

Electron Beam-Induced Writing of Nanoscale Iron Wires on a Functional Metal Oxide

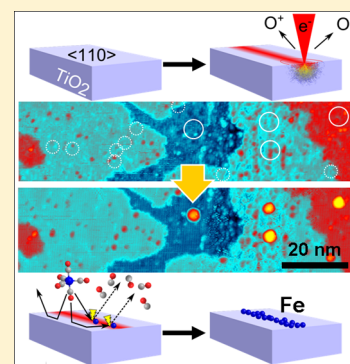
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S Supporting Information

ABSTRACT: Electron beam-induced surface activation (EBISA) has been used to grow wires of iron on rutile $\text{TiO}_2(110)-(1 \times 1)$ in ultrahigh vacuum. The wires have a width down to ~ 20 nm and hence have potential utility as interconnects on this dielectric substrate. Wire formation was achieved using an electron beam from a scanning electron microscope to activate the surface, which was subsequently exposed to $\text{Fe}(\text{CO})_5$. On the basis of scanning tunneling microscopy and Auger electron spectroscopy measurements, the activation mechanism involves electron beam-induced surface reduction and restructuring.



INTRODUCTION

Surfaces of rutile TiO_2 have for many years been used as model systems to explore the physics and chemistry associated with the varied applications of the material.^{1,2} Of particular interest has been the surface science associated with light harvesting processes such as photocatalysis and dye-sensitized photovoltaics.^{1,2} Applications in the field of molecular electronics have also been suggested, with TiO_2 being employed as a dielectric substrate. This has motivated studies of metal wire growth as interconnects using physical vapor deposition.³ Moreover, electron-induced surface modification of $\text{TiO}_2(110)$ ^{4–6} is of interest as a means of patterning the substrate for wire formation or to template an array of functional molecules. There is a related interest in the generation of memristor elements on TiO_2 .⁷

In this work we examine the potential of electron beam-induced surface activation (EBISA)⁸ to deposit metallic structures, including wires, on rutile $\text{TiO}_2(110)$ in ultrahigh vacuum (UHV). The latter environment is essential for the analysis methods employed as well as to avoid contamination from residual gases.⁹ EBISA is a technique that is part of a more general approach termed focused electron beam-induced processing (FEBIP), which also includes the related technique of electron beam-induced deposition (EBID).^{10–13}

A step-by-step explanation of the recently developed EBISA process is given in Figure 1. The substrate (a) is first irradiated with energetic electrons, resulting in local activation of the surface by electron-stimulated desorption of O atoms (b).^{6,14} In

a second step, precursor molecules are decomposed upon contact with the activated area, resulting in a deposition of nonvolatile material (c, d). In addition, the initial deposit grows in size because of autocatalytic decomposition of precursor molecules (e, f).⁸ The autocatalytic growth has been demonstrated for iron pentacarbonyl, $\text{Fe}(\text{CO})_5$, on SiO_x in UHV^{8,15,16} and recently also for $\text{Co}_2(\text{CO})_8$, indicating the general nature of this process for metalcarbonyls.^{17,18} Here we demonstrate EBISA-induced interconnect growth on $\text{TiO}_2(110)$. Moreover, additional insight into the mechanism of the EBISA process has been achieved using local Auger electron spectroscopy (AES) to gain chemical information and scanning tunneling microscopy (STM) to provide atomic scale images.

EXPERIMENTAL SECTION

All experiments were performed in an Omicron UHV Multiscanlab at room temperature. The base pressure of the instrument was $< 2 \times 10^{-10}$ mbar. The instrument includes an electron column for scanning electron microscopy (SEM) with spatial resolution better than 3 nm. This column is also used in combination with a hemispherical electron energy analyzer for local Auger electron spectroscopy (AES) and scanning Auger microscopy (SAM) with a spatial resolution better than 10 nm.

