

Ionization cross-sections for positron collisions with N₂^{*}

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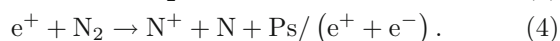
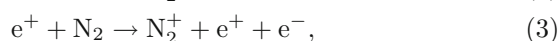
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Abstract. The total and direct ionization cross-sections for positron impact on molecular nitrogen have been measured in the energy range from 5 eV to 850 eV. The results are compared with other experimental and theoretical determinations.

1 Introduction

In recent years, a focus in the investigations of positron (e^+) interactions with atoms and molecules has been ionization, a reaction that at low and intermediate energies is often even more significant for positrons than for electrons (e.g. [1,2]). Recent measurements include cross-sections that are differential in energy and/or angle (e.g. [3,4]), and those relating to ionization accompanied by the excitation of positronium (Ps) [5] and/or of target ion [6]. In the present work, the total and direct ionization cross-sections for positron impact on molecular nitrogen at energies in the range 5–850 eV are presented and compared with other experimental and theoretical results.

The main processes by which a nitrogen molecule can be ionized by a low energy positron are annihilation, positronium formation, direct ionization and dissociative ionization (which itself can proceed via direct ionization or positronium formation), as expressed respectively by:



Whilst the annihilation channel is open at all energies, the ground-state Ps-formation threshold is given by $E_{Ps} = E_i - 6.8 = 8.78$ eV where E_i is the first ionization energy

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of N₂ and 6.8 eV the binding energy of ground state Ps. The corresponding total ionization cross-section is usually expressed as:

$$Q_i^t = Q_{Ps} + Q_i^+ + Q_i^{diss} + \sum HO \quad (5)$$

where Q_{Ps} is the Ps formation cross-section, Q_i^+ the direct single ionization cross-section, Q_i^{diss} the dissociative ionization cross-section, and $\sum HO$ the sum over all higher order processes (such as double and transfer ionization) as well as annihilation, since the latter is considered to give generally a negligible contribution to the total ionization cross-section, as discussed in reference [7].

2 Experimental apparatus and method

Details of the present experimental set-up can be found in references [8,9]; those elements specific to this work, along with a brief description, are outlined below and illustrated in Figure 1. A beam of approximately 10^4 slow positrons s^{-1} , generated by moderating the fast β^+ output of a ²²Na source using an annealed W-mesh moderator, is constrained radially by a magnetic field of 100 Gauss. By applying a positive potential V_m to the moderator, the peak energy of the beam can be varied according to $E_+ \approx eV_m + 2$ eV. The transmission of unwanted components, for example unmoderated positrons or secondary electrons, is substantially reduced by transporting the beam through a Wien filter and a cylindrical electrode held at -500 V prior to the interaction cell. This consists of a hemispherical polished aluminium gas cell from which ions are extracted perpendicularly to the beam axis using a lens held at -500 V and are detected using a channel electron multiplier (CEM) with the cone held at -2800 V. In the present set-up, no N⁺ signal was resolved and, in accordance with an ion optics simulation of the system which indicated a poor efficiency for the extraction of energetic dissociation products, considered negligible.

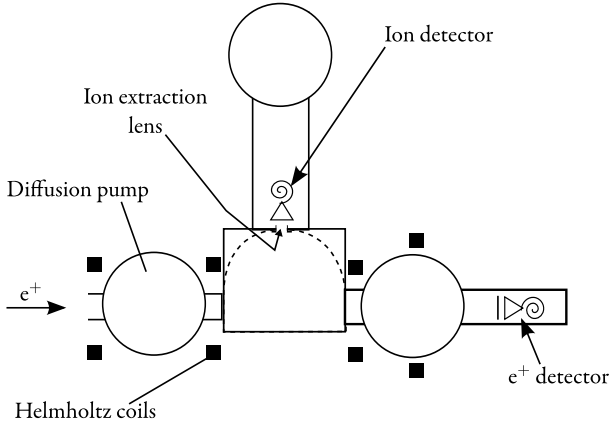


Fig. 1. Interaction region of the beamline showing the positioning of the detectors.

The total yield of ions is defined as:

$$Y_i^t(E) = \frac{N_i - B_i}{N_+ - B_+} \quad (6)$$

where N_i and N_+ are the measured number of ions and positrons respectively, and B_i and B_+ represent the associated backgrounds measured by biasing off the slow portion of the positron beam. The absolute scale of the corresponding cross-section is set according to

$$Q_i^t(E) = \frac{1}{nl_{\text{eff}}} \frac{\epsilon_i^+ Y_i^t}{\epsilon_i^-}, \quad (7)$$

where the ϵ_i^+ and ϵ_i^- are the positron- and ion-detector efficiencies and nl_{eff} is the effective areal gas density, established by the normalization of Q_i^t for Ar (determined under identical conditions to the present measurements) to that of reference [10]. The absolute scale thus set yielded values found to agree with $Q_i^t(e^-)$ at high energies. The latter result is in agreement with the Born approximation according to which the total and direct ionization cross-sections for e^+ and e^- should merge at sufficiently high energies. The direct ionization cross-section (Q_i^+) is measured simultaneously to Q_i^t via the corresponding yield Y_i^+ , determined by counting coincidences between the ion- and the positron-detectors,

$$Q_i^+(E) = \frac{1}{nl_{\text{eff}}} \frac{Y_i^+}{\epsilon_i^+ \epsilon_c} = \frac{Q_i^t}{Y_i^t} \frac{Y_i^+}{\epsilon_i^+ \epsilon_c t_g^4}, \quad (8)$$

where ϵ_c is the efficiency for the coincidence system and t_g the transmission coefficient of a grid. The normalization procedure was checked by measuring the equivalent electron cross-section using the same apparatus [6].

