

Significant Excess of Electronlike Events in the MiniBooNE Short-Baseline Neutrino Experiment

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(Dated: October 30, 2018)

The MiniBooNE experiment at Fermilab reports results from an analysis of ν_e appearance data from 12.84×10^{20} protons on target in neutrino mode, an increase of approximately a factor of two over previously reported results. A ν_e charged-current quasielastic event excess of 381.2 ± 85.2 events (4.5σ) is observed in the energy range $200 < E_{\nu}^{QE} < 1250$ MeV. Combining these data with the $\bar{\nu}_e$ appearance data from 11.27×10^{20} protons on target in antineutrino mode, a total ν_e plus $\bar{\nu}_e$ charged-current quasielastic event excess of 460.5 ± 99.0 events (4.7σ) is observed. If interpreted in a two-neutrino oscillation model, $\nu_{\mu} \rightarrow \nu_e$, the best oscillation fit to the excess has a probability of 21.1%, while the background-only fit has a χ^2 probability of 6×10^{-7} relative to the best fit. The MiniBooNE data are consistent in energy and magnitude with the excess of events reported by the Liquid Scintillator Neutrino Detector (LSND), and the significance of the combined LSND and MiniBooNE excesses is 6.0σ . A two-neutrino oscillation interpretation of the data would require at least four neutrino types and indicate physics beyond the three neutrino paradigm. Although the data are fit with a two-neutrino oscillation model, other models may provide better fits to the data.

Evidence for short-baseline neutrino anomalies at an $L/E_{\nu} \sim 1$ m/MeV, where E_{ν} is the neutrino energy and L is the distance that the neutrino traveled before detection, comes from both neutrino appearance and disappearance experiments. The appearance anomalies include the excess of ν_e and $\bar{\nu}_e$ charge-current quasielastic (CCQE) events observed by the LSND [1] and MiniBooNE [2, 3] experiments, while the disappearance anomalies, although not completely consistent, include the deficit of ν_e and $\bar{\nu}_e$ events observed by reactor [4] and radioactive-source experiments [5]. As the masses and

mixings within the 3-generation neutrino matrix have been attached to solar and long-baseline neutrino experiments, more exotic models are typically used to explain these anomalies, including, for example, 3+N neutrino oscillation models involving three active neutrinos and N additional sterile neutrinos [6–14], resonant neutrino oscillations [15], Lorentz violation [16], sterile neutrino decay [17], sterile neutrino nonstandard interactions [18], and sterile neutrino extra dimensions [19]. This Letter presents improved MiniBooNE ν_e and $\bar{\nu}_e$ appearance results, assuming two-neutrino oscillations with probability

$P = \sin^2(2\theta) \sin^2(1.27\Delta m^2 L/E)$, where θ is the mixing angle, Δm^2 (eV^2/c^4) is the difference in neutrino mass eigenstates squared, L (m) is the distance traveled by the neutrino, and E (MeV) is the neutrino energy.

The booster neutrino beam (BNB) at Fermilab delivers to the MiniBooNE experiment a flux of neutrinos and antineutrinos that is simulated using information from external measurements [20]. The BNB is produced by 8 GeV protons from the Fermilab booster interacting on a beryllium target inside a magnetic focusing horn. Depending on the polarity of the horn, either π^+ are focused and π^- are defocused to produce a fairly pure beam of ν_μ , or π^- are focused and π^+ are defocused to produce a somewhat pure beam of $\bar{\nu}_\mu$. In neutrino mode, the ν_μ , $\bar{\nu}_\mu$, ν_e , and $\bar{\nu}_e$ flux contributions at the detector are 93.5%, 5.9%, 0.5%, and 0.1%, respectively, while in antineutrino mode, the flux contributions are 15.7%, 83.7%, 0.2%, and 0.4%, respectively. The ν_μ and $\bar{\nu}_\mu$ fluxes peak at approximately 600 MeV and 400 MeV, respectively.

The MiniBooNE detector is described in detail in reference [21]. The detector consists of a 12.2 m diameter sphere filled with 818 tonnes of pure mineral oil (CH_2) and is located 541 m from the beryllium target. The detector is covered by 1520 8-inch photomultiplier tubes (PMTs), where 1280 PMTs are in the interior detector region and 240 PMTs are located in the optically isolated outer veto region. Charged particles produced by neutrino interactions in the mineral oil emit both directed Cherenkov light and isotropic scintillation light that is detected by the PMTs. Event reconstruction [22] and particle identification make use of the hit PMT charge and time information, and the reconstructed neutrino energy, E_ν^{QE} , is estimated from the measured energy and angle of the outgoing muon or electron, assuming the kinematics of CCQE scattering [23].

From 2002-2017, the MiniBooNE experiment has collected a total of 11.27×10^{20} protons on target (POT) in antineutrino mode, 12.84×10^{20} POT in neutrino mode, and a further 1.86×10^{20} POT in a special beam-off target mode to search for sub-GeV dark matter [24]. The neutrino sample has approximately doubled in size since the previous publication [3]. The published neutrino-mode data correspond to 6.46×10^{20} POT, while 6.38×10^{20} POT were obtained in 2016 and 2017. During the 15 years of running, the BNB and MiniBooNE detector have been stable to within 2% in neutrino energy.

The analysis is optimized to measure ν_e and $\bar{\nu}_e$ induced CCQE events, and the event reconstruction [22] and selection are identical to the previous analysis [3]. The average selection efficiency is $\sim 20\%$ ($\sim 0.1\%$) for ν_e -induced CCQE events (ν_μ -induced background events) generated over the fiducial volume. The fraction of CCQE events in antineutrino mode that are from wrong-sign neutrino events was determined from the angular distributions of muons created in CCQE interactions and by measuring CC single π^+ events [25].

The predicted but unconstrained ν_e and $\bar{\nu}_e$ CCQE background events for the neutrino energy range $200 < E_\nu^{QE} < 1250$ MeV are shown in Table I for both neutrino mode and antineutrino mode [26]. See appendices for more information on backgrounds. The upper limit of 1250 MeV corresponded to a small value of L/E and was chosen by the collaboration before unblinding the data in 2007. The lower limit of 200 MeV is chosen because we constrain the ν_e events with the CCQE ν_μ events and our CCQE ν_μ event sample only goes down to 200 MeV, as we require a visible Cherenkov ring from the muon. The estimated sizes of the intrinsic ν_e and gamma backgrounds are based on MiniBooNE event measurements and uncertainties from these constraints are included in the analysis. The intrinsic $\nu_e/\bar{\nu}_e$ background from muon decay is directly related to the large sample of observed $\nu_\mu/\bar{\nu}_\mu$ events, as these events constrain the muons that decay in the 50 m decay region. This constraint uses a joint fit of the observed $\nu_\mu/\bar{\nu}_\mu$ and $\nu_e/\bar{\nu}_e$ events, assuming that there are no substantial $\nu_\mu/\bar{\nu}_\mu$ disappearance oscillations. The other intrinsic ν_e background component, from kaon decay, is constrained by fits to kaon production data and SciBooNE measurements [27]. The intrinsic ν_e background from pion decay (1.2×10^{-4} branching ratio) and hyperon decay are very small. Other backgrounds from mis-identified ν_μ or $\bar{\nu}_\mu$ [28, 29] events are also constrained by the observed CCQE sample.