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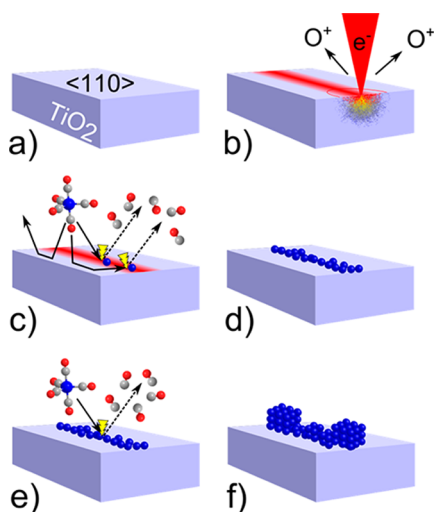


Figure 1. Electron beam-induced surface activation (EBISA) and the secondary growth process. The pristine (110) surface of a TiO_2 (rutile) crystal (a) is exposed to a focused 15 keV electron beam (b), leading to the localized release of oxygen via electron-stimulated desorption processes. The activated surface mediates the dissociation of $\text{Fe}(\text{CO})_5$ molecules (c) into volatile CO molecules and a primary deposit of iron atoms (d). Further supply of $\text{Fe}(\text{CO})_5$ increases the deposit size via an autocatalytic decomposition reaction (e). Under the given reaction conditions, the growth process produces very pure, cubic iron crystallites (f).

For AES, a beam energy of 15 keV and a current of 3 nA were used, while all electron exposures for SEM and lithography were done at an electron beam energy of 15 keV and a probe current of 400 pA. The lithographic processes were controlled via a homemade lithography application based on LabVIEW 8.6 (National Instruments) and a high-speed DAC PCIe-card (M2i.6021-exp, Spectrum GmbH, Germany). For line exposures (single pixel width), a single sweep with a step size of 12 nm was applied. The line dose d_L is calculated as $d_L = I_{PE} \times t_{\text{dwell}} / \Delta x$, where I_{PE} is the beam current, t_{dwell} the dwell time at every position, and Δx the step size along the line. The electron dose was controlled via the respective dwell times. The square patterns were exposed in a single sweep, meander-like pattern at a step size of 6 nm with the area dose d_A being defined as $d_A = I_{PE} \times t_{\text{dwell}} / (\Delta x)^2$. The corresponding area dose d_A for a given line dose d_L can be estimated by dividing d_L by the step size.

The backscattered electron (BSE) exit areas and coefficient were simulated with the program Casino V2.42.¹⁹ The diameter corresponds to an area of which 99% of BSEs are emitted. For the simulations, a primary electron (PE) number of 10^7 was applied. The beam diameter was set to 3 nm.

An STM scanner can be inserted between the pole piece of the SEM column and the sample for in situ STM measurements. Tip positioning is done using SEM imaging. W tips were used in this work, negatively biased relative to the sample. ST micrographs were acquired using Matrix V3.0 (Omicron) and evaluated using the latest versions of WSxM²⁰ (Nanotec Electronica, Madrid, Spain).

Rutile $\text{TiO}_2(110)$ samples (PiKem and Crystec) were prepared by repeated cycles of Ar^+ sputtering and annealing to 1160–1180 K until a sharp (1×1) low-energy electron diffraction (LEED) pattern was obtained with no detectable contaminants in AES. The purity of the $\text{Fe}(\text{CO})_5$ precursor gas (ACROS Organics, 99.5%) was confirmed using a mass

spectrometer. The precursor gas was dosed through a nozzle to about 12 mm from the $\text{TiO}_2(110)$ surface at an estimated local pressure of $\sim 9 \times 10^{-6}$ mbar, based on simulations with the software GIS Simulator (version 1.5).^{21,22}

