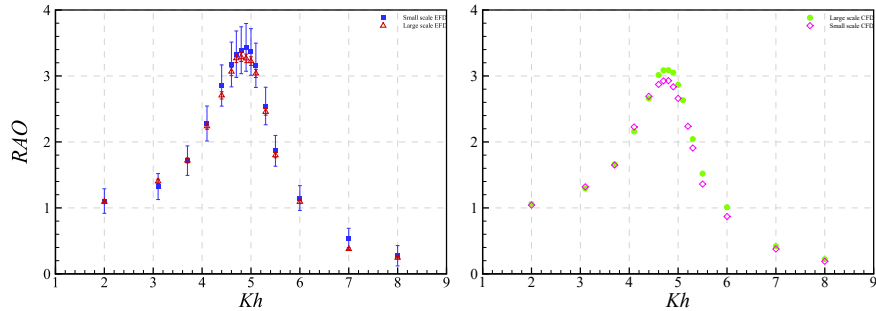


Figure 8: Comparison of RAO obtained from tank test(EFD) and CFD simulation for both large and small scales. PTO damping is set to 4 orifices open to the air.

336 The RAO of the small scale tank test suggests a larger response around the  
 337 peak response frequencies compared with the large scale tank test, however, the  
 338 large uncertainty covers the large scale results which makes the comparison less  
 339 reliable. On the other hand, the CFD simulation suggests that the response of  
 340 the small scale OWC is smaller than that of the large scale OWC.

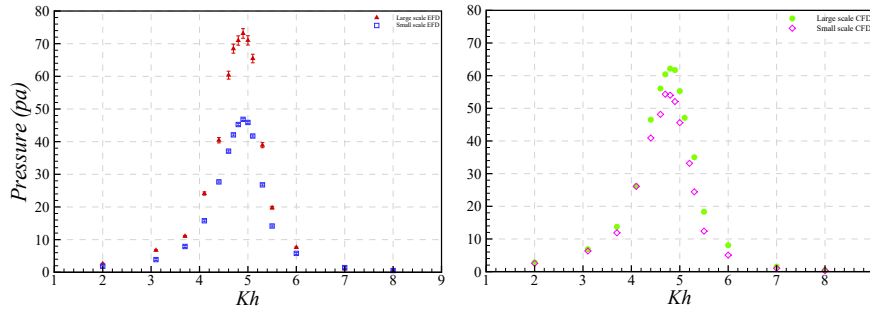


Figure 9: Comparison of pressure amplitude obtained from tank test(EFD) and CFD simulation for both large and small scales. PTO damping is set to 4 orifices open to the air.

341 The pressure amplitude obtained by the tank test and the CFD simulation  
 342 both suggest the small scale OWC has a smaller pressure amplitude around the  
 343 peak response frequency. According to the Froude scaling rule, the pressure  
 344 amplitude should scale with the geometric scale factor. It can be deduced from  
 345 Figure 9 that the small scale pressure amplitude is smaller than the pressure  
 346 amplitude of the large scale after extrapolation to the large scale. Comparing  
 347 the results between the tank test and the CFD simulation, it can be seen that  
 348 the peak pressure amplitude of the large scale simulation is smaller than that  
 349 of the tank test while the small scale simulation has a higher peak pressure  
 350 amplitude. The smaller pressure amplitude of the large scale simulation can be  
 351 explained by the NWT dissipation, which results in, the wave amplitude arriving  
 352 at the OWC device with smaller than the specified wave amplitude. The larger  
 353 pressure amplitude of the small scale CFD simulation will be explained in section  
 354 4.3.

