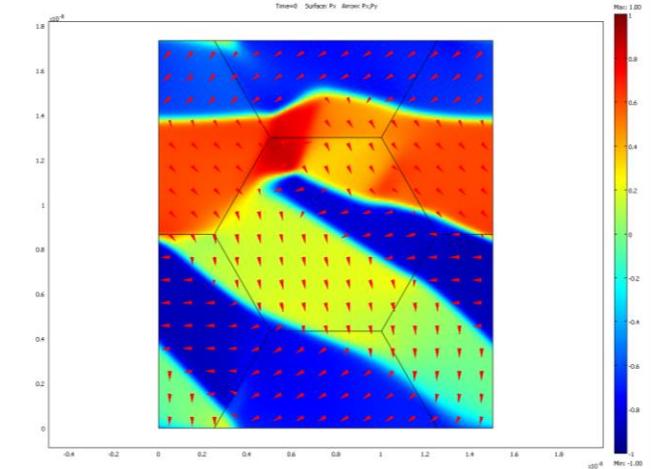
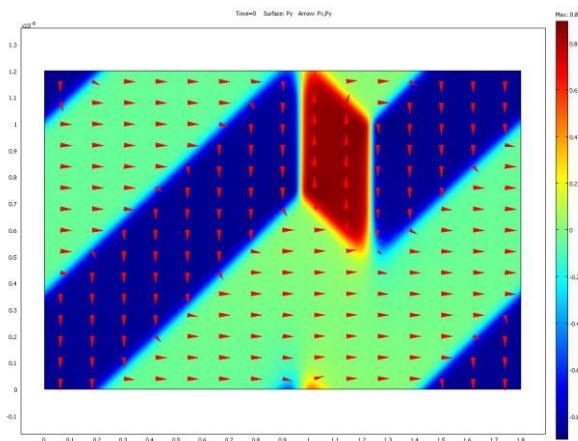
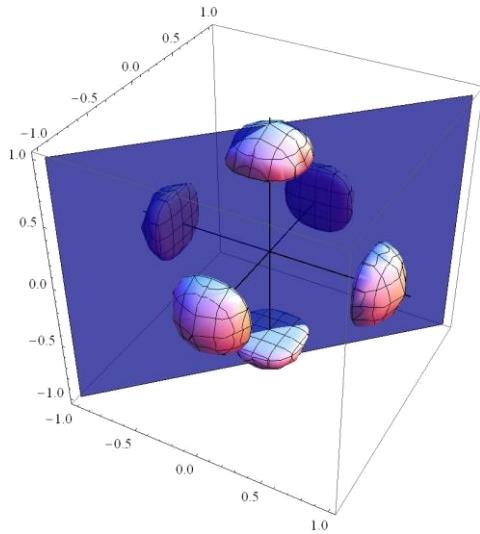


Phase-field modeling of ferroelectric materials in the context of a multiscale simulation chain

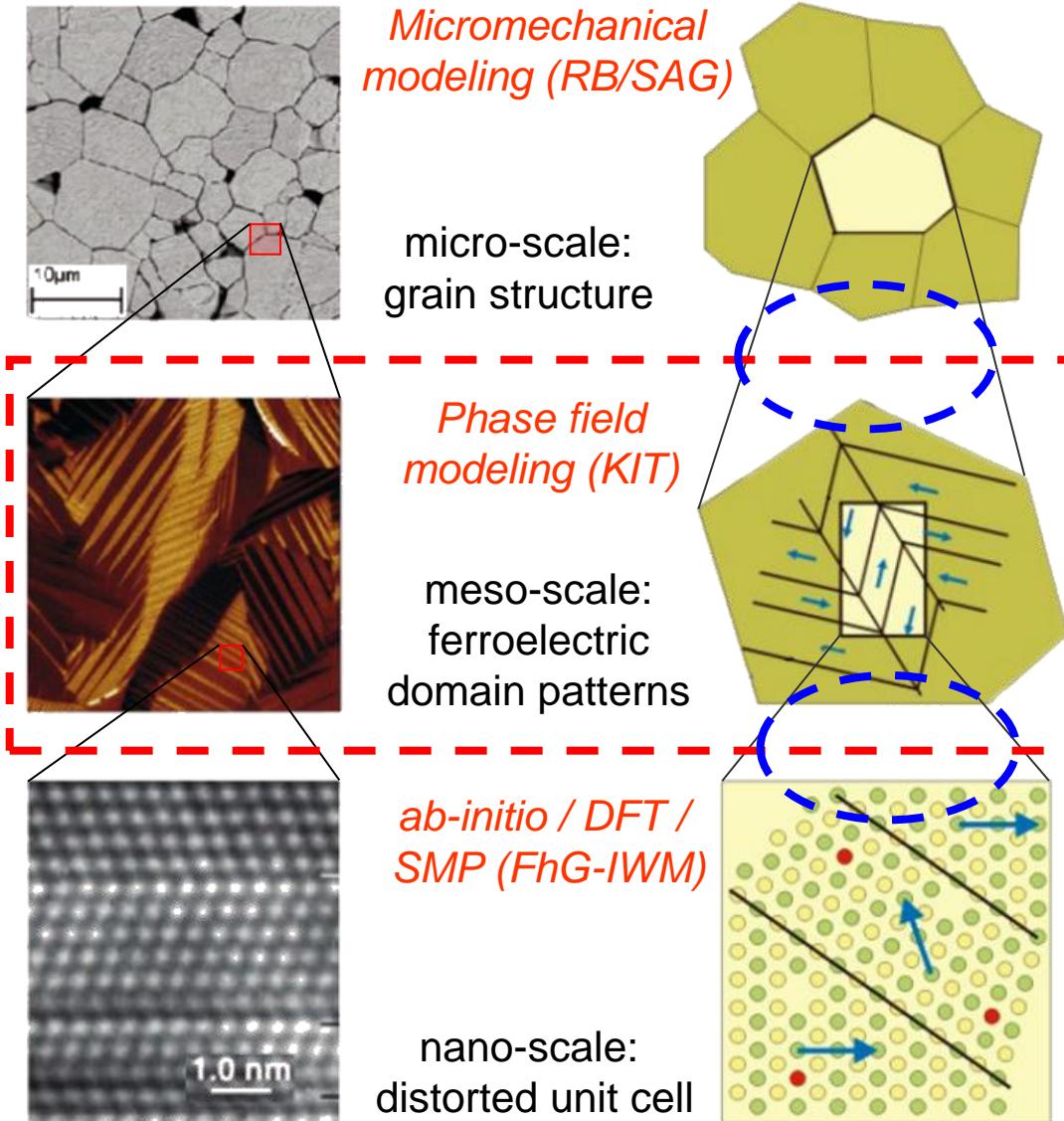
Benjamin Völker, Magalie Huttin and Marc Kamlah