RESULTS AND DISCUSSION

The first goal of this work was to extend the use of the EBISA technique, which thus far has only been reported for silica surfaces, to a different oxide material, namely a $\text{TiO}_2(110)$ -(1×1) surface. For that purpose, line patterns and square patterns ($2 \times 2 \mu\text{m}^2$) were irradiated with a focused electron beam. $\text{Fe}(\text{CO})_5$ was subsequently introduced to allow growth of iron structures at the irradiated positions. The gas was pumped off after 270 min, corresponding to an exposure of about 1.1×10^5 langmuir (1 langmuir is 10^{-6} Torr s). Figure 2 shows SEM images of the resulting Fe line deposits (panels a–d) and Fe square deposits (panels e–h). The images show a strong increase in the amount of deposited iron with electron dose. AES measurements 24 h after fabrication (see Figure S1 in Supporting Information) indicate the purity of the iron structures is >90 atom %, which is similar to that achieved on SiO_2 substrates.^{8,15,16} Those iron deposits on SiO_2 exhibit a low room-temperature resistivity of $88 \mu\Omega \text{ cm}$ as well as ferromagnetic behavior.²³ These properties can be expected to be similar on TiO_2 , potentially making the corresponding iron deposits suitable interconnects.

At higher SEM magnification we observe that the structures are comprised of crystalline iron nanocubes. The density of the cubes within the line and square deposits varies with the applied electron dose, whereas the size of the individual cubes at a given gas exposure seems to be similar. With increasing gas exposure the size of the clusters increases because of autocatalytic growth. For increased electron doses, the cubes start to merge into continuous patches, as can be seen for the 1.1 C/cm^2 square deposit (Figure 2g) or the line deposits (Figure 2a,b). Higher electron doses mark the start of pronounced proximity effects, i.e., unintended deposits near the irradiated patterns due to electron scattering.^{11,16,24} Pronounced proximity effects are present at the 3.3 C/cm^2 square irradiation doses (Figure 2h) and the higher line doses (Figure 2c,d). For EBISA, all of the observed proximity effects can be attributed solely to BSE and associated secondary electrons (SE_{II}).¹⁶ Proximity effects from forward scattered electrons (FSE) are not observed in EBISA because of the absence of structure growth during the electron beam exposure step.¹⁶

The observation of homogeneously scattered cubes in the nonirradiated surface regions at low electron irradiation doses (Figure 2a,b,e) points to the presence of nucleation points on the $\text{TiO}_2(110)$ surface that were not induced by the electron beam (see also Supporting Information, Figure S2). The defects are oxygen vacancies, which typically are present at an initial concentration of 5% ML (1 monolayer is the number of surface unit cells). The vacancies react with water in the residual vacuum to form the second major form of defects, bridging hydroxyls (one water molecule reacts with one vacancy to produce two hydroxyls).²⁵

Turning to the EBISA mechanism, we can first look at the effect of the 15 keV electron beam on the reduction state of the surface. Figure 3a shows Auger electron spectra for nonirradiated (red line, “pristine”) and strongly irradiated (blue line, “reduced”) areas of a freshly prepared $\text{TiO}_2(110)$ surface. The most prominent features are the Ti_{LMM} and Ti_{LMV} peaks at