355 Both the tank test and the CFD simulation suggest that the small scale  
 356 capture factor is smaller than that of the large scale. Table 4 lists the measured  
 357 maximum and minimum possible capture factor for the large and the small  
 358 scale tank test. Based on the maximum and minimum possible value, the small  
 359 scale test results under predict the large scale capture factor by about 24.5%  
 360 to 37.6%. The calculation based on the measured value suggests that the small

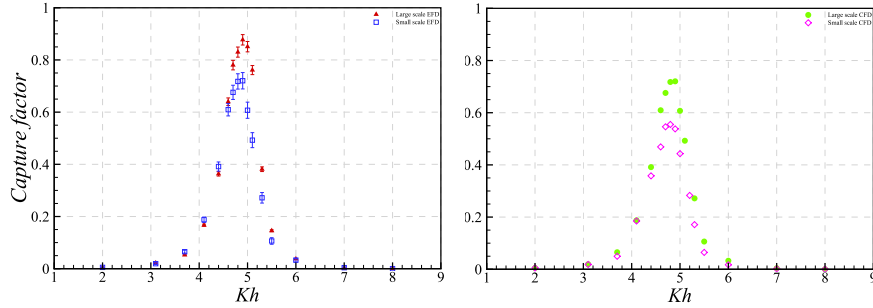


Figure 10: Comparison of capture factor obtained from tank test(EFD) and CFD simulation for both large and small scales. PTO damping is set to 4 orifices open to the air.

361 scale under predicts the large scale result by about 31%. On the other hand, the  
 362 CFD simulation results indicate that the small scale simulation under predicts  
 363 the large scale simulation by about 22.9%.

Table 4: Tank test peak capture factor and possible minimum and maximum value for the large and small scale OWC.

	Mearsured peak Cf	Uncertainty at peak Cf	Maximum possible Cf	Minimum possible Cf
Large scale	0.877	0.039	0.916	0.839
Small scale	0.602	0.031	0.633	0.571

#### 364 4.3. PTO scaling

365 The modelled PTO system: the orifice plate, as mentioned previously, has  
 366 a quadratic pressure-flow rate characteristic as shown by Figure 11 and can be  
 367 described by

$$p = \Lambda q^2 \quad (9)$$

368 the damping coefficient  $\Lambda$  is a real number describing the relationship be-  
 369 tween the pressure and the volume flow rate.

370 As indicated by Figure 11, apart from the small scale tank test, all the other  
 371 pressure and volume flow rate amplitude relationship are very close to each

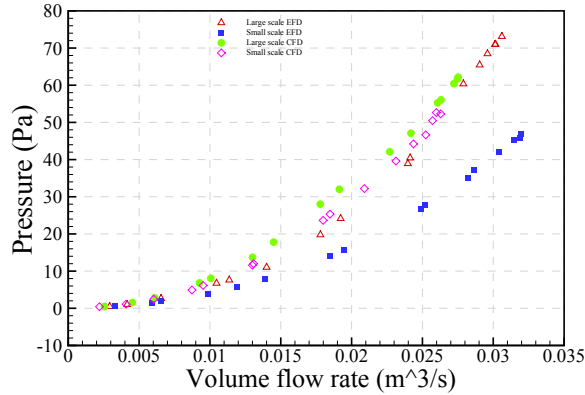


Figure 11: Pressure amplitude and volume flow rate amplitude relationship, here the small scale result are extrapolated to the large scale according to the Froude scaling rule. Each data point corresponding to a single frequency simulation/tank test.

372 other. This suggests that the damping applied in the small scale tank test is  
 373 different from the damping used in the CFD simulation.

374 Following the Froude scaling rule,  $\Lambda$  should scale with  $s^{-4}$ , yielding 1/81  
 375 in the present work, here  $s$  is the scale factor defined as the ratio between  
 376 the geometry dimension of the large scale and the geometry dimension of the  
 377 small scale. The scale factor is 3 in the present study. Figure 12 compared  
 378 the damping ratio (defined as  $\Lambda_{small}/\Lambda_{large}$ ) for the tank test and the CFD  
 379 simulation. This explains why the small scale CFD simulation has a larger  
 380 pressure amplitude than the small scale tank test.

