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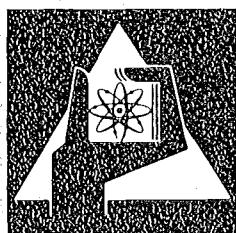
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**FANAC - A Shape-Analysis Program for
Resonance Parameter Extraction from
Neutron Capture Data for Light and
Medium-Weight Nuclei**

F. H. Fröhner



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FANAC - A Shape Analysis Program for Resonance
Parameter Extraction from Neutron Capture Data for
Light and Medium-Weight Nuclei

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Abstract

A least-squares shape analysis program is described which is used at the Karlsruhe Nuclear Research Center for the extraction of resonance parameters from high-resolution capture data. The FORTRAN program was written for light to medium-weight or near-magic target nuclei whose cross sections are characterized on one hand by broad s-wave levels with negligible Doppler broadening but pronounced multi-level interference, on the other hand by narrow p-, d- ... wave resonances with negligible multi-level interference but pronounced Doppler broadening. Accordingly the Reich-Moore multi-level formalism without Doppler broadening is used for s-wave levels, and a single-level description with Doppler broadening for p-, d- ... wave levels. Calculated capture yields are resolution broadened. Multiple-collision events are simulated by Monte Carlo techniques. Up to five different time-of-flight capture data sets can be fitted simultaneously for samples containing up to ten isotopes. Input and output examples are given and a FORTRAN list is appended.

FANAC - Ein Multiniveau-Formanalysen-Programm zur Resonanzparameter-Bestimmung aus Neutroneneinfangdaten für leichte und mittelschwere Kerne

Zusammenfassung

Ein nach der Methode der kleinsten Quadrate arbeitendes Formanalysenprogramm wird beschrieben, welches am Kernforschungszentrum Karlsruhe zur Bestimmung von Resonanzparametern aus Neutroneneinfangdaten hoher Auflösung verwendet wird. Dieses FORTRAN-Programm wurde für leichte bis mittelschwere oder fastmagische Kerne geschrieben, deren keV-Querschnitte charakterisiert sind einerseits durch breite s-Wellen-Resonanzen mit starker Multiniveau-Interferenz bei vernachlässigbarer Doppler-Verbreiterung, andererseits durch sehr schmale p-, d- ... Wellen-Resonanzen mit vernachlässigbarer Multiniveau-Interferenz bei starker Doppler-Verbreiterung. Dementsprechend dient der Reich-Moore-Vielniveau-Formalismus ohne Doppler-Verbreiterung zur Beschreibung der s-Wellen-Resonanzen, während die p-, d- ... Wellen-Resonanzen durch Einniveau-Formeln mit Doppler-Verbreiterung beschrieben sind. Die gerechneten Einfangausbeuten werden auflösungsverbreitert. Vielfachstöße werden mit Hilfe der Monte-Carlo-Methode simuliert. An Einfangdaten von bis zu fünf verschiedenen Flugzeitmessungen kann gleichzeitig angepaßt werden für Proben bestehend aus bis zu zehn verschiedenen Isotopen. Ein- und Ausgabebeispiele sowie eine FORTRAN-Liste sind beigelegt.

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1. Main Characteristics of the Code

In the present report a computer program is described which was developed for shape analysis of data on neutron capture by structural materials. The program, FANAC, extracts resonance parameters from high-resolution capture yields measured with the time-of-flight method and calculates neutron capture cross sections corrected for experimental effects such as self-shielding, multiple scattering, instrumental resolution and detector efficiency. It permits determination of up to 20 cross section parameters by simultaneously fitting calculated capture yield curves to experimental data from up to 5 time-of-flight measurements that may differ with respect to sample thickness, flight path or other experimental characteristics. The measured and calculated capture yield data and the resulting cross sections are plotted with a general-purpose plotting subroutine in use at Karlsruhe (subroutine PLOTA, Ref. 1).

The methods employed are similar to those of the programs FANAL (for shape analysis of transmission data, Ref. 2) and TACASI (for single-level analysis of transmission areas, capture areas and self-indication ratios, Ref. 4). The main advantage over the area analysis code TACASI consists in the possibility to treat many resonances simultaneously, the more reliable interpretation of incompletely resolved multiplets and the multi-level cross section formalism which allows a better description of multiple scattering. FANAC fits obtained between 6 and 165 keV for various iron and nickel isotopes are shown e.g. in Ref. 3 .

Normally one derives neutron resonance capture cross sections from capture yield data, measured with the time-of-flight method by detection of the prompt gamma radiation that is emitted after each capture event. The count rate observed in a narrow flight time interval with neutron flux ϕ , after dead-time and background correction, can be written as

$$c = \phi \gamma \epsilon \quad (1)$$

where y , called the capture yield, is the probability that an incoming neutron is captured, and ϵ is the detector efficiency. The data reduction consists of stripping off ϕ and ϵ to get y and then to extract from y the capture cross section. For thin samples $y \approx n\sigma_\gamma$, where n is the sample thickness in atoms/b and σ_γ the radiative-capture cross section. The flux ϕ can be measured with a reference sample having a sufficiently well known yield (e.g. gold); the efficiency (or efficiency ratio in case of a relative measurement) is determined by the detector characteristics.

In practice one measures the resolution-broadened capture yield

$$\bar{y}(E) = \int r(E, E') y(E') dE', \quad (2)$$

where $r(E, E')dE'$ is the probability that capture events induced by neutrons with energies E' in dE' are registered as if the energy were E .

A further complication arises especially for light- and medium weight nuclei from the fact that their scattering cross sections are very much larger than their capture cross sections. As a consequence self-shielding and multiple-scattering corrections are quite important for practical sample thicknesses and y cannot be taken as simply $n\sigma_\gamma$, but must be calculated as a more complicated functional of the total and capture cross sections.

Strictly speaking, the cross sections are Doppler-broadened by the thermal motion of the sample atoms. For light and medium-weight nuclides, however, Doppler broadening of typical s-wave resonances can be neglected. For the narrow p-, d- ... wave levels, on the other hand, Doppler broadening is important. Consequently, the FANAC program treats instrumental resolution according to Eq. (2) but neglects Doppler broadening for s-wave levels. It is therefore applicable only to resonance data where

the Doppler width $\Delta = \sqrt{4EkT/A}$ is much smaller than the width of the typical s-wave resonances. Narrow p-, d- ... wave levels are Doppler broadened.

The s-wave cross sections are parametrized with an R-matrix multi-level formula, whereas multi-level interference effects are neglected for narrow (p-, d-wave) levels. The program starts by calculating cross sections and capture yields from approximate values of the parameters. These starting values are then improved by application of the least-squares method (cf. e.g. Ref. 4,5). In order to make this method applicable the problem is linearized by Taylor expansion with respect to the cross section parameters and truncation after the linear terms. The solution of the linearized problem is thus an approximation which can be improved by iteration. The program iterates until the number of iterations reaches a prescribed limit or until the sum of squared deviations, χ_k^2 , differs by less than a given small fraction ϵ from that of the preceding step, χ_{k-1}^2 ,

$$\left| \frac{\chi_k^2 - \chi_{k-1}^2}{\chi_k^2} \right| < \epsilon, \quad (3)$$

where

$$\chi_k^2 = \sum_i \left(\frac{\eta_i - \bar{y}_i^{(k)}}{\delta \eta_i} \right)^2; \quad (4)$$

- η_i : i-th measured capture yield,
- $\delta \eta_i$: uncertainty of η_i ,
- $\bar{y}_i^{(k)}$: i-th calculated capture yield computed from parameters of k-th iteration.

The parameters to which χ^2 is most sensitive are

- resonance energies E_0 ;
- radiation widths Γ_γ and neutron widths Γ_n of strongly scattering (typically s-wave) levels with $\Gamma_n \gg \Gamma_\gamma$;
- the capture peak area parameter $g_J \Gamma_n \Gamma_\gamma / \Gamma$ for narrow (typically p-, d- or f-wave) levels with unknown $g_J \Gamma_n$. If $g_J \Gamma_n$ is known one can get $g_J \Gamma_\gamma$.

2. Formulae

2.1 Cross sections for $l = 0$

In order to explain the approximations employed in FANAC the relevant cross section formalism is briefly reviewed in this section.

The total and reaction cross sections for a single nuclide and $l = 0$ (s-wave interactions) can be written as

$$\sigma_{T0} = 2\pi\lambda_n^2 \sum_J g_J (1 - \text{Re } U_{nn}^J), \quad (5)$$

$$\sigma_{x0} = \pi\lambda_n^2 \sum_J g_J |\delta_{nx} - U_{nx}^J|^2 \quad (6)$$

where $2\pi\lambda_n$ is the neutron wave length in the center-of-mass system, δ_{nx} the Kronecker symbol and $g_J = (1/2)(2J+1)/(2I+1)$ the spin factor, with J and I the compound and target spin quantum numbers, respectively. The subscripts of the collision matrix elements U_{nn}^J and U_{nx}^J refer to reaction channels (n : elastic

scattering, x: arbitrary reaction). For $l = 0$ the spin sum consists of only one or two terms ($|I-1/2| \leq J \leq I+1/2$). According to R-matrix theory (Ref. 6) one can write

$$U_{CC'}^J = \Omega_C [(1-iK)^{-1} (1+iK)]_{CC'} \Omega_{C'} \quad (7)$$

with

$$K_{CC'} = \frac{1}{2} \sum_{\lambda} \frac{\Gamma_{\lambda C}^{1/2} \Gamma_{\lambda C'}^{1/2}}{E_{\lambda} - E} \quad (8)$$

$$\Omega_C = e^{-ik_C a_C} \quad (9)$$

where a_C is the R-matrix channel radius and the level sum in Eq. (8) runs over all (∞) s-wave levels with spin J . The partial-width amplitudes

$$\Gamma_{\lambda C}^{1/2} = (2k_C a_C)^{1/2} \gamma_{\lambda C} \quad (10)$$

are real quantities and vary with energy as $k_C^{1/2}$, where $k_C^2 = 1/\lambda_C^2 = 2m_C(E-E_C)/\hbar^2$ (m_C : reduced mass, E_C : reaction threshold, $\gamma_{\lambda C}$: energy-independent reduced width amplitude). The $\Gamma_{\lambda C}^{1/2}$ are to be understood as having the sign of $\gamma_{\lambda C}$.

In practice the level sum in Eq. (8) must be restricted to a finite number of resonances, typically those in a given energy interval. The levels outside this interval (distant levels) give rise to a K-matrix component, K^0 say, that is smooth inside the interval. We assume that direct reactions are negligible. In this case K^0 is diagonal and one can account for the distant levels omitted in Eq. (8) by using, instead of Eqs. (9) and (10),

$$\Omega_c = \exp[-i(k_c a_c - \arctan K_{cc}^0)] \quad (11)$$

$$\Gamma_{\lambda c}^{1/2} = (2k_c a_c)^{1/2} \frac{\gamma_{\lambda c}}{|1 - iK_{cc}^0|} \quad (12)$$

with

$$K_{cc}^0 = k_c a_c \sum_{\lambda}^{\prime} \frac{\gamma_{\lambda c}^2}{E_{\lambda} - E} \quad , \quad (13)$$

the prime indicating that the sum contains only terms from outside the interval ΔE (Ref. 6). Although the $\gamma_{\lambda c}$ and E_{λ} outside ΔE are mostly unknown one can estimate K_{cc}^0 by means of level-statistical theory:

Introducing the optical-model pole strength, $s_c = \langle \gamma_{\lambda c}^2 \rangle_{\lambda} / D_c$ (D_c : mean level spacing), one can replace the sum by an integration,

$$\sum_{\lambda}^{\prime} \frac{\gamma_{\lambda c}^2}{E_{\lambda} - E} \approx P \int_{-\infty}^{\infty} dE' \frac{s_c(E')}{E' - E} - P \int_{\bar{E} - \Delta E/2}^{\bar{E} + \Delta E/2} dE' \frac{s_c(E')}{E' - E} \quad , \quad (14)$$

where $P \int$ indicates Cauchy's principal value, whereas ΔE and \bar{E} are length and mid-point of the energy interval excluded from the sum. The first integral is the so-called distant-level parameter, usually denoted by R_c^{∞} , that describes the influence of the energy variation of s_c (including the vanishing of s_c below the ground state energy). Furthermore one can neglect the energy variation of s_c and R_c^{∞} in the interval if ΔE is small compared to the spacing of optical-model size resonances (~ 2 MeV). Then

$$K_{cc}^0(E) \approx k_c a_c R_c^{\infty}(\bar{E}) + 2k_c a_c s_c(\bar{E}) \arctan \frac{E - \bar{E}}{\Delta E/2} \quad . \quad (15)$$

The contribution of distant levels can thus be estimated from the level-statistical parameters R_C^∞ and s_C or, equivalently, from the effective nuclear radius

$$R_C' = a_C (1 - R_C^\infty) \quad (16)$$

and the s-wave strength function

$$S_{CO} = 2k_C a_C s_C \sqrt{\frac{1 \text{ eV}}{E - E_C}} \quad (17)$$

After thus reducing the number of levels to a manageable size one can reduce the dimensionality of the matrix K. Following Reich and Moore (Refs. 7,8) one applies the Teichmann-Wigner channel elimination prescription (Ref. 6) to the photon channels and then introduces the approximation

$$\sum_{C \in \gamma} \Gamma_{\lambda C}^{1/2} \Gamma_{\mu C}^{1/2} \approx \delta_{\lambda \mu} \sum_{C \in \gamma} \Gamma_{\lambda C} = \delta_{\lambda \mu} \Gamma_{\lambda \gamma}, \quad (18)$$

where the summation is over photon channels and $\Gamma_{\lambda \gamma}$ is the total radiation width. This is justified by the usually large number of radiative transitions that can de-excite the compound state and by the fact that the $\Gamma_{\lambda C}^{1/2}$ are symmetrically distributed around zero so that the cross terms in the sum tend to cancel, in contrast to the terms with $\lambda = \mu$. The result is the Reich-Moore prescription

$$\begin{aligned} U_{CC'}^J &= \Omega_C [(1-iK)^{-1} (1+iK)]_{CC'} \Omega_{C'} \\ &= \Omega_C [2(1-iK)^{-1} - 1]_{CC'} \Omega_{C'} \end{aligned} \quad (19)$$

$$K_{cc'} = \frac{1}{2} \sum_{\lambda} \frac{\Gamma_{\lambda c}^{1/2} \Gamma_{\lambda c'}^{1/2}}{E_{\lambda} - E - i\Gamma_{\lambda \gamma} / 2}, \quad (20)$$

valid inside the interval of explicitly given resonances ($\bar{E} - \Delta E/2 < E_{\lambda} < \bar{E} + \Delta E/2$) and for particle channels ($c, c' \notin \gamma$). All cross sections except those for neutron capture can now be obtained from Eqs. (5), (6), while the capture cross sections must be calculated as

$$\sigma_{\gamma 0} = \sigma_{T0} - \sum_{x \notin \gamma} \sigma_{x0} = \pi \lambda_n^2 \sum_J g_J (1 - \sum_{x \notin \gamma} |U_{nx}^J|^2). \quad (21)$$

The last equation follows from the unitarity of the collision matrix. The Reich-Moore formalism as a genuine R matrix formalism guarantees this unitarity: The restriction (18) for the partial radiation widths affects neither the symmetry nor the reality of the K-matrix. The same is true for the approximation (11), (15) to the distant levels.

In the FANAC program s-wave cross sections are calculated according to the Eqs. (5), (6), (11), (15), (19), (20), (21). The number of particle channels is restricted to 2 (elastic channel plus at most one inelastic channel), so that $1-iK$ is at most a 2x2 matrix which is easily inverted.

2.2 Cross sections for $l \geq 1$

The cross sections for partial waves with $l \geq 1$ are described in FANAC by

$$\sigma_{T1} = \sum_{\lambda} (\sigma_0 \psi)_{\lambda} + 4\pi \lambda_n^2 \cdot 3 \sin^2(k_n R_1' - \arctan k_n R_1'), \quad (22)$$

$$\sigma_{x1} = \sum_{\lambda} \left(\frac{\Gamma_x}{\Gamma} \sigma_o \psi \right)_{\lambda} \quad (x=n, n', \gamma), \quad (23)$$

where

$$\sigma_o = 4\pi\lambda_n^2 g_J \frac{\Gamma_n}{\Gamma}, \quad (24)$$

$$\psi = \int_{-\infty}^{\infty} dE' \frac{\exp[-(E-E')^2/\Delta^2]}{\Delta\sqrt{\pi}} \frac{\Gamma^2/4}{(E'-E_o)^2 + \Gamma^2/4}, \quad (25)$$

$$\Delta = \sqrt{\frac{4kTE_o}{A}}. \quad (26)$$

Thus levels associated with $l \geq 1$ are treated as Doppler-broadened, non-interfering resonances. The total cross section σ_{T1} contains also the potential-scattering term for p-wave scattering. Potential scattering for $l \geq 2$ and potential-resonance interference for $l \geq 1$ is neglected. The quantity kT in Eq. (26) is the Lamb-corrected (Ref. 9) sample temperature in energy units, E_o the resonance energy and A the nuclear mass divided by the neutron mass.

The cross sections are then

$$\sigma_T = \sigma_{T0} + \sigma_{T1}, \quad (27)$$

$$\sigma_x = \sigma_{x0} + \sigma_{x1} \quad (x=n, n'), \quad (28)$$

$$\sigma_{\gamma} = \sigma_T - \sigma_n - \sigma_{n'}. \quad (29)$$

The cross section parameters are (for each isotope, spin and parity)

- the potential-scattering parameters

R_J^i and S_{OJ} for the s-wave channel(s) (cf. Eqs. (16), (17)),
 R_1^i for the p-wave (Eq. (22)),

- the resonance parameters

E_λ , $\Gamma_{\lambda n}$, $\Gamma_{\lambda n'}$, $\Gamma_{\lambda \gamma}$ and, if $\Gamma_n \neq 0$, the relative signs
 $\text{sgn}(\Gamma_{\lambda n}^{1/2} \Gamma_{\lambda n'}^{1/2}) = \text{sgn}(\gamma_{\lambda n} \gamma_{\lambda n'})$.

For non-monotopic samples the isotopic cross sections must, of course, be multiplied by the appropriate abundances (and, in the case of σ_γ , by the detection efficiencies) and summed over all isotopes present in the sample, as explained in the next section.

2.4 Capture yields

The capture yield can be written as a collision series,

$$Y = Y_0 + Y_1 + Y_2 + \dots \quad (30)$$

where the subscripts 0, 1, 2, ... refer to the number of scattering collisions after which capture occurs. In order to be more specific we need

- the probability $\exp(-n\sigma_T)$ for a neutron to traverse without interaction a layer of material of thickness n (atoms/b);
- the probability $dn(d\sigma_n/d\Omega)d\Omega$ that scattering into a solid-angle element $d\Omega$ occurs within an infinitesimal sample layer dn ;
- the corresponding probability $dn\sigma_\gamma$ for capture in a layer dn .

We can then write, for a sample of uniform thickness n , and cross sections σ_T , $d\sigma_n/d\Omega$, σ_γ for the incident energy,

$$\begin{aligned}
 y_0 &= \int_0^n e^{-n'\sigma_T} dn' \sigma_\gamma \\
 y_1 &= \int_0^n e^{-n'\sigma_T} dn' \int_{4\pi} d\Omega \frac{d\sigma_n}{d\Omega} \int_0^{n_1} e^{-n_1'\sigma_{T1}} dn_1' \sigma_{\gamma 1} \\
 y_2 &= \int_0^n e^{-n'\sigma_T} dn' \int_{4\pi} d\Omega \frac{d\sigma_n}{d\Omega} \int_0^{n_1} e^{-n_1'\sigma_{T1}} dn_1' \int_{4\pi} d\Omega_1 \frac{d\sigma_{n1}}{d\Omega_1} \int_0^{n_2} e^{-n_2'\sigma_{T2}} dn_2' \sigma_{\gamma 2} \\
 &\text{etc.} \qquad \qquad \qquad (31)
 \end{aligned}$$

where the subscripts 1,2,... refer to the number of preceding scattering collisions. The cross sections σ_{Tk} , $\sigma_{\gamma k}$, $d\sigma_{nk}/d\Omega_k$ depend on the neutron energy after k scattering collisions, which in turn depends on the angles and nuclear masses involved in all preceding collisions. Similarly the maximum material thickness n_k that a neutron can traverse without interaction after the k -th scattering collision depends on the spatial and angular coordinates of all k preceding collisions.

Performing the last integration in each term y_k one finds

$$\begin{aligned}
 y_0 &= (1 - e^{-n\sigma_T}) \frac{\sigma_\gamma}{\sigma_T} \\
 y_1 &= (1 - e^{-n\sigma_T}) \frac{\sigma_n}{\sigma_T} \left\langle (1 - e^{-n_1\sigma_{T1}}) \frac{\sigma_{\gamma 1}}{\sigma_{T1}} \right\rangle_1 \\
 y_2 &= (1 - e^{-n\sigma_T}) \frac{\sigma_n}{\sigma_T} \left\langle (1 - e^{-n_1\sigma_{T1}}) \frac{\sigma_{n1}}{\sigma_{T1}} \left\langle (1 - e^{-n_2\sigma_{T2}}) \frac{\sigma_{\gamma 2}}{\sigma_{T2}} \right\rangle_2 \right\rangle_1 \\
 &\text{etc.} \qquad \qquad \qquad (32)
 \end{aligned}$$

where the brackets

$$\langle \dots \rangle_k \equiv \int_0^{n_k \sigma_{Tk}} dn'_k \sigma_{Tk} \frac{e^{-n'_k \sigma_{Tk}}}{1 - e^{-n_k \sigma_{Tk}}} \int \frac{d\Omega_k}{4\pi} \frac{d\sigma_{nk}}{d\Omega_k} \dots \quad (33)$$

denote averages over all possible scattering angles θ_k and azimuths ϕ_k ($d\Omega_k = d(\cos\theta_k)d\phi_k$) and all possible places for the k -th collision, the frequency distributions on the right-hand side of Eq. (33) being properly normalized to unity. The upper limit $n_k \sigma_{Tk}$ for the spatial distribution of interaction points along the neutron trajectory is just the maximum number of mean free paths that a neutron can traverse before the k -th collision, i.e. the distance between the $(k-1)$ -th collision point (or, for $k=1$, the entrance point) and the sample surface. The quantities occurring in Eq. (32) can be interpreted as "thin sample" scattering and capture yields $n_k \sigma_{nk}$ and $n_k \sigma_{\gamma k}$, corrected for beam attenuation (self-shielding) by the factors $(1 - \exp(-n_k \sigma_{Tk})) / (n_k \sigma_{Tk})$.

Since each average $\langle \dots \rangle_k$ implies three integrations as shown by Eq. (33) the dimensionality of the integrals to be calculated for y_k increases rapidly with k , and only the first-collision yield is a simple function of cross sections and sample thickness.

Already the second-collision yield looks rather awkward even in the simple case of infinite-slab geometry:

$$y_1 = (1 - e^{-n\sigma_T}) \frac{\sigma_n}{\sigma_T} \int_{-1}^1 d\mu p(\mu) \left(1 - \frac{1 - e^{-t-t'}}{t-t'} \frac{t}{1 - e^{-t}} \right) \quad (34)$$

with

$$t \equiv \text{sgn } \mu \cdot n\sigma_T, \quad t' \equiv \text{sgn } \mu \cdot \frac{n\sigma_T l}{\mu} \quad (35)$$

where μ is the cosine of the scattering angle and $p(\mu)d\mu$ the probability for μ in $d\mu$.

Higher-order terms in the collision expansion are increasingly more complicated functionals of the cross sections σ_T , σ_n , σ_γ .