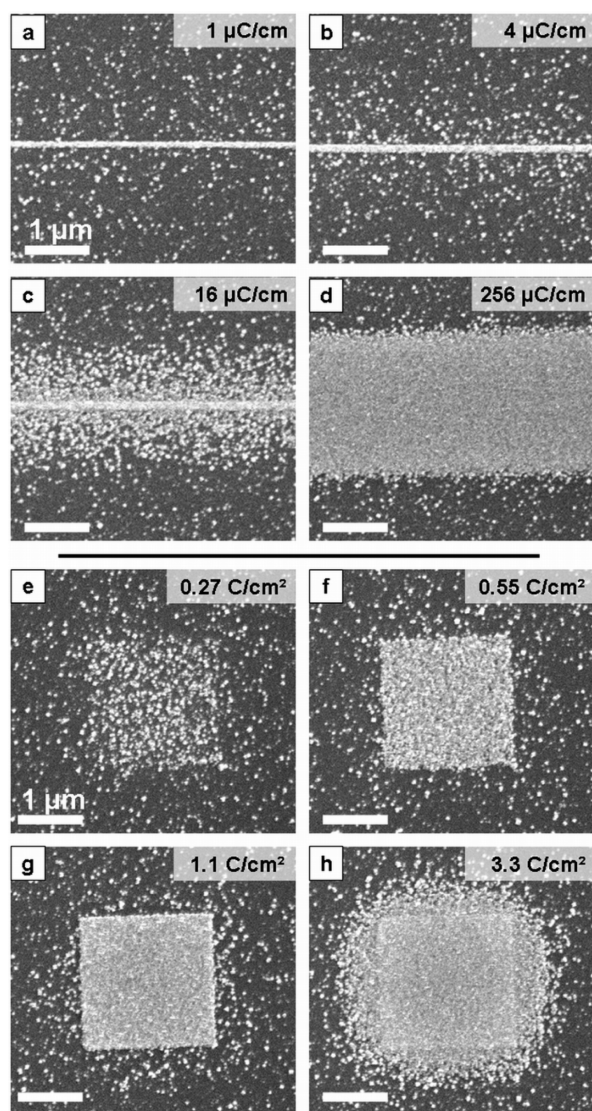


Figure 2. Line (a–d) and $2 \times 2 \mu\text{m}^2$ square deposits (e–h) fabricated by EBISA on $\text{TiO}_2(110)$ using $\text{Fe}(\text{CO})_5$ as a precursor (approximately 1.1×10^5 langmuir). Primary electron (PE) doses are indicated on the top right of each panel. The line width (fwhm, SEM intensity) increases with PE dose range from approximately 150 nm (a) to 2.6 μm (d). The central line is a result of activation by PEs, while the fringe is caused by BSE proximity effects. For the square fields, deposit density increases from a loose grain assembly at 0.27 C/cm^2 via a close-packed assembly at 0.55 C/cm^2 to a fused grain deposit of near ideal geometry at 1.1 C/cm^2 . At even higher doses, proximity effects also cause a fringe growth (3.3 C/cm^2). Note that scattered iron clusters are found all over the surface without electron exposure, a fact which is attributed to active defects created during sample preparation.

383 and 413–419 eV, respectively, and the O_{KLL} peaks at 491 and 511 eV. It is immediately apparent that the oxygen signal for the strongly irradiated surface (blue) is considerably smaller than that for the pristine surface (red); the Ti signals show the opposite trend, albeit not as pronounced. This behavior evidences the electron beam-induced desorption of oxygen from the surface and is in line with that expected from electron-stimulated desorption (ESD) results from TiO_2 at lower electron beam energies.^{6,14} In addition, the change of the peak shape of the Ti_{LMV} signal allows for insight into the chemical state of the Ti atoms of the surface. As discussed by Nishigaki²⁶

and Göpel,²⁷ the LMV Auger peak consists of two components: 413 eV for Ti^{4+} and 419 eV for $\text{Ti}^{<4+}$. In a simplified picture, for Ti^{4+} , in the complete absence of a Ti 3d valence electron, the valence electron in the LMV Auger process stems from an interatomic transition of an O 2p electron with a higher binding energy; on the contrary, for $\text{Ti}^{<4+}$, at least a partial electron charge remains in the Ti 3d level (at lower binding energy than the O 2p level) and thus participates in an intra-atomic Auger transition at higher kinetic energy. Therefore, the observed intensity shift from the low- to the high-kinetic energy contribution of the Ti_{LMV} peak is a direct indication of a reduction of the Ti^{4+} by the electron beam. The evolution of the LMV peak shape upon increasing electron irradiation is illustrated in Figure 3b from top (pristine, red) to bottom (reduced, blue).

In Figure 3c, the normalized ($I_{\text{pristine}} = 1$) peak areas of the Ti_{LMM} (black) and O_{KLL} (orange) signals and the O:Ti peak ratio (green) are plotted versus the primary electron dose (note the logarithmic scale). The apparent O:Ti peak ratio is normalized so that the value for the pristine surface is set to 2, i.e., $\text{TiO}_{2.0}$. At doses lower than 0.1 C/cm^2 , the peak ratio does not change, i.e., the applied dose is not sufficient to induce a reduction of the surface detectable in AES. After crossing the 0.1 C/cm^2 threshold, the oxygen intensity and the O:Ti ratio drop rapidly until a fairly stable lower limit at ~ 0.9 is approached at doses larger than $10^3 \text{ C}/\text{cm}^2$. The corresponding relative oxygen intensity is 0.6 of the initial value of the pristine surface. The obtained O:Ti ratio of 0.9 indicates an average surface composition of $\text{TiO}_{0.9}$. Note that the chemical composition of the surface is likely to be nonuniform after the reduction process, so the detected Auger signal is a convolution of different Ti oxidation states and chemical surrounding. In addition, decreased damping due to the loss of oxygen atoms from the topmost layers has to be taken into consideration when assessing the Ti_{LMM} intensity. Nevertheless, the O:Ti ratio serves as an indicator for the degree of electron beam-induced reduction of the surface.