SMASIS10, Sept. 29th, 2010
Philadelphia, PA

Institute for Materials Research (IMF II)



Motivation: virtual material development for ferroelectrics



→ need for multiscale-approach
→ BMBF-Project COMFEM:
Computer based multiscale modeling for virtual development of polycrystalline ferroelectric materials (esp. PZT)

Project partners:

- Fraunhofer IWM Freiburg (FhG-IWM)
- Robert Bosch GmbH (RB)
- Siemens AG (SAG)
- PI Ceramic AG (PIC)
- CeramTec AG (CT)
- TU Hamburg-Harburg (TUHH)

Our aim:

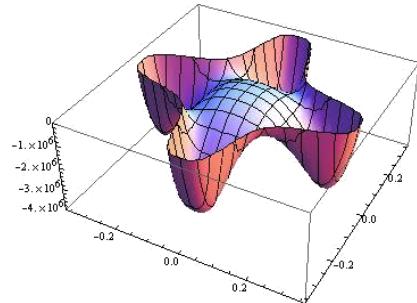
→ development of two interfaces in simulation chain

Thermodynamically motivated phase-field theory

Aim: calculation of ferroelectric domain patterns on meso-scale

- Helmholtz free energy function contains all crystallographic and boundary information

$$\psi(P_i, P_{i,j}, \epsilon_{ij}, D_i)$$



- state variables: partial derivatives with respect to natural variables

$$\sigma_{ji} = \frac{\partial \psi}{\partial \epsilon_{ij}} \quad E_i = \frac{\partial \psi}{\partial D_i}$$

- temporal and spatial evolution of polarization (order parameter):
time-dependent Ginzburg-Landau-equation

- domain switching caused by minimization of free energy

$$\left(\frac{\partial \psi}{\partial P_{i,j}} \right)_{,j} - \frac{\partial \psi}{\partial P_i} = \beta_{ij} \dot{P}_j$$

Formulation of the phase-field model's free energy

6th order free energy

$$\psi = \frac{1}{2} G_{ijkl} P_{i,j} P_{k,l}$$

$$+ \frac{1}{2} \alpha_{ij} P_i P_j + \frac{1}{2} \alpha_{ijkl} P_i P_j P_k P_l + \frac{1}{6} \alpha_{ijklmn} P_i P_j P_k P_l P_m P_n$$

$$+ q_{ijkl} \epsilon_{ij} P_k P_l + \frac{1}{2} c_{ijkl} \epsilon_{ij} \epsilon_{kl}$$

$$+ \frac{1}{2\kappa_0} (D_i - P_i)(D_i - P_i)$$

Main parts of energy function:

- gradient term
- Landau energy
- electromechanical coupling term
- elastic energy term
- electric field energy

Adjustment of 6th order free energy

Interface ab-initio / phase field modeling: Adjustment of parameters

→ ab-initio: piezoelectric coefficients
 (input) dielectric permittivity

elastic stiffness

spontaneous strain

spontaneous polarization

domain wall energy (90°/180°)

domain wall thickness (90°/180°)

d_{ijk}	
K_{ij}	
C_{ijkl}	
ϵ_S	
P^S	
$\Upsilon_{90/180}$	
$\xi_{90/180}$	

Ginzburg-Landau-theory:
 15 parameters (6th order)

a_{ijklmn}
 q_{ijkl}
 c_{ijkl}
 G_{ijklj}

$$\begin{aligned} \psi = & \frac{1}{2} G_{ijkl} P_{i,j} P_{k,l} \\ & + \frac{1}{2} \alpha_{ij} P_i P_j + \frac{1}{2} \alpha_{ijkl} P_i P_j P_k P_l + \frac{1}{6} \alpha_{ijklmn} P_i P_j P_k P_l P_m P_n \\ & + q_{ijkl} \epsilon_{ij} P_k P_l + \frac{1}{2} c_{ijkl} \epsilon_{ij} \epsilon_{kl} \\ & + \frac{1}{2\kappa_0} (D_i - P_i)(D_i - P_i) \end{aligned}$$

- adjustment method has been developed
- applied to PTO and PZT

B. Völker, P. Marton, C. Elsässer, M. Kamlah: "Multiscale Modeling of ferroelectric materials: a transition from the atomic level to phase-field modeling". Continuum Mechanics and Thermodynamics, submitted on Sept. 3rd, 2010

Adjustment of 6th order free energy: Results for PTO and PZT

		PTO		PZT	
unit		first-principles data (input)	phase-field model (adjusted)	first-principles data (input)	phase-field model (adjusted)
DFT	P_0	[C/m ²]	0.88 0.04209* −0.007388*	0.88 0.04209 −0.007388	0.58 0.012039* −0.0017946*
	κ_{33}		$17\kappa_0$	$17\kappa_0$	$18\kappa_0$
	κ_{11}		$54\kappa_0$	$54\kappa_0$	$76\kappa_0$
	C_{11}	[Pa]	342×10^9	342×10^9	361×10^9
	C_{12}	[Pa]	131×10^9	131×10^9	115×10^9
	C_{44}	[Pa]	108×10^9	108×10^9	91×10^9
	d_{33}	[C/m]	2.46×10^{-11}	1.42×10^{-11}	1.57×10^{-11}
	d_{31}	[C/m]	-8.04×10^{-12}	-2.52×10^{-12}	-4.32×10^{-12}
	d_{15}	[C/m]	1.72×10^{-11}	1.72×10^{-11}	1.53×10^{-12}
	$\gamma_{DFT,180}$	[mJ/m ²]	112	173	96
SMP	$\gamma_{DFT,90}$	[mJ/m ²]	24	71	-
	$\xi_{DFT,180}$	[m]	4.5×10^{-10}	4.5×10^{-10}	6.7×10^{-10}
	$\xi_{DFT,90}$	[m]	5.4×10^{-10}	5.4×10^{-10}	-
	$\gamma_{SMP,180}$	[mJ/m ²]	156	156	-
	$\gamma_{SMP,90}$	[mJ/m ²]	64	64	36
	$\xi_{SMP,180}$	[m]	3.9×10^{-10}	3.9×10^{-10}	-
	$\xi_{SMP,90}$	[m]	4.9×10^{-10}	4.9×10^{-10}	6.6×10^{-10}

atomistic input:

DFT: density functional theory

SMP: shell-model potential

(*P. Marton and C. Elsässer,
IWM Freiburg*)

generally good agreement, but:

- not enough degrees of freedom for piezoelectric coefficients
- only cubic elastic behavior taken into account

Formulation of the phase-field model's free energy

6th order free energy

$$\psi = \frac{1}{2} G_{ijkl} P_{i,j} P_{k,l}$$

$$+ \frac{1}{2} \alpha_{ij} P_i P_j + \frac{1}{2} \alpha_{ijkl} P_i P_j P_k P_l + \frac{1}{6} \alpha_{ijklmn} P_i P_j P_k P_l P_m P_n$$

$$+ q_{ijkl} \epsilon_{ij} P_k P_l + \frac{1}{2} c_{ijkl} \epsilon_{ij} \epsilon_{kl}$$

$$+ \frac{1}{2\kappa_0} (D_i - P_i)(D_i - P_i)$$

Main parts of energy function:

- gradient term
- Landau energy
- electromechanical coupling term
- elastic energy term
- electric field energy

additional terms

$$+ f_{ijklmn} \epsilon_{ij} \epsilon_{kl} P_m P_n + g_{ijklmn} \epsilon_{ij} P_k P_l P_m P_n$$

[Su,Landis2007], for BaTiO₃

benefit of additional terms:

more degrees of freedom
for adjustment process:

- ➔ f-term: tetragonal elastic behavior
- ➔ g-term: independent adjustment of d_{ijk}

Y. Su, C. M. Landis: "Continuum thermodynamics of ferroelectric domain evolution: Theory, finite element implementation, and application to domain wall pinning". Journal of the Mechanics and Physics of Solids, 55 (2007), 280–305

Improvement of adjustment process – additional energy terms

Additional elastic energy term

$$\psi_{\text{elast}}(\epsilon_{ij}, P_i) = c_{ijkl}\epsilon_{ij}\epsilon_{kl} + f_{ijklmn}\epsilon_{ij}\epsilon_{kl}P_mP_n$$

From [Su,Landis2007], for BaTiO₃

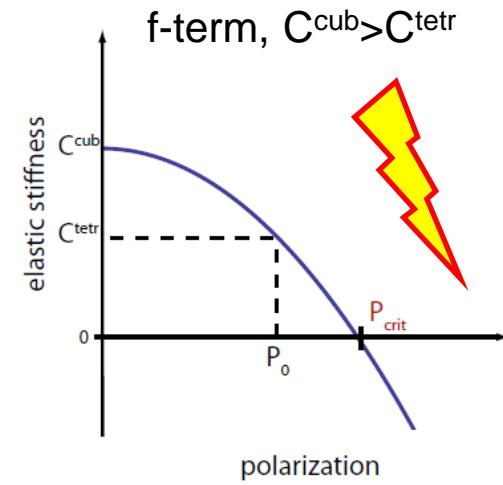
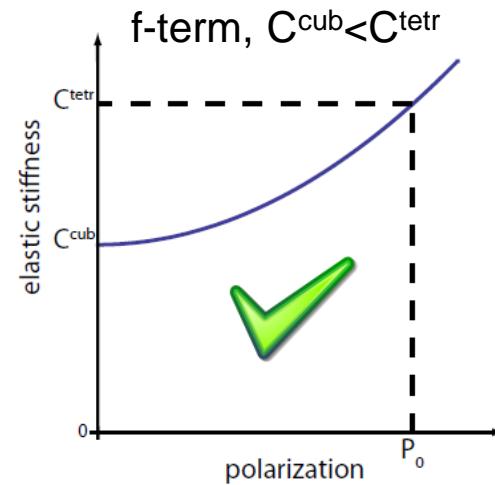
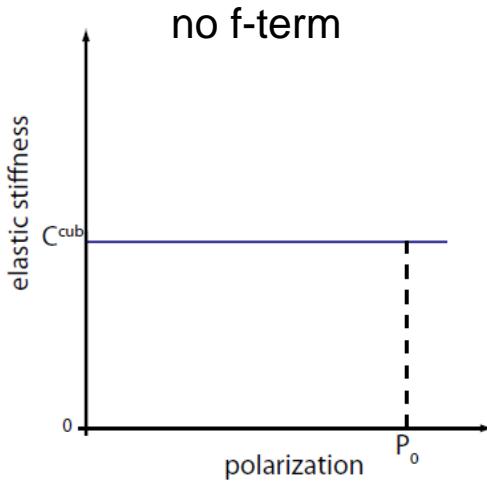
Idea: elastic stiffness depends on polarization $C_{ijkl} = \frac{\partial}{\partial \epsilon_{kl}}\sigma_{ij} = \frac{\partial^2}{\partial \epsilon_{ij} \partial \epsilon_{kl}}\psi(P_i, \epsilon_{ij})$

