

Spatiotemporal coding of inputs for a system of globally coupled phase oscillators

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We investigate the spatiotemporal coding of low amplitude inputs to a simple system of globally coupled phase oscillators with coupling function $g(\phi) = -\sin(\phi + \alpha) + r \sin(2\phi + \beta)$ that has robust heteroclinic cycles (slow switching between cluster states). The inputs correspond to detuning of the oscillators. It was recently noted that globally coupled phase oscillators can encode their frequencies in the form of spatiotemporal codes of a sequence of cluster states [P. Ashwin, G. Orosz, J. Wordsworth, and S. Townley, *SIAM J. Appl. Dyn. Syst.* **6**, 728 (2007)]. Concentrating on the case of $N=5$ oscillators we show in detail how the spatiotemporal coding can be used to resolve all of the information that relates the individual inputs to each other, providing that a long enough time series is considered. We investigate robustness to the addition of noise and find a remarkable stability, especially of the temporal coding, to the addition of noise even for noise of a comparable magnitude to the inputs.

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I. INTRODUCTION

Coupled phase oscillators provide a simple mechanism for modeling collective activity such as synchronization between clusters of neurons (see, for example, [1]). By taking the simplest dynamical model of an oscillator, they give tractable models for neural systems that can be derived rigorously in the limit of weak coupling, and which give insight into a very wide range of neural applications.

This paper investigates a simple system of this form with mean field (global) coupling of nearly identical phase oscillators, through a coupling function that robustly gives slow switching [2] in the case of identical oscillators—the “ideal” system. The dynamics of the ideal system consists of attractors that are robust heteroclinic attractors between cluster states, robust to perturbations in the coupling function [3].

The dynamics of heteroclinic networks gives examples of *winnerless competition models* that are proposed to be good models for certain neural systems [4,5], in particular, those that involve conversion of a steady multidimensional input into a persistent spatiotemporal code. This motivates the current paper; we investigate how small inputs to the system of coupled oscillators [6] result in a spatiotemporal coding of those inputs. The inputs are implemented as perturbations to the natural frequencies of the oscillators. The coding can be used to classify the input ordering (via the spatial coding, i.e., the ordering of visits to a number of cluster states) and the input magnitudes (via the temporal coding, namely, the residence times at each of the cluster states). We completely characterize the spatial coding in terms of its combinatorics and successfully predict magnitudes of inputs via the temporal coding. We also investigate the influence of noise on the resulting code and show (a) that the spatial coding is very resistant to noise as long as it has a lower magnitude than the inputs, and (b) that the entropy of the spatial coding increases in a very sharp manner as the noise magnitude increases beyond the input magnitude.

We consider a set of globally coupled phase oscillators of the form

$$\dot{\theta}_n = \omega_n + \frac{1}{N} \sum_{m=1}^N g(\theta_n - \theta_m) + \eta \omega_n(t), \quad n = 1, \dots, N, \quad (1)$$

where θ_n is the phase of the n th oscillator such that $\theta_n \in [0, 2\pi)$ and $g(\phi + 2\pi) = g(\phi)$ for all $\phi \in \mathbb{R}$. ω_n is the *natural frequency* of the n th oscillator, and $w_n(t)$ is *uncorrelated white noise* such that the associated random walk has a unit growth of variance per unit time. The quantity η is the *noise strength* and we include independent and identically distributed (i.i.d.) additive noise independently to each phase, with zero mean, such that the standard deviation scales as η per unit time.

We will restrict our discussion to coupling of the form

$$g(x) = -\sin(x + \alpha) + r \sin(2x + \beta), \quad (2)$$

where $\alpha = 1.7$, $\beta = -2.0$, and $r = 0.2$ [6], although qualitatively the dynamics are similar for an open set of nearby parameter values.

The particular form (2) includes all Fourier components up to the second harmonic and includes dynamics that are not robustly observable for fewer than this number of components. The coupling (1) was investigated in [6] where it was shown that for a system of five oscillators the only attractors for this coupling involve switching between a total of 30 different periodic orbits that have a clustering into two pairs of two oscillators and a singleton; we describe this as (2, 1, 2) clustering. More generally, in systems of $2k+1$ globally coupled phase oscillators, this coupling leads to dynamics in which the model becomes partially synchronized with oscillators in $(k, 1, k)$ clusters. These attractors consist of $\frac{(2k+1)!}{(k!)^2}$ unstable manifolds, each with k branches leading to different, symmetrical unstable $(k, 1, k)$ clusters, forming a robust heteroclinic network.

The coupling (2) is chosen to give nontrivial dynamical clustering behavior that is robust for such systems. The case $\beta = 0$ has been studied by others [7,8,2,3] and also exhibits robust heteroclinic attractors, although between two-cluster

TABLE I. An enumeration of the set of possible (2,1,2) cluster states for a fully symmetrical system with $N=5$ oscillators. The state $s_1=S_{33211}$ corresponds to the phases clustered as $(\phi_1+\xi, \phi_1+\xi, \phi_2+\xi, \phi_3+\xi, \phi_3+\xi)$.