The gamma background from neutral-current (NC) π^0 production and $\Delta \rightarrow N\gamma$ radiative decay [30, 31] are constrained by the associated large two-gamma sample (mainly from Δ production) observed in the MiniBooNE data, where π^0 measurements [32] are used to constrain the π^0 background. The π^0 background measured in the first and second neutrino data sets were found to be consistent, resulting in a lower statistical background uncertainty for the combined data. Other neutrino-induced single gamma production processes are included in the theoretical predictions, which agree well with the MiniBooNE estimates [30, 33]. Single-gamma backgrounds from external neutrino interactions (“dirt” backgrounds) are estimated using topological and spatial cuts to isolate the events whose vertices are near the edge of the detector and point towards the detector center [34]. With the larger data set, the background from external neutrino interactions is now better determined to be approximately 7% larger, but with smaller uncertainty, than in the previous publication [3]. A new technique to measure or constrain the gamma and dirt backgrounds based on event timing relative to the beam is in development.

Systematic uncertainties are determined by considering the predicted effects on the ν_μ , $\bar{\nu}_\mu$, ν_e , and $\bar{\nu}_e$ CCQE rates from variations of uncertainty parameters. The parameters include uncertainties in the neutrino and antineutrino flux estimates, uncertainties in neutrino cross sections, most of which are determined by in situ cross-section measurements at MiniBooNE [28, 32], uncertain-

TABLE I: *The expected (unconstrained) number of events for the $200 < E_{\nu}^{QE} < 1250$ MeV neutrino energy range from all of the backgrounds in the ν_e and $\bar{\nu}_e$ appearance analysis before using the constraint from the CC ν_{μ} events. Also shown are the constrained background, as well as the expected number of events corresponding to the LSND best fit oscillation probability of 0.26%, assuming oscillations at large Δm^2 . The table shows the diagonal-element systematic plus statistical uncertainties, which become substantially reduced in the oscillation fits when correlations between energy bins and between the electron and muon neutrino events are included. The antineutrino numbers are from a previous analysis [3].*

Process	Neutrino Mode	Antineutrino Mode
ν_{μ} & $\bar{\nu}_{\mu}$ CCQE	73.7 ± 19.3	12.9 ± 4.3
NC π^0	501.5 ± 65.4	112.3 ± 11.5
NC $\Delta \rightarrow N\gamma$	172.5 ± 24.1	34.7 ± 5.4
External Events	75.2 ± 10.9	15.3 ± 2.8
Other ν_{μ} & $\bar{\nu}_{\mu}$	89.6 ± 22.9	22.3 ± 3.5
ν_e & $\bar{\nu}_e$ from μ^{\pm} Decay	425.3 ± 100.2	91.4 ± 27.6
ν_e & $\bar{\nu}_e$ from K^{\pm} Decay	192.2 ± 41.9	51.2 ± 11.0
ν_e & $\bar{\nu}_e$ from K_L^0 Decay	54.5 ± 20.5	51.4 ± 18.0
Other ν_e & $\bar{\nu}_e$	6.0 ± 3.2	6.7 ± 6.0
Unconstrained Bkgd.	1590.6 ± 176.9	398.2 ± 49.7
Constrained Bkgd.	1577.8 ± 85.2	398.7 ± 28.6
Total Data	1959	478
Excess	381.2 ± 85.2	79.3 ± 28.6
0.26% (LSND) $\nu_{\mu} \rightarrow \nu_e$	463.1	100.0

ties from nuclear effects, and uncertainties in detector modeling and reconstruction. A covariance matrix in bins of E_{ν}^{QE} is constructed by considering the variation from each source of systematic uncertainty on the ν_e and $\bar{\nu}_e$ CCQE signal and background, and the ν_{μ} and $\bar{\nu}_{\mu}$ CCQE prediction as a function of E_{ν}^{QE} . This matrix includes correlations between any of the ν_e and $\bar{\nu}_e$ CCQE signal and background and ν_{μ} and $\bar{\nu}_{\mu}$ CCQE samples, and is used in the χ^2 calculation of the oscillation fits.

Table I also shows the expected number of events corresponding to the LSND best fit oscillation probability of 0.26%, assuming oscillations at large Δm^2 . LSND and MiniBooNE have the same average value of L/E, but MiniBooNE has a larger range of L/E. Therefore, the appearance probabilities for LSND and MiniBooNE should not be exactly the same at lower L/E values.

Fig. 1 shows the E_{ν}^{QE} distribution for ν_e CCQE data and background in neutrino mode for the total 12.84×10^{20} POT data. Each bin of reconstructed E_{ν}^{QE} corresponds to a distribution of “true” generated neutrino energies, which can overlap adjacent bins. In neutrino mode, a total of 1959 data events pass the ν_e CCQE event selection requirements with $200 < E_{\nu}^{QE} < 1250$ MeV, compared to a background expectation of $1577.8 \pm 39.7(stat.) \pm 75.4(syst.)$ events. The excess is then 381.2 ± 85.2 events or a 4.5σ effect. Note that the 162.0 event excess in the first 6.46×10^{20} POT data is approximately 1σ lower than the average excess, while the 219.2 event excess in the second 6.38×10^{20} POT data is approximately 1σ higher than the average ex-

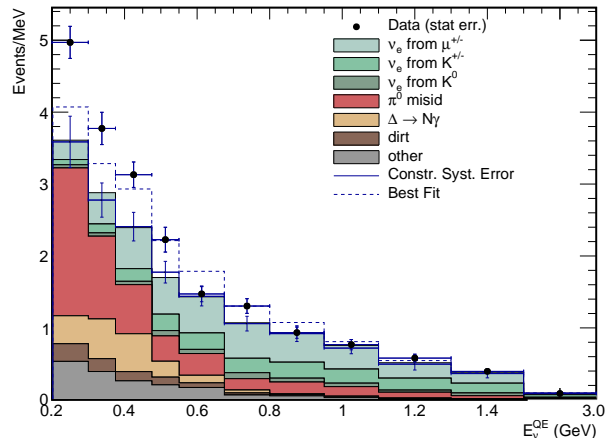


FIG. 1: The MiniBooNE neutrino mode E_{ν}^{QE} distributions, corresponding to the total 12.84×10^{20} POT data, for ν_e CCQE data (points with statistical errors) and background (histogram with systematic errors). The dashed curve shows the best fit to the neutrino-mode data assuming two-neutrino oscillations. The last bin is for the energy interval from 1500-3000 MeV.

cess. Fig. 2 shows the excess events in neutrino mode from the first 6.46×10^{20} POT data and the second 6.38×10^{20} POT data (top plot). Combining the MiniBooNE neutrino and antineutrino data, there are a total of 2437 events in the $200 < E_{\nu}^{QE} < 1250$ MeV energy region, compared to a background expectation of $1976.5 \pm 44.5(stat.) \pm 88.5(syst.)$ events. This corresponds to a total ν_e plus $\bar{\nu}_e$ CCQE excess of 460.5 ± 99.0 events with respect to expectation or a 4.7σ excess. Fig. 2 (bottom plot) shows the total event excesses as a function of E_{ν}^{QE} in both neutrino mode and antineutrino mode. The dashed curves show the two-neutrino oscillation predictions at the best-fit point ($\Delta m^2 = 0.041$ eV², $\sin^2 2\theta = 0.92$), as well as at a point within 1σ of the best-fit point ($\Delta m^2 = 0.4$ eV², $\sin^2 2\theta = 0.01$).

