

SIMPLE MODEL FOR THE CALCULATION OF THE COEFFICIENT OF SELF-DIFFUSION IN A LIQUID

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The coefficient of self-diffusion in three-dimensional classical liquid is computed approximately from the hierarchy of kinetic equations for the time-correlation functions (TCF).

The coefficient of self-diffusion D_s of a molecule can be obtained from the TCF $\pi(t) = \langle \mathbf{p}_1(0) \mathbf{p}_1(t) \rangle / \langle \mathbf{p}_1^2 \rangle$ of momentum $\mathbf{p}_1(t)$ at time t of a chosen particle in a fluid by $D_s = (kT/m) \lim_{z \rightarrow +0} \tilde{\pi}(z)$, where $\tilde{\pi}(z) = \int_0^\infty dt \cdot \exp(-zt) \pi(t)$, where m is the mass of a particle and kT the thermal energy. Here we shall obtain the formula for D_s from the small- z behavior of the hierarchy of the equations for TCF.

Zwanzig was the first to succeed in deriving a kinetic equation from the Liouville equation [1]: $\pi'(t) = - \int_0^t d\tau K(\tau) \pi(t - \tau)$, where the memory function is

$$K(\tau) = \langle \mathbf{p}_1 \hat{\mathcal{L}} \exp\{i \mathcal{L}_{22}^{(1)} \tau\} \hat{\mathcal{L}} \mathbf{p}_1 \rangle / \langle \mathbf{p}_1^2 \rangle, \quad \mathcal{L}_{22}^{(1)} = P \hat{\mathcal{L}} P, \quad \hat{\mathcal{L}} = -i \hat{L}, \quad \hat{\mathcal{L}} = \sum_{j=1}^N \frac{1}{m} \mathbf{p}_j \nabla_j - \sum_{i \neq j=1}^N \nabla_j u(j, i) \nabla_{\mathbf{p}_j},$$

$$P = 1 - \Pi, \quad \Pi = \frac{\mathbf{p}_1 \langle \mathbf{p}_1 \rangle}{\langle \mathbf{p}_1^2 \rangle},$$

Π is the projection operator, \hat{L} is the ordinary Liouville operator and $u(i, j)$ is the interparticle potential. If we use the identity [2]

$$\exp\{t(A+B)\} = \exp(tA) + \int_0^t du \exp\{(t-u)A\} B \exp\{u(A+B)\},$$

for the arbitrary operators A and B , we can obtain an expansion of $K(t)$:

$$K(t) = \omega^2 f(t) + \sum_{n=1}^{\infty} (-1)^{n+1} \int_0^t dt_1 \int_0^{t_1} dt_2 \dots \int_0^{t_{n-1}} dt_n \pi'(t-t_1) \pi'(t_1-t_2) \dots \pi'(t_{n-1}-t_n) \pi''(t_n),$$

where $\omega^2 = \langle F_1^2 \rangle / \langle \mathbf{p}_1^2 \rangle$, $f(t) = \langle F_1 \exp(i \hat{\mathcal{L}} t) F_1 \rangle / \langle F_1^2 \rangle$, $F_1 = - \sum_{j>1}^N \nabla_1 u(1, j)$ is the total force on a chosen particle.

Now one can write the Laplace transform of $K(t)$ in the form

$$\tilde{K}(z) = \omega^2 \tilde{f}(z) + z \sum_{n=1}^{\infty} \{1 - z \tilde{\pi}(z)\}^{n+1}. \quad (1)$$

Then for $z \rightarrow +0$ we obtain $\tilde{K}(z) \approx \omega^2 \tilde{f}(z)$ and $\tilde{\pi}(z) \approx \{z + \omega^2 \tilde{f}(z)\}^{-1}$. Following Zwanzig [1] we may write the kinetic equation for the TCF $f(t)$: $f'(t) = - \int_0^t d\tau \cdot V(\tau) f(t - \tau)$, where

$$V(\tau) = \langle F_1 \hat{\mathcal{L}} \exp\{i \mathcal{L}_{22}^{(2)} \tau\} \hat{\mathcal{L}} F_1 \rangle / \langle F_1^2 \rangle, \quad \mathcal{L}_{22}^{(2)} = Q \hat{\mathcal{L}} Q, \quad Q = 1 - R, \quad R = \frac{F_1 \langle F_1 \rangle}{\langle F_1^2 \rangle},$$