2.4 Monte Carlo calculation of multiple-collision yields

The Monte Carlo method is best suited to calculate multi-dimensional integrals like those in Eqs. 32. In our case one has to simulate a sufficient number of multiple-collision events by sampling the neutron beam profile for the entry point and then, for successive collisions, the path length before the collision from

$$p(s')ds' = \frac{e^{-s'}}{1-e^{-s}} ds' \quad (0 < s' < s), \quad (36)$$

the c.m.s. scattering angle θ_c from

$$p(\mu_c)d\mu_c = \frac{d\mu_c}{2} \quad (-1 \leq \mu_c \equiv \cos\theta_c \leq 1), \quad (37)$$

the azimuth ϕ_c from

$$p(\phi_c)d\phi_c = \frac{1}{2\pi} d\phi_c \quad (0 \leq \phi_c < 2\pi). \quad (38)$$

In writing down Eq. 37 we assumed isotropic c.m.s. scattering. In this case the simplest sampling technique is applicable to all three distributions: One samples a distribution $p(x)dx$ by generating a random number ρ between 0 and 1 and solving the equation

$$\rho = \int_0^x p(x')dx' \quad (39)$$

for x . From the sampled c.m.s. quantities μ_c and ϕ_c one gets the corresponding quantities in the lab system,

$$\mu = \frac{A\mu_C + 1}{\sqrt{A^2 + 2A\mu_C + 1}} = \cos\theta, \quad (40)$$

$$\phi = \phi_C \quad (41)$$

and the new energy

$$E' = E \frac{A^2 + 2A\mu_C + 1}{(A+1)^2} \quad (42)$$

Although Eqs. 40-42 are valid only for target nuclei at rest and purely elastic scattering they are used in FANAC for target nuclei in thermal motion and also for inelastic collisions. Now Monte Carlo tests with the TACASI code (Ref. 4) showed that in the absence of inelastic processes the stationary-target approximation does not lead to significant errors in multiple-collision capture yield calculations even for very thick samples, except perhaps for very light sample nuclei (see Ref. 4). Application of Eqs. 40-42 to inelastic collisions, however, has no justification other than convenience and the fact that for light and medium-weight nuclei inelastic scattering in the keV region is usually much less probable than elastic scattering. The FANAC code must therefore be used with caution if both inelastic scattering and multiple collisions are important, i.e. if $\sigma_n \gtrsim \sigma_n \gg \sigma_\gamma$ and $n\sigma_T \gtrsim 1$.

Once the point of collision and the angles are established for the k-th collision one can calculate E' and find the corresponding cross sections $\sigma_{Tk} = \sigma_T(E')$, $\sigma_{nk} = \sigma_n(E')$ and $\sigma_{\gamma k} = \sigma_\gamma(E')$. The sample thickness n_k to the surface is obtained as follows:

The polar angle θ was introduced above as the lab scattering angle. This implies use of a reference frame which has its polar axis parallel to the pre-collision velocity of the neutron. For the calculation of n_k , however, a lab reference frame is needed in which the sample surface has a simple description. We therefore take as the lab system a reference frame which has its z-axis coincident with the axis of the cylindrical disc sample. The transformation between both reference frames can be established as follows.

Let the neutron directions before and after the collision be given by the unit vectors $\vec{\Omega}$ and $\vec{\Omega}'$. In order to calculate the co-ordinates of $\vec{\Omega}'$ from those of $\vec{\Omega}$ and from the scattering angles θ and ϕ we introduce an intermediate reference frame which has its z-axis parallel to $\vec{\Omega}$, so that

$$\vec{\Omega} = (0, 0, 1), \quad (43)$$

$$\vec{\Omega}' = (\sin\theta \cos\phi, \sin\theta \sin\phi, \cos\theta). \quad (44)$$

In the lab system we introduce polar co-ordinates by writing (see Fig. 1)

$$\vec{\Omega} = (\Omega_x, \Omega_y, \Omega_z) = (\sin\zeta \cos\eta, \sin\zeta \sin\eta, \cos\zeta) \quad (45)$$

Now a unit vector can be brought from the position $(0, 0, 1)$ into the position $(\sin\zeta \cos\eta, \sin\zeta \sin\eta, \cos\zeta)$ by two successive rotations (Fig. 2):

- (1) rotation about the y'' -axis, through an angle ζ ,

$$\begin{pmatrix} x' \\ y' \\ z' \end{pmatrix} = \begin{pmatrix} \cos\zeta & 0 & \sin\zeta \\ 0 & 1 & 0 \\ -\sin\zeta & 0 & \cos\zeta \end{pmatrix} \cdot \begin{pmatrix} x'' \\ y'' \\ z'' \end{pmatrix}, \quad (46)$$

(2) rotation about the z'' -axis, through an angle η ,

$$\begin{pmatrix} x \\ y \\ z \end{pmatrix} = \begin{pmatrix} \cos\eta & -\sin\eta & 0 \\ \sin\eta & \cos\eta & 0 \\ 0 & 0 & 1 \end{pmatrix} \cdot \begin{pmatrix} x' \\ y' \\ z' \end{pmatrix}. \quad (47)$$

Combining the two transformations and expressing the polar coordinates of $\vec{\Omega}$ by its Cartesian coordinates,

$$\begin{aligned} \cos\zeta &= \Omega_z, & \cos\eta &= \frac{\Omega_x}{\sqrt{1-\Omega_z^2}}, \\ \sin\zeta &= \sqrt{1-\Omega_z^2}, & \sin\eta &= \frac{\Omega_y}{\sqrt{1-\Omega_z^2}}, \end{aligned} \quad (48)$$

we find eventually the lab coordinates of $\vec{\Omega}'$

$$\begin{pmatrix} \Omega'_x \\ \Omega'_y \\ \Omega'_z \end{pmatrix} = \begin{pmatrix} \Omega_x \Omega_z (1-\Omega_z^2)^{-1/2} & -\Omega_y (1-\Omega_z^2)^{-1/2} & \Omega_x \\ \Omega_y \Omega_z (1-\Omega_z^2)^{-1/2} & \Omega_x (1-\Omega_z^2)^{-1/2} & \Omega_y \\ -(1-\Omega_z^2)^{+1/2} & 0 & \Omega_z \end{pmatrix} \begin{pmatrix} \sin\theta \cos\phi \\ \sin\theta \sin\phi \\ \cos\theta \end{pmatrix} \quad (49)$$

If $\Omega_z^2 = 1$ this expression fails and is to be replaced by

$$\begin{pmatrix} \Omega'_x \\ \Omega'_y \\ \Omega'_z \end{pmatrix} = \begin{pmatrix} \sin\theta \cos\phi \\ \sin\theta \sin\phi \\ \cos\theta \end{pmatrix} \Omega_z. \quad (50)$$

Now we can calculate $n_k \sigma_{Tk}$, the number of mean free paths in the direction $\vec{\Omega}$ between the last collision point, (x_0, y_0, z_0) , and the sample surface. The equations of the neutron trajectory are

$$\frac{x-x_0}{\Omega_x} = \frac{y-y_0}{\Omega_y} = \frac{z-z_0}{\Omega_z} , \quad (51,52)$$

Let the surfaces of the cylindrical-disc sample be given by

$$z = 0 \quad (\text{front face}), \quad (53)$$

$$z = n \quad (\text{back face}), \quad (54)$$

$$x^2 + y^2 = r^2 \quad (\text{cylinder surface}), \quad (55)$$

where all distances are expressed in units of atoms/b.

The point where the neutron trajectory intersects the cylinder is defined by the solution of Eqs. 51, 52 and 55,

$$\begin{aligned} x &= x_0 + \Omega'_x d_0 \\ y &= y_0 + \Omega'_y d_0 \\ z &= z_0 + \Omega'_z d_0 \end{aligned} \quad (56)$$

$$\text{with } d_0 = \frac{\sqrt{b^2 + ac - b}}{a} , \quad (57)$$

$$\text{and } a = \Omega_x^2 + \Omega_y^2 , \quad (58)$$

$$b = x_0 \Omega_x + y_0 \Omega_y , \quad (59)$$

$$c = r^2 - x_0^2 - y_0^2 . \quad (60)$$

The quantity d_0 is just the distance to the cylinder surface. This must be compared to the distance to the front face,

$$d_1 = -\frac{z_0}{\Omega_z}, \quad (61)$$

if $\Omega_z < 0$, and to the distance to the back face,

$$d_2 = \frac{n-z_0}{\Omega_z}, \quad (62)$$

if $\Omega_z > 0$, the smaller quantity being the distance n_k (in atoms/b) that the neutron must traverse before it can escape from the sample:

$$n_k = \begin{cases} \min(d, d_1) & \text{if } \Omega_z < 0, \\ d_0 & \text{if } \Omega_z = 0, \\ \min(d, d_2) & \text{if } \Omega_z > 0. \end{cases} \quad (63)$$

One can now compute the quantities

$$[1 - \exp(-n_k \sigma_{Tk})] \cdot \sigma_{nk} / \sigma_{Tk} \quad \text{and} \quad [1 - \exp(-n_k \sigma_{Tk})] \cdot \sigma_{\gamma k} / \sigma_{Tk}$$

and eventually the products

$$\prod_{j=0}^{k-1} (1 - e^{-n_j \sigma_{Tj}}) \frac{\sigma_{nj}}{\sigma_{Tj}} (1 - e^{-n_k \sigma_{Tk}}) \frac{\sigma_{\gamma k}}{\sigma_{Tk}} \quad (64)$$

whose average over many such simulated multiple-collisions events constitutes the Monte Carlo estimate of y_k .

2.5 Detector efficiencies

For non-monotopic samples the cross sections from which the capture yields must be calculated are actually abundance-weighted sums over all isotopes present in the sample, $\sigma_T = \sum a_i \sigma_{Ti}$, $\sigma_X = \sum a_i \sigma_{Xi}$, with the exception of the capture cross section for which the weighting must also account for the usual isotope dependence of the capture detector efficiency (see e.g. Ref. 10). The required quantity is not the yield itself but rather the "observable yield", $y\epsilon \equiv \sum y_i \epsilon_i$. It is obtained if the cross section for detected capture, $\sigma_Y \epsilon \equiv \sum a_i \sigma_{Yi} \epsilon_i$, is used instead of $\sigma_Y = \sum a_i \sigma_{Yi}$. For nuclei with low level density (in the light-to-medium mass range or near closed shells) capture γ -ray spectra fluctuate significantly from resonance to resonance. The efficiencies ϵ_i depend therefore on energy if the capture detector is sensitive to these fluctuations. To the extent that level-level interference may be neglected for the capture cross section one can assume a different γ -ray spectrum and hence a different efficiency associated with each resonance. The efficiency to be used at a given energy is then

$$\epsilon_i = \frac{\sum_{\lambda} \epsilon_{i\lambda} \sigma_{\gamma\lambda}}{\sum_{\lambda} \sigma_{\gamma\lambda}}, \quad (65)$$

where each $\sigma_{\gamma\lambda}$ is a single-level Breit-Wigner resonance term for the i -th isotope.

2.6 Resolution function

The resolution function can usually be taken as Gaussian,

$$r(E, E') dE' = \frac{1}{W\sqrt{\pi}} e^{-(E-E')^2/W^2} dE, \quad (66)$$

with
$$W^2 = E^2 \left[c_1 \left(\frac{\Delta L}{L} \right)^2 + c_2 \left(\frac{\Delta t}{L} \right)^2 E \right] \quad (67)$$

where δL is the thickness of the capture sample and Δt the observed half width (FWHM) of the γ peak in the time-of-flight spectrum resulting from photons produced in the accelerator target together with each neutron pulse. This form of the resolution function accounts for

- (1) finite burst width of the accelerator-pulsed neutron source (Δt)
- (2) finite channel width of the flight time analyzer (Δt),
- (3) flight path differences due to finite sample thickness (ΔL),
- (4) effects of detector size such as path differences of capture gamma rays and scintillation photons (Δt),
- (5) electronic time jitter (Δt).

The constants $c_1 = 2/3$ and $c_2 = (60.259 \text{ MeV}^{1/2} \text{ ns m}^{-1})^{-2}$ guarantee the correct variance.

3. The Code

The modular structure of the FANAC code is shown in Fig. 3. The main program is essentially a series of CALL statements for the various subroutines, with a loop on iterations. The iterative least-squares procedure stops when either the maximum number of iterations specified in the input is completed or chi-squared remains essentially unchanged in successive iterations, as explained in Sect. 1, Eqs. 3,4. Then input for another similar job can be read in.

3.1 Subroutines

Although comments are generously used throughout the code (see Appendix) a brief description of the various subroutines may help to understand their functions.

CARDIN

CARDIN is called by the main program to read card input and to print sample and time-of-flight run specifications. Spin factors and compound spins are calculated. Then subroutine PAROUT (see below) is called to print the cross section parameter input. In the present version of the FANAC code CARDIN also calls a special subroutine, ETEC, that corrects truncation errors in channel energies caused by the data acquisition system. If the input energies are free of such errors one can eliminate subroutine ETEC and the associated CALL and WRITE statements in CARDIN (cards 2240 to 2390).

INDEX

This subroutine provides, for each isotope and compound spin, subscripts structuring the one-dimensional arrays containing the cross section parameters and their uncertainties:

MP(I,J) is the subscript of the first of two potential-scattering parameters (the two are distant-level strength function and effective nuclear radius),

MR(I,J) is the subscript of the first resonance parameter (there are four for each resonance, E_0 , Γ_n , Γ_n' , and Γ_γ)

for the I-th isotope and the J-th compound spin.

Cross section parameters are counted;

MX is their total number,

MA is the number of those which must be adjusted.

CONIJ

This subroutine calculates certain constants for each isotope (I):

D(I) is a constant which multiplied by \sqrt{E} gives the Doppler width Δ (see Eq. 26),
 CT(I,J) and CG(I,J) give $2\pi\lambda^2 a g_J$ and $\pi\lambda^2 a g_J$ after division by the neutron energy (a: abundance, g_j : spin factor)

MEV

MEV converts energies and resonance widths from keV to MeV and calculates the adjusted quantities $\Gamma_n^{1/2}$ and $\Gamma_{n'}^{1/2}$ from the input quantities Γ_n and $\Gamma_{n'}$, taking over the signs so that $\text{sgn } \Gamma_n^{1/2} = \text{sgn } \Gamma_n$, $\text{sgn } \Gamma_{n'}^{1/2} = \text{sgn } \Gamma_{n'}$.

CONE

This subroutine, called by the main program, calculates for each iteration the coefficients

$$A_{\mu\nu} = \sum_i \frac{Y_{i,\mu}}{\delta\eta_i} \frac{Y_{i,\nu}}{\delta\eta_i}, \quad (69)$$

$$c_\mu = \sum_i \frac{Y_{i,\mu}}{\delta\eta_i} \frac{\eta_i - Y_i}{\delta\eta_i} \quad (70)$$

of the normal equations

$$\sum_\nu A_{\mu\nu} (x_\nu - x_{0\nu}) = c_\mu \quad (71)$$

for the linearized least-squares problem, where

- $x_{0\nu}$ is the initial approximation to the ν -th adjusted parameter,
- x_ν the improved value to be calculated,
- Y_i the i -th calculated capture yield
- $Y_{i,\mu}$ its derivative with respect to the μ -th adjusted parameter,
- η_i the i -th observed capture yield and
- $\delta\eta_i$ its uncertainty.

The y_i and $y_{i,\mu}$ are calculated by numeric broadening of yields and yield derivatives obtained from the subroutines PRY, SEY and MUY. Strictly speaking the derivatives $y_{i,\mu}$ are taken as those of the primary yields because of the difficulty to get derivatives from Monte Carlo calculations. This is tantamount to neglect of errors and uncertainties in calculated multiple collision yields, so that

$(\eta_i - y_i) / \delta \eta_i = (\eta_{oi} - y_{oi}) / \delta \eta_i$. In other words, measured yields are reduced in each step to primary yields by subtraction of calculated multiple-collision contributions. The resulting primary yields are then fitted by parameter adjustment.

The required resolution function values come from the subroutines TGAUSS or CHISQ4. Numeric integration according to Simpson's rule is performed by subroutine SIMP. Furthermore CONE prints a table of measured and calculated yields ((for details see Sect. 3.3 below)).

PRY

PRY produces unbroadened primary capture yields (y_0 in Eq. 32) and derivatives for the calling subroutine CONE. First an equidistant energy grid is established on the basis of the smallest resolution and Doppler width encountered. The grid is chosen so that for each incident energy E at least five integrand points are available between $E-\Delta$ and $E+\Delta$ for the resolution broadening, Δ being the Doppler width. If the resolution width and/or the energy interval to be analysed is very much larger than Δ it may happen that the number of grid points exceeds the available storage (201 points within six resolution widths, a total of 2048 grid points). In this case the Doppler width is artificially increased until the narrow resonances are sufficiently broad to allow an adequate description with the available number of grid points. This leaves s-wave results unaffected but p-wave results only to the extent that sample-thickness effects are small. It should be realized that such an artificial raise in temperature tends to increase the primary yield because the self-shielding factor $(1 - \exp(-n\sigma_T)) / (n\sigma_T)$ across a narrow

resonance increases. On the other hand multiple-collision yields are reduced because the sample becomes more transparent in the peak region. Thus a partial compensation could be expected even in cases where the sample is not "thin" ($n\sigma_T \ll 1$) at resonance. For narrow levels, whose width is much smaller than the average neutron energy loss per collision ($\Delta^2 + (\Gamma/2)^2 \ll (2E/A)^2$), the self-shielding effect predominates, however, and the peak area parameters $g_n \Gamma_\gamma / \Gamma$ calculated with artificially raised temperature are somewhat too small. If the limitation to 2048 internal grid points causes the difficulties one can avoid temperature readjustment by reducing the analysis interval, treating only one resonance at a time in the extreme case.

The internal grid extends beyond the boundaries of the analysis region in order to allow proper resolution broadening and multiple-collision treatment at these boundaries.

After having established the internal grid PRY calculates observable primary yields including efficiencies and their derivatives for all grid points. The required cross sections and cross section derivatives come from subroutine XSECT, efficiency weighting factors (Eq.65) from subroutine EFFI (see below).

SEY

SEY, called from CONE, calculates secondary yields according to Eq. 34 for all internal grid points. Since Eq. 34 is approximate and neglects effects of lateral sample extension and higher-order collisions the results are only used as weights for importance sampling in the Monte Carlo subroutine MUY.

MUY

MUY simulates multiple-collision events and calculates multiple-collision yields for all internal grid energies.

The principle of importance sampling is applied. This means that small capture yield contributions are computed with lower statistical accuracy (less neutron histories) than large contributions, in order to keep costs low. One tries to maximize the accuracy of $\sum_k y_k$ with the constraint $\sum_k N_k = N$, where the subscript $k=0,1,2,\dots$ gives the number of completed scattering collisions, N_k is the number of simulations for the k -th order capture yield y_k and N the total number of simulations. With the plausible assumption that the accuracy of y_k is proportional to $1/\sqrt{N_k}$ one gets, by solving the extremum problem,

$$N_{k+1} : N_k = Y_{k+1} : Y_k \quad (72)$$

Accordingly the initial number of histories for each grid point is taken as

$$N_1 = \frac{Y_1}{\max y_1} N_0 \quad (73)$$

where N_0 , the maximum number of simulations per grid point, is specified in the input and pertains to the energy point where y_1 has its maximum (In order to get sufficient statistics N_1 is set equal to 10 if Eq. 73 gives less). For higher-order collisions ($k>1$) N_{k+1} can be estimated with the geometric series approximation

$$Y_{k+1} : Y_k = Y_k : Y_{k-1} \quad (74)$$

as

$$N_{k+1} = N_k \frac{Y_k}{Y_{k-1}}, \quad \text{if } Y_k \leq Y_{k-1} \quad (75)$$

For $y_k > y_{k-1}$ one encounters difficulties since not enough pre-collision angles, coordinates and energies are stored. In order to avoid "splitting" /Ref. 11/ we set

$$N_{k+1} = N_k \quad , \quad \text{if } y_k > y_{k-1} \quad . \quad (75)$$

MUY starts the simulations of N_1 multiple-collision events by sampling for each one the profile of the incident beam: The neutron density in the beam is taken as constant up to a certain fraction (EDGE) of the sample radius, as zero beyond. (In the present version EDGE is set equal to 0.8). After having thus established the points of incidence MUY samples coordinates and angles for N_1 first collisions. Scattering is taken as isotropic in the center-of-mass system as mentioned before. The new energies (Eq. 42) and the corresponding cross sections are found next, the latter by linear interpolation between stored values. The new direction cosines (Eq. 40,41,49) and the number of mean free paths, $n_1 \sigma_{T1}$, between the collision points and the sample surface (Eq. 56-63) are then calculated. Finally the scattered and captured fractions of the incident neutrons can be computed. Averaging the latter over all N_1 first collisions one gets the desired Monte Carlo estimate of y_1 .

After establishing, on the basis of y_1 , the number N_1 of second collisions to be simulated /Eq.74/ the computation proceeds in exact analogy to second, then to third etc. collisions until, after k collisions, N_{k+1} drops below 0.5, or until the 20th collision is completed.

For the relatively short flight paths used in Van-de-Graaff measurements the times spent by neutrons inside the sample may constitute noticeable fractions of the total registered flight times. Therefore MUY keeps track of these time delays and deposits calculated yields in correspondingly shifted

energy bins, applying some additional spreading so that the statistical fluctuations of the Monte Carlo results are smoothed out to a certain degree.

A last point should be mentioned. For the multiple-collision simulation the atomic weight of all sample nuclei is taken as the same, namely as that of the first isotope in the input. The small errors introduced in lab scattering angles and energies (cf. Eqs. 41,43) are usually tolerable in view of the other approximations invoked. If hydrogen or other very light nuclei contribute significantly to multiple scattering, however, it may be necessary to take the predominant scatterer as the first isotope in the input.

For users of the TACASI code (Ref. 4) it may be pointed out that the method adopted in MUY corresponds to the "FS-mode" (finite sample, stationary target nuclei) in TACASI.

EFFI

Subroutine EFFI, called by PRY in case not all resonance efficiencies are equal, provides energy-dependent efficiencies according to Eq. 65.

XSECT

Subroutine XSECT, called by PRY, calculates the required cross sections, cross section derivatives and transmission values for a given energy. Cross sections for the pure first isotope are also calculated for purpose of impurity correction. The formalism utilized for partial waves with $l=0$ and $l>1$ is described in Sects. 2.1 and 2.2 above. The inverse of the matrix $1-iK$ (cf. Eq. 19) and its derivatives are furnished by the subroutine KMAT.

KMAT

Subroutine KMAT, called by XSECT, calculates K-matrix elements (Eq. 20), their derivatives and the elements of $(1-iK)^{-1}$.

CHOBAN

The solution of the system of normal equations (71) for the improved parameter vector requires inversion of the matrix ($A_{\mu\nu}$). This is accomplished by the subroutine CHOBAN. The method employed is that of Choleski and Banachiewicz as described e.g. in Ref. 12. The result is the covariance matrix ($B_{\mu\nu}$).

ADJ

The improved parameter vector is calculated in subroutine ADJ by multiplication of the covariance matrix ($B_{\mu\nu}$) into the deviation vector c_{μ} (Eq. 70) and addition of the resulting adjustment vector to the unadjusted parameter vector. The uncertainties of the adjusted parameters are calculated as

$$\delta x_{\mu} = \sqrt{B_{\mu\mu} \frac{\chi^2}{N-M}} \quad (76)$$

where χ^2 is the sum of squared deviations over all N utilized data points and M the number of adjusted parameters.

For strong multiple-collision capture the convergence could be improved if the adjustment vector was reduced by 20-30 % in each step. In the present FANAC version the reduction factor (FUDGE) is set equal to 0.75 (see DATA statement).

KEV

The subroutine KEV converts energies and resonance widths back to keV prior to printing and plotting.

YPLOT

Plots of the results are prepared by subroutine YPLOT which calls a standard plotting subroutine PLOTA in use at the Karlsruhe nuclear research center (Ref. 1). The meaning of the arguments of PLOTA is explained on comment cards (see appendix) in order to facilitate replacement of PLOTA by equivalent plotting subroutines elsewhere.

PAROUT

At the start and after each iteration the cross section parameters and associated information such as abundances, atomic weights, spin quantum numbers etc. are printed by PAROUT as illustrated in Sect. 3.3 below.

SIMP

The subroutine SIMP performs integrations according to Simpson's rule. It is utilized in the resolution-broadening calculation in CONE.