381 It is clear that the damping ratio of the tank test is far from the Froude-  
 382 scaled value. This is because the small scale PTO was not scaled correctly due to  
 383 the errors and uncertainties. The uncertainty in the damping ratio for the tank  
 384 test is enormous due to the large uncertainty in the small scale OWC elevation  
 385 measurement. On the other hand, the damping ratio of the CFD simulation is  
 386 close to the theory ratio compared with the tank test result. However, it is still  
 387 smaller than the theoretical value.

388 Figure 13 illustrates the relationship between the damping coefficient  $\Lambda$  and  
 389 the orifice Reynolds number for the large and small scale CFD simulation. The

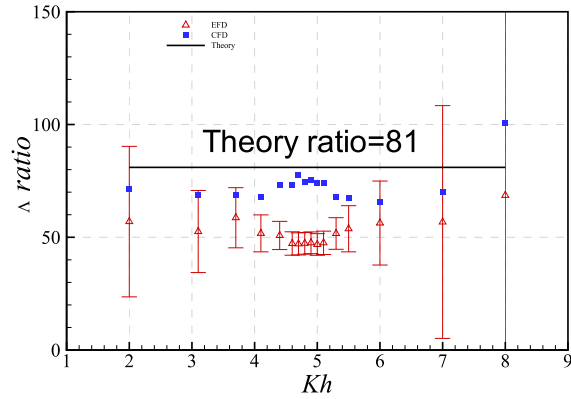


Figure 12: Comparison of the  $\Lambda$  ratio between the tank test and the CFD simulation.

390 orifice Reynolds number is defined as

$$Re = \frac{D \cdot U}{\nu_{air}} \quad (10)$$

391 where  $D$  is the characteristic length of the orifice plate and is defined as 2 times  
 392 the orifice diameter (This characteristic length is decided to be the diameter of  
 393 an orifice whose area is equivalent to the total area of the four orifice),  $U$  is the  
 394 mean air velocity through the orifices, calculated by dividing the OWC volume  
 395 flow rate by the total orifice area and  $\nu$  is the dynamic viscosity of air.

396 Figure 13 suggests that for the large scale, the damping coefficient  $\Lambda$  is  
 397 smaller at low Reynolds number and increases with increasing Reynolds number.  
 398 With further increased Reynolds number (up to about  $6.5E5$ ),  $\Lambda$  is found to  
 399 reduce to some extent and tends to stabilize. The small scale simulation has a  
 400 much smaller Reynolds number and it seems like that the extrapolated small  
 401 scale  $\Lambda$  falls into the low Reynolds number region of the large scale  $\Lambda$ , suggesting  
 402 the small scale  $\Lambda$  experienced Reynolds number effect.

#### 403 4.4. CFD simulation of a larger scale OWC

404 Keeping the air chamber of the OWC device the same height as previous two  
 405 scale simulations, a further extrapolation of three times larger than the large  
 406 scale simulation is carried out. Figure 14 shows the comparison of the mean

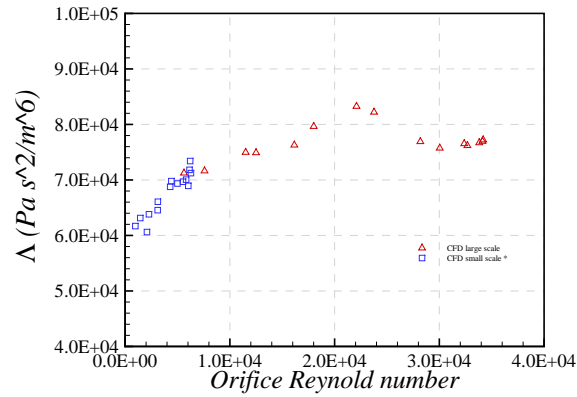


Figure 13: Reynolds effect on the  $\Delta$ . Here the  $\Delta$  of the small scale simulation is extrapolated to the full scale according to the Froude scaling rule while the Reynolds number is kept at small scale.