The Auger spectra indicate that above 0.1 C/cm^2 the stoichiometry of the surface changes gradually from TiO_2 to $\text{TiO}_{0.9}$ at $10^3 \text{ C}/\text{cm}^2$. As for the accompanying morphology change, Figure 4 shows the effect of the electron beam as viewed by STM. Prior to electron beam exposure, the STM image shows [001] direction bright rows separated by 0.65 nm that correspond to Ti atoms.^{1,28} A line exposure corresponding to an electron area dose of approximately 60–100 C/cm^2 (electron energy, 15 keV; beam current, 400 pA; step size, 12 nm; dwell time, 150 ms; line dose, 50 $\mu\text{C}/\text{cm}$) results in an O:Ti ratio of 1.1–1.2. The exposed region can be clearly identified in the image in Figure 4 as an area with a disordered appearance and a strongly increased number of protrusions (bright spots) arranged within a vertical line with a width of approximately $20 \pm 5 \text{ nm}$ (indicated by dotted yellow lines); this width is ~ 7 times larger than the nominal beam diameter of 3 nm (Gauss profile, 20/80 criterion), indicating the well-known broadening due to SE and BSE proximity effects. The apparent surface roughness (rms) of the exposed area (0.12 nm) in STM is double that of the nonirradiated surface. The left-hand side of the image in Figure 4 resembles the as-prepared surface, with the characteristic [001] direction Ti rows visible (see also Supporting Information, Figure S3). However, there are additional defects in the form of isolated protrusions and in-row depressions (white arrows) that can be attributed to proximity effects, as the density is higher than on the as-

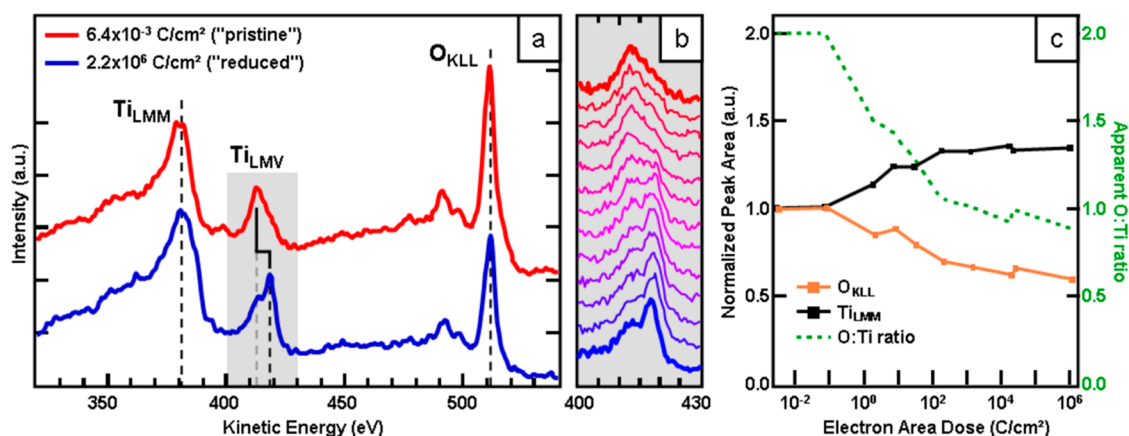


Figure 3. (a) AE spectra of pristine (red) and electron irradiated, reduced (blue) TiO_2 samples in nonderivative mode. The loss of oxygen KLL intensity is a clear indication of oxygen loss, i.e., ESD, while the intensity shift of the titanium LMV peak from lower- to higher-kinetic energy is an indication of a shift from fully oxidized Ti^{4+} toward more reduced states like Ti^{3+} or even lower. The gradual shift of the peak shape with increasing electron dose is shown from top to bottom in (b). (c) Quantitative evaluation of the normalized Ti_{LMM} (black) and O_{KLL} (orange) peak areas and apparent O:Ti ratio (green) versus applied electron dose. The apparent O:Ti ratio (green) is set to 2:1 (TiO_2) for very low electron doses, i.e., a pristine substrate. Electron stimulated desorption causes a loss of oxygen atoms, i.e., signal intensity, which simultaneously leads to an increase in the relative intensity of the Ti_{LMM} signal.