3 Results

Figure 2 shows the present results for the total ionization cross-section for N_2^+ production compared with the electron data of reference [13] for the non-dissociative

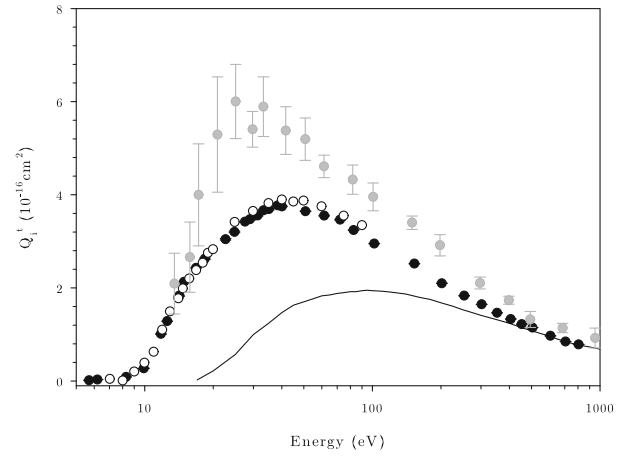


Fig. 2. Measurements of the total non-dissociative ionization cross-section (Q_i^t) for positron impact on N_2 : present measurements (black circles) and [11] (grey circles), compared with those including dissociative ionization [12] (hollow circles). The corresponding non-dissociative ionization cross-section by electron impact [13] (solid line) is also shown.

component of the cross-section only. The positron cross-section, peaking at 40–50 eV with a magnitude of around $3.8 \times 10^{-16} \text{ cm}^2$, exceeds that for electrons, predominantly due to Ps formation, up to approximately 600 eV where they merge. Corresponding determinations of Q_i^t by [11] display a peak at some 20 eV lower energy with a 50–60% larger magnitude. Also shown are the results of reference [12] which comprise, unlike the present measurement, possible contributions from dissociative (as well as non-dissociative) ionization. The close correspondence between the two sets below 20 eV, and the slight excess of the latter in comparison with the present results above this energy, might imply a cross-section of the order of $0.2 \times 10^{-16} \text{ cm}^2$ for N^+ production (the threshold for which is at 24.29 eV), a factor of approximately 3.5 lower than than the peak value for electrons [13]. The main dissociative pathway for N^+ production in single ionization is via the $N_2^+(C^2\Sigma_u^+)$ state, the creation of which requires the removal of an inner-valence electron [14]. The suppression of the latter type of reactions by a positron has been previously observed (e.g. [6]), although it must be noted that a peak value of approximately $0.8 \times 10^{-16} \text{ cm}^2$ has also been reported for positrons by reference [11].

As for Q_i^t , Figure 3 shows the present direct ionization cross-section for N_2^+ production compared with the corresponding electron data of reference [13] which it undercuts both at low and high energies. Whilst the former feature is typical and thought to arise mainly from competition with Ps formation, the latter (amounting to approximately 10%) is not and might originate from systematic uncertainties in the normalization [6]. A fair shape agreement may be discerned among all experimental determinations and differences in absolute scale may arise from either the inclusion of dissociation (in the case of Ref. [12]) or normalization methods in comparison with [11]. Indeed, the discrepancy of $\sim 20\%$ between the present data and

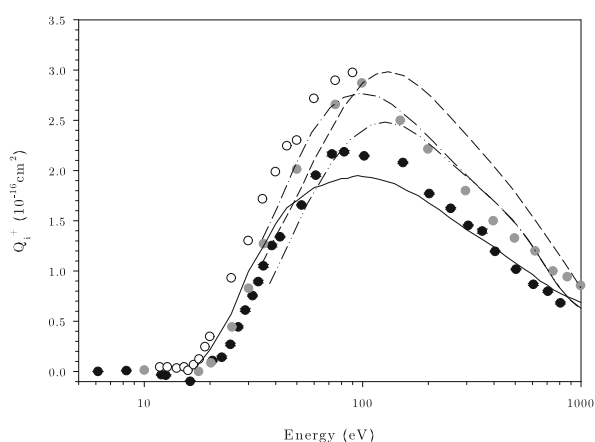


Fig. 3. Measurements of the direct non-dissociative ionization cross-section for positron impact on N_2 : present measurements (black circles) and [11] (grey circles), compared with those including dissociative ionization [12] (hollow circles). Also shown are the corresponding non-dissociative ionization cross-section by electron impact [13] (solid line) and the theoretical results of reference [15] corresponding to various models (CPE: short dash; ES: dash dot; TS: dash double dot), as discussed in the text.

that of [11] would be essentially removed if the latter were normalized to the more recent electron data of reference [13] rather than of reference [16]. All three experimental determinations broadly agree on the position of the peak at an energy of ~ 100 eV. Also shown in the figure are the Distorted Wave Born Approximation calculations of reference [15] for the production of N_2^+ ions in which the wavefunctions of the ejected electron and of the incident and scattered positron are calculated in a spherically averaged potential created by the nuclei and the bound electrons (denoted by ES in the figure). Although the present results are generally lower than theory, we note that these theoretical results were improved relative to the CPE model (where the wavefunctions were approximated by Coulomb or plane waves) and total screening (TS) model where the screened potentials have been used for calculating wavefunctions for both the incoming and outgoing particles. The results of the latter two models are also shown in the figure.

4 Conclusions and outlook

In conclusion, measurements of Q_i^t and Q_i^+ for N_2^+ production by positrons impact on N_2 have been performed from 5 eV to 850 eV and compared with previous determinations for both positrons and electrons. A new gas cell has been designed with improved ion extraction properties

which will allow in future the sensitive investigation of dissociative ionization channels.

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