407 captured power for the three different scales and Table 5 summarises the results  
 408 and relative differences.

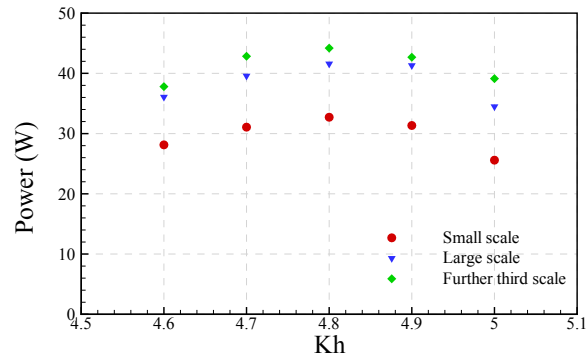


Figure 14: Comparison of the mean captured power for different scales. Power of the small scale and the large scale are extrapolated to the further third scale according to Froude scaling rule.

409 The large scale simulation under-predicts the values obtained for the further  
 410 three-times scale by an average of 7% and the small scale simulation under-  
 411 predicts the further third scale by an average of 28%. Figure 15 plots the peak  
 412 mean power captured against the Reynolds number. Here the Reynolds number

Table 5: Mean power captured for the three different scales. Here the large scale and small scale power are extrapolated to the further 3 times larger scale using Froude scaling rule. The relative difference are calculated based on the further 3 times larger scale.

Kh	4.6	4.7	4.8	4.9	5.0
Further 3 times larger scale (W)	37.80	42.84	44.19	42.68	39.13
large scale (W)	36.06	39.57	41.58	41.31	34.48
small scale (W)	28.13	31.07	32.72	31.35	25.59
Relative difference (large scale)	-5%	-8%	-6%	-3%	-12%
Relative difference (small scale)	-26%	-27%	-26%	-27%	-35%

413 is defined according to [12] as

$$Re = \frac{A\omega D_{out}}{\nu_{water}} \quad (11)$$

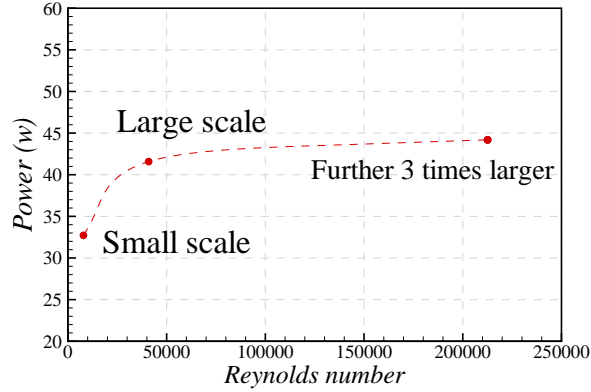


Figure 15: Peak mean captured power against Reynolds number.

414 As indicated by the trend line (dashed line) and Table 5, the small scale  
 415 simulation experienced significant Reynolds number effect. On the other hand,  
 416 the large scale simulation results were less affected by the Reynolds number.  
 417 Judging by the trend line, it seems like Sheng’s [12] recommendation of the  
 418 critical Reynolds number (of the order of  $10^5$ ) is a good estimation.

## 419 5. Discussion and conclusion

420 This paper presents tank test and CFD simulation of two different scale  
421 OWC type WEC. A CFD simulation three times larger again is presented for  
422 further investigation of the scale effect.

423 The comparison between the tank test and CFD simulation of the cases  
424 without the modelled PTO suggest that, in spite of the uncertainty in the draft  
425 of tank test, the CFD simulation can predict the scale effect quite accurately.  
426 The small scale CFD simulation under-predicts the large scale peak RAO by  
427 about 15.3% while the tank tank test suggests a 12.4% difference.