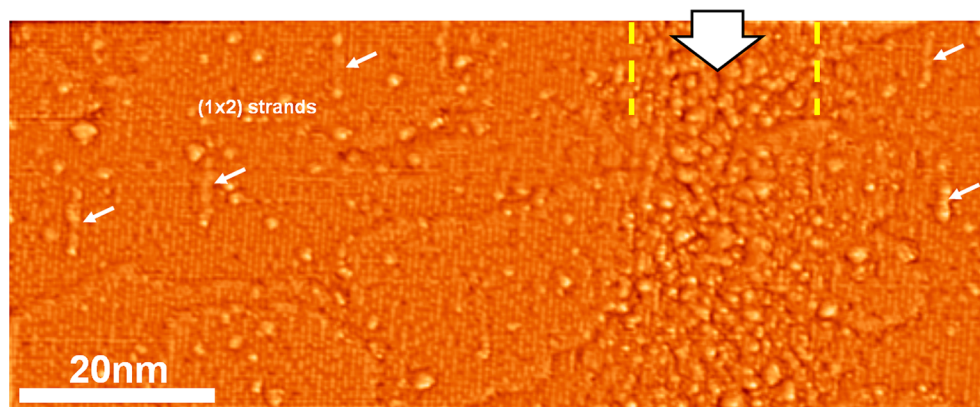


Figure 4. High resolution ST micrograph ($I = 280 \text{ pA}$, $U = 1.2 \text{ V}$, moderately high-pass filtered) of an electron beamline irradiation ($50 \mu\text{C/cm}$), marked by the large arrow and yellow lines. The surface on the left shows only minor damage; the typical row pattern of the $\text{TiO}_2(110)$ surface is preserved. Defects include (1×2) reconstructed sections (small arrows), dark defects in the row structure, and isolated protrusions. On the right, a disordered section running top to bottom was identified as the result of a line irradiation with a primary electron dose of $50 \mu\text{C/cm}$. The (1×1) structure is barely retained in the highly damaged area; no characteristic defects can be identified in the irradiated area. The loss of long- and short-range order corresponds to an electron-induced surface amorphization.

prepared surface. Apart from the increase in mean roughness, the electron beam-induced disorder does not cause a significant increase in average apparent height. This indicates the absence of unintentionally formed deposits from residual gases.

The center of the SEM image in Figure 5a clearly shows a significant electron beam-induced darkening. In previous studies we showed that this darkening of irradiated areas is an intrinsic feature of electron beamed oxide surfaces and is not indicative of contamination.⁸ These results indicate that the modifications to the $\text{TiO}_2(110)$ morphology and subsequent reaction with $\text{Fe}(\text{CO})_5$ are a result of electron beam-induced oxygen loss.

Panels b and c in Figure 5 show the STM images of the irradiated $50 \mu\text{C/cm}$ line pattern before and after $\text{Fe}(\text{CO})_5$ dosing (total exposure: ~ 10 langmuir), respectively. After the gas exposure, the formation of iron deposits in the form of clusters is observed, primarily in the regions irradiated with electrons (indicated by yellow lines). Larger clusters appear as slightly asymmetric protrusions in the image, a consequence of

the applied scanning speed. These clusters have a typical diameter of $3.2 \pm 0.3 \text{ nm}$, including tip convolution, and an apparent height of $1.2 \pm 0.1 \text{ nm}$ (see Supporting Information, Figure S2). This corresponds to clusters containing 650–1000 Fe atoms.