$s_1=S_{33211}$	$s_2=S_{11233}$	$s_3=S_{31231}$	$s_4=S_{13213}$	$s_5=S_{31213}$
$s_6=S_{13231}$	$s_7=S_{13321}$	$s_8=S_{31123}$	$s_9=S_{13123}$	$s_{10}=S_{31321}$
$s_{11}=S_{33121}$	$s_{12}=S_{11323}$	$s_{13}=S_{11332}$	$s_{14}=S_{33112}$	$s_{15}=S_{31312}$
$s_{16}=S_{13132}$	$s_{17}=S_{13312}$	$s_{18}=S_{31132}$	$s_{19}=S_{21133}$	$s_{20}=S_{23311}$
$s_{21}=S_{23131}$	$s_{22}=S_{21313}$	$s_{23}=S_{21331}$	$s_{24}=S_{23113}$	$s_{25}=S_{32113}$
$s_{26}=S_{12331}$	$s_{27}=S_{12313}$	$s_{28}=S_{32131}$	$s_{29}=S_{12133}$	$s_{30}=S_{32311}$

states rather than three-cluster states. The precise dynamics that we consider is described in detail in the following section.

In Sec. II we examine the robust clustering dynamics of the proposed model in the ideal system of identical oscillators. We then investigate the effect of applying fixed inputs to the system by detuning individual oscillators. Inputs of this fashion force the clustering to change in a predictable manner as the system explores part of the heteroclinic network, such that one can completely predict the sequence of states that a given trajectory will pass through. There is a predictable relationship between the relative strength of the inputs and the time spent near specific cluster states. In this section we describe a method of interpreting output from our proposed model as a sequence of codes that represent the spatial properties of each cluster that the system spends time near (spatial coding) as well as a sequence representing the amount of time the system spends near each state (temporal coding). Combining these codes generates a spatiotemporal coding for the system that can be used to completely resolve information about relative differences between inputs.

In the case where the strength of each input is fixed and unequal the branch followed when leaving a cluster state is predetermined. In Sec. III we investigate this case where the network collapses to a network of two individual components, which both result in a periodic cycle of cluster states of length 6. After discussing this theory we justify our work by reference to numerical simulations. In Sec. III C we explore the inverse problem—determining the relative strengths of the inputs to the system from a given length of spatiotemporal coding. For an unknown uniform detuning we show how one can exactly determine the detuning by examining the spatiotemporal code.

In many neural systems noise plays an important role. In Sec. IV we consider the entropy of code generated by our model in the presence of both uniform detuning and random inputs in the form of noise. We observe that the system is highly stable in the presence of noise up to a magnitude equal to the magnitude of the inputs and characterize the amount of information from the spatial coding that is lost (entropy) due to the noise present in the system.

Finally, in Sec. V we briefly explore the dynamics for situations where there are $N>5$ oscillators. While there are many similarities to the simple case of $N=5$ oscillators, the network of cluster states grows exceedingly quickly, hence the problem becomes much more complex. We also discuss the surprising resilience that the proposed system shows to noise.

II. SPATIOTEMPORAL CODING OF INPUTS

We consider the smallest number of oscillators that can generate nontrivial robust switching between cluster state dynamics, namely, the case with five oscillators ($N=5$, $k=2$) [6]. For $N=4$ oscillators there can be robust switching, but the network will be trivial [8].

A. Network of cluster states

For the fully symmetric case with no input, there is a uniform natural frequency across all of the oscillators ($\omega_i = \omega_c$ for $n=1, \dots, N$) and no noise ($\eta=0$). This case is described in [6] as having an attractor composed of 30 saddle periodic orbits (states) in phase space, each of which has a one-dimensional unstable manifold whose branches limit to two other symmetrically related states. For no input, the system will approach this network of states, progressively staying longer near each state. In the presence of noise ($\eta>0$) the system will explore a random path on the network.

We refer to an l -cluster state of type (N_1, \dots, N_l) as a state where each oscillator has one of l different relative phases ϕ_1, \dots, ϕ_l , up to addition of the same constant to all phases. For each r there are N_r oscillators with phase ϕ_r , up to addition of a constant, such that $\sum_{r=1}^l N_r = N$ [6]. For example, the oscillator phases of a (2, 1, 2) cluster state are grouped into one of three values, of the form $(\phi_1, \phi_1, \phi_2, \phi_3, \phi_3)$ or some permutation thereof. For an examination of the existence and stability of general cluster states in such a system, see [9].

For the (2, 1, 2) clusters we can enumerate all 30 permutations as the states s_k , $k=1, \dots, 30$ listed in Table I. We write $s_k = S_{j_1^{k_1} \dots j_5^{k_5}}$ to denote the state where the l th oscillator is in the $j_l^{k_l}$ th cluster. For example, we say $\theta \in s_1 = S_{11233}$ if and only if $(\theta_1, \dots, \theta_5) = (\phi_1, \phi_1, \phi_2, \phi_3, \phi_3) + \xi(1, 1, 1, 1, 1)$ for some $\phi_1 < \phi_2 < \phi_3 < \phi_1 + 2\pi$.

As outlined in [6] each of the cluster states in Table I has a one-dimensional unstable manifold that connects to two other such states. We describe each close approach to a saddle cluster state as an *epoch* [10].