VOIGT

The subroutine VOIGT is an adaptation of the subroutine PSIXI of the group-constant program MIGROS (Ref. 13). It calculates the symmetric and asymmetric Doppler-broadened line shape functions (Voigt profiles)

$$\psi(x, \beta) = \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} \frac{dx'}{\beta} \exp \left[-\left(\frac{x-x'}{\beta}\right)^2 \right] \frac{1}{1+x'^2} , \quad (77)$$

$$\phi(x, \beta) = \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} \frac{dx'}{\beta} \exp \left[-\left(\frac{x-x'}{\beta}\right)^2 \right] \frac{x'}{1+x'^2} \quad (77')$$

that are needed in subroutine XSECT for the description of narrow resonances as described in Sect. 2.2 above, and for calculation of the associated cross section derivatives. The calculation is based on the representation

$$\psi(x, \beta) = \frac{2}{\sqrt{\pi}} \int_0^{\infty} dy e^{-y^2} \frac{1+x^2+\beta^2 y^2}{(1-x^2+\beta^2 y^2)^2+4x^2} , \quad (78)$$

$$\phi(x, \beta) = \frac{2}{\sqrt{\pi}} \int_0^{\infty} dy e^{-y^2} \frac{1+x^2-\beta^2 y^2}{(1-x^2+\beta^2 y^2)^2+4x^2} , \quad (78')$$

where $y = (x' - x) / \beta$. The numerical integration is performed according to the trapezoidal rule with $\Delta y = 0.7$. Unless

$$\beta < \frac{\Delta y}{\pi} \quad \text{and} \quad \frac{x^2 - 1}{\beta^2} + \frac{2\pi}{\beta \Delta y} > 25, \quad (79)$$

which means unless Doppler broadening is very small ($\Delta \lesssim 0.1 \Gamma$) and the energy very far from the peak, the following correction terms are added:

$$\delta \psi = \frac{2\sqrt{\pi}}{\beta} \exp\left(-\frac{x^2 - 1}{\beta^2} - \eta\right) \frac{\cos \xi (e^{-\eta} - \cos \zeta) - \sin \xi \sin \zeta}{(e^{-\eta} - \cos \zeta)^2 + \sin^2 \zeta} \quad (80)$$

$$\delta \phi = -\frac{2\sqrt{\pi}}{\beta} \exp\left(-\frac{x^2 - 1}{\beta^2} - \eta\right) \frac{\sin \xi (e^{-\eta} - \cos \zeta) + \cos \xi \sin \zeta}{(e^{-\eta} - \cos \zeta)^2 + \sin^2 \zeta} \quad (80')$$

$$\text{with } \xi = \frac{2x}{\beta^2}, \quad \eta = \frac{2\pi}{\beta \Delta y}, \quad \zeta = \frac{2\pi x}{\beta \Delta y}. \quad (81)$$

TGAUSS

This subroutine yields values of a truncated Gaussian distribution, normalized to unity, for the resolution-broadening calculation. The truncation limits are at $\pm 2\sqrt{2}$ standard deviations from the center.

CHISQ4

This subroutine is included in the present description to show the use of a skew resolution function, in this case of a χ^2 -distribution with four degrees of freedom. It was actually applied to certain Fe-56 data, as shown by the call statement in subroutine CONE (the atomic weight, 56, was used as the criterion). With the information given in the comments it is easy to change the subroutine so as

to produce χ^2 -distributions with other degrees of freedom, i.e. different skewness.

RANDOM

This function calls a (pseudo-)random-number generator (RANDU for IBM installations, for instance) and yields random numbers between 0 and 1. If necessary, RANDU can be replaced by an equivalent random number generator.

3.2 Input description

All numeric input must be given as FORTRAN-readable floating-point numbers. The present version of the program uses card input. As a rule 10 card columns are reserved for each number (the exception being the last two items on "resonance cards", see below). Within this field the number can be placed arbitrarily in F format, for E format the usual rules apply. This structure permits complete omission of resonances or of data points from the input. One can easily change the number of resonances or data points by simply adding or removing input cards. Isotopic or spin reassignment of a given resonance requires nothing but repositioning of one "resonance card" in the input deck.

All energies must be given in keV, all channel radii in fm, all flight-time increments in ns, all lengths (flight paths, sample dimensions) in m. If other units are preferred it is easy to change the corresponding input and output formats by means of FORTRAN scale factors (along with the units in the printed table headings).

Input formats are essentially the same as for the FANAL code (Ref. 2) so that the same "isotope", "potential-scattering" and "resonance cards" (see below) can be used

for both codes. This has proved very convenient for parallel shape analysis of both capture and transmission data.

Each potential-scattering or resonance parameter is followed by an initial "uncertainty". If this uncertainty is set equal to zero the associated parameter is treated as constant, otherwise the parameter is adjusted in each iterative step.

The fact that the same input structure is used both in FANAL und FANAC is the reason that certain parameters such as distant-level strength functions or effective radii can be declared as adjustable although this is not practical in capture data analysis since capture yields depend on these quantities only weakly through self-shielding factors and multiple-collision yields.

The card input is structured as follows.

1st Card (Title Card)

Columns 1 - 80 Arbitrary alphameric text. This text appears on the print-out and also on the plots.

2nd Card (Analysis and Iteration Characteristics)

Columns 1 - 10 Lower boundary E_{\min} (keV),
" 11 - 20 upper boundary E_{\max} (keV) of the range of explicitly treated resonances, i.e. of the range for which the extracted parameters (including the distant-level strength function and the effective nuclear radius, Eqs. 15-17) are valid. Note that ΔE and \bar{E} in Eq. 15 are calculated as $\Delta E = E_{\max} - E_{\min}$, $\bar{E} = (E_{\max} + E_{\min})/2$.
" 21 - 30 Effective temperature T(K), cf. Eq. 26, taken as the same for all isotopes.
" 31 - 40 Largest relative variation $\delta\chi^2/\chi^2$ between successive iterations (cf. Eq. 3) which is considered as sufficient to terminate the iterative process and to declare convergence achieved (values of order 1 to 5 % were found to be reasonable).
" 41 - 50 Maximum number of iterations to be followed through.

3rd Card (Isotope Card)

The isotope which contributes most to multiple-collision capture should come first, as its atomic mass is used to determine all energy losses and lab angles.

Columns 1 - 10 Isotopic abundance a of the first sample nuclide (i.e. fraction of sample atoms belonging to this nuclide).

Columns	11 - 20	Atomic mass of first sample nuclide divided by neutron mass (it is usually sufficient to use simply the nucleon number A).
"	21 - 30	Nuclear spin quantum number I of the first nuclide.
"	31 - 40	Effective nuclear radius R'_1 (fm) for p-wave scattering of the first nuclide (see Eq. 22).

4th Card (Potential-Scattering Card)

Columns	1 - 10	Effective s-wave strength function $S_{I+1/2}$ of first nuclide. ($S_{I+1/2}$ and $R'_{I+1/2}$ determine the influence of distant levels on s-wave scattering, see Eqs. 15-17).
"	11 - 21	Initial uncertainty ^{+) of $S_{I+1/2}$.}
"	21 - 30	Effective radius $R'_{I+1/2}$ (fm) of first nuclide (see Eq. 16).
"	31 - 40	Initial uncertainty ^{+) of $R'_{I+1/2}$.}
"	41 - 50	Threshold $E_{I+1/2}$ (keV) for first inelastic channel, first isotope, compound spin I+1/2 (see text below Eq. 10).

5th Card (Resonance Card)

Columns	1 - 10	Resonance energy E_0 (keV), lab system,
"	11 - 20	Initial uncertainty ^{+) of E_0,}
"	21 - 30	Neutron width Γ_n (keV),
"	31 - 40	Initial uncertainty ^{+) of E_0,}
"	41 - 50	Partial width for inelastic scattering $\Gamma_{n'}$ (keV) with the sign of $\gamma_n \gamma_{n'}$ (see text after Eq. 10),
"	51 - 60	Initial uncertainty ^{+) of $\Gamma_{n'}$,}
"	61 - 70	Radiation width Γ_γ (keV).
"	71 - 75	Initial uncertainty ^{+) of Γ_γ.}

^{+) The preceding parameter is kept constant if the uncertainty is zero, otherwise it is adjusted.}

Columns 76 - 80 Detection efficiency ϵ (see Eq. 65 where this quantity is denoted by $\epsilon_{i\lambda}$), may include gamma ray absorption, gamma spectrum fraction above detector threshold, intrinsic detector efficiency etc.

As a rule the smallest partial width should be varied (if necessary), otherwise convergence may be bad.

A similar "resonance card" follows for each s-wave resonance with spin $I+1/2$ belonging to the first isotope. If this isotope has non-zero ground state spin, s-wave neutrons can excite also resonances with spin $I-1/2$. For this second possible compound spin a similar set of cards must be prepared, consisting of at least a "potential-scattering card" and optionally a number of "resonance cards". If, for a given isotope and compound spin, no "resonance cards" are encountered, only the potential-scattering cross section is calculated.

So-called p-, d-, ... wave levels, i.e. resonances which can only be excited by neutrons with $l \geq 1, l \geq 2, \dots$, must be treated as levels of additional (pseudo-) isotopes, with vanishing s-wave scattering ($S_J = R'_J = 0$). For each occurring spin factor g_J there must be one pseudo-isotope, with I (ground state spin quantum number) set equal to zero and the abundance set equal to $g_J a$ (a : true abundance). In Sect. 4 below this is illustrated for p- and d-wave levels of $^{56}\text{Fe} + n$ ($g_J = 1, 2$ and 3). For most levels with $l \geq 1$, however, spins are unknown and only the combination $g_J \Gamma_n \Gamma_\gamma / \Gamma$ can be estimated from the capture data. One must then use average values for g_J and Γ_γ to calculate the corresponding Γ_n and use these figures as input (see the example in Sect. 4). This is quite adequate if self-shielding and (multiple) scattering are unimportant ($n\sigma_T \ll 1$) or dominated by s-wave scattering ($\sigma_{T1} \ll \sigma_{T0}$, Eq. 27). In this case the capture yields depend

mainly on $g_J \Gamma_n \Gamma_\gamma / \Gamma$, i.e. on the "unshielded" resonance area. For thick samples and relatively small s-wave cross sections, however, it can make a difference whether a resonance with given peak area is treated as narrow and high (large g_J) or as broad and low (small g_J).

If the capture sample contains other isotopes (sample impurities, oxygen in oxides, hydrogen in adsorbed moisture etc.) completely analogous input cards must be prepared for each nuclide: an "isotope card" followed by a "potential-scattering card" plus "resonance cards" (optional) for $I+1/2$, then for $I-1/2$ if $I>0$. Up to ten isotopes and pseudo-isotopes can be handled.

The cards specifying the cross section parameters are followed by other cards carrying information for up to five time-of-flight runs. For each run one needs a "sample card", a "time-of-flight card", "capture yield data cards" and one blank card:

Sample Card

Columns	1 - 10	Sample thickness n (atoms/b), i.e. total number of sample atoms per barn.
	11 - 20	Sample radius r (atoms/b), see Eq. 55.

Time-of-flight Card

Columns	1 - 10	Flight path L (m), to mid-plane of sample.
"	11 - 20	Sample thickness ΔL (m), see Eq. 67.
"	21 - 30	Full width at half maximum of gamma peak, Δt (ns), see Eq. 67.
"	31 - 40	maximum number N_0 of Monte Carlo multiple-collision event simulations per energy point, see Eq. 73.
"	41 - 50	Time shift (ns) to correct for small time-zero channel differences between time-of-flight runs (can usually be left blank).

Capture Yield Card

Columns	1 - 10	Energy E_1 (keV)
	11 - 20	Capture yield divided by sample "thickness", η_1/n (b).
	21 - 30	Uncertainty $\delta \eta_1/n$ (b).
	31 - 40	Energy E_{i+1} (keV).
	41 - 50	Capture yield divided by sample "thickness", η_{i+1}/n (b).
	51 - 60	Uncertainty $\delta \eta_{i+1}/n$ (b).

Thus two data points including their uncertainties can be put on each "capture yield card". Up to 512 experimental points (total for all TOF runs) can be simultaneously fitted. It is to be noted that energies are given explicitly, i.e. regularity or equidistance on a time-of-flight scale is not required (in contrast to the FANAL code, Ref. 2). Pairs of data points not to be used for the fit can thus be taken out quite easily.

The last capture yield card from a given time-of-flight run (which may contain a single data point) must be followed by a blank card. Two blank cards signal the end of the whole input.

Input for other calculations may follow, i.e. problems can be stacked.

In addition to the card input there are a few quantities which are specified by DATA or other statements in the program:

- the adjustment reduction factor FUDGE in subroutine ADJ (DATA statement),
- the ratio of beam radius to sample radius, EDGE, in subroutine MUY (DATA statement),
- the minimum number of first collisions per grid energy in subroutine MUY (card no. 11830).
- The constants A,B,C in subroutine CHISQ4 refer to 4 degrees of freedom as explained on comment cards in CHISQ4.

3.3 Output description

The output consists of print-out and plots.

The print-out shows first the contents of "title", "isotope", "potential-scattering", "resonance" and "time-of-flight" cards and some additional derived information such as compound spins. This is followed by tables of measured and calculated capture cross section data (observed yields divided by sample "thickness" y/n and capture cross sections σ_{γ}). These quantities are calculated for the sample composition and include detector efficiencies. Subsequently the values of the squared-error sum χ^2 and of Gauss' error adjustment factor $\sqrt{\chi^2/(N-M)}$ are printed (N: number of data prints, M: number of adjusted parameters). For a good fit the error adjustment factor should be close to 1.

After that one gets a table with the adjusted and constant parameters and their uncertainties. The latter are the square roots of the corresponding diagonal elements of the covariance matrix, multiplied by the error adjustment factor. The covariance matrix elements result from the experimental uncertainties by normal error propagation (in linear approximation), the error adjustment factor characterizes the goodness of the fit.

If the maximum number of iterative steps exceeds 1 a similar printout (yield table plus improved parameters) is obtained for each completed step. If the number is set equal to 1 the cross section of the pure first isotope is also calculated. In this way the code can be used to generate and plot Doppler-broadened cross sections corrected for all experimental effects such as sample impurities, instrumental resolution, detector efficiency, self-shielding, and multiple scattering.

For each iterative step a plot is produced which contains all experimental points and their error bars together with the calculated values in curve form. The text of the title card appears under each plot. If only one iteration is requested a second plot is obtained showing the Doppler-broadened capture cross section of the pure first isotope.

For the last set of parameters neither capture yield nor χ^2 values are calculated, printed or plotted. If the convergence criterion, Eq. 3, with reasonably chosen ϵ is satisfied there should be no significant change with respect to the preceding step.

4. Example

Fig. 4 shows the input cards for a realistic fitting problem which illustrates most features of FANAC:

Data from two measurements, taken with the same Fe_2O_3 sample enriched to 99.7 % ^{56}Fe but with slightly different resolution, were to be analyzed between 21.6 and 41.1 keV (for experimental details see Ref. 3, 14). The range of explicitly treated ^{56}Fe -resonances with $l=0$ (including a subthreshold level at -3.9 keV) was taken as -6 to 100 keV, the effective temperature as 300 K and the convergence threshold (ϵ in Eq. 3) as 3 %.

The cards for s-wave levels are followed by those for narrow resonances with $l \geq 1$ which are represented as levels of three fictitious isotopes with target spin zero and abundance of $0.997 = 99.7\%$ (for $g_J=1$), $2 \cdot 0.997 = 1.994$ (for $g_J=2$) and $3 \cdot 0.997 = 2.991$ (for $g_J=3$), as explained in the input description (Sect. 3.2). Potential scattering for $l=1$ is represented by an effective nuclear radius of 5.4 fm of the first fictitious isotope.

The only important impurity was oxygen. Its smooth cross section is specified by a potential-scattering card without resonance cards following. Its relative abundance is 1.5 corresponding to the stoichiometric ratio in Fe_2O_3 . This means that the sample thickness must be given in Fe-atoms per barn. Alternatively one could have given the total number of all (iron and oxygen) nuclei per barn in the sample along with relative abundances of $(2/5) \cdot 0.997 \cdot g_J$ for ^{56}Fe and $3/5$ for oxygen.

Sample, run and capture yield cards for the utilized two time-of-flight runs follow.

Fig. 5-7 show the cross section parameter input and the tables with measured and calculated yields and cross sections (inclusive of detector efficiencies) as they are printed by the computer. Fig. 8 shows the corresponding plots. Convergence, i.e. less than 5 % variation of χ^2 , was formally reached after four iterations. The final value $\sqrt{\chi^2/(M-N)} = 1.288$ is close enough to unity to indicate that input data, input errors and mathematical model (cross section formalism, multiple-collision treatment, resolution function) are reasonably consistent with each other. A slightly better χ^2 might be achieved with a more exact resolution function, but no significant change in resonance parameters is expected.

A few words on input preparation may be appropriate.

First estimates of the resonance energies were obtained directly from the capture data by inspection. The starting value for the radiation width of the broad s-wave resonance near 27.6 keV was taken from the "barn book" (Ref. 15). The peak area parameters $g_J \Gamma_n \Gamma_\gamma / \Gamma$ of the narrow levels were estimated from the data by equating the observed peak areas to the theoretical thin-sample expression,

$$\delta E \sum_i \frac{\eta_i}{n} = \frac{a\epsilon}{n} 2\pi\lambda^2 g_J \frac{\Gamma_n \Gamma_\gamma}{\Gamma} \quad (82)$$

where the sum extends over the peak, δE is the spacing of the data points near the peak, χ^2 is calculated at resonance and the other symbols have the same meaning as before (a : abundance, ϵ : detection efficiency for the resonance). From a Bayesian analysis of available values of $g_J \Gamma_n$, $g_J \Gamma_\gamma$ and $g_J \Gamma_n \Gamma_\gamma / \Gamma$ for $^{56}\text{Fe}+n$ the radiation widths for p- and d- wave resonances were known to be close to 0.33 and 0.6 eV, respectively. The same type of analysis was used to select spins and parities by Monte Carlo sampling. The resulting g_J and Γ_γ were then used to calculate the input values of Γ_n from the estimated $g_J \Gamma_n \Gamma_\gamma / \Gamma$. For each observed resonance the smaller of the two partial widths was adjusted. The efficiency factors assigned to individual resonances were estimated from observed capture gamma ray spectra. They also include an estimated correction for quasi-prompt capture of resonance-scattered neutrons in the immediate vicinity of the sample (cf. Refs. 3,10).

The print-out shows that the temperature was raised artificially to 6368 K in order to avoid difficulties with the equidistant internal energy grid as explained in Sect. 3.1 (subroutine PRY). One can treat the narrow resonances separately to avoid or at least reduce this automatic temperature change. Taking for example only the data between 22.1 and 23.2 keV, i.e. analyzing only the peak near 22.75 keV, one gets a much smaller change to only 1171 K. Nevertheless, the resulting $g_J \Gamma_n \Gamma_\gamma / \Gamma$ is essentially unchanged.

A similar insensitivity exists with respect to the correct spin factor: if all spin factors are set equal to 1 (i.e. all narrow levels are assigned to the same fictitious isotope with $a = 0.997$) the results are essentially unchanged. The reason is the same as before: In our example s-wave scattering of ^{56}Fe and ^{16}O predominates even near the peaks of the narrow levels. Therefore changes in g_J or in the temperature, which otherwise would affect the peak area through the beam-attenuation

factor $(1 - \exp(-n\sigma_T)) / (n\sigma_T)$, have no great influence. In other cases, however, where s-wave scattering does not predominate, spin and (unless $\Gamma \ll \Delta$) temperature effects may be important.

Multiple-collision contributions in this example are seen to be of the order of few percent near the peaks of the narrow resonances but quite substantial near the broad peak near 27.6 keV and in the valleys between resonances.

The CPU time for the whole job (3 iterations, 12 adjusted parameters, a maximum of 30 Monte Carlo multiple-collision event simulations per grid point) was 3 min 19 sec on an IBM/370-168 computer. The memory space required was 476 k bytes.

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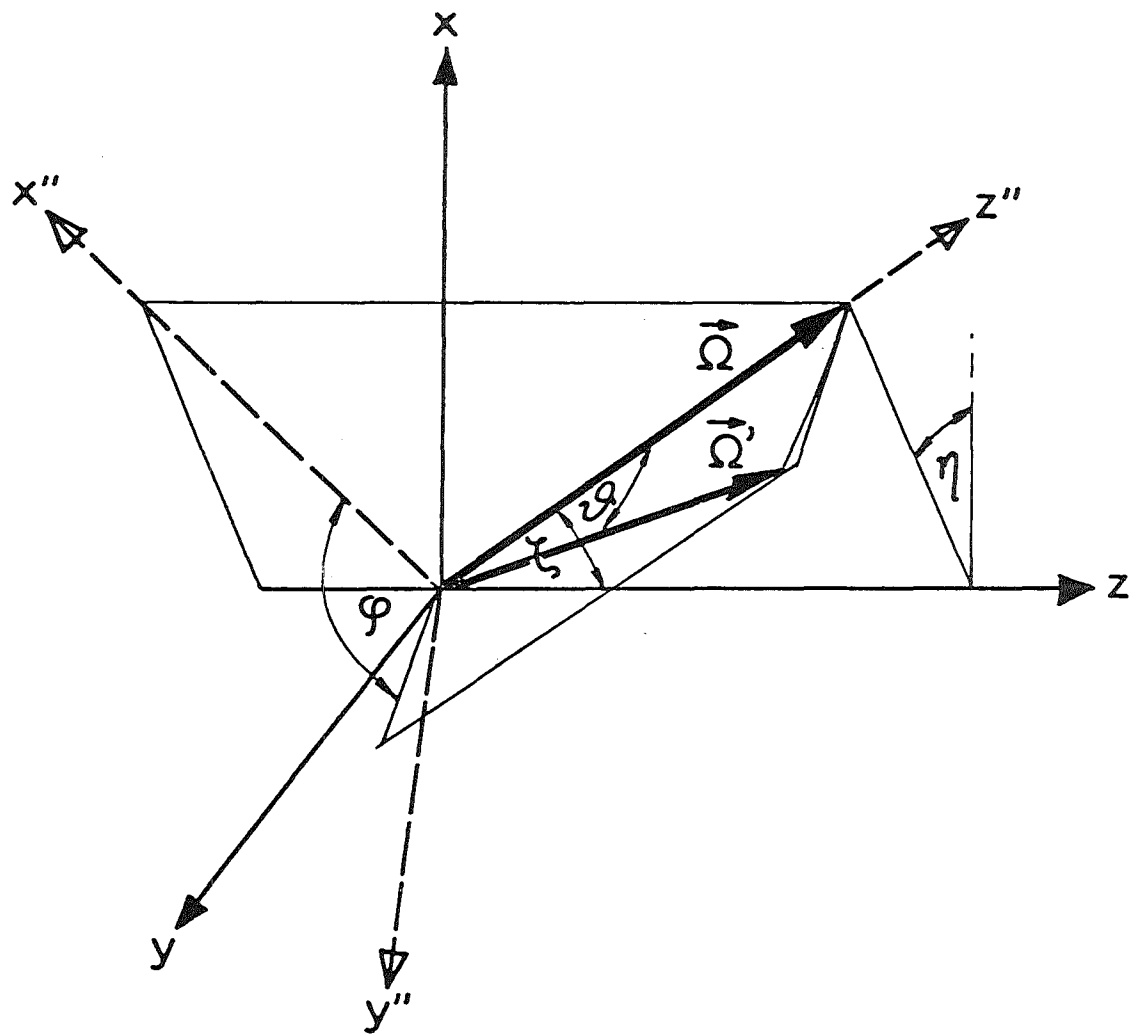


Figure 1. Neutron directions $\vec{\Omega}$, $\vec{\Omega}'$ (before and after collision, resp.) and angles ζ , η , ϑ , φ , and their positions relative to the laboratory system (x, y, z) and to the system $S'' (x'', y'', z'')$, in which $\vec{\Omega}$ has the coordinates $(0, 0, 1)$.