P=0: cubic elastic properties C^{cub}

P=P₀: tetragonal elastic properties C^{tetr}

} can be adjusted independently

→ works fine for BaTiO₃, but problematic for DFT predictions of PbTiO₃ and PZT



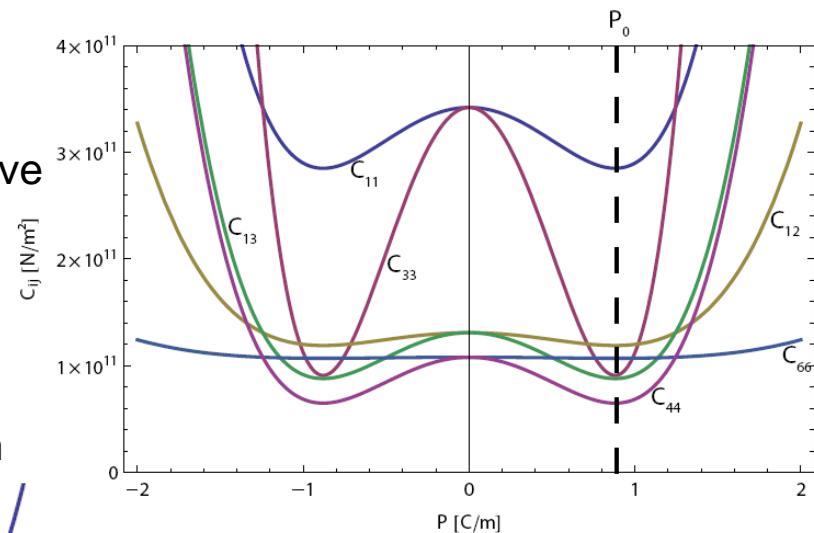
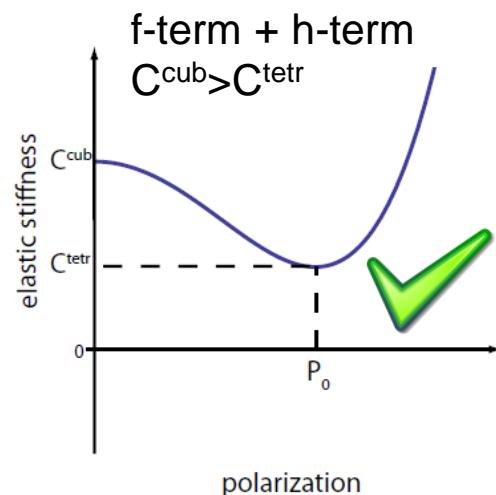
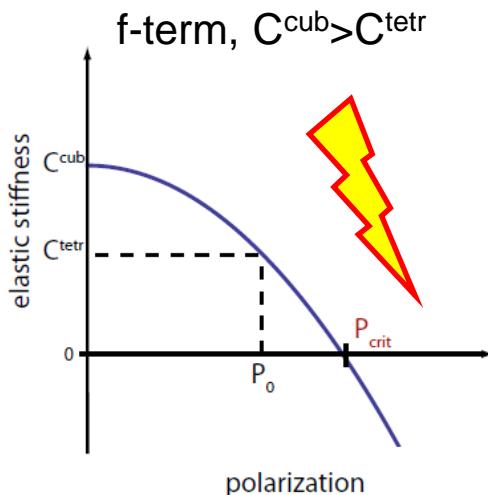
Improvement of adjustment process – additional energy terms

Suggestion: additional elastic energy term

$$\psi_{\text{elast}}(\epsilon_{ij}, P_i) = c_{ijkl}\epsilon_{ij}\epsilon_{kl} + f_{ijklmn}\epsilon_{ij}\epsilon_{kl}P_mP_n + h_{ijklmnr}\epsilon_{ij}\epsilon_{kl}P_mP_nP_rP_s$$

h-term necessary, when $C^{\text{cub}} > C^{\text{tetr}}$

- h-term: ensures elastic stiffness to remain positive
- one possibility: $C(P)$ has minimum at $P=P_0$



Extended free energy

$$\psi = \frac{1}{2} G_{ijkl} P_{i,j} P_{k,l}$$

$$\begin{aligned}
 & + \frac{1}{2} \alpha_{ij} P_i P_j + \frac{1}{2} \alpha_{ijkl} P_i P_j P_k P_l + \frac{1}{6} \alpha_{ijklmn} P_i P_j P_k P_l P_m P_n + \frac{1}{8} \alpha_{ijklmnr} P_i P_j P_k P_l P_m P_n P_r P_s \\
 & + q_{ijkl} \epsilon_{ij} P_k P_l + \frac{1}{2} c_{ijkl} \epsilon_{ij} \epsilon_{kl} + f_{ijklmn} \epsilon_{ij} \epsilon_{kl} P_m P_n + g_{ijklmn} \epsilon_{ij} P_k P_l P_m P_n \quad [\text{Su,Landis2007}] \\
 & + \frac{1}{2\kappa_0} (D_i - P_i) (D_i - P_i) + h_{ijklmnr} \epsilon_{ij} \epsilon_{kl} P_m P_n P_r P_s
 \end{aligned}$$

Additional free energy terms:

Landau energy:

P^8 -term

180° domain wall adjustment

$P_i^4 P_j^4$ -term

90° domain wall adjustment

Elastic energy:

f-term

tetragonal elastic behavior

h-term

necessary if $C^{\text{cub}} > C^{\text{tetr}}$

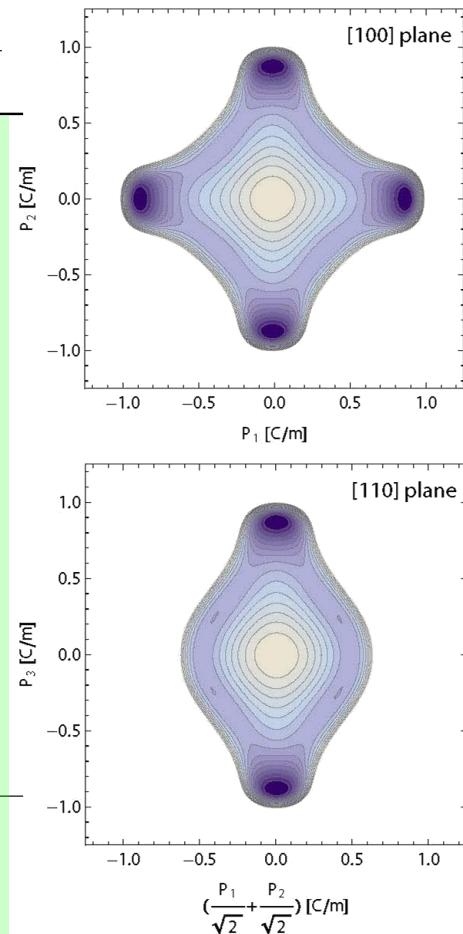
coupling energy:

g-term

piezoelectric coefficients

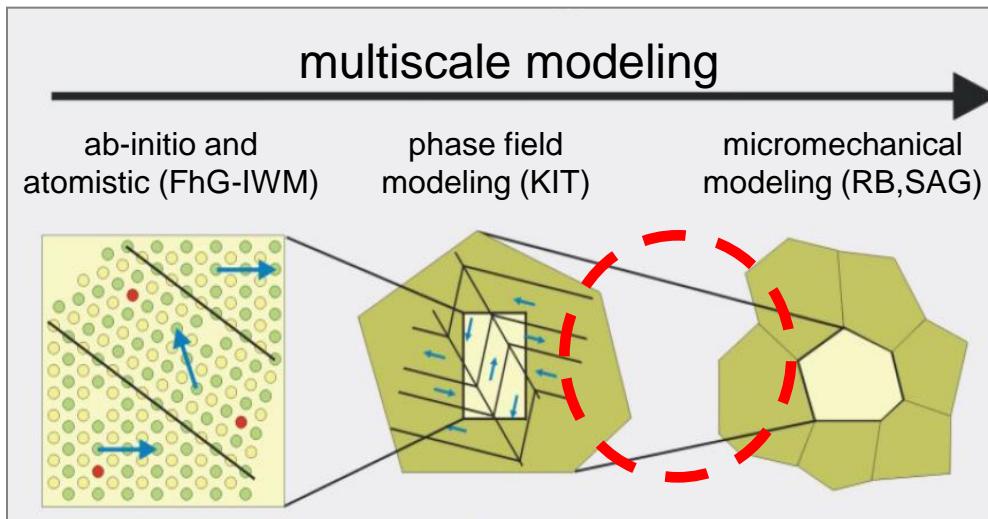
Results: Adjustment of 8th order free energy + additional terms

unit	PTO		PZT		$P_2 [C/m]$
	first-principles data (input)	phase-field model (adjusted)	first-principles data (input)	phase-field model (adjusted)	
P_0 [C/m ²]	0.88	0.88	0.58	0.58	
e_{\parallel}	0.04209	0.04209	0.012039	0.012039	
e_{\perp}	-0.007388	-0.007388	-0.0017946	-0.0017946	
κ_{33}	$17\kappa_0$	$17\kappa_0$	$18\kappa_0$	$18\kappa_0$	
κ_{11}	$54\kappa_0$	$54\kappa_0$	$76\kappa_0$	$76\kappa_0$	
C_{11}^{cub} [Pa]	342×10^9	342×10^9	361×10^9	361×10^9	
C_{12}^{cub} [Pa]	131×10^9	131×10^9	115×10^9	115×10^9	
C_{44}^{cub} [Pa]	108×10^9	108×10^9	91×10^9	91×10^9	
C_{11}^{tetr} [Pa]	285×10^9	285×10^9	327×10^9	327×10^9	
C_{33}^{tetr} [Pa]	91×10^9	91×10^9	178×10^9	178×10^9	
C_{12}^{tetr} [Pa]	119×10^9	119×10^9	110×10^9	110×10^9	
C_{13}^{tetr} [Pa]	88×10^9	88×10^9	107×10^9	107×10^9	
C_{44}^{tetr} [Pa]	65×10^9	65×10^9	73×10^9	73×10^9	
C_{66}^{tetr} [Pa]	108×10^9	108×10^9	92×10^9	92×10^9	
d_{33} [C/m]	2.46×10^{-11}	2.46×10^{-11}	1.57×10^{-11}	1.57×10^{-11}	
d_{31} [C/m]	-8.04×10^{-12}	-8.04×10^{-12}	-4.32×10^{-12}	-4.32×10^{-12}	
d_{15} [C/m]	1.72×10^{-11}	1.72×10^{-11}	1.53×10^{-12}	1.53×10^{-12}	
γ_{180} [mJ/m ²]	112	208	96	96	
γ_{90} [mJ/m ²]	24	24	$(36)^{\dagger}$	36	
ξ_{180} [m]	4.5×10^{-10}	4.5×10^{-10}	6.7×10^{-10}	6.7×10^{-10}	
ξ_{90} [m]	5.4×10^{-10}	5.4×10^{-10}	$(4.9 \times 10^{-10})^{\dagger}$	4.9×10^{-10}	



PZT: complete agreement between atomistic input and adjusted phase-field model
 PTO: only 180° domain wall energy too high, otherwise complete agreement

Second interface: phase-field - micromechanics



Input for micromechanical model:

- domain effective material parameters:
 d_{ijk}^{eff} C_{ijkl}^{eff} κ_{ij}^{eff}
- irreversible switching behavior

FE-Implementation:

[Su,Landis2007]

degrees of freedom per node: $u_i, P_i, \phi \rightarrow$ independent variables $\epsilon_{ij}, P_i, P_{i,j}, E_i$

Weak form: $\int_V (\sigma_{ji}\delta\epsilon_{ij} - D_i\delta E_i + \eta_i\delta P_i + \xi_{ji}\delta P_{i,j}) dV = \int_S (t_i\delta u_i - \omega\delta\phi) dS$

subdomain (volume) terms

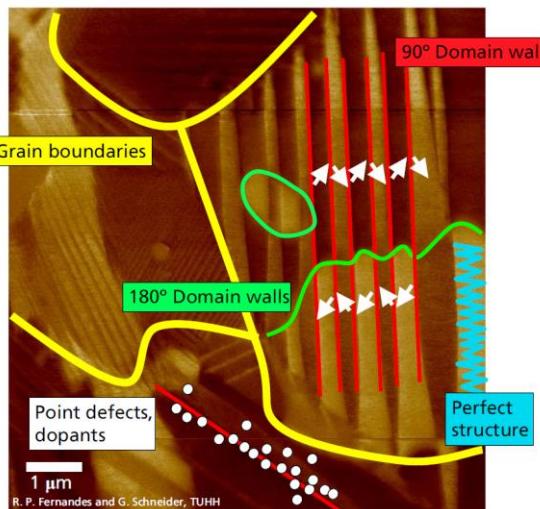
boundary terms

→ direct implementation of weak form in COMSOL Multiphysics

Aim: investigation of typical bulk domain structures

How to obtain typical domain configurations?