More precisely, if $s_i = S_{j_1^{i_1} \dots j_5^{i_5}}$ and $s_k = S_{j_1^{k_1} \dots j_5^{k_5}}$, then the unstable manifold of s_i connects to the state s_k on advancing the phase of the p th oscillator by a small amount, precisely $O_p(s_i) = s_k$, where

$$O_p(s_i) = s_k \Leftrightarrow \begin{cases} k = i & (\text{if } j_p^i \neq 3) \\ j_l^k = 3 & \text{if } j_l^i = 1 \quad (\text{if } j_p^i = 3) \\ j_l^k = 1 & \text{if } j_l^i = 2 \quad (\text{if } j_p^i = 3) \\ j_l^k = 1 & \text{if } j_l^i = 3 \text{ and } l \neq p \quad (\text{if } j_p^i = 3) \\ j_l^k = 2 & \text{if } j_l^i = 3 \text{ and } l = p \quad (\text{if } j_p^i = 3). \end{cases} \quad (3)$$

From the signs of the eigenvalues of the linearization we infer that perturbations that separate the oscillators in group 1 will decay over time. Consequently, we describe the cluster of oscillators in the group labeled by 1's as the *stable cluster*. Conversely, perturbations that separate the oscillators in group 3 will grow with time. For this reason, we describe the cluster of oscillators in group 3 as the *unstable cluster*. Investigating the dynamics around the other eigenvectors shows that there are no other perturbations that grow with time.

Using the function O_p the dynamics on the network of cluster states can be represented as a directed graph, where each cluster state is represented by a node, and each branch of its unstable manifold is represented by a directed edge.

Considering both the information about O_p and the states as in Table I, one can verify that the graph of cluster states collapses to two components, each of which has an eventual period 6 terminating cycle listed in Table I, as shown in [6].

B. Response of the network to inputs (detuning)

Let us write

$$\omega_n = \omega_0 + \delta_n,$$

where $\delta_n \ll 1$ and ω_0 are constants; we can choose $\delta_1 = 0$ without loss of generality. Describe the *amplitude of detuning* for oscillator n as δ_n . We describe the response of the network to differences in ω_n (steady inputs) by using the operator O_p . In particular, for any given state s_i we expect that s_k follows s_i precisely if $\omega_p > \omega_q$, where p, q are the two oscillators in the unstable cluster and $O_p(s_i) = s_k$.

Suppose that the system is at, or near, a cluster state s_i and its unstable cluster consists of the oscillators p, q . For convenience define the *detuning in the unstable cluster* to be

$$\delta_u = \omega_p - \omega_q. \quad (4)$$

If $\delta_u > 0$ then the next expected state will be $O_p(s_i)$, while if $\delta_u < 0$ the next expected state will be $O_q(s_i)$. Without loss of generality we can always assume that p, q in the unstable cluster are chosen such that $\delta_u \geq 0$.

Moreover, we expect from [6,11] that the time of residence $\tau(s_i)$ near a given state s_i is approximately given by

$$\tau(s_i) \approx -\frac{1}{\lambda} \ln(\delta_u) + K \quad (5)$$

at leading order $\delta_u \rightarrow 0$, where $\lambda > 0$ is the unstable eigenvalue corresponding to the rate of growth of any perturbations that break the unstable cluster, and K is a fixed time related to the minimum transition time along a connection between states.

Note that Eq. (5) implies that we can infer the detuning of the unstable cluster from the time of residence near that cluster

$$\delta_u \approx e^{-\lambda(\tau-K)}. \quad (6)$$

C. Coding by sequence of states and residence times

The sequence of cluster states, and the residence time near each cluster state, can be considered as a spatiotemporal coding sequence that represents the dynamical state of the system. We represent a trajectory near the heteroclinic network as a sequence that consists of a numeric index to reference the cluster state that the system was near during an epoch and the amount of time spent near that state—the *residence time* or the *epoch length*.

One can fairly easily determine whether a solution is near one of the cluster states in Table I. On a trajectory $(\theta_1(t), \dots, \theta_N(t))$ close to the network of cluster states s_i we choose a tolerance ξ such that $|s_k - s_j| > 2\xi$ for $k \neq j$ and then define the spatial state at time t as follows:

$$C(t) = \begin{cases} k: & |\theta(t) - s_k| < \xi \\ 0 & \text{otherwise.} \end{cases}$$

i.e., $C(t) = k$ if and only if the trajectory is close to s_k at time t . $C(t) = 0$ implies that the trajectory is not near any classifiable state; this usually means that it is in transit between two such states [12]. We define an increasing sequence of times t_n such that the n th epoch starts at time t_n if the following holds:

$$C(t_n) \neq 0,$$

$$C(t) \neq C(t_n) \quad \text{for } t < t_n \quad \text{with } |t_n - t| \text{ arbitrarily small,}$$

$$C(t) = C(t_n) \quad \text{for } t_n < t < t_n + T_{\min},$$

and there are no epochs between t_{n-1} and t_n , where $T_{\min} > 0$ is a fixed “minimum residence time.” This enables us to classify a trajectory as a discrete time sequence of states $\{C_n\}$ associated with the n th epoch by

$$C_n = C(t_n),$$

where because of the instability of each state, $C_n \neq C_{n-1}$. This sequence is simply the measured itinerary of states and will depend on the precise initial conditions, the parameters, and any input (detuning) or noise added to the system. We call this the *spatial coding* sequence of the trajectory. Note that the structure of the network in Fig. 1 will limit the possible sequences that can appear as a spatial coding.