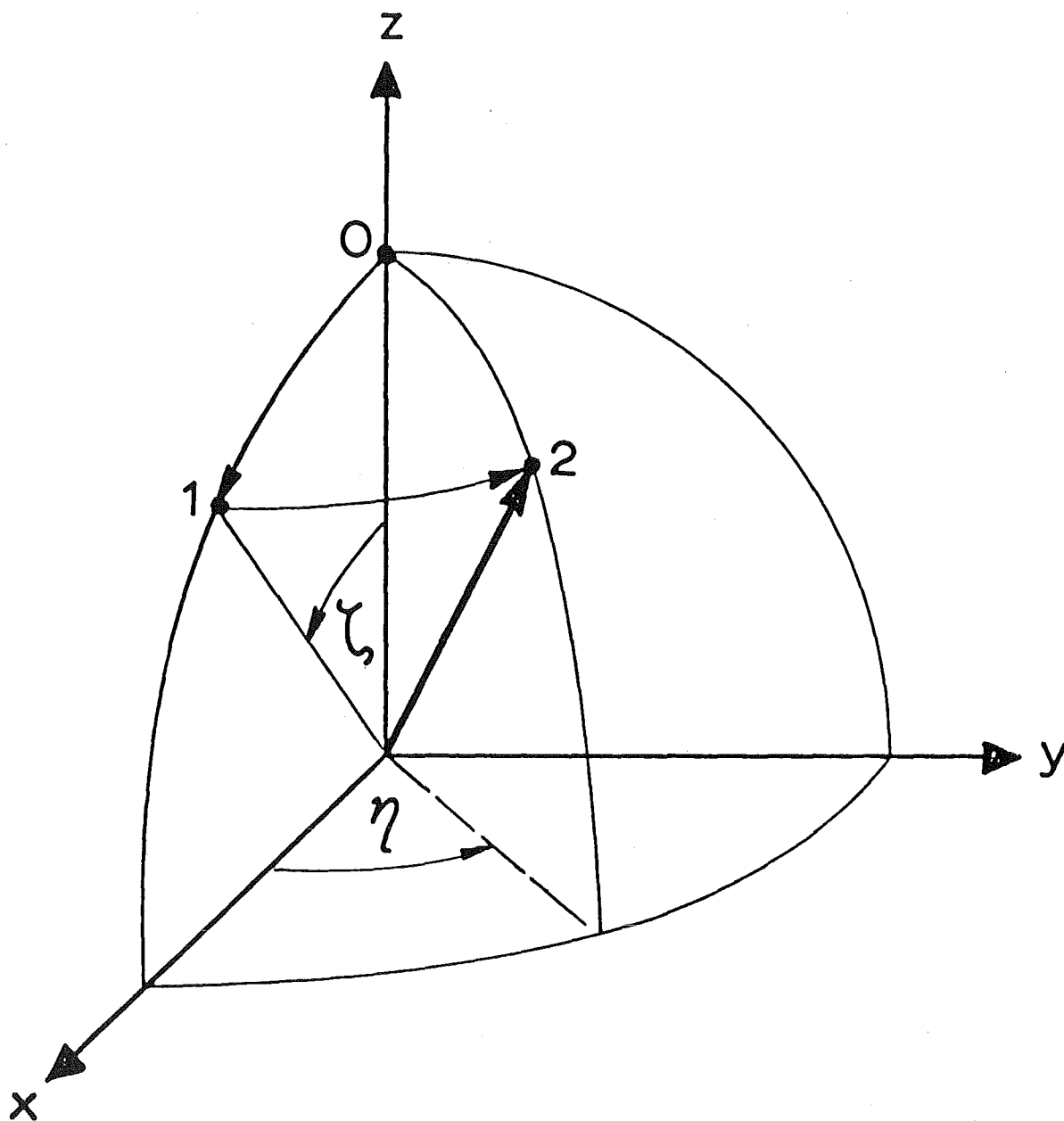


Figure 2. Two-step rotation of a vector from the position $(0, 0, 1)$ into the position $(\sin\zeta \cos\eta, \sin\zeta \sin\eta, \cos\zeta)$.

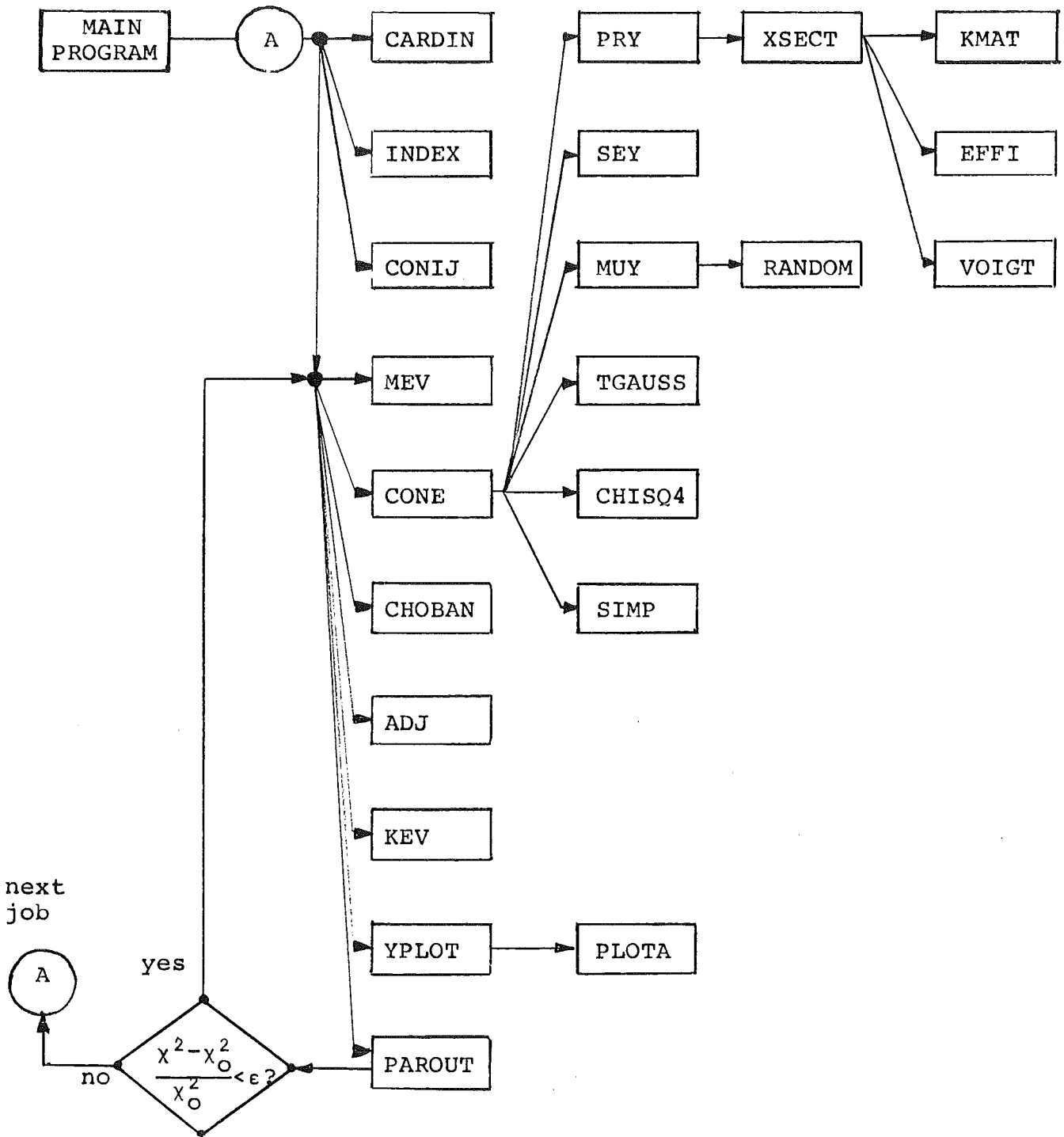


Fig. 3 FANAC program structure

FE-56(N,GAMMA), 0.009924 FE-AT./B, FANAC TEST PROBLEM, 11.5.78

EXPLICIT-RESONANCE RANGE: -100.000 ... 100.000 KEV
 EFFECTIVE TEMPERATURE: 300.0 KELVIN
 MAXIMUM TOLERABLE CHI**2 CHANGE: 0.0500 (RELATIVE)
 MAX. NUMBER OF ITERATIVE STEPS: 5.

SPECIFICATIONS OF UTILIZED MEASUREMENTS:

=====

MEAS. NO.	SAMPLE THICKNESS (NUCLEI/B)	SAMPLE RADIUS (NUCLEI/B)	FLIGHT PATH (M)	SAMPLE THICKNESS (M)	CHANNEL WIDTH (NS)	FWHM GAMMA PEAK (NS)	MONTE CARLO HISTORIES	TIME SHIFT CORRECTION (NS)
1	9.924E-03	1.200E-01	2.0500	0.0070	2.000	7.000	30.	0.0
2	9.924E-03	1.200E-01	2.0500	0.0070	2.000	5.500	30.	0.0

CROSS SECTION PARAMETERS:

=====

ABUN-DANCE	ATOMIC HEIGHT	TARGET SP IN	P-WAVE RADIUS (FM)	COMP. SPIN	S-WAVE STRENGTH FUNCTION /UNCERT.	S-WAVE RADIUS (FM) /UNCERT.	INEL. THRESH. (KEV)	RESONANCE ENERGY (KEV) /UNCERT.	PARTIAL WIDTHS FOR EL. SCATT. (KEV) /UNCERT.	INEL. SCATT. (KEV) /UNCERT.	CAPTURE (KEV) /UNCERT.	DETECT. EFFIC.
0.9970	56.0	0.0	0.0	0.5	2.500E-04 0.0	6.100 0.0	0.0	-3.900 0.0	5.200E-01 0.0	0.0 0.0	1.300E-03 0.0	0.970
								27.700 1.000	1.400E+00 0.0	0.0 0.0	1.450E-03 0.100E+01	0.970
								73.980 0.0	5.400E-01 0.0	0.0 0.0	6.300E-04 0.0	0.930
								83.650 0.0	1.300E+00 0.0	0.0 0.0	5.000E-04 0.0	0.970
0.9970	56.0	0.0	5.400	0.5	0.0 0.0	0.0	0.0	36.600 1.000	3.000E-03 0.0	0.0 0.0	3.300E-04 0.100E+01	0.830
1.9940	56.0	0.0	0.0	0.5	0.0 0.0	0.0 0.0	0.0	34.100 1.000	1.000E-03 0.0	0.0 0.0	3.300E-04 0.100E+01	0.870
								36.300 1.000	1.300E-04 0.100E+01	0.0 0.0	6.000E-04 0.0	0.850
2.99.0	56.0	0.0	0.0	0.5	0.0 0.0	0.0 0.0	0.0	22.800 1.000	7.000E-05 0.100E+01	0.0 0.0	6.000E-04 0.0	0.870
								24.500 1.000	1.000E-05 0.100E+01	0.0 0.0	6.000E-04 0.0	0.870
1.5000	.6.0	0.0	5.300	0.5	0.0 0.0	5.3 0.0	6052.000					

EFFECTIVE TEMPERATURE : 8249. DEG. KELVIN,
 RESULTING DOPPLER WIDTH: 31.133 EV AT 19.092 KEV

Fig. 5 - Beginning of FANAC print-out: content of input cards excepting capture yields. The last two lines are a warning that the effective temperature was raised artificially in order to reduce the number of internal grid points (see text).

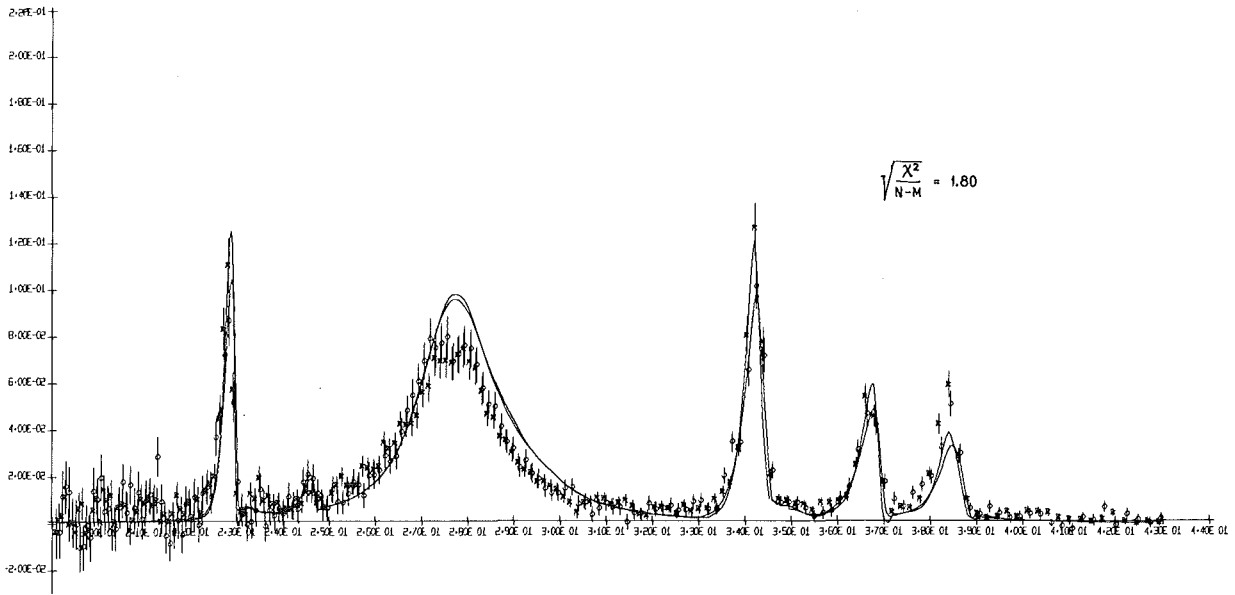
CROSS SECTION PARAMETERS:

=====

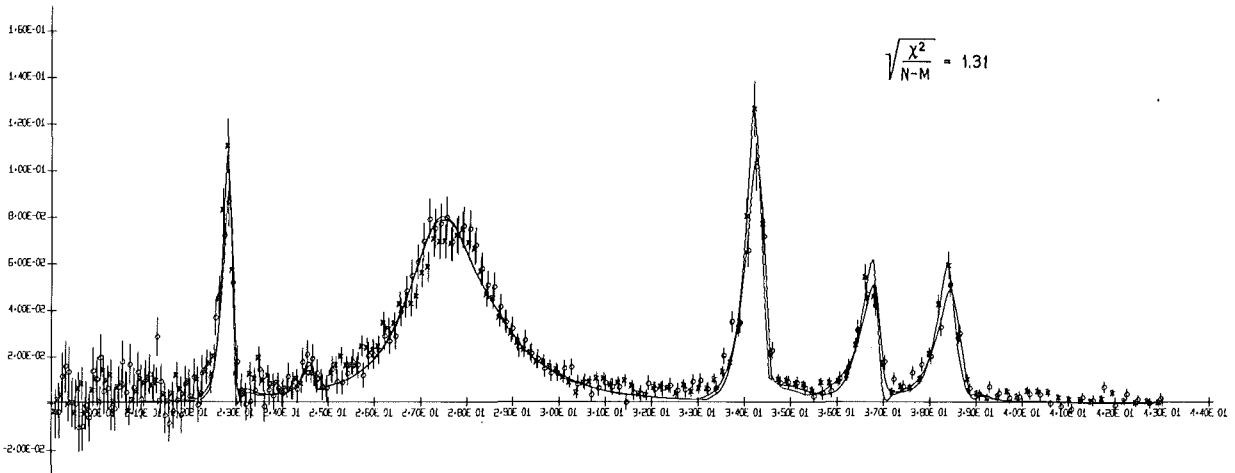
ABUN- DANCE	ATOMIC WEIGHT	TARGET SPIN	P-WAVE RADIUS (FM)	COMP. SPIN	S-WAVE STRENGTH FUNCTION /UNCERT.	S-WAVE RADIUS (FM) /UNCERT.	INEL. THRESH. (KEV)	RESONANCE ENERGY (KEV) /UNCERT.	PARTIAL WIDTHS FOR			DETECT. EFFIC.
									EL. SCATT. (KEV) /UNCERT.	INEL. SCATT. (KEV) /UNCERT.	CAPTURE (KEV) /UNCERT.	
0.9970	56.0	0.0	0.0	0.5	2.500E-04 0.0	6.100 0.0	0.0	-3.900 0.0	5.200E-01 0.0	0.0 0.0	1.300E-03 0.0	0.970
								27.478 0.038	1.400E+00 0.0	0.0 0.0	1.224E-03 0.414E-04	0.970
								73.980 0.0	5.400E-01 0.0	0.0 0.0	6.300E-04 0.0	0.930
								83.650 0.0	1.300E+00 0.0	0.0 0.0	5.000E-04 0.0	0.970
0.9970	56.0	0.0	5.400	0.5	0.0 0.0	0.0 0.0	0.0	36.641 0.008	3.000E-03 0.0	0.0 0.0	3.366E-04 0.222E-04	0.830
1.9940	56.0	0.0	0.0	0.5	0.0 0.0	0.0 0.0	0.0	34.118 0.006	1.000E-03 0.0	0.0 0.0	3.795E-04 0.238E-04	0.870
								38.278 0.007	2.349E-04 0.185E-04	0.0 0.0	6.000E-04 0.0	0.850
2.9910	56.0	0.0	0.0	0.5	0.0 0.0	0.0 0.0	0.0	22.754 0.005	6.705E-05 0.461E-05	0.0 0.0	6.000E-04 0.0	0.870
								24.488 0.037	7.365E-06 0.207E-05	0.0 0.0	6.000E-04 0.0	0.870
1.5000	16.0	0.0	5.300	0.5	0.0 0.0	5.3 0.0	6051.988					

CONVERGENCE AFTER 3 ITERATION(S).

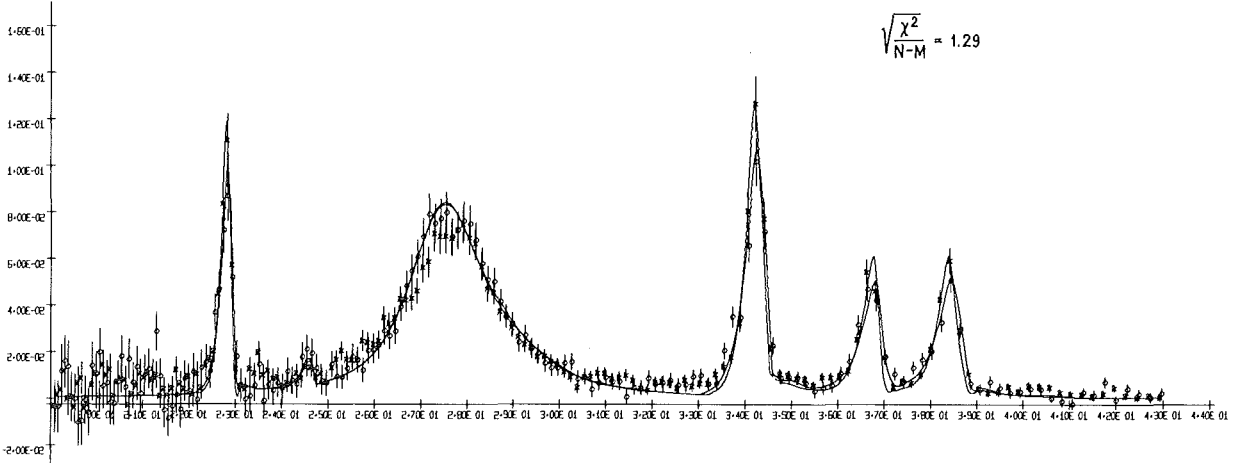
Fig. 7 - Print-out of final parameters after convergence has been declared. It should be noticed that for the narrow resonances (with $l > 1$, members of level sequences with zero s-wave radius) only the product $g\Gamma_n/\Gamma_\gamma$ (with g absorbed in the abundance as far as input is concerned) is meaningful, not Γ_n or Γ_γ separately.



HEB-00001 FE-56(N-GAMMA), 0.003824 FE-AT-1/3, FANUC TEST PROBLEM, 11-S-



HEB-00002 FE-56(N-GAMMA), 0.003824 FE-AT-1/3, FANUC TEST PROBLEM, 11-S-



HEB-00003 FE-56(N-GAMMA), 0.003824 FE-AT-1/3, FANUC TEST PROBLEM, 11-S-

Fig. 8 - Plots produced along with the print-out shown in Figs. 5-7 illustrating progress of fit. Top corresponds to input parameters. The two solid lines correspond to the two experimental runs with slightly different resolution.

APPENDIX

Listing of FANAC code

```

C      FANAC, MAIN PROGRAM                                000010
C                                                        000020
C      NOTE TO NON-KFK USERS:                            000030
C                                                        000040
C      PLOTA(X,Y,N,NT,NP,NH,I,NS,NR,XMAX,XMIN,SY,YMAX,YMIN,SY,TEXT,IB) 000050
C      IS A STANDARD PLOTTER SUBROUTINE IN USE AT KFK (KARLSRUHE) 000060
C      WHICH MUST BE REPLACED BY AN EQUIVALENT PLOTTER PACKAGE 000070
C      ELSEWHERE THE ARGUMENTS ARE EXPLAINED IN SUBROUTINE YPLOT. 000080
C                                                        000090
C      RANDU(IX,IY,RANDOM) EMPLOYED IN SUBROUTINE MUY      000100
C      IS A RANDOM-NUMBER GENERATOR (IBM SUBROUTINE) YIELDING 000110
C      VARIATE VALUES UNIFORMLY DISTRIBUTED IN THE INTERVAL 0 ... 1 000120
C      AND MUST BE REPLACED BY AN EQUIVALENT IF UNAVAILABLE. 000130
C                                                        000140
C      FANAC SUMMARY:                                     000150
C                                                        000160
C      VERSION : NOVEMBER 1977                            000170
C      PROGRAMMING LANGUAGE: FORTRAN IV                   000180
C      PURPOSE  : SHAPE ANALYSIS OF NEUTRON CAPTURE DATA. 000190
C                EXTRACTION OF RESONANCE PARAMETERS.     000200
C                CALCULATION OF TRUE CROSS SECTION.       000210
C      METHOD    : SIMULTANEOUS LEAST-SQUARES FIT TO      000220
C                SEVERAL SETS OF TIME-OF-FLIGHT DATA    000230
C                (TAKEN E.G. WITH DIFFERENT SAMPLES OR   000240
C                FLIGHT PATHS).                           000250
C      FORMALISM : MULTI-LEVEL R-MATRIX FORMULA WITH     000260
C                1 ELASTIC AND 1 INELASTIC NEUTRON      000270
C                CHANNEL PER COMPOUND SPIN AND PARITY.   000280
C                CAPTURE CHANNELS ARE ELIMINATED WITH    000290
C                TEICHMANN-WIGNER REDUCTION METHOD        000300
C                FOLLOWING REICH AND MOORE.               000310
C      CORRECTIONS : DOPPLER BROADENING IS APPLIED ONLY TO 000320
C                LEVEL SEQUENCES WITH ZERO S-WAVE RADIUS 000330
C                (P-,D-...WAVE LEVELS) AS FANAC WAS      000340
C                DEVELOPPED FOR STRUCTURAL MATERIALS     000350
C                FOR WHICH DOPPLER BROADENING OF S-WAVE  000360
C                LEVELS IS NEGLIGIBLE.                   000370
C                YIELDS INCLUDE SELF-SHIELDING AND      000380
C                MULTIPLE-SCATTERING CORRECTIONS.       000390
C                RESOLUTION BROADENING WITH GAUSSIAN    000400
C                OR CHI-SQUARE DISTRIBUTION IS FOLDED IN. 000410
C                AN EFFICIENCY FACTOR CAN BE APPLIED TO 000420
C                EACH RESONANCE TO TAKE ACCOUNT OF GAMMA 000430
C                SPECTRUM FLUCTUATIONS.                 000440
C      DOCUMENTATION : F. H. FROEHNER, REPORT KFK 2145 (1973) 000450
C                                                        000460
C      SUBSCRIPT CONVENTIONS IN FANAC:                   000470
C      I (MAX. 10) LABELS ISOTOPES,                      000480
C      J (MAX. 2)  "  COMPOUND SPINS,                    000490
C      K (MAX. 2048) "  MEASURED DATA POINTS,           000500
C      L (MAX. 48)  "  RESONANCES,                      000510
C      M (MAX. 200) "  CROSS SECTION PARAMETERS          000520
C                  (BOTH FIXED AND ADJUSTED),          000530
C      N (MAX. 5)  "  TIME-OF-FLIGHT MEASUREMENTS.      000540
C                                                        000550
C      NOTE: FOR EACH SAMPLE NUCLIDE ONE MUST PREPARE INPUT FOR 000560
C      1 ISOTOPE WITH NONZERO S-WAVE RADIUS FOR ALL S-WAVE LEVELS, 000570
C      000580