In addition to the clusters observed within the yellow lines in Figure 5, some deposits are also found in nonirradiated areas to both sides of the irradiated line. These clusters originate from growth due to the intrinsic surface defects in $\text{TiO}_2(110)$ (see above), as well as from growth due to BSE proximity effect-induced defects. The BSE exit range of TiO_2 , here defined as the distance from the point of impact that includes 99% of all emitted BSEs (r_{99}), was extracted from MC simulations and amounts to about 1050 nm. All of the scanned area in Figure 5c is well within this exit range. The BSEs ($E > 50 \text{ eV}$ by definition) are able to induce ESD of oxygen atoms within the Knotek–Feibelman framework ($E_{\text{min,ESD}} = \sim 34 \text{ eV}$),¹⁴ thus creating isolated defects.

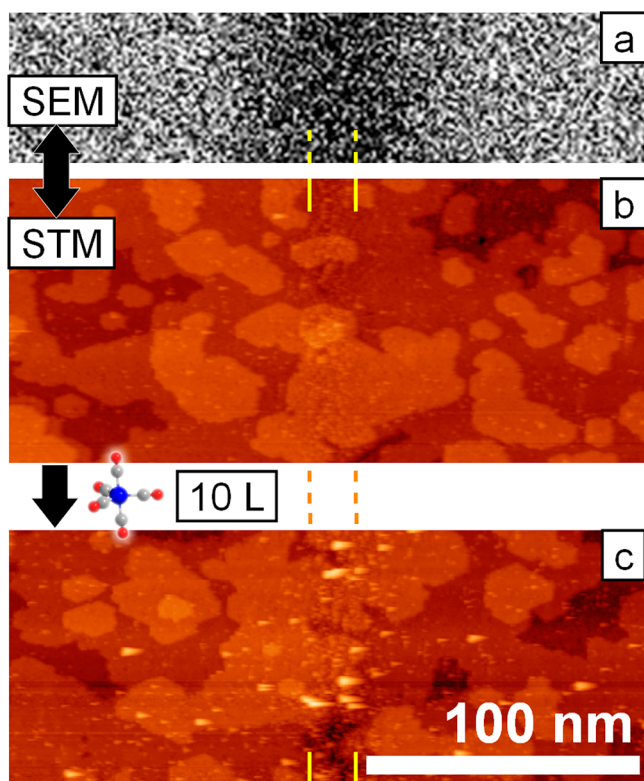


Figure 5. SEM scan (a) and STM topography (b) of the $\text{TiO}_2(110)$ surface after line irradiation with $50 \mu\text{C}/\text{cm}$ and after initial Fe growth from $\text{Fe}(\text{CO})_5$ (c). The line position is indicated by yellow lines. The strongly contrast-enhanced scanning electron micrograph (a) shows a dark feature at the position of the line irradiation, conventionally attributed to carbon deposition from residual gases. STM reveals a high degree of disorder at the line position, while the surrounding surface remains fairly unaffected by the irradiation. Upon exposure to about 10 langmuir of $\text{Fe}(\text{CO})_5$, clusters form predominantly at the irradiated surface area (c). Cluster formation next to the line exposure is attributed to defects induced by backscattered electrons, i.e., common proximity effects.

SUMMARY

Iron nanowires and other nanostructures have been grown on $\text{TiO}_2(110)-(1 \times 1)$ by electron beam-induced activation and subsequent dissociative adsorption of $\text{Fe}(\text{CO})_5$. The mechanism of the EBISA process appears to involve reduction and restructuring of the substrate due to electron stimulated desorption of oxygen. This work has demonstrated the potential of EBISA to write nanoscale interconnects on a functional dielectric substrate and opens up the possibility of electronic circuit fabrication at the nanoscale.

ASSOCIATED CONTENT

Supporting Information

Local Auger electron spectrum of an iron structure prepared by EBISA on TiO_2 (Figure S1), higher resolution STM data of a line irradiation before and after Fe cluster growth (Figure S2), and STM images of the freshly prepared rutile $\text{TiO}_2(110)$ surface (Figure S3). This material is available free of charge via the Internet at <http://pubs.acs.org>.

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The authors declare no competing financial interest.

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