428 When the modelled PTO is considered, very precise scaling of the orifice  
429 at such scales is difficult for tank test. The damping provided by the orifice is  
430 extremely sensitive to the size of the orifice at the small scale tank test. With  
431 only 0.07 mm difference in the orifice diameter (hence 0.28 mm difference in  
432 total since 4 orifices are used.) and 0.167 mm difference in the orifice plate  
433 thickness, the volume flow rate and pressure relationship changed significantly  
434 as suggested by Figure 11. In contrast, the CFD simulation is not restricted  
435 by these practical scaling issues. The small scale CFD simulation provides a  
436 similar volume flow rate and pressure relationship as the large scale simulation.  
437 The CFD results indicate that the small scale simulation underestimated the  
438 large scale peak capture factor by about 22.9% while the tank test suggests the  
439 small scale tank test under-predicts the large scale by between 24.5% to 37.5%  
440 considering the uncertainty. The relative difference between the small scale tank  
441 test and the large scale tank test without considering the uncertainty is about  
442 31%. The difference between the large scale and the small scale tank test is  
443 anticipated to be smaller than 31% if the orifice of the small scale tank test  
444 were able to be perfectly scaled.

445 The discrepancy between the large scale tank test of the OWC with PTO and  
446 the CFD simulation is mainly due to the dissipation introduced by the NWT.  
447 CFD simulations have uncertainties introduced by the mesh, the choice of time  
448 step etc. A careful study of mesh and time step size impact should be carried



449 out to examine the errors and uncertainties for accurate simulation of the OWC  
450 device. However, in this study, same numerical settings and mesh strategies  
451 are adopted for all those different scales, it is assumed to be reasonable to  
452 assume that the errors and uncertainties introduced are unidirectional and have  
453 similar relative effect on the final output. Therefore, it will probably not affect  
454 the comparison between different scales significantly. In order to accurately  
455 simulate the performance of a WEC, it is suggested to calibrate the NWT in  
456 advance and adjust the input wave in such a way so that the wave arriving at  
457 the device is equal to the required value.

458 The performance of the modelled PTO system (the orifice plate) is affected  
459 by the orifice Reynolds number as indicated by Figure 13. The orifice Reynolds  
460 number is dependent on the motion of the OWC which is not known in ad-  
461 vance. Therefore, it is recommended to check the Reynolds number effect on  
462 the damping coefficient  $\Lambda$  afterwards to check whether the performance of the  
463 orifice is strongly affected by the Reynolds number. If so, the orifice Reynolds  
464 number should be reported along with the final result.

465 A further three-times scaled up CFD simulation result indicates that the  
466 large scale simulation used here is not affected by the Reynolds number signifi-  
467 cantly. Judging by the trend line (Figure 15), Sheng's [12] recommendation of  
468 critical Reynolds number seems to be a good choice. It should be noted here  
469 that for the further three-times scaled CFD simulation, the results may be more  
470 affected by the air compressibility which would need further investigation. For  
471 example, perform CFD simulation with a more realistic compressible air model  
472 such as real gas model.

473 CFD simulations of the three different scale OWC in current study required  
474 similar amount of computation resource since mesh and numerical settings were  
475 scaled accordingly. On the other hand, the cost of the tank test increased with  
476 the scale (mainly introduced by the cost of the facilities and model.). However,  
477 the cost of the small scale tank test is in fact lower than that of the small scale  
478 CFD simulations. It is still not cost-effective to investigate an OWC type WEC  
479 at small scale (about 1:100 scale of a full scale device) using CFD simulation at

480 current stage.

## 481 **6. Future work**

482 Although the effect of side wall reflection is not reported in this work, pri-  
483 mary tank test results indicate that a model breadth to the tank width ratio  
484 of 0.2 is not enough to ignore the tank width effect for current OWC device.  
485 Detailed work on the tank width effect will be reported in the future.

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