In addition to the spatial coding sequence of epochs we define a temporal coding as a sequence of residence times by

$$T_n = \sup\{t > 0: C(t_n) = C(t_n + s) \text{ for all } 0 < s < t\}.$$

Note that $T_n > T_{\min}$ but the coding sequence is typically unaffected by T_{\min} as the value of T_{\min} is relatively small. T_n is the length of the time interval for which $C(t) = C_n$ for a given n . The *spatiotemporal code* is then the sequence

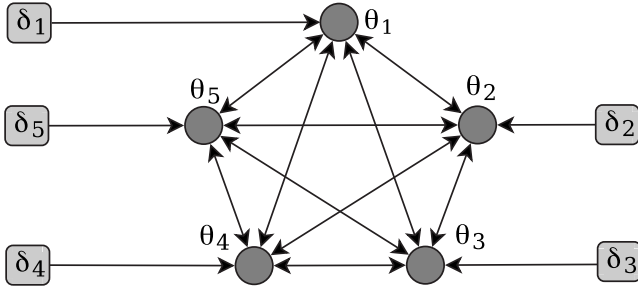


FIG. 1. A system of $N=5$ phase oscillators whose natural frequency is ω_i and phase is θ_i . The oscillators are coupled with mean-field (global) coupling where each oscillator receives one-dimensional input in the form of detuning its natural frequency. Each oscillator is coupled equally to all other oscillators by the coupling function $g(x)=-\sin(x+\alpha)+r \sin(2x+\beta)$.

$$\{(C_n, T_n): n = 1, 2, 3, \dots\}.$$

An example showing a trajectory and its spatiotemporal code is illustrated in Fig. 2.

III. DETUNING IN THE ABSENCE OF NOISE

The dynamics of the case where noise or impulses drive the system has been explored in some detail [6]; the case where we apply a vector of constant inputs to the system, specifically, by altering the natural frequency of each oscillator by a fixed amount, deserves further investigation. We expand upon [6] by considering not only the spatial coding generated by such a system, but also the temporal coding.

A. Uniform detuning

The case of uniform detuning

$$\omega_n = \omega_0 + (n - 1)\delta, \tag{7}$$

where δ is constant, was briefly considered in [6]. By examining the ordering of the unstable clusters, each state will have a preferred direction of switching (4) determined by the choice of p, q in the unstable cluster such that $\delta_u > 0$. Moreover, uniform detuning implies that

$$\delta_u = k\delta, \tag{8}$$

where $k \in \{1, 2, 3, 4\}$. This means that each node on the graph shown in Fig. 3 has only one outgoing directed edge and so all directed paths on the graph become eventually periodic after a possible initial transient. In fact, as shown in [6], the graph reduces to two components each of which has an eventual period 6 coding sequence. Moreover, the residence times are determined by δ_u which, by Eq. (8) means that the residence times can be used to determine δ .

Now consider the period 6 coding sequences that appear after the initial transient states have been passed in the presence of uniform detuning. We show how one can determine these sequences with a combinatorial method. After examining the possible sequences, one can verify that the only two periodic sequences for uniform detuning will be as listed in Table II.

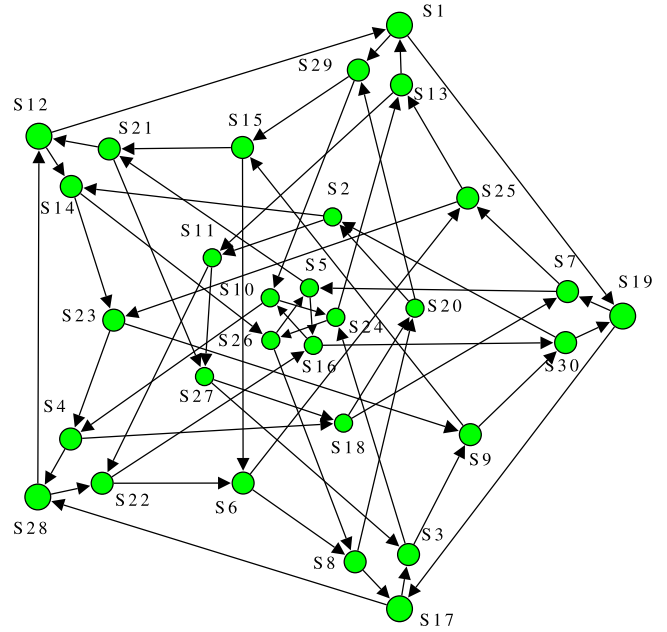


FIG. 2. (Color online) A graphical representation of the robust heteroclinic network generated by the system depicted in Fig. 1 where each node s_i ($i=1, \dots, 30$) represents a saddle point in phase space. Each directed edge leaving state s_i represents an individual branch of the unstable manifold where differences in the frequencies will drive the system along a particular edge from a given node.

The period 6 sequences shown in Table II are generated for uniform detuning (8) such that $\omega_1 < \omega_2 < \omega_3 < \omega_4 < \omega_5$ but any such detuning sequence can be considered to have the same effect by permutation.

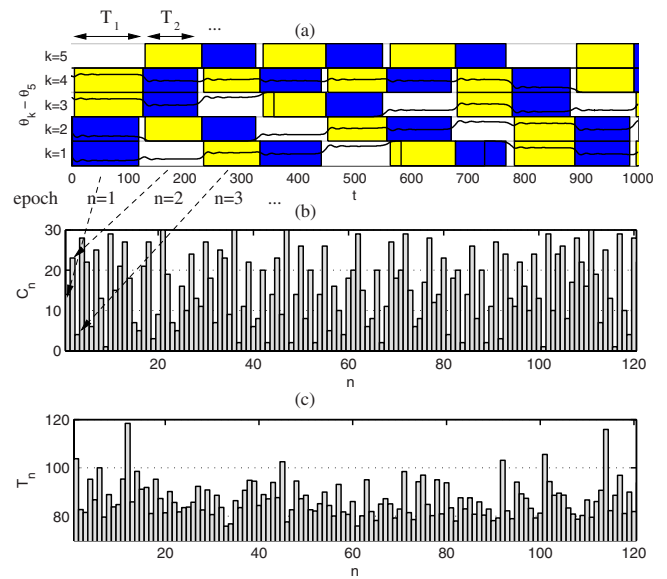


FIG. 3. (Color online) A spatiotemporal code generated by the trajectory where $\omega_{n+1}=\omega_n$ for $n=1, \dots, 5$ and $\eta=10^{-7}$. Panel (a) represents the trajectory being converted into a spatio-temporal code. Panel (b) shows the state sequence C_n over a set of epochs n , where C_n is a reference to state code. Panel (c) shows the residence time sequence T_n representing the time spent in a given epoch.

TABLE II. A combinatorial representation of the two possible period 6 cycles that the system will collapse to given uniform detuning (7), $\delta \neq 0$ and $\eta = 0$.

Epoch	j_1^k	j_2^k	j_3^k	j_4^k	j_5^k	C_n	Epoch	j_1^k	j_2^k	j_3^k	j_4^k	j_5^k	C_n
n	3	1	1	3	2	18	n	3	1	3	1	2	15
$n+1$	1	3	3	2	1	7	$n+1$	1	3	2	3	1	6
$n+2$	3	1	2	1	3	5	$n+2$	3	1	1	2	3	8
$n+3$	1	3	1	3	2	16	$n+3$	1	3	3	1	2	17
$n+4$	3	1	3	2	1	10	$n+4$	3	1	2	3	1	3
$n+5$	1	3	2	1	3	4	$n+5$	1	3	1	2	3	9

B. Numerical simulations

Consider a sequence of residence times for a given uniform detuning with an amplitude of detuning δ . From long numerical simulations using high resolution data integrated using a fourth order Runge-Kutta method followed by an Euler step for the noise, it can be seen that the residence time at each state is close to one of four discrete values; $T_n \in \{\alpha_1, \alpha_2, \alpha_3, \alpha_4\}$, where $\alpha_1 < \alpha_2 < \alpha_3 < \alpha_4$. More specifically, after an initial transient period of t duration 1000, T_n repeats every six epochs, but takes only four different values (see Fig. 4).

Consider now the values of $\{T_n\}$ for a range of values of δ , after the system has passed through the transient period. For each δ we see the same two possible coding sequences for C_n dependent on initial conditions. There is a logarithmic relationship between δ and the residence times as shown in Fig. 5.

C. Inverse problem—determining the detuning from a spatiotemporal code

For computational purposes it is necessary to understand the process by which we could determine the input to a system given a length of spatiotemporal code. We will consider the spatial and temporal coding sequence separately.

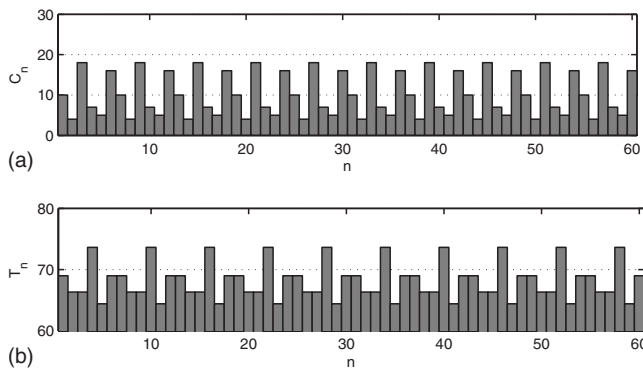


FIG. 4. The spatiotemporal code (C_n, T_n), as a function of epoch number n , from a numerical simulation with uniform detuning (7) for $\delta = 10^{-7}$ and $\eta = 0$ after an initial transient period has been allowed to decay. The data shows one of the two period 6 cycles shown in Table II. Notice that while the system passes through six different states, there are only four possible values for the residence times.

1. Decoding the spatial coding

As shown in Sec. III B, a uniform detuning will always result in a period 6 cycle in the spatial coding sequence. For the given state in the spatial coding C_n , consider the oscillators p, q in the unstable cluster such that $j_p^n = j_q^n = 3$. Note then, that C_{n+1} is equal to either $O_p(C_n)$ or $O_q(C_n)$. These two cases can be used to compare ω_p and ω_q as follows:

$$O_p(C_n) = C_{n+1} \Leftrightarrow \omega_p > \omega_q,$$

$$O_q(C_n) = C_{n+1} \Leftrightarrow \omega_q > \omega_p.$$

Given a number of transitions, every oscillator that arrives in an unstable cluster will be subject to a comparison. Consider a graph where each node represents an oscillator and edges between those nodes represent a comparison between those oscillators (both oscillators are in the unstable cluster at the time of a transition). This graph, shown in Fig. 6(a), is bipartite and has a minimum spanning tree of length 4 [see Fig. 6(b)]. This shows that we need to examine only four transitions to determine an ordering in which we can state that two oscillators form the “slow” group and three oscillators form the “fast” group. We call this the *slow-fast partition* and use this to order ω_i as follows:

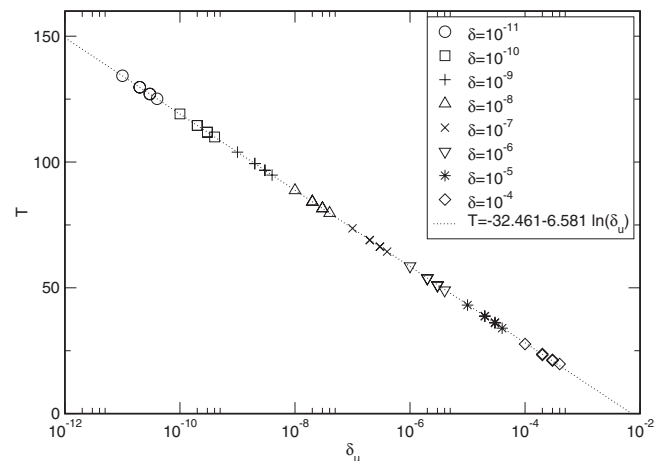


FIG. 5. A plot showing residence time T against δ_u . Slope fits to $y = A + B \ln(x)$, where $A = -32.461$ and $B = -6.581 \ln(x)$.

TABLE III. Decoding the spatiotemporal coding $\{C_n, T_n\}$ for a trajectory on the period 6 loop with $\eta = 0$. (a) shows an analysis of the data from one of the cycles shown in Table II (b) is from the simulation described in Fig. 7. At each epoch the detuning δ_u between the oscillators in the third column is compared; the switch is determined by which of these has greater frequency. The residence time can be used to estimate the size of δ_u by $\delta_u \approx \exp[(32.461 + T)/-6.58]$; this is shown in the fourth column. The final column gives the actual detuning; this is approached on taking averages over many visits to the same state.

Epoch	C_n	T_n	Compares	Inferred δ_u	Actual δ_u
(a)					
n	S_{31132}	66.3	$\omega_1 < \omega_4$	3.03×10^{-7}	3×10^{-7}
$n+1$	S_{13321}	73.7	$\omega_2 < \omega_3$	9.84×10^{-8}	1×10^{-7}
$n+2$	S_{31213}	64.4	$\omega_1 < \omega_5$	4.04×10^{-7}	4×10^{-7}
$n+3$	S_{13132}	69.0	$\omega_2 < \omega_4$	2.01×10^{-7}	2×10^{-7}
$n+4$	S_{31321}	69.0	$\omega_1 < \omega_3$	2.01×10^{-7}	2×10^{-7}
$n+5$	S_{13213}	66.3	$\omega_2 < \omega_5$	3.03×10^{-7}	3×10^{-7}
(b)					
n	S_{12331}	57.3	$\omega_3 < \omega_4$	1.19×10^{-6}	12×10^{-7}
$n+1$	S_{31123}	62.9	$\omega_1 < \omega_5$	5.08×10^{-7}	5×10^{-7}
$n+2$	S_{13312}	60.8	$\omega_2 < \omega_3$	6.99×10^{-7}	7×10^{-7}
$n+3$	S_{32131}	56.3	$\omega_1 < \omega_4$	1.39×10^{-6}	14×10^{-7}
$n+4$	S_{11323}	66.2	$\omega_3 < \omega_5$	3.08×10^{-7}	3×10^{-7}
$n+5$	S_{33112}	59.2	$\omega_1 < \omega_2$	8.92×10^{-7}	9×10^{-7}

Table III illustrates the process of recreating the input data from a given length of spatiotemporal coding. The values in the columns C_n and T_n are determined by inspection of the spatiotemporal coding. The fourth column shows which two oscillators are in the unstable cluster during this epoch, and hence, which two oscillators are compared to determine the length of the current epoch and the value of C_n for the next epoch. Now estimate the value of δ_u for the compared oscillators by $\delta_u \approx \exp[(32.461 + T)/-6.58]$. The final column shows the actual value of δ_u for comparison. The estimate of δ_u approaches the value of δ_u on taking averages over many visits to the same state.

IV. DETUNING IN THE PRESENCE OF NOISE

Adding noise to the system adds a level of uncertainty to the trajectory of the system. We evaluate how robust our model is for $\eta \neq 0$ and investigate how much information about the detuning can be extracted, dependent on the level of noise introduced to the system.

A. Influence of noise on the spatial coding

Consider the spatial coding sequence C_n and the residence time sequence T_n for a particular trajectory (Fig. 8). Let $p_i \in (p_1, \dots, p_{30})$ and $\sum_{i=1}^{30} p_i = 1$ be the relative frequency that, in a given epoch, the system is at state s_i ; i.e.,

$$p_i = \lim_{N \rightarrow \infty} \frac{\mathcal{N}\{1 \leq n \leq N : C_n = i\}}{N}.$$

Let $P_{i,j}$ be the relative frequency that the state sequence advances from state s_i to s_j defined as

$$P_{i,j} = \lim_{N \rightarrow \infty} \frac{\mathcal{N}\{n \leq N : c_n = i, c_{n+1} = j\}}{\mathcal{N}\{n \leq N : c_n = i\}}.$$

Then it can be verified that, for any i ,

$$\sum_{j=1}^{30} P_{i,j} = 1,$$

and P can be interpreted as a transition matrix.

Some examples of transition matrices from numerical simulations with constant uniform detuning and increasing noise are depicted in Fig. 9.

We observe that the spatial coding is approximately a Markov chain; the transitions from epoch n to $n+1$ are ap-

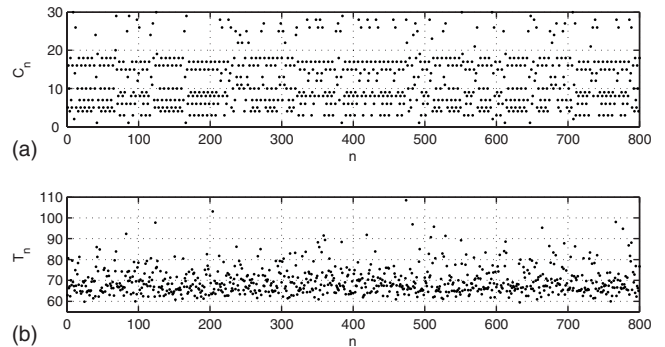


FIG. 8. A spatiotemporal code generated by a numerical simulation over 800 epochs as in Fig. 3 where there is noise present; $\eta/\delta=5.0$. Note there is an apparently random switching between the two possible cycles of length six, plus some occasional longer deviations. This is the activity we see in Fig. 9(c).

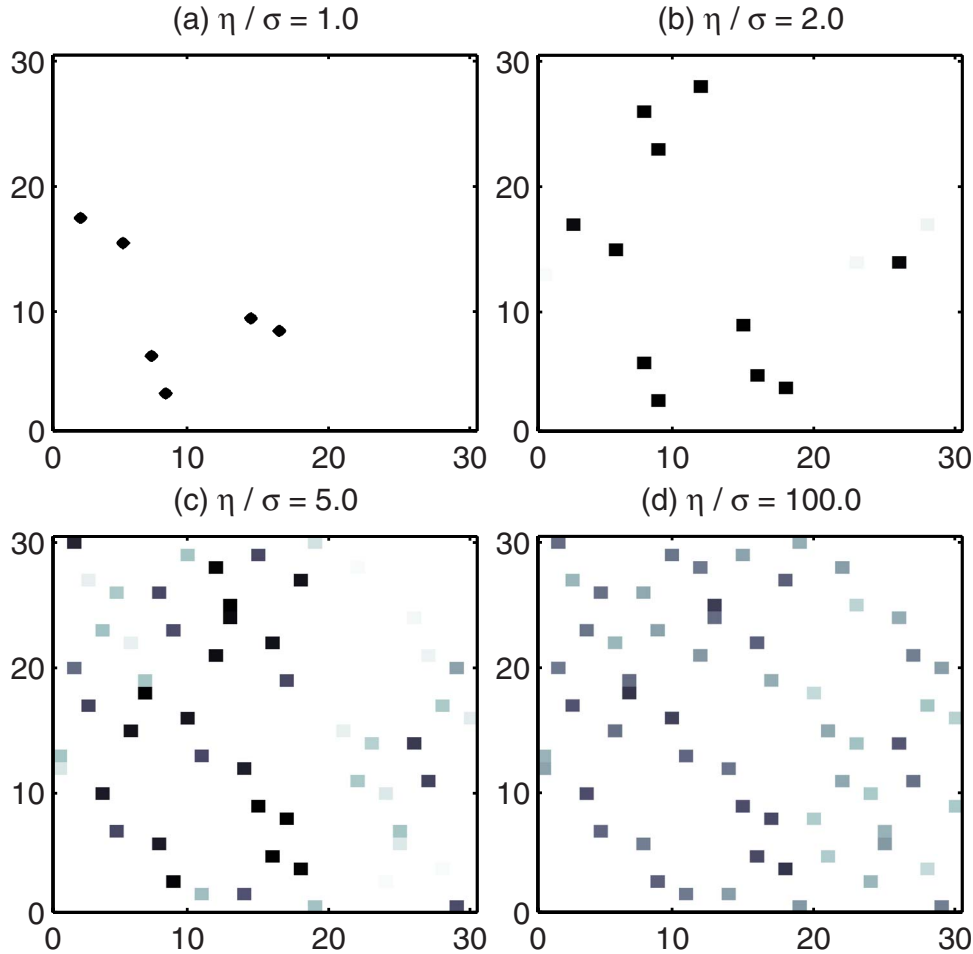


FIG. 9. (Color online) $P_{i,j}$ for long simulations with uniform detuning $\delta=10^{-7}$ and varying noise η . The horizontal axis represents i , and the vertical axis represents j . The dark squares indicate a transition probability close to 1 while the white squares indicate a transition probability close to 0. As before, an initial transient is disregarded. Panel (a) ($\eta/\delta=1.0$) corresponds to the case where noise apparently has no effect, as seen in Fig. 4. Panel (b) ($\eta/\delta=2.0$) shows the effect of low noise that is sufficient to occasionally nudge the system between the two stable period 6 cycles present for this detuning. Panel (c) ($\eta/\delta=5.0$) corresponds to the coding from Fig. 8. Panel (d) ($\eta/\delta=100.0$) corresponds to the case where detuning is comparatively weak and all edges are traversed.

parently independent of those for epoch $n-1$ to n . The entropy of the spatial coding is given by [13]

$$\xi = - \sum_{i,j} p_i P_{i,j} \ln P_{i,j}.$$

Note $\xi=0$ corresponds to an eventually periodic trajectory, while increasing ξ is associated with an increasing range of possibilities of the future trajectory.

For fixed $\delta=10^{-7}$ and varying $0.01 < \frac{\eta}{\delta} < 100$, Fig. 10 shows that the entropy for the system increases from zero up to the theoretical maximum for the network in Fig. 2, namely, $0 \leq \xi \leq \ln 2$. Specifically, the entropy determines the average likelihood that, when the system is at state C_n , the next state C_{n+1} is not that which is predicted by the ordering of the natural frequencies of the oscillators.

For $\eta/\delta < 1.0$ and $\xi=0$ the system will follow the period 6 coding sequence indefinitely. $\eta/\delta=2.0$ corresponds to Fig. 9(b), where the coding sequence deviates from the period 6 cycle on occasion, but quickly returns to the cycle by a predetermined route. $\eta/\delta=100.0$ corresponds to Fig. 9(d),

where the detuning has a minimal effect and the system is almost completely driven by noise. In summary, this shows that the spatial coding of the system is very robust to noise, up to a strength comparable to the order of the amplitude of detuning.

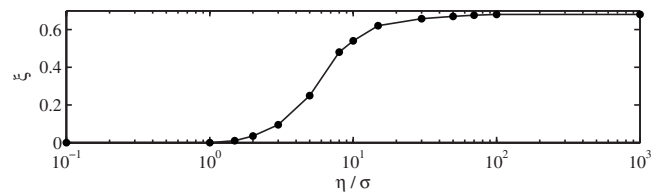


FIG. 10. The entropy ξ in the spatial coding of the system as a function of η/δ from numerical simulations. At $\eta/\delta < 0.5$, $\xi \approx 0$ —there is little uncertainty in the path that the system will follow. The entropy ξ increases quickly between $0.5 < \eta/\delta < 10.0$, where it is close to the theoretical entropy of an unbiased random walk on Fig. 2, namely, $\xi = \ln(2)$.

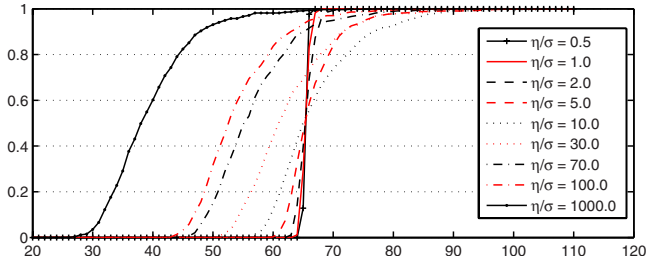


FIG. 11. (Color online) Each data set shows the cumulative frequency $f(T)$ of n such that $C_n=18$ and $T_n < T$ from the simulations depicted in Fig. 10, more precisely $f(T)\mathcal{N}\{n: C_n=18 \text{ and } T_n \leq T\}/\mathcal{N}\{n: C_n=18\}$. As η/δ increases the variance in the sequence of residence times T_n increases and the mean epoch length decreases.

B. Influence of noise on the temporal coding

For $\eta=0$ and $\delta>0$ we expect the sequence T_n to be dependent on the detuning δ_u of the unstable cluster for state C_n , given by Eq. (5). Long simulations were run for a number of different η/δ and the distribution of T_n obtained is shown in Fig. 11. From this figure we see that as η/δ increases, the variance of T_n increases, while the mean decreases slowly for $\eta/\delta < 10$ and more rapidly for $\eta/\delta > 10$ (for comparison, see [11]).

Examining the mean and variance of the distribution of residence times T_n shows that the temporal coding of the system is also robust to noise up to $\eta/\delta \approx 10$. In fact, the mean of the distribution does not significantly change even when the strength of the noise is of the same order as the amplitude of detuning.

V. DISCUSSION

This paper demonstrates that the itinerary between cluster states for a robust heteroclinic cycle for the system of five coupled phase oscillators (1) can be used to encode all information about the “inputs” to the system in terms of the detunings between oscillators. Examining the itinerary (spatial coding) and the residence times at each of these states (temporal coding) gives enough information to obtain all frequency differences between oscillators. The spatial coding gives a partition of the oscillators into two slow and three fast oscillators, while the average residence times allow one to estimate the frequency differences. At each cluster state

only two of the oscillators are in the unstable cluster. This enables one to estimate the frequency difference as an initial perturbation that grows exponentially. It is quite surprising that there is not more information loss; for example, if the coupling were such that an in-phase solution were stable, the response to very low amplitude inputs would be linear and very small, unlike the situation here, where very low amplitude inputs result in a large response.

For larger populations of phase oscillators, very similar behavior can be observed except that the coding becomes much more difficult to describe in detail owing to a combinatorial explosion in the number of cluster states potentially involved, the larger number of switching possibilities involved, and the unstable manifold of each cluster state being higher dimensional. For instance, using the same coupling as in Eq. (2) for $N=9$ oscillators one can find a robust attracting heteroclinic cycle involving over 600 cluster states joined by four-dimensional manifolds [14]. In this case the dynamics in the presence of detuning does not compare all oscillators; there is, as in the case $N=5$, also an eventually period 6 itinerary but this compares only the fastest five oscillators. Nonetheless, we conjecture that there are other coupling functions for higher N cases that will robustly and fully encode the phase differences by involving other cluster states.

We have illustrated, by computation of the entropy of the resulting spatiotemporal sequence, that the effect of noise on the spatiotemporal code generation is surprisingly small up to quite high amplitude noise. In particular, the entropy of the spatial coding remains almost zero up to a signal to noise ratio of 10:1.

Winnerless competition, where an attractor is a heteroclinic network, has been used to describe a variety of neural systems, most particularly in rate-based competition models [5,15]. It is particularly useful for describing the dynamics of neural processes that lead to sequence generation and/or spatiotemporal code generation. The model considered here can be considered as an example of winnerless competition that is not a rate model. In our model, cluster states play the same role as subspaces for rate models where one or more of the neural units has a zero rate. It illustrates that winnerless competition can occur in phase oscillator models and, at least for small numbers of oscillators, that the spatiotemporal coding is efficient and invertible. We suggest that these models can give useful insight to the dynamics of neural assemblies where once again, “instability makes sense” [16].

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