```

```

C      AND, IF P-,D-...WAVE LEVELS ARE PRESENT,          000590
C      1 ISOTOPE WITH ZERO S-WAVE RADIUS FOR NON-S-WAVE LEVELS. 000600
C      THE LATTER ARE TREATED WITHOUT MULTI-LEVEL OR POTENTIAL/ 000610
C      RESONANCE INTERFERENCE.                             000620
C      THIS APPROXIMATION MAY FAIL ABOVE 300-400 KEV.       000630
C                                                         000640
0001      COMMON                                           000650
          HI,GI,Z11,Z21,Z22,IZ,CHISQ,CHISQ0,              000660
          EI,IX,J,JX(11),KX(5),L,LX(11,2),M,MA,MX,N,NX,KKX,MP(11,2),MR(11,2), 000670
          BTITLE(20),ZIT,EPS,E1,E2,TEFF,H(11),AG(11),SPIN(11),RF(11), 000680
          4S(11,2),CS(11,2),X(200),DLX(200),ES(11,2),EFF(11,2,200), 000690
          5FP(5),DLFP(5),TC(5),TB(5),ZH(5),EE(5,2),Y(512),DLY(5,2), 000700
          6XN(6),XR(6),TMX,AL(10),D(10),CT(10,2),CG(10,2),M1,M2,M0, 000710
          7BLE,FMN,ST(2048),DLST(2048),DST(20),SG(2048),DLSG(2048),DSG(20), 000720
          8SC(2048),T(2048),YY(2048),DYY(2048,20),ZZ(512),Z(512),DZ(20),EFC, 000730
          9F(201),RF(201)                                  000740
0002      COMMON /ABC/ A(20,20),B(20,20),C(20)           000750
0003      DOUBLE PRECISION A,B                             000760
0004      COMMON/II/ SG11(2048),DLSG11(2048)              000770
0005      COMPLEX HI,GI,Z11,Z21,Z22,DZ11,DZ21,DZ22       000780
C                                                         000790
0006      1 CHISQ=0.                                       000800
0007      CALL CARDIN                                       000810
0008      CALL INDEX                                       000820
0009      CALL CONTJ                                       000830
0010      IZX=ZIT                                          000840
0011      DO 2 IZ=1,IZX                                     000850
0012      CHISQ0=CHISQ                                     000860
0013      CHISQ=0.                                         000870
0014      CALL MEV                                         000880
0015      CALL CONE                                         000890
0016      CALL CHORAN(MA,A,B)                               000900
0017      CALL ADJ(IZ,MX,MA,KX(NX),CHISQ,X,DLX,B,C)       000910
0018      CALL KEV                                         000920
0019      CALL YPLOT                                        000930
0020      CALL PAROUT                                       000940
0021      VCHISQ=((CHISQ-CHISQ0)/CHISQ)**2                000950
0022      DLCHSQ=SQRT(VCHISQ)                              000960
0023      IF(DLCHSQ.LT.EPS)GO TO 3                          000970
0024      WRITE(6,100)IZ                                    000980
0025      100 FORMAT(//' AFTER',I3,' ITERATION(S) NO CONVERGENCE YET.'//) 000990
0026      2 CONTINUE                                       001000
0027      GO TO 1                                           001010
0028      3 WRITE(6,101)IZ                                   001020
0029      101 FORMAT(//' CONVERGENCE AFTER',I3,' ITERATION(S).'//) 001030
0030      GO TO 1                                           001040
0031      END                                              001050

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0001      C          SUBROUTINE CARDIN                      C01080
          C          C                                C01070
          C          C                                C01080
          C          C                                C01090
          C          C          CARDIN READS THE INPUT FROM CARDS          C01100
0002      C          COMMON                                C01110
          C          C                                C01120
          C          C          1HI,GI,Z11,Z21,Z22,IZ,CHISQ,CHISQ0,          C01130
          C          C          2I,IX,J,JX(11),KX(5),L,LX(11,2),M,MA,MX,N,NX,KKX,MP(11,2),MR(11,2),          C01140
          C          C          3TITLE(20),ZIT,EPS,E1,E2,TEFF,H(11),AG(11),SPIN(11),RP(11),          C01150
          C          C          4G(11,2),CS(11,2),X(200),DLX(200),ES(11,2),EFF(11,2,200),          C01160
          C          C          5FP(5),DLFP(5),TC(5),TB(5),ZH(5),EE(512),Y(512),DLY(512),          C01170
          C          C          6XN(6),XR(6),TMX,AL(10),D(10),CT(10,2),CS(10,2),M1,M2,MD,          C01180
          C          C          7DLE,EMN,ST(2048),DLST(2048),DST(20),SG(2048),DLSG(2048),DSG(20),          C01190
          C          C          8SC(2048),T(2048),YY(2048),DYY(2048,20),ZZ(512),Z(512),DZ(20),EFC,          C01200
          C          C          9F(201),RF(201)                                C01210
0003      C          DIMENSION TS(5)                        C01220
0004      C          COMPLEX HI,GI,Z11,Z21,Z22,DZ11,DZ21,DZ22          C01230
0005      C          I=0                                    C01240
0006      C          N=0                                    C01250
0007      C          MX=0                                    C01260
0008      C          HI=(0.,.5)                            C01270
0009      C          GI=(0.,1.)                            C01280
0010      C          1 READ(5,100,END=999)TITLE            C01290
0011      C          100 FORMAT(20A4)                      C01300
0012      C          WRITE(6,101)TITLE                     C01310
0013      C          101 FORMAT(1H1//30X,20A4//)           C01320
0014      C          READ(5,102)E1,E2,TEFF,EPS,ZIT         C01330
0015      C          102 FORMAT(5E10.5)                   C01340
0016      C          WRITE(6,103)E1,E2,TEFF,EPS,ZIT       C01350
0017      C          103 FORMAT(//                          C01360
          C          1' EXPLICIT-RESONANCE RANGE: ',F8.3,' ...',F8.3,' KEV'/          C01370
          C          2' EFFECTIVE TEMPERATURE: ',F8.1,' KELVIN'/          C01380
          C          3' MAXIMUM TOLERABLE CHI**2 CHANGE: ',F8.4,' (RELATIVE)'/          C01390
          C          4' MAX. NUMBER OF ITERATIVE STEPS: ',F8.0)          C01400
0018      C          2 I=I+1                                C01410
          C          READ ISOTOPE CARD (I-TH ISOTOPE)      C01420
          C          H(I) : ABUNDANCE,                    C01430
          C          AG(I) : ATOMIC WEIGHT,                C01440
          C          SPIN(I): TARGET SPIN,                C01450
          C          RP(I) : EFFECTIVE RADIUS FOR P-WAVE SCATTERING (FM). C01460
0019      C          READ(5,104)H(I),AG(I),SPTN(I),RP(I)   C01470
0020      C          104 FORMAT(4E10.5)                    C01480
0021      C          3 G(I,1)=.5*(1.+1./(2.*SPIN(I)+1.))   C01490
0022      C          G(I,2)=1.-G(I,1)                     C01500
0023      C          IF(SPIN(I).EQ.0.)JX(I)=1              C01510
0024      C          IF(SPIN(I).GT.0.)JX(I)=2              C01520
0025      C          CS(I,1)=SPIN(I)+.5                   C01530
0026      C          CS(I,2)=SPIN(I)-.5                   C01540
0027      C          L=0                                    C01550
0028      C          J=1                                    C01560
0029      C          MN=MX+1                                C01570
0030      C          MX=MN+1                                C01580
          C          READ POTENTIAL-SCATTERING CARD (I-TH ISOTOPE, J-TH SPIN) C01590
          C          X(MN) : S-WAVE STRENGTH FUNCTION,     C01600
          C          DLX(MN): INITIAL UNCERTAINTY OF X(MN), C01610
          C          X(MX) : EFFECTIVE RADIUS FOR S-WAVE SCATTERING (FM), C01620

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C          DLX(MX): INITIAL UNCERTAINTY OF X(MX),          001630
C          ES(I,J): 1ST INELASTIC THRESHOLD (KEV).        001640
C          NOTE: PUT INITIAL UNCERTAINTY =0. FOR FIXED,    001650
C                   >0. FOR ADJUSTED PARAMETERS.        001660
0031      READ(5,105)(X(M),DLX(M),M=MN,MX),ES(I,J)        001670
0032      105 FORMAT(5E10.5)                              001680
0033      4 L=L+1                                          001690
0034      MN=MX+1                                         001700
0035      MX=MN+3                                         001710
C          READ RESONANCE CARD                            001720
C          X(MN) : RESONANCE ENERGY (KEV),              001730
C          X(MN+1): NEUTRON WIDTH (KEV),                 001740
C          X(MN+2): PARTIAL WIDTH FOR INELASTIC SCATTER (KEV) 001750
C          X(MX) : RADIATION WIDTH (KEV)                 001760
C          DLX(M) : INITIAL UNCERTAINTY OF X(M)          001770
C          EFF(I,J,L): CAPTURE DETECTION EFFICIENCY      001780
0036      READ(5,106)(X(M),DLX(M),M=MN,MX),EFF(I,J,L)    001790
0037      106 FORMAT(7E10.5,2E5.3)                       001800
C          CHECK CARD TYPE                               001810
0038      IF(X(MX).NE.0.) GO TO 4                         001820
0039      IF(J.LT.JX(I)) GO TO 5                          001830
0040      IF(DLX(MN).GE.1.) GO TO 6                       001840
0041      GO TO 7                                         001850
C          LAST CARD WAS POTENTIAL-SCATTERING CARD      001860
0042      5 LX(I,J)=L-1                                   001870
0043      MX=MX-2                                         001880
0044      J=2                                             001890
0045      ES(I,J)=X(MN+2)                                001900
0046      L=0                                             001910
0047      GO TO 4                                         001920
C          LAST CARD WAS ISOTOPE CARD                   001930
0048      6 LX(I,J)=L-1                                   001940
0049      MX=MX-4                                         001950
0050      I=I+1                                          001960
0051      H(I) =X(MN)                                    001970
0052      AG(I) =DLX(MN)                                 001980
0053      SPIN(I)=X(MN+1)                                001990
0054      RP(I) =DLX(MN+1)                               002000
0055      GO TO 3                                         002010
C          LAST CARD WAS SAMPLE CARD                   002020
C          XN(N): SAMPLE THICKNESS (NUCLEI/B)           002030
C          XR(N): SAMPLE RADIUS (NUCLEI/R)             002040
0056      7 LX(I,J)=L-1                                   002050
0057      MX=MX-4                                         002060
0058      IX=I                                            002070
0059      N=N+1                                           002080
0060      XN(N)=X(MN)                                     002090
0061      XR(N)=DLX(MN)                                  002100
0062      KH=0                                            002110
C          READ SPECIFICATIONS OF N-TH MEASUREMENT      002120
C          FP(N) : FLIGHT PATH (M)                     002130
C          DLFP(N): SAMPLE THICKNESS (M)                002140
C          TC(N) : CHANNEL WIDTH (NS)                   002150
C          TR(N) : FWHM OF GAMMA PEAK (NS)              002160
C          ZH(N) : NUMBER OF MONTE CARLO HISTORIES PER ENERGY 002170
C          TS(N) : TIME SHIFT (NS) TO COMPENSATE ZERO-TIME ERRORS 002180
0063      12 READ(5,107)FP(N),DLFP(N),TC(N),TR(N),ZH(N),TS(N) 002190
0064      107 FORMAT(6E10.5)                              002200

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0065      8 KN=KH+1                                C02210
0066      KH=KN+1                                C02220
C          READ DATA CARD                      C02230
C          EF(K) : EXPERIMENTAL ENERGY (KEV)   C02240
C          Y(K) : EXPERIMENTAL YIELD           C02250
C          DLY(K): EXPERIMENTAL YIELD UNCERTAINTY (%) C02260
0067      READ(5,108)(EE(K), Y(K),DLY(K),K=KN,KH) C02270
0068      108 FORMAT(6E10.5)                     C02280
C          WAS THIS LAST DATA CARD?           C02290
C          NOTE: LAST DATA CARD IS FOLLOWED BY A BLANK CARD (1 FOR EACH C02300
C          MEASUREMENT). A 2ND BLANK CARD SIGNALS END OF PROBLEM. C02310
C          INPUT FOR ANOTHER PROBLEM MAY FOLLOW. C02320
0069      IF(DLY(KH).NE.0.)GO TO 8               C02330
0070      DO 9 K=KN,KH                           C02340
0071      IF(DLY(K).NE.0.)GO TO 9               C02350
0072      KX(N)=K-1                              C02360
0073      GO TO 10                               C02370
0074      9 CONTINUE                             C02380
0075      10 IF(KX(N).EQ.KH-2)GO TO 11           C02390
0076      READ(5,109)DUMMY                       C02400
0077      109 FORMAT(E10.5)                     C02410
0078      11 N=N+1                              C02420
0079      READ(5,110)XN(N),XR(N)                C02430
0080      110 FORMAT(2E10.5)                    C02440
0081      KH=XK(N-1)                             C02450
C          END OF INPUT?                       C02460
0082      IF(XN(N).GT.0.)GO TO 12               C02470
0083      NX=N-1                                C02480
C          CALCULATE ABSOLUTE ERROR           C02490
0084      KH=XK(NX)                             C02500
0085      DO 13 K=1,KH                          C02510
0086      DLY(K)=DLY(K)*ABS(Y(K))*0.1           C02520
0087      13 CONTINUE                             C02530
C          CHECK WHETHER ALL CAPTURE EFFICIENCIES ARE EQUAL C02540
0088      LL=0                                  C02550
0089      EFC=1.                                C02560
0090      DO 14 I=1,IX                          C02570
0091      JH=JX(I)                              C02580
0092      DO 15 J=1,JH                          C02590
0093      IF(LX(I,J).EQ.0)GO TO 15              C02600
0094      LH=LX(I,J)                            C02610
0095      DO 16 L=1,LH                          C02620
0096      LL=LL+1                              C02630
0097      IF(LL.EQ.1)EFC=EFF(I,J,L)             C02640
0098      IF(LL.GT.1.AND.EFF(I,J,L).NE.EFC)EFC=0. C02650
0099      16 CONTINUE                             C02660
0100      15 CONTINUE                             C02670
0101      14 CONTINUE                             C02680
C          CORRECT TRUNCATION ERRORS OF ENERGIES AND PRINT MEASURED DATA C02690
C          C02700
C          C02720
C          WRITE RUN INFORMATION                C02730
0102      WRITE(6,113)                           C02740
0103      113 FORMAT(1H0///4H SPECIFICATIONS OF UTILIZED MEASUREMENTS: / C02750
          . 4.H =====// 126H C02760
          1 MMENT. SAMPLE SAMPLE FLIGHT SAMPLE CHANNEL C02770
          2 FWHM MONTE TIME SHIFT /126H C02780
          3 NO. THICKNESS RADIUS PATH THICKNESS WIDTH C02790

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      4  GAMMA PEAK CARLO      CORRECTION      /126H  002800
      5  (NUCLEI/B) (NUCLEI/B) (M)          (M)      (NS)  002810
      6  (NS) HISTORIES (NS)                /)      002820
0104      DO 51 N=1,NX
0105      WRITE (6,114)N,XN(N),XR(N),FP(N),DLFP(N),TC(N),TB(N),ZH(N),TS(N)  002830
0106      5: CONTINUE
0107      114 FORMAT(I3,1PF16.2,1PF12.3,0PF10.4,F12.4,F10.3,F12.3,F12.0,F12.3)  002840
      C      WRITE CRUSS SECTION PARAMETERS 002850
      C
0108      CALL PAROUT 002860
      C
0109      IF (IX.GT.10.OR.NX.GT.5.OR.KX(NX).GT.512)WRITE(6,115)IX,NX,KX(NX)  002870
0110      115 FORMAT(1H //' TOO MUCH INPUT:', I6, ' ISOTOPES, '/ 002880
      1      ' ', I6, ' MEASUREMENTS, '/ 002890
      2      ' ', I6, ' DATA POINTS.' ) 002900
0111      IF (IX .GT. 10)IX = 10 002910
0112      IF (NX .GT. 5)NX = 5 002920
0113      IF (KX(NX).GT.512)KX(NX)=512 002930
0114      RETURN 002940
0115      999 STOP 002950
0116      END 002960

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C
C
0001      SUBROUTINE ETEC(EE,Y,DLY,FP,TS,KX,NX)
C
C          ETEC CORRECTS NONLINEARITIES OF THE TIME-OF-FLIGHT SCALE
C          CAUSED BY TOO FEW DIGITS USED FOR THE INPUT ENERGIES
C          BY A LINEAR LEAST-SQUARES FIT TO THE FLIGHT TIMES.
C          EE: ENERGIES (KEV), FIRST TRUNCATED, THEN CORRECTED
C          Y: CAPTURE YIELD DATA (B)
C          DLY: UNCERTAINTIES (B)
C          FP: FLIGHT PATH
C          TS: TIME SHIFT CORRECTION
C          KX: ARRAY BOUNDARIES FOR Y AND DLY
C          NX: NUMBER OF TIME-OF-FLIGHT RUNS
C
0002      DIMENSION EE(512),Y(512),DLY(512),FP(5),TS(5),KX(5)
0003      WRITE(6,100)
0004      100 FORMAT(1HC//' INPUT DATA TO BE FITTED: '//
C          1 ' * ====='///
C          2 ' TRUNCATED   CORRECTED   CAPTURE YIELD'//
C          3 ' ENERGY     ENERGY     DATA   '//
C          4 ' (KEV)       (KEV)       (B)     ')
0005      KH=0
0006      DO 3 N=1,NX
0007      KN=KH+1
0008      KH=KX(N)
0009      CG=TS(N)/(72.3*FP(N)*SQRT(1000.))
0010      WRITE(6,101)
0011      101 FORMAT(1H /)
0012      A11=0.
0013      A12=0.
0014      A22=0.
0015      C1 =0.
0016      C2 =0.
0017      DO 1 K=KN,KH
0018      FK=K
0019      ETA=1./SQRT(EE(K))
0020      A11=A11+1.
0021      A12=A12+FK
0022      A22=A22+FK**2
0023      C1 =C1+ETA
0024      C2 =C2+ETA*FK
0025      1 CONTINUE
0026      DET=A11*A22-A12*A12
0027      B11= A22/DET
0028      B12=-A12/DET
0029      B22= A11/DET
0030      X1=B11*C1+B12*C2+CG
0031      X2=B12*C1+B22*C2
0032      DO 2 K=KN,KH
0033      FK=K
0034      YK=X1+FK*X2
0035      EU=EE(K)
0036      EE(K)=1./YK**2
0037      WRITE(6,102)EU,EE(K),Y(K),DLY(K)
0038      102 FORMAT(F8.3,F12.3,F 2.6,2H+-,F8.6)
0039      2 CONTINUE
0040      3 CONTINUE
0041      RETURN
0042      END

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C
C
0001 SUBROUTINE PAROUT
C
C PAROUT PRINTS THE CROSS SECTION PARAMETERS
C
0002 COMMON
      HI,GI,Z11,Z21,Z22,IZ,CHISQ,CHISQ0,
      ZI,IX,J,JX(11),KX(5),L,LX(11,2),M,MA,MX,N,NX,KKX,MP(11,2),MR(11,2),
      ATITLE(20),ZIT,EPS,E1,E2,TEFF,H(11),AG(11),SPIN(11),RP(11),
      4G(11,2),CS(11,2),X(200),DLX(200),ES(11,2),EFF(11,2,200),
      5FP(5),DLFP(5),TC(5),TB(5),ZH(5),EE(512),Y(512),DLY(512),
      6XN(6),XR(6),TMX,AL(10),D(10),CT(10,2),CG(10,2),M1,M2,M0,
      7DLE,EMN,ST(2048),DLST(2048),DST(20),SG(2048),DLSG(2048),DSG(20),
      8SC(2048),T(2048),YY(2048),DYY(2048,20),ZZ(512),Z(512),DZ(20),EFC,
      9F(201),RF(201)
0003 COMPLEX HI,GI,Z11,Z21,Z22,DZ11,DZ21,DZ22
C
0004 MH=MX
0005 WRITE(6,99)
0006 99 FORMAT(1H1//' CROSS SECTION PARAMETERS: '/
      ' ====='/)
0007 WRITE(6,100)
0008 100 FORMAT(
      1 127H ABUN- ATOMIC TARGET P-WAVE COMP. S-WAVE
      2 S-WAVE INEL. RESONANCE PARTIAL WIDTHS FOR
      3 DETECT. /127H DANCE WEIGHT SPIN RADIUS SPIN STRENGTH
      4 RADIUS THRESH. ENERGY EL. SCATT. INEL. SCATT. CAPTURE
      5 EFFIC. /127H (FM) FUNCTION
      6 (FM) (KEV) (KEV) (KEV) (KEV) (KEV)
      7 /127H /UNCERT. /UNCERT. /UNCERT. /UNCERT. /UNCERT.
      8 /)
0009 MX=0
0010 DO 1 I=1,IX
0011 J=1
0012 IF(LX(I,1).EQ.0)GO TO 2
0013 MN=MX+1
0014 MX=MN+5
0015 WRITE(6,101)H(I),AG(I),SPIN(I),RP(I),CS(I,J),X(MN),X(MN+1),ES(I,J)
      1,X(MN+2),X(MN+3),X(MN+4),X(MX),EFF(I,J,1),(DLX(M),M=MN,MX)
0016 101 FORMAT(F7.4,F8.1,F7.1,F9.3,F8.1,1PE12.3,OPF8.3,F9.3,F12.3,1P3E12.3
      1,OPF8.3/39X,E12.3,F8.3,9X,F12.3,3E12.3/)
0017 GO TO 3
0018 2 MN=MX+1
0019 MX=MN+1
0020 WRITE(6,102)H(I),AG(I),SPIN(I),RP(I),CS(I,1),X(MN),X(MX),ES(I,1)
      1,DLX(MN),DLX(MX)
0021 102 FORMAT(F7.4,F8.1,F7.1,F9.3,F8.1,1PE12.3,OPF8.3,F9.3/39X,E12.3,F8.3
      1/)
0022 GO TO 4
0023 3 IF(LX(I,J).LE.1)GO TO 4
0024 LMX=LX(I,J)
0025 DO 5 L=2,LMX
0026 MN=MX+
0027 MX=MN+5
0028 WRITE(6,103)(X(M),M=MN,MX),FFF(I,J,L),(DLX(M),M=MN,MX)
0029 103 FORMAT(68X,F12.3,1P3E12.3,OPF8.3/68X,F12.3,3E12.3/)
0030 5 CONTINUE

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0031      4 IF(J.EQ.JX(I))GO TO 1      004170
0032      J=2                          004180
0033      IF(LX(I,2).EQ.0)GO TO 6      004190
0034      MN=MX+1                       004200
0035      MX=MN+5                       004210
0036      WRITE(6,104)CS(I,2),X(MN),X(MN+1),ES(I,2),X(MN+2),X(MN+3),X(MN+4), 004220
      IX(MX),EFF(I,J,1),(DLX(M),M=MN,MX)
0037      104 FORMAT(31X,F8.1,1PE12.3,0PF8.3,F9.3,F12.3,1P3E12.3,0PF8.3/39X,E12. 004240
      13,F8.3,9X,F12.3,3E12.3/)
0038      GO TO 3                       004250
0039      6 MN=MX+1                     004260
0040      MX=MN+1                       004270
0041      WRITE(6,105)CS(I,2),X(MN),X(MX),ES(I,2),DLX(MN),DLX(MX)      004280
0042      105 FORMAT(31X,F8.1,1PE12.3,0PF8.3,F9.3/39X,E12.3,F8.3/)      004300
0043      1 CONTINUE                    004310
0044      MX=MN                          004320
0045      IF(IZ.GT.0)RETURN              004330
0046      RETURN                        004340
0047      END                          004350
```

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C
C
0001 C SUBROUTINE INDEX 004360
C 004370
C INDEX PROVIDES, FOR EACH ISOTOPE AND COMPOUND SPIN, SUBSCRIPTS 004390
C STRUCTURING THE CROSS SECTION PARAMETER ARRAYS X AND DLX. 004400
C MP(I,J): BEGIN OF POTENTIAL-SCATTERING PARAMETERS, 004410
C MR(I,J): BEGIN OF RESONANCE PARAMETERS 004420
C FOR THE I-TH ISOTOPE AND THE J-TH COMPOUND SPIN, 004430
C MX: TOTAL NUMBER OF PARAMETERS, 004440
C MA: NUMBER OF ADJUSTED PARAMETERS. 004450
C 004460
0002 C COMMON 004470
C IHI,GI,Z11,Z21,Z22,IZ,CHISQ,CHISQ0, 004480
C 21,IX,J,JX(11),KX(5),L,LX(11,2),M,MA,MX,N,NX,KKX,MP(11,2),MR(11,2), 004490
C 3TITLE(20),ZIT,EPS,E1,E2,TEFF,H(11),AG(11),SPIN(11),RP(11), 004500
C 4G(1,2),CS(11,2),X(200),DLX(200),ES(11,2),EFF(11,2,200), 004510
C 5FP(5),DLFP(5),TC(5),TB(5),ZH(5),EE(512),Y(512),DLY(512), 004520
C 6XN(6),XR(6),TMX,AL(10),D(10),CT(10,2),CG(10,2),M1,M2,M0, 004530
C 7DLB,EMN,ST(2048),DLST(2048),DST(20),SG(2048),DLSG(2048),DSG(20), 004540
C 8SC(2048),T(2048),YY(2048),OYY(2048,20),ZZ(512),Z(512),DZ(20),EFC, 004550
C 9F(201),RF(201) 004560
0003 C COMPLEX HI,GI,Z11,Z21,Z22,DZ11,DZ21,DZ22 004570
C 004580
0004 C MSUM=1 004590
0005 DO 1 I=1,IX 004600
0006 JH=JX(I) 004610
0007 DO 1 J=1,JH 004620
0008 MP(I,J)=MSUM 004630
0009 MR(I,J)=MSUM+2 004640
0010 1 MSUM=MR(I,J)+4*LX(I,J) 004650
0011 MX=MSUM-1 004660
0012 MA=0 004670
0013 DO 2 M=1,MX 004680
0014 IF(DLX(M).GT.0.)MA=MA+1 004690
0015 2 CONTINUE 004700
0016 IF(MX.GT.200.OR.MA.GT.20)WRITE(6,100)MX,MA 004710
0017 100 FORMAT(1H//' TOO MANY PARAMETERS:',I6,' ALL TOGETHER,'/ 004720
C ' ',I6,' ADJUSTED ONES.')
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```