Simulation of a whole grain ($\text{Ø} \sim \mu\text{m}$): not possible!



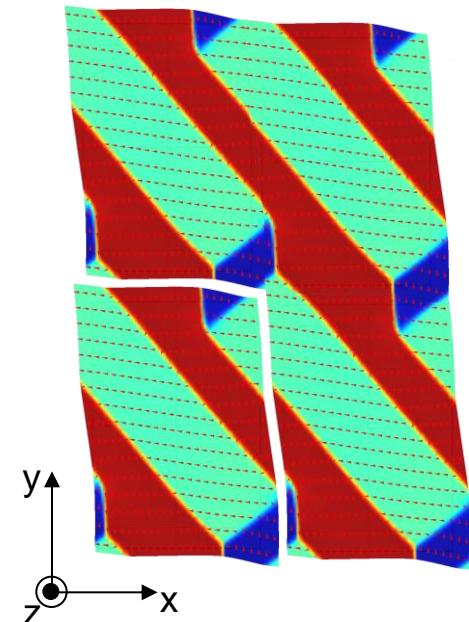
real domain structure: no knowledge about pinning, boundaries, ...

- 1) bulk behavior: periodic boundary conditions required
- 2) stabilize configuration: apply global strain

- investigate “typical” domain structures:
- monodomain
 - ideal 90° domain stack
 - defect-free bulk domain structures
 - influence of charge defects and grain boundaries

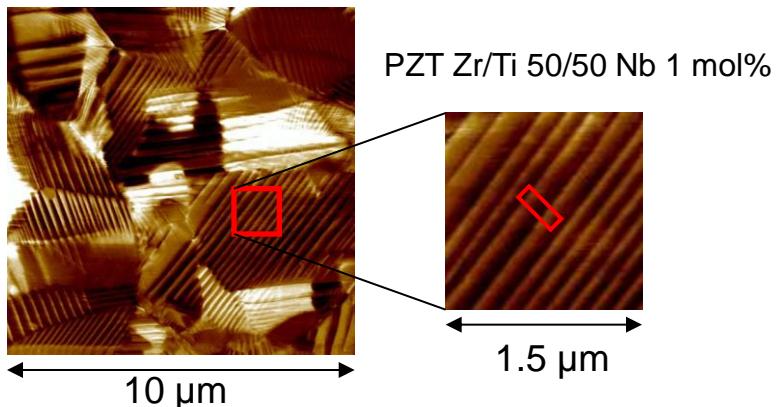
FE-model:

- 2D
- DOF: $P_x, P_y, P_z, u_x, u_y, u_z, \Phi$
- x/y: periodic boundary conditions (P_i, u_i, Φ)
- z-direction: plain strain
- reasonable mesh density: 5-6 nodes / nm

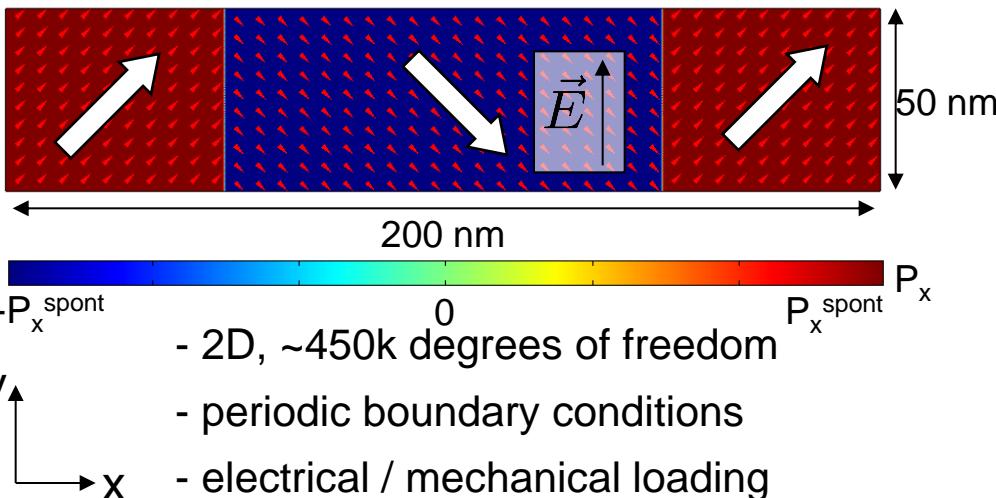


Example 1: Investigation of 90° domain stacks

From PFM-experiments: typical domain width ~100-200 nm [Fernandéz/Schneider,TUHH]



Phase field model: 90° domain stack

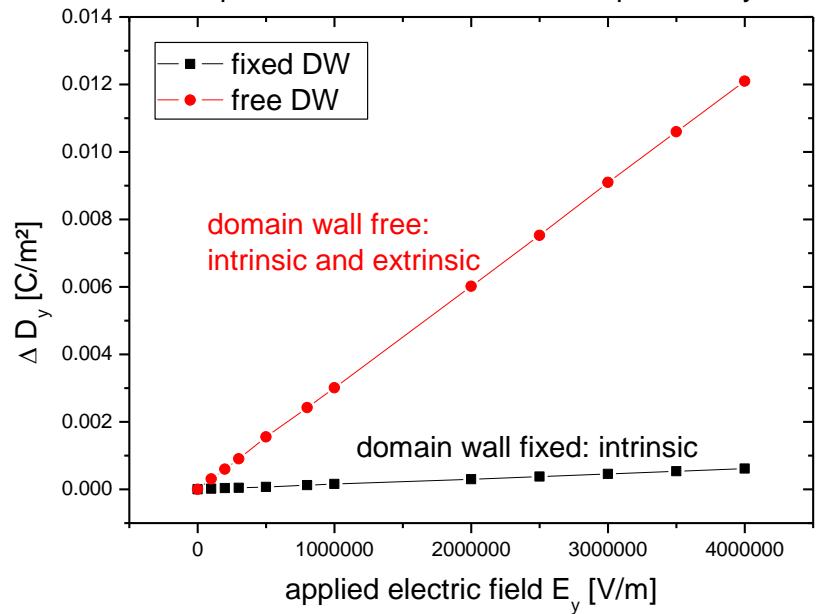


Example: 90°- stack, electrical loading (Y-direction)