A two-neutrino model is assumed for the MiniBooNE oscillation fits in order to compare with the LSND data. However, the appearance neutrino experiments appear to be incompatible with the disappearance neutrino experiments in a 3+1 model [10, 12], and other models [15–19] may provide better fits to the data. The oscillation parameters are extracted from a combined fit of the observed E_{ν}^{QE} event distributions for muonlike and electronlike events using the full covariance matrix described previously in the full energy range $200 < E_{\nu}^{QE} < 3000$ MeV. The fit assumes the same oscillation probability for both the right-sign ν_e and wrong-sign $\bar{\nu}_e$, and no ν_{μ} , $\bar{\nu}_{\mu}$, ν_e , or $\bar{\nu}_e$ disappearance. Using a likelihood-ratio technique [3], the confidence level values for the fitting statistic, $\Delta\chi^2 = \chi^2(point) - \chi^2(best)$, as a function of oscillation parameters, Δm^2 and $\sin^2 2\theta$, is de-

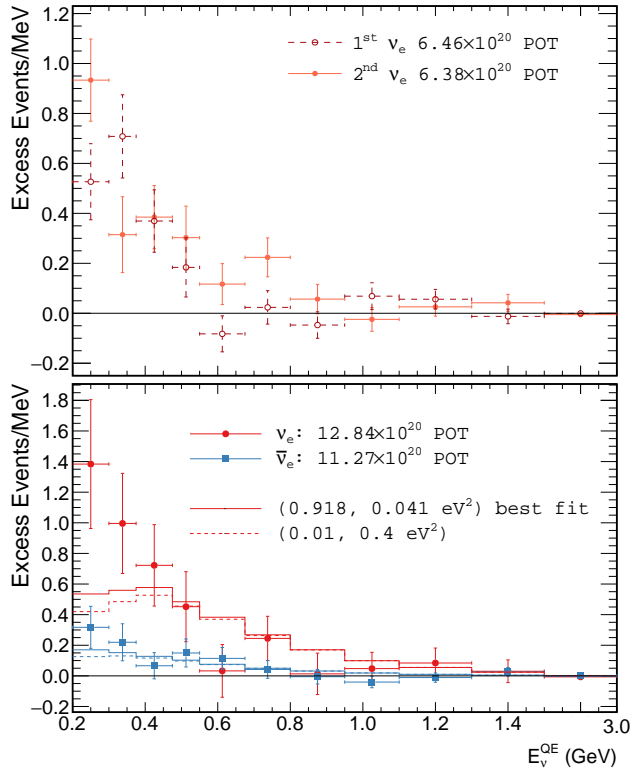


FIG. 2: The top plot shows the MiniBooNE event excesses in neutrino mode as a function of E_{ν}^{QE} from the first 6.46×10^{20} POT data and the second 6.38×10^{20} POT data. The bottom plot shows the total event excesses in both neutrino mode and antineutrino mode, corresponding to 12.84×10^{20} POT and 11.27×10^{20} POT, respectively. The solid (dashed) curve is the best fit (1σ fit point) to the neutrino-mode and antineutrino-mode data assuming two-neutrino oscillations. The last bin is for the energy interval from 1500-3000 MeV. Error bars include only statistical uncertainties for the top plot and both statistical and correlated systematic uncertainties for the bottom plot.

terminated from frequentist, fake data studies. The fake data studies also determine the effective number of degrees of freedom and probabilities. With this technique, the best neutrino oscillation fit in neutrino mode occurs at $(\Delta m^2, \sin^2 2\theta) = (0.039 \text{ eV}^2, 0.84)$, as shown in Fig. 3. The χ^2/ndf for the best-fit point in the energy range $200 < E_{\nu}^{QE} < 1250$ MeV is 9.9/6.7 with a probability of 15.5%. The background-only fit has a χ^2 probability of 0.06% relative to the best oscillation fit and a $\chi^2/ndf = 24.9/8.7$ with a probability of 0.21%. Fig. 3 shows the MiniBooNE closed confidence level (C.L.) contours for ν_e appearance oscillations in neutrino mode in the $200 < E_{\nu}^{QE} < 3000$ MeV energy range.

Nuclear effects associated with neutrino interactions on carbon can affect the reconstruction of the neutrino energy, E_{ν}^{QE} , and the determination of the neutrino oscillation parameters [35]. These effects were studied previ-

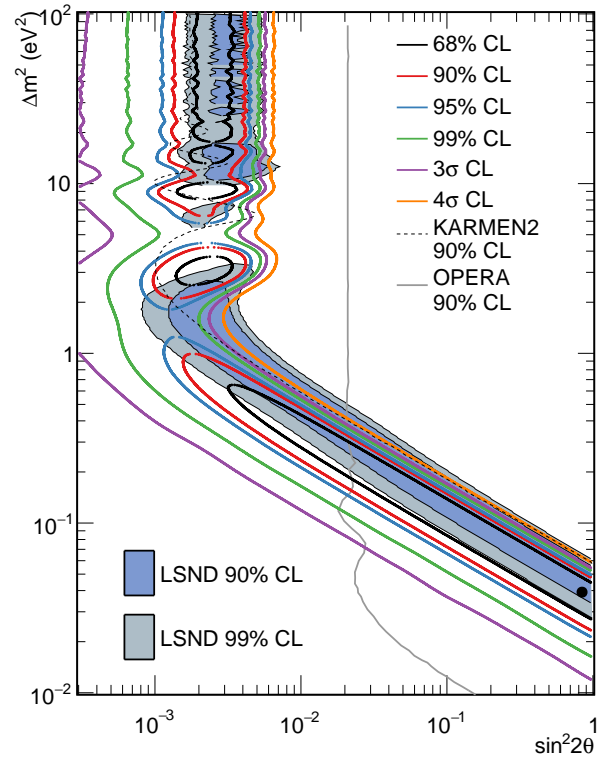


FIG. 3: MiniBooNE allowed regions in neutrino mode (12.84×10^{20} POT) for events with $200 < E_{\nu}^{QE} < 3000$ MeV within a two-neutrino oscillation model. The shaded areas show the 90% and 99% C.L. LSND $\bar{\nu}_{\mu} \rightarrow \bar{\nu}_e$ allowed regions. The black point shows the MiniBooNE best fit point. Also shown are 90% C.L. limits from the KARMEN [37] and OPERA [38] experiments.

ously [3, 36] and were found to not affect substantially the oscillation fit. In addition, they do not affect the gamma background, which is determined from direct measurements of NC π^0 and dirt backgrounds.