C
C
0001      SUBROUTINE MEV
C
C          MEV CONVERTS ENERGIES FROM KEV TO MEV
C
0002      COMMON
          HI,GI,Z11,Z21,Z22,IZ,CHISQ,CHISQ0,
          CI,IX,J,JX(11),KX(5),L,LX(11,2),M,MA,MX,N,NX,KKX,MP(11,2),MR(11,2),
          BTITLE(20),ZIT,EPS,E1,E2,TEFF,H(11),AG(11),SPIN(11),RP(11),
          4G(11,2),CS(11,2),X(200),DLX(200),ES(11,2),EFF(11,2,200),
          5FP(5),DLFP(5),TC(5),TB(5),ZH(5),EE(512),Y(512),DLY(512),
          6XN(6),XR(6),TMX,AL(10),D(10),CT(10,2),CG(10,2),M1,M2,MO,
          7DLE,EMN,ST(2048),DLST(2048),DST(20),SG(2048),DLSG(2048),DSG(20),
          8SC(2048),T(2048),YY(2048),DYY(2048,20),ZZ(512),Z(512),DZ(20),EFC,
          9F(201),RF(201)
0003      COMPLEX HI,GI,Z11,Z21,Z22,DZ11,DZ21,DZ22
C
0004      IF(IZ.GT.1)GO TO 1
0005      E1=E1*.001
0006      E2=E2*.001
0007      DO 2 I=1,IX
0008      JH=JX(I)
0009      DO 2 J=1,JH
0010      ES(I,J)=ES(I,J)*.001
0011      IF(LX(I,J).EQ.0)GO TO 2
0012      M1=MR(I,J)
0013      M2=MR(I,J)+4*LX(I,J)-4
0014      DO 2 M=M1,M2,4
0015      X(M)=X(M)*.001
0016      X(M+1)=SIGN(SQRT(.001*ABS(X(M+1))),X(M+1))
0017      X(M+2)=SIGN(SQRT(.001*ABS(X(M+2))),X(M+2))
0018      X(M+3)=X(M+3)*.001
0019      2 CONTINUE
0020      KH=KX(NX)
0021      DO 3 K=1,KH
0022      3 EE(K)=.001*EE(K)
0023      IF(IZ.EQ.1)RETURN
0024      EMN=EMN*.001
0025      DLE=DLE*.001
0026      RETURN
0027      END

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C
C
0001      SUBROUTINE CONIJ                                005200
C
C          CONIJ YIELDS CONSTANTS FOR THE CROSS SECTION CALCULATION 005210
C
C          COMMON                                         005220
0002      COMMON                                         005230
          HI,GI,Z11,Z21,Z22,IZ,CHISQ,CHISQ0,              005240
          ZI,IX,J,JX(11),KX(5),L,LX(11,3),M,MA,MX,N,NX,KKX,MP(11,2),MR(11,2), 005250
          3TITLE(20),ZIT,EPS,E1,E2,TEFF,H(11),AG(11),SPIN(11),RP(11),      005260
          4G(11,2),CS(11,2),X(200),DLX(200),ES(11,2),EFF(11,2,200),       005270
          5FP(5),DLFP(5),TC(5),TB(5),ZH(5),EE(512),Y(512),DLY(512),       005280
          6XN(6),XR(6),TMX,AL(10),D(10),CT(10,2),CG(10,2),M1,M2,M0,      005290
          7BLE,EMN,ST(2048),DLST(2048),DST(20),SG(2048),DLSG(2048),DSG(20), 005300
          8SC(2048),T(2048),YY(2048),DYY(2048,20),ZZ(512),Z(512),DZ(20),EFC, 005310
          9F(201),RF(201)                                           005320
0003      COMPLEX HI,GI,Z11,Z21,Z22,DZ11,DZ21,DZ22          005330
C
0004      DO 1 I=1,IX                                         005340
0005      AL(I)=1.+1./AG(I)                                    005350
0006      CO=H(I)*AL(I)**2                                    005360
0007      D(I)=SQRT(344.68*TEFF/AG(I))*0.00001              005370
0008      JH=JX(I)                                           005380
0009      DO 2 J=1,JH                                        005390
0010      CT(I,J)=CO*G(I,J)                                  005400
0011      CG(I,J)=-5*CT(I,J)                                 005410
0012      2 CONTINUE                                         005420
0013      1 CONTINUE                                         005430
0014      RETURN                                             005440
0015      END                                               005450

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C
C
C001 SUBROUTINE CONE                                005480
C
C          CONE YIELDS THE COEFFICIENTS OF THE NORMAL EQUATIONS  005490
C
C
C002 COMMON                                         005500
C          IHT,GI,Z11,Z21,Z22,IZ,CHISQ,CHISQ0,          005510
C          2I,IX,J,JX(11),KX(5),L,LX(11,2),M,MA,MX,N,NX,KKX,MP(11,2),MR(11,2),  005520
C          3TITLE(20),ZIT,EPS,E1,E2,TEFF,H(11),AG(11),SPIN(11),RP(11),      005530
C          4G(11,2),CS(11,2),X(200),DLX(200),ES(11,2),EFF(11,2,200),      005540
C          5FP(5),DLFP(5),TC(5),TB(5),ZH(5),EE(512),Y(512),DLY(512),      005550
C          6XN(6),XR(6),TMX,AL(10),D(10),CT(10,2),CG(10,2),M1,M2,MO,      005560
C          7DLE,EMN,ST(2048),DLST(2048),DST(10),SG(2048),DLSG(2048),DSG(20),  005570
C          8SC(2048),T(2048),YY(2048),DYY(2048,20),ZZ(512),Z(512),DZ(20),EFC,  005580
C          9F(201),RF(201)                                         005590
C003 COMMON /ABC/ A(20,20),B(20,20),C(20)              005600
C004 DOUBLE PRECISION A,B                                005610
C005 COMMON/SEYCD/YS(2048)                               005620
C006 COMPLEX HI,GI,Z11,Z21,Z22,DZ11,DZ22,DZ22          005630
C007 DIMENSION FS(201),DYC(201,20)                    005640
C
C          INITIALIZE
C008 DO 1 M=1,MA                                         005650
C009 C(M)=0.                                             005660
C010 DO 1 MM=1,MA                                       005670
C011 1 A(M,MM)=0.                                       005680
C012 KH=0                                               005690
C
C          BEGIN LOOP OF MEASUREMENTS
C013 DO 2 N=1,NX                                       005700
C          GET PRIMARY, SECONDARY AND MULTIPLE-COLLISION YIELDS  005710
C
C          CALL PRY
C014 CALL SEY                                           005720
C015 CALL MUY                                           005730
C016                                                    005740
C
C          CONSTANTS FOR RESOLUTION WIDTH
C017 V1=(DLFP(N)/FP(N))**2/.5                          005750
C018 V2=(TB(N)/(60.25*FP(N)))**2                      005760
C          WRITE TABLE HEADING
C019 WRITE(6,101)IZ,N,XN(N),FP(N)                    005770
C020 101 FORMAT('1'/' ITERATION STEP',I3,' TIME-OF-FLIGHT RUN NO.',I3//  005780
C          118H SAMPLE THICKNESS:,.PE10.3,9H NUCLEI/B /      005790
C          218H FLIGHT PATH      :.1PE10.3,2H M /           005800
C          WRITE(6,102)
C021 102 FORMAT(                                        005810
C          1' NEUTRON      MEASURED          CALCULATED   FROM MULT.   CALC  005820
C          2ULATED' /                                005830
C          3' ENERGY     YIELD-OVER-THICKNESS VALUES     COLLISIONS  CRDS  005840
C          4S SECTION' /                                005850
C          5' (KEV)       (MILLI-BARN)          (MILLI-BARN) (PERCENT)  (MIL  005860
C          6LI-BARN)' /)                                005870
C          BEGIN ENERGY LOOP
C023 KN=KN+1                                           005880
C024 KH=KX(N)                                           005890
C025 DO 3 KK=KN,KH                                       005900
C026 XL=(EE(KK)-EMN)/OLE+1.                             005910
C027 L=XL                                               005920

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0028      DD=XL-FLOAT(L)                                006050
0029      RW=EE(KK)*SQRT(VI+V2*EE(KK))                 006060
0030      QD=DLE/RW                                     006070
0031      KR=3./QD                                       006080
0032      STL =ST(L) +DD*(ST(L+1) -ST(L) )            006090
0033      SGL =SG(L) +DD*(SG(L+1) -SG(L) )            006100
0034      DLSGL=DLSG(L)+DD*(DLSG(L+1)-DLSG(L))        006110
0035      TL  =T(L) +DD*(T(L+1) -T(L) )              006120
C          SELF-SHIELDING EFFECT ON MEAN PATH LENGTH IN SAMPLE 006130
0036      DELXL=(2./(XN(N)*STL)-(1.+TL)/(1.-TL))*DLFP(N)/FP(N)*EE(KK)/DLE 006140
0037      XL=XL+DELXL                                    006150
0038      L=XL                                           006160
0039      DD=XL-FLOAT(L)                                006170
0040      IF(KR+L.GT.KKX)KR=KKX-L                       006180
0041      IF(KR.GT.100)WRITE(6,100)IZ,KR                006190
0042      100 FORMAT(2H IZ,17H: ITERATION, KR =,16,'READJUSTED TO 100' ) 006200
0043      IF(KR.GT.100)KR=100                            006210
C          PREPARE INTEGRANDS FOR RESOLUTION BROADENING 006220
0044      K1=101-KR                                      006230
0045      K2=101+KR                                      006240
C                                                    006250
0046      IF(AG(1).NE.56.)                               006260
C          CALL TGAUSS(QD,R,F,K2)                       006270
0047      IF(AG(1).EQ.56.)                               006280
C          CALL CHISQ4(QD,R,F,K2)                       006290
C                                                    006300
0048      DO 4 K=K1,K2                                    006310
0049      LK=L+K-101                                     006320
0050      F(K) =RF(K)*(YY(LK)+DD*(YY(LK+1)-YY(LK)))    006330
0051      FS(K) =RF(K)*(YS(LK)+DD*(YS(LK+1)-YS(LK)))    006340
0052      DO 4 M=1,MA                                     006350
0053      DYC(K,M)=RF(K)*(DYY(LK,M)+DD*(DYY(LK+1,M)-DYY(LK,M))) 006360
0054      4 CONTINUE                                     006370
C          RESOLUTION BROADENING:                       006380
C                                                    006390
0055      CALL SIMP(F ,K1,K2,Z(KK))                     006400
0056      CALL SIMP(FS,K1,K2,ZS)                        006410
C                                                    006420
0057      ZS=ZS/Z(KK)                                    006430
0058      Z(KK)=Z(KK)/XN(N)                              006440
0059      ZDEV=(Y(KK)-Z(KK))/DLY(KK)                    006450
C          RESOLUTION BROADENING OF DERIVATIVES:       006460
0060      DO 5 M=1,MA                                    006470
C                                                    006480
0061      CALL SIMP(DYC(1,M),K1,K2,DZ(M))               006490
C                                                    006500
0062      DZ(M)=DZ(M)/DLY(KK)/XN(N)                    006510
C          COEFFICIENTS OF THE NORMAL EQUATIONS        006520
0063      C(M)=C(M)+DZ(M)*ZDEV                          006530
0064      DO 5 MM=1,M                                     006540
0065      5 A(M,MM)=A(M,MM)+DZ(M)*DZ(MM)               006550
0066      CHISQ=CHISQ+ZDEV**2                            006560
C          PRINT TABLE ENTRY                          006570
0067      WRITE(6,103)EE(KK),Y(KK),DLY(KK),Z(KK),ZS,SGL,DLSGL 006580
0068      103 FORMAT(' ',3PF7.3,F10.3,' +',F8.3,F13.3,2PF13.3,3PF13.3,' +',F8.3) 006590
0069      WRITE(6,104)                                    006600
0070      104 FORMAT('+',10X,'_',17X,'_')               006610
0071      3 CONTINUE                                     006620

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0072      C      END OF ENERGY LOOP
          C      2 CCNTINUE
0073      C      END OF LOOP OF MEASUREMENTS
0074      C      RETURN
          C      END
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C
C
0001      SUBROUTINE PRY
C
C          PRY YIELDS THE PRIMARY YIELD, ITS DERIVATIVES AND THE
C          TRANSMISSION
C          E: NEUTRON ENERGY
C          K: ENERGY-GRID LABEL
C
0002      COMMON
          IHI,GI,Z11,Z21,Z22,IZ,CHISG,CHISQO,
          2I,IX,J,JX(11),KX(5),L,LX(11,2),M,MA,MX,N,NX,KKX,MP(11,2),MR(11,2),
          3TITLE(20),ZIT,EPS,E1,E2,TEFF,H(11),AC(11),SPIN(11),RP(11),
          4G(11,2),CS(11,2),X(200),DLX(200),ES(11,2),EFF(11,2,200),
          5FP(5),DLFP(5),TC(5),TB(5),ZH(5),EF(512),Y(512),OLY(512),
          6XN(6),XR(6),TMX,AL(10),D(10),CT(10,2),CG(10,2),M1,M2,MO,
          7DLE,EMN,ST(2048),DLST(2048),DST(20),SG(2048),DLSG(2048),DSG(20),
          8SC(2048),T(2048),YY(2048),DYY(2048,20),ZZ(512),Z(512),DZ(20),EFC,
          9F(201),RF(201)
0003      COMMON /ABC/ A(20,20),B(20,20),C(20)
0004      DOUBLE PRECISION A,B
0005      COMMON/ I1/ SG11(2048),DLSG11(2048)
0006      COMPLEX HI,GI,Z11,Z21,Z22,DZ11,DZ21,DZ22
C
0007      KH=KX(N)
0008      KN=1
0009      IF(.GT.1)KN=KX(N-1)+1
C
0010      EMIN=EE(KN)
0011      EMAX=EE(KH)
0012      V1=(DLFP(N)/FP(N))**2/1.5
0013      V2=(TB(N)/(63.259*FP(N)))**2
0014      RW1=EMIN*SQRT(V1+V2*EMIN)
0015      RW =EMAX*SQRT(V1+V2*EMAX)
0016      ET=2.*EMAX *XR(N)*DLFP(N)/(XN(N)*FP(N))
0017      EMN=EMIN-3.*RW1
0018      EMX=EMAX+3.*RW +ET
C
0019      IS VALIDITY INTERVAL WIDE ENOUGH?
0020      IWARN=0
0021      HINT=.5*(EMX-EMN)
0022      EBAR=.5*(EMX+EMN)
0023      IF(E1.LE.EBAR-HINT)GO TO 22
0024      IWARN=1
0025      F1=EBAR-HINT
0026      22 IF(E2.GE.EBAR+HINT)GO TO 23
0027      IWARN=1
0028      E2=EBAR+HINT
0029      23 IF(IWARN.EQ.1)WRITE(6,101)E1,E2
101  FORMAT(///
1' WARNING: VALIDITY RANGE READJUSTED,/'
2' PARAMETERS ARE NOW VALID FROM',3PF8.3,' TG',3PF8.3,' KEV'/)
C
C          TAKE INTERNAL GRID INTERVAL (DLE) AS HALF THE SMALLEST DOPPLER
C          WIDTH. IF MORE THAN 2048 INTERNAL GRID POINTS RESULT, OR IF
C          MORE THAN 101 GRID POINTS FALL WITHIN THRICE THE LARGEST
C          RESOLUTION WIDTH, THE DOPPLER WIDTH (I. E. THE TEMPERATURE)
C          MUST BE INCREASED.
0030      DLE=.5*D(1)*SQRT(EMIN)
0031      DLEJ=DLE
    
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0032      DLE1=3.*KX/1.1.
0033      DLE2=(EMX-EMN)/2.144.
0034      DLE=AMAX1(DLE0,DLE1,DLE2)
0035      IF(DLE.LE.DLE0)GO TO 2
0036      QD=DLE/DLE0
0037      TFFF=TFFF*QD**2
0038      DO 1 I=1,IX
0039      1 D(I)=D(I)*QD
0040      DW=D(1)*SQRT(EMIN)*1000000.
0041      ER QD=EMIN*1000.
0042      WRITE(6,103)TFFF,DW,EL000
0043      100 FORMAT(25H0EFFECTIVE TEMPERATURE : ,F7.0,13H DEG. KELVIN,
1/      25H RESULTING DOPPLER WIDTH: ,F7.3, 6H EV AT,F9.3,4H KEV/)
0044      2 KXX=(EMX-EMN)/DLE+1.
C      BEGIN ENERGY LOOP
0045      E=EMN-DLE
0046      DO 11 K=1,KXX
0047      E=E+DLE
C      CALCULATE CROSS SECTIONS AND DERIVATIVES
C
0048      CALL XSECT(E,K)
C
C      TRANSMISSION:
0049      T(K)=EXP(-XN(N)*ST(K))
C      PRIMARY YIELD:
0050      YY(K)=(1.-T(K))*SG(K)/ST(K)
C      DERIVATIVES:
0051      DO 3 M=1,MA
0052      DYY(K,M)=YY(K)*{DSG(M)/SG(K)-DST(M)/ST(K)}
1      +XN(N)*SG(K)*T(K)*DST(M)/ST(K)
0053      3 CONTINUE
C      CALCULATE CROSS SECTION ERRORS
0054      VST=}.
0055      VSG=}.
0056      IF(IZ.EQ.1)GO TO 5
0057      DO 5 MM=1,MA
0058      DO 4 ML=1,MA
0059      VST=VST+DST(MM)*B(MM,ML)*DST(ML)
0060      VSG=VSG+DSG(MM)*B(MM,ML)*DSG(ML)
0061      4 CONTINUE
0062      5 CONTINUE
0063      6 DLST(K)=SQRT(VST)
0064      DLSG(K)=SQRT(VSG)
C      TREAT FIRST ISOTDPE SEPARATELY
0065      VSG11=}.
0066      IF(IZ.EQ.1)GO TO 10
0067      IH=1
0068      IF(I(2).EQ.H(1))IH=2
0069      JH=JX(IH)
0070      MX1=MR(IH,JH)+4*LX(IH,JH)-1
0071      MA1=}
0072      DO 7 M=1,MX1
0073      IF(DLX(M).GT.0.)MA1=MA1+1
0074      7 CONTINUE
0075      DO 9 MM=1,MA1
0076      DO 8 ML=1,MA1
0077      VSG11=VSG11+DSG(MM)*B(MM,ML)*DSG(ML)
0078      8 CONTINUE

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```
0079          9 CONTINUE
0080          10 DLSG11(K)=SQRT(VSG11)/H(1)
0081          11 CCNTINUE
0082             RETURN
0083             END
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C
C
0001 SUBROUTINE XSECT(E,K)
C
C      XSECT CALCULATES CROSS SECTIONS AND DERIVATIVES.
C      E: NEUTRON ENERGY
C      K: ENERGY GRID SUBSCRIPT (INTERNAL OR "FINE" GRID)
C      S-WAVE CROSS SECTIONS ARE CALCULATED WITH THE REICH-MOORE
C      FORMALISM WITHOUT DOPPLER BROADENING.
C      HIGHER-ORDER PARTIAL-WAVE CROSS SECTIONS ARE CALCULATED
C      WITH DOPPLER-BROADENED SINGLE-LEVEL EXPRESSIONS WITHOUT
C      RESONANCE/POTENTIAL INTERFERENCE.
C
0002 COMMON
      IHI,GI,Z11,Z21,Z22,IZ,CHISO,CHISOO,
      ZI,IX,J,JX(11),KX(5),L,LX(11,2),M,MA,MX,N,NX,KKX,MP(11,2),MR(11,2),
      BTITLE(20),ZIT,EPS,C1,E2,TEFF,H(11),AG(11),SPIN(11),RP(11),
      4G(11,2),CS(11,2),X(20),DLX(200),ES(11,2),EFF(11,2,200),
      5FP(5),DLFP(5),TC(5),TB(5),ZH(5),EE(512),Y(512),CLY(512),
      6XM(6),XR(6),TMX,AL(10),D(10),CT(10,2),CQ(10,2),M1,M2,MQ,
      7DLE,EMN,ST(2048),DLST(2048),DST(20),SG(2048),DLSG(2048),DSG(20),
      8SC(2048),T(2048),YY(2048),DYY(2048,20),ZZ(512),Z(512),DZ(20),EFC,
      9F(20),RF(20)
0003 COMMON/GQSR/DR11(20),DR21(20),DR22(20)
0004 COMMON/I1/SGI1(2048),DLSGI1(2048)
0005 COMPLEX HI,GI,Z11,Z21,Z22,DZ11,DZ21,DZ22
0006 COMPLEX DR11,DR21,DR22,U11,DU11(20),U21,DU21(20),EX1,ZZZ
C      FUNCTION DEFINITION
0007 ABSQ(ZZZ)=REAL(ZZZ)**2+AIMAG(ZZZ)**2
C
0008 M=0
0009 ST(K)=0.
0010 SG(K)=0.
0011 SC(K)=0.
0012 DO 50 I=1,MA
0013   DST(I)=0.0
0014   DSG(I)=0.0
0015 50 CONTINUE
0016   PLQ2=1.3019/F
C      POTENTIAL-SCATTERING PHASE FACTOR, PREPARATION:
0017   AR1=(2.*E-E2-E1)/(E2-E1)
0018   AN1=AR1*0.95
0019   ATGH=.5*ALOG((1.+AR1)/(1.-AR1))
0020   AR2=SQRT(1.E6*E)*ATGH
0021   DO 1 I=1,IX
C      DOPPLER WIDTH
0022   DW=D(I)*SQRT(E)
C      CALCULATE POTENTIAL SCATTERING FOR P-WAVE
0023   XK1=.21969*SQRT(E)/AL(I)
0024   XC=XK1*FP(I)
0025   X1=XJ-ATAN(XC)
0026   SP=H(I)*PLQ2*AL(I)**2*6.*SIN(X1)**2
0027   ST(K)=ST(K)+SP
0028   JH=JX(I)
0029   DO 2 J=1,JH
0030     M1=MP(I,J)+2
0031     M2=MP(I,J)-2+4*LX(I,J)
C      S- OR P-WAVE LEVEL SEQUENCE?