Domain wall (DW):

- 1) artificially completely fixed
- 2) free

slope: domain effective dielectric permittivity



domain wall free:
intrinsic and extrinsic

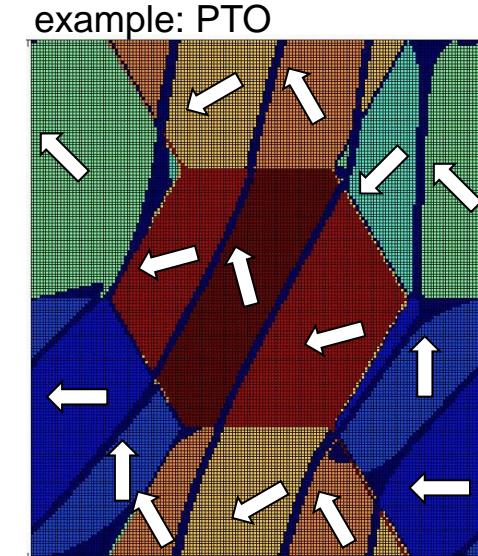
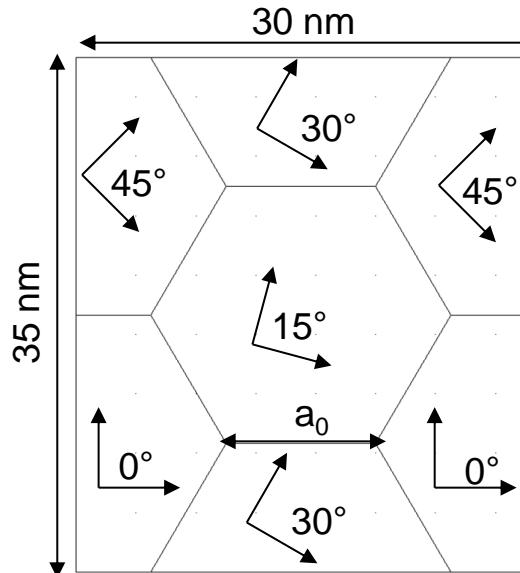
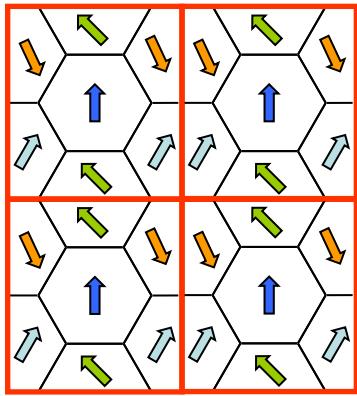
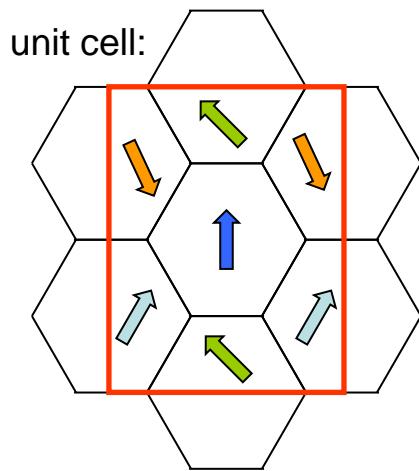
domain wall fixed: intrinsic

→intrinsic/ extrinsic piezoelectric effect
→(reversible) DW motion

small signal behavior: reversible domain wall motion identified as governing process

Example 2: Influence of grain boundaries

Motivation: Influence of polarization orientation mismatch at grain boundaries on domain structure



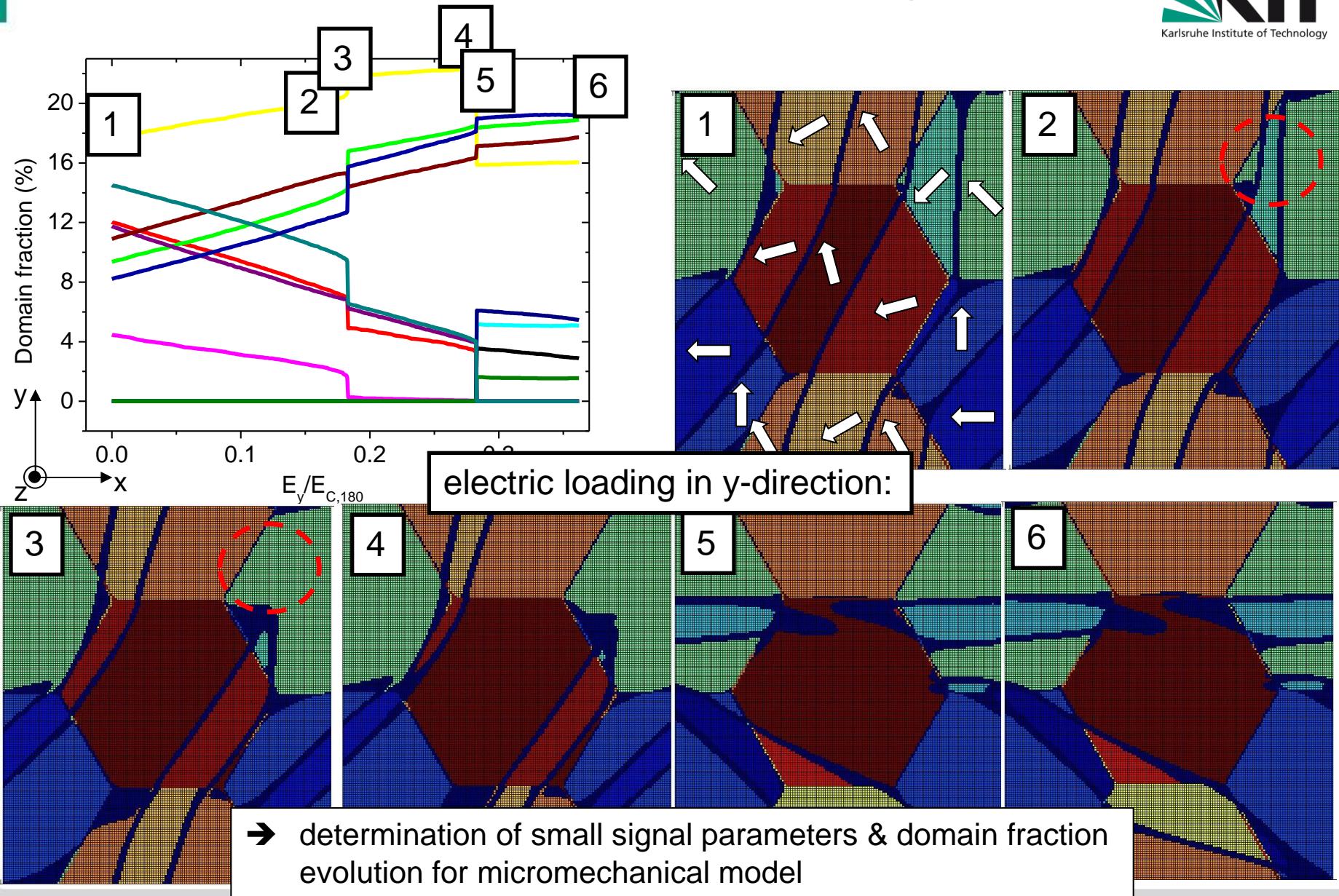
→ simple model:

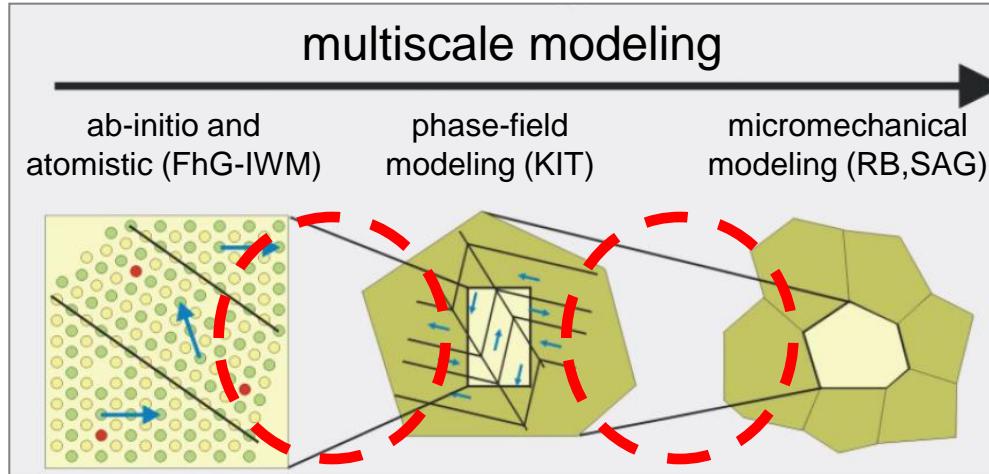
- allows for different polarization directions
- can be continued periodically

example:

- 4 "grains" rotated between 0° and 45° around (001)-axis
- $a_0 = 10\text{nm}$

Reversible DW motion and irreversible switching



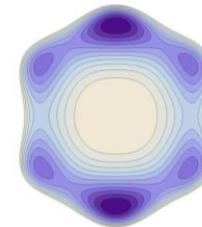
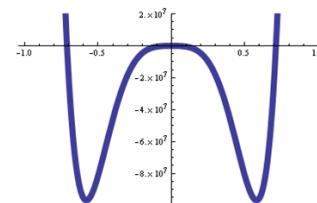
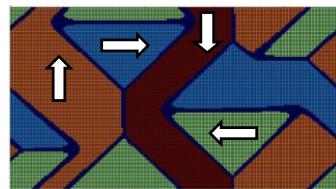


Interface ab-initio / phase-field

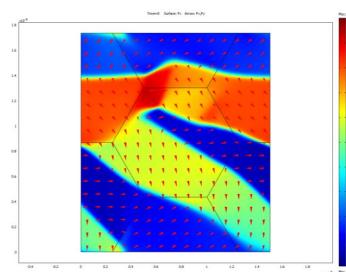
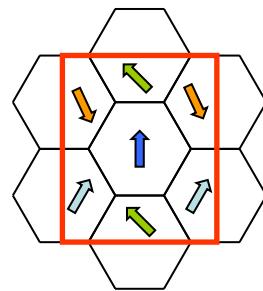
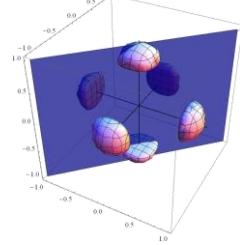
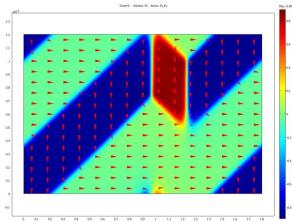
- new approach for adjustment of energy function parameters *solely* based on results of atomistic calculations
- additional energy term introduced enabling tetragonal elastic behavior in PTO and PZT
- successfully applied to PTO and PZT

Interface phase-field / micromechanics

- FE-implementation in COMSOL Multiphysics, including periodic boundary conditions
- intensive investigation of typical bulk domain structures
- computation of small signal parameters, can be transferred to micromechanical model



Thanks for your attention!



Literature:

[Su,Landis2007]

Yu Su, Chad M. Landis: "Continuum thermodynamics of ferroelectric domain evolution: Theory, finite element implementation, and application to domain wall pinning". Journal of the Mechanics and Physics of Solids, 55 (2007), 280–305

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