Fig. 4 shows the MiniBooNE allowed regions in both neutrino mode and antineutrino mode [3] for events with $200 < E_{\nu}^{QE} < 3000$ MeV within a two-neutrino oscillation model. For this oscillation fit the entire data set is used and includes the 12.84×10^{20} POT data in neutrino mode and the 11.27×10^{20} POT data in antineutrino mode. As shown in the figure, the MiniBooNE 1σ allowed region lies mostly within the LSND 90% C.L. band, which demonstrates good agreement between the LSND and MiniBooNE signals. Also shown are 90% C.L. limits from the KARMEN [37] and OPERA [38] experiments. The KARMEN2 90% C.L. limits are outside the MiniBooNE 95% C.L. allowed region, while the OPERA 90% C.L. limits disfavor the MiniBooNE allowed region below approximately 0.3 eV^2 . The best combined neutrino oscillation fit occurs at $(\Delta m^2, \sin^2 2\theta) = (0.041 \text{ eV}^2, 0.92)$. The χ^2/ndf for the best-fit point in the energy range $200 < E_{\nu}^{QE} < 1250$ MeV is 19.4/15.6 with a probability of 21.1%, and the background-only fit has a χ^2 probab-

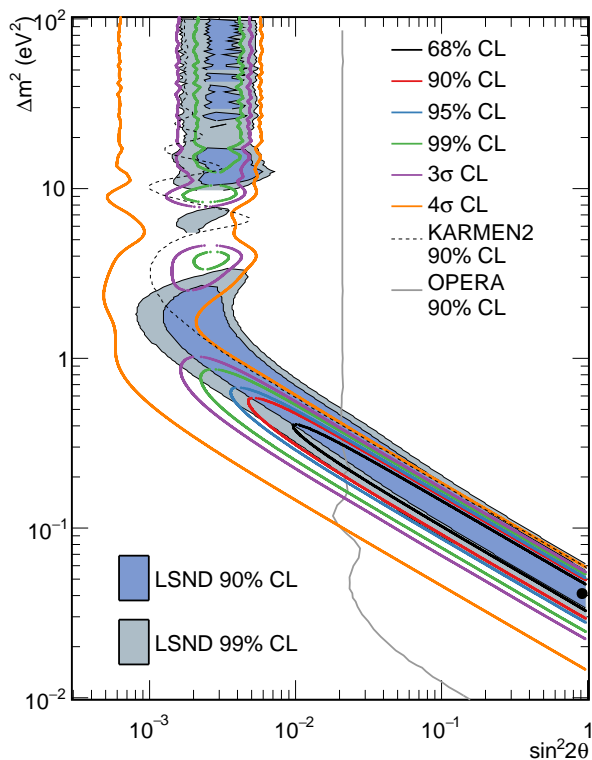


FIG. 4: MiniBooNE allowed regions for a combined neutrino mode (12.84×10^{20} POT) and antineutrino mode (11.27×10^{20} POT) data sets for events with $200 < E_\nu^{QE} < 3000$ MeV within a two-neutrino oscillation model. The shaded areas show the 90% and 99% C.L. LSND $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ allowed regions. The black point shows the MiniBooNE best fit point. Also shown are 90% C.L. limits from the KARMEN2 [37] and OPERA [38] experiments.

ity of 6×10^{-7} relative to the best oscillation fit and a $\chi^2/ndf = 47.1/17.3$ with a probability of 0.02%.

Fig. 5 compares the L/E_ν^{QE} distributions for the MiniBooNE data excesses in neutrino mode and antineutrino mode to the L/E distribution from LSND [1]. The error bars show statistical uncertainties only. As shown in the figure, there is agreement among all three data sets. Assuming two-neutrino oscillations, the curves show fits to the MiniBooNE data described above. Fitting both MiniBooNE and LSND data, by adding LSND L/E data as additional terms, the best fit occurs at $(\Delta m^2, \sin^2 2\theta) = (0.041 \text{ eV}^2, 0.96)$ with a $\chi^2/ndf = 22.4/22.4$, corresponding to a probability of 42.5%. The MiniBooNE excess of events in both oscillation probability and L/E spectrum is, therefore, consistent with the LSND excess of events. The significance of the combined LSND (3.8σ) [1] and MiniBooNE (4.7σ) excesses is 6.0σ , which is obtained by adding the significances in quadrature, as the two experiments have completely different neutrino energies, neutrino fluxes, reconstructions, backgrounds, and systematic uncertainties.

In summary, the MiniBooNE experiment observes a

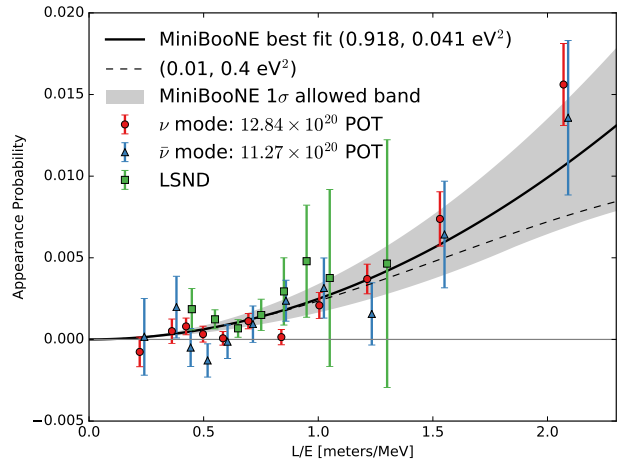


FIG. 5: A comparison between the L/E_ν^{QE} distributions for the MiniBooNE data excesses in neutrino mode (12.84×10^{20} POT) and antineutrino mode (11.27×10^{20} POT) to the L/E distribution from LSND [1]. The error bars show statistical uncertainties only. The curves show fits to the MiniBooNE data, assuming two-neutrino oscillations, while the shaded area is the MiniBooNE 1σ allowed band. The best-fit curve uses the reconstructed neutrino energy, E_ν^{QE} , for the MiniBooNE data. The dashed curve shows the example 1σ fit point.

total ν_e CCQE event excess in both neutrino and antineutrino running modes of 460.5 ± 99.0 events (4.7σ) in the energy range $200 < E_\nu^{QE} < 1250$ MeV. The MiniBooNE allowed region from a two-neutrino oscillation fit to the data, shown in Fig. 4, is consistent with the allowed region reported by the LSND experiment [1]. On the other hand, a two-neutrino oscillation interpretation of the data would require at least four neutrino types and indicate physics beyond the three neutrino paradigm. The significance of the combined LSND and MiniBooNE excesses is 6.0σ . All of the major backgrounds are constrained by in situ event measurements, so nonoscillation explanations would need to invoke new anomalous background processes. Although the data are fit with a two-neutrino oscillation model, other models may provide better fits to the data. The MiniBooNE event excess will be further studied by the Fermilab short-baseline neutrino (SBN) program [39].

We acknowledge the support of Fermilab, the Department of Energy, and the National Science Foundation, and we acknowledge Los Alamos National Laboratory for LDRD funding.

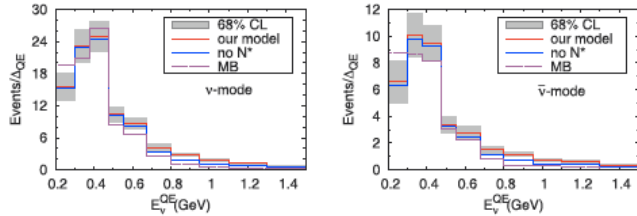


FIG. 6: A comparison of the single gamma background estimate from reference [31] with that of MiniBooNE, where good agreement is obtained.

Appendix: Background Determination & Data vs Monte Carlo Comparisons

Almost all of the backgrounds in the electron-neutrino candidate event sample are determined directly from MiniBooNE data [26]. The muon-neutrino charged-current quasi-elastic (CCQE) observed data sample allows the background determination of both mis-identified ν_μ CCQE events from pion decay [28] and ν_e CCQE events from muon decay, as the neutrinos come from the same parent pion and ν_e and ν_μ cross sections are the same from lepton universality after correcting for charged lepton mass effects. The ν_e CCQE background from kaon decay was determined from external measurements and confirmed by data from the SciBooNE experiment [27]. In addition, the neutral current (NC) π^0 data sample allows the background determination of both NC π^0 events [32] and single gamma events from $\Delta \rightarrow N\gamma$ decays. Various theoretical estimates [30, 31] have confirmed the MiniBooNE single gamma background estimate. Fig. 6 shows a comparison of the single gamma background estimate from reference [31] with that of MiniBooNE, where good agreement is obtained. Single-gamma backgrounds from external neutrino interactions (“dirt” backgrounds) are estimated using topological and spatial cuts to isolate the events whose vertices are near the edge of the detector and point towards the detector center [34]. These estimates have been confirmed by preliminary measurements of the absolute event time reconstruction of electron-neutrino candidate events, where a fit to the dirt event background using timing agrees within 10% with the background estimate using topological and spatial cuts.

In order to demonstrate that the MiniBooNE background estimates are reliable, various comparisons between the neutrino data, corresponding to 12.84×10^{20}

protons on target (POT), and the Monte Carlo simulation have been performed to check and confirm the accuracy of the simulation. Fig. 7 shows an absolute comparison of the π^0 reconstructed mass distribution between the data and the simulation for NC π^0 events. Excellent agreement is obtained, and the ratio of the number of data events (42,483) to the number of Monte Carlo events (42,530) in the mass range from 80 to 200 MeV/ c^2 is equal to 0.999. Fig. 8 shows an absolute comparison of the reconstructed neutrino energy distribution for CCQE events between the data and the simulation. Excellent agreement is also obtained, and the ratio of the number of data events (232,096) to the number of Monte Carlo events (236,145) is equal to 0.983.

In order to check the particle identification (PID) cuts, Figs. 9, 10, and 11 show comparisons between the data and simulation for the electron-muon likelihood distribution, the electron-pion likelihood distribution, and the gamma-gamma mass distribution. In each figure, distributions are shown after successive cuts are applied: no PID cut, electron-muon likelihood cut, electron-muon plus electron-pion likelihood cuts, and electron-muon plus electron-pion likelihood cuts and a gamma-gamma mass cut. The last plot in each figure shows distributions with the final event selection. The vertical lines in the figures show the range of energy-dependent cut values. Good agreement between the data and the simulation is obtained outside the cut values, while an excess of events is observed inside the cut values. Figs. 12 and 13 show the momentum and gamma-gamma opening angle distributions after successive cuts are applied. Good agreement is obtained between the data and Monte Carlo simulation for the no PID cut distributions, while event excesses are observed after the final event selection. These five plots also demonstrate that sidebands show good agreement between the data and simulation.

Fig. 14 shows the visible energy (Evis) and $\cos \theta_e$ (Uz) distributions for the electron-neutrino candidate events in neutrino mode (top) and antineutrino mode (bottom). Also shown in the figures are the expectations from all known backgrounds and from the oscillation best fit. These distributions are important because the reconstructed neutrino energy, E_ν^{QE} , is determined from Evis and Uz. A check on the spatial reconstruction is shown in Fig. 15, where the spatial reconstruction in the data is compared to the Monte Carlo simulation. As shown in the figure, the event excess is evenly distributed up to the 5m radius cut. A further check of the electron efficiency was obtained from the reconstruction of electron events in MiniBooNE that originated from the off-axis NUMI beam (P. Adamson et al., Phys. Rev. Lett. **102**, 211801 (2009)), as the intrinsic ν_e background was approximately ten times higher in the NUMI beam than in the BNB. The measured electron rate agreed with the simulation within errors.

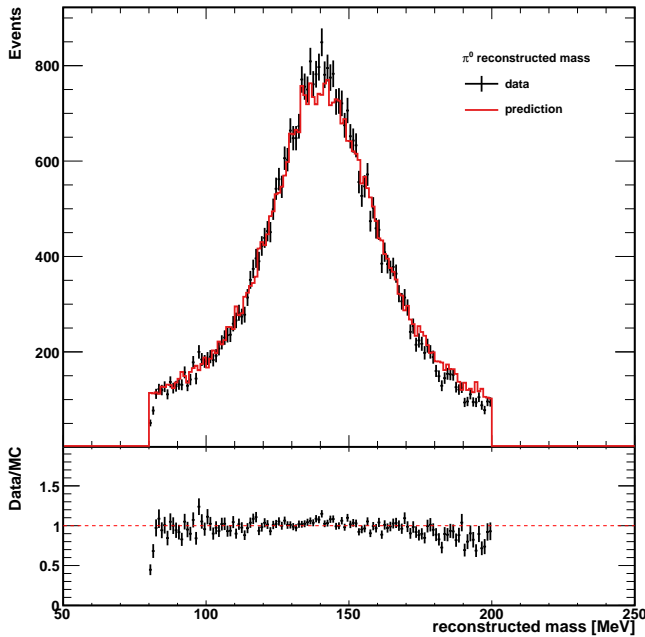


FIG. 7: An absolute comparison of the π^0 reconstructed mass distribution between the neutrino data (12.84×10^{20} POT) and the simulation for NC π^0 events (top). Also shown is the ratio between the data and Monte Carlo simulation (bottom). The error bars show only statistical uncertainties.

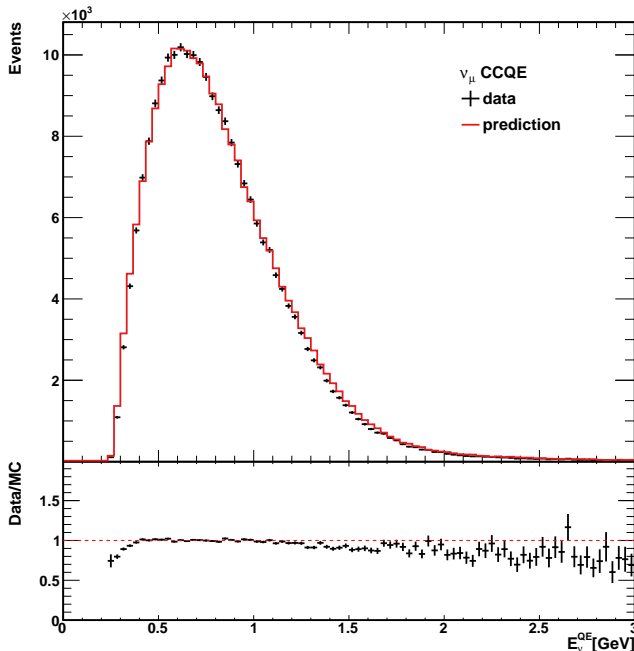


FIG. 8: An absolute comparison of the reconstructed neutrino energy distribution for CCQE events between the neutrino data (12.84×10^{20} POT) and the simulation (top). Also shown is the ratio between the data and Monte Carlo simulation (bottom). The error bars show only statistical uncertainties.

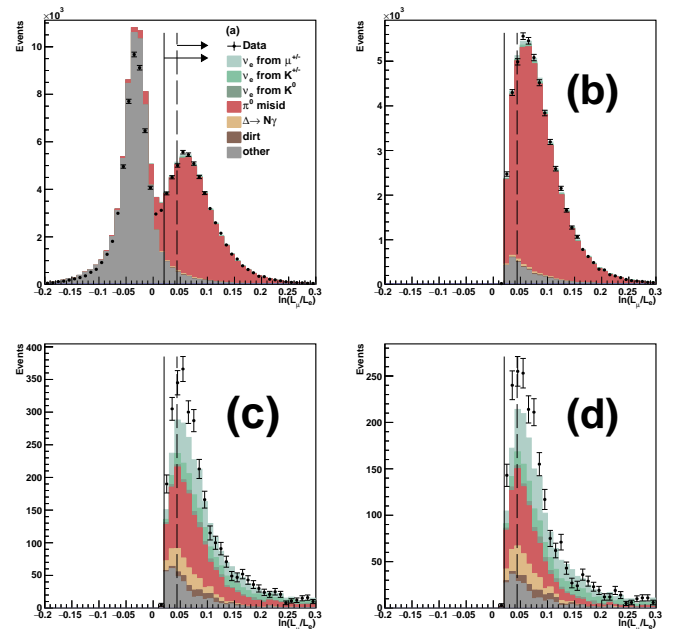


FIG. 9: Comparisons between the data and simulation for the electron-muon likelihood distribution after successive cuts are applied: (a) no PID cut, (b) electron-muon likelihood cut, (c) electron-muon plus electron-pion likelihood cuts, and (d) electron-muon plus electron pion likelihood cuts plus a gamma-gamma mass cut. The vertical lines in the figures show the range of energy-dependent cut values. The error bars show only statistical uncertainties.

Appendix: Stability Checks

Many checks have been performed on the data, including beam and detector stability checks. Fig. 16 shows the total number of neutrino events observed per 10^{17} POT over the lifetime of MiniBooNE in neutrino mode, antineutrino mode and beam-dump mode. The neutrino mode event rate of 100 events per 10^{17} POT has been stable to $< 2\%$ over the 15 year MiniBooNE running period. This is within the expected errors from time variations in BNB performance, such as target/horn change, beam rate monitoring, etc. A small change in the detector energy response between the first and second neutrino data set has been corrected by increasing the measured energy in the second data set by 2%. About half of the energy change is from PMT failures in the intervening years, and the remainder is within the detector response error from gain variations, oil properties, etc. With this energy correction, the first and second data sets are found to agree well. Fig. 17 compares the reconstructed ν_μ CCQE energy distributions for the second data set in 2016 and 2017 to the first data set, where good agreement is obtained. Likewise, Fig. 18 shows that the π^0 mass distribution has also not changed.

Appendix: Comparing New and Old Neutrino Data

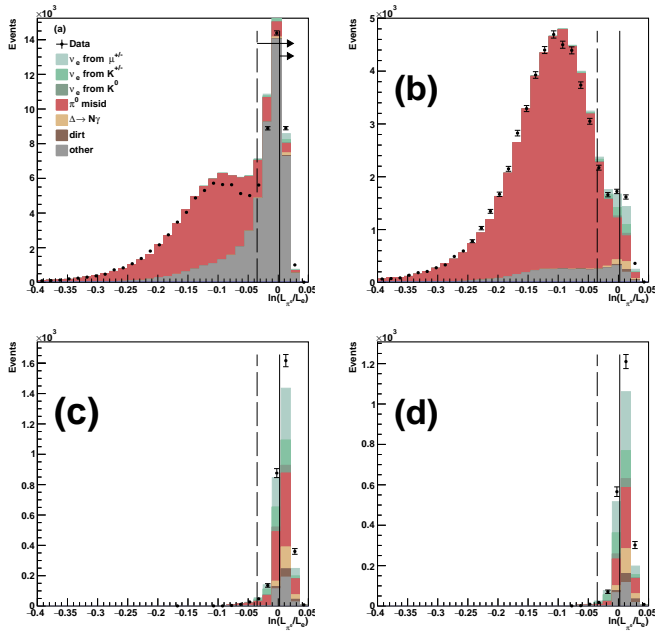


FIG. 10: Comparisons between the data and simulation for the electron-pion likelihood distribution after successive cuts are applied: (a) no PID cut, (b) electron-muon likelihood cut, (c) electron-muon plus electron-pion likelihood cuts, and (d) electron-muon plus electron pion likelihood cuts plus a gamma-gamma mass cut. The vertical lines in the figures show the range of energy-dependent cut values. The error bars show only statistical uncertainties.

Figs. 19 and 20 show the E_ν^{QE} distribution for ν_e CCQE data and background in neutrino mode over the full available energy range for the first 6.46×10^{20} POT data set and the second 6.38×10^{20} POT data set. Fig. 21 shows the ν_e CCQE data and background in antineutrino mode. Each bin of reconstructed E_ν^{QE} corresponds to a distribution of “true” generated neutrino energies, which can overlap adjacent bins. Note that the 162.0 event excess in the 6.46×10^{20} POT data is approximately 1σ lower than the average excess, while the 219.2 event excess in the 6.38×10^{20} POT data is approximately 1σ higher than the average excess. In antineutrino mode, a total of 478 data events pass the ν_e CCQE event selection requirements with $200 < E_\nu^{QE} < 1250$ MeV, compared to a background expectation of $398.7 \pm 20.0(stat.) \pm 20.5(syst.)$ events. The excess is then 79.3 ± 28.6 events or a 2.8σ effect.

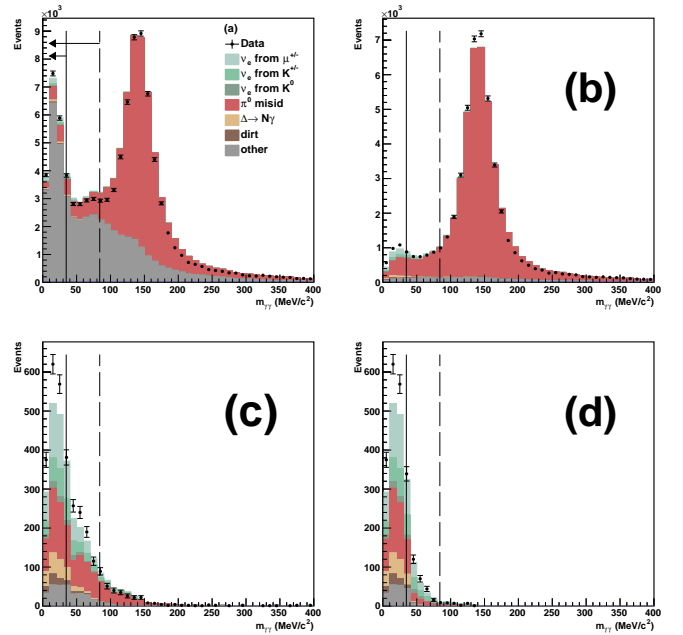


FIG. 11: Comparisons between the data and simulation for the gamma-gamma mass distribution after successive cuts are applied: (a) no PID cut, (b) electron-muon likelihood cut, (c) electron-muon plus electron-pion likelihood cuts, and (d) electron-muon plus electron pion likelihood cuts plus a gamma-gamma mass cut. The vertical lines in the figures show the range of energy-dependent cut values. The error bars show only statistical uncertainties.

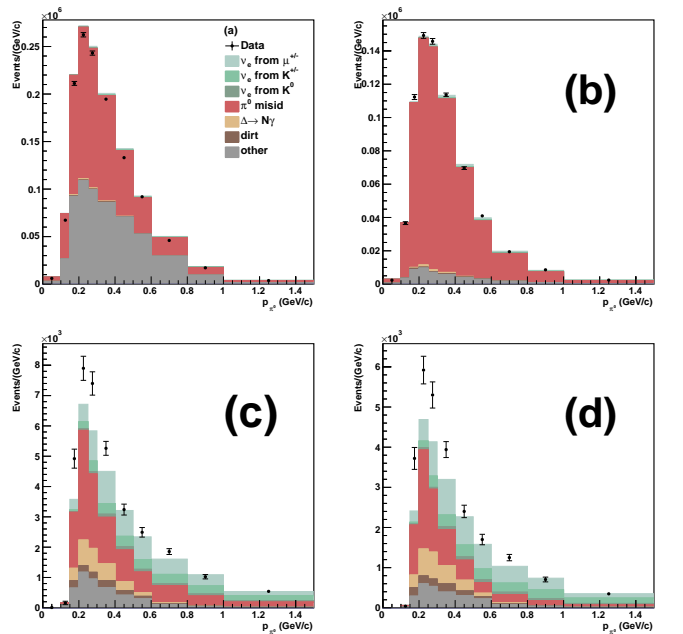


FIG. 12: Comparisons between the data and simulation for the momentum distribution after successive cuts are applied: (a) no PID cut, (b) electron-muon likelihood cut, (c) electron-muon plus electron-pion likelihood cuts, and (d) electron-muon plus electron pion likelihood cuts plus a gamma-gamma mass cut. The error bars show only statistical uncertainties.

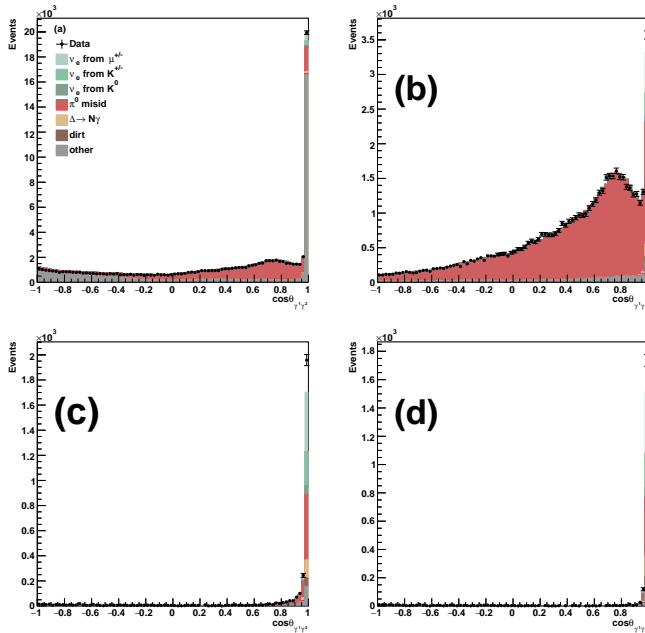


FIG. 13: Comparisons between the data and simulation for the gamma-gamma opening angle distribution after successive cuts are applied: (a) no PID cut, (b) electron-muon likelihood cut, (c) electron-muon plus electron-pion likelihood cuts, and (d) electron-muon plus electron pion likelihood cuts plus a gamma-gamma mass cut. The event excess occurs almost entirely for opening angles less than 13 degrees. The error bars show only statistical uncertainties.

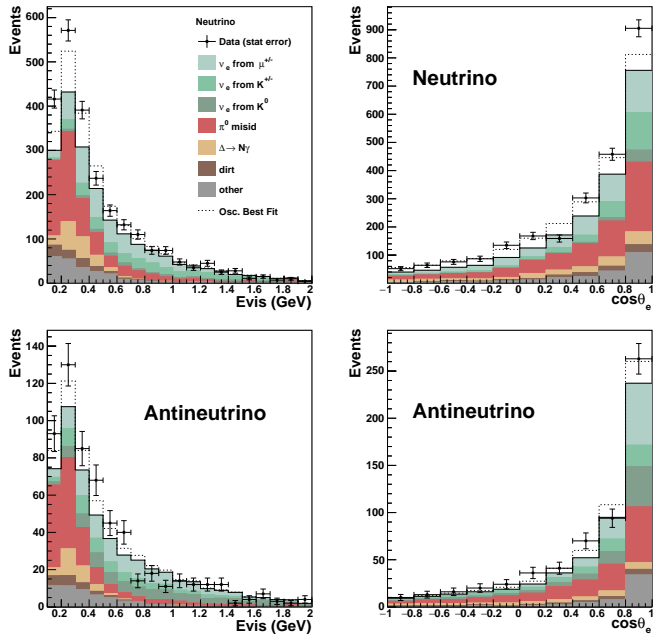


FIG. 14: The visible energy (Evis) and $\cos\theta_e$ (U_z) distributions for the electron-neutrino candidate events in neutrino mode (top) and antineutrino mode (bottom). (The error bars show only statistical uncertainties.) Also shown in the figure are the expectations from all known backgrounds and from the oscillation best fit.

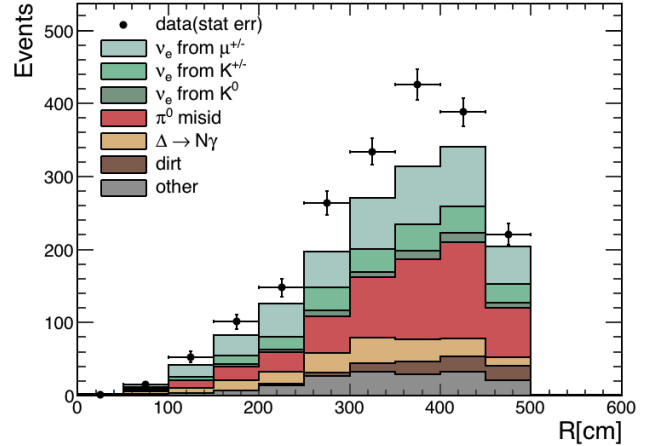


FIG. 15: The radius reconstruction in the data is compared to the Monte Carlo simulation. The event excess is evenly distributed up to the 5m radius cut.

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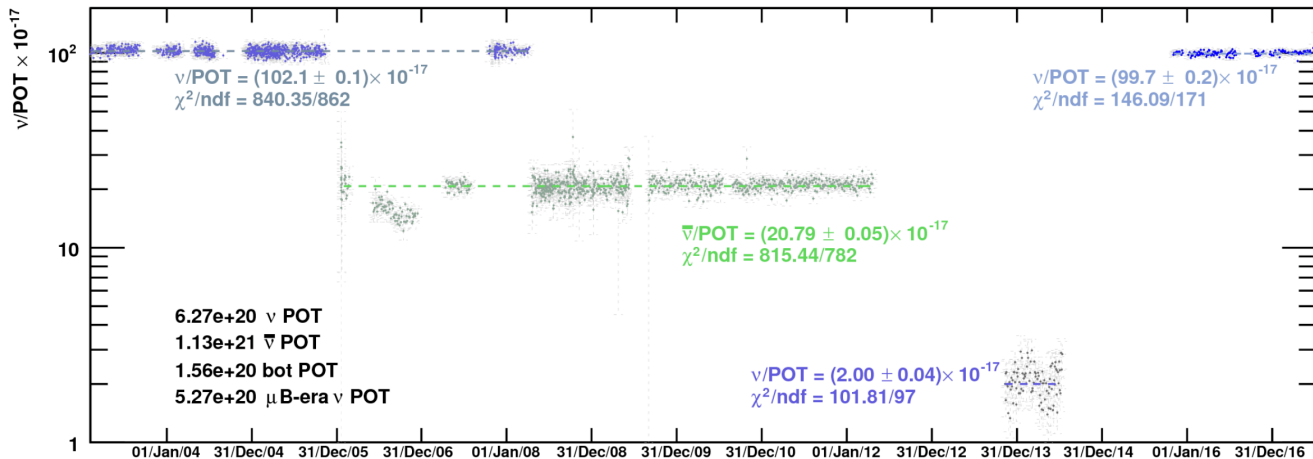


FIG. 16: The total number of neutrino events observed per 10^{17} POT over the lifetime of MiniBooNE in neutrino mode, antineutrino mode and beam-dump mode.

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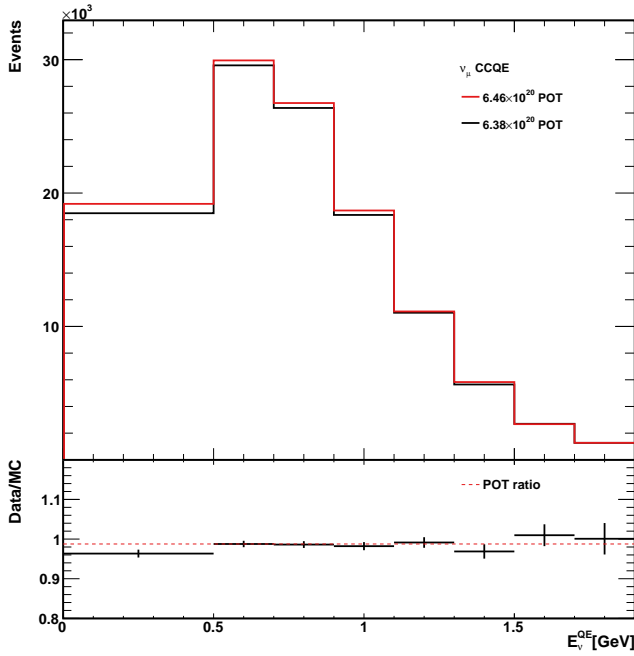


FIG. 17: The top plot shows a comparison between the reconstructed ν_μ CCQE energy distributions for the second data set in 2016 and 2017 (6.38×10^{20} POT) to the first data set (6.46×10^{20} POT). The bottom plot shows the ratio of the second data set to the first data set.

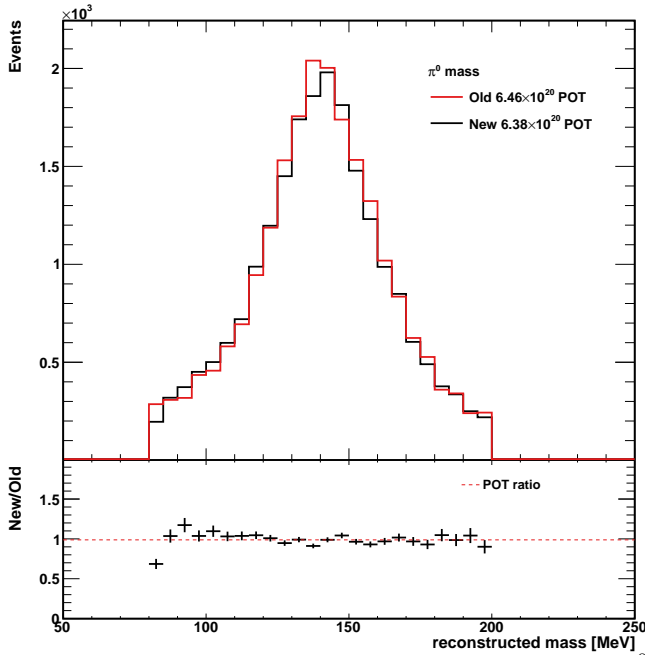


FIG. 18: The top plot shows a comparison between the π^0 mass distributions for the second data set in 2016 and 2017 (6.38×10^{20} POT) to the first data set (6.46×10^{20} POT). The bottom plot shows the ratio of the second data set to the first data set.

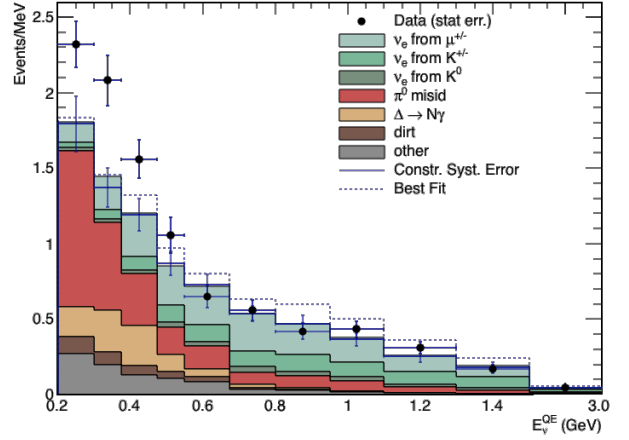


FIG. 19: The neutrino mode E_ν^{QE} distributions, corresponding to the first 6.46×10^{20} POT data set, for ν_e CCQE data (points with statistical errors) and background (histogram with systematic errors).

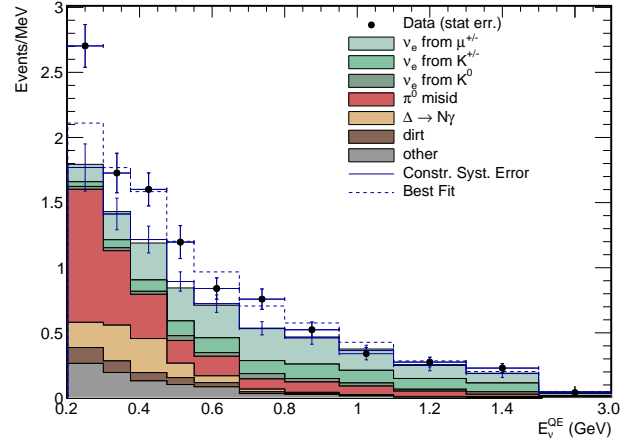


FIG. 20: The neutrino mode E_ν^{QE} distributions, corresponding to the second 6.38×10^{20} POT data set, for ν_e CCQE data (points with statistical errors) and background (histogram with systematic errors).

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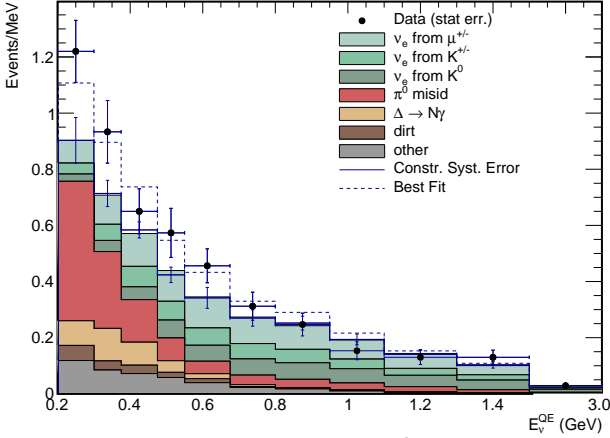


FIG. 21: The antineutrino mode E_{ν}^{QE} distributions, corresponding to the published 11.27×10^{20} POT data, for ν_e CCQE data (points with statistical errors) and background (histogram with systematic errors).

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