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0032          IF(X(M1-1).EQ.0..AND.DLX(M1-1).EQ.0.)GO TO 7          008450
C              S=WAVE LEVELS:          008460
0033          XT=CT(I,J)*PLQ2          008470
0034          XG=CG(I,J)*PLQ2          008480
C              008490
0035          CALL KMAT(E)          008500
0036          CALL EFFI(E,K,EFK)          008510
C              008520
C              POTENTIAL-SCATTERING PHASE FACTOR:          008530
0037          AR3=X(M1-2)*AR2          008540
0038          XI1=-XK1*X(M1-1)+ATAN(AR3)          008550
0039          EX1=CEXP(2.*GI*XI1)          008560
C              COLLISION MATRIX ELEMENTS:          008570
0040          U11=EX1*(2.*Z11-1.)          008580
0041          U21=          2.*Z21          008590
C              TOTAL AND CAPTURE CROSS SECTION          008600
0042          ST(K)=ST(K)+XT*(1.-REAL(U11))          008610
0043          SC0=XG*(1.-ABSQ(U11)-ABSQ(U21))          008620
0044          SC(K)=SC(K)+SC0          008630
0045          SG(K)=SG(K)+SC0*EFK          008640
0046          IF(SC(K).LT.1.E-7*ST(K))SC(K)=1.E-7*ST(K)          008650
0047          IF(SG(K).LT.1.E-7*ST(K))SG(K)=1.E-7*ST(K)          008660
C              DERIVATIVES WITH RESPECT TO POTENTIAL-SCATTERING PARAMETERS          008670
0048          MD=M          008680
0049          IF(DLX(M1-2).EQ.0.)GO TO 8          008690
0050          M=M+1          008700
0051          ABL= AR2/(1.+AR3**2)          008710
0052          DU11(M)=2.*GI*ABL*U11          008720
0053          8 IF(DLX(M1-1).EQ.0.)GO TO 9          008730
0054          M=M+1          008740
0055          ABL=-XK1          008750
0056          DU11(M)=2.*GI*ABL*U11          008760
0057          9 IF(LX(I,J).EQ.0.)GO TO 10          008770
C              DERIVATIVES WITH RESPECT TO RESONANCE PARAMETERS          008780
0058          DO 11 MM=M1,M2,4          008790
0059          DO 11 ML=1,4          008800
0060          MK=MM+ML-1          008810
0061          IF(DLX(MK).LE.0.)GO TO 11          008820
0062          M=M+1          008830
0063          DU11(M)=GI*(Z11**2*DR11(M)+2.*Z11*Z21*DR21(M)+Z21**2*DR22(M))*EX1          008840
0064          DU21(M)=Z21*Z11*DR11(M)+(Z21**2+Z11*Z22)*DR21(M)+Z21*Z22*DR22(M)          008850
0065          DU21(M)=DU21(M)*GI          008860
0066          11 CONTINUE          008870
0067          10 IF(M.EQ.MD)GO TO 2          008880
0068          ML=MD+1          008890
0069          DO 12 MM=ML,M          008900
0070          DST(MM)=-XT*REAL(DU11(MM))          008910
0071          DSG(MM)=-XT*REAL(DU11(MM)*CONJG(U11)+DU21(MM)*CONJG(U21))*EFK          008920
0072          12 CONTINUE          008930
0073          GO TO 2          008940
C              P-WAVE LEVELS:          008950
0074          7 IF(LX(I,J).EQ.0)GO TO 2          008960
0075          CT1=CT(I,J)*PLQ2*2.          008970
0076          L=0          008980
C              BEGIN RESONANCE LOOP          008990
0077          DO 3 MM=M1,M2,4          009000
0078          L=L+1          009010
0079          GN=X(MM+1)**2          009020

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0080      GX=X(MM+2)**2
0081      GT=GN+GX+X(MM+3)
0082      QGT=EFF(I,J,L)*X(MM+3)/GT
0083      MC=M
0084      IF(DLX(MM ).NE.0.)M=M+1
0085      IF(DLX(MM+1).NE.0.)M=M+1
0086      IF(DLX(MM+2).NE.0.)M=M+1
0087      IF(DLX(MM+3).NE.0.)M=M+1
          C      APPROXIMATE 1/E-WIDTH OF DOPPLER BROADENED RESONANCE
0088      HW=SQRT(DW**2+GT**2/2.77259)
          C      IS LEVEL FAR AWAY AND HENCE NEGLIGIBLE?
0089      IF(ABS(E-X(MM)).GT.6.*HW)GO TO 3
0090      XV=2.*(E-X(MM))/GT
0091      BV=2.*DW/GT
          C
0092      (ALL VOIGT(XV,BV,PSI,PHI))
          C
0093      SQ=CT1*GN/GT
0094      ST1=SQ*PSI
0095      SG1=ST1*QGT
          C      CROSS SECTIONS
0096      ST(K)=ST(K)+ST1
0097      SG(K)=SG(K)+SG1
0098      SC(K)=SC(K)+SG1/EFF(I,J,L)
          C      ANY DERIVATIVES NEEDED?
0099      IF(M.EQ.M0)GO TO 3
0100      M=MC
0101      A0=-4./(GT*RV**2)
0102      A1=A0*( PHI-XV*PSI)
0103      A2=A0*(1.-PSI-XV*PHI)
          C      DERIVATIVES
0104      IF(DLX(MM ).EQ.0.)GO TO 6
0105      M=M+1
0106      DST(M)=SQ*A1
0107      DSG(M)=DST(M)*QGT
0108      6 IF(DLX(MM+1).EQ.0.)GO TO 5
0109      M=M+1
0110      DST(M)=SQ*(2.*PSI/GN+A2)*X(MM+1)
0111      DSG(M)=DST(M)*QGT-2.*SG1*X(MM+1)/GT
0112      5 IF(DLX(MM+2).EQ.0.)GO TO 4
0113      M=M+1
0114      DST(M)=SQ*A2*X(MM+2)
0115      DSG(M)=DST(M)*QGT-2.*SG1*X(MM+2)/GT
0116      4 IF(DLX(MM+3).EQ.0.)GO TO 3
0117      M=M+1
0118      DST(M)=SQ*A2*.5
0119      DSG(M)=DST(M)*QGT+(ST1*EFF(I,J,L)-SG1)/GT
0120      3 CONTINUE
          C      (END OF RESONANCE LOOP)
0121      2 CONTINUE
          C      (END OF SPIN LOOP)
0122      IF(ZIT.LE.1..AND.AG(I).EQ.AG(1))SGI1(K)=SC(K)/H(1)
0123      1 CONTINUE
          C      (END OF ISOTOPE LOOP)
0124      RETURN
0125      END

```

```

C
C
0001      C      SUBROUTINE KMAT(E)
C
C      KMAT CALCULATES THE ELEMENTS OF THE MATRIX R=2*K
C      AND INVERTS THE MATRIX I-I*K = I-I*R/2.
C
0002      COMMON
      HI,G1,Z11,Z21,Z22,IZ,CHISQ,CHISQ0,
      ZI,IX,J,JX(11),KX(5),L,LX(11,2),M,MA,N,NX,KKX,MP(11,2),MR(11,2),
      3TITLE(20),ZIT,EPS,E1,E2,TEFF,H(11),AG(11),SPIN(11),RP(11),
      4G(11,2),CS(11,2),X(200),DLX(200),ES(11,2),EFF(11,2,200),
      5FP(5),DLFP(5),TC(5),TB(5),ZH(5),EE(512),Y(512),DLY(512),
      6XN(6),XR(6),TMX,AL(10),D(10),CT(10,2),CG(10,2),M1,M2,MO,
      7DLF,EMN,ST(2048),DLST(2048),DST(20),SG(2048),DLSG(2048),DSG(20),
      8SC(2048),T(2048),YY(2048),DYY(2048,20),ZZ(512),Z(512),DZ(20),EFC,
      9F(201),RF(201)
0003      COMMON/GOSR/DR11(20),DR21(20),DR22(20)
0004      COMPLEX HI,G1,Z11,Z21,Z22,DZ11,DZ21,CZ22
      I,FE,F1,F2,EX1,EX2,U11,U21,DU11(20),DU21(20),DET,R11,R21,R22
0005      COMPLEX DR11,DR21,DR22
C
C      R11=(0.,0.)
0006      R21=(0.,0.)
0007      R22=(0.,0.)
0008      IF(LX(I,J).EQ.0)GO TO 2
0009      (GO TO 2 IF THERE ARE NO RESONANCES)
C
0010      MD=M
0011      IF(DLX(M1-2).NE.0.)M=M+1
0012      IF(DLX(M1-1).NE.0.)M=M+1
0013      DO 3 MM=M1,M2,4
0014      SQ1=(E/ABS(X(MM)))**.25
0015      IF(E.LE.ES(I,J))SQ2=0.
0016      IF(E.GT.ES(I,J))SQ2=((E-ES(I,J))/ABS(X(MM)-ES(I,J)))**.25
0017      FE=1./(X(MM)-E-HI*X(MM+3))
0018      W1=SQ1*X(MM+1)
0019      W2=SQ2*X(MM+2)
0020      F1=W1*FE
0021      F2=W2*FE
0022      R11=R11+F1*W1
0023      R21=R21+F2*W1
0024      R22=R22+F2*W2
0025      IF(DLX(MM).EQ.0.)GO TO 4
0026      M=M+1
0027      DR11(M)=-F1*F1
0028      DR21(M)=-F2*F1
0029      DR22(M)=-F2*F2
0030      4 IF(DLX(MM+1).LE.0.)GO TO 5
0031      M=M+1
0032      DR11(M)=F1*SQ1*2.
0033      DR21(M)=F2*SQ1
0034      DR22(M)=0.
0035      5 IF(DLX(MM+2).EQ.0.)GO TO 6
0036      M=M+1
0037      DR11(M)=0.
0038      DR21(M)=SQ2*F1
0039      DR22(M)=SQ2*F2*2.
0040      6 IF(DLX(MM+3).LE.0.)GO TO 3

```

```

0041 M=M+1 010160
0042 DR11(M)=HI*F1*F1 010170
0043 DR21(M)=HI*F2*F1 010180
0044 DR22(M)=HI*F2*F2 010190
0045 3 CONTINUE 010200
0046 M=MC 010210
0047 2 DET=(1.-HI*R11)*(1.-HI*R22)-(HI*R21)**2 010220
0048 Z11=(1.-HI*R22)/DET 010230
0049 Z21=( HI*R21)/DET 010240
0050 Z22=(1.-HI*R11)/DET 010250
0051 RETURN 010260
0052 END 010270

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C 010280
C
0001 SUBROUTINE EFFI(E,K,EFK) 010290
C 010300
C EFFI YIELDS THE AVERAGE DETECTOR EFFICIENCY (EFK) 010310
C FOR THE I-TH ISOTOPE AND THE J-TH COMPOUND SPIN 010320
C 010330
0002 COMMON 010340
01HI,G1,Z11,Z21,Z22,IZ,CHISQ,CHISQ0, 010350
2I,IX,J,JX(11),KX(5),L,LX(11,2),M,MA,MX,N,NX,KKX,MP(11,2),MR(11,2), 010360
3TITLE(20),ZIT,EPS,E1,E2,TEFF,H(11),AG(11),SPIN(11),RP(11), 010370
4G(11,2),CS(11,2),X(200),DLX(200),ES(11,2),EFF(11,2,200), 010380
5FP(5),DLFP(5),TC(5),TB(5),ZH(5),EE(512),Y(512),OLY(512), 010390
6XN(6),XR(6),TMX,AL(10),D(10),CT(10,2),CG(10,2),M1,M2,M0, 010400
7DLE,EMN,ST(2048),DLST(2048),DST(20),SG(2048),DLSG(2048),DSG(20), 010410
8SC(2048),T(2048),YY(2048),DYY(2048,20),ZZ(512),Z(512),DZ(20),EFC, 010420
9F(201),RF(201) 010430
0003 COMPLEX HI,G1,Z11,Z21,Z22,DZ11,DZ21,DZ22 010440
C 010450
C ANY RESONANCES? 010460
0004 IF(LX(I,J).EQ.0)RETURN 010470
0005 EFK=EFC 010480
C SAME EFFICIENCY FOR ALL RESONANCES? 010490
0006 IF(EFC.NE.0.)RETURN 010500
0007 SUM=. 010510
0008 SOM=. 010520
0009 L=0 010530
0010 M1=MP(I,J)+2 010540
0011 M2=MP(I,J)-2+4*LX(I,J) 010550
0012 DC 1 MM=M1,M2,4 010560
0013 L=L+1 010570
0014 GN=X(MM+1)**2*SQRT(E/ABS(X(MM))) 010580
0015 GX=X(MM+2)**2 010590
0016 GT=GN+GX+X(MM+3) 010600
0017 TM=GN*X(MM+3)/(4.*(E-X(MM))**2+GT**2) 010610
0018 SUM=SUM+TM*EFF(I,J,L) 010620
0019 1 SOM=SOM+TM 010630
0020 EFK=SUM/SOM 010640
0021 RETURN 010650
0022 END 010660

```

```

C
C
0001 SUBROUTINE VOIGT(X,B,PSI,PHI)
C
C      CALCULATION OF SYMMETRIC VOIGT PROFILE
C      X : DEVIATION FROM RESONANCE ENERGY IN UNITS OF HALF
C          THE HALF WIDTH
C      B : DOPPLER WIDTH IN THE SAME UNITS
C      PSI : SYMMETRIC VOIGT PROFILE
C      PHI : ASYMMETRIC " "
C
C      YY=Y*Y, GF=(2*DLY/SQRT(PI))*EXP(-YY) GIVEN FOR DLY=0.7
C      C1=DLY/SQRT(PI),C2=DLY/PI FOR DLY=0.7, C3=2*SQRT(PI)
C
0002 DIMENSION YY(8),GF(8)
0003 DATA YY/.49,1.96,4.41,7.84,12.25,17.64,24.01,31.36/
0004 DATA GF/.483893E1,.111259E0,.960101E-2,.310948E-3,
0005 DATA C1/.394933/C2/.222817/C3/3.544908/
C
0006 XX=X*X
0007 BB=B*B
0008 PSI=C1/(1.+XX)
0009 PHI=PSI
0010 DO 1 N=1,8
0011 YYPB=YY(N)*BB
0012 A0=(1.-XX+YYPB)**2+4.*XX
0013 A1=(1.+XX+YYPB)*GF(N)/A0
0014 A2=(1.+XX-YYPB)*GF(N)/A0
0015 PSI=PSI+A1
0016 1 PHI=PHI+A2
0017 PHI=PHI*X
0018 IF(B.LE.C2)RETURN
0019 ETA=2./(C2*B)
0020 EXPC=(XX-1.)/BB+ETA
0021 IF(EXP0.GT.25.)RETURN
0022 XI=2.*X/BB
0023 ZETA=X*ETA
0024 ACOS=COS(XI)
0025 RSIN=SIN(XI)
0026 CCOS=EXP(-ETA)-COS(ZETA)
0027 DSIN=SIN(ZETA)
0028 F=C3*EXP(-EXP0)/((CCOS**2+DSIN**2)*B)
0029 PSI=PSI+F*(ACOS*CCOS-BSIN*DSIN)
0030 PHI=PHI-F*(BSIN*CCOS+ACOS*DSIN)
0031 RETURN
0032 END

```

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010670
010680
010690
010700
010710
010720
010730
010740
010750
010760
010770
010780
010790
010800
010810
010820
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010900
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010960
010970
010980
010990
011000
011010
011020
011030
011040
011050
011060
011070
011080
011090
011100
011110
011120

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C
C
0001      SUBROUTINE SEY                                011130
C
C          SEY YIELDS THE SECONDARY CAPTURE YIELD FOR AN INFINITE SLAB
C          THE TARGET NUCLEI ARE TAKEN TO BE AT REST.    011140
C
C
0002      COMMON
C          HI,GI,Z11,Z21,Z22,IZ,CHISO,CHISQO,            011150
C          ZI,IX,J,JX(11),KX(5),L,LX(11,2),M,MA,MX,N,NX,KKX,MP(11,2),MR(11,2), 011160
C          STITLE(20),ZIT,EPS,E1,E2,TEFF,H(11),AG(11),SPIN(11),RP(11),      011170
C          4G(11,2),CS(11,2),X(200),DLX(200),ES(11,2),EFF(11,2,200),      011180
C          5FP(5),DLFP(5),TC(5),TB(5),ZH(5),EE(512),Y(512),ELY(512),      011190
C          6XN(6),XR(6),TMX,AL(10),D(10),CT(10,2),CG(10,2),M1,M2,M0,      011200
C          7OLE,EMN,ST(2048),DLST(2048),DST(20),SG(2048),DLSG(2048),DSG(20), 011210
C          8SC(2048),T(2048),YY(2048),DYY(2048,20),ZZ(512),Z(512),DZ(20),EFC, 011220
C          9F(201),RF(201)
C          COMMON/SEYCO/YS(2048)                                011230
0003      COMMON/SEYM/FS(2048),PG(2048),PN(2048),FX,AA    011240
0004      COMPLEX HI,GI,Z11,Z21,Z22,DZ11,DZ21,DZ22        011250
0005      DIMENSION U(5),W(5)                               011260
0006      DATA U/.013046736,.067468317,.160295215,.28330230,.42556283/    011270
0007      DATA W/.033335672,.074725675,.109543181,.13463336,.14776211/    011280
C
C          SECONDARY CAPTURE YIELD (INFINITE SLAB APPROXIMATION)
C
0009      ISDT=1                                           011290
0010      FX=C.                                           011300
0011      AA=(1.+AG(1))**2                                011310
0012      E=EMN-DLE                                       011320
0013      DO 1 K=1,KKX                                    011330
0014      E=E+OLE                                         011340
0015      S1=C.                                           011350
0016      S2=0.                                           011360
0017      GX=XN(N)*ST(K)                                  011370
0018      T2=EXP(-GX)                                     011380
C
C          INTEGRATION BY 10-POINT GAUSSIAN QUADRATURE
C
0019      DO 2 J=1,2                                       011390
0020      DO 2 I=1,5                                       011400
0021      XX=(1.-U(I))*FLOAT(J-1)+U(I)*FLOAT(2-J)        011410
0022      ANU=XN(N)                                       011420
0023      6 SQ=SQRT(AG(1)**2-1.+XX**2)                    011430
0024      EX=E*(XX+SQ)**2/4A                                011440
0025      WX=(XX+SQ)**2/AG(1)/SQ                          011450
0026      XL=(EX-E)/OLE+1.                                  011460
0027      L=XL                                             011470
0028      IF(L.GT.1)GO TO 4                                011480
0029      SST=ST(1)                                        011490
0030      SSG=SG(1)                                        011500
0031      GO TO 5                                          011510
0032      4 DIFF=AMOD(XL,1.)                                011520
0033      SST=ST(L)+DIFF*(ST(L+1)-ST(L))                  011530
0034      SSG=SG(L)+DIFF*(SG(L+1)-SG(L))                  011540
0035      5 G1=W*SSG/SST                                    011550
0036      TAD=SST/XX                                       011560
0037      IF(GX-XN(N)*TAD.EQ.0.)GO TO 10                 011570
0038      G2=(EXP(-ANU*TAD)-T2*EXP(-(ANU-XN(N))*TAD))/(GX-XN(N)*TAD)*G1    011580
0039      GO TO 71                                         011590
0040      10 G2=EXP(-ANU*TAD)                              011600

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0041	70	S1=S1+G1*W(I)	011700
0042		S2=S2+G2*W(I)	011710
0043		IF(ANU.EQ.3.)GO TO 2	011720
0044		ANU=0.	011730
0045		XX=-XX	011740
0046		GO TO 6	011750
0047	2	CONTINUE	011760
	C	SECONDARY CAPTURE YIELD	011770
0048		SNI=ST(K)-SC(K)	011780
0049		FS(K)=XN(N)*SNI /2.*((1.-T2)/GX*S1-S2)	011790
0050	1	IF(FS(K).GT.FX)FX=FS(K)	011800
0051		I=ISDT	011810
0052		RETURN	011820
0053		END	011830

```

C
C
0001      SUBROUTINE MUY                                011840
C
C          MUY YIELDS THE MULTIPLE-COLLISION CONTRIBUTION TO THE
C          CAPTURE YIELD. THERMAL MOTION OF THE TARGET NUCLEI IS
C          NEGLECTED. THE SAMPLE IS TAKEN AS A CYLINDRICAL DISK.
C          THE INCIDENT BEAM IS ASSUMED COAXIAL WITH THE SAMPLE,
C          OF CONSTANT FLUX DENSITY AND WITH THE RADIUS EQUAL TO
C          EDGE*XR(N) (XR(N) IS THE SAMPLE RADIUS).
C
0002      COMMON
C          1HI,G1,Z11,Z21,Z22,IZ,CHISQ,CHISQ0,
C          2I,IX,J,JX(11),KX(5),L,LX(11,2),M,MA,MX,N,NX,KKX,MP(11,2),MR(11,2),
C          3TITLE(20),ZIT,EPS,E1,E2,TEFF,H(11),AG(11),SPIN(11),FP(11),
C          4G(11,2),CS(11,2),X(200),DLX(200),ES(11,2),EFF(11,2,200),
C          5FP(5),DLFP(5),TC(5),TB(5),ZH(5),EE(512),Y(512),DLY(512),
C          6XN(6),XR(6),TMX,AL(10),D(10),CT(10,2),CG(10,2),M1,M2,MD,
C          7DLE,EMN,ST(2048),DLST(2048),DST(20),SG(2048),DLSG(2048),DSG(20),
C          8SC(2048),T(2048),YY(2048),DYY(2048,20),ZZ(512),Z(512),DZ(20),EFC,
C          9F(201),RF(201)
C          COMMON/SEYCC/YS(2048)
0003      COMMON/SEYM/FS(2048),PG(2048),PN(2048),FX,AA
0004      COMPLEX HI,G1,Z11,Z21,Z22,DZ11,DZ21,DZ22
0005      DIMENSION NC(2048)
0006      DIMENSION EM(500),WG(500),WN(500),U(500),V(500),W(500),
0007      1XM(500),YM(500),ZM(500),SK(500)
0008      DATA EDGE/.5/
C
C          JSPIN=J
0009      CALCULATE NEEDED FUNCTIONS
C
0010      DO 11 K=1,KKX
0011      YS(K)=C.
0012      PG(K)=SG(K)/ST(K)
0013      11 PN(K)=1.-SC(K)/ST(K)
C          BEGIN ENERGY LOOP
0014      E=EMH-DLE
0015      DO 50 KE=1,KKX
0016      E=E+DLE
0017      WN)=(1.-T(KE))*PN(KE)
C          INITIAL NUMBER OF HISTORIES (IMPORTANCE SAMPLING)
0018      NH=ZH(N)*FS(KE)/FX+.5
0019      NHX=ZH(N)
0020      IF(NH.GT.NHX)NH=NHX
0021      IF(NH.LT. 10)NH=10
0022      CK=2.*DLFP(1)/FP(1)*E/DLE/XN(1)
0023      DO 13 J=1,NH
0024      EM(J)=E
0025      WN(J)=WN0
C          INITIAL DIRECTION;
0026      W(J)=1.
C          ADDITIONAL PATH LENGTH
0027      SK(J)=J.
C          COORDINATES OF FIRST COLLISION
C          (BEAM RADIUS SMALLER THAN SAMPLE RADIUS)
0028      XM(J)=XR(N)*SQRT(RANDOM(D))*EDGE
0029      YM(J)=C.
0030      13 ZM(J)=-ALOG(1.-WN0/PN(KE)*RANDOM(D))/ST(KE)

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0031          FK0=YY(KL)          012410
0032          SQM=J.              012420
                                012430
C          BEGIN COLLISION LOOP 012440
0033          DO 14 K=1,21        012450
0034          YH=NH                012460
                                012470
C          BEGIN HISTORY LOOP   012480
0035          DO 15 J=1,NH        012490
                                012500
C          COSINE OF CMS SCATTERING ANGLE
0036          CTHC=2.*RANDOM(D)-1. 012510
0037          QU=1.+AG(1)**2+2.*AG(1)*CTHC 012520
                                012530
C          LAB SCATTERING ANGLE  012540
0038          CTH=(1.+AG(1)*CTHC)/SQRT(QU) 012550
0039          STH=SQRT(1.-CTH**2)  012560
                                012570
C          NEW LAB ENERGY       012580
0040          EM(J)=EM(J)*QU/AA    012590
                                012600
C          NEW TOTAL CROSS SECTION, SCATTERING AND CAPTURE PROBABILITY
0041          XL=(EM(J)-EMN)/DLEF+1. 012610
0042          IF(XL.GT.1.001)GO TO 16 012620
0043          SST=ST(L)            012630
0044          PPN=PN(L)            012640
0045          PPG=PG(L)            012650
0046          GO TO 19             012660
0047          L=XL                 012670
0048          DD=XL-FLQAT(L)        012680
0049          SST=ST(L)+DD*(ST(L+1)-ST(L)) 012690
0050          PPN=PN(L)+DD*(PN(L+1)-PN(L)) 012700
0051          PPG=PG(L)+DD*(PG(L+1)-PG(L)) 012710
                                012720
C          AZIMUTH                012730
0052          19 PHI=6.28318*RANDOM(D) 012740
0053          CPH=COS(PHI)          012750
0054          SPH=SIN(PHI)         012760
                                012770
C          NEW DIRECTION COSINES  012780
0055          IF(W(J)**2.LT.0.999)GO TO 21 012790
0056          U(J)=STH*CPH*W(J)    012800
0057          V(J)=STH*SPH*W(J)    012810
0058          W(J)=CTH*W(J)        012820
0059          GO TO 22             012830
0060          21 RHO=SQRT(1.-W(J)**2) 012840
0061          UNEW=CTH*U(J)+STH*(W(J)*U(J)*CPH-V(J)*SPH)/RHO 012850
0062          V(J)=CTH*V(J)+STH*(V(J)*W(J)*CPH+U(J)*SPH)/RHO 012860
0063          W(J)=CTH*W(J)-STH*CPH*RHO 012870
0064          U(J)=UNEW             012880
                                012890
C          DISTANCE TO SAMPLE SURFACE 012900
C          DC: DISTANCE TO CYLINDER 012910
C          DP: DISTANCE TO PLANE   012920
0065          22 DC=2.*XR(N)        012930
0066          DP=DC                012940
0067          IF(W(J).LT.-0.999)DP=-ZM(J)/W(J) 012950
0068          IF(W(J).GT.+0.999)DP=(XN(N)-ZM(J))/W(J) 012960
0069          IF(W(J).LT.-0.9999.OR.W(J).GT.+0.9999)GO TO 23 012970
0070          B1=U(J)**2+V(J)**2     012980
0071          B2=U(J)*XM(J)+V(J)*YM(J) 012990
0072          B3=XR(N)**2-XM(J)**2-YM(J)**2 013000
0073          IF(B2.GE.0.)DC=B3/(B2+SQRT(B2**2+B1+B3)) 013010
0074          IF(B2.LT.0.)DC=(SQRT(B2**2+B1+B3)-B2)/B1 013020
0075          23 DIS=AMIN1(DC,DP)    013030
0076          DS=AMIN1(DIS*SST,67.) 013040
0077          IF(DS.LT.1.E-07) DS=1.E-07 013050

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0078      C      INTERACTING FRACTION      01299.
0079      WI=1.-EXP(-DS)      01300.
0080      WG(J)=WN(J)*WI*PPG      01301.
0081      WN(J)=WN(J)-WI*PPG      01302.
0081      C      KY: SUBSCRIPT OF GRID POINT WHERE YS IS TO BE ADJUD TO YY      01303.
0081      KY=K2-INT(CK-(SK(J)+(1.+DS-DS/WI)/SST)+.5)      01304.
0082      C      SMOOTHING OF MONTE CARLO CONTRIBUTION:      01305.
0082      DELS=1.      01306.
0083      IF(DS.LT.*.5.)DELS=DS*.2887      01307.
0084      K1=CK+DELS/SST+.5      01308.
0085      IF(K1.LT.3)K1=3      01309.
0086      WGY=WG(J)/(YH*FLNAT(2*K1+1))      01310.
0087      K2=KY+K1      01311.
0088      K1=KY-K1      01312.
0089      IF(K2.LT.1.OR.K1.GT.KKX)GO TO 15      01313.
0090      IF(K1.LT.1)K1=1      01314.
0091      IF(K2.GT.KKX)K2=KKX      01315.
0092      DO 12 KY=K1,K2      01316.
0093      12 YS(KY)=YS(KY)+WGY      01317.
0094      C      COORDINATES OF NEXT COLLISION      01318.
0095      S=-ALOG(1.-WI*RANDOM(D))/SST      01319.
0096      XM(J)=XM(J)+S*U(J)      01320.
0097      YM(J)=YM(J)+S*V(J)      01321.
0098      ZM(J)=ZM(J)+S*W(J)      01322.
0099      D=SQR(XM(J)**2+YM(J)**2)      01323.
0099      IF(D).LE.XR(N)GO TO 24      01324.
0100      C      CORRECT POSSIBLE TRUNCATION ERRORS      01325.
0100      XM(J)=XM(J)/D*.99999      01326.
0101      YM(J)=YM(J)/D*.99999      01327.
0102      IF(ZM(J).GT.XN(N))ZM(J)=*.99999*XN(N)      01328.
0103      24 SK(J)=SK(J)+S      01329.
0104      15 CONTINUE      01330.
0105      C      END OF HISTORY LOOP      01331.
0106      SUM=1.      01332.
0106      DO 29 J=1,NH      01333.
0107      29 SUM=SUM+WG(J)      01334.
0108      C      AVERAGE CAPTURE PROBABILITY      01335.
0108      FK=SUM/YH      01336.
0109      SOM=SOM+FK      01337.
0110      NC(KE)=K+1      01338.
0111      C      NEW NUMBER OF HISTORIES      01339.
0112      NNH=Z(N)*FK*.2/FK0/FX+.5      01340.
0113      IF(NNH.LT.NH)NNH=NNH      01341.
0114      IF(NNH.LT.1)GO TO 51      01342.
0115      FK0=FK      01343.
0116      14 CONTINUE      01344.
0116      C      END OF COLLISION LOOP      01345.
0116      52 CONTINUE      01346.
0117      C      END OF ENERGY LOOP      01347.
0117      DO 3 K=1,KKX      01348.
0118      3 YY(K)=YY(K)+YS(K)      01349.
0119      J=JSPIN      01350.
0120      RETURN      01351.
0121      END      01352.

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C
C
0001      FUNCTION RANDOM(KRRR)
C
C          (PSEUDO-)RANDOM NUMBER GENERATOR.
C
0002      DATA IY/32767/
0003      1 IX=IY
0004      CALL RANDU(IX,IY,RANDOM)
0005      IF(RANDOM.GT.2.)RETURN
0006      IY=32767
0007      GO TO 1
0008      END

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C
C
0001      SUBROUTINE CHISQ4(X,Y,K)
C
C          CHISQ4 YIELDS VALUES OF A CHI-SQUARED DISTRIBUTION WITH N
C          DEGREES OF FREEDOM
C          X: INCREMENT OF INDEPENDENT VARIABLE DIVIDED BY
C             SQR(2*VARIANCE)
C          Y: DEPENDENT VARIABLE MULTIPLIED BY INCREMENT
C          K: SUBSCRIPT AT UPPER LIMIT (SUBSCRIPT AT CENTER OF GRAVITY
C             IS 101)
C          A: N/2
C          B: SQR(N)*X
C          C: EXP(-N/2)/GAMMA(N/2)*B
C
0002      DIMENSION Y(201)
0003      A=2.
0004      B=2.*X
0005      C=.1353353*B
0006      D=EXP(B)
0007      E=1./D
C
C          GO FROM CENTER OF GRAVITY TOWARDS SMALLER J
0008      XI=A-B
0009      F=C/E
0010      DO 1 J=1,J,K
0011      XI=XI+B
0012      F=F*E
0013      1 Y(J)=F*XI**(A-1.)
C          GO FROM CENTER OF GRAVITY TOWARDS LARGER J UNTIL DISTRIBUTION
C          DROPS TO ZERO
0014      I1=A/B
0015      XI=A
0016      F=C
0017      DO 2 I=1,I1
0018      J=101-I
0019      XI=XI-B
0020      F=F*D
0021      2 Y(J)=F*XI**(A-1.)
C          COMPLETE BY FILLING WITH ZEROS
0022      I1=I1+1
0023      I2=K-101
0024      DO 3 I=I1,I2
0025      J=101-I
0026      3 Y(J)=0.
0027      RETURN
0028      END

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C
C
0001      SUPROUTINE TGAUSS(X,Y,K)
C
C      TGAUSS YIELDS VALUES OF A TRUNCATED GAUSSIAN, NORMALIZED TO 1
C      IF TRUNCATED AT TWICE THE 1/E-WIDTH (K*X=2)
C      X: INCREMENT OF INDEPENDENT VARIABLE DIVIDED BY 1/E-WIDTH,
C      Y: DEPENDENT VARIABLE TIMES INCREMENT,
C      K: SUBSCRIPT AT UPPER LIMIT (SUBSCRIPT AT MEDIAN IS 101)
C
0002      DIMENSION Y(201)
0003      Y(101)=0.566342*X
0004      F1=EXP(-X*X)
0005      F2=F1*F1
0006      DO 1 I=102,K
0007      Y(I)=Y(I-1)*F1
0008      J=202-I
0009      Y(J)=Y(I)
0010      1 F1=F1*F2
0011      RETURN
0012      END

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C
C
0001      SUBROUTINE SIMP(Y,M,N,SUM)
C
C      SIMPSON'S RULE:
C
0002      DIMENSION Y(201)
0003      SUM=0.
0004      K1=M+1
0005      K2=N-1
0006      DO 1 K=K1,K2,2
0007      1 SUM=SUM+Y(K-1)+4.*Y(K)+Y(K+1)
0008      SUM=SUM/3.
0009      RETURN
0010      END

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C
C
0001      C      SUBROUTINE CHOBAN(MX,A,B)
C
C      MATRIX INVERSION BY THE "SQUARE-ROOT" METHOD OF
C      CHOLESKI AND BANACHIEWICZ,
C      MX: RANK
C      A: ORIGINAL MATRIX
C      B: INVERTED MATRIX
C
C      DOUBLE PRECISION A(20,20),B(20,20),U(20,20),SUM
C      CONSTRUCT "SQUARE-ROOT" MATRIX U
0003      DO 2 N=1,MX
0004      DO 2 M=N,MX
0005      SUM=0.
0006      DO 1 K=1,N
0007      1 IF(K.LT.N)SUM=SUM+U(M,K)*U(N,K)
0008      IF(M.EQ.N)U(N,N)=DSQRT(A(N,N)-SUM)
0009      2 IF(M.GT.N)U(M,N)=(A(M,N)-SUM)/U(N,N)
C      CALCULATE INVERSE B
0010      DO 4 NN=1,MX
0011      N=MX-NN+1
0012      DO 4 MM=NN,MX
0013      M=MX-MM+1
0014      SUM=0.
0015      DO 3 K=M,MX
0016      3 IF(K.GT.M)SUM=SUM-U(K,M)*B(K,N)
0017      IF(M.EQ.N)SUM=SUM+1./U(M,M)
0018      B(M,N)=SUM/U(M,M)
0019      4 IF(M.NE.N)B(N,M)=B(M,N)
0020      RETURN
0021      END

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C
C
0001      SUBROUTINE ADJ(IZ,MX,MA,K,CHISQ,X,DLX,B,C)
C
C          CALCULATION OF IMPROVED CROSS SECTION PARAMETERS BY
C          MULTIPLICATION OF COVARIANCE MATRIX B INTO DEVIATION VECTOR C.
C
0002      DIMENSION X(200),DLX(200),C(20)
0003      DOUBLE PRECISION B(20,20)
0004      DATA FUDGE/.75/
C          CALCULATE GAUSS' ERROR ADJUSTMENT FACTOR
0005      CF=SQRT(CHISQ/FLD(4*(K-MA)))
0006      WRITE(6,100)CHISQ,CF,IZ
0007      100 FORMAT(1H //' CHI SQAED:           ',1PE10.3/
1          ' ERROR ADJUSTMENT FACTOR:',1PE10.3/
2          ' ITERATION STEP:           ',1I0//)
0008      M=0
0009      DO 2 MM=1,MX
0010      IF(DLX(MM).LE.0.)GO TO 2
0011      M=M+1
0012      XM=C.
0013      DO 1 MN=1,MA
0014      1 XM=XM+B(M,MN)*C(MN)
0015      XM=XM*FUDGE
0016      X(MM)=X(MM)+XM
0017      DLX(MM)=DSQRT(B(M,M))
0018      DLX(MM)=DLX(MM)*CF
0019      2 CONTINUE
0020      RETURN
0021      END

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C
C
0001      SUBROUTINE KEV                                015060
C
C          KEV CONVERTS ENERGIES FROM MEV TO KEV      015070
C
C
0002      COMMON                                        015080
C          CHI,GI,Z11,Z21,Z22,IZ,CHISO,CHISCO,      015090
C          ZI,IX,J,JX(11),KX(5),L,LX(11,2),M,MA,    015100
C          MX,N,NX,KKX,MP(11,2),MR(11,2),          015110
C          3TITLE(20),ZIT,EPS,E1,E2,TEFF,H(11),AG(11),SP1N(11),RP(11),
C          4G(11,2),CS(11,2),X(20),DLX(200),ES(11,2),EFF(11,2,200),
C          5FP(5),DLFP(5),TC(5),TB(5),ZH(5),EE(512),Y(512),DLY(512),
C          6XN(6),XR(6),TMX,AL(10),D(10),CT(10,2),CG(10,2),M1,M2,MO,
C          7DLE,EMN,ST(2048),DLST(2048),DST(20),SG(2048),DLSG(2048),DSG(20),
C          8SC(2048),T(2048),YY(2048),DYY(2048,20),ZZ(512),Z(512),DZ(20),EFC,
C          9F(201),RF(201)                            015170
C
0003      COMPLEX HI,GI,Z11,Z21,Z22,DZ11,DZ21,CZ22   015180
C
0004      DO 1 I=1,IX                                  015190
0005      JH=JX(I)                                     015200
0006      DO 1 J=1,JH                                  015210
0007      ES(I,J)=ES(I,J)*1000.                      015220
0008      IF(LX(I,J).EQ.0)GO TO 1                     015230
0009      M1=MR(I,J)                                  015240
0010      M2=MR(I,J)+4*LX(I,J)-4                     015250
0011      DO 1 M=M1,M2,4                              015260
0012      X(M)=X(M)*1000.                             015270
0013      X(M+1)=X(M+1)*ABS(X(M+1))*1000.            015280
0014      X(M+2)=X(M+2)*ABS(X(M+2))*1000.            015290
0015      X(M+3)=X(M+3)*1000.                       015300
0016      IF(DLX(M).GT.0.)DLX(M)=DLX(M)*1000.       015310
0017      IF(DLX(M+1).GT.0.)DLX(M+1)=DLX(M+1)*SQRT(ABS(X(M+1))*0.01)*2000.
0018      IF(DLX(M+2).GT.0.)DLX(M+2)=DLX(M+2)*SQRT(ABS(X(M+2))*0.01)*2000.
0019      IF(DLX(M+3).GT.0.)DLX(M+3)=DLX(M+3)*1000.
0020      1 CONTINUE                                  015330
0021      KH=KX(NX)                                   015340
0022      DO 2 K=1,KH                                  015350
0023      2 EE(K)=EE(K)*1000.                         015360
0024      EMN=1000.*EMN                               015370
0025      DLE=1000.*DLE                               015380
0026      RETURN                                       015390
0027      END                                         015400

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C
C
0001 C SUBROUTINE YPLOT 015470
C
C YPLOT PREPARES PLOTS OF MEASURED AND CALCULATED CAPTURE YIELDS 015480
C AND, FOR IZ=J, A CROSS SECTION PLOT FOR THE PURE 1ST ISOTOPE. 015490
C 015500
C 015510
C 015520
C PLOTA(X,Y,N,NT,NP,NH,I,NS,NR,XMAX,XMIN,SX,YMAX,YMIN,SY,TEXT,ID) 015530
C IS A KFK PLOT SUBROUTINE WITH THE FOLLOWING ARGUMENTS: 015540
C 015550
C 1. X ARRAY OF ABSCISSAE 015560
C 2. Y ARRAY OF ORDINATES 015570
C 3. N NUMBER OF CO-ORDINATE PAIRS 015580
C 4. NT=1 PLOT POINT SYMBOLS 015590
C =2 DRAW LINE 015600
C =3 DRAW LINE WITH POINT SYMBOLS 015610
C 5. NP CHOOSE NP-TH POINT SYMBOL (FROM A LIST) IF NT=1 OR 3 015620
C 6. NH=1 HEIGHT OF POINT SYMBOL 0.12 IN. 015630
C =2 HEIGHT OF POINT SYMBOL 0.16 IN. 015640
C =3 HEIGHT OF POINT SYMBOL 0.24 IN. 015650
C 7. I=1 LINEAR INTERPOLATION (FOR NT=2 OR 3) 015660
C =2 QUADRATIC "" (FOR NT=2 OR 3) 015670
C =3 CUBIC "" (FOR NT=2 OR 3) 015680
C 8. NS SPACING: EVERY NS-TH POINT IS TO BE MARKED 015690
C (FOR NT=3) 015700
C 9. NR=0 DRAW (NTD EXISTING PLOT 015710
C =1 BEGIN NEW PLOT, (XMAX-XMIN)/(YMAX-YMIN) = 1 015720
C =2 --- = 2 015730
C =3 --- = 3 015740
C =4 --- = 4 015750
C >=5 --- = 1.5 015760
C (NH,XMAX,XMIN,SX,YMAX,YMIN,SY,TEXT NEED NOT BE 015770
C SPECIFIED FOR NR=0) 015780
C 10. XMAX MAXIMAL ABSCISSA 015790
C 11. XMIN MINIMAL ABSCISSA 015800
C 12. SX X-INCREMENT CORRESPONDING TO 0.01 IN. 015810
C 13. YMAX MAXIMAL ORDINATE 015820
C 14. YMIN MINIMAL ORDINATE 015830
C 15. SY Y-INCREMENT CORRESPONDING TO 0.01 IN. 015840
C 16. TEXT FIGURE CAPTION, 60 ALPHAMERIC CHARACTERS 015850
C 17. ID FIGURE NUMBER 015860
C 015870
0002 C COMMON 015880
C 015890
C IHI,G1,Z11,Z21,Z22,IZ,CHISQ,CHISQ0, 015900
C 2I,IX,J,JX(11),KX(5),L,LX(11,2),M,MA,MX,N,NX,KKX,MP(11,2),MR(11,2), 015910
C BTITLE(21),ZIT,EPS,E1,E2,TEFF,H(11),AG(11),SPIN(11),RP(11), 015920
C 4G(11,2),CS(11,2),X(21),DLX(200),ES(11,2),EFF(11,2,200), 015930
C 5FP(5),DLFP(5),TC(5),TB(5),ZH(5),EC(512),Y(512),OLY(512), 015940
C 6XN(6),XR(6),TMX,AL(10),D(10),GT(10,2),CG(10,2),M1,M2,MO, 015950
C 7OLE,EMN,ST(2048),DLST(2048),DST(20),SG(2048),DLSG(2048),DSG(20), 015960
C 8SC(2048),T(2048),YY(2048),DYY(2048,20),ZZ(512),Z(512),DZ(20),EFC, 015970
C 9F(201),RF(201) 015980
0003 C COMMON/I1/ SGI1(2048),DLSGI1(2048) 015990
0004 C DIMENSION EF(2),YF(2),EK(2048) 016000
0005 C COMPLEX HI,G1,Z11,Z21,Z22,DZ11,DZ21,DZ22 016010
C 016020
0006 C ENERGY RANGE OF PLOTS (SAME FOR 1ST AND 2ND PLOT): 016030
C KH=KX(1) 016040

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0007          EMIN=EE(1)                016040
0008          EMAX=EE(KH)                016050
0009          IF(NX.EC.1)GO TO 2        016060
0010          DO 1 N=2,NX                016070
0011          KN=KX(N-1)+1              016080
0012          KH=KX(N)                   016090
0013          EMIN=AMIN1(EMIN,EE(KN))    016100
0014          EMAX=AMAX1(EMAX,EE(KH))    016110
0015          1 CONTINUE                  016120
0016          2 EBER=(EMAX-EMIN)*1.001   016130
0017          ILG=ALOG10(EBER)           016140
0018          IF(EBER.LT.1.)ILG=ILG-1   016150
0019          DEK=10.0**ILG              016160
0020          EBER=EBER/DEK              016170
0021          BER=1.0                    016180
0022          IF(EBER.LE.5.)BER=5.       016190
0023          IF(EBER.LE.2.5)BER=2.5     016200
0024          IF(EBER.LE.1.25)BER=1.25  016210
0025          BER=BER*DEK                016220
0026          EMIN=FLOAT(INT(25.*EMIN/BER))*BER/25. 016230
0027          EMAX=EMIN+BER              016240
0028          SE=0.0005*BER              016250
C          CAPTURE YIELD RANGE OF 1ST PLOT: 016260
0029          YMAX=0.                    016270
0030          KH=0                       016280
0031          DO 4 N=1,NX                 016290
0032          KN=KH+1                    016300
0033          KH=KX(N)                   016310
0034          DO 3 KK=KN,KH              016320
0035          IF(Y(KK)+DLY(KK).GT.YMAX)YMAX=Y(KK)+DLY(KK) 016330
0036          3 CONTINUE                  016340
0037          4 CONTINUE                  016350
0038          YBER=YMAX*22./25.*1.2      016360
0039          ILG=ALOG10(YBER)           016370
0040          IF(YBER.LT.1.)ILG=ILG-1   016380
0041          DEK=10.0**ILG              016390
0042          YBER=YBER/DEK              016400
0043          BER=11.                    016410
0044          IF(YBER.LE.5.5)BER=5.5     016420
0045          IF(YBER.LE.4.4)BER=4.4     016430
0046          IF(YBER.LE.2.2)BER=2.2     016440
0047          IF(YBER.LE.1.1)BER=1.1     016450
0048          BER=BER*DEK                016460
0049          YMIN=-BER*3./22.           016470
0050          YMAX= BER                   016480
0051          SY=0.001*BER*25./22.       016490
0052          KH=0                       016500
0053          DO 8 N=1,NX                 016510
0054          KN=KH+1                    016520
0055          KH=KX(N)                   016530
C          NP: NUMBER OF POINTS TO BE PLOTTED 016540
0056          NP=KH-KN+1                 016550
C          PLOT CALCULATED YIELDS AS CURVE (1ST PLOT) 016560
0057          IF(N.EQ.1)CALL PLOTA(EE ,Z ,NP,2,0,1,3,0,2,EMAX,EMIN,SE, 016570
          IYMAX,YMIN,SY,TITLE,IZ)        016580
0058          IF(N.GT.1)CALL PLOTA(EE(KN),Z(KN),NP,2,0,1,3,0,0,EMAX,EMIN,SE, 016590
          IYMAX,YMIN,SY,TITLE,IZ)        016600
C          PLOT MEASURED POINTS (1ST PLOT) 016610

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0059      IF(N.LE.3)NS=N-1      016620
0060      IF(N.EJ.2)NS=4      016630
0061      IF(N.GE.4)NS=N+3      016640
0062      CALL PLCTA(EI(KN),Y(KN),NP,1,NS,1,0,0,0,EMAX,EMIN,SE,ZMAX,ZMIN,SZ,
      10,J)      016650
      C      016660
      C      PLOT ERROR BARS OF MEASURED POINTS (1ST PLOT)      016670
0063      DO 5 KK=KN,KH      016680
0064      EF(1)=EE(KK)      016690
0065      EF(2)=EE(KK)      016700
0066      YF(1)=Y(KK)+DLY(KK)      016710
0067      YF(2)=Y(KK)-DLY(KK)      016720
0068      IF(YF(2).LT.YMIN)YF(2)=YMIN      016730
0069      CALL PLCTA(EF,YF,2,2,0,1,1,0,0,EMAX,EMIN,SE,YMAX,YMIN,SY,0,0)      016740
0070      5 CONTINUE      016750
0071      8 CONTINUE      016760
0072      IF(ZIT.GE.1.)RETURN      016770
      C      PLOT CAPTURE CROSS SECTION OF PURE FIRST ISOTOPE (2ND PLOT)      016780
      C      CROSS SECTION RANGE OF 2ND PLOT:      016790
0073      EK(1)=EMN      016800
0074      DO 6 K=2,KKX      016810
0075      EK(K)=EK(K-1)+DLE      016820
0076      6 CONTINUE      016830
0077      K1=(EMIN-EK(1))/DLE+2.      016840
0078      K2=(EK(KKX)-EMAX)/DLE      016850
0079      K2=KKX-K2-1      016860
0080      IF(K1.LT. 1)K1= 1      016870
0081      IF(K2.GT.KKX)K2=KKX      016880
0082      NP=K2-K1+1      016890
0083      SMIN=0.      016900
0084      SMAX=0.      016910
0085      DO 7 K=K1,K2      016920
0086      IF(SG11(K).GT.SMAX)SMAX=SG11(K)      016930
0087      7 CONTINUE      016940
0088      SBER=SMAX*1.2      016950
0089      ILG=ALOG10(SBER)      016960
0090      IF(SBER.LT.1.)ILG=ILG-1      016970
0091      DEK=1).C**ILG      016980
0092      SBER=SBER/DEK      016990
0093      BER=12.5      017000
0094      IF(SBER.LE.6.25)BER=6.25      017010
0095      IF(SBER.LE.5. )BER=5.      017020
0096      IF(SBER.LE.2.5 )BER=2.5      017030
0097      IF(SBER.LE.1.25)BER=1.25      017040
0098      SMAX=BER*DEK      017050
0099      SS=0.001*SMAX      017060
0100      CALL PLCTA(EK(K1),SG11(K1),NP,2,0,1,3,0,2,EMAX,EMIN,SE,
      ISMAX,SMIN,SS,TITLE,IZ)      017070
0101      RETURN      017080
0102